

# Locating the critical end point of QCD

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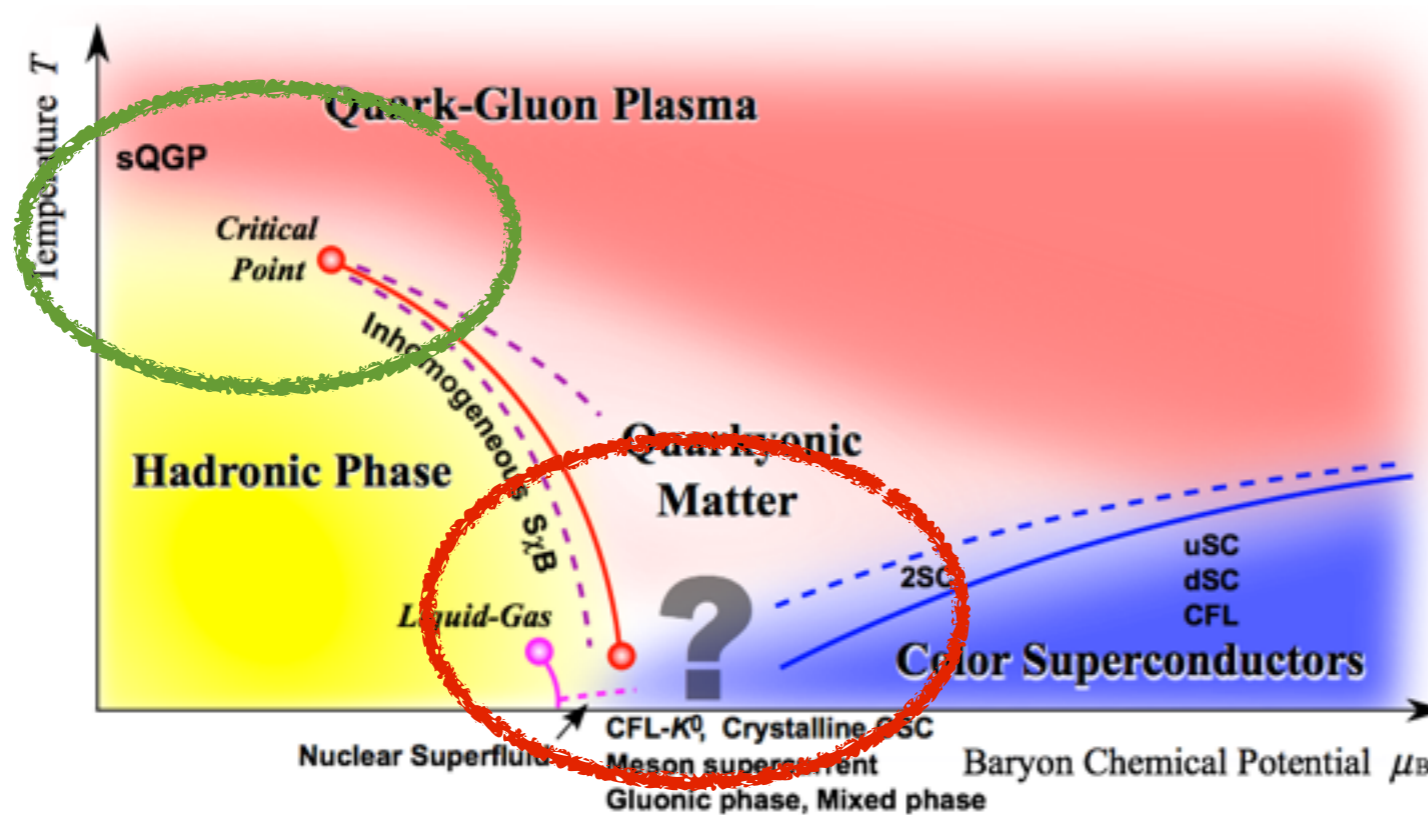


with Jan Luecker and Christian Wezbacher, arXiv:1405.4762

and Leonard Fister, Axel Maas, Jan Pawłowski

# QCD phase transitions I

Fukushima, Hatsuda, Rept. Prog. Phys. 74 (2011) 014001



DSEs at low T, large  $\mu$ :

Müller, Buballa, Wambach, EPJA 49 (2013),  
PLB 727 (2013) 240

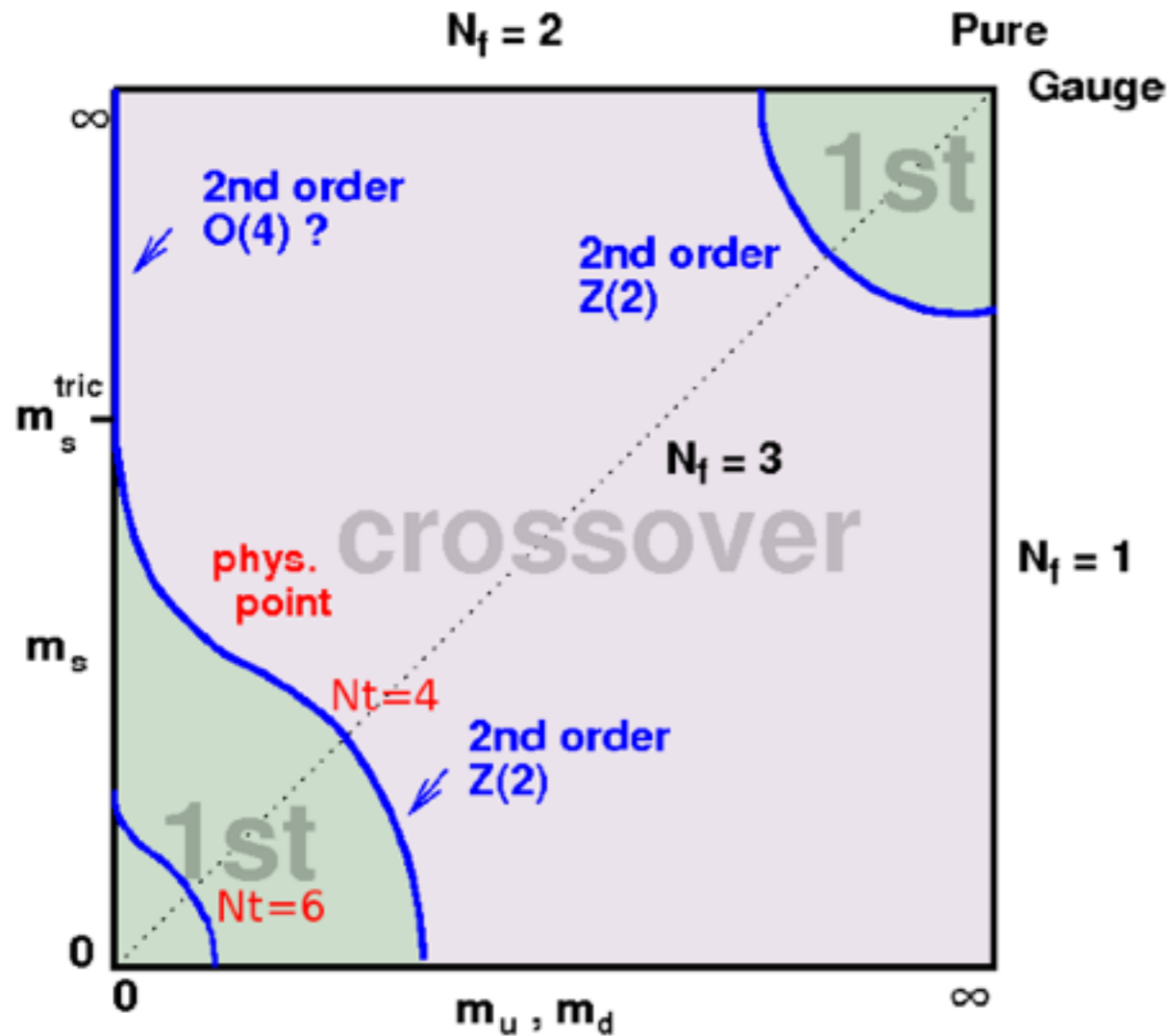
- Chiral limit ( $M_{\text{weak}} \rightarrow 0$ ): order parameter chiral condensate

$$\langle \bar{\Psi} \Psi \rangle \sim \text{Tr} \int S$$

- Static quarks ( $M_{\text{weak}} \rightarrow \infty$ ): order parameter Polyakov-loop

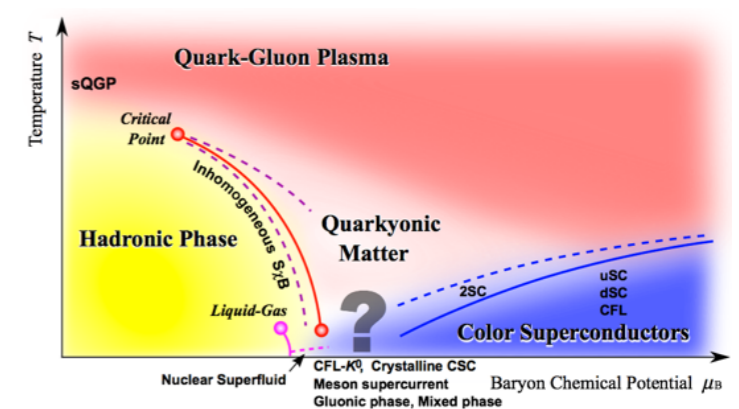
$$\langle |L| \rangle \sim e^{-F_q/T}$$

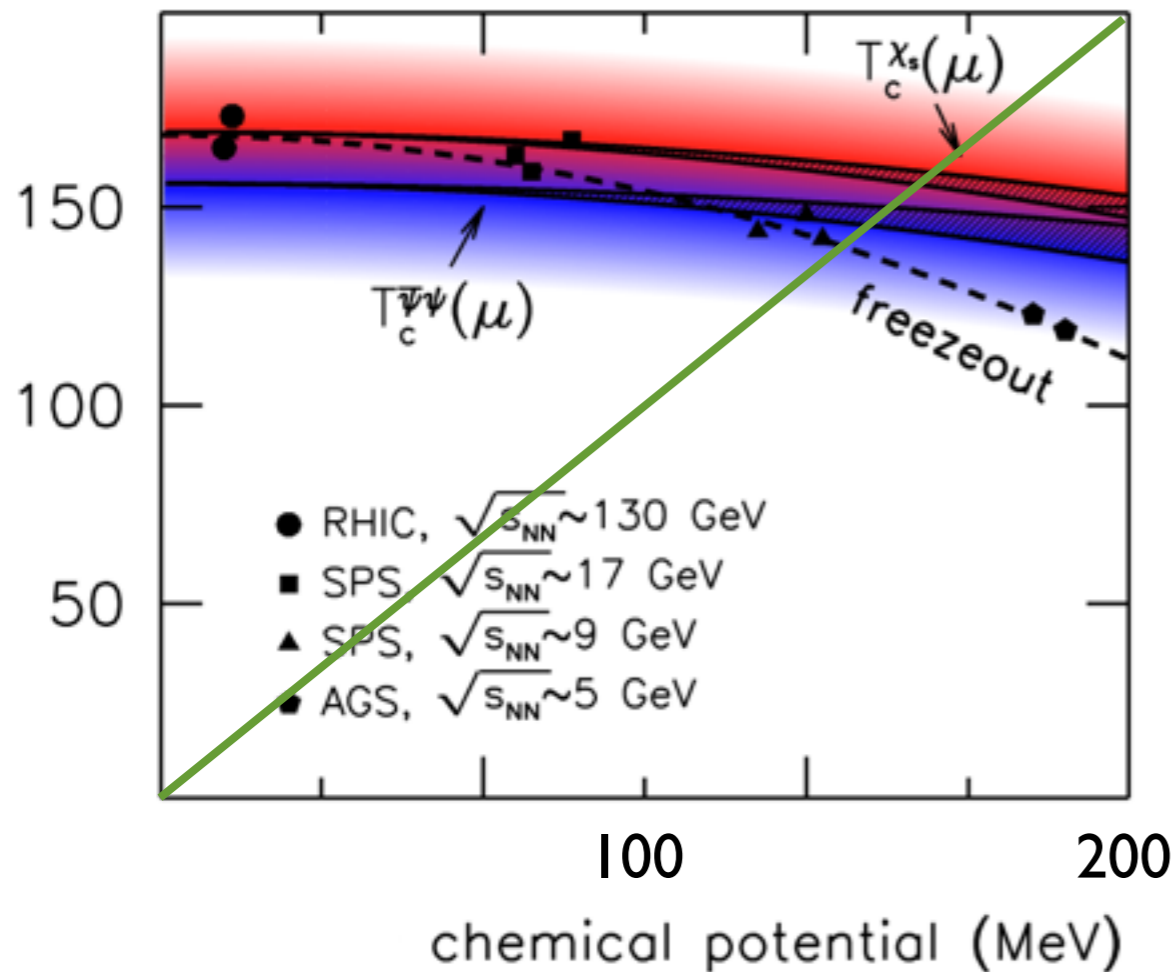
# QCD phase transitions II



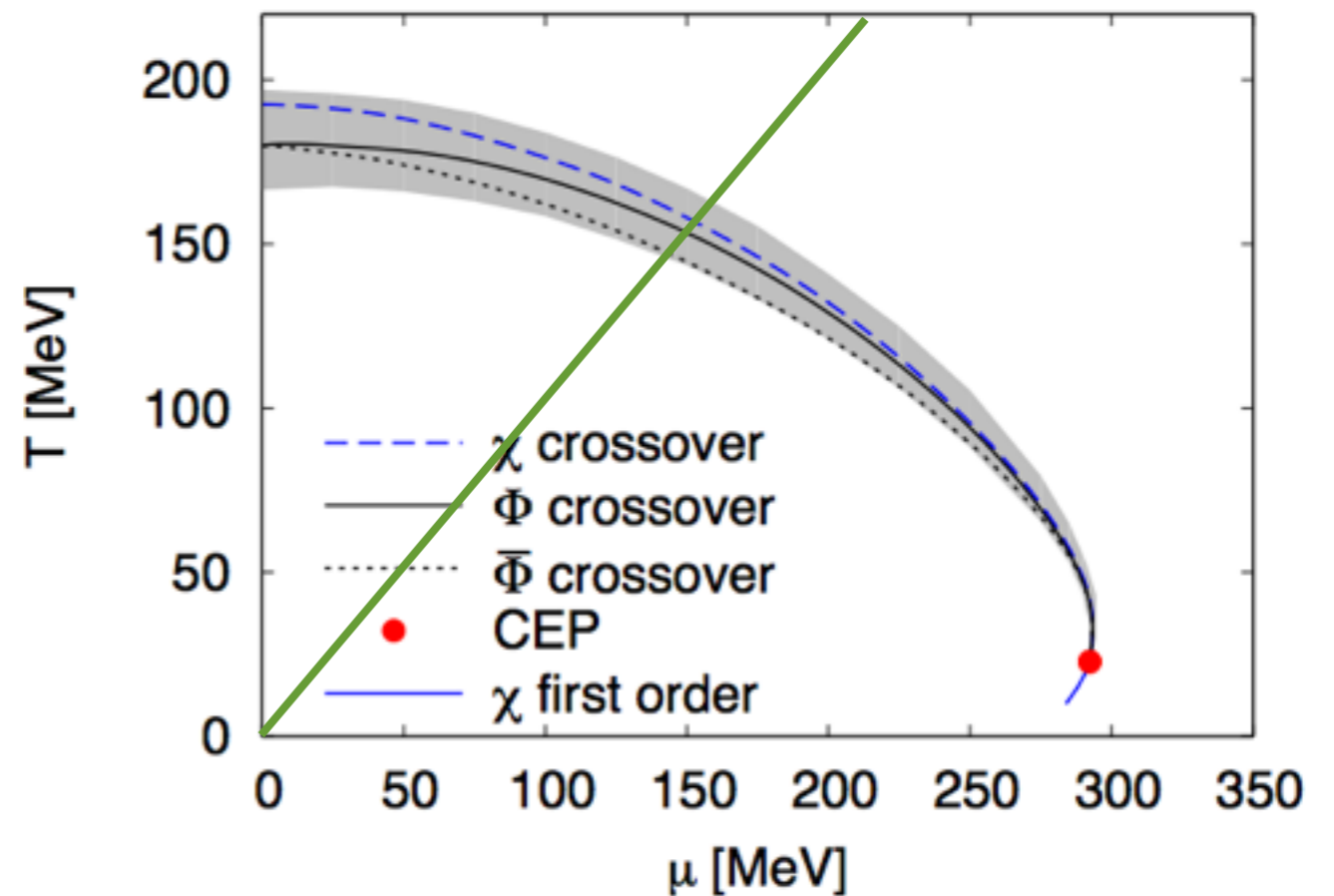
Plot: O. Philipsen

Is this happening ??





Endrodi, Fodor, Katz, Szabo, JHEP 1104 (2011) 001



Herbst, Pawłowski, Schaefer, PLB 696 (2011) 58

- Lattice extrapolation reliable for  $\mu/T \leq 1$
- No CEP for small chemical potential
- PQM plus RG-methods (functional methods)

# QCD order parameters from propagators

$$\text{---} \overset{-1}{\bullet} \text{---} = \text{---} \overset{-1}{\text{---}} \text{---} - \text{---} \overset{\text{---}}{\text{---}} \text{---}$$

Chiral order parameter:

$$\langle \bar{\Psi} \Psi \rangle = Z_2 N_c \text{Tr}_D \frac{1}{T} \sum_{\omega} \int \frac{d^3 p}{(2\pi)^3} S(\vec{p}, \omega)$$

Deconfinement:

- dressed Polyakov loop

$$\Sigma = - \int_0^{2\pi} \frac{d\varphi}{2\pi} e^{-i\varphi} \langle \bar{\Psi} \Psi \rangle_{\varphi}$$

Synatschke, Wipf, Wozar, PRD 75, 114003 (2007)  
 Bilgici, Bruckmann, Gattringer, Hagen, PRD 77 094007 (2008)  
 CF, PRL 103 052003 (2009)

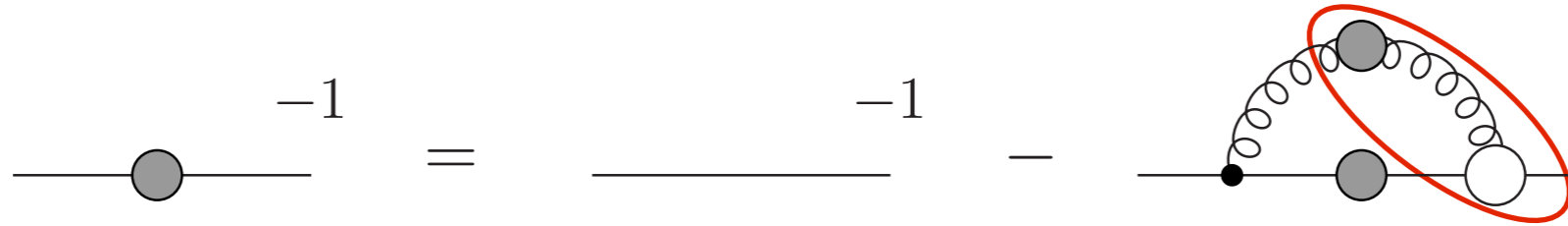
- Polyakov loop potential

$$L = \frac{1}{N_c} \text{Tr} e^{ig\beta A_0}$$

$$\frac{\delta(\Gamma - S)}{\delta A_0} = \frac{1}{2} \left( \text{---} \text{---} \text{---} \text{---} \right) - \left( \text{---} \text{---} \text{---} \text{---} \right) - \left( \text{---} \text{---} \text{---} \text{---} \right) - \frac{1}{6} \left( \text{---} \text{---} \text{---} \text{---} \right) + \left( \text{---} \text{---} \text{---} \text{---} \right)$$

Braun, Gies, Pawłowski, PLB 684, 262 (2010)  
 Braun, Haas, Marhauser, Pawłowski, PRL 106 (2011)  
 Fister, Pawłowski, PRD 88 045010 (2013)  
 CF, Fister, Luecker, Pawłowski, arXiv:1306.6022

# The DSE for the quark propagator



$$[S(p)]^{-1} = [-i\not{p} + M(p^2)]/Z_f(p^2)$$

Input:

- dressed Gluon propagator
- dressed Quark-Gluon-Vertex

Two strategies: I. use **model** for gluon and vertex

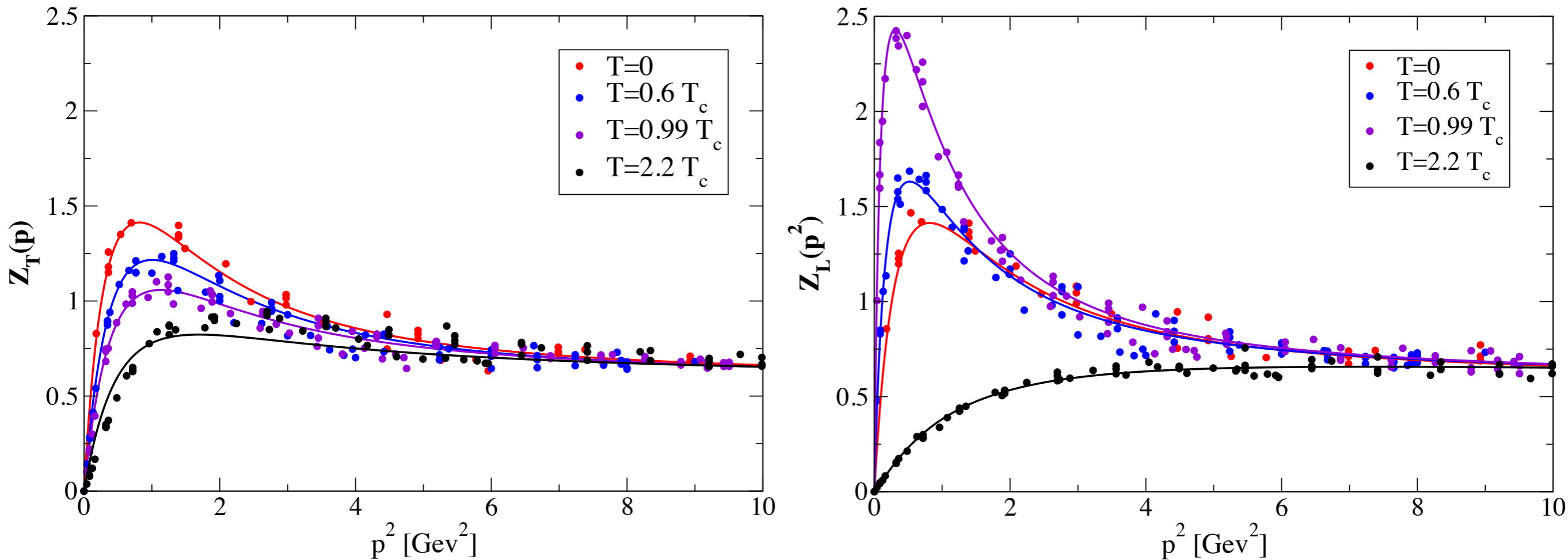
Qin, Chang, Chen, Liu and Roberts, PRL 106 (2011) 172301  
Gutierrez, Ahmad, Ayala, Bashir and Raya, JPG 41 (2014) 075002

- ok for first insights
- not good enough for systematic study

II. **determine gluon and vertex** explicitly

# Glue at finite temperature ( $T \neq 0$ )

T-dependent gluon propagator from quenched lattice simulations:



- Crucial difference between magnetic and electric gluon
- Maximum of electric gluon near  $T_c$

Cucchieri, Maas, Mendes, PRD 75 (2007)

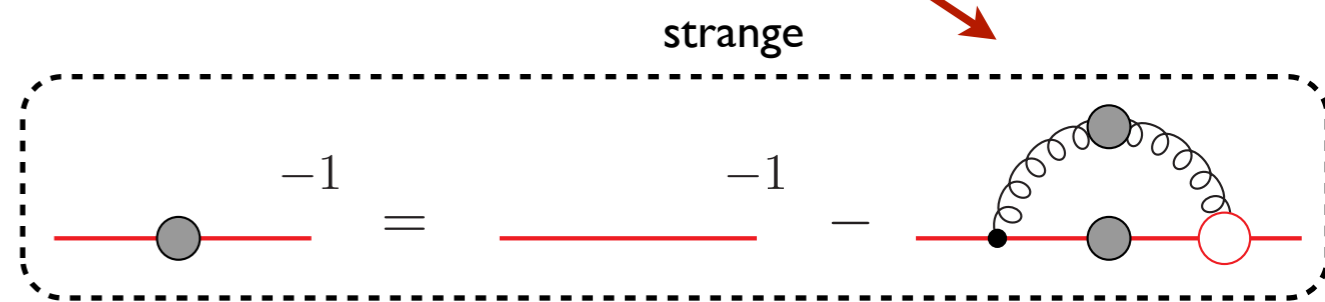
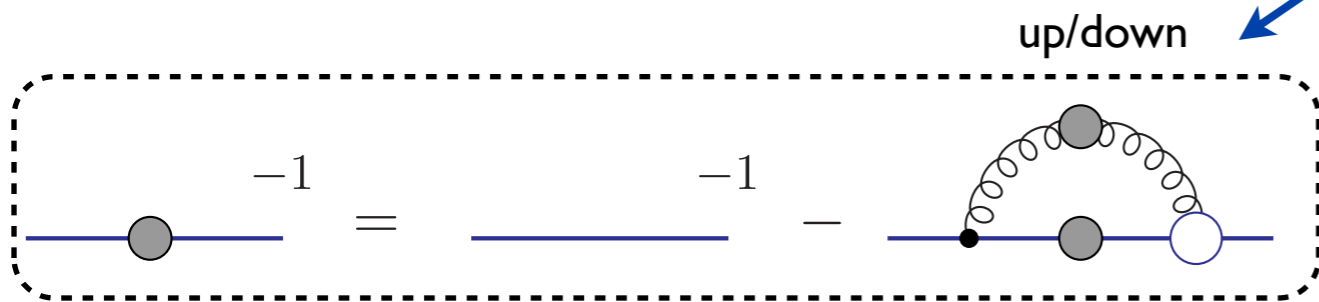
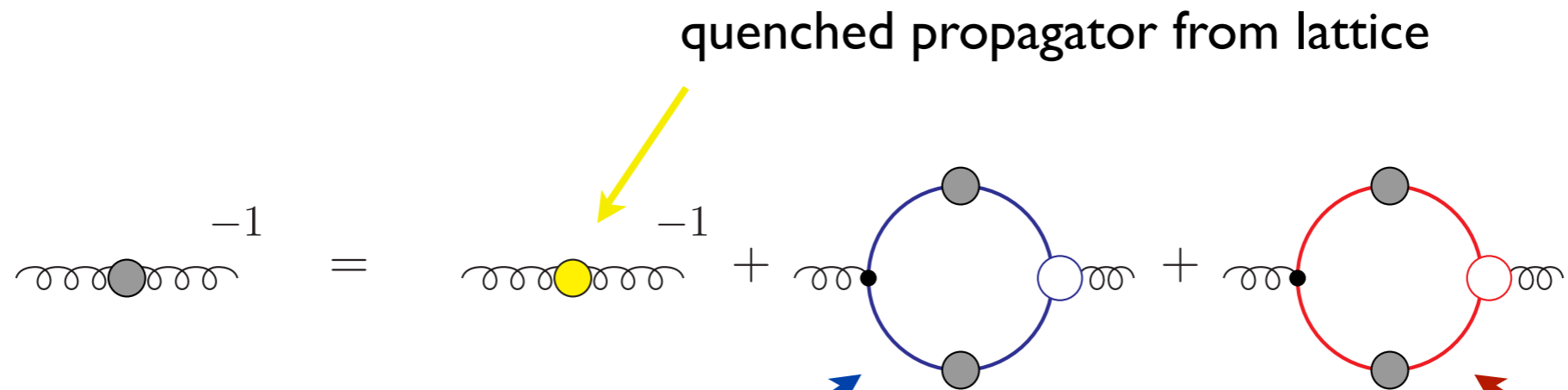
CF, Maas, Mueller, EPJC 68 (2010)

Cucchieri, Mendes, PoS FACESQCD 007 (2010)

Aouane, Bornyakov, Ilgenfritz, Mitryushkin, Muller-Preussker and Sternbeck, PRD 85 (2012) 034501

FRG: Fister, Pawlowski, arXiv:1112.5440

# $N_f=2+1$ -QCD with DSEs

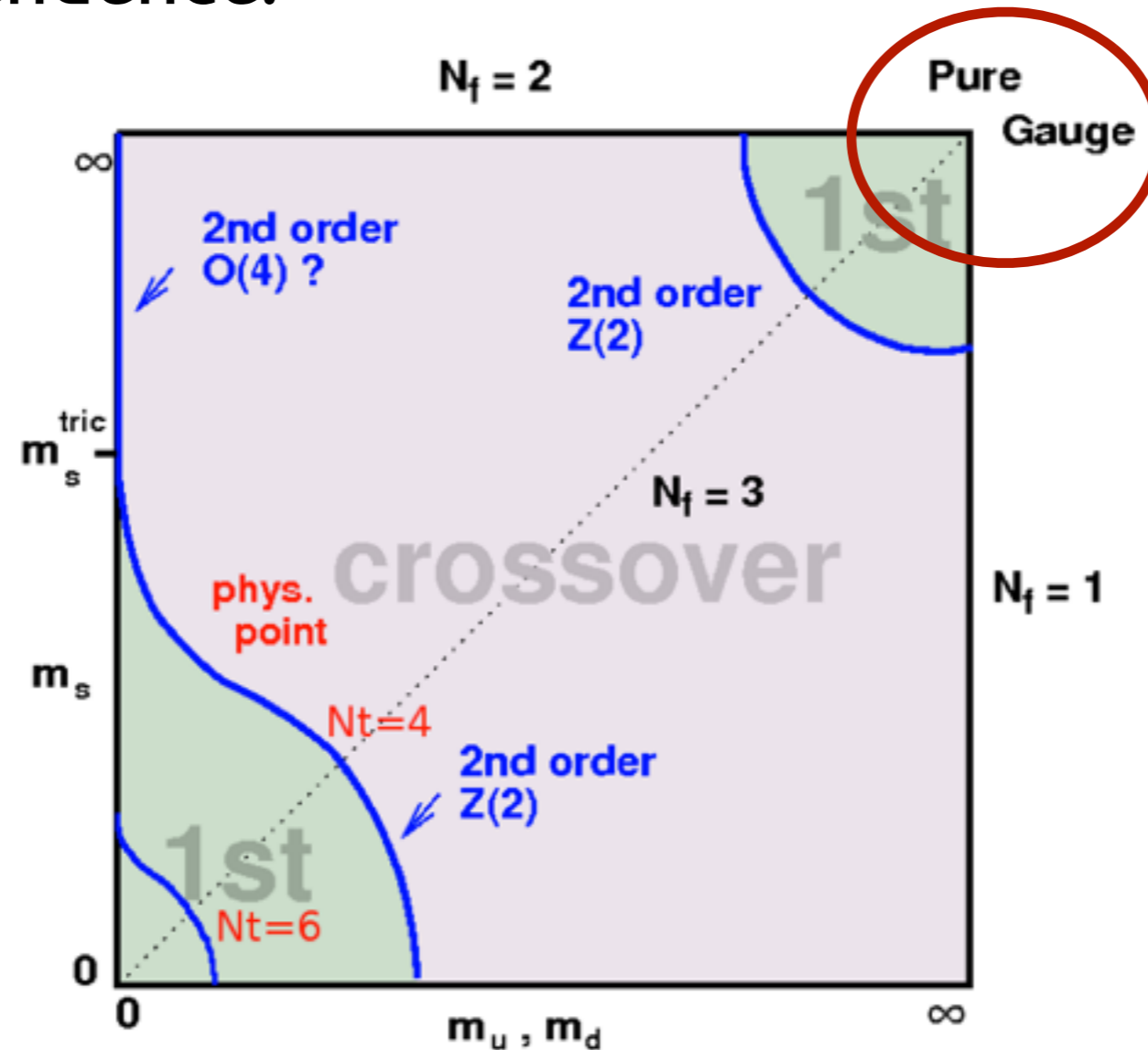


- quenched: without quark-loop
- $N_f=2$ : isospin symmetry
- $N_f=2+1$ : solve coupled system of  $2+3+3$  equations



# QCD phase transition: heavy quark limit/quenched

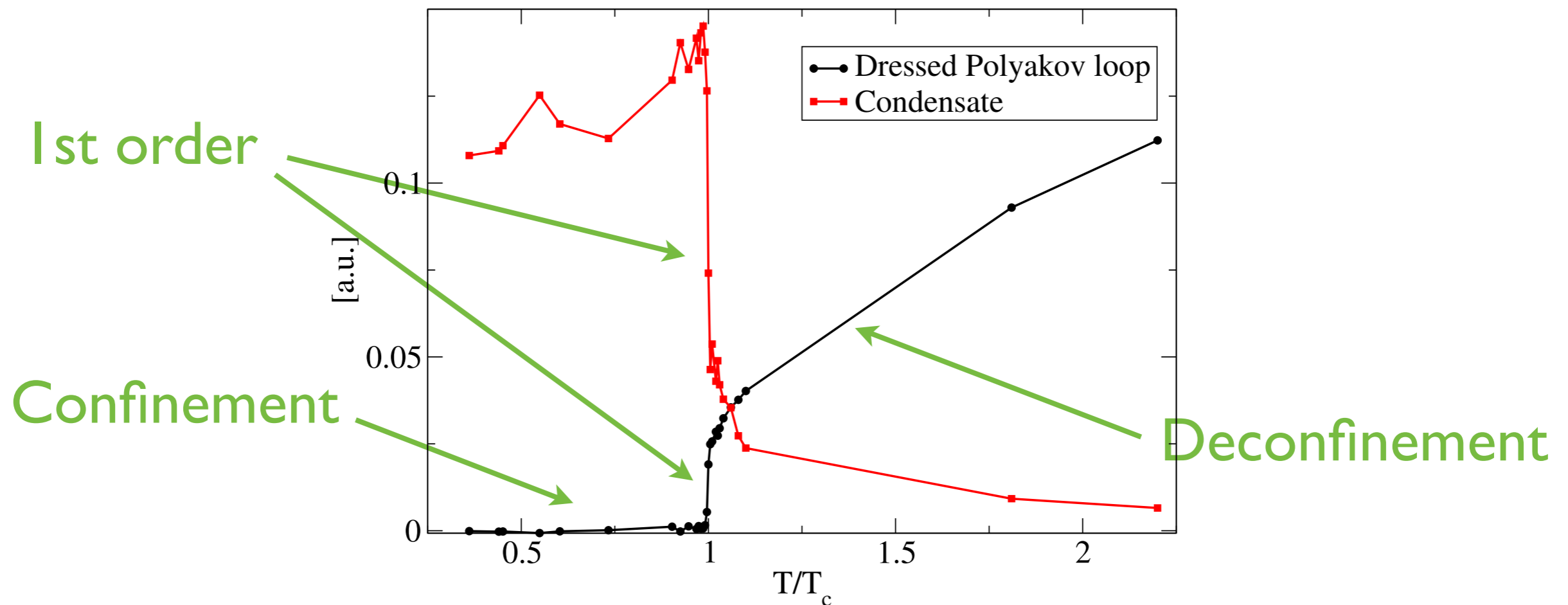
Quark mass dependence:



- Expect: Transitions controlled by deconfinement
- $SU(2)$  second order,  $SU(3)$  first order

# Transition temperatures, quenched

## quenched DSE: SU(3)

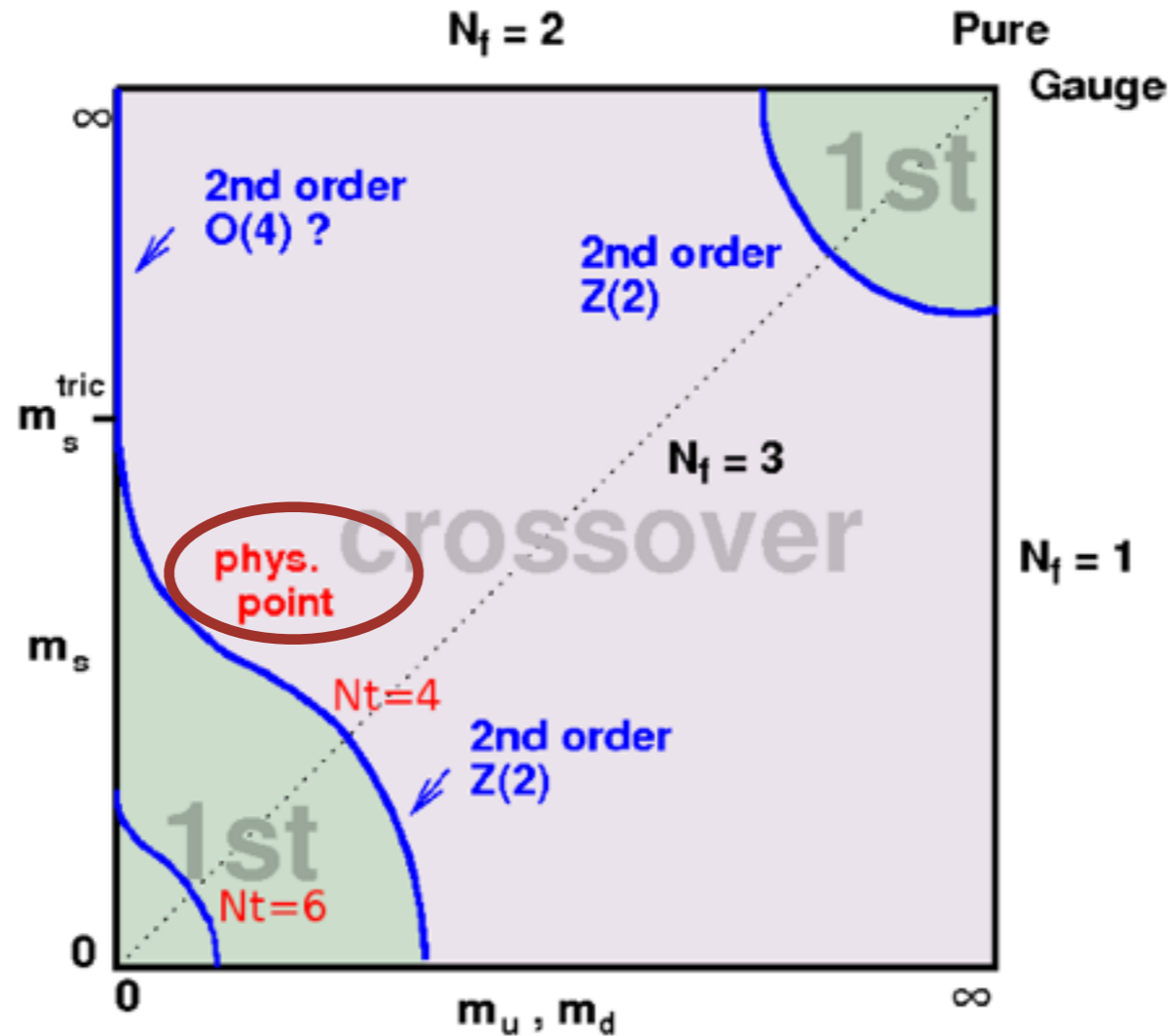


Luecker, CF, Prog.Part.Nucl.Phys. 67 (2012) 200-205  
CF, Maas, Mueller EPJC 68 (2010)

- SU(2):  $T_c \approx 305$  MeV
- SU(3):  $T_c \approx 270$  MeV
- $T \leq T_c$ : increasing condensate due to electric part of gluon

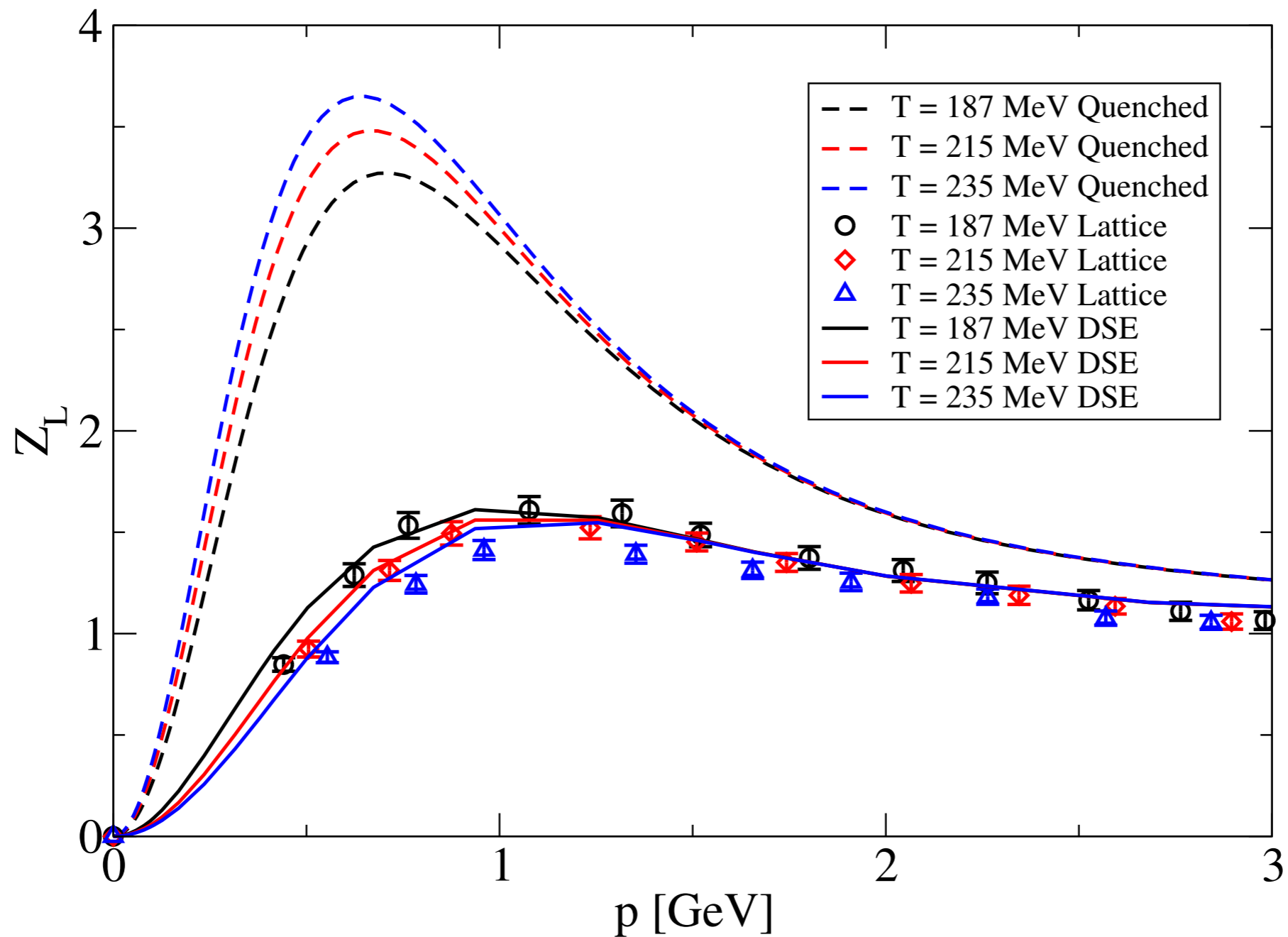
cf. Buividovich, Lushevskaya, Polikarpov, PRD 78 (2008) 074505  
cf. Braun, Gies, Pawłowski, PLB 684 (2010) 262.

# QCD phase transitions: $N_f=2+1$



- Physical up/down and strange quark masses
- Transition controlled by chiral dynamics
- at  $\mu=0$ : compare to available lattice results

# Unquenched Gluon DSE vs Lattice

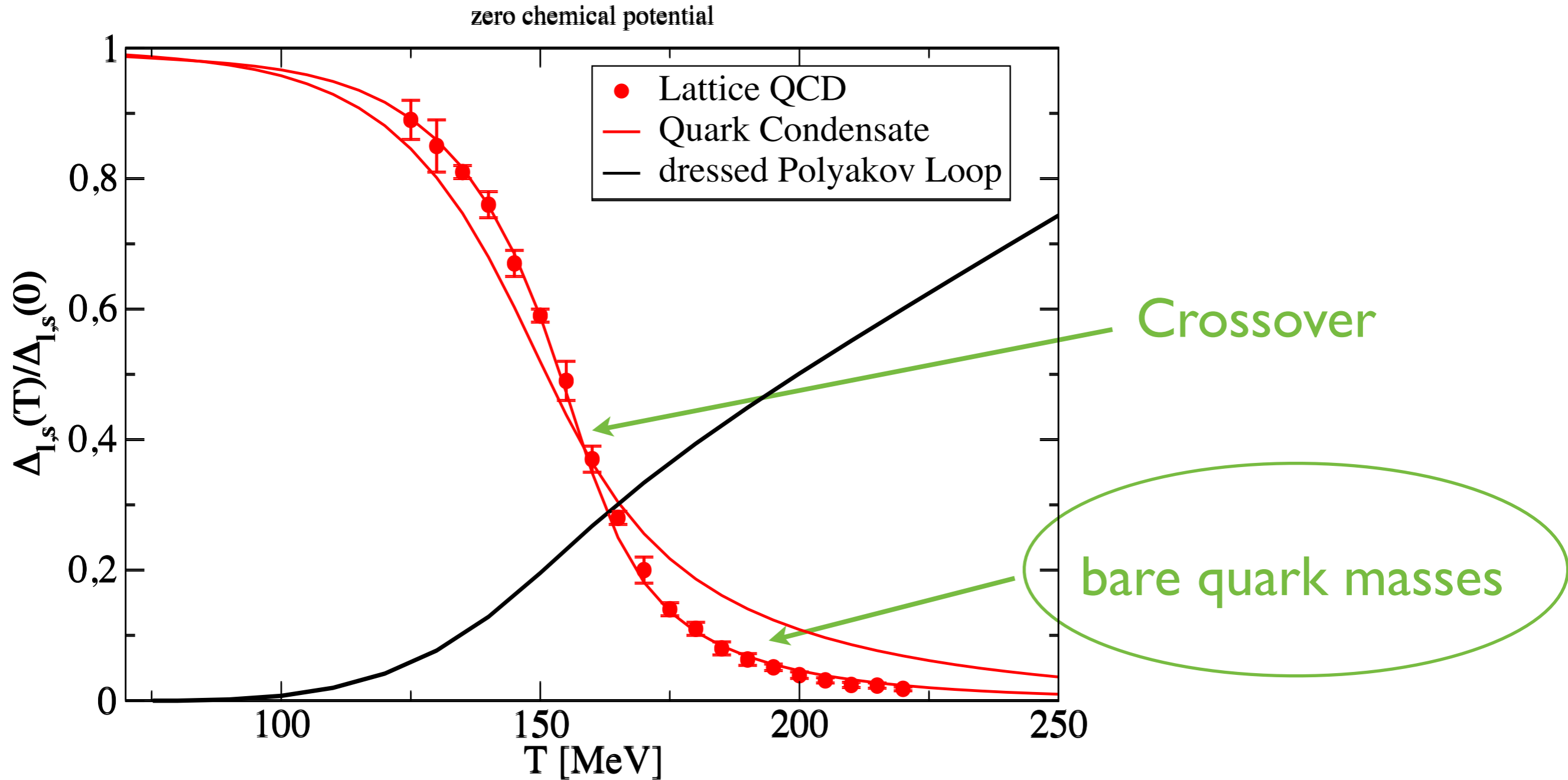


● quantitative agreement: **DSE prediction verified by lattice**

DSE: CF, Luecker, PLB 718 (2013) 1036 [arXiv:1206.5191]

Lattice: Aouane, Burger, Ilgenfritz, Muller-Preussker and Sternbeck, arXiv:1212.1102

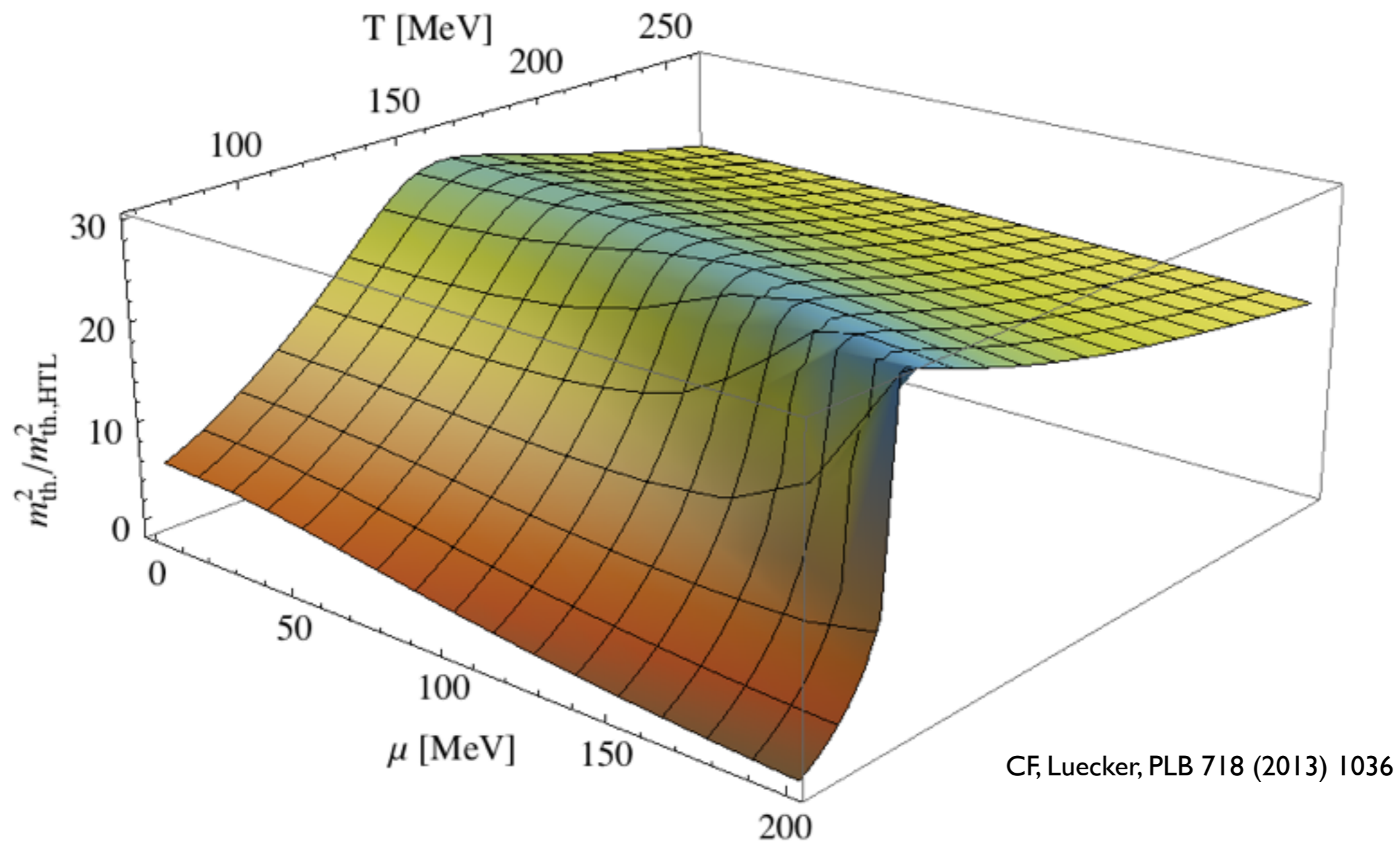
# $N_f=2+1$ , zero chemical potential



Lattice: Borsanyi *et al.* [Wuppertal-Budapest Collaboration], JHEP 1009(2010) 073  
DSE: CF, Luecker, PLB 718 (2013) 1036, CF, Luecker, Welzbacher, arXiv:1405.4762

● quantitative agreement

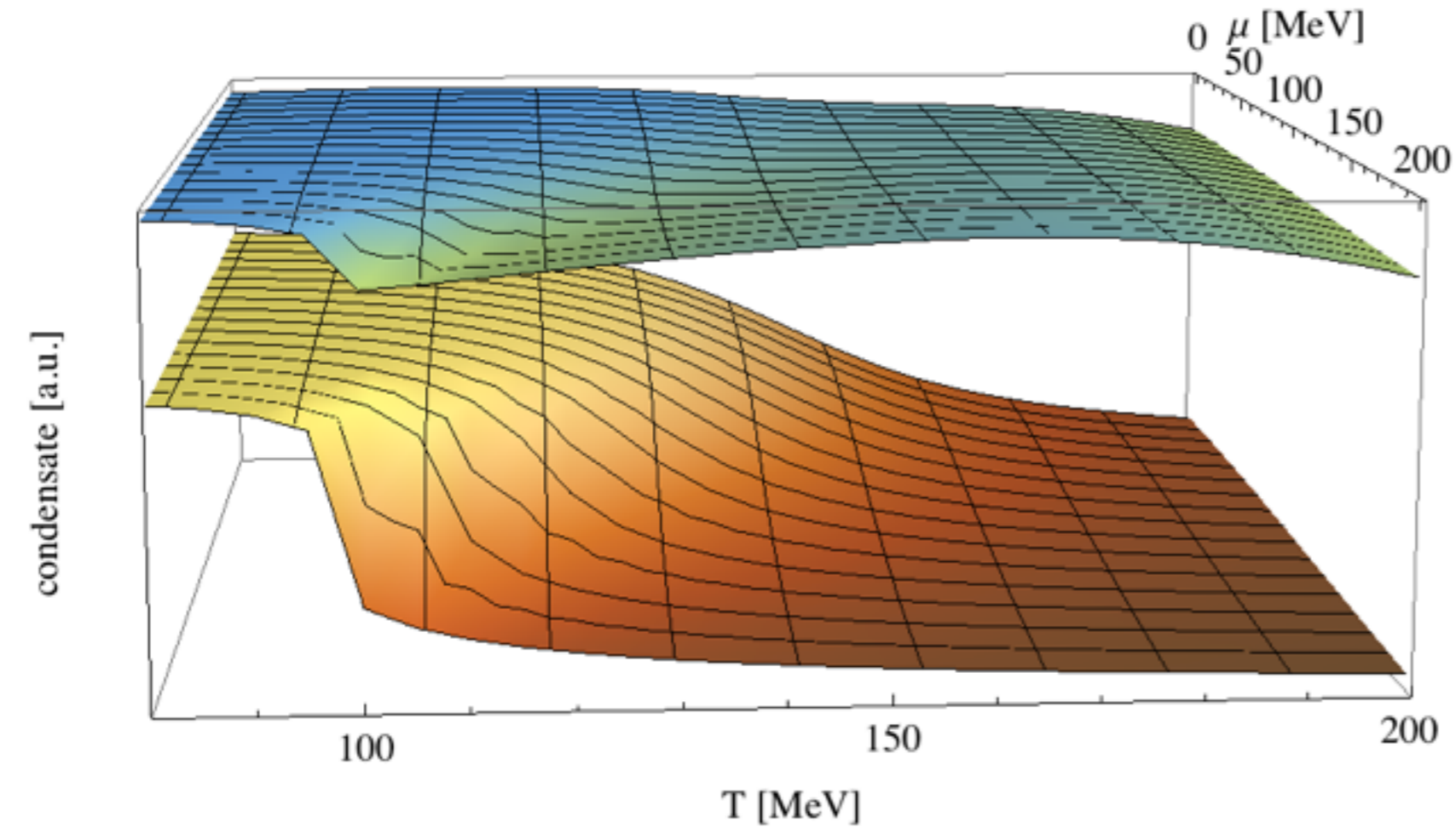
# $N_f=2+1$ : thermal electric gluon mass



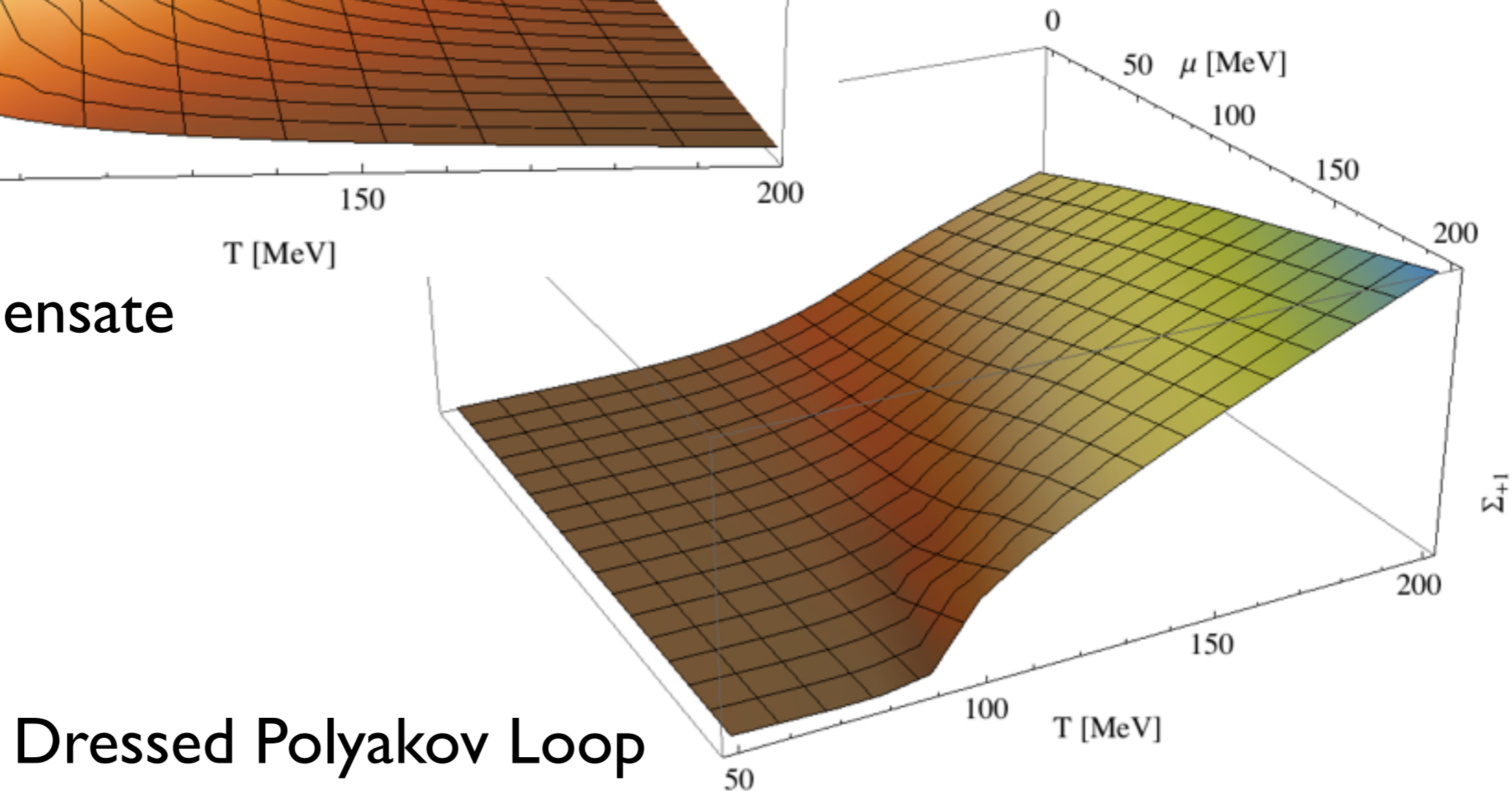
- large temperatures: behaviour as expected from HTL
- first order transition at large chemical potential



# $N_f=2+1$ : Condensate and dressed Polyakov Loop



Quark condensate



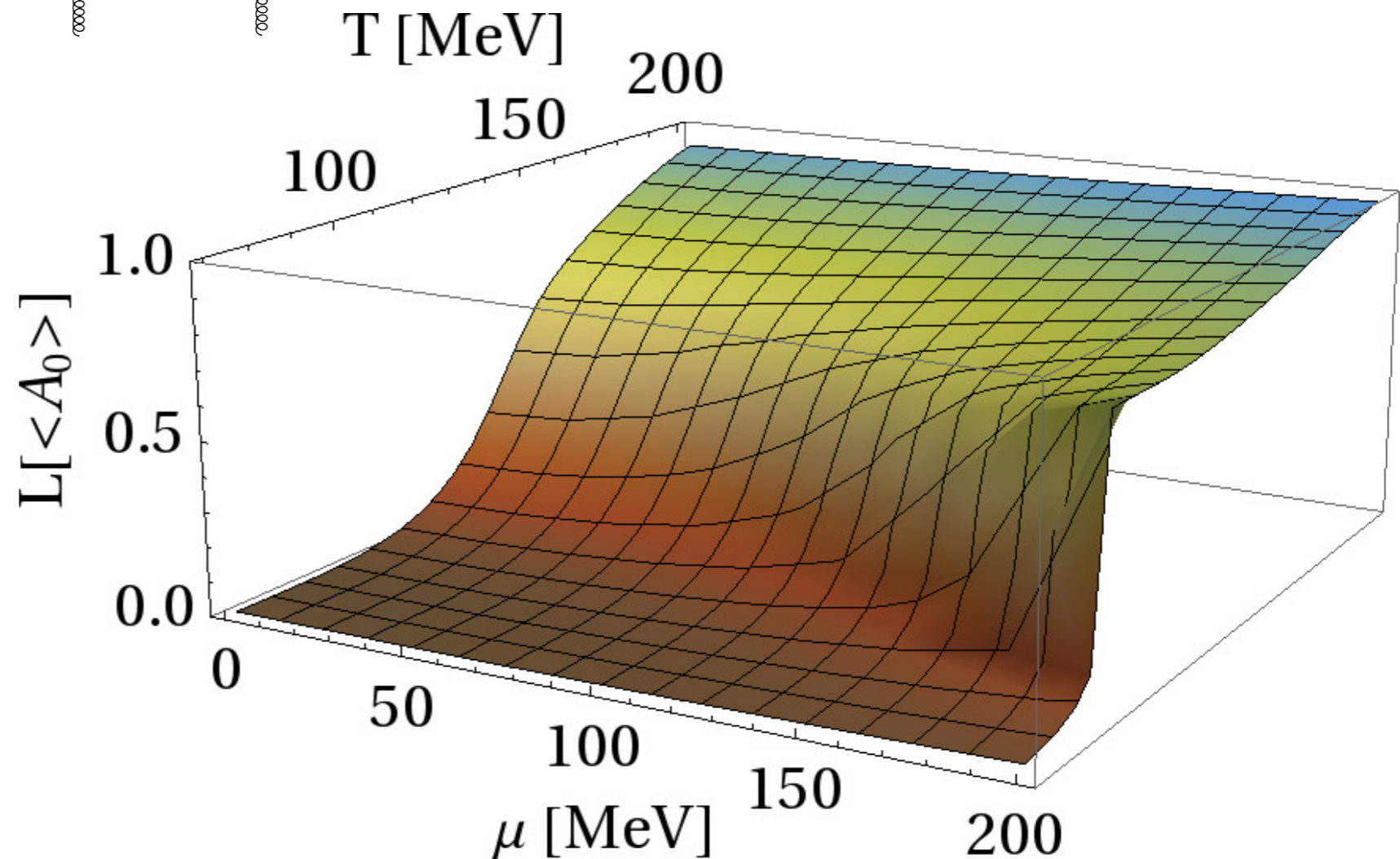
Dressed Polyakov Loop

# $N_f=2+1$ : Polyakov loop potential at finite $\mu$

$$\frac{\delta(\Gamma - S)}{\delta A_0} = \frac{1}{2} \left[ \text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} - \frac{1}{6} \left( \text{Diagram 4} + \text{Diagram 5} \right) \right]$$

## Polyakov-Loop

$$L = \frac{1}{N_c} \text{tr} e^{ig \int A_0}$$



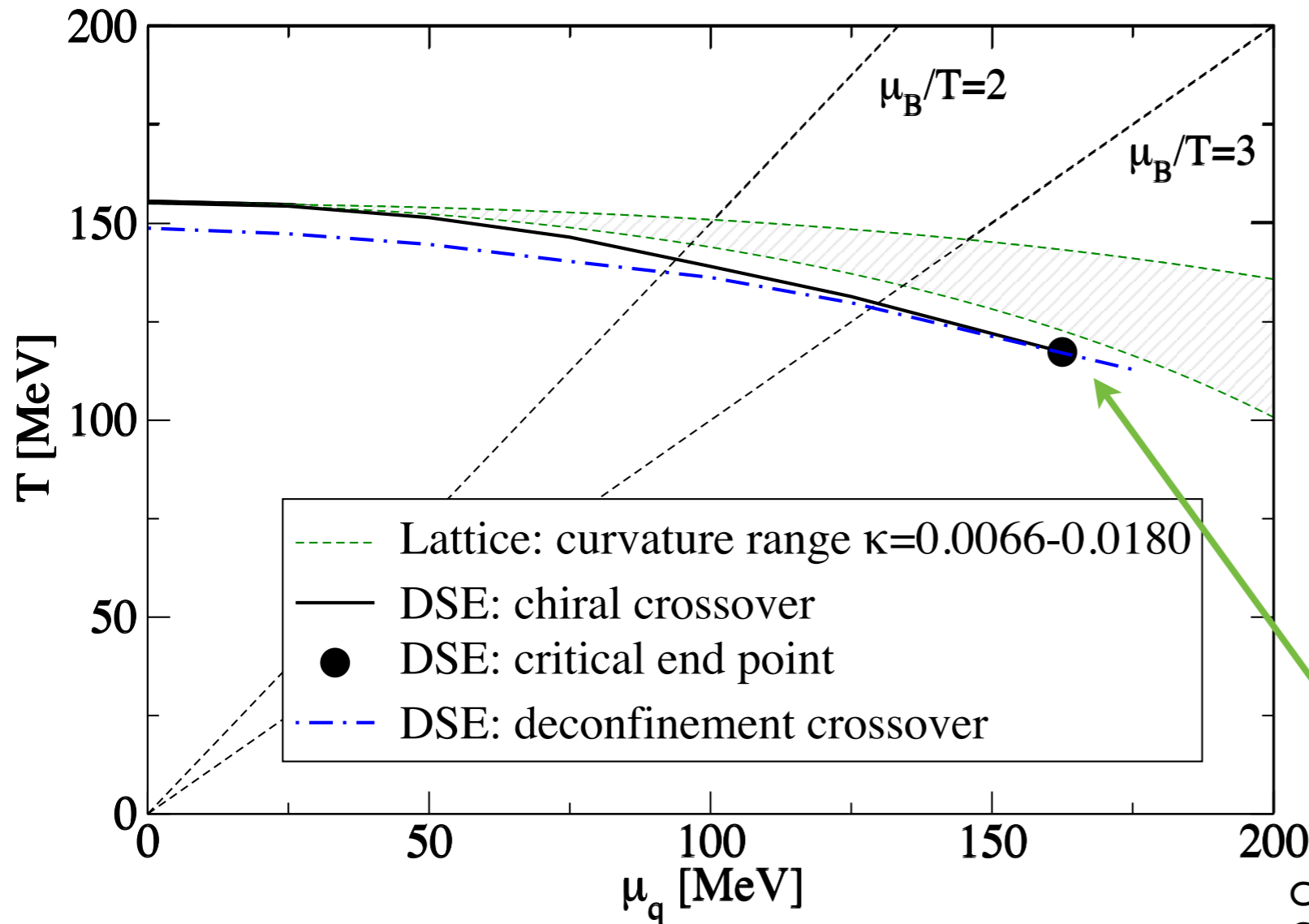
CF, Fister, Luecker, Pawłowski, PLB in press, arXiv:1306.6022

- evaluated from Polyakov-Loop potential
- important input for P-models: PQM, PNJL !

Herbst, Mitter, Pawłowski, Schaefer, Stiele, arXiv:1308.3621



# $N_f=2+1$ : Polyakov loop and phase diagram



Extrapolated curvature from lattice

Kaczmarek et al. PRD 83 (2011) 014504,  
Endrodi, Fodor, Katz, Szabo, JHEP 1104 (2011) 001  
Cea, Cosmai, Papa, PRD 89 (2014) 074512

CEP at large  $\mu$

CF, Luecker, PLB 718 (2013) 1036,  
CF, Fister, Luecker, Pawlowski, PLB in press, arXiv:1306.6022  
CF, Luecker, Welzbacher, arXiv:1405.4762

● no CEP at  $\mu_B/T < 2$  in agreement with lattice and FRG

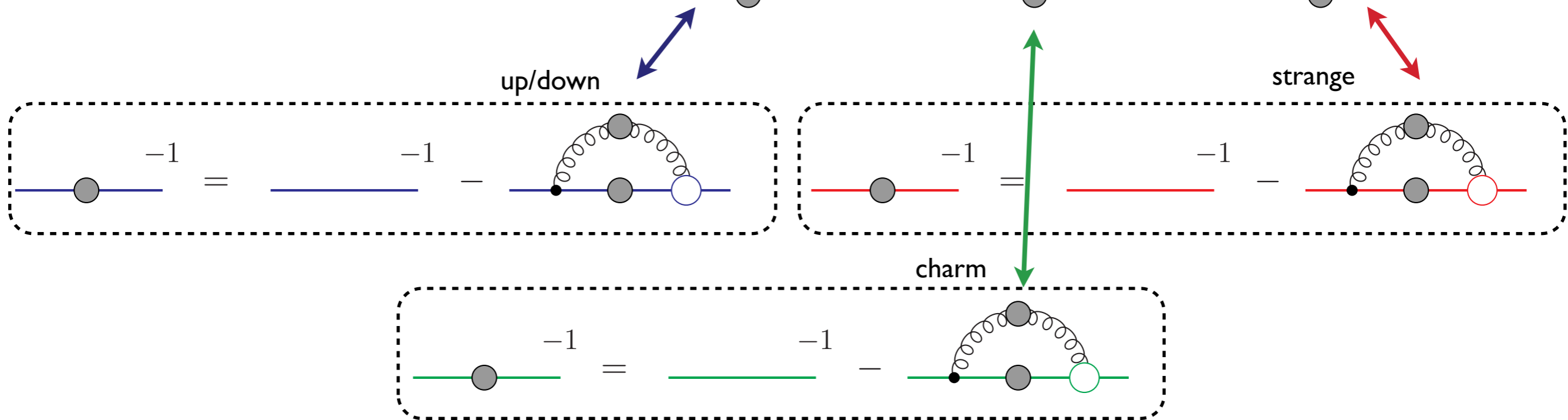
de Forcrand, Philipsen, JHEP 0811 (2008) 012; Nucl Phys. B642 (2002) 290-306  
Endrodi, Fodor, Katz, Szabo, JHEP 1104 (2011) 001  
Herbst, Pawlowski, Schaefer, PRD 88 (2013) 014007

Caveat: baryon effects missing...

$N_c=2$ : Brauner, Fukushima and Hidaka, PRD 80 (2009) 74035  
Strodthoff, Schaefer and Smekal, PRD 85 (2012) 074007

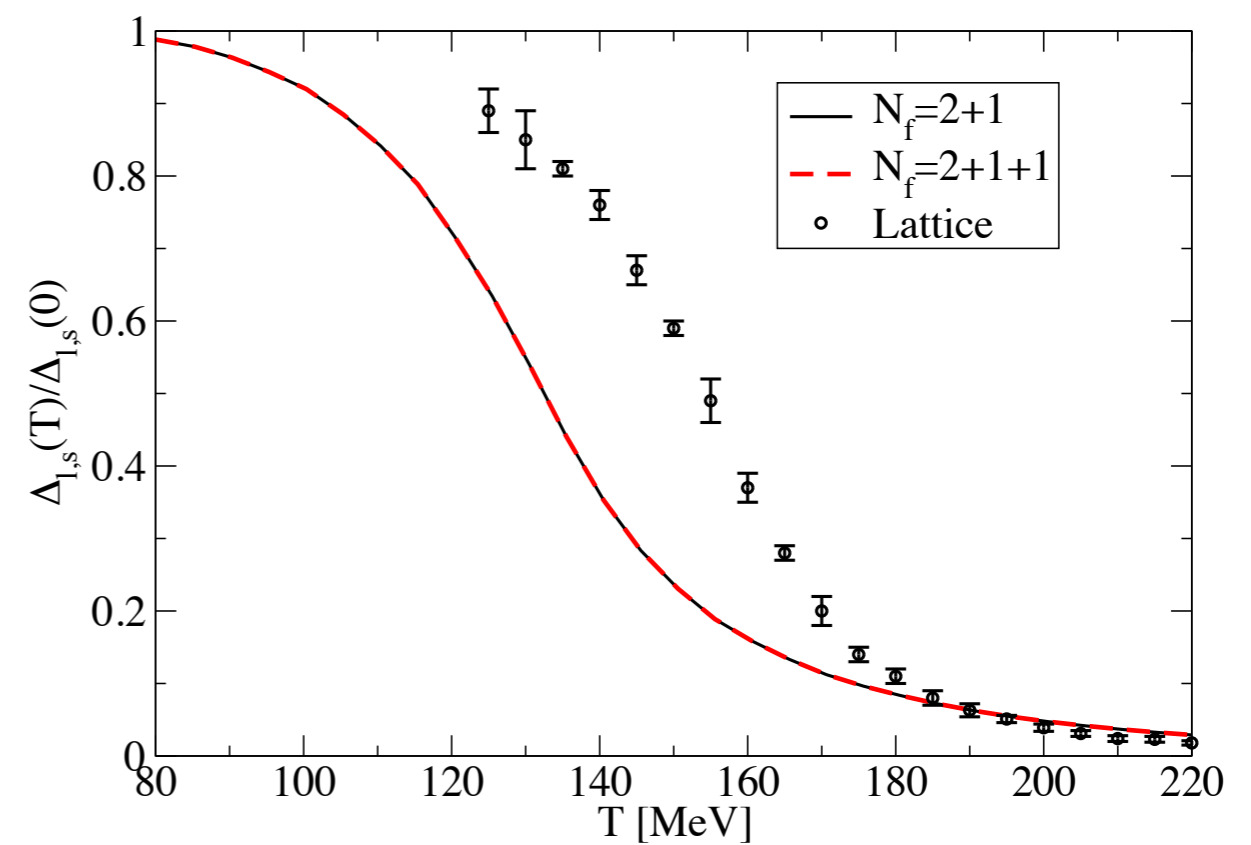
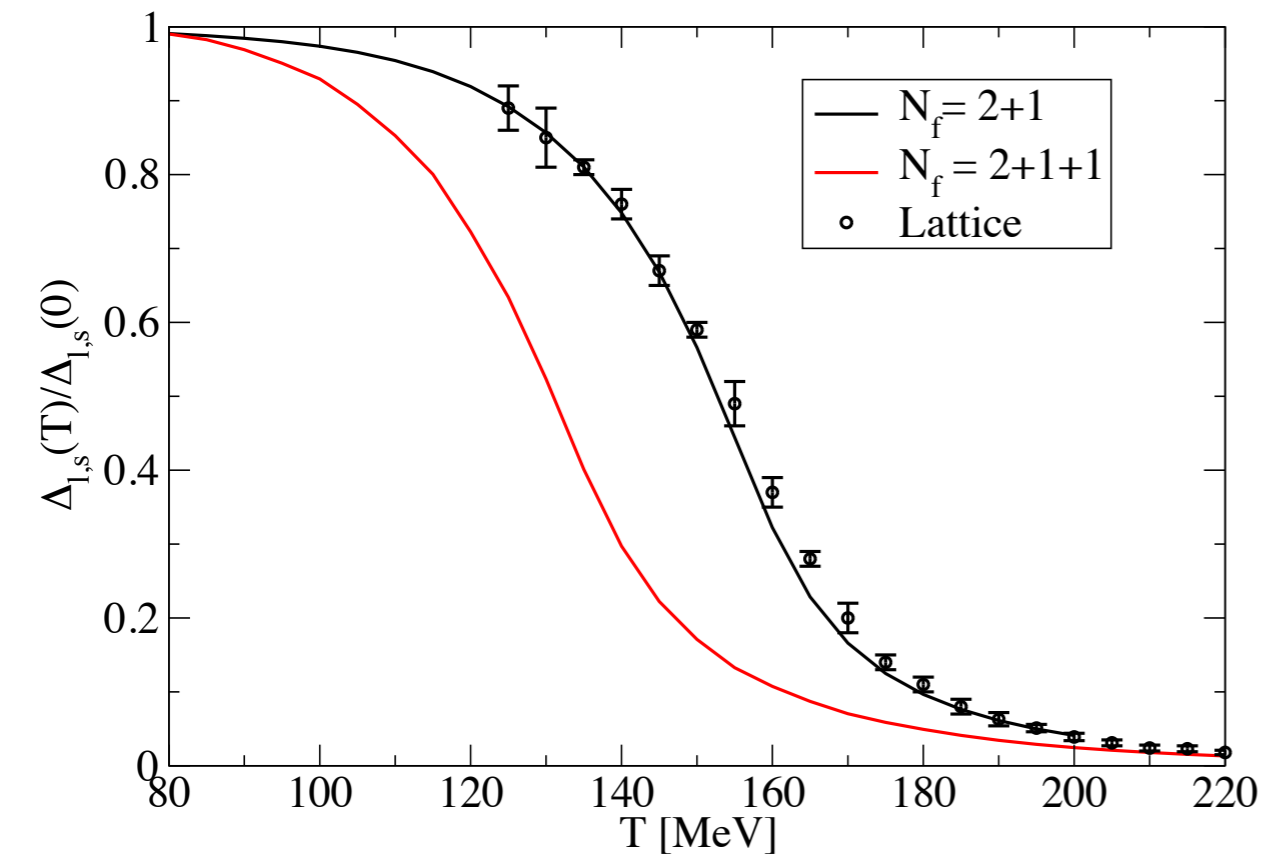
# $N_f=2+1+1$ -QCD with DSEs

$$\begin{array}{c}
 \text{gluon} \\
 \text{---} \bullet \text{---} \\
 -1
 \end{array}
 =
 \begin{array}{c}
 \text{gluon} \\
 \text{---} \bullet \text{---} \\
 -1
 \end{array}
 +
 \begin{array}{c}
 \text{quark loop} \\
 \text{---} \bullet \text{---} \\
 -1
 \end{array}
 +
 \begin{array}{c}
 \text{quark loop} \\
 \text{---} \bullet \text{---} \\
 -1
 \end{array}
 +
 \begin{array}{c}
 \text{quark loop} \\
 \text{---} \bullet \text{---} \\
 -1
 \end{array}$$



- Physical up/down, strange and **charm quark masses**
- Transition controlled by chiral dynamics
- *no lattice or model results available yet*

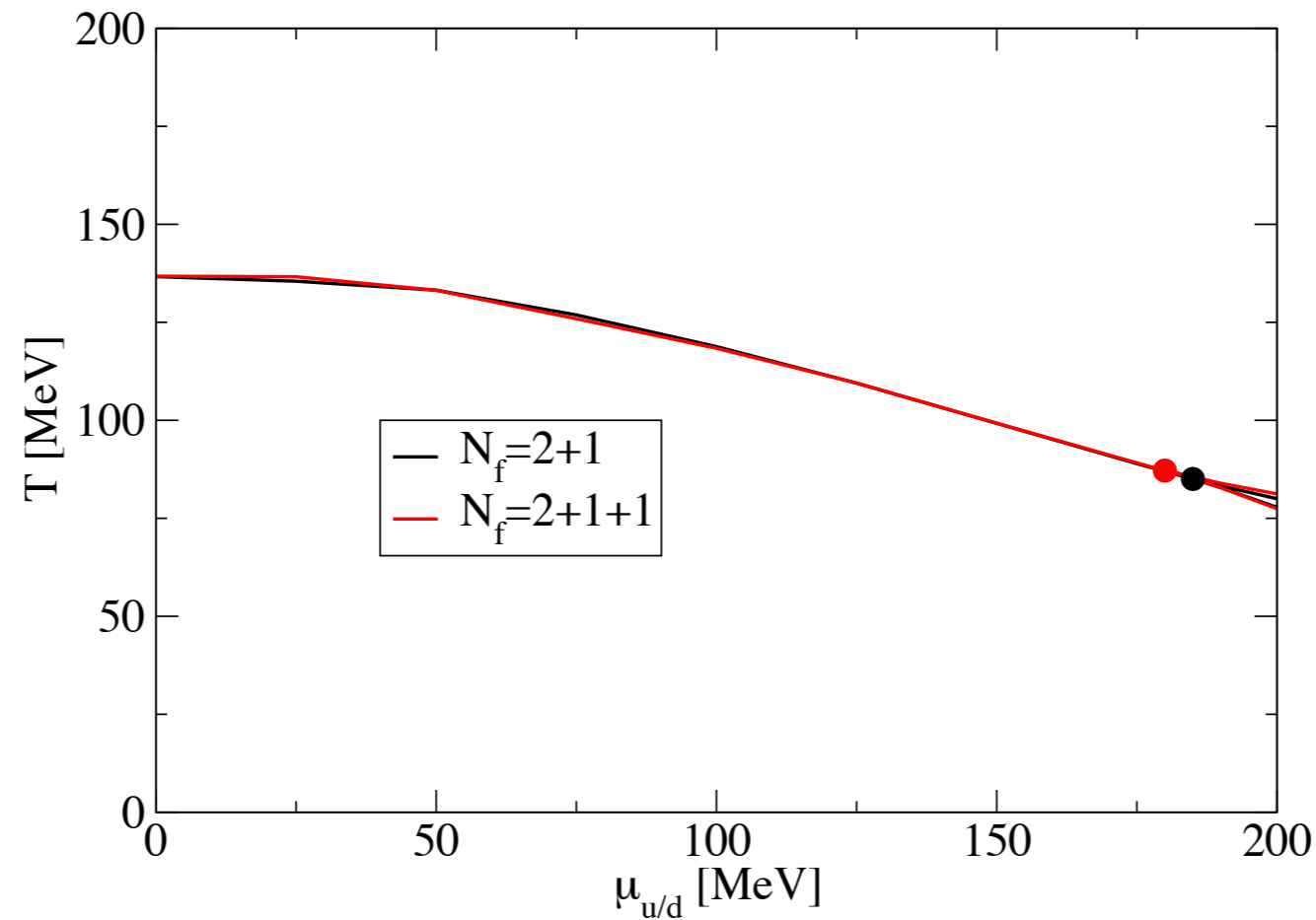
# $N_f=2+1+1$ -QCD with DSEs



CF, Luecker, Welzbacher, arXiv:1405.4762

- Left: Interaction fixed:  $T_{PC}$  decreases by  $O(20 \text{ MeV})$
- Right: **Physics fixed ( $m_\pi, f_\pi$ ):  $T_{PC}$  similar**

# $N_f=2+1+1$ -QCD with DSEs



CF, Luecker, Welzbacher, arXiv:1405.4762

- Physics fixed ( $m_\pi, f_\pi$ ):  $T_C$  similar
- Charm quark has no influence on QCD phase diagram

# Summary

- Temperature dependent gluon propagator
  - characteristic behaviour of electric gluon
  - ‘melting’ of magnetic gluon with temperature
- Deconf.  $T_{pc}$  from dressed Polyakov-loop/Polyakov-loop potential
- QCD with finite chemical potential (beyond mean field)
  - backreaction of quarks onto gluons important
  - $N_f=2+1$  and  $N_f=2+1+1$ : CEP at  $\mu_c/T_c > 1$

Work in progress: include baryons...  
include magnetic field...

Mueller, Bonnet, CF, PRD in press, arXiv:1401.1647