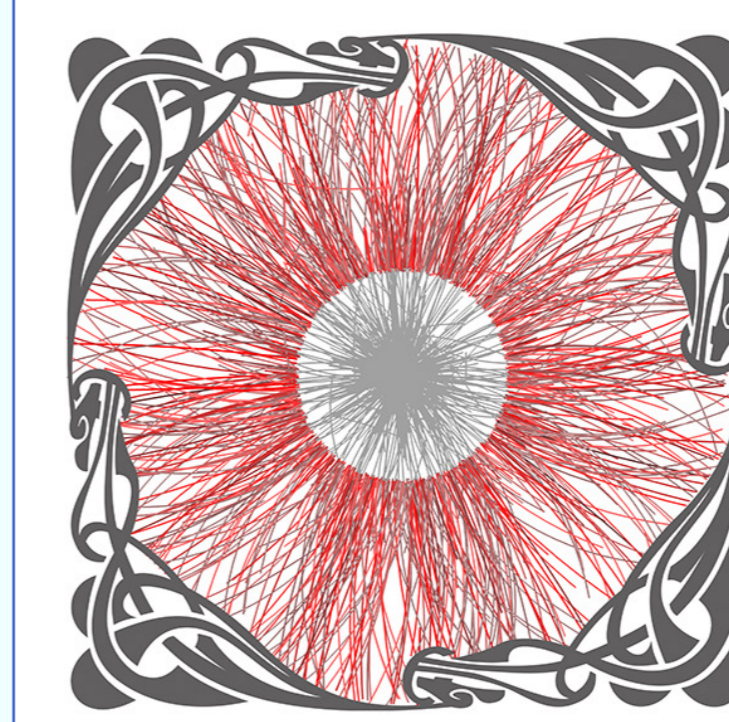
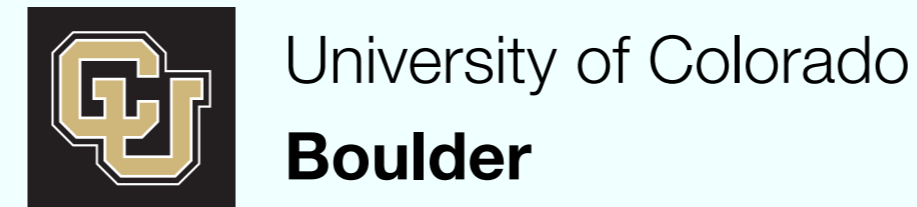


Inferring nuclear modification of charm and beauty hadrons from simulated electron data

Unfolding non-photonic electron observables

Andrew M. Adare

University of Colorado
for the PHENIX collaboration



XXIV
QUARK
MATTER
DARMSTADT
2014

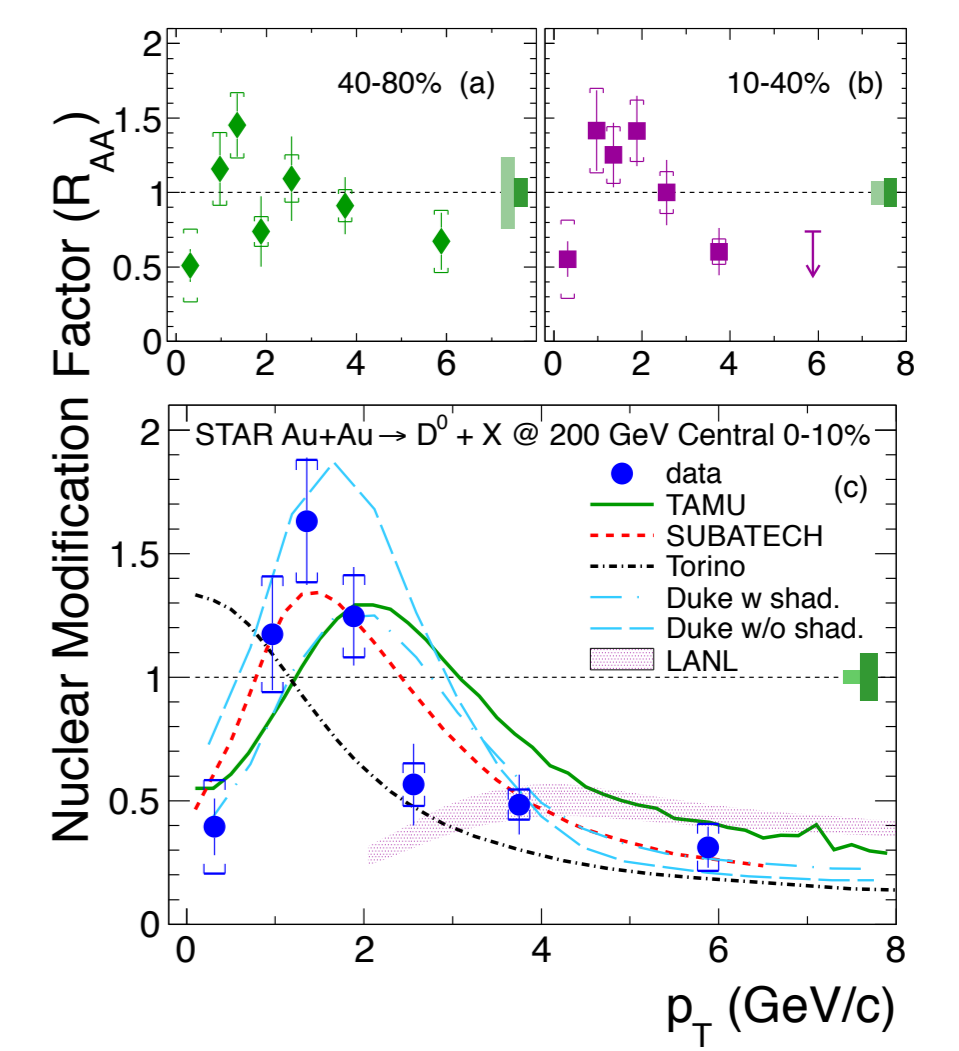
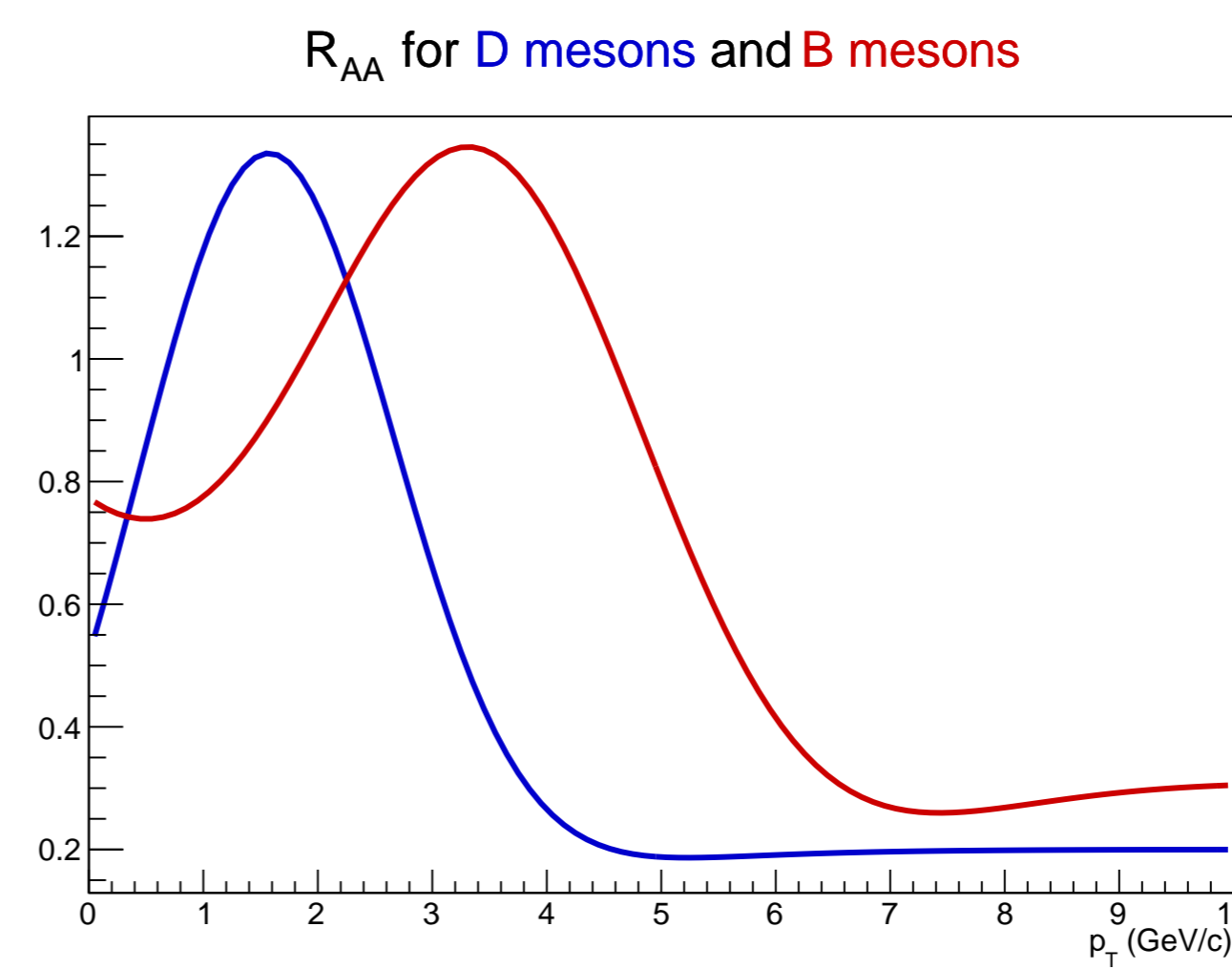
Nuclear modification of heavy-flavor hadrons

The goal of this poster is to demonstrate that it is possible to infer the modification of heavy-flavor hadrons from electron observables, namely (a) the displaced vertices of tracks identified as electrons and (b) electron invariant yields vs p_T .

In this exercise, heavy-flavor hadron p_T spectra are generated by pythia, then modified by a blast-wave inspired estimate of R_{AA} [PLB 557 (2003) 26-32] (near right). Their decay electrons are used as the unfolding input "data".

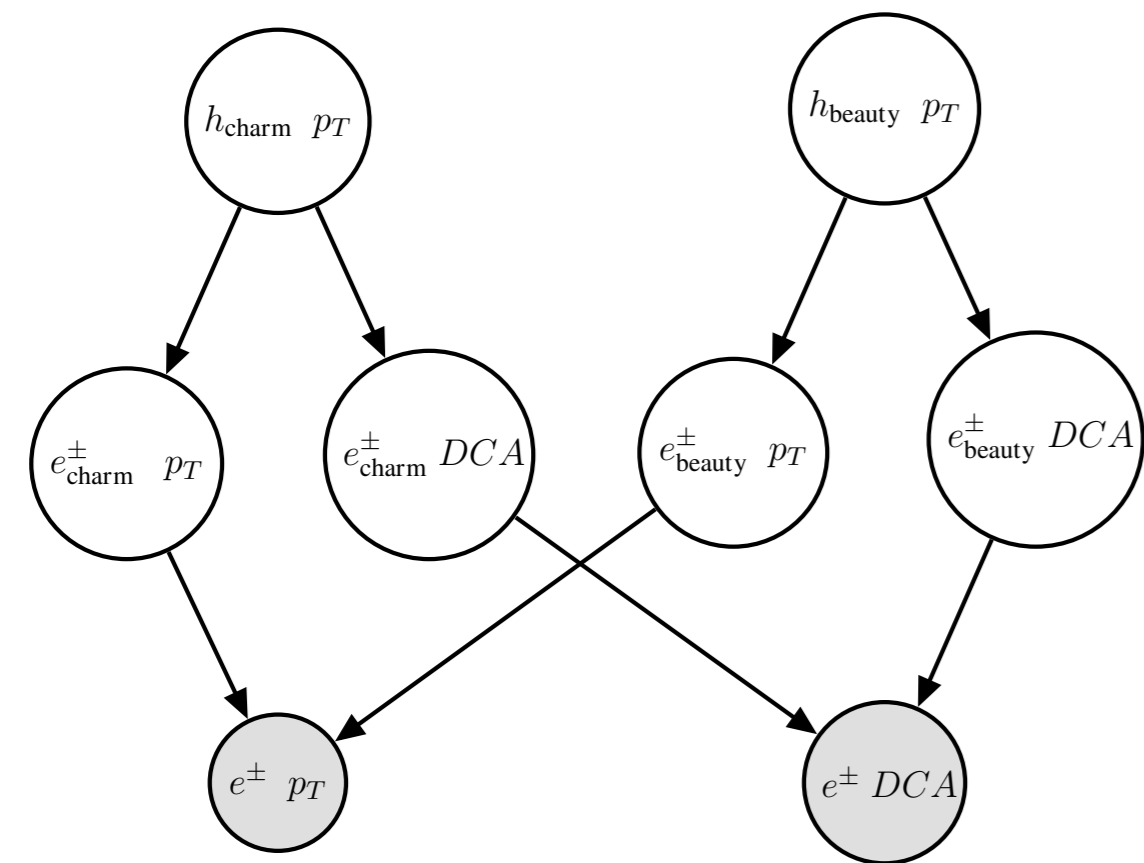
This form for the modification has gained support from a recent result by the STAR collaboration [1404.6185] (far right).

The PYTHIA generator also provides the heavy-flavor decay kinematics, represented as matrices of decay probabilities. This information, along with the (modified) electron data, forms a linear system that can be solved using Bayesian inference.



A generative model for heavy-flavor decays

Each node in this directed graph represents a probability distribution. The arrows represent this directed conditional dependence.



The objective is to infer the distributions over the latent variables $h_{\text{charm}} p_T$ and $h_{\text{beauty}} p_T$ from the (shaded) electron observables.

Bayesian unfolding using displaced vertices and spectra

See 1201.4612v4 by G. Choudalakis. His terminology is used here:

- T** truth vector (length N_t). $\tilde{\mathbf{T}}$ is the modeled truth (e.g. from MC).
- R** reconstructed N_r -vector (again, e.g. in MC).
- D** measured data.

Bayes' theorem says $p(\mathbf{T}|\mathbf{D}) \propto L(\mathbf{D}|\mathbf{T}) \cdot \pi(\mathbf{T})$. In words: the posterior probability \propto the likelihood \times the prior probability.

The problem amounts to assuming $\pi(\mathbf{T})$ and computing $L(\mathbf{D}|\mathbf{T})$:

$$L(\mathbf{D}|\mathbf{T}) = \prod_{r=1}^{N_r} \frac{D_r^{D_r}}{D_r!} e^{-R_r} \quad (\text{for Poisson data}) \quad (1)$$

The result is **not** a spectrum of points with covariance. Instead, an N_r -dimensional posterior probability is obtained.

In each t bin, a 1D posterior is marginalized (integrated) from $p(\mathbf{T}|\mathbf{D})$.

Bayesian unfolding: implementation

First, the mapping of $\mathbf{T} \rightarrow \mathbf{R}$ is established, yielding a truth spectrum $\tilde{\mathbf{T}}$ and a matrix M containing $P(r|t)$ values. $P(r|t)$ is the probability for an object from bin t to be reconstructed in bin r .

Also, select $\pi(\mathbf{T})$. Non-constant prior \Rightarrow bias. Regularization!

Then:

- 1 A trial \mathbf{T} point is pulled from an N_t -dimensional sampling volume
- 2 $\mathbf{R} = M\mathbf{T}$
- 3 $\pi(\mathbf{T})L(\mathbf{D}|\mathbf{T})$ (or, in practice, the log) is computed from $\pi(\mathbf{T})$, \mathbf{R} and \mathbf{D} (eq. 1)
- 4 \mathbf{T} and $L(\mathbf{D}|\mathbf{T})$ are stored (TTree)
- 5 Repeat 1-4 until $p(\mathbf{T}|\mathbf{D})$ is well sampled
- 6 Marginalize: project $p(\mathbf{T}|\mathbf{D})$ to 1D posteriors $p_t(T_t|\mathbf{D})$

The trickiest part is step 1.

Sampling $L(\mathbf{D}|\mathbf{T})\pi(\mathbf{T})$ in N_t dimensions

The initial sampling volume must be large enough to enclose the "answer" conservatively.

On the other hand, the hyper-volume grows enormously as the boundaries are expanded. Grid sampling and uniform MC can quickly become prohibitive.

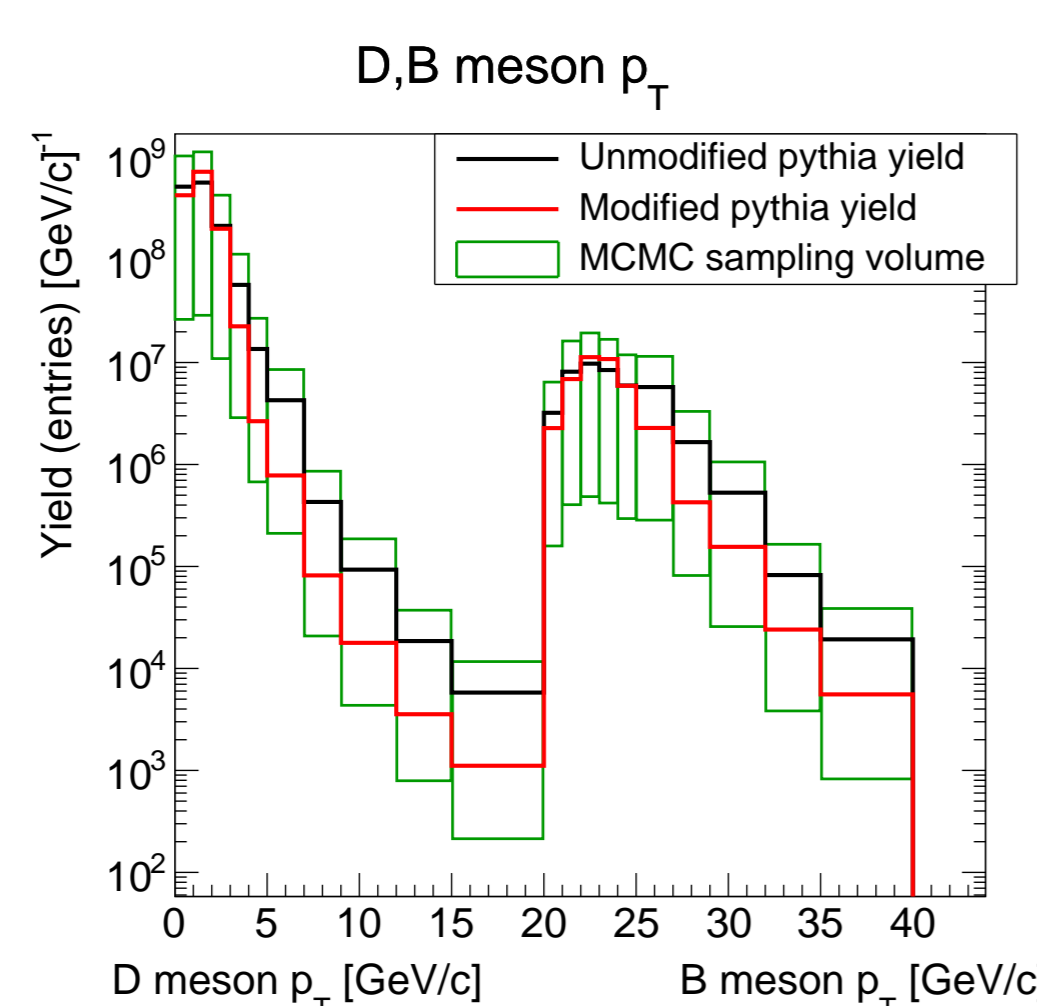
The solution is **Markov chain Monte Carlo** (MCMC). It is **ergodic** (it visits the whole space).

Moreover, it samples in **direct proportion** to $p(\mathbf{T}|\mathbf{D})$!

A **sketch** of the Metropolis-Hastings algorithm with a uniform $\pi(\mathbf{T})$:

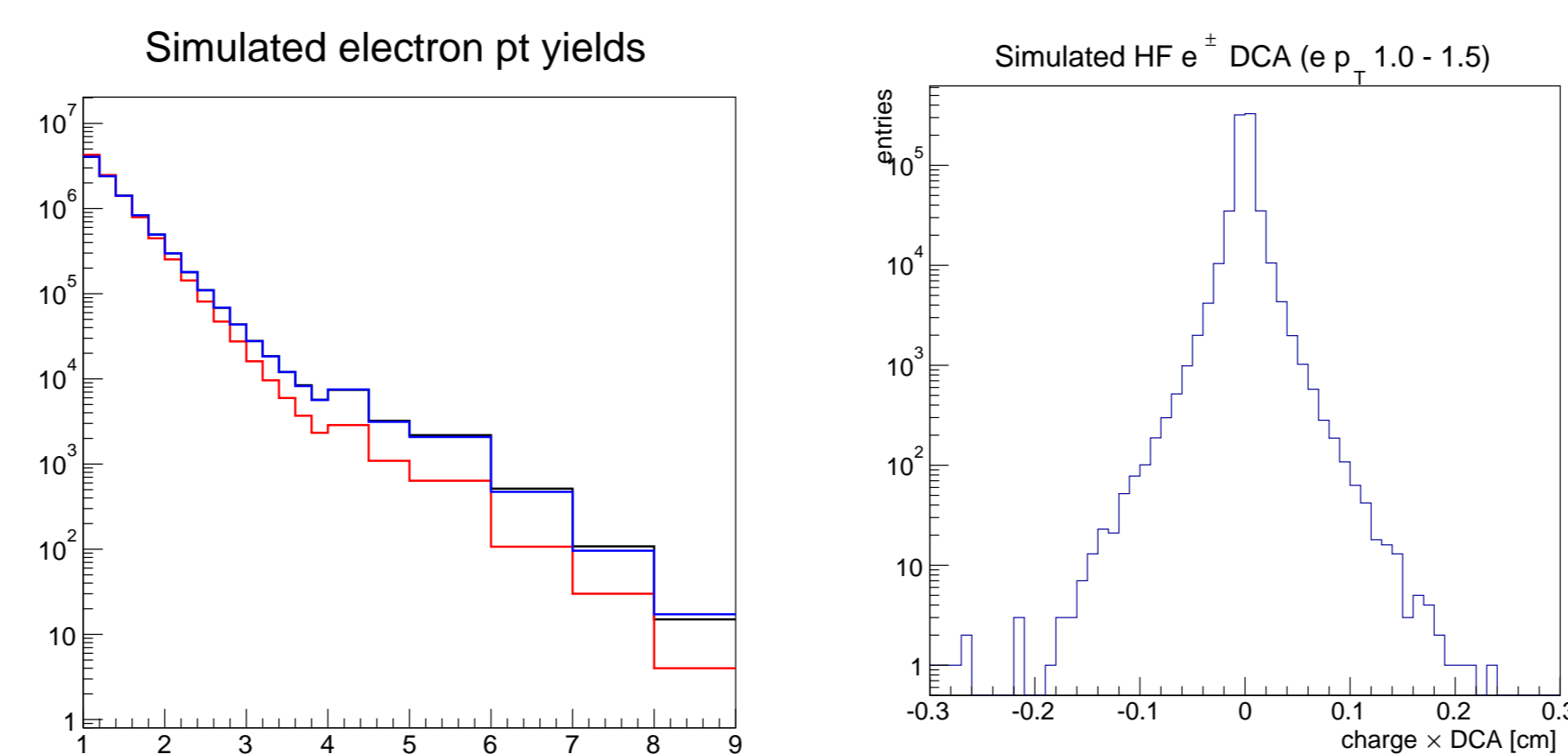
- Start with $\mathbf{T}_0 = \tilde{\mathbf{T}}$, and pick a large "hyperbox" around it. Save L_0 .
- Propose a new point \mathbf{T}_1 near \mathbf{T}_0 . Compute L_1 .
- Is it better ($L_1 > L_0$)? Keep it, update $\mathbf{T}_0 \leftarrow \mathbf{T}_1$, and resample.
- If not, \mathbf{T}_1 gets a second chance. Roll the dice again and accept it with a probability L_1/L_0 .
- Repeat...

After equilibration, a Markov chain has randomly toured the whole box, but has climbed to the highest-likelihood regions most often.



Simulated PYTHIA data

The simulation setup includes 10 million electrons sampled from a distribution obtained from the PYTHIA generator. The samples are independent of those used to model the heavy-flavor decays (see below).



Blue/black: unmodified; Red: $\times R_{AA}$

Combining datasets (Electron spectra + DCA)

This problem is a simultaneous unfolding of displaced vertices and electron spectra. This involves computing the joint likelihood for each monte carlo sample and comparing it to data.

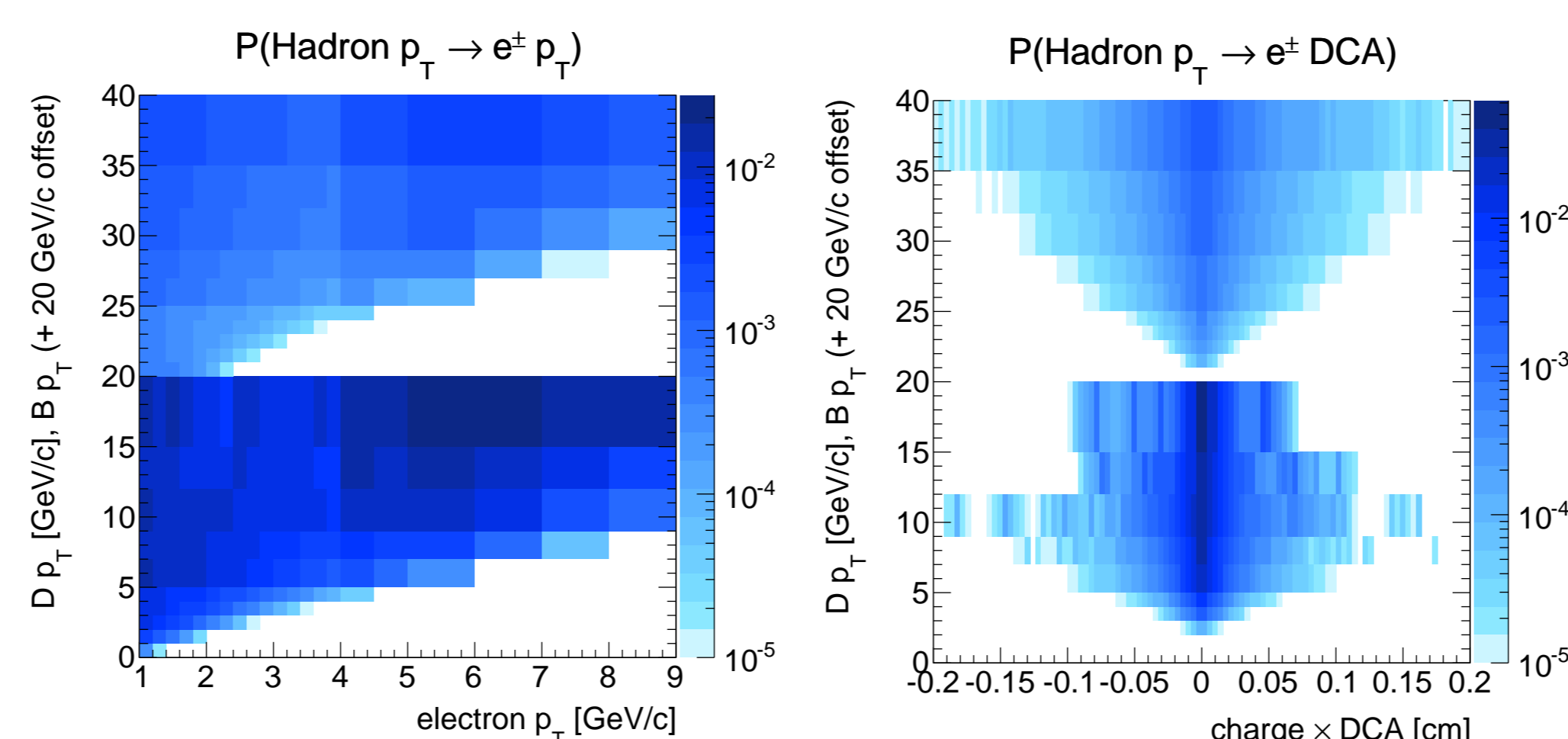
$$L(n|\mu) = \prod_{i=1}^6 w_i \prod_{j=1}^{N_{DCA}} \text{Pois}(n_{i,j}^{DCA} | \mu_{i,j}^{DCA}) \times \prod_{k=1}^{N_{e^{\pm} p_T}} w_k \text{Pois}(n_k^{e^{\pm} p_T} | \mu_k^{e^{\pm} p_T}) \quad (2)$$

Where $w_i = 1/12$ and $w_k = 1/2$.

If the efficiencies of the DCA samples vs. $e^{\pm} p_T$ are unavailable, each trial μ can be scaled to match the integral of the DCA distribution. This removes the dependence on $\|\tilde{\mu}\|_2$, and only the shape of the DCA "guess" matters.

D,B $p_T \rightarrow e^{\pm} p_T$ and DCA matrices

These matrices represent the probability for a heavy-flavor hadron at a given p_T to decay to an electron at a given p_T and/or DCA.



x axis: $e^{\pm} p_T$ bin

y axis: $D p_T$ (0-10) and $B p_T$ (10-20). $D, B p_T$ both 0-10 GeV/c

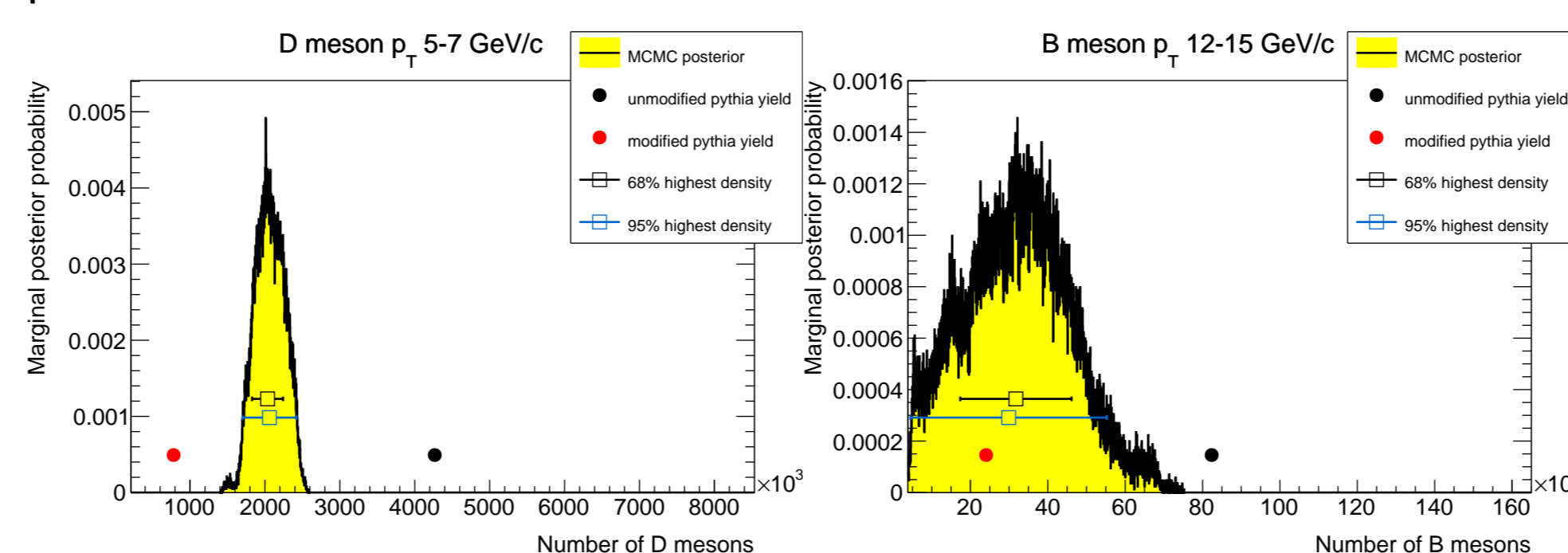
Regularization

Due to statistical fluctuations in the data and ill-conditioned transfer matrices, unbiased unfolding results typically exhibit large variances. The problem grows with model complexity (i.e. number of free parameters).

To deal with this, a prior distribution is included that penalizes results whose ratio to the initial guess has a large total curvature (second derivative). The regularization strength is an adjustable parameter, requiring careful study and transparent disclosure.

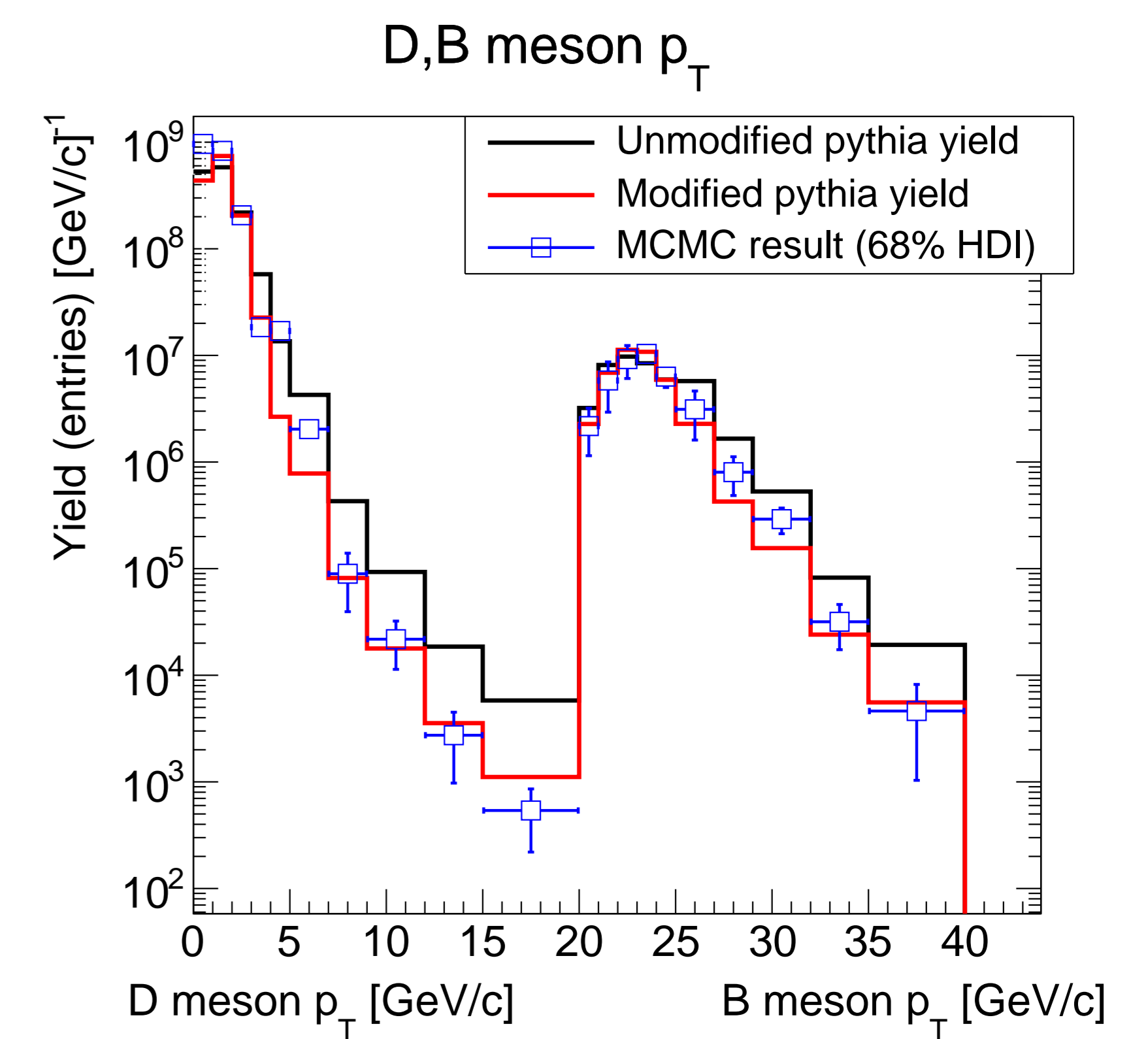
Unfolding results: examples

The output of the sampling algorithm is a 20-dimensional posterior probability distribution. Two marginal distributions from the joint posterior are shown here.



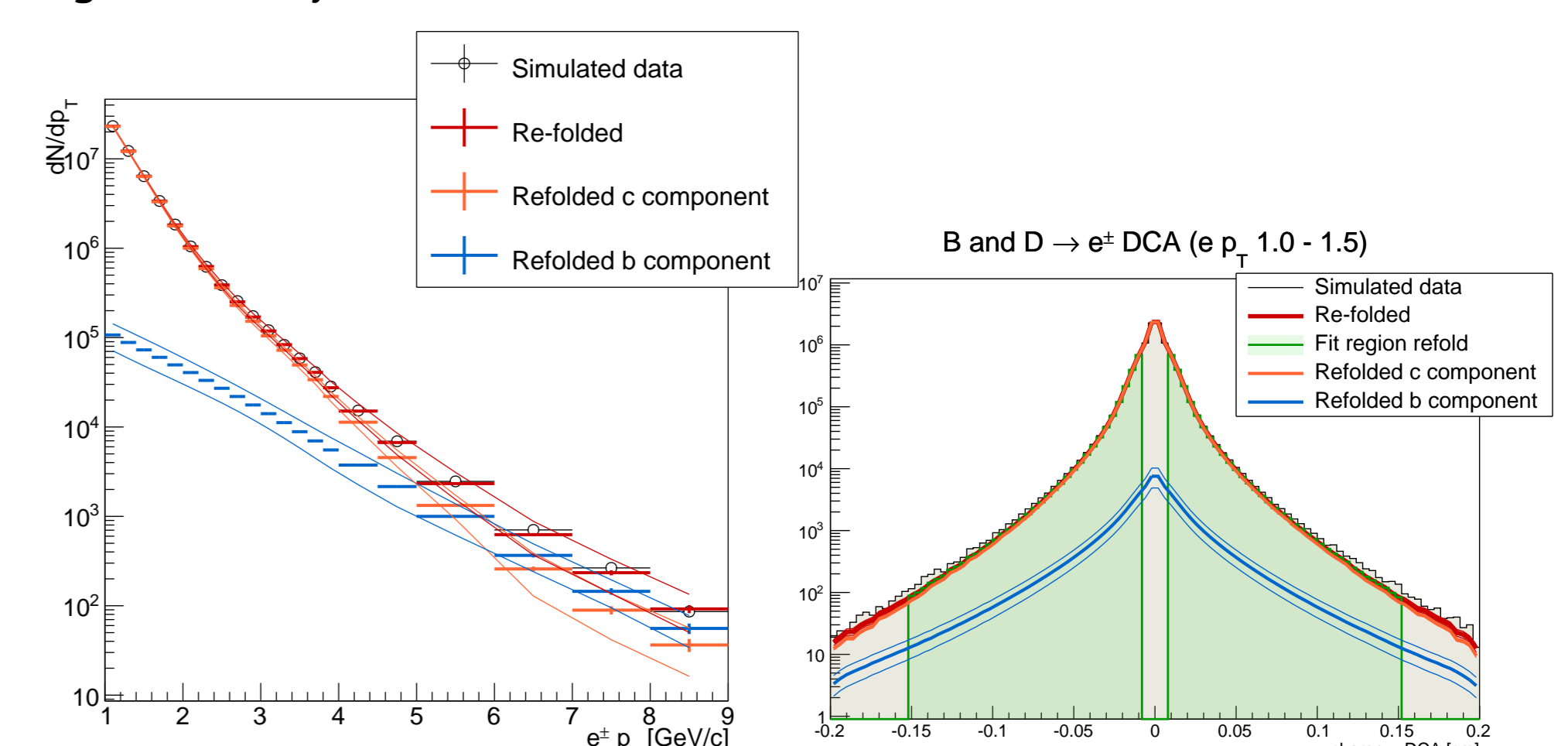
Unfolding results: summary

The shortest interval containing 68% of the samples is shown. This interval is used to summarize the distribution for each hadron p_T dimension, as shown here:



Re-folding the results

A necessary requirement is that the output from the unfolding calculation, when "re-folded" (i.e. multiplied by the decay matrices), agree closely with the observed data.

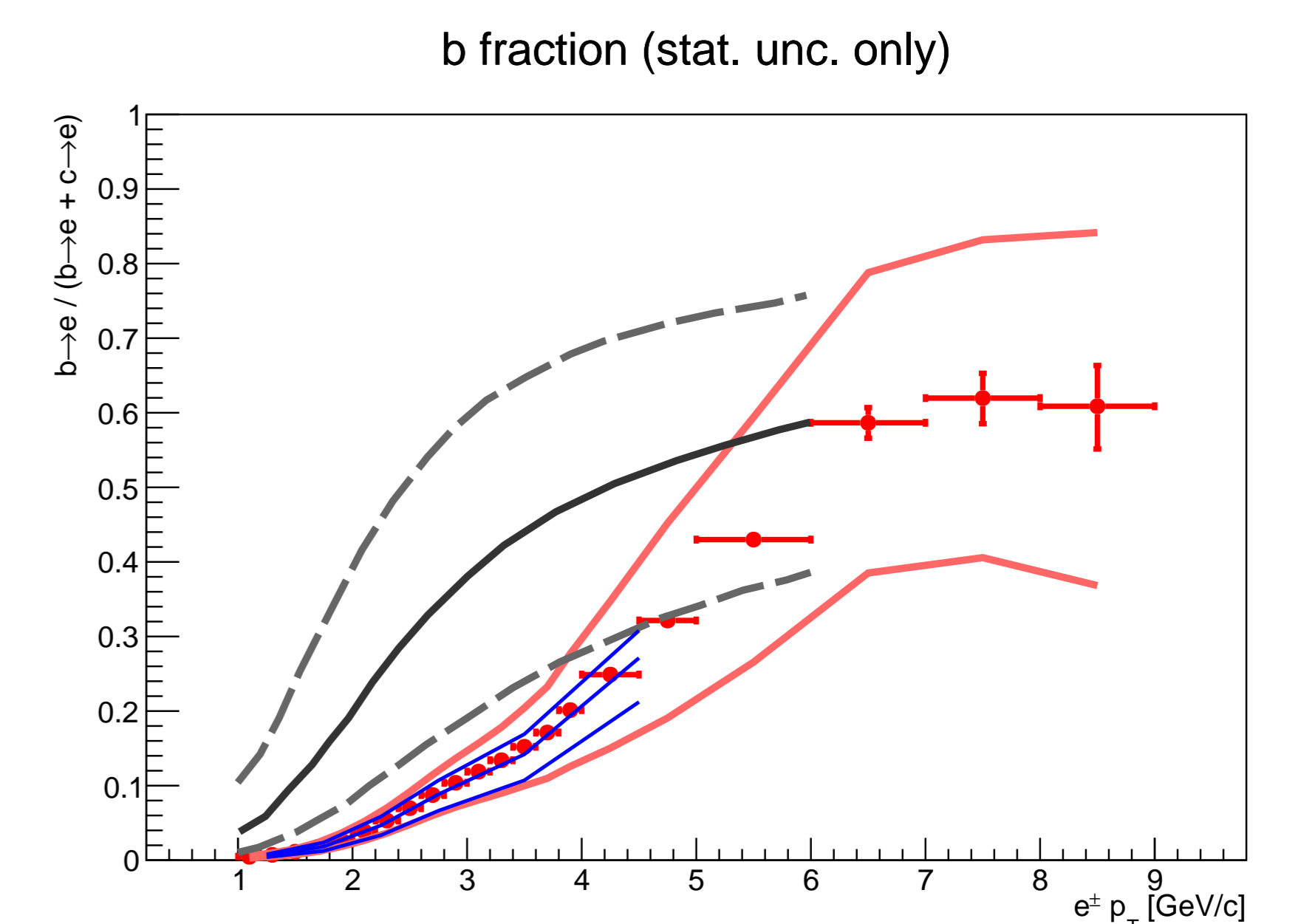


This should be the case for both datasets, and the result should be self-consistent. This self-consistency is shown below.

Beauty fraction vs. electron p_T

The red curve is from the re-folded electron p_T spectrum, and the blue curve is from the re-folded DCA distributions.

The black curve is the result of a FONLL calculation [PRC 84, 044905 (2011)].



Summary

We leave you with the following points:

- Even with no "knowledge" of R_{AA} in the unfolding matrices, the unfolding algorithm approximately recovers the correct modified hadron spectrum when given a noisy, modified electron dataset.
- The result is not perfect. Regularization is required to impose some degree of smoothness on the result. This is a bias introduced by the experimenter based on prior expectations.
- The agreement of a re-folded result with the input dataset is a necessary, but not sufficient requirement for an accurate result.

We're doing this with real data too! Keep an eye out for a PHENIX publication.