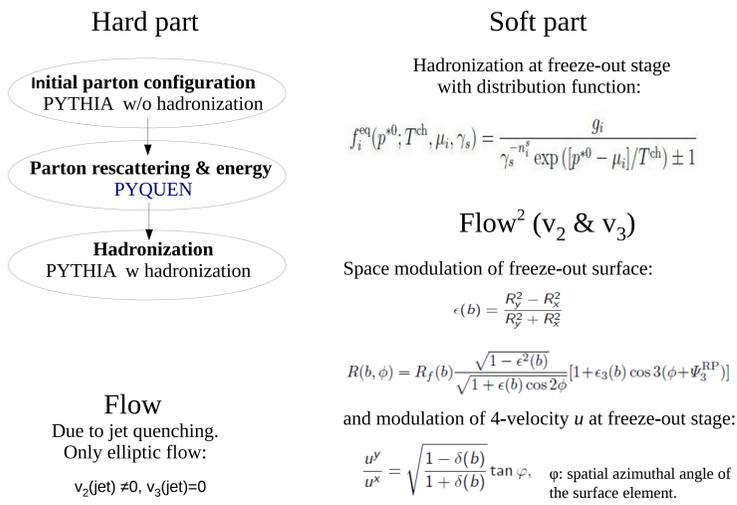


Dihadron angular correlations in PbPb collisions with HYDJET++ model

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The hybrid model HYDJET++, which includes soft and hard physics, is employed for the analysis of azimuthal anisotropy harmonics and dihadron angular correlations measured in PbPb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. The soft part of the model represents a thermal hadron production at the freeze-out hypersurface in accordance with hydrodynamical calculations. The possible triangular shape fluctuation of the initial overlap density of the colliding nuclei was implemented in HYDJET++ by the modulation of the final freeze-out hypersurface with the appropriate triangular coefficient, which results in triangular flow v_3 . Along with elliptic flow v_2 , it generates higher order flow coefficients, as well as a specific structure of dihadron angular correlations on relative azimuthal angle in a broad range of relative pseudorapidities ($\Delta\phi, \Delta\eta$). The comparison of model results with the LHC data on short- and long-range angular correlations is presented for different collisions centralities and transverse momenta intervals.

HYDJET++ model¹



Interference of v_2 and v_3 flow harmonics leads to appearance of higher order harmonics.

- I.P. Lokhtin, L. Malinina, S. Petrushanko, A. Snigirev, I. Arsene, K. Tywoniuk, Comp. Phys. Comm. 180 (2009) 779, <http://cern.ch/lokhtin/hydjet++>
- L.V. Bravina et al., Eur. Phys. J. C74 (2014) 2807

Dihadron correlations

The signal and background function:

$$S(\Delta\varphi, \Delta\eta) = \frac{1}{N_{\text{trig}}} \frac{d^2 N^{\text{same}}}{d\Delta\varphi d\Delta\eta} \quad B(\Delta\varphi, \Delta\eta) = \frac{1}{N_{\text{trig}}} \frac{d^2 N^{\text{mixed}}}{d\Delta\varphi d\Delta\eta}$$

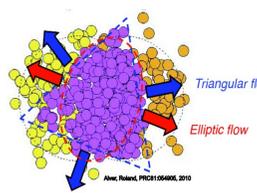
Definition 1: $C(\Delta\phi, \Delta\eta) \equiv \frac{N^{\text{mixed}}}{N^{\text{same}}} \times \frac{N^{\text{same}}(\Delta\phi, \Delta\eta)}{N^{\text{mixed}}(\Delta\phi, \Delta\eta)}$

$\Delta\varphi = \varphi_1 - \varphi_2$
 $\Delta\eta = \eta_1 - \eta_2$
 N_{trig} : number of trigger particles.

Definition 2: $\frac{1}{N_{\text{trig}}} \frac{d^2 N}{d\Delta\varphi d\Delta\eta} = B(0, 0) \frac{S(\Delta\varphi, \Delta\eta)}{B(\Delta\varphi, \Delta\eta)}$

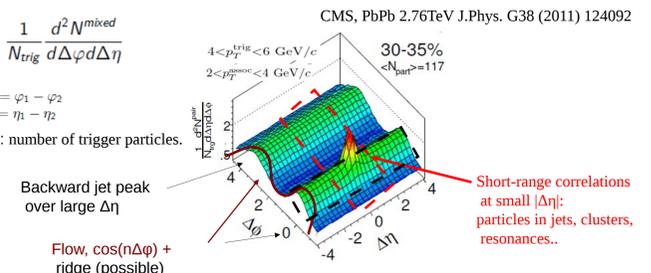
Flow and dihadron correlations:

$$\frac{dN}{d\varphi} = \frac{1}{2\pi} (1 + \sum 2v_n(p_t) \cos(n[\varphi - \Psi_n]))$$

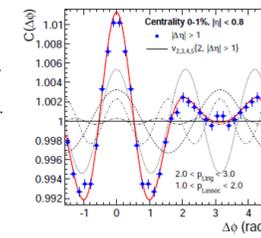


$$\frac{dN^{\text{pairs}}}{d\Delta\phi} \propto 1 + \sum_{n=1}^{\infty} 2V_{n\Delta}(p_T^1, p_T^2) \cos(n\Delta\phi)$$

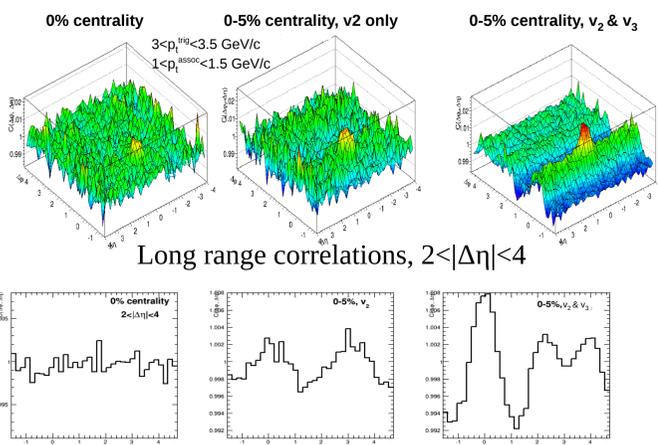
$$V_{n\Delta} = \langle \cos(n\Delta\varphi) \rangle$$



K. Aamodt, et al., (ALICE), Phys. Rev. Lett. 107 (2011) 032301

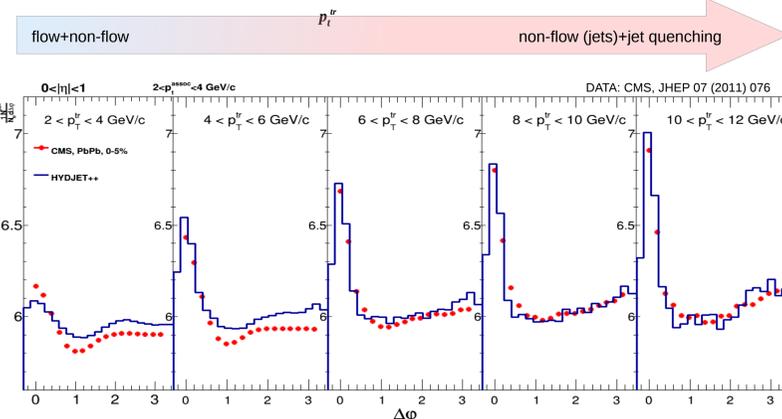


Dihadron correlations in HYDJET++ model



- Long-range correlations appear due to flow.
- v_3 leads to double-peak structure at away side over $\Delta\phi$.

Short range correlations, $|\Delta\eta| < 1$



Test of factorization: $V_{\Delta n} \approx v_n\{2\}(p_T^{\text{trig}}) * v_n\{2\}(p_T^{\text{assoc}})$

At low p_t one neglects non-flow and estimates:

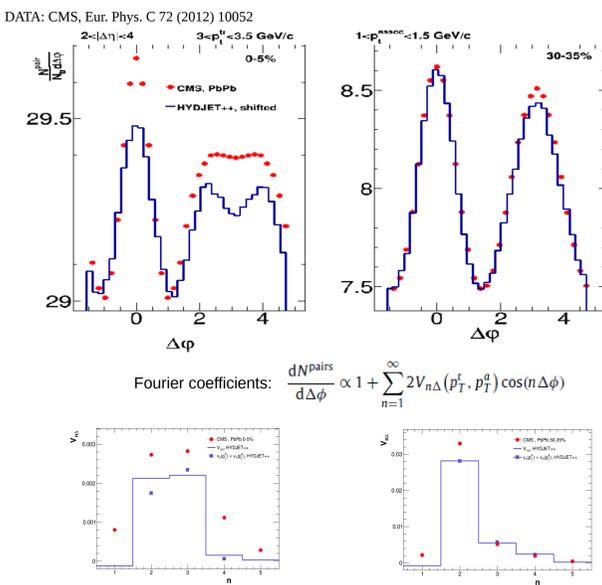
$$v_n(p_t^{\text{low}}) = \sqrt{V_{\Delta n}(p_t^{\text{low}}, p_t^{\text{low}})}$$

$$v_n(p_t) = V_{\Delta n}(p_t, p_t^{\text{low}}) / v_n(p_t^{\text{low}}) \quad (1)$$

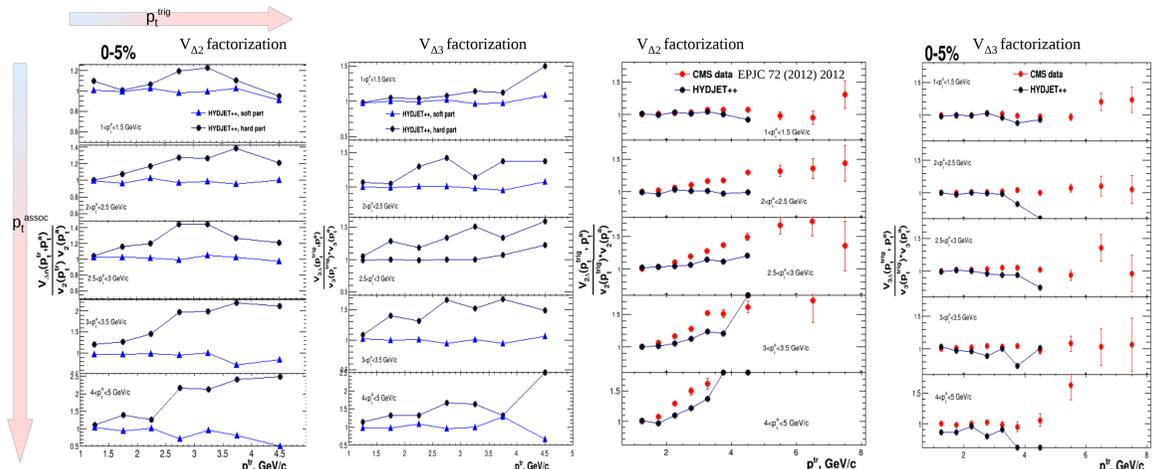
(here v_n is effectively equivalent to $v_n\{2\}$)

It is shown¹ that jet-correlations from PYTHIA simulations of pp collisions also approximately factorize.

I. D. Kikola, L. Yi, S. Esumi, F. Wang, W. Xie, Phys. Rev. C 86 (2012) 014901.



- Central collisions: $V_{\Delta 4}, V_{\Delta 5}$ coefficients are much less than in data, $V_{\Delta n} \approx v_n(p_T^{\text{trig}}) * v_n(p_T^{\text{assoc}})$ (+- non-flow).
- Mid-peripheral collisions: $V_{\Delta n}$ coefficients are described well.



$V_{\Delta n}$ factorization only for soft and only for hard part of the model. $v_n(p_T^{\text{trig}}), v_n(p_T^{\text{assoc}})$ are found by Eq. (1).

$V_{\Delta n}$ factorization for data (Eur. Phys. C 72 (2012) 10052) and the model.

- The correlations of soft particles factorize well at all p_t .
- The correlations of hard (jet) particles doesn't factorize at $p_t > 2$ GeV/c.
- Factorization for $V_{\Delta 3}$ holds both in data and in model in a large p_t range.
- Factorization for $V_{\Delta 2}$ breaks in data at lower p_t , compared to HYDJET++. This could be due to v_2 flow fluctuations, not present in the model.

CONCLUSIONS:

- In HYDJET++ single-particle flow v_2 & v_3 are tuned to describe data at low p_t . The resulting from v_2 & v_3 interference higher order flow harmonics describes data at semi-central collisions well but underestimate data at more central events.
- Similarly it has been found that dihadron correlations (both long-range and short range) are described well at semi-central collisions.

- The factorization of Fourier coefficients of dihadron correlation onto single-particle flow coefficients at low p_t may be fulfilled also in a presence of non-flow (or other than flow sources of correlations).