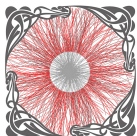


Elliptic flow and nuclear modification factor within a partonic transport model

Florian Senzel

with J. Uphoff, O. Fochler, C. Wesp, Z. Xu and C. Greiner
based on arXiv:1401.1364

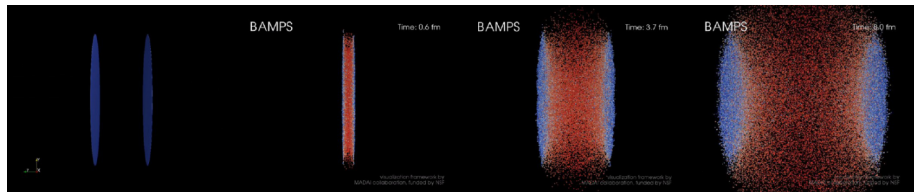


XXIV
QUARK
MATTER
DARMSTADT
2014

Darmstadt, 19.05.2014

Outline

- 1 Motivation
- 2 The partonic transport model BAMPS
- 3 Recent results about the...
 - ... nuclear modification factor R_{AA}
 - ... elliptic flow v_2
 - ... momentum imbalance A_J of reconstructed jets

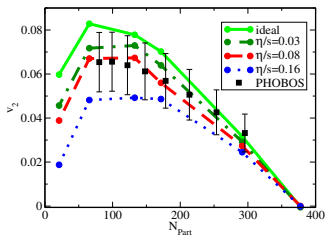
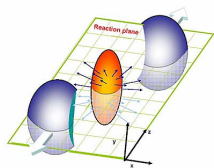


Visualization by Jan Uphoff
Visualization framework courtesy MADAI collaboration
funded by the NSF under grant NSF-PHY-09-41373

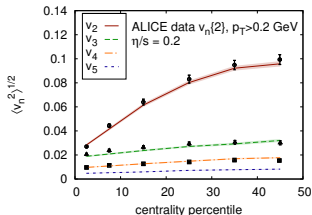
Collectivity of the bulk regime: Elliptic flow v_2

Fourier decomposition of particle spectra

$$\frac{d^3N}{p_t dp_t dy d\phi} (p_t, y, \phi) = \frac{1}{2\pi} \frac{d^2N}{p_t dp_t dy} [1 + 2v_2(p_t, y) \cos(2\phi) + \dots]$$



by Romatschke, Phys.Rev.Lett. 99, (2007)



by Gale et al., Phys.Rev.Lett. 110 (2013)

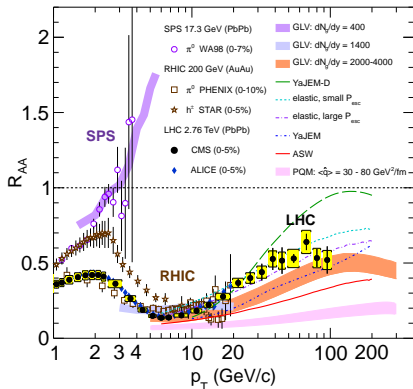
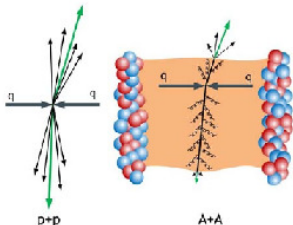
State-of-the-art

✓ Well described by relativistic (viscous) hydrodynamics

Jet quenching: Nuclear modification factor R_{AA}

Suppression of inclusive particle spectra

$$R_{AA} = \frac{d^2 N_{AA}/dp_t dy}{N_{bin} d^2 N_{pp}/dp_t dy}$$



by CMS Collaboration, Eur. Phys. J. C (2012)

State-of-the-art

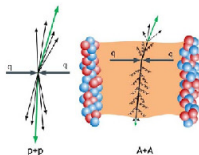
✓ Well described by perturbative quantum chromodynamics

Our question:

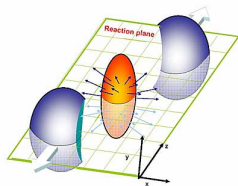
Can perturbative QCD interactions explain in a common framework

both the **high p_t** and the **bulk medium** regime

and thereby give microscopical insight into the QGP?



pQCD?



The partonic transport model BAMPS

BAMPS $\hat{=}$ Boltzmann Approach to Multi-Parton Scattering¹

Numerical solver for the (3+1)D Boltzmann transport equation for partons on the mass-shell:

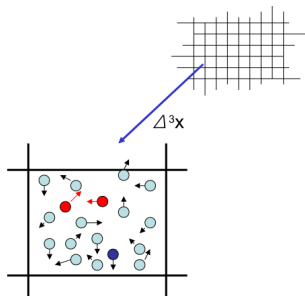
$$\frac{\partial f}{\partial t} + \frac{\mathbf{p}}{E} \frac{\partial f}{\partial \mathbf{r}} = C_{2 \rightarrow 2} + C_{2 \leftrightarrow 3}$$

- Massless particles (gluons & quarks)
- Discretized space ΔV and time Δt :

$$P_{2 \rightarrow 2} = v_{rel} \sigma_{2 \rightarrow 2} \frac{\Delta t}{\Delta V} \quad P_{2 \rightarrow 3} = v_{rel} \sigma_{2 \rightarrow 3} \frac{\Delta t}{\Delta V}$$

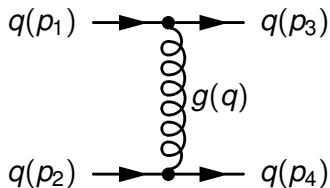
$$P_{3 \rightarrow 2} = \frac{l_{3 \rightarrow 2}}{8E_1 E_2 E_3} v_{rel} \frac{\Delta t}{\Delta V^2}$$

- Test-particles ansatz N_{test}



¹Xu and Greiner, Phys.Rev.C71 (2005); Xu and Greiner, Phys.Rev.C76 (2007)

Implemented processes - elastic collisions



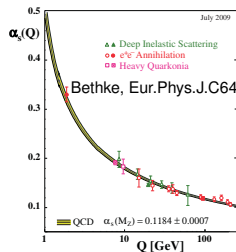
Screened ME with running coupling

$$|\overline{\mathcal{M}}_{X \rightarrow Y}|^2 = C_{X \rightarrow Y} 64\pi^2 \alpha_s^2(t) \frac{s^2}{[t - m_D^2(\alpha_s(t))]^2}$$

$$\text{with } m_D^2(\alpha_s(t)) = d_G \pi \alpha_s(t) \int \frac{d^3 p}{(2\pi)^3} \frac{1}{p} (N_c f_g + N_f f_q)$$

LO pQCD cross-sections

$$\begin{aligned} gg &\rightarrow gg \\ gg &\rightarrow q\bar{q} \\ q\bar{q} &\rightarrow gg \quad \text{and} \quad q\bar{q} \rightarrow q'\bar{q}' \\ qg &\rightarrow qg \quad \text{and} \quad \bar{q}g \rightarrow \bar{q}g \\ q\bar{q} &\rightarrow q\bar{q} \\ qq &\rightarrow qq \quad \text{and} \quad \bar{q}\bar{q} \rightarrow \bar{q}\bar{q} \\ qq' &\rightarrow qq' \quad \text{and} \quad q\bar{q}' \rightarrow q\bar{q}' \end{aligned}$$



Uphoff, Fochler, Xu, Greiner: Phys.Rev.C84 (2011)

Implemented processes - radiative processes

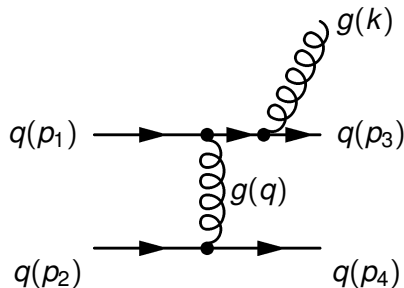
Improved Gunion-Bertsch ME

$$|\overline{\mathcal{M}}_{X \rightarrow Y+g}|^2 = |\overline{\mathcal{M}}_{X \rightarrow Y}|^2 48\pi\alpha_s(k_\perp^2) (1 - \bar{x})^2 \times \left[\frac{\mathbf{k}_\perp}{k_\perp^2} + \frac{\mathbf{q}_\perp - \mathbf{k}_\perp}{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2 + m_D^2(\alpha_s(k_\perp^2))} \right]^2$$

$$\text{with } \bar{x} = k_\perp e^{|\gamma|} / \sqrt{s}$$

2 → 3 processes

$$\begin{aligned} gg &\rightarrow ggg \\ qg &\rightarrow qgg \quad \text{and} \quad \bar{q}g &\rightarrow \bar{q}gg \\ q\bar{q} &\rightarrow q\bar{q}g \\ qq &\rightarrow qqg \quad \text{and} \quad \bar{q}\bar{q} &\rightarrow \bar{q}\bar{q}g \\ qq' &\rightarrow qq'g \quad \text{and} \quad \bar{q}\bar{q}' &\rightarrow \bar{q}\bar{q}'g \end{aligned}$$

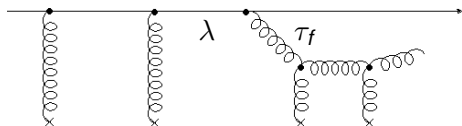


- Gunion, Bertsch: Phys.Rev.D25 (1982)
- Fochler, Uphoff, Xu, Greiner: Phys.Rev.D88 (2013)

Effective modeling of the LPM effect

Issue

Coherence effects within a **semi-classical** approach are not trivial.



Effective method

Parent parton is not allowed to scatter before emitted gluon is formed:

$$|\mathcal{M}_{2 \rightarrow 3}|^2 \rightarrow |\mathcal{M}_{2 \rightarrow 3}|^2 \Theta(\lambda - X_{\text{LPM}} \tau_f)$$

$X_{\text{LPM}} = 0$ No LPM suppression

$X_{\text{LPM}} = 1$ Only independent scatterings (forbids too many emissions)

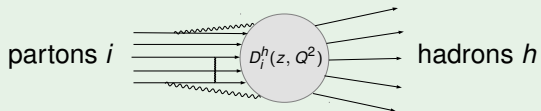
$X_{\text{LPM}} \in (0; 1)$ Allows effectively some interference effects

From partons to hadronic observables

“High p_t ” observables

- Folding with fragmentation functions $D_i^h(z, Q^2)$,

$$\frac{d^2 N^h}{dp_t dy}(p_t^h) = \sum_i \int_{z_{min}}^1 dz \frac{d^2 N^i}{dp_t dy}\left(\frac{p_t^h}{z}\right) D_i^h(z, Q^2) \quad \text{with } z = \frac{p_t^h}{p_t^i}$$

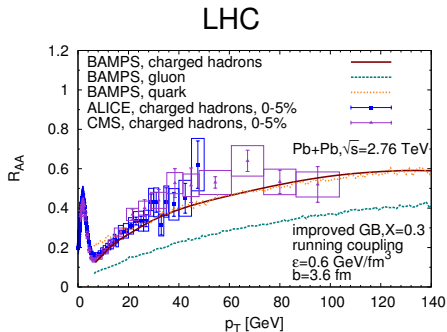
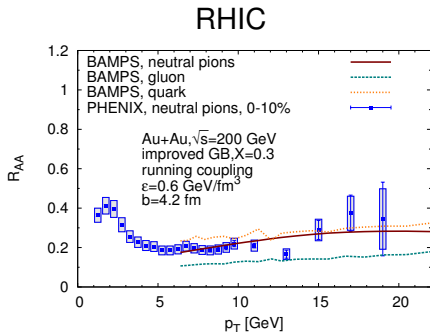


Fragmentation functions by e.g. Albino, Kniehl, Kramer: Nucl.Phys.B803(2008)

“Low p_t ” observables

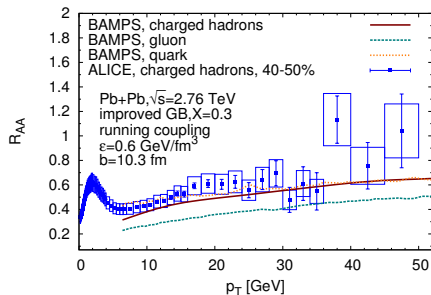
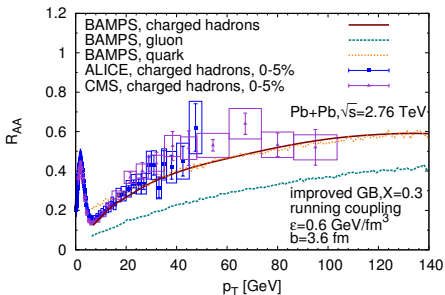
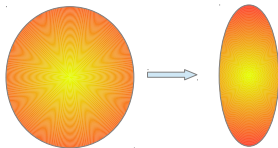
- Microscopic hadronization processes are unknown.
- Integrated flow should not be sensitive to phase transition.

Nuclear modification factor R_{AA} of central HI-collisions

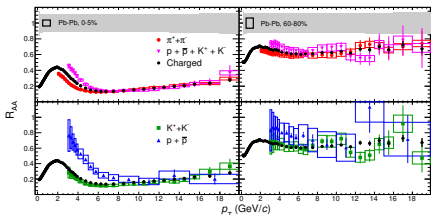
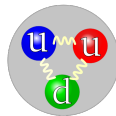
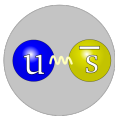


- PYTHIA initial conditions distributed by Glauber model.
- After fixing the LPM parameter $X_{LPM} = 0.3$ by comparing to RHIC data, BAMPS describes the R_{AA} also at LHC.
- Suppression caused by the interplay between the improved GB matrix element and the microscopic running coupling.

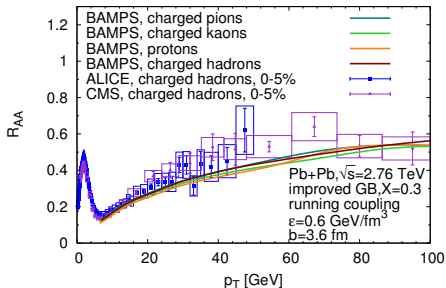
R_{AA} of peripheral HI-collisions



R_{AA} of central HI collisions for different hadron species

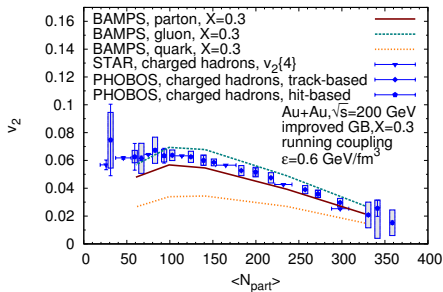


ALICE Collaboration, arXiv:1401.1250

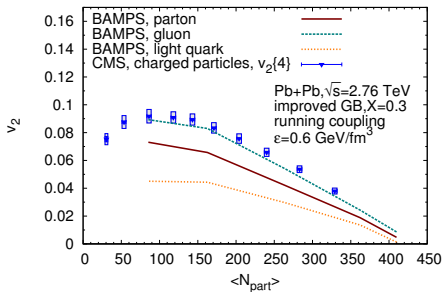


Integrated elliptic flow v_2 (N_{part})

RHIC



LHC



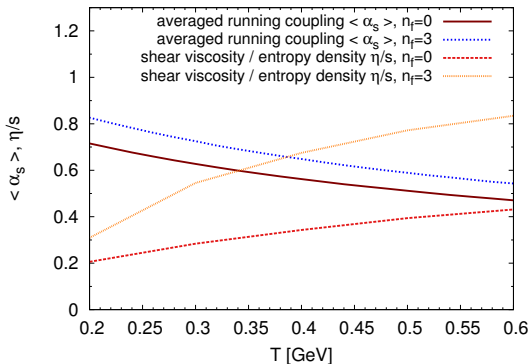
Same setup with LPM parameter $X_{\text{LPM}} = 0.3$ leads to a sizable elliptic flow built up in the partonic phase.

Attention

No hadronization for bulk \Rightarrow No hadronic after-burner \Rightarrow Missing 10-20%?!

Macroscopic quantities from microscopic interactions

Static medium



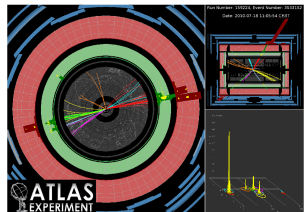
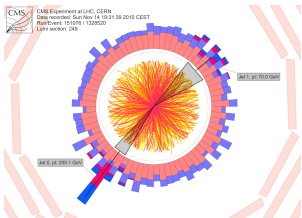
Shear viscosity ratio η/s

- Reason for large flow: small shear viscosity over entropy ratio
- Calculated with Green-Kubo formalism
- Recent viscous hydro: $\eta/s = 0.2$

Running coupling $\alpha_s(T)$

- Temperature dependent coupling by microscopically evaluated interactions.

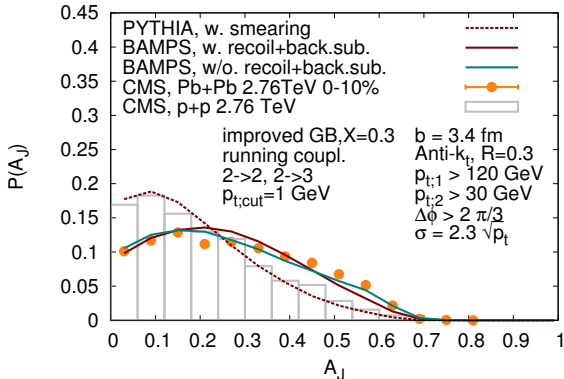
Reconstructed jets in heavy-ion collisions



Momentum imbalance A_J of reconstructed di-jets

Definition

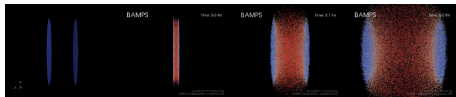
$$A_J = \frac{p_{t; \text{Leading Jet}} - p_{t; \text{Subleading Jet}}}{p_{t; \text{Leading Jet}} + p_{t; \text{Subleading Jet}}}$$



FS, Fochler, Uphoff, Xu, Greiner: arXiv:1309.1657

Conclusions

- Partonic transport provides means for...
 - exploring dynamics of the QGP evolution based on pQCD processes.
 - exploring different observables within a common framework.
- Realistic suppression of jets both at RHIC and LHC.
- Sizable collective flow within the medium by microscopic pQCD cross sections.



Future plans:

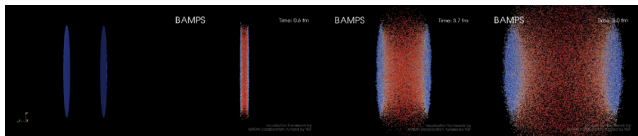
- How does a revisited modeling of the LPM effect change the energy loss and its path-length dependence?
- Systematic studies!

Conclusions

- Partonic transport provides means for...
 - exploring dynamics of the QGP evolution based on pQCD processes.
 - exploring different observables within a common framework.
- Realistic suppression of jets both at RHIC and LHC.
- Sizable collective flow within the medium by microscopic pQCD cross sections.

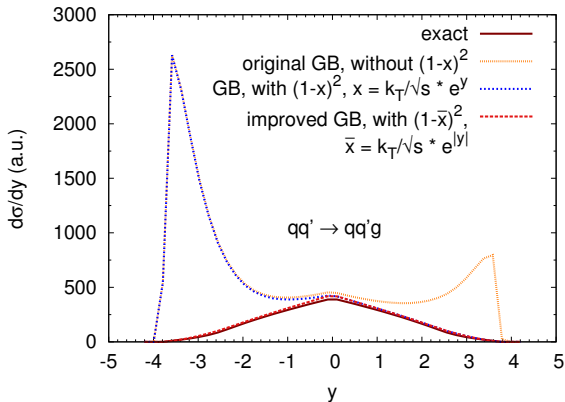
BAMPS@QM2014

- Heavy quarks:
Talk by J. Uphoff
- Mach cones:
Poster by I. Bouras
- Reconstructed jets:
Poster by FS



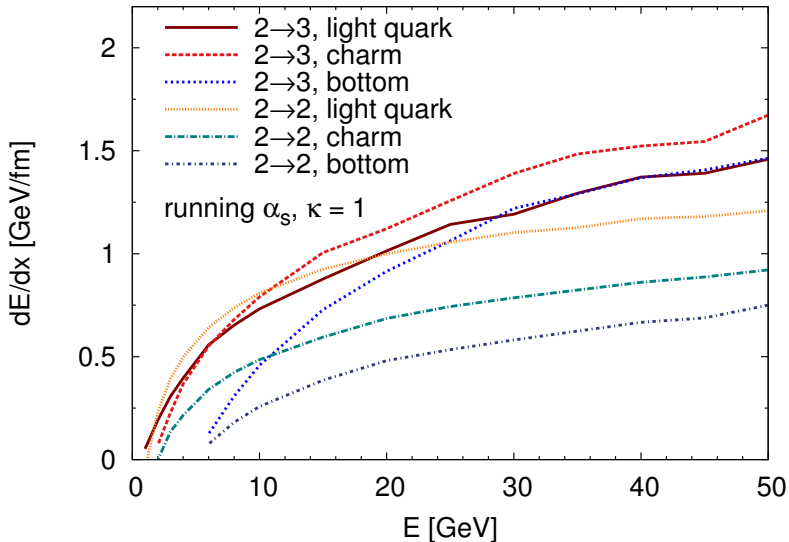
Backup slides

Improved Gunion-Bertsch matrix element



- Infrared screening for both GB and exact: $\theta(\text{cut}) = \theta(p_i p_j - \lambda)$.
- Integration both in GB coordinates and in standard phase space with numeric δ -functions.

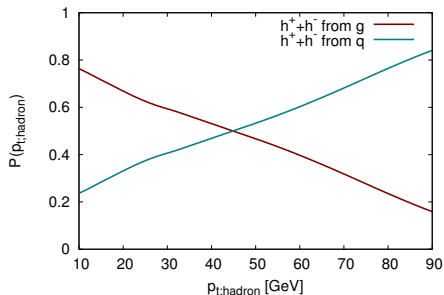
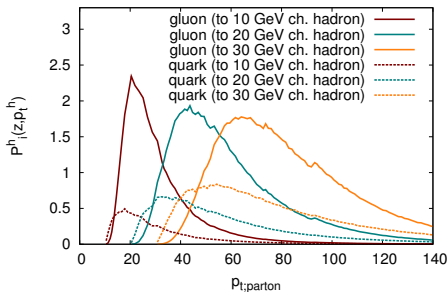
Differential energy loss in a static medium



Closer look on the role of fragmentation

Probability for hadron h with p_t^h out of parton i with $p_t^i = p_t^h/z$

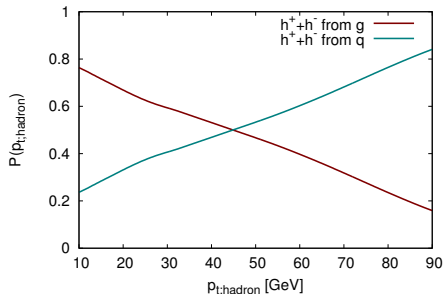
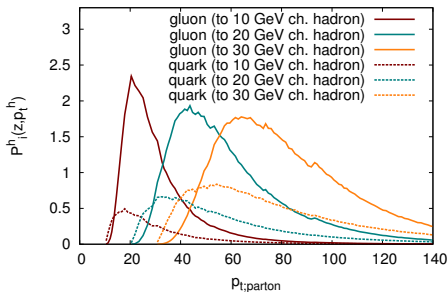
$$P^{i \rightarrow h}(z, p_t^h) = \frac{1}{\frac{d^2 N^h}{dp_t dy}(p_t^h)} \frac{d^2 N^i}{dp_t dy}\left(\frac{p_t^h}{z}\right) D_i^h(z, Q^2)$$



Example: R_{AA} for hadrons with $p_t^h = 30$ GeV

Hadrons with $p_t^h = 30$ GeV stem...

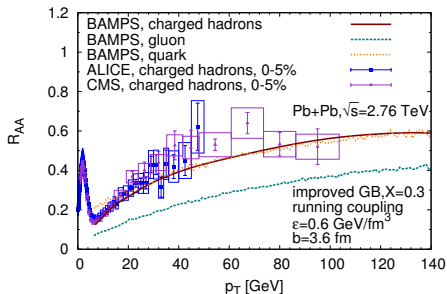
- ... mainly from ≈ 60 GeV gluon and ≈ 45 GeV quark.
- ... $\approx 60\%$ from gluons and $\approx 40\%$ from quarks. 6



Example: R_{AA} for hadrons with $p_t^h = 30$ GeV

Hadrons with $p_t^h = 30$ GeV stem...

- ... mainly from ≈ 60 GeV gluon and ≈ 45 GeV quark.
- ... $\approx 60\%$ from gluons and $\approx 40\%$ from quarks. 6



$$R_{AA}^h(30 \text{ GeV}) = 0.4 R_{AA}^g(60 \text{ GeV}) + 0.6 R_{AA}^q(45 \text{ GeV}) \approx 0.3$$

Example: Shower event with first 100 recoil partons

