

The “ripples” on relativistically expanding fluid with applications in heavy-ion collisions

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Abstract

We present **complete** and **analytic sound wave** solutions on top of both **Bjorken** flow and **Hubble** flow background. We visualize the evolution patterns for **Gaussian** perturbation.

Motivation

- ▶ Hydrodynamics works well for QGP in relativistic heavy-ion collisions.
- ▶ Fluctuations, including initial and thermal ones, play important role in heavy-ion collisions.
- ▶ We study analytically how fluctuation evolves on top of known background.
- ▶ For high energy collisions at LHC and RHIC, the initial expansion is close to Bjorken flow, while the later expansion is close to Hubble flow.

Working Frame

- ▶ hydro equation: $T^{\mu\nu}_{;\mu} = 0$.
- ▶ ideal hydro: $T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu - g^{\mu\nu}p$.
- ▶ E.o.S. $\epsilon = c_s^{-2}p$.
- ▶ perturbation: $p = p_0 + p_1$, $u^\mu = u_0^\mu + u_1^\mu$.
- ▶ keep only 1st-order term \Rightarrow linearize hydro equation.
- ▶ We solve the equations in the most convenient coordinate.

Coordinate

Bjorken:

$$\begin{aligned} t &= \tau_z \cosh \eta_z \\ z &= \tau_z \sinh \eta_z \\ x &= \rho \cos \phi \\ y &= \rho \sin \phi \end{aligned}$$

background flow:

$$\begin{aligned} p_0 &\propto \tau_z^{-4/3} \\ u_0^\mu &= (1, 0, 0, 0) \end{aligned}$$

Hubble:

$$\begin{aligned} t &= \tau_r \cosh \eta_r \\ z &= \tau_r \sinh \eta_r \cos \theta \\ x &= \tau_r \sinh \eta_r \sin \theta \cos \phi \\ y &= \tau_r \sinh \eta_r \sin \theta \sin \phi \end{aligned}$$

$$\begin{aligned} p_0 &\propto \tau_r^{-4} \\ u_0^\mu &= (1, 0, 0, 0) \end{aligned}$$

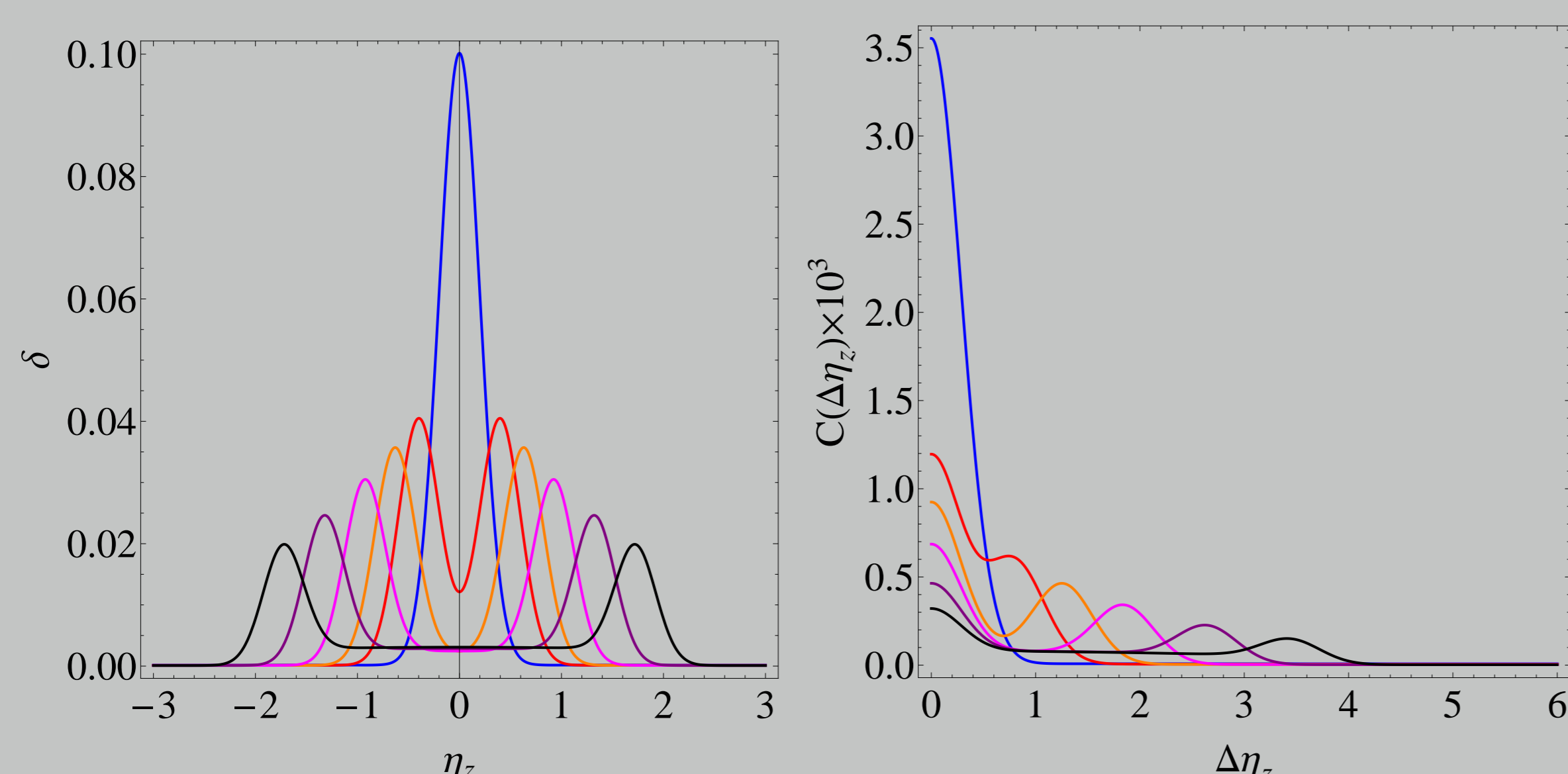
(I) Perturbation on Top of Longitudinal Bjorken Wave

- ▶ We take Gaussian initial perturbation $\frac{\xi}{\sqrt{2\pi\sigma}} e^{-\frac{(\eta-\eta')^2}{2\sigma^2}}$ as an example.
- ▶ Evolution for pressure and velocity

$$\begin{aligned} \frac{p_1}{p_0} &= \frac{\xi}{2\pi} \left(\frac{\tau'}{\tau}\right)^{\frac{1}{3}} \int_{-\infty}^{\infty} dk e^{-\frac{\sigma^2 k^2}{2}} \cos[k(\eta - \eta')] \cos\left[c_s \sqrt{k^2 - \frac{1}{3}} \ln \frac{\tau}{\tau'}\right] \\ u_1^\eta &= \frac{\xi}{4\pi} \left(\frac{\tau'}{\tau}\right)^{\frac{4}{3}} \int_{-\infty}^{\infty} dk e^{-\frac{\sigma^2 k^2}{2}} \frac{\sin[k(\eta - \eta')]}{k} \times \\ &\quad \left(\frac{1}{2} \cos\left[c_s \sqrt{k^2 - \frac{1}{3}} \ln \frac{\tau}{\tau'}\right] - \frac{\sqrt{3}}{2} \sqrt{k^2 - \frac{1}{3}} \sin\left[c_s \sqrt{k^2 - \frac{1}{3}} \ln \frac{\tau}{\tau'}\right] \right) \end{aligned}$$

- ▶ We define $\delta \equiv p_1/p_0$, and we can also calculate correlations

$$C(\Delta\eta_z) = \int \delta(\eta_z) \delta(\eta_z + \delta\eta_z) d\eta_z.$$

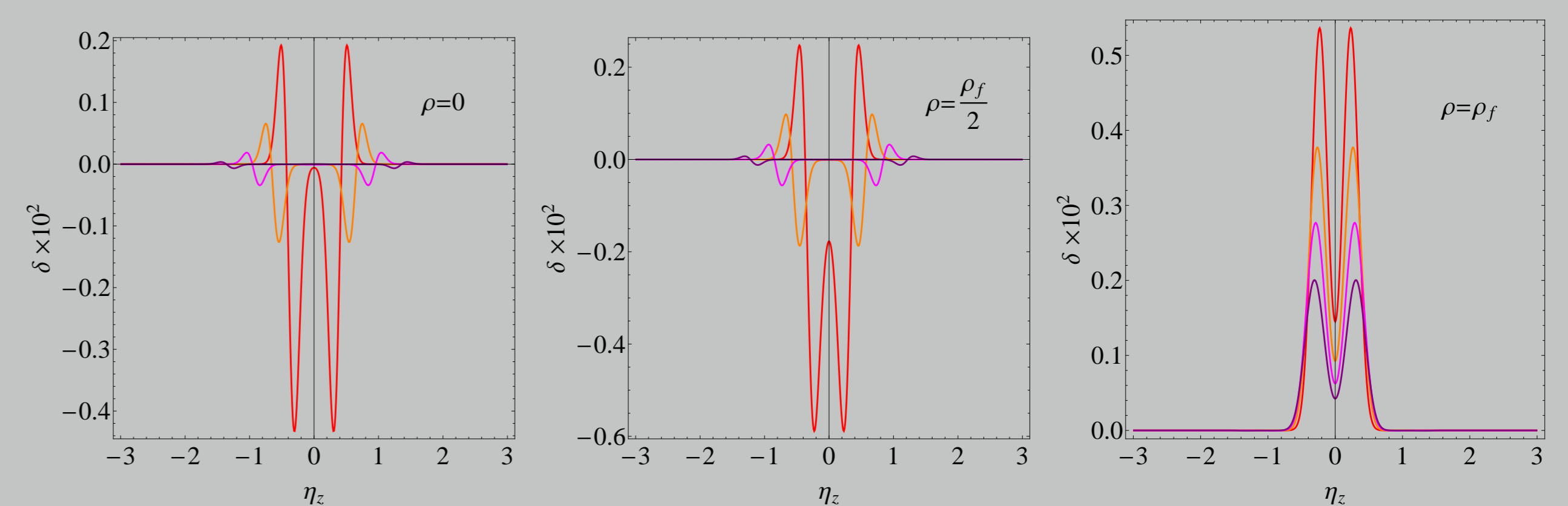


- ▶ curves in different color mean quantities at different τ .
- ▶ sound spreads out inside sound horizon and fills up the interval region.
- ▶ correlation pattern shows a peak associated with the sound horizon generates together with the short-range peak.

(II) Perturbation on Top of General Bjorken Wave

- ▶ Evolution of Gaussian initial perturbation $\frac{\xi}{(2\pi)^{3/2}\sigma^3} e^{-\frac{\eta^2 + \rho^2/\tau'^2}{2\sigma^2}}$:

$$\frac{p_1}{p_0} = \int_0^\infty d\omega \int_{-\infty}^\infty dk e^{ik\bar{\eta}} J_0(\omega\bar{\rho}) W_{\omega,k}(\tau)$$

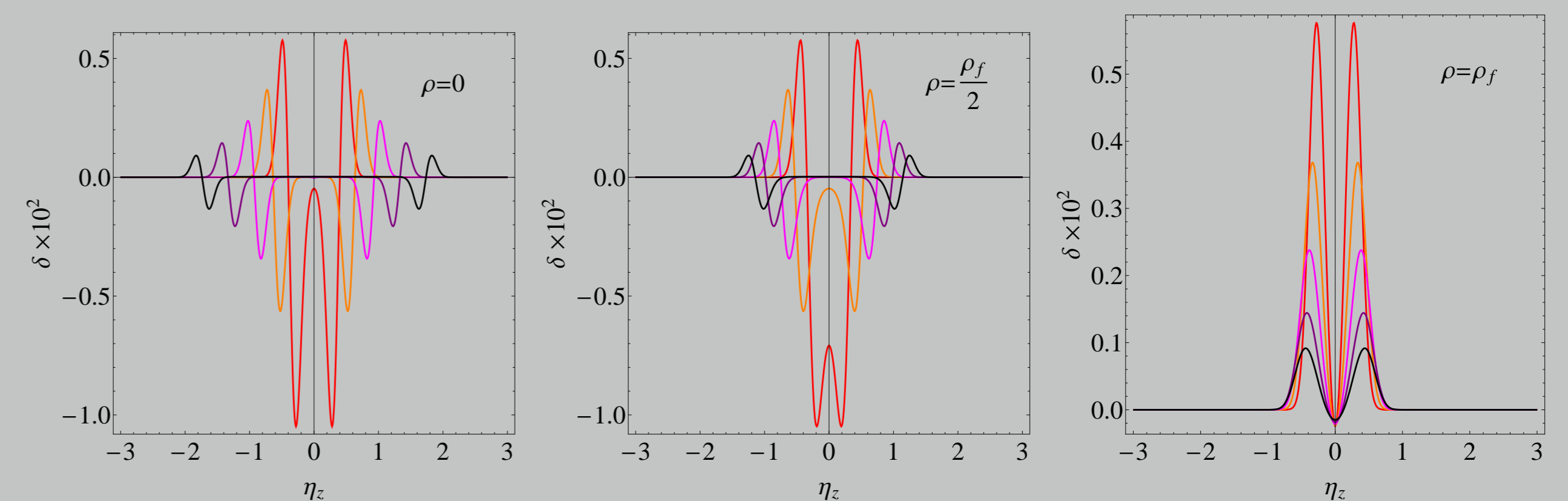


- ▶ $\rho_f = c_s(\tau - \tau')$ is the transverse wave-front.
- ▶ wave at transverse wave-front has nearly vanishing longitudinal velocity.

(III) Perturbation on Top of Hubble Wave

- ▶ Evolution of Gaussian initial perturbation $\frac{\xi}{(2\pi)^{3/2}\sigma^3} e^{-\frac{\eta^2}{2\sigma^2}}$:

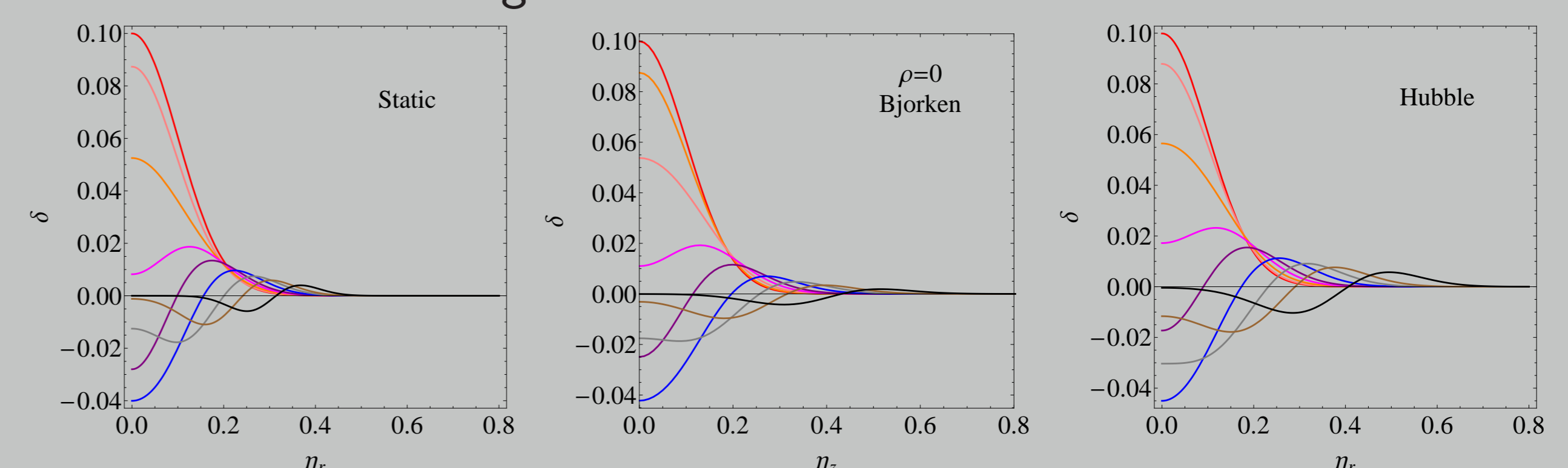
$$\frac{p_1}{p_0} = \frac{\xi}{4\pi^2} \int_{-\infty}^{\infty} \frac{\sin(k\sigma^2)}{k\sigma^2 e^{-\sigma^2/2}} e^{-\frac{\sigma^2 k^2}{2}} \frac{\sin(k\bar{\eta})}{\sinh(\bar{\eta})} \cos\left[\beta_k \ln \frac{\tau}{\tau'}\right] k dk.$$



- ▶ $\rho_f = \tau_r \sinh[c_s \ln(\tau/\tau')]$ is the transverse wave-front.
- ▶ negative perturbation is found in both 3D case, but not in longitudinal Bjorken case.

Early Evolution

- ▶ We visualize early evolution of Gaussian initial perturbation on top of static, Bjorken and Hubble background flow:



- ▶ Early evolution are similar to the static background as the initial local flow in the generating spot is weak.
- ▶ Background independent negative perturbation is seen in 3D flow.

Conclusion & Outlook

- ▶ Generally derive fluctuation propagation in different background.
- ▶ Explicitly show propagation for Gaussian perturbation.
- ▶ Visualize the pattern of evolution.
- ▶ Further work on the application in heavy ion collision

References

J. I. Kapusta *et al*, Phys. Rev. C **85**, 054906 (2012)