



More strange hadrons from QCD thermodynamics and strangeness freeze-out in HIC

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BNL-Bi-CCNU Collaboration:

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Content

arXiv:1404.6511

- (1) **Definitions and motivations:**
cumulants and correlations of conserved charge fluctuations, (lattice) QCD vs HRG
- (2) **Evidence for experimentally not yet observed open strange hadrons** and their importance for bulk thermodynamics
- (3) **Implications for freeze-out conditions**

arXiv:1404.4043

- **Evidence for experimentally not yet observed open charm hadrons**
- **Melting of open charm hadrons**

⇒ see Poster by S. Sharma

Definitions

Expansion of the pressure:

$$\frac{p}{T^4} = \sum_{i,j,k=0}^{\infty} \frac{1}{i!j!k!} \chi_{ijk,0}^{BQS} \left(\frac{\mu_B}{T}\right)^i \left(\frac{\mu_Q}{T}\right)^j \left(\frac{\mu_S}{T}\right)^k$$

$X = B, Q, S$: conserved charges

Lattice

$$\chi_n^X = \left. \frac{\partial^n [p/T^4]}{\partial (\mu_X/T)^n} \right|_{\mu_X=0}$$

generalized susceptibilities

⇒ only at $\mu_X = 0$!

Experiment

$$\begin{aligned} VT^3 \chi_2^X &= \langle (\delta N_X)^2 \rangle \\ VT^3 \chi_4^X &= \langle (\delta N_X)^4 \rangle - 3 \langle (\delta N_X)^2 \rangle^2 \\ VT^3 \chi_6^X &= \langle (\delta N_X)^6 \rangle \\ &\quad - 15 \langle (\delta N_X)^4 \rangle \langle (\delta N_X)^2 \rangle \\ &\quad + 30 \langle (\delta N_X)^2 \rangle^3 \end{aligned}$$

cumulants of net-charge fluctuations

$$\delta N_X \equiv N_X - \langle N_X \rangle$$

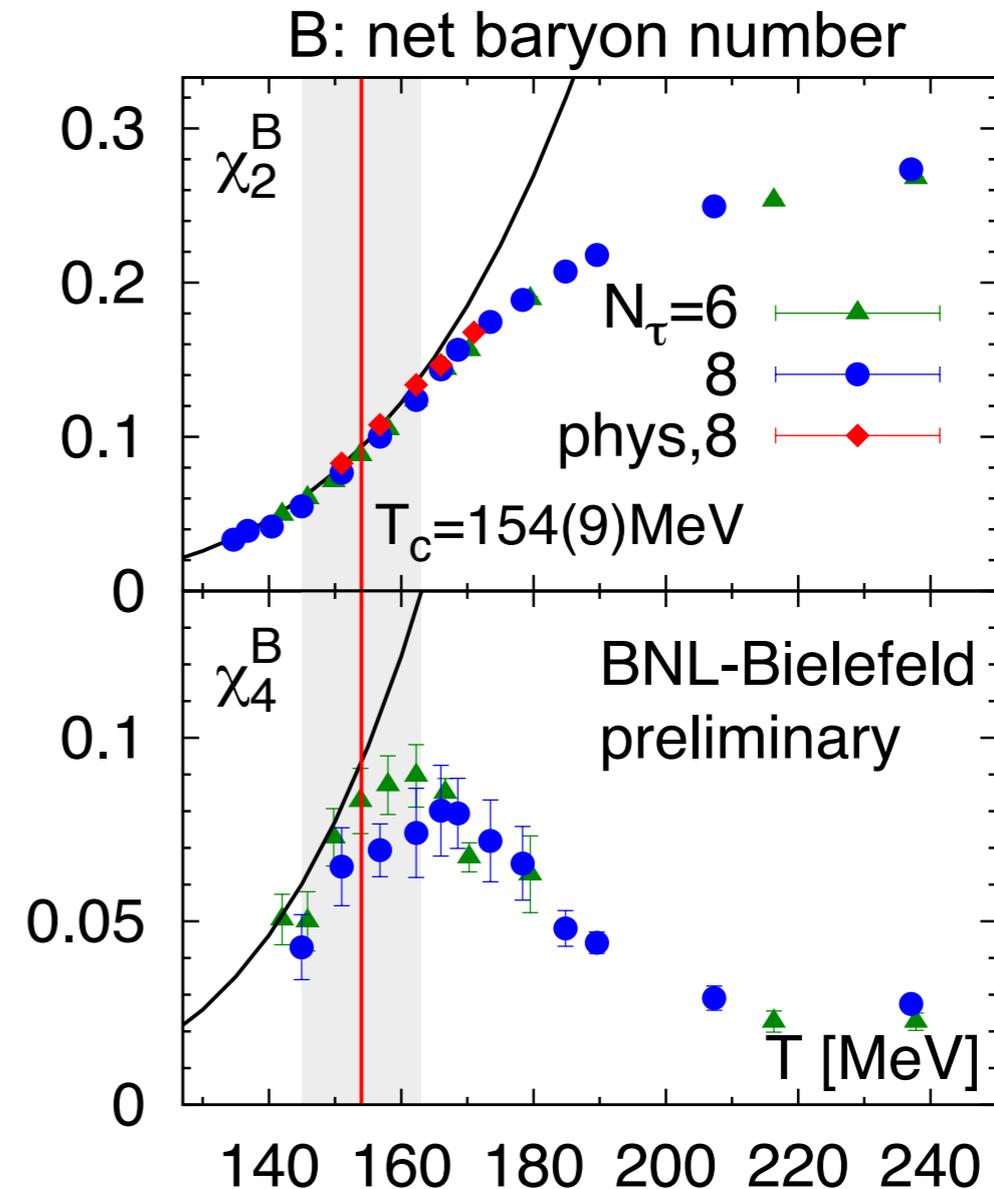
⇒ only at freeze-out $(\mu_f(\sqrt{s}), T_f(\sqrt{s}))$!

(lattice) QCD vs HRG

B,S diagonal cumulants up to 4th order

Lattice details

- (2+1)-flavor of highly improved staggered fermions (HISQ-action)
- a set of different lattice spacings ($N_\tau = 6, 8, 12$)
- two different pion masses: $m_\pi = 140, 160 \text{ MeV}$
- high statistics: $(10 - 16) \times 10^3$ configurations

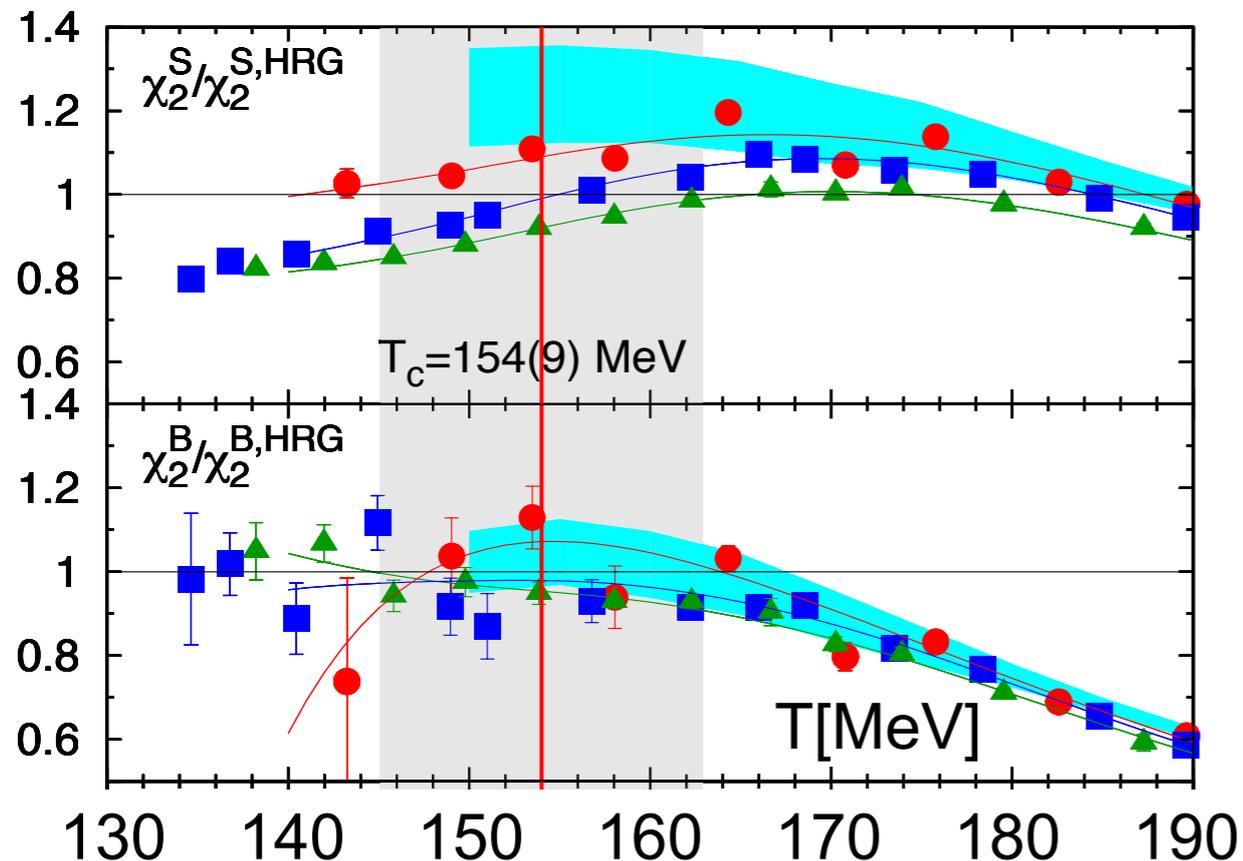


⇒ statistical and systematical errors are under control

⇒ In general: find good agreement with HRG model for $T < 155 \text{ MeV}$

(lattice) QCD vs HRG

However, some notable differences in the strangeness sector



● $N_\tau = 12$ HISQ-Action
■ $N_\tau = 8$
▲ $N_\tau = 6$

cyan bands indicate continuum extrapolations

HotQCD, PRD 86 (2012) 034509
[arXiv: 1203.0784]

HRG:

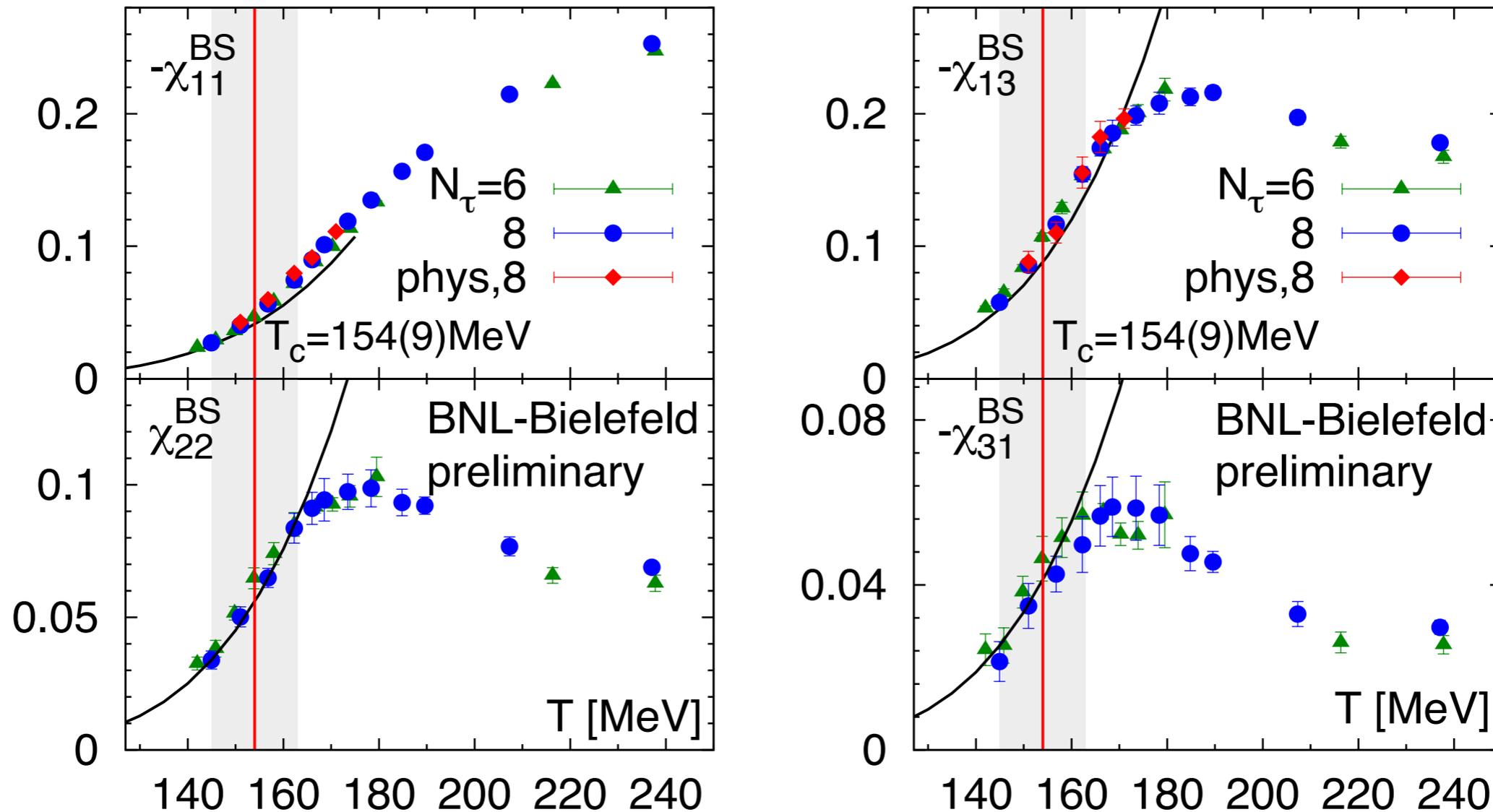
$$p^{\text{HRG}} = \sum_{i \in \text{mesons}} p_{m_i}^M(T, \mu_S) + \sum_{i \in \text{baryons}} p_{m_i}^B(T, \mu_B, \mu_S)$$

hadron masses from PDG up to 2.5 GeV, and

$p_{m_i}^{M/B}$ pressure of free bosonic/fermionic quantum gas

(lattice) QCD vs HRG

Also apparent in B,S off-diagonal cumulants



⇒ overshooting of HRG in the crossover region and below

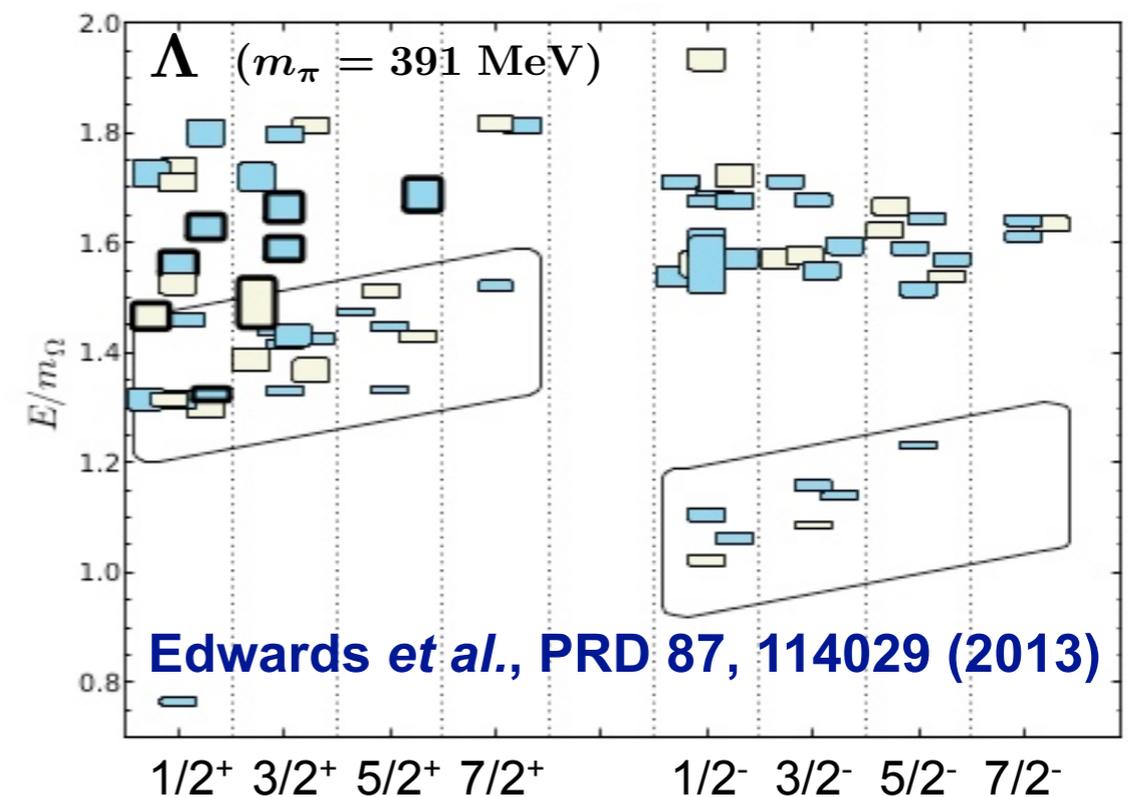
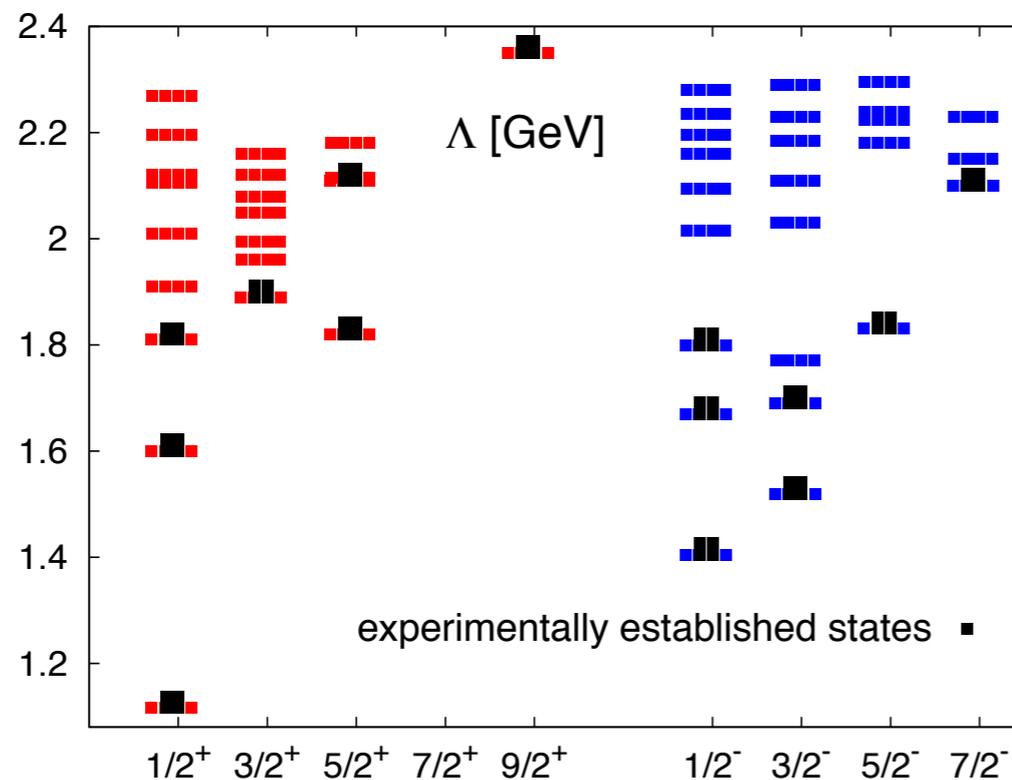
the QM-HRG

Are differences due to missing states in the PDG?

- Obvious in the charm sector (See Poster by S. Sharma)
- How large could be the effect of missing states in the strange sector?

⇒ construct **QM-HRG**, including additional states predicted by Quark-Model

- Use mesonic states from: **S. Capstick and N. Isgur, PRD 34, 2809 (1986).**
- Use baryonic states from: **D. Ebert *et al.*, PRD 79, 114029 (2009)**



- Similar to the spectrum of strange baryons on the lattice

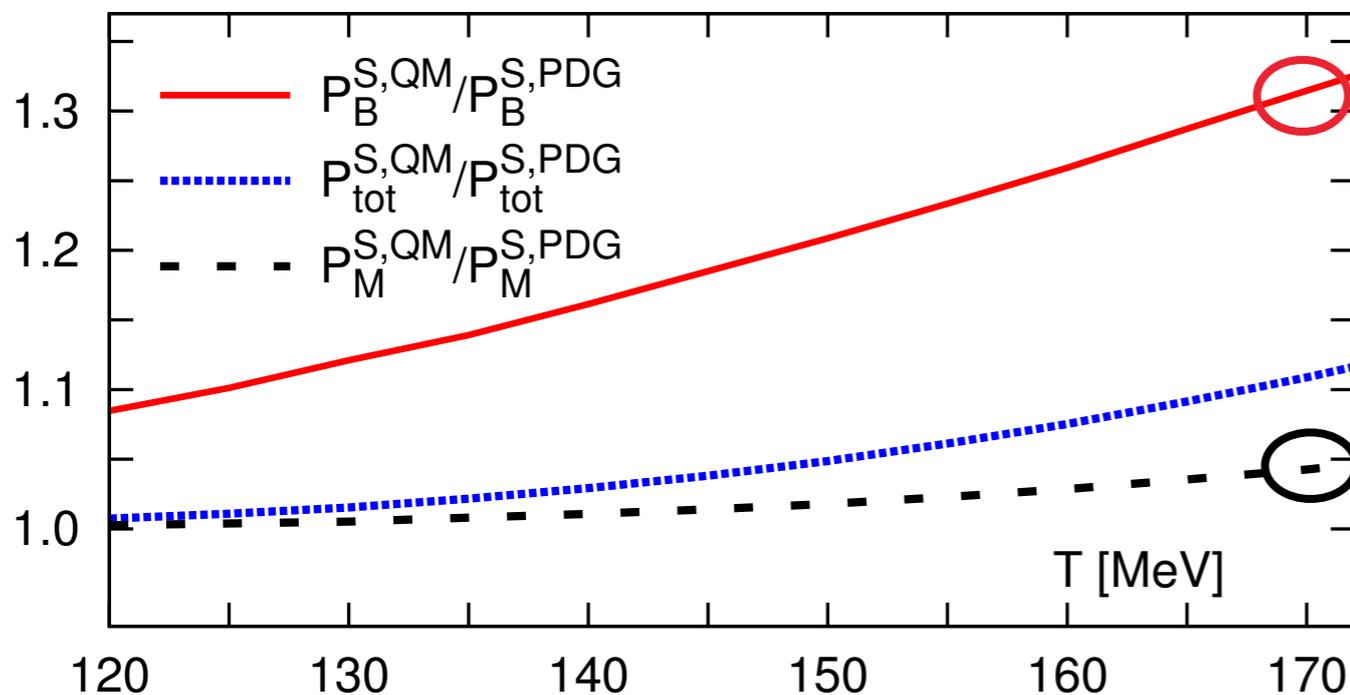
QM-HRG vs PDG-HRG

Partial pressure of strange mesons and baryons:

- Boltzmann approximation is used here and in the following

$$P_{M/B}^{S,X}(T, \vec{\mu}) = \frac{T}{2\pi^2} \sum_{i \in X} g_i \left(\frac{m_i}{T} \right)^2 K_2(m_i/T) \cosh(B_i \hat{\mu}_B + Q_i \hat{\mu}_Q + S_i \hat{\mu}_S)$$

with $X = \text{PDG, QM}$ and $\hat{\mu}_q = \mu_q/T$, $q = B, Q, S$



⇒ open strange baryon sector experimentally much less known, additional baryons contribute up to 30% at $T=170$ MeV

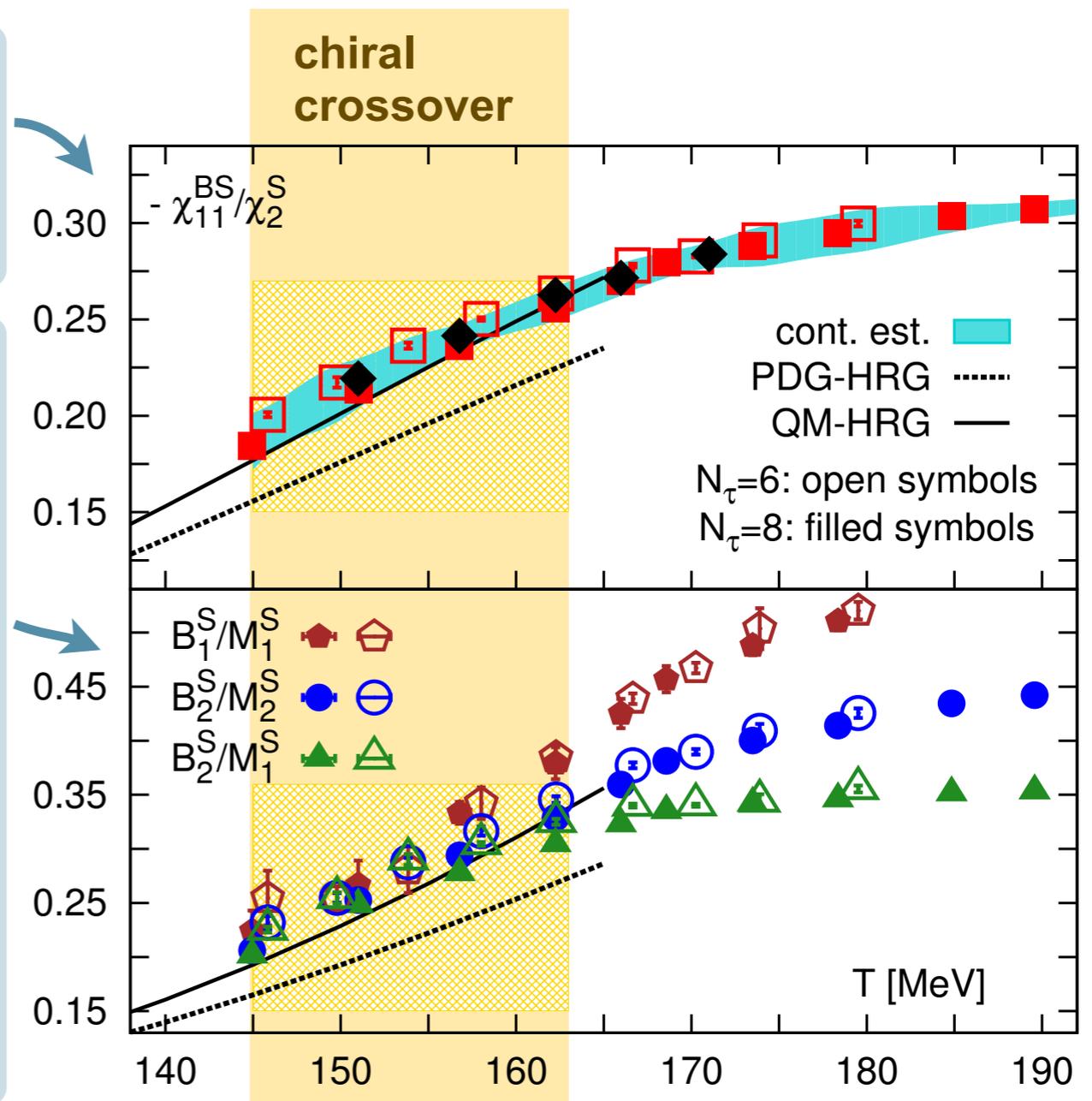
⇒ open strange meson sector experimentally well known

Evidence for more strange hadrons

- **BS-correlation** χ_{11}^{BS}
at low T: weighted sum of partial pressure of strange baryons
- Different **linear combinations** of $\chi_2^S, \chi_4^S, \chi_{11}^{BS}, \chi_{31}^{BS}, \chi_{22}^{BS}, \chi_{13}^{BS}$ are used to project onto partial pressure of strange baryons (B_i^S) and mesons (M_i^S) in the hadronic phase, e.g.

$$B_1^S = -\frac{1}{6}(11\chi_{11}^{BS} + 6\chi_{22}^{BS} + \chi_{13}^{BS})$$

$$B_2^S = \frac{1}{12}(\chi_4^S - \chi_2^S) - \frac{1}{3}(4\chi_{11}^{BS} - \chi_{13}^{BS})$$



⇒ **QM-PDG provides more accurate description of lattice data**

⇒ Re-confirmation of our previous findings [[PRL 111,082301](#)]: onset of melting of open strange hadrons consistent with chiral crossover

Implications for freeze-out conditions

- in order to make contact to experiment: apply strangeness neutrality constraint $n_S = 0$ (and iso-spin $n_Q/n_B = 0.4$)

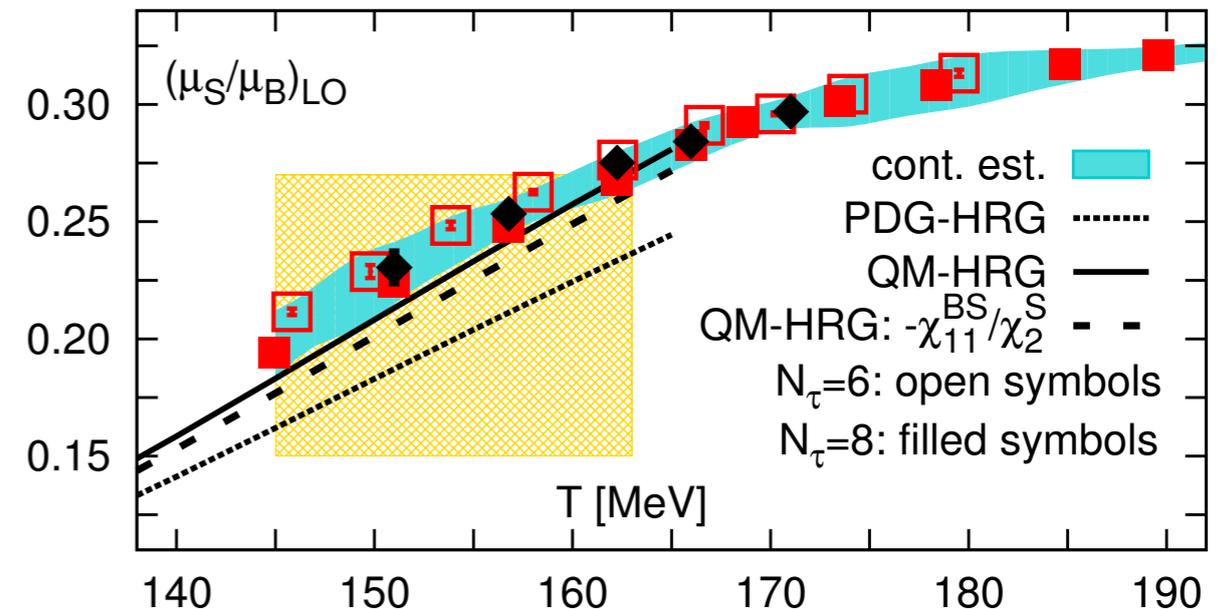
$\Rightarrow T, \mu_B, \mu_S, \mu_Q$ not independent

(Lattice) QCD: unique expansion

$$(\mu_S/\mu_B)(T) = s_1(T) + s_3(T)\hat{\mu}_B^2 + \mathcal{O}(\hat{\mu}_B^4)$$

$$\Rightarrow \left(\frac{\mu_S}{\mu_B}\right)_{LO} \equiv s_1(T) = -\frac{\chi_{11}^{BS}}{\chi_2^S} - \frac{\chi_{11}^{QS}}{\chi_2^S} \frac{\mu_Q}{\mu_B}$$

- small corrections from $\mu_Q > 0$
- neglect NLO contribution, known to be small for $\mu_B \lesssim 200 \text{ MeV}$ [[PRL 111,082301](#)]



HRG: $\mu_S(T, \mu_B)$ depends on relative abundance of open strange baryons and mesons

additional strange baryons
 \Rightarrow larger μ_S
 (at fixed T, μ_B)

Implications for freeze-out conditions

Experiment:

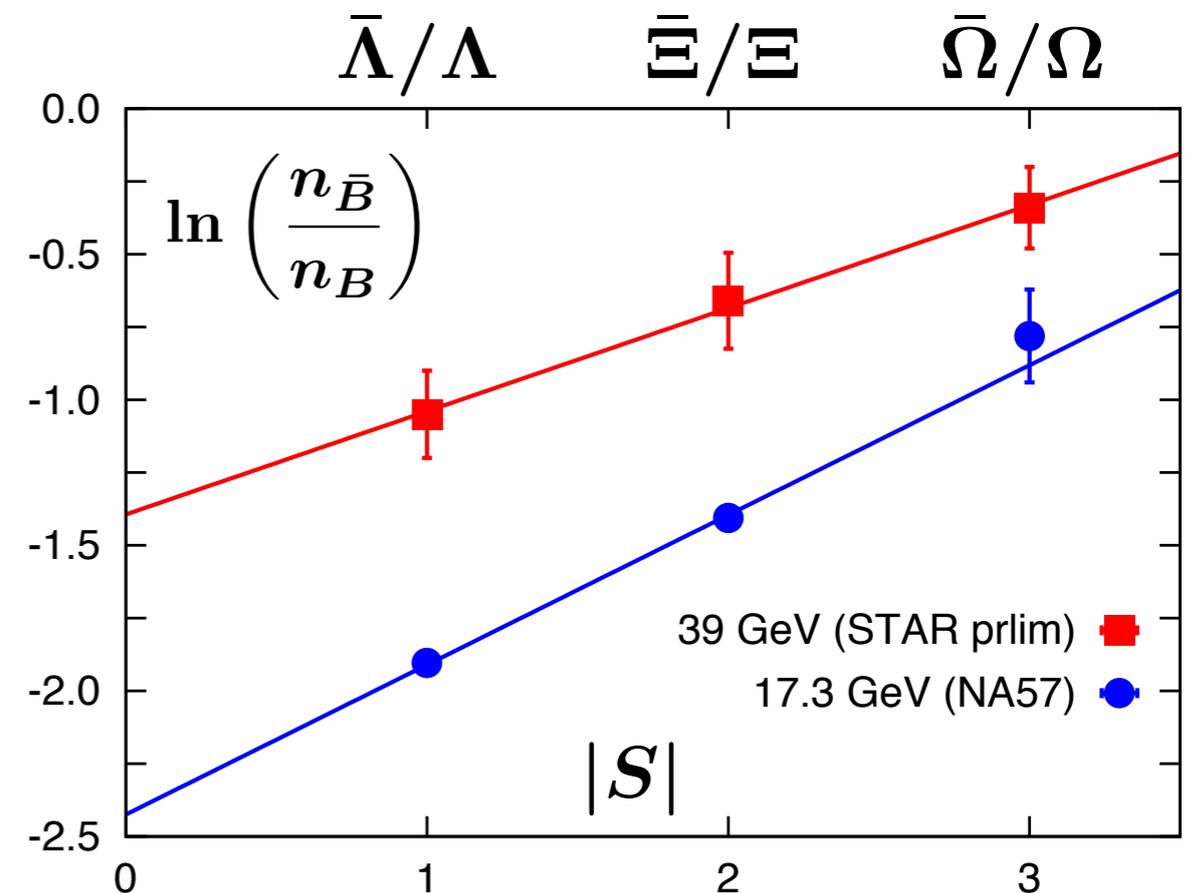
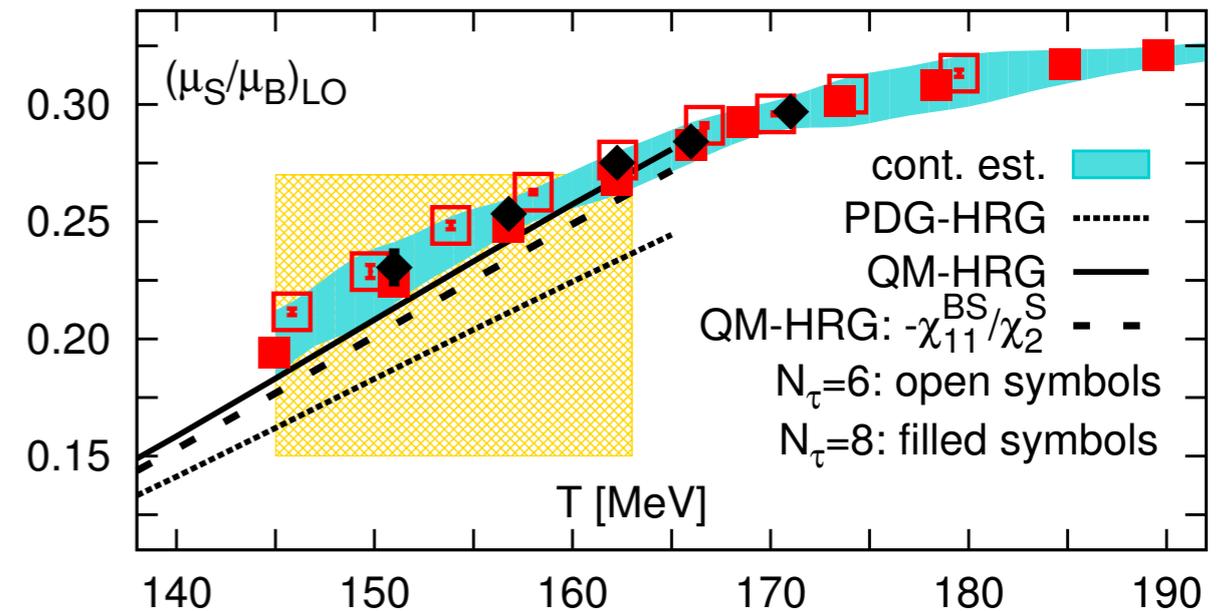
relative yields of strange anti-baryons to baryons at freeze-out are controlled by freeze-out parameter (T^f, μ_B^f, μ_S^f)

$$\frac{n_{\bar{B}}}{n_B} = \exp \left\{ -\frac{2\mu_B^f}{T^f} + \frac{2\mu_S^f}{T^f} |S| \right\}$$

$$= \exp \left\{ -\frac{2\mu_B^f}{T^f} \left(1 - \frac{\mu_S^f}{\mu_B^f} |S| \right) \right\}$$

$\Rightarrow \mu_B^f/T^f$ and μ_S^f/μ_B^f can be obtained by fitting experimentally measured values of $\bar{\Lambda}/\Lambda$, $\bar{\Xi}/\Xi$ and $\bar{\Omega}/\Omega$

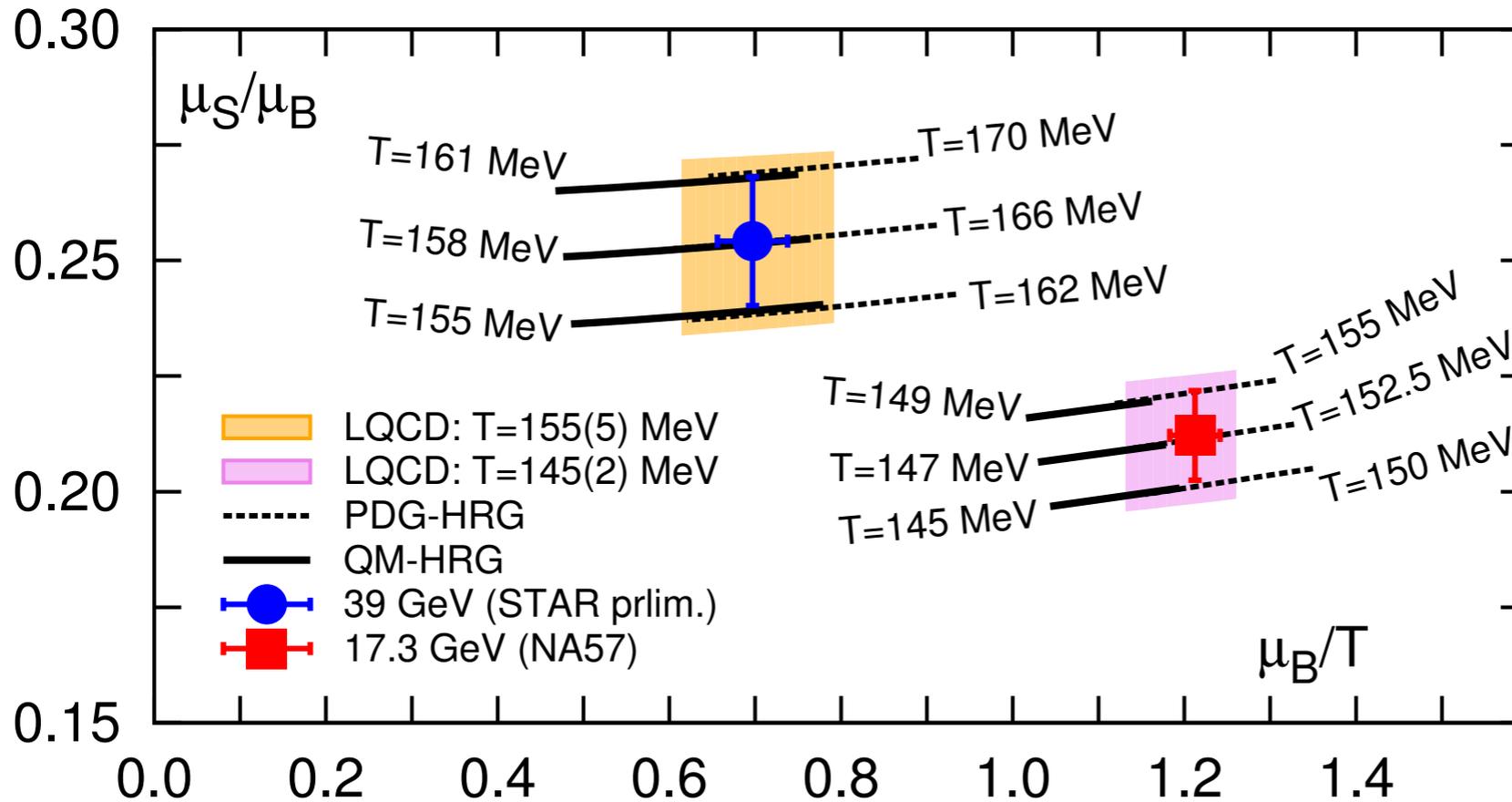
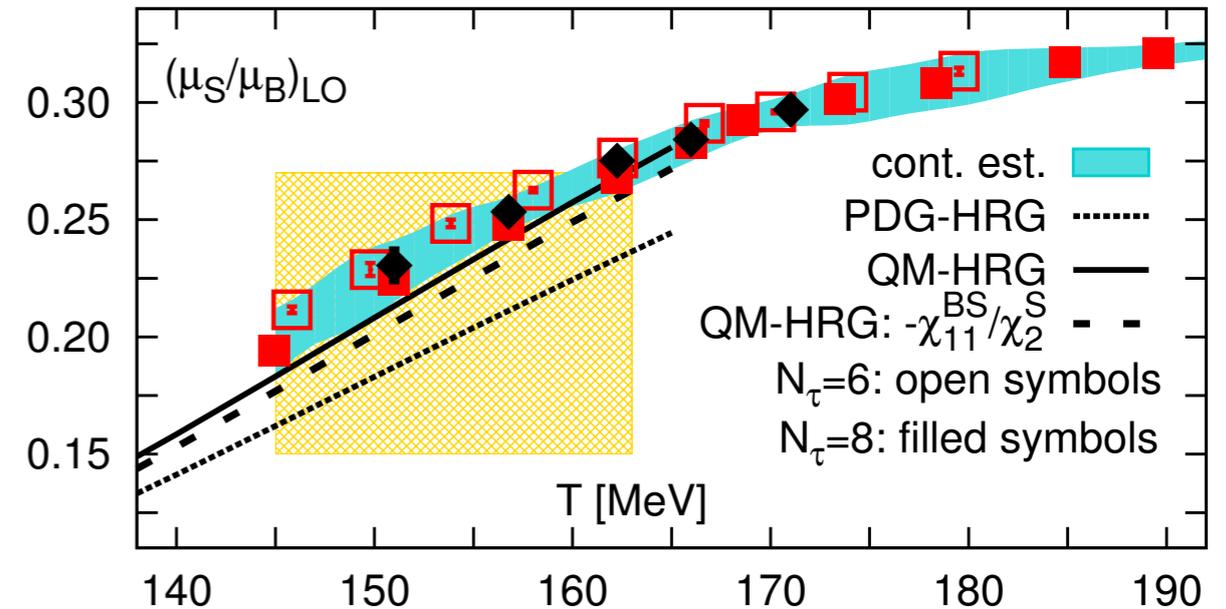
\Rightarrow fit function independent of details of open strange hadron spectrum



Implications for freeze-out conditions

Comparison:

vary T^f in order to match μ_B^f/T^f
and μ_S^f/μ_B^f



\Rightarrow QM-PDG in agreement with (lattice) QCD
 \Rightarrow QM-HRG yields 5-8 MeV smaller freeze-out temperature than PDG-HRG

Summary and Conclusions

- correlations of net baryon number and net strangeness are sensitive probes of the strangeness content of the system
- partial pressure of strange baryons from QM-HRG and PDG-HRG differ by about 30% at $T=170$ MeV
 - ⇒ additional strange baryons are thermodynamically relevant in the crossover region
- in the crossover region the QM-HRG provides a more accurate description of (lattice) QCD w.r.t. the conventionally used PDG-HRG
 - ⇒ evidence for additional, experimentally not yet observed, strange baryons
- presence of additional strange hadrons get imprinted in the yields of ground state strange hadrons
 - ⇒ significant reductions of freeze-out temperature of (5-8) MeV