

Phase diagram of lattice QCD in auxiliary field Monte-Carlo method in the strong coupling region

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Abstract: We study the QCD phase diagram in the strong coupling region by using the auxiliary field Monte-Carlo method. We have the sign problem in both AFMC and monomer-dimer-polymer (MDP) simulations when we include fluctuation effects. Although we have the sign problem, we could investigate the QCD phase diagram in AFMC, which is consistent with MDP simulations in the strong coupling limit. Preliminary results seem to suggest that finite coupling effects make the sign problem more severe than in the strong coupling limit.

Introduction

The sign problem in QCD phase diagram

- MC simulation in lattice QCD

$$Z = \int \mathcal{D}[\bar{\psi}, \psi, U] \exp[-\bar{\psi} \mathcal{D} \psi - S_G] \quad T$$

$$= \int \mathcal{D}[U] \det \mathcal{D} \exp[-S_G]$$

- Weight of MC integral becomes complex at finite real μ

- γ_5 Hermiticity

$$\gamma_5 \mathcal{D}(\mu) \gamma_5 = [\mathcal{D}(-\mu^*)]^\dagger$$

$$\det[\mathcal{D}(\mu)] = [\det \mathcal{D}(-\mu^*)]^*$$

- Partition function and sign problem

We calculate Z in the eigen states $|\psi\rangle$ of $H \rightarrow$ semi-positive definite \rightarrow no sign problem

We adopt states which are close to the eigen states of H

Hadronic degrees of freedom \rightarrow Reducing sign problem (?)

Strong coupling lattice QCD

- $1/g^2$ expansion in lattice QCD
- Integrating out link & Grassmann variables analytically
- MC integral by Hadronic DOF \rightarrow Reduced sign problem

"The sign problem" with fluctuations

- Monomer-Dimer-Polymer simulations

W. Unger, Ph. de Forcrand, J. Phys. G: Nucl. Part. Phys. 38 124190 (2011)

- Auxiliary field Monte-Carlo method

A. Ohnishi, T. I. and T. Z. Nakano: PoS LATTICE2012, 088 (2012). T. I., T. Z. Nakano, and A. Ohnishi arXiv:1311.5352

- We study the QCD phase diagram

in the strong coupling region

including fluctuation effects

by using auxiliary field Monte-Carlo (AFMC) method.



$$Z = \text{Tr} e^{-\beta H}$$

$$S_{\text{LQCD}} = S_{\text{F}} + S_{\text{G}}$$

Gluon field \rightarrow Link variables

$$U_\mu(x) \approx \exp(i g A_\mu)$$

Gluon action \rightarrow Plaquette action

$$S_G = \frac{2N_c}{g^2} \sum_{\text{plaq.}} \left[1 - \frac{1}{N_c} \text{Re tr } U_{\mu\nu}(n) \right]$$

$$F_{\mu\nu}^2 \sim U_\nu^\dagger(n) \frac{1}{g^2} U_\nu(n+\hat{\mu}) U_\mu(n) U_\mu^\dagger(n+\hat{\nu})$$

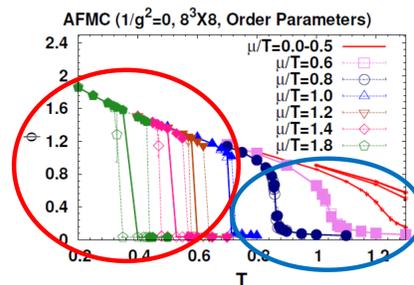
Results

- Reservations:

Chiral limit, Jack knife errors, All results in lattice unit

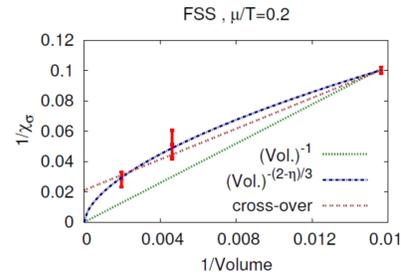
- Behavior of order parameters

- Chiral condensate ϕ



Would-be 1st at high μ Would-be 2nd at low μ

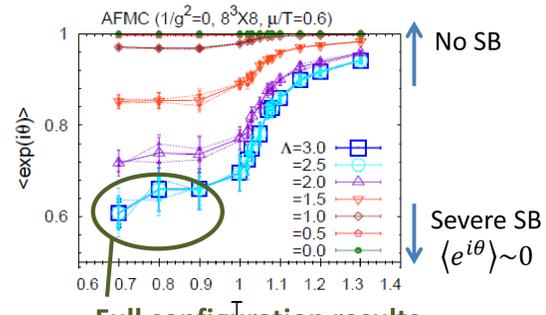
- Scaling behavior of chiral susceptibility at low μ



Would-be 2nd or cross over

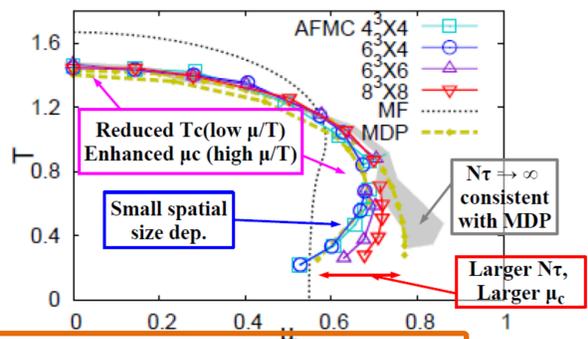
- Sign problem (SB)

- Average phase factor, $\langle e^{i\theta} \rangle$



Full configuration results

- The QCD phase diagram in the strong coupling limit



- The would-be 1st order at high μ and the would be 2nd at low μ

- NLO results

- Temporal corrections

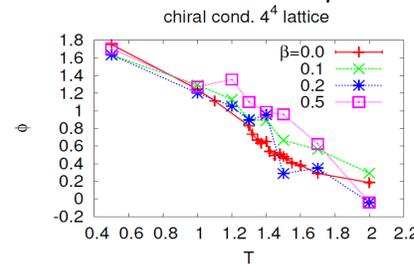
- Might be severe weight cancellation (SB) due to complex auxiliary fields.

$$\Delta \alpha_x^- = C_\tau \left\{ -i\omega_{x+\hat{j}}^* - i\omega_{x-\hat{j}}^* - (\epsilon\Omega^*)_{x+\hat{j}} - (\epsilon\Omega^*)_{x-\hat{j}} \right\}$$

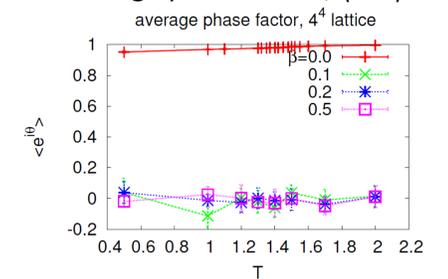
$$\Delta \alpha_x^+ = C_\tau \left\{ i\omega_{x+\hat{j}} + i\omega_{x-\hat{j}} + (\epsilon\Omega)_{x+\hat{j}} + (\epsilon\Omega)_{x-\hat{j}} \right\}$$

- Results of only spatial corrections

- Chiral condensate ϕ



- Average phase factor, $\langle e^{i\theta} \rangle$



- Reducing sign problem?

Effective mass term

$$m_x = m + \frac{1}{4N_c} \sum_j (\sigma + i\epsilon\pi) x_{\pm j}$$

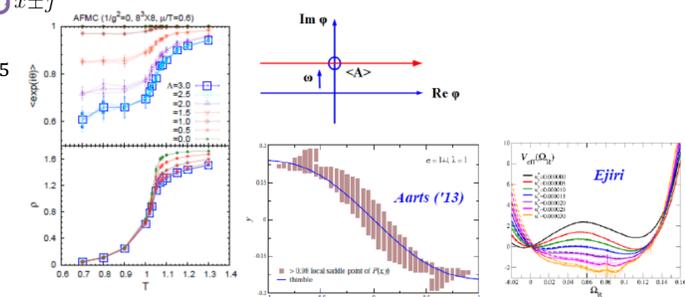
$$\epsilon = (-1)^{x_0 + \dots + x_d} \sim \gamma_5$$

- Imaginary terms do not contribute as long as low momentum mode contribute \rightarrow idea 1

- Idea 1: Cutoff or Gauss integral of high momentum modes

- Idea 2: Change the integral path

- Idea 3: Combination of Fugacity exp. or Histogram method



Formalism

Strong coupling lattice QCD action

- Unrooted staggered fermion, anisotropic lattice, strong coupling limit

- $1/d$ expansion, U_j integration N.Bilic et al. (1992), G. Faldt et al.(1986)

$$S_{\text{SCL-LQCD}} = \frac{1}{2} \sum_{\nu=0}^d \sum_x [\eta_{\nu,x}^+ \bar{\chi}_x U_{\nu,x} \chi_{x+\hat{\nu}} - \eta_{\nu,x}^- (\text{H.C.})] \quad \text{Kinetic}$$

$$+ m \sum_x \bar{\chi}_x \chi_x \quad \text{Mass}$$

$$\eta_{\nu,x}^\pm = \left[\gamma e^{\pm i\mu a_\tau}, (-1)^{x_1 + \dots + x_{\nu-1}} \right]$$

- Bosonization: Extended HS transformation

$$M_x M_{x+j} \rightarrow M_x \text{---} \times \text{---} M_{x+j}$$

$$\sigma_x = -\langle \bar{\chi}_x \chi_x \rangle$$

$$\pi_x = \langle i\epsilon \bar{\chi}_x \chi_x \rangle$$

$$\epsilon = (-1)^{x_0 + \dots + x_d} \sim \gamma_5$$

Notice!

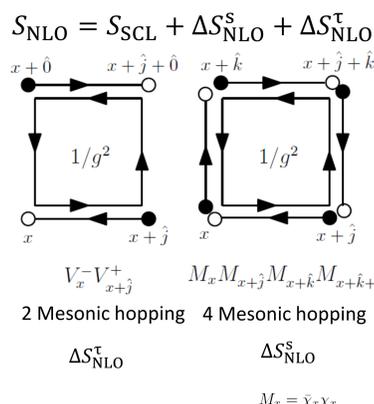
$$\exp[\alpha AB] = \int \mathcal{D}[\phi, \varphi] \exp[-\alpha(\phi^2 + \varphi^2 + (A+B)\varphi - i(A-B)\phi)]$$

- Auxiliary Field Monte-Carlo

$$Z = \int \mathcal{D}[\sigma_{\mathbf{k},\tau}, \pi_{\mathbf{k},\tau}] e^{-S_{\text{eff}}^{\text{AF}}}$$

NLO effective action

- Spatial plaquette terms
 - 8-fermion fields \rightarrow 4-fermion fields \rightarrow 2-fermion fields
 - Effective mass modification
- Temporal plaquette terms
 - 4-fermion fields \rightarrow 2-fermion fields
 - Effective chem. pot. + wave fn. ren.



We investigated chiral phase transition in AFMC and deduced phase transition order from the behavior of the chiral condensate and finite size scaling analysis. We also show the average phase factor in AFMC, which originate from bosonization procedure, and some NLO effects.

While we have the sign problem, we could show the QCD phase diagram in the strong coupling limit, which is consistent with MDP. Several prescriptions to weaken the sign problem were also discussed.