

Beam energy scan using a 3+1D viscous hydro+cascade model

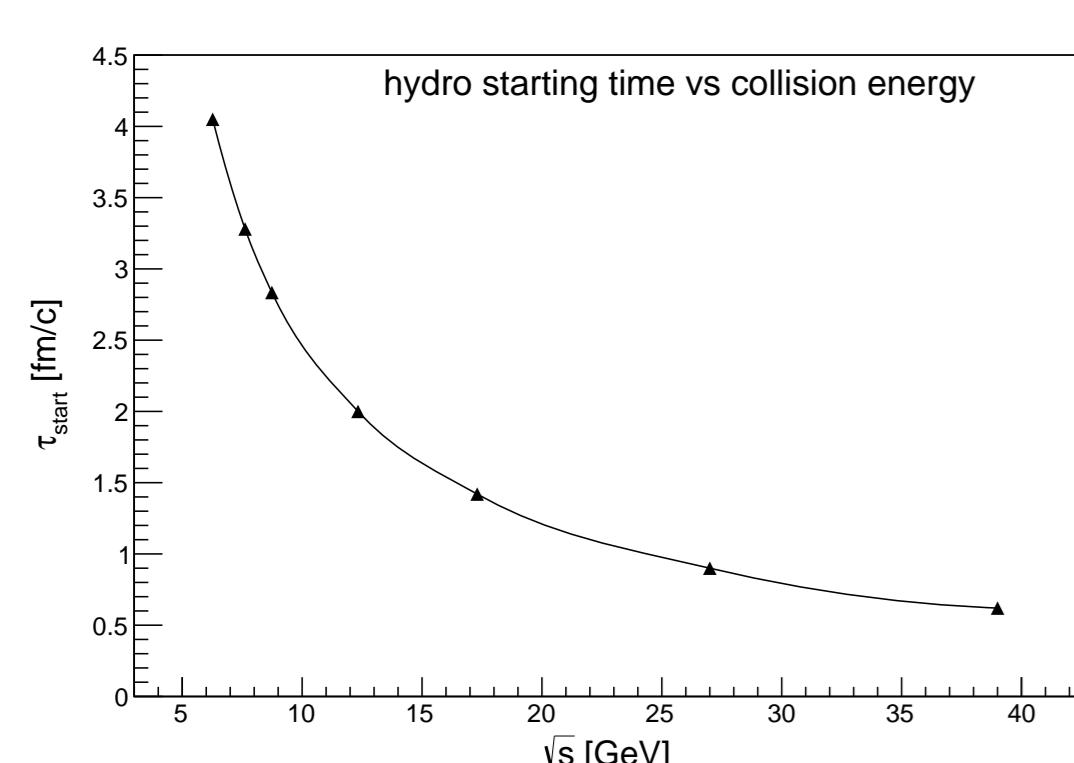
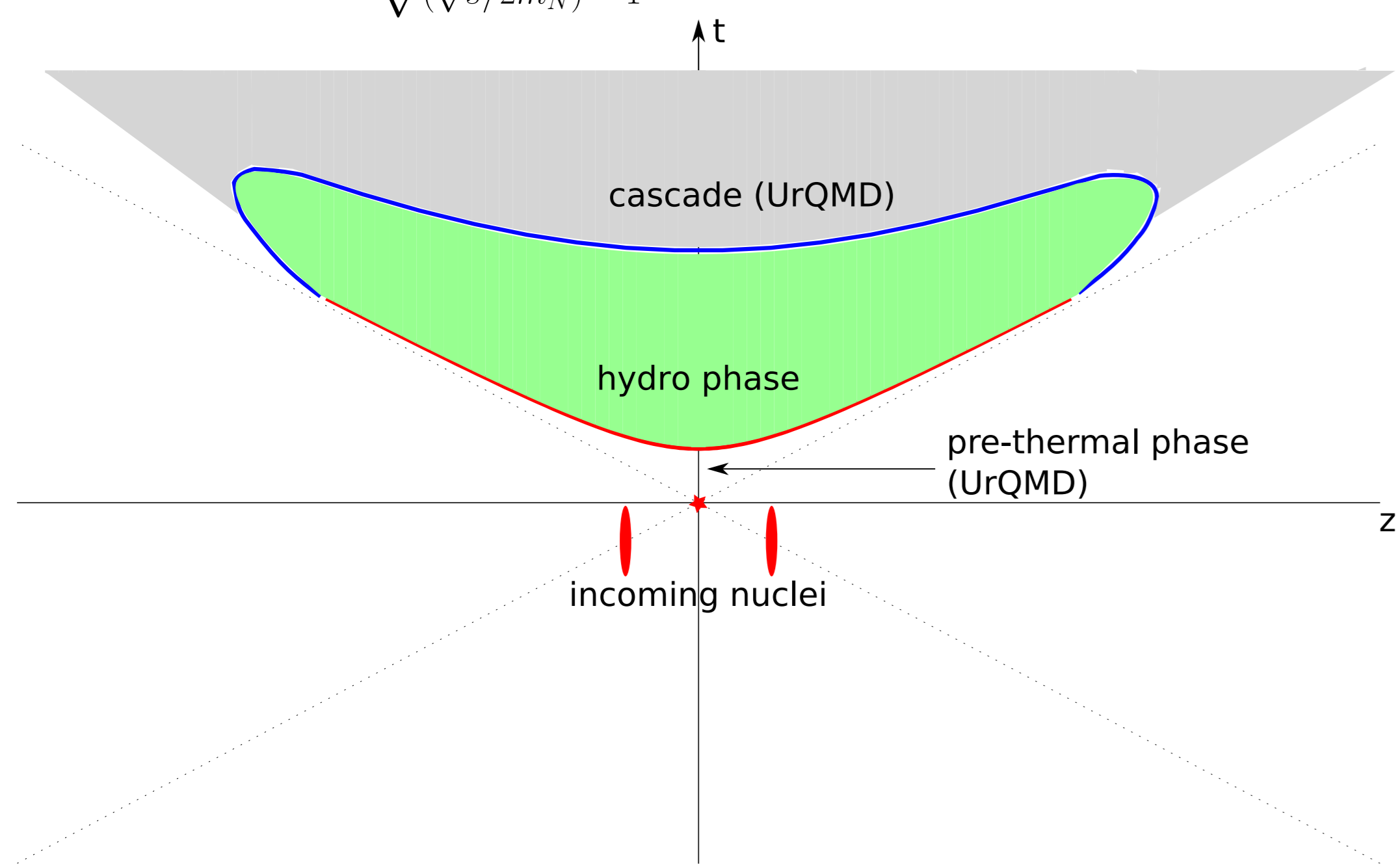
The model

Cascade-hydro-cascade approach:

- Initial state: UrQMD cascade [1]
- Hydrodynamic phase: numerical 3+1D hydro solution via original relativistic viscous hydro code [2]
- Hadronic cascade: UrQMD

Initial conditions for hydrodynamic evolution from UrQMD

Switch from UrQMD to fluid at Bjorken proper time $\tau = \sqrt{t^2 - z^2} = \tau_0$, where $\tau_0 = \frac{2R}{\gamma v_z} = \frac{2R}{\sqrt{(\sqrt{s}/2m_N)^2 - 1}}$. Switching surface is the red curve



Initial $\{T^{0\mu}, N_b^0, N_q^0\}$ of fluid = averaged or event-by-event $\{T^{0\mu}, N_b^0, N_q^0\}$ of particles
Initial shear stress tensor $\pi^{\mu\nu} = 0$.

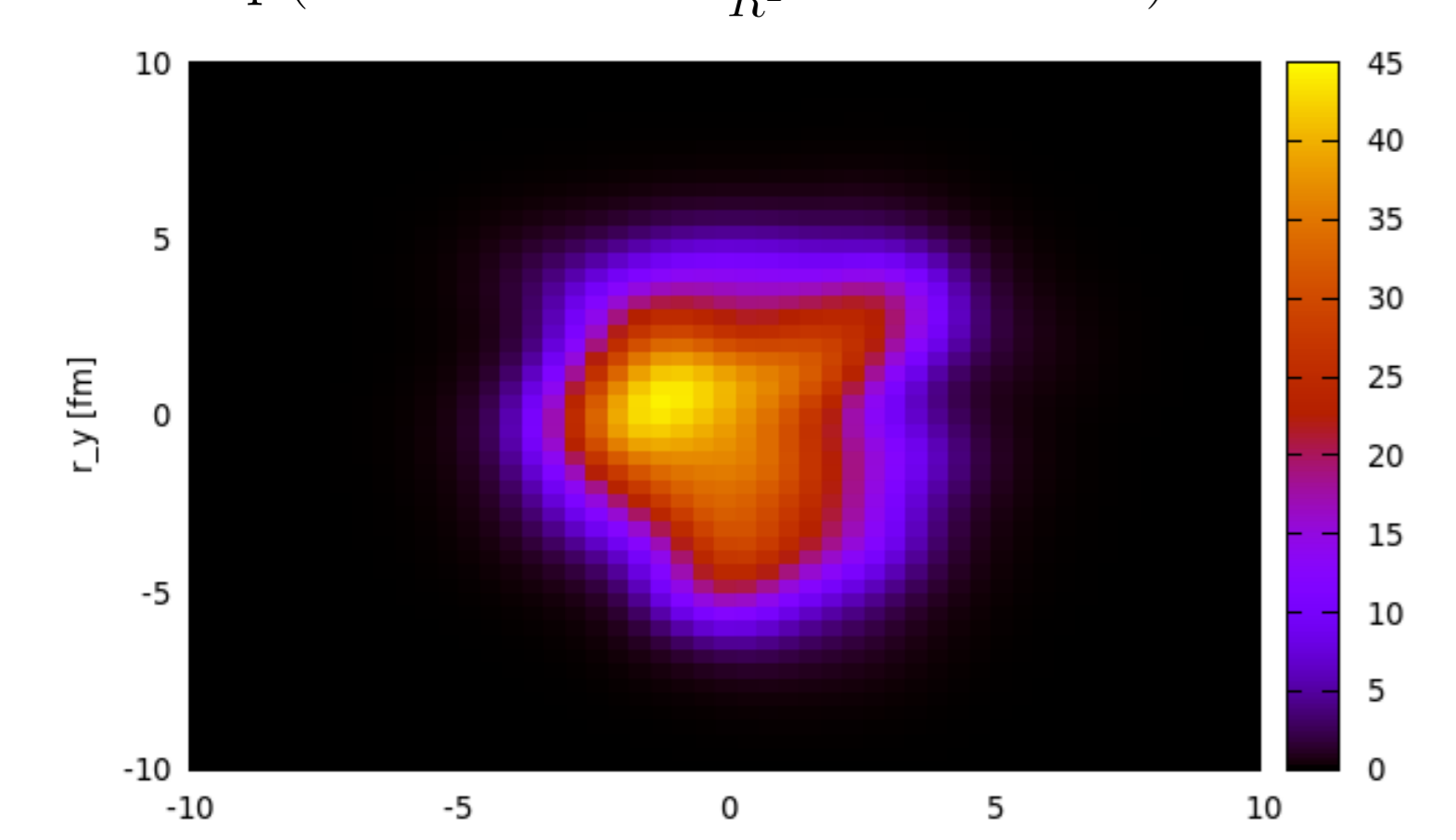
Option 1) Averaged initial state

Initial particle distribution is taken as an average over $\propto 10^4$ UrQMD simulations of initial state. No smoothing involved.

Option 2) Fluctuating initial state

Fluctuating, but smoothed initial state [6]:

$E \propto \exp(-\frac{(x-x_{part})^2 + (y-y_{part})^2 + \gamma_z^2(z-z_{part})^2}{R^2})$, where $R = 1.4$ fm



Fluctuating IC with $R = 1.4$ fm yields 23% larger average $dS/dy(y=0)$ than the averaged IC.

Equation of state

The equation of state from Chiral model [3] is used, which agrees qualitatively with lattice QCD results for zero baryon density. However the EoS is also constructed for finite (large) baryon densities.

Hydrodynamic phase

Numerical 3+1D relativistic viscous hydro solution in Israel-Stewart formalism and Milne coordinates is used. The evolutionary equations for shear stress tensor are:

$$\langle u^\gamma \partial_\gamma \pi^{\mu\nu} \rangle = -\frac{\pi^{\mu\nu} - \pi_{NS}^{\mu\nu}}{\tau_\pi} - \frac{4}{3} \pi^{\mu\nu} \partial_\gamma u^\gamma$$

▷ Bulk viscosity $\zeta = 0$, charge diffusion=0

▷ Shear relaxation time ansatz used: $\tau_\pi = 3\eta/(sT)$

Fluid→particle transition

$\epsilon = \epsilon_{sw} = 0.5$ GeV/fm³ (blue curve):

$\{T^{0\mu}, N_b^0, N_q^0\}$ of hadron-resonance gas = $\{T^{0\mu}, N_b^0, N_q^0\}$ of fluid

▷ Cooper-Frye prescription for hadron sampling:

$$p^0 \frac{d^3 n_i}{d^3 p} = \sum f_{i,eq}(x, p) \left[1 + (1 \mp f_{eq}) \frac{p_\mu p_\nu \pi^{\mu\nu}}{2T^2(\epsilon + p)} \right] p^\mu \Delta \sigma_\mu$$

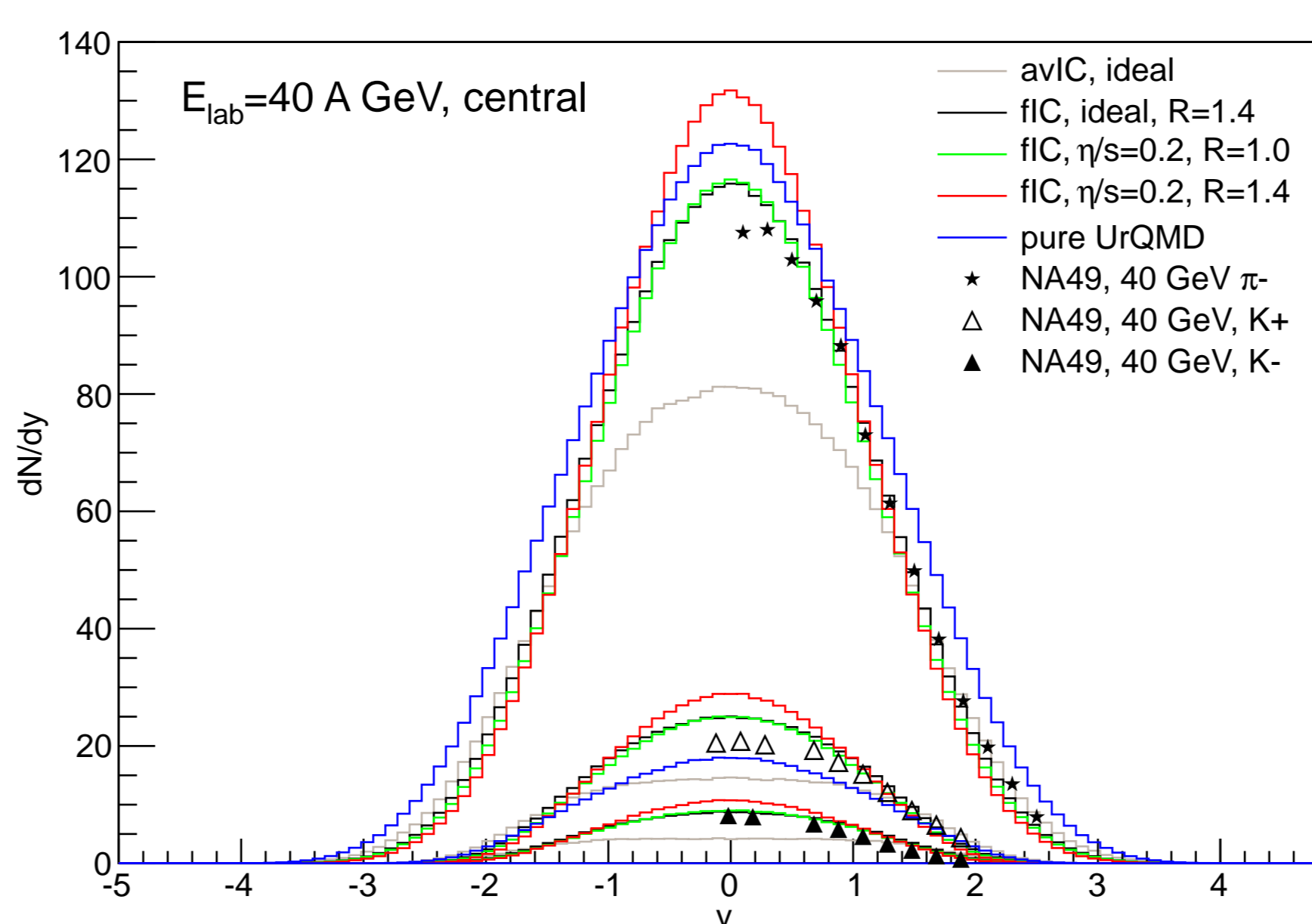
▷ Cornelius subroutine [4] to compute $\Delta \sigma_i$ on transition hypersurface.

▷ UrQMD cascade is employed after particlization surface.

Abstract

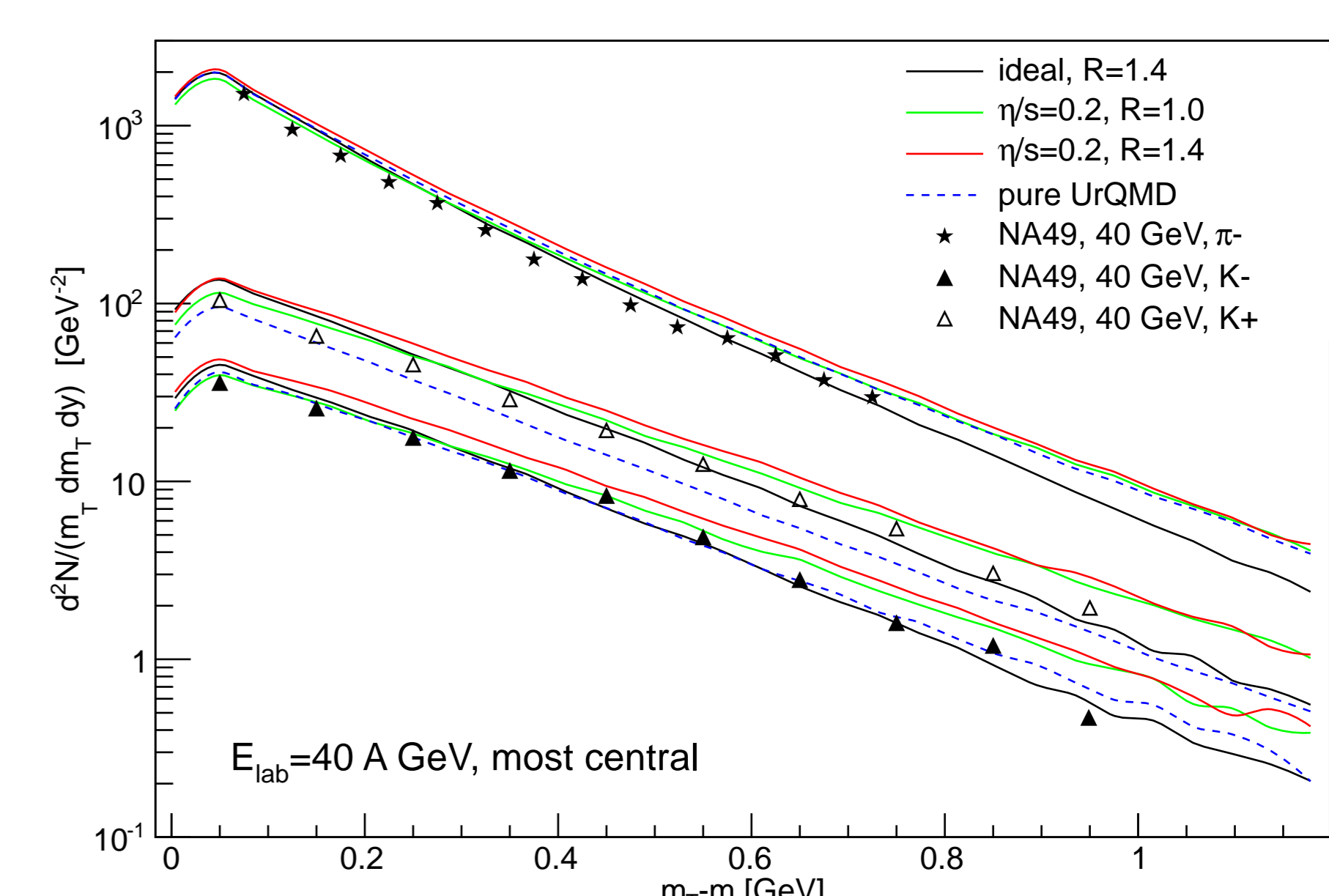
We apply a 3+1D viscous hydro+cascade model for A+A collisions at RHIC Beam Energy Scan energies ($\sqrt{s} = 7.7 - 39$ GeV), as well as for SPS energy points. We show how the results are sensitive to the shear viscosity in hydrodynamic phase and estimate η/s for Au+Au collisions in RHIC BES.

Rapidity distributions



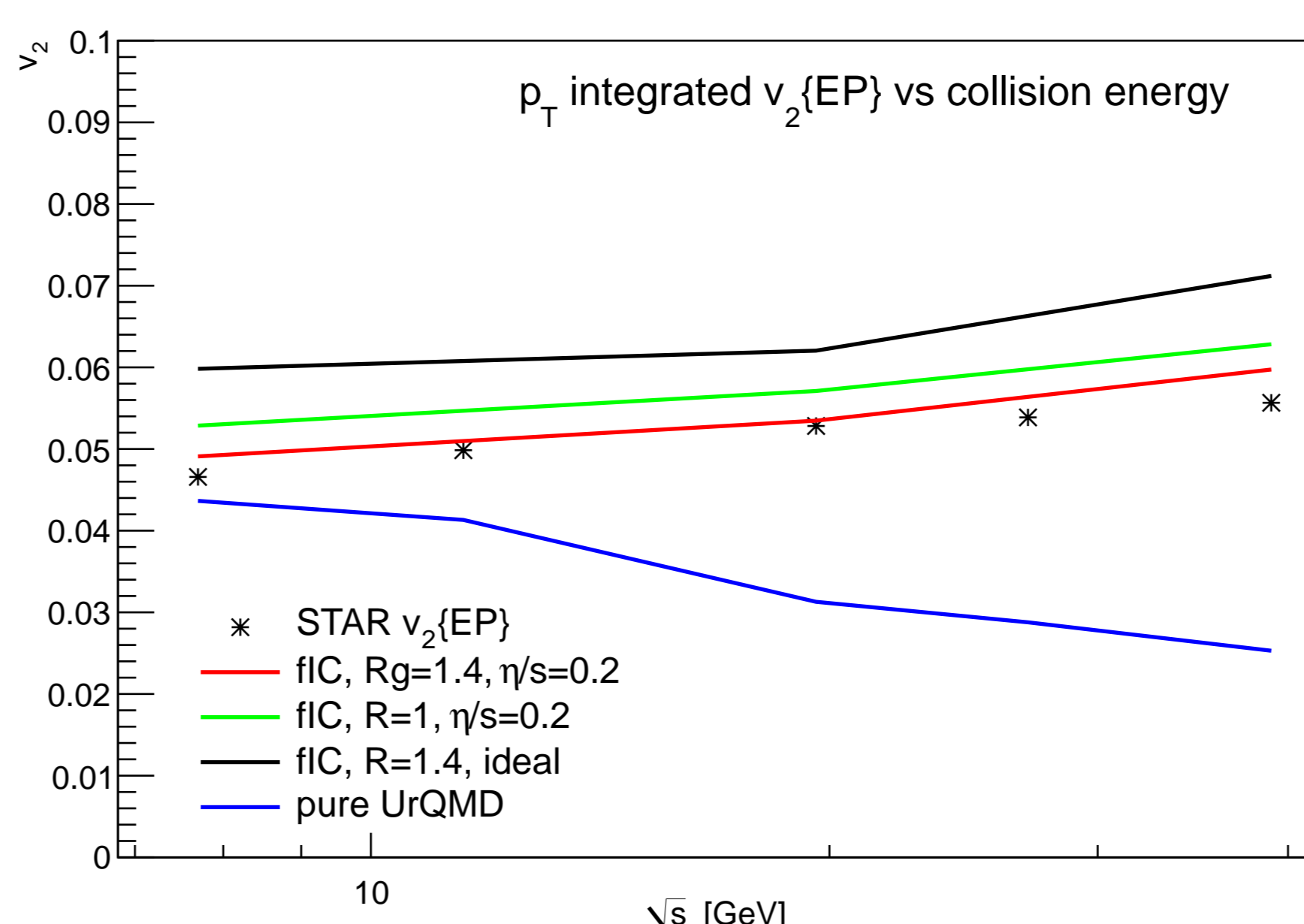
- shear viscosity in hydrodynamic phase makes overall expansion more spherical, bringing extra energy in transverse expansion at midrapidity. This increases both the multiplicity at midrapidity and the effective temperature of p_T spectra

Transverse momentum spectra



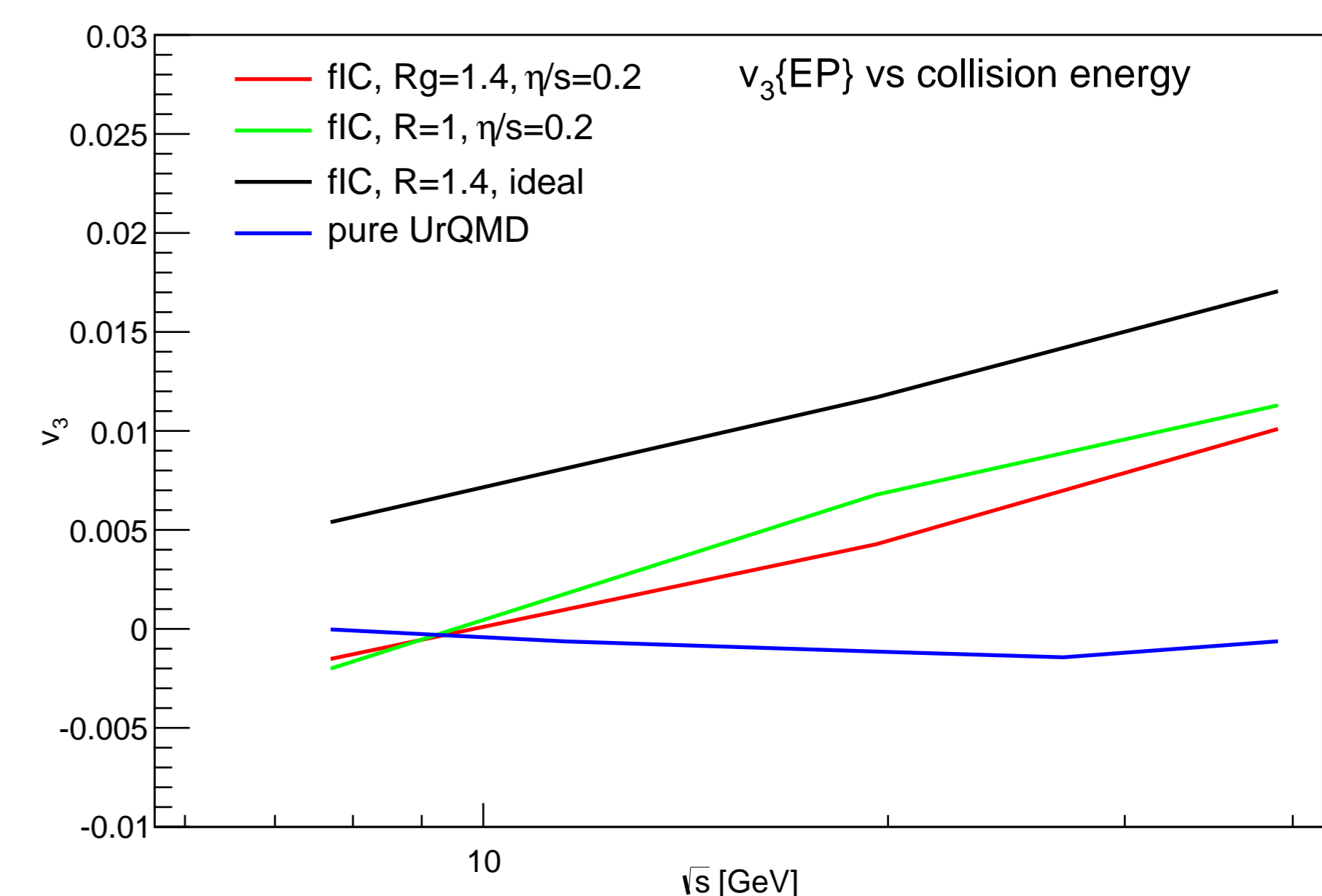
- larger initial entropy for smoothed fluctuating initial state leads to larger final multiplicity
- broader Gaussian smearing of the initial state has an effect on observables, which is qualitatively similar to the increase of shear viscosity in hydrodynamic phase.

p_T integrated elliptic flow at BES energies



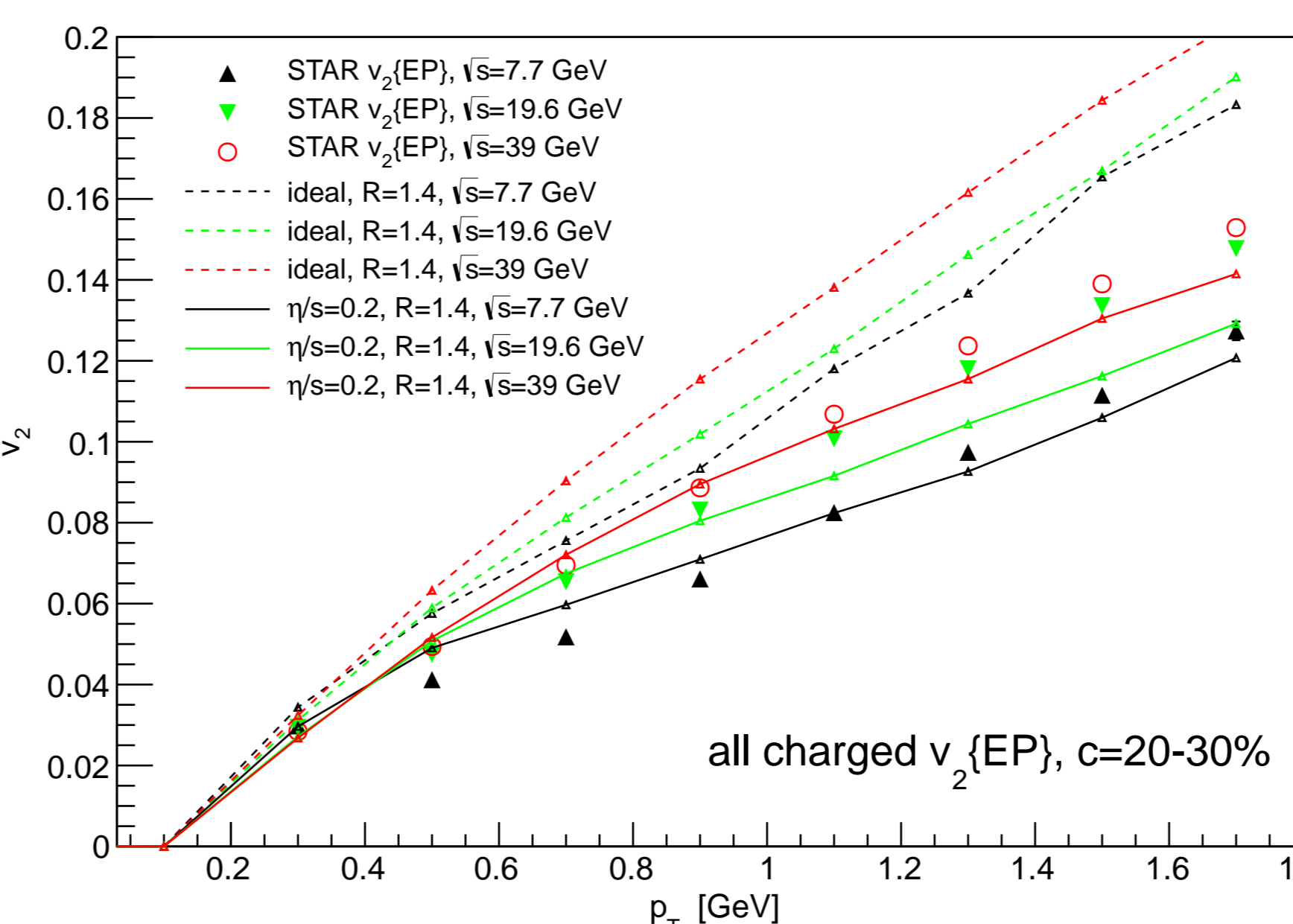
- Comparison to the experimental v_2 from RHIC BES suggests $\eta/s \geq 0.2$ in hydro phase for BES collision energy range

p_T integrated triangular flow

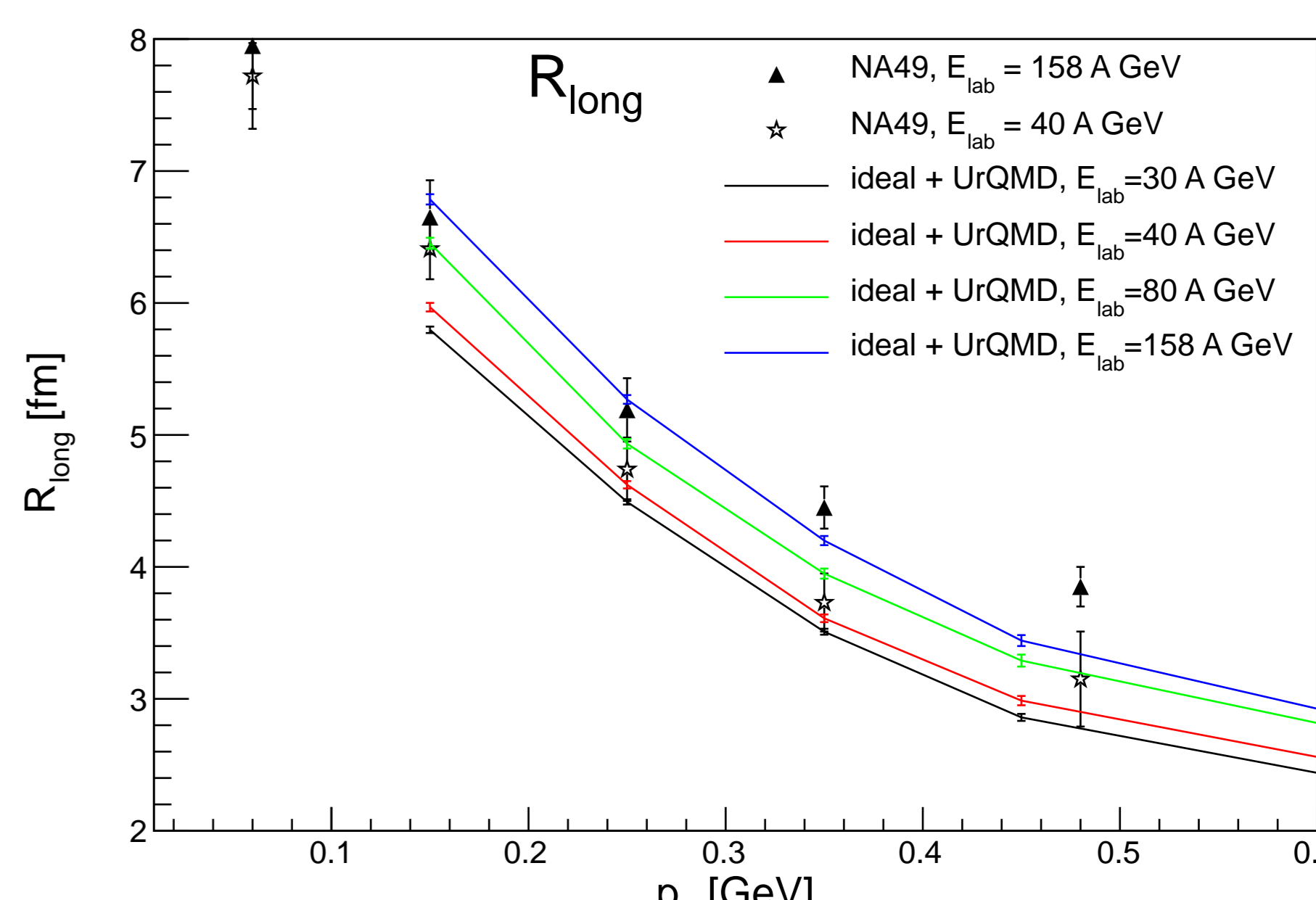


- v_3 is strongly sensitive to the shear viscosity in hydro phase
- pure cascade (UrQMD) shows $v_3 \approx 0$

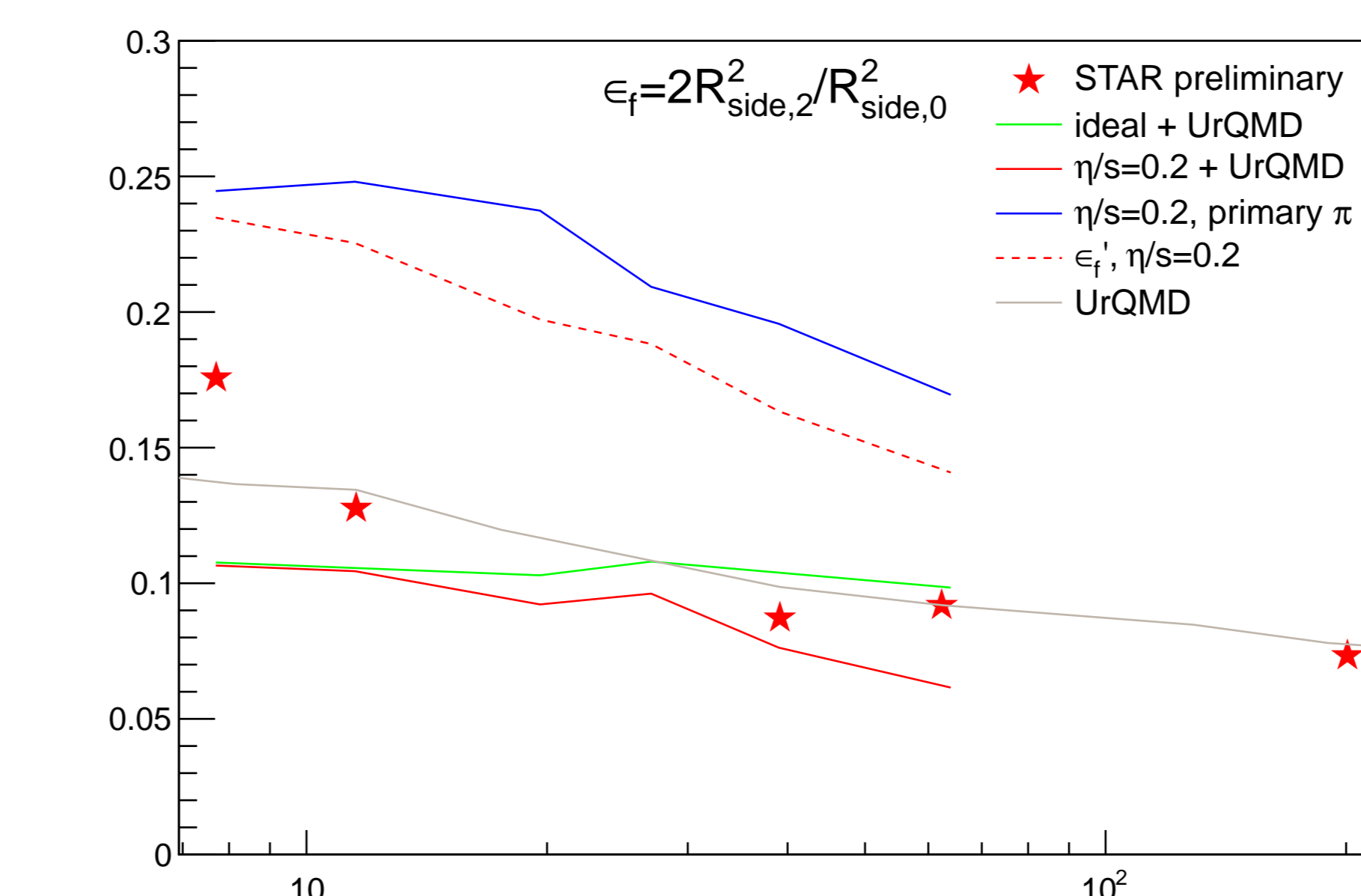
p_T differential elliptic flow



Femtoscopic radius R_{long} (averaged IC only)



- $\eta/s = 0.2$ in hydro phase systematically increases R_{long} by 5-10%, does not affect R_{out}/R_{side} .



Freezeout eccentricity from aZHBT (averaged IC only) for 10-30% central Au+Au, $p_T = 0.15 \dots 0.6$ GeV
Dashed curve: $\epsilon' = \frac{\int (y^2 - x^2) u^\mu d\sigma_\mu}{\int (y^2 + x^2) u^\mu d\sigma_\mu}$ (directly from the freezeout geometry)

References

- [1] S.A. Bass et al., *Prog. Part. Nucl. Phys.* **41** 255-369, 1998
- [2] Iu. Karpenko, P. Huovinen, M. Bleicher, arXiv:1312.4160
- [3] J. Steinheimer, S. Schramm and H. Stocker, *J. Phys. G* **38** 35001, 2011
- [4] P. Huovinen and H. Petersen, *Eur.Phys.J. A* **48** 171, 2012
- [5] N.S. Amelin et al, *Phys.Rev.C* **74**:064901, 2006
- [6] H. Petersen et al., *Phys.Rev. C* **78** 044901, 2008