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Quark production, Bose-Einstein condensates and thermalization of the quark-gluon plasma

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In this talk, I would like to report the results of our recent work on the thermalization of gluons and N_f flavors of massless quarks and antiquarks in a spatially homogeneous

system. First, we give two coupled transport equations for gluons and quarks (and antiquarks), which are derived within the diffusion approximation of

the Boltzmann equation with only $2\leftrightarrow 2$ processes included in the collision term. These transport equations are solved

numerically in order to study the thermalization of the quark-gluon plasma. Next, we discuss three different patterns of the thermalization of the quark-gluon system. At initial time, we assume that no quarks or antiquarks are present

and we choose the gluon distribution in the form $f = f_0 \theta \left(1 - \frac{p}{Q_s}\right)$ with Q_s the saturation momentum and f_0 a constant. The subsequent evolution of systems

may, or may not, lead to the formation of a (transient) Bose condensate, depending on the value of f_0 . The three patterns of thermalization are as follows: (a) thermalization

without gluon Bose-Einstein condensates (BEC) for $f_0 \leq f_{0t}$,

(b) thermalization with transient BEC for $f_{0t} < f_0 \le f_{0c}$ and (c) thermalization with BEC for $f_{0c} < f_0$. Here, the values of f_{0t} and f_{0c} depend on N_f . When $f_0 \ge 1 > f_{0c}$, the formation of BEC starts at a finite time $t_c \sim \frac{1}{(\alpha_s f_0)^2} \frac{1}{Q_s}$. We also find that the equilibration time

for $N_f = 3$ is typically about 5 to 6 times longer than that for $N_f = 0$ at the same Q_s .

On behalf of collaboration:

None

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