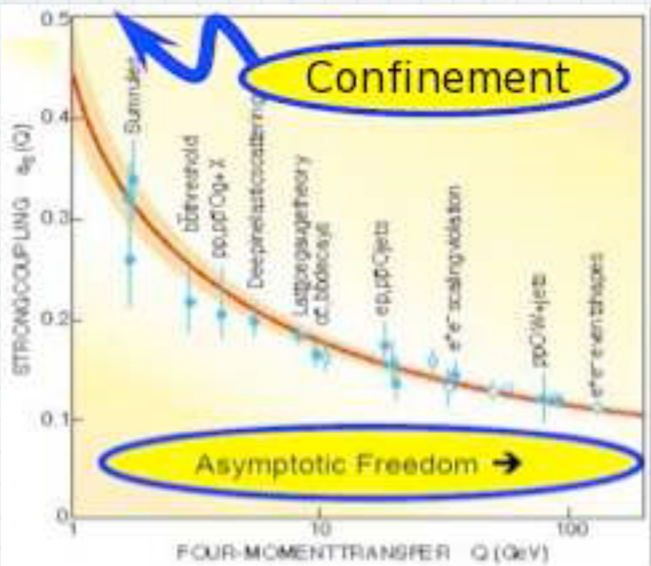
Phase diagram, fluctuations, thermodynamics and hadron chemistry

Observables and concepts

Claudia Ratti
University of Torino and INFN Torino (ITALY)

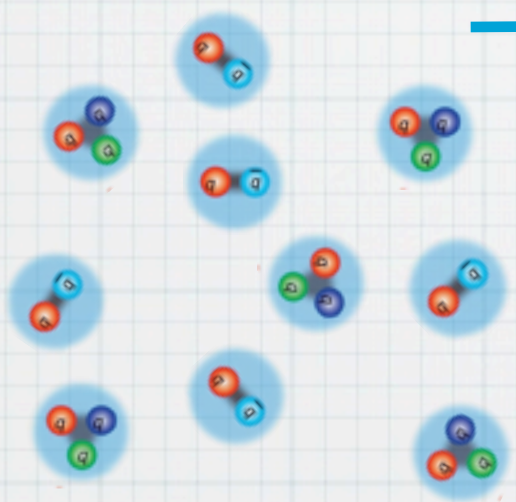
Quark Matter Student day - 18/05/2014

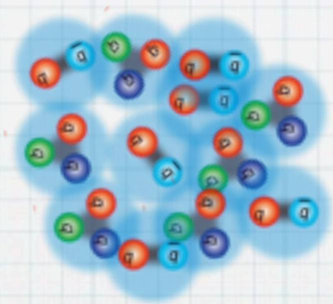
QCD Thermodynamics



* Confinement
  * At large distances the effective coupling is large
  * Free quarks are not observed in nature
* Asymptotic freedom
  * At short distances the effective coupling decreases
  * Quarks and gluons appear to be quasi-free

$T,\ \mu_B$





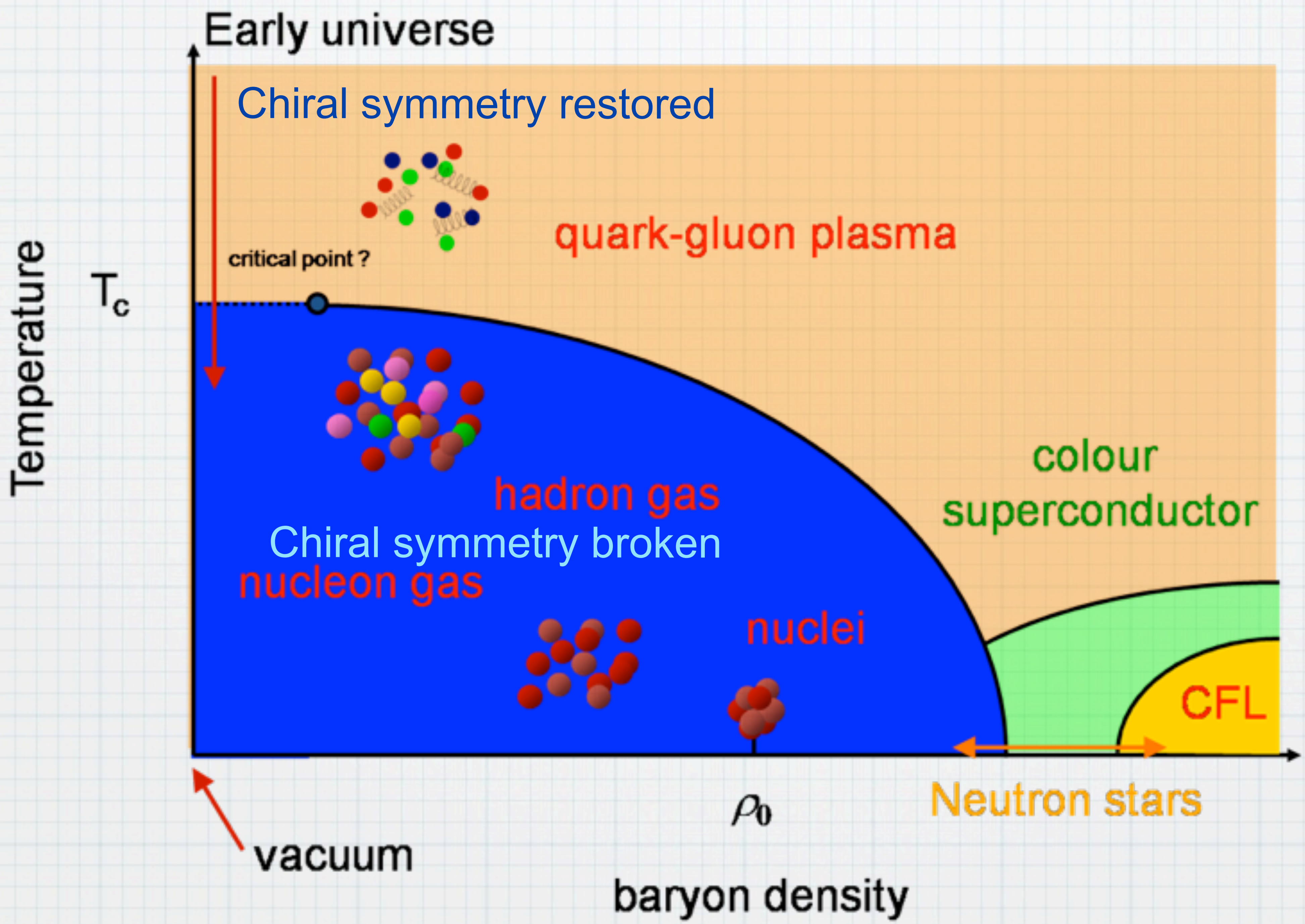
Hadron Gas

Quark-Gluon Plasma

Chiral Symmetry: broken

Chiral Symmetry: restored

QCD Phase Diagram



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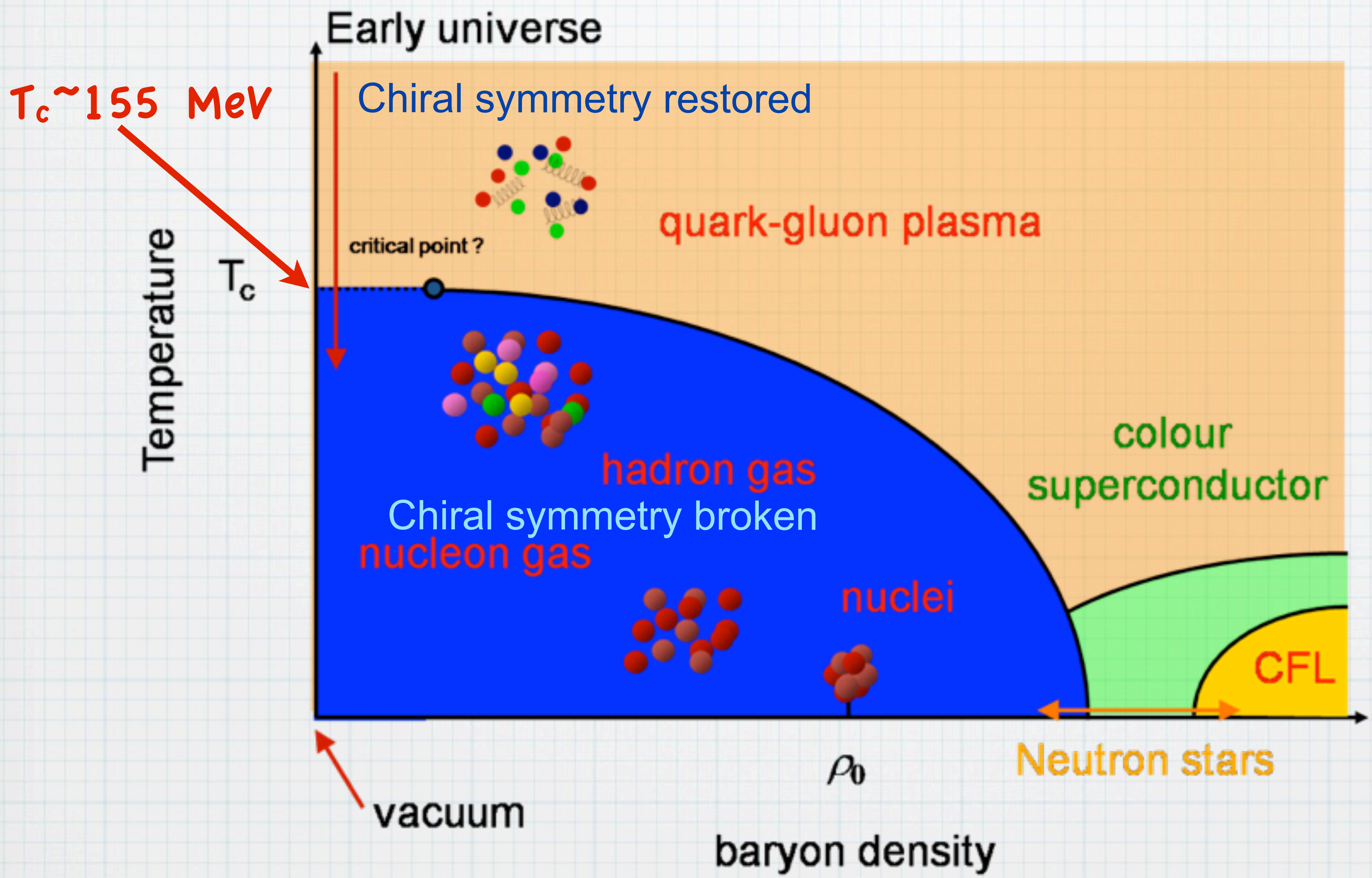

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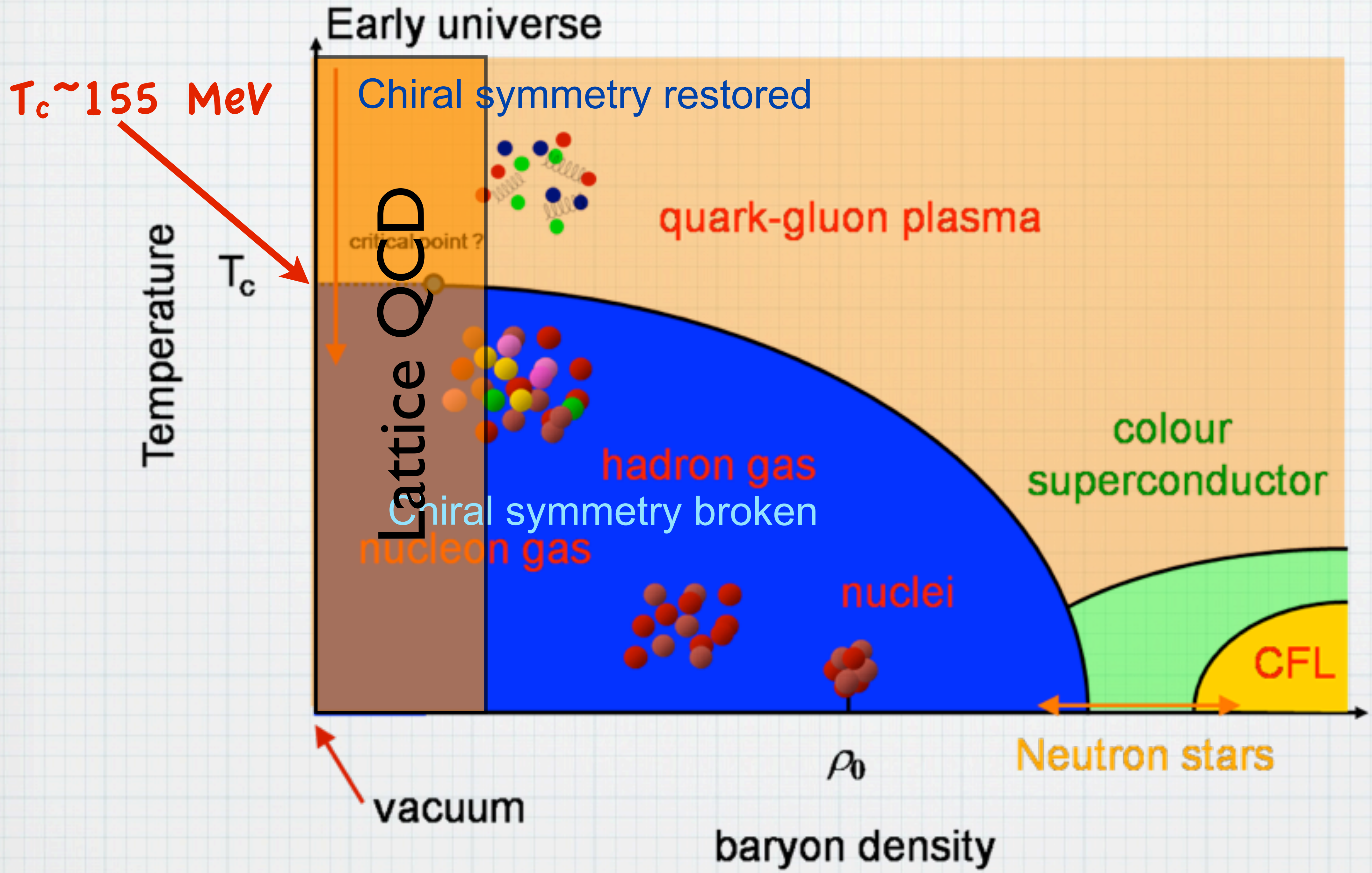

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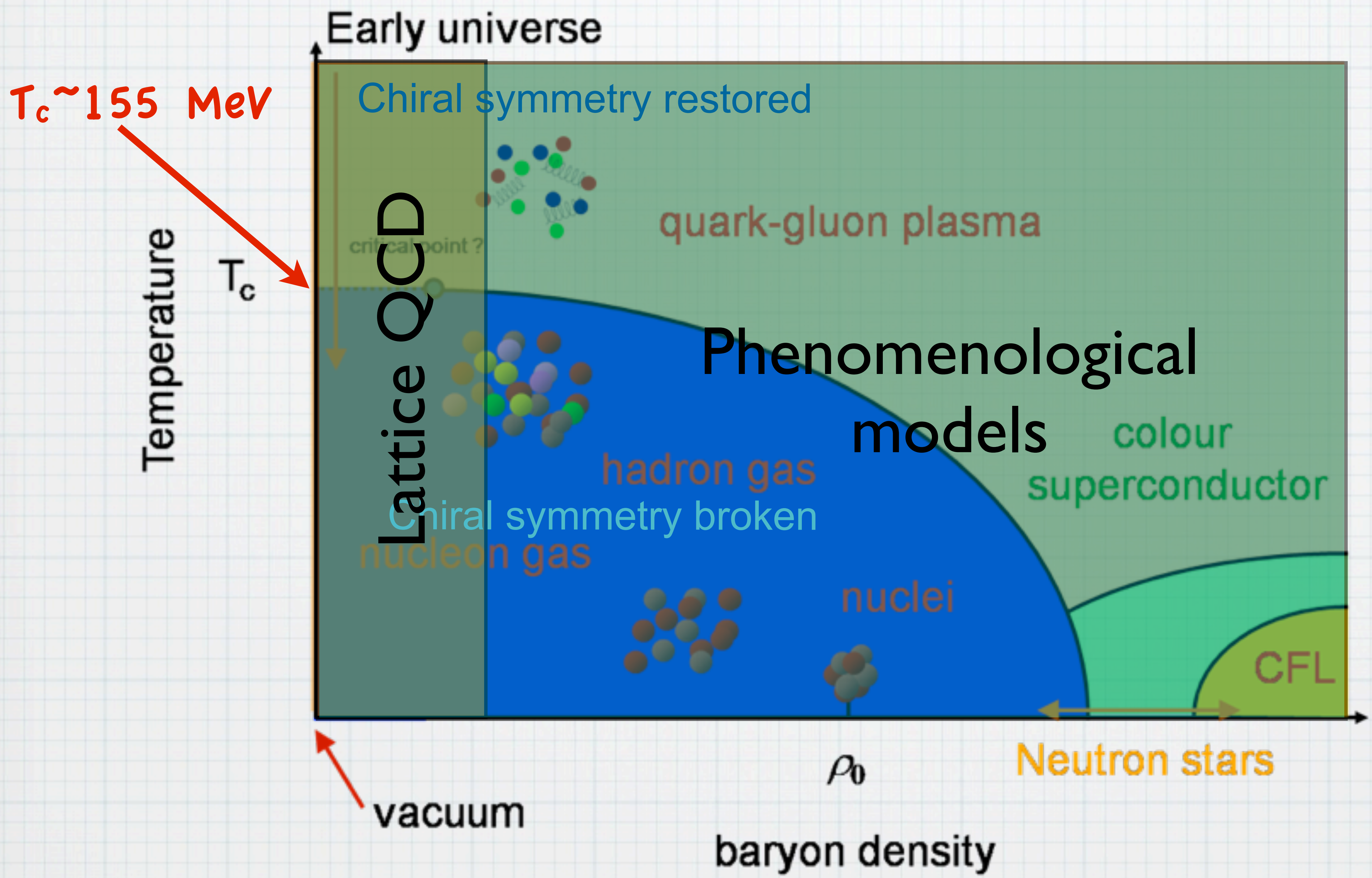


Lattice QCD

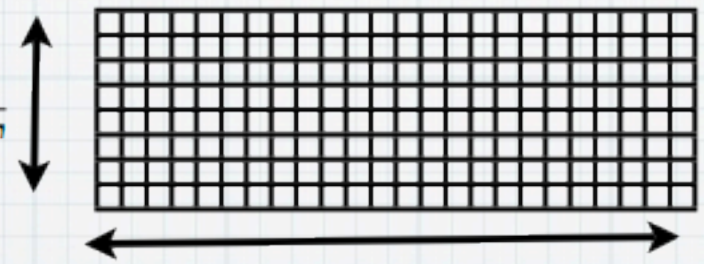
- Analytic or perturbative solutions **in low-energy QCD** are hard or impossible due to the **highly nonlinear nature** of the strong force

- Lattice QCD: well-established **non-perturbative approach** to solving QCD

- Solving QCD on a grid of points in space and time

- The lattice action is the parameterization used to discretize the Lagrangian of QCD on a space-time grid

$$N_t = \frac{1}{aT}$$



$$N_s = L/a$$

- From the partition function Z, knowledge of all the thermodynamics

$$F = -T \ln Z \,,$$

$$p = \frac{\partial (T \ln Z)}{\partial V} \,,$$

$$S = \frac{\partial (T \ln Z)}{\partial T} \,,$$

$$\bar{N}_i = \frac{\partial (T \ln Z)}{\partial \mu_i} \,,$$

$$E = -pV + TS + \mu_i \bar{N}_i$$

Sign problem

* The QCD path integral is computed by **Monte Carlo algorithms** which samples field configurations with a weight proportional to the **exponential of the action**

$$Z(\mu_B, T) = \mathrm{Tr}\left(e^{-\frac{H_{\mathrm{QCD}} - \mu_B N_B}{T}}\right) = \int \mathcal{D}U e^{-S_G[U]} \det M[U, \mu_B]$$

$\det M[\mu_B]$ **complex** $\Longrightarrow$ **Monte Carlo simulations are not feasibile.**

* If the action is **complex**, its exponential is **oscillating**: it cannot be used as a probability

* This is the reason why lattice QCD simulations cannot presently be performed at finite chemical potential

* Possible solutions:

  - ➡ Taylor expansion around μ_B=0
  - ➡ Imaginary chemical potential
  - ➡ Reweigthing techinque

All valid at small chemical potentials

Phase transitions and order parameters

- We want to study the transition from hadrons to the QGP: deconfinement and chiral symmetry restoration

- A phase transition is the transformation of a thermodynamic system from one phase or state of matter to another

- During a phase transition of a given medium certain properties of the medium change, often discontinuously, as a result of some external conditions

- The measurement of the external conditions at which the transformation occurs is called the phase transition point

- Order parameter: some observable physical quantity that is able to distinguish between two distinct phases

- We need to find observables which allow us to distinguish between confined/deconfined system and between chirally broken/restored phase

Polyakov loop: deconfinement

* We consider a system of gluons in which we put a heavy quark-antiquark pair as a probe

$$\langle \Phi \rangle \sim e^{-F/T}$$

➡ How much energy F is needed to extract the heavy quark from the system?

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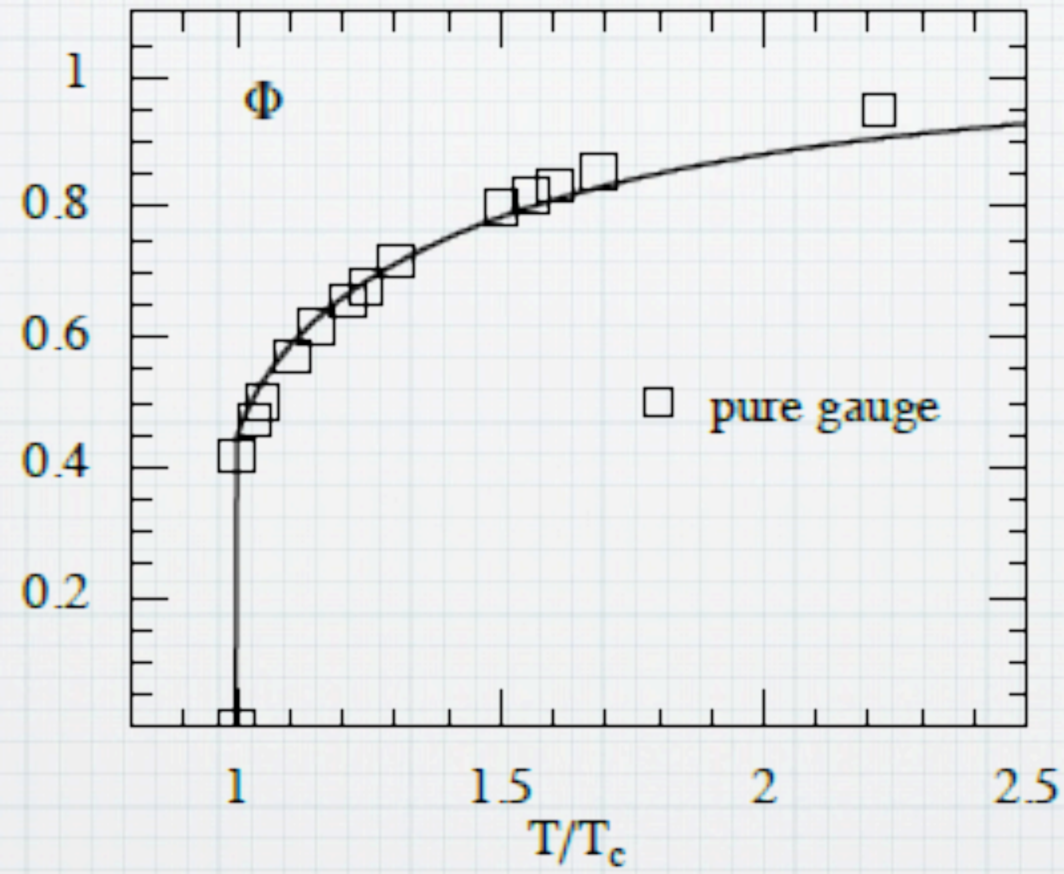

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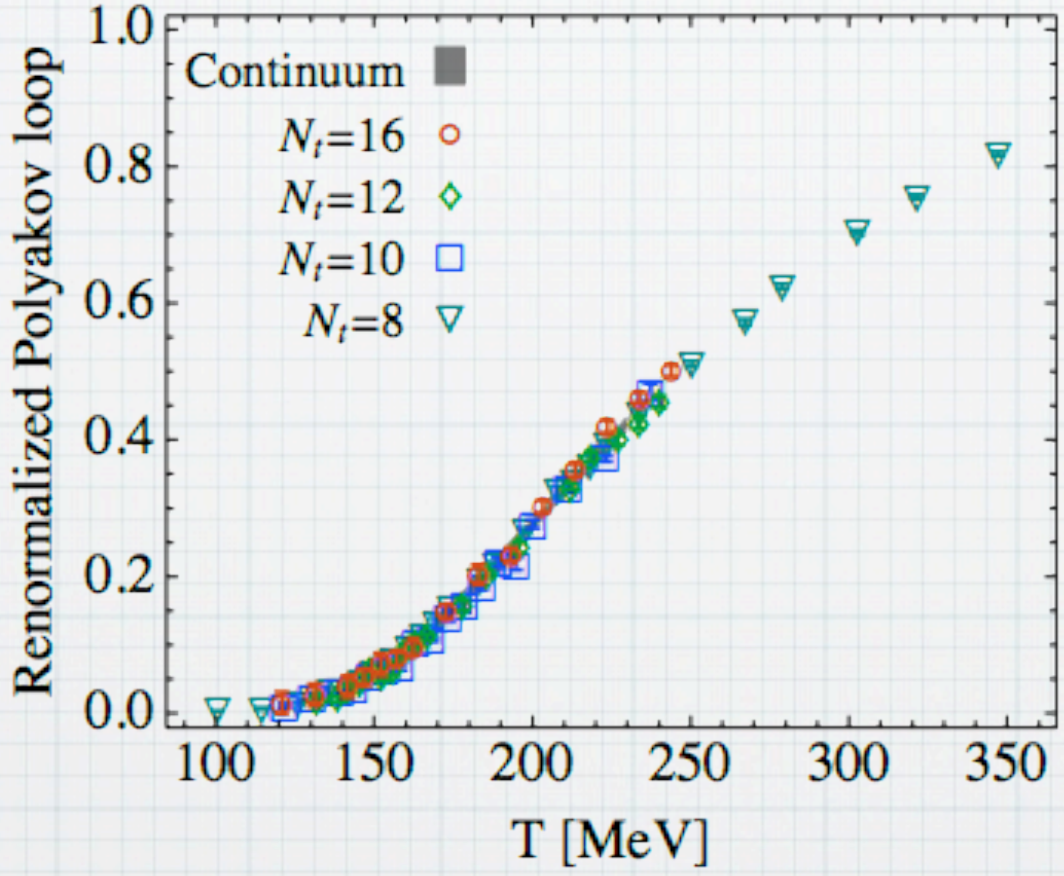

For QCD with physical quark masses the transition is a smooth crossover

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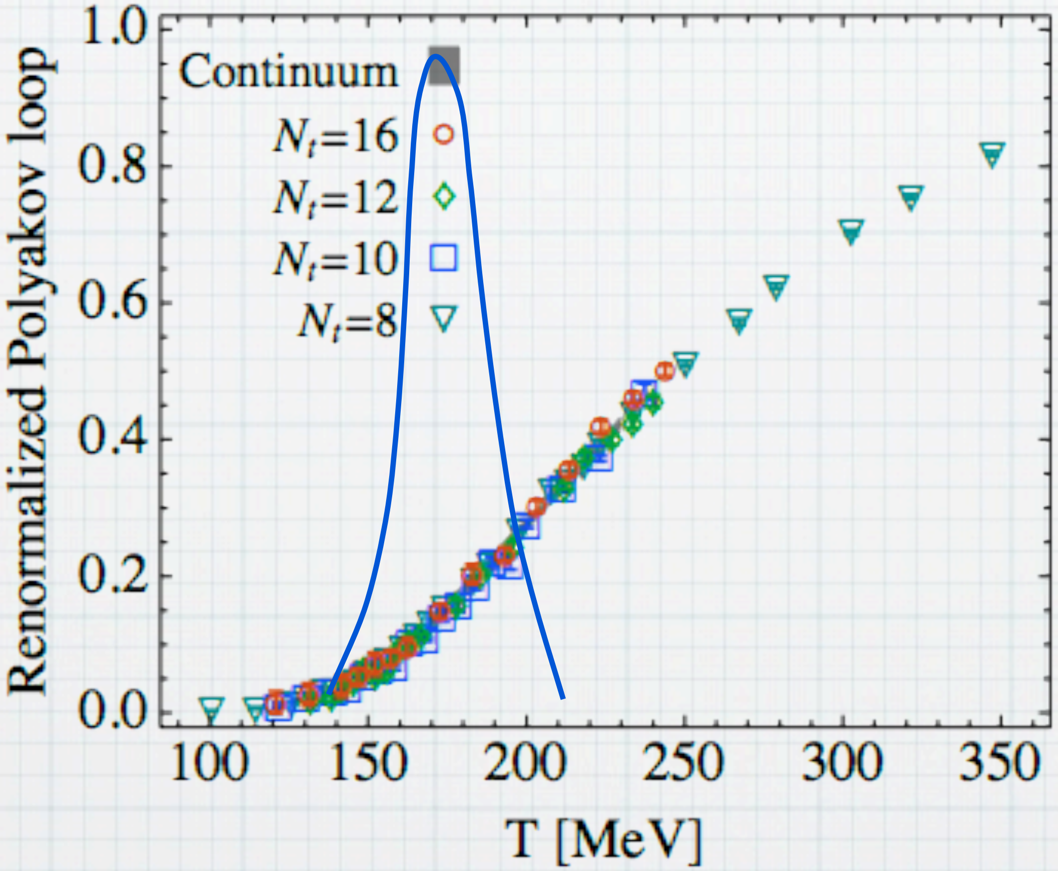

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Polyakov loop: order parameter for deconfinement

Chiral condensate: chiral transition

* The chiral condensate $\langle\bar{\psi}\psi\rangle$ is the vacuum expectation value of the operator $\bar{\psi}\psi$.

* The magnitude of the constituent quark mass is proportional to it

  ➡ Even if the “bare” quark mass in the QCD Lagrangian is small, they develop a constituent one, through interaction with the chiral condensate

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Chirally broken system
Large effective quark mass

$$\langle\bar{\psi}\psi\rangle \neq 0$$

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Chirally broken system Large effective quark mass $\langle\bar{\psi}\psi\rangle \neq 0$	Chirally restored system Small effective quark mass $\langle\bar{\psi}\psi\rangle = 0$

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Chirally broken system	Chirally restored system
Large effective quark mass	Small effective quark mass
$\langle\bar{\psi}\psi\rangle \neq 0$	$\langle\bar{\psi}\psi\rangle = 0$

Chiral condensate: order parameter for chiral phase transition

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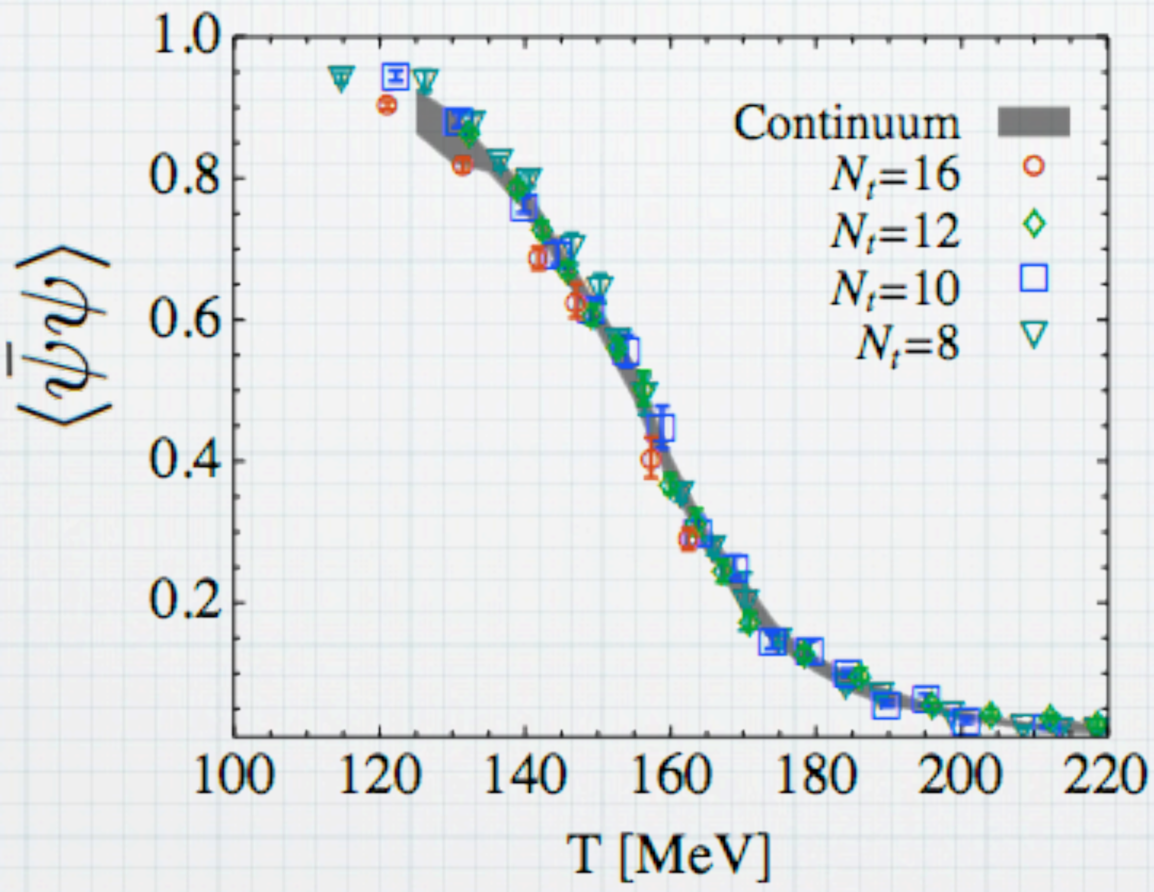

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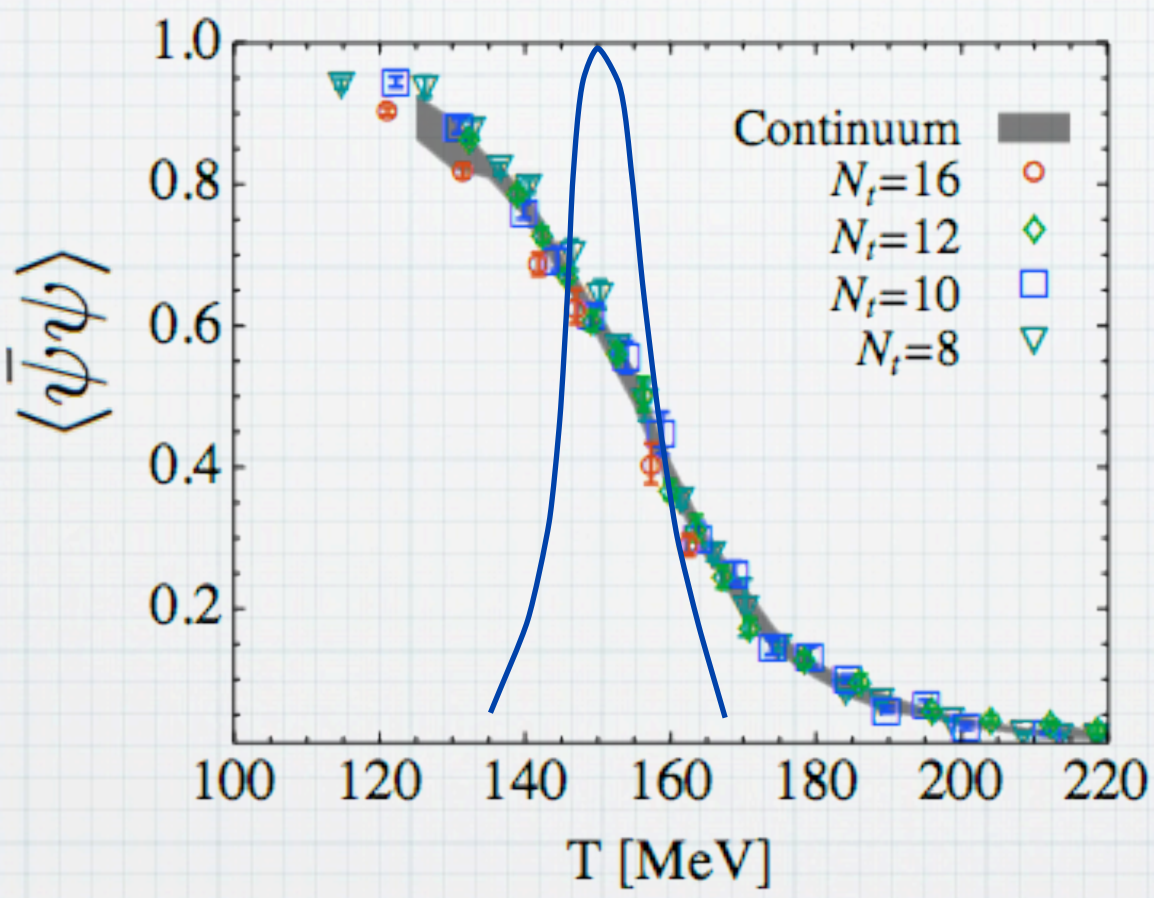

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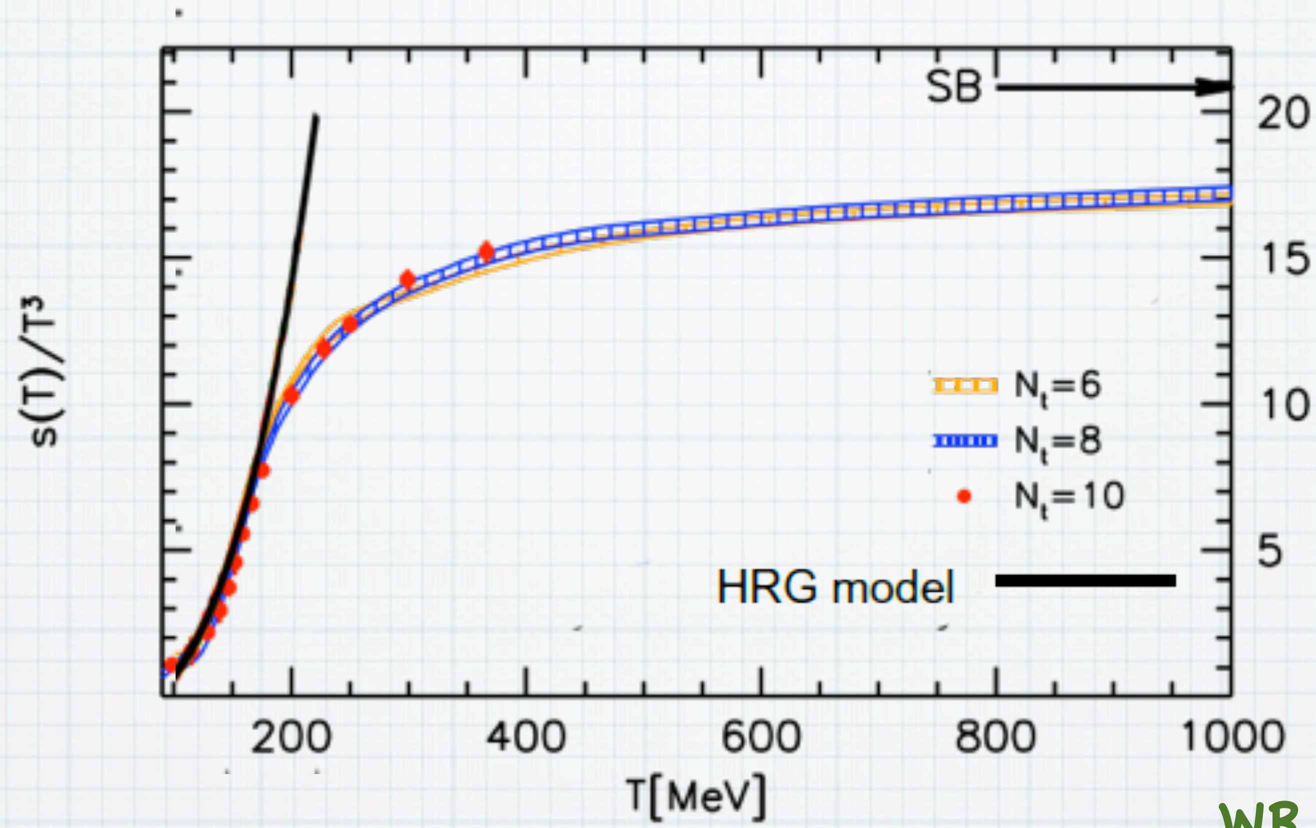
Transition from QCD Thermodynamics



WB Collab., JHEP (2010)

* s/T^3 indicates the number of particle species

* Rapid rise = liberation of degrees of freedom

* Compare to an ideal gas of quarks and gluons

$$s = \frac{4g}{\pi^2} T^3$$

* This gives us an idea of how strong the interaction is

What happens below T_c?

* At low T and μ_B=0, QCD thermodynamics is dominated by **pions**

* as T increases, heavier hadrons start to contribute

* Their mutual interaction is suppressed:

$$n_i n_k \sim \exp[-(M_i + M_k)/T]$$

* **Interacting** hadronic matter in the ground state can be well approximated by a **non-interacting** gas of **hadronic resonances**

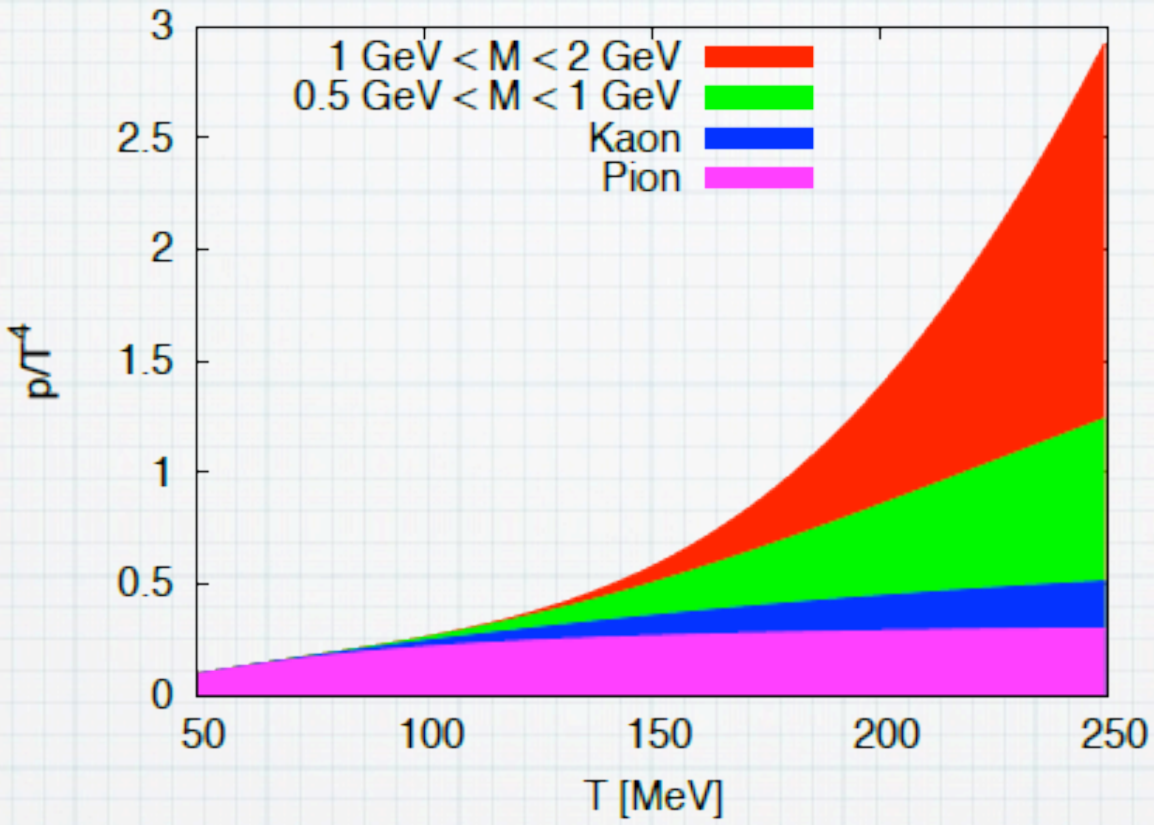
$$p^{HRG}/T^4 = \frac{1}{VT^3} \sum_{i \in \, mesons} \ln \mathcal{Z}^M_{m_i}(T, V, \mu_{X^a}) + \frac{1}{VT^3} \sum_{i \in \, baryons} \ln \mathcal{Z}^B_{m_i}(T, V, \mu_{X^a}) \, ,$$

with $\ln \mathcal{Z}^{M/B}_{m_i} = \mp \frac{V d_i}{2\pi^2} \int_0^\infty dk k^2 \ln(1 \mp z_i e^{-\varepsilon_i/T})$, $\varepsilon_i = \sqrt{k^2 + m_i^2}$,

$z_i = \exp\left((\sum_a X_i^a \mu_{X^a})/T \right)$ and X^a are all conserved charges.

R. Hagedorn, N. Cabibbo and G. Parisi

How many resonances do we include?

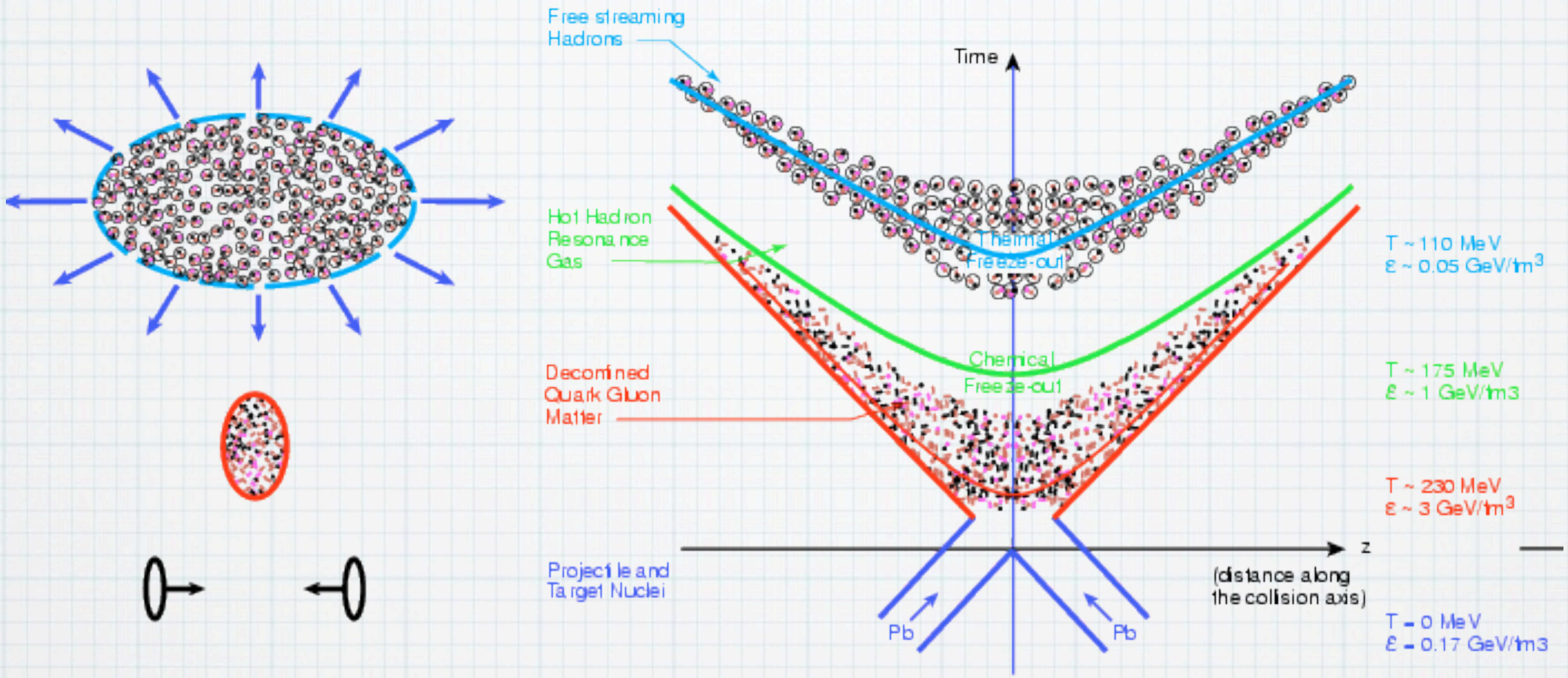


* With different mass cut-offs we can separate the contribution of different particles

* Known resonances up to M=2.5 GeV

* ~170 different masses ⟷ 1500 resonances

Evolution of a heavy-ion collision

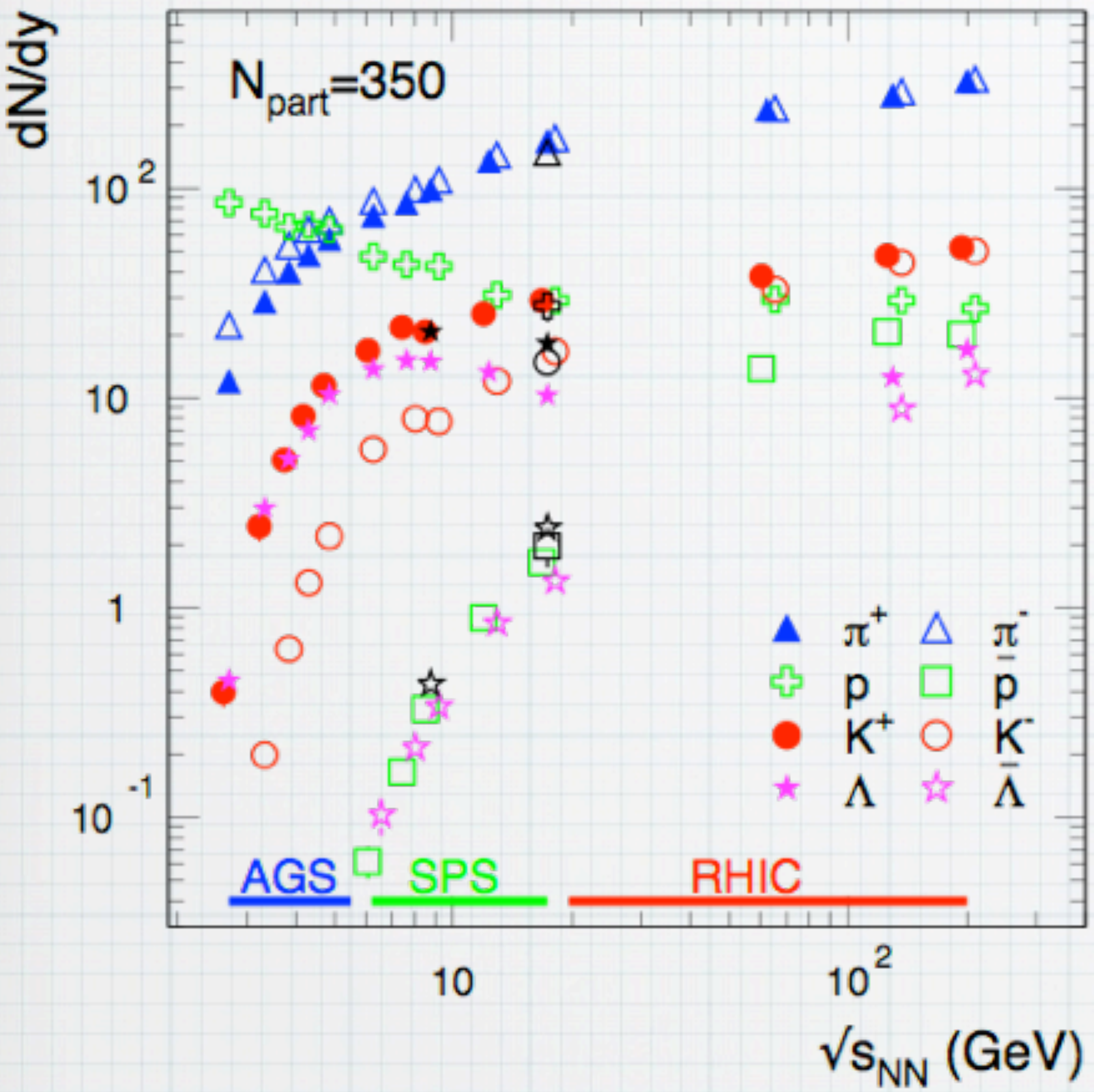


* **Chemical freeze-out**: inelastic reactions cease: the chemical composition of the system is fixed (particle yields and fluctuations)

* **Kinetic freeze-out**: elastic reactions cease: spectra and correlations are frozen (free streaming of hadrons)

* Hadrons reach the detector

Hadron yields



* **E=mc²**: lots of particles are created

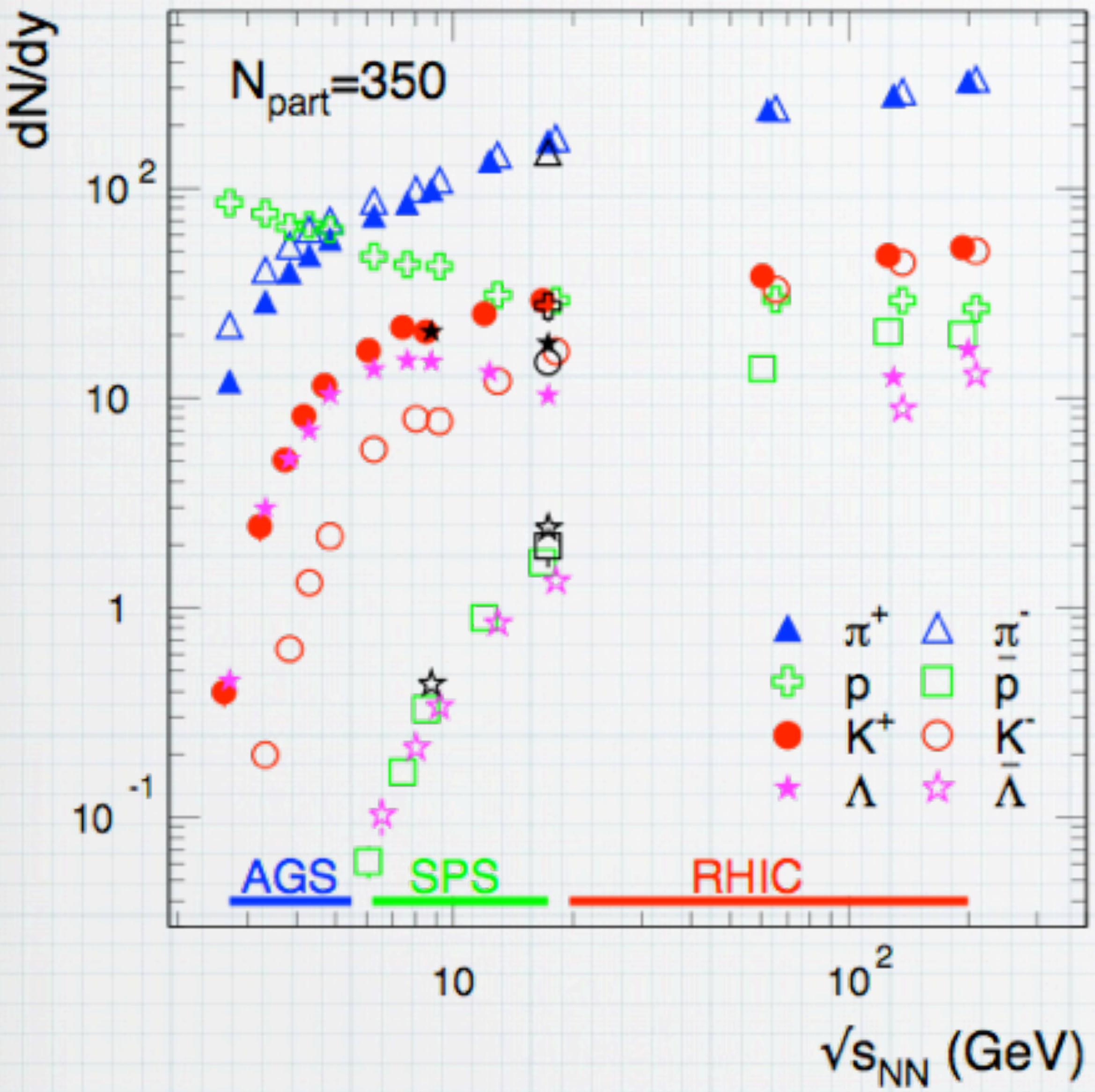
* **Particle counting** (average over many events)

* Take into account:
  * detector inefficiency
  * missing particles at low p_T
  * decays

* **HRG model**: test hypothesis of hadron abundancies in equilibrium

$$N_i = -T\frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i V}{2\pi^2}\int_0^\infty \frac{p^2 \mathrm{d}p}{\exp[(E_i-\mu_i)/T]\pm 1}$$

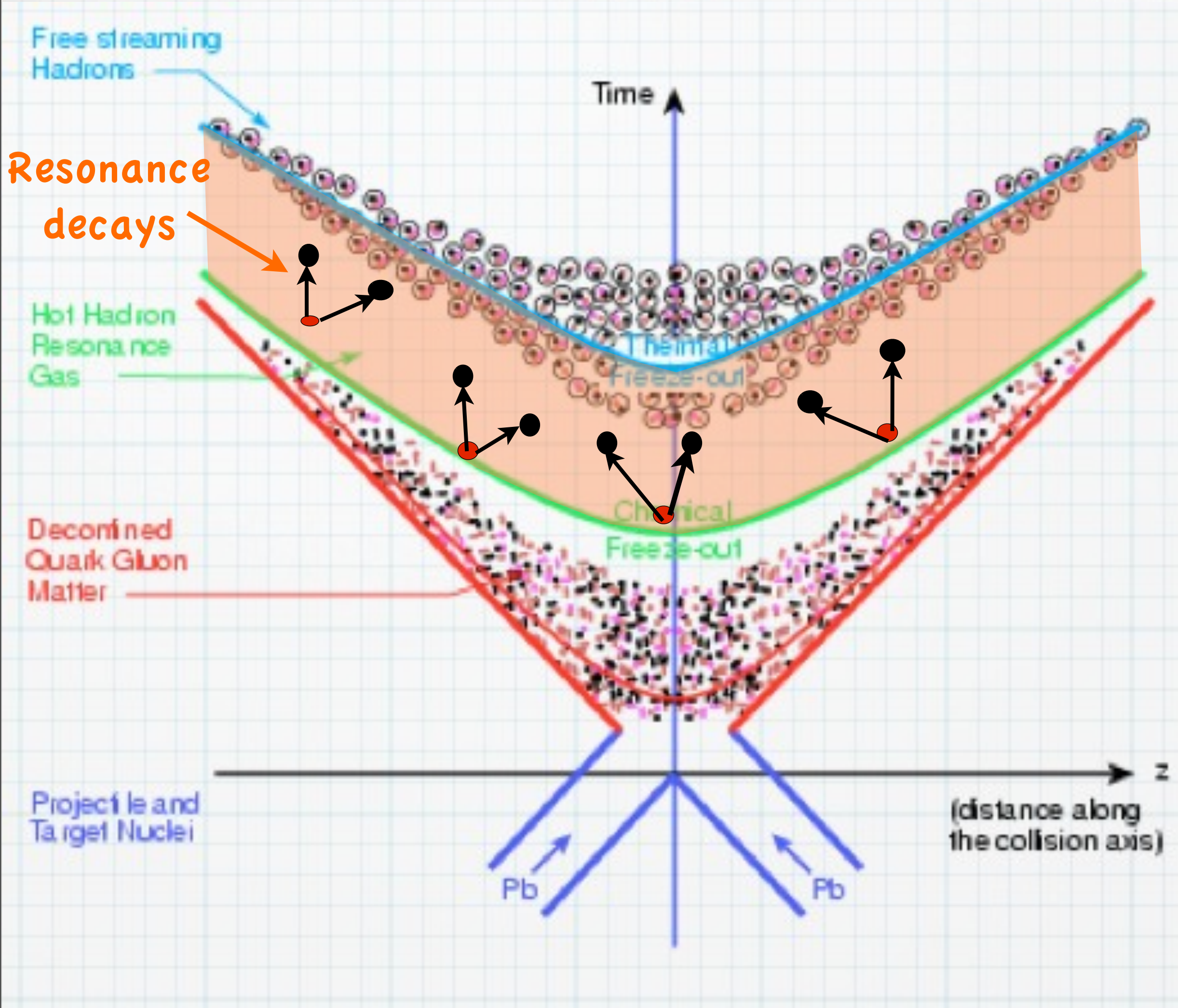
Hadron yields



- **E=mc²**: lots of particles are created
- **Particle counting** (average over many events)
- Take into account:
  - detector inefficiency
  - missing particles at low p_T
  - decays

- **HRG model**: test hypothesis of hadron abundancies in equilibrium
- We need:
  - a complete hadron spectrum
  - control the hadron fraction from decays

Decays



* Most hadrons are subject to **strong and electromagnetic decays**

$$\Delta \rightarrow p(n) + \pi,\ \rho \rightarrow \pi + \pi$$

* e.g. pions: **1/4** primordial, **3/4** from strong decays

* Weak decays can be treated too:

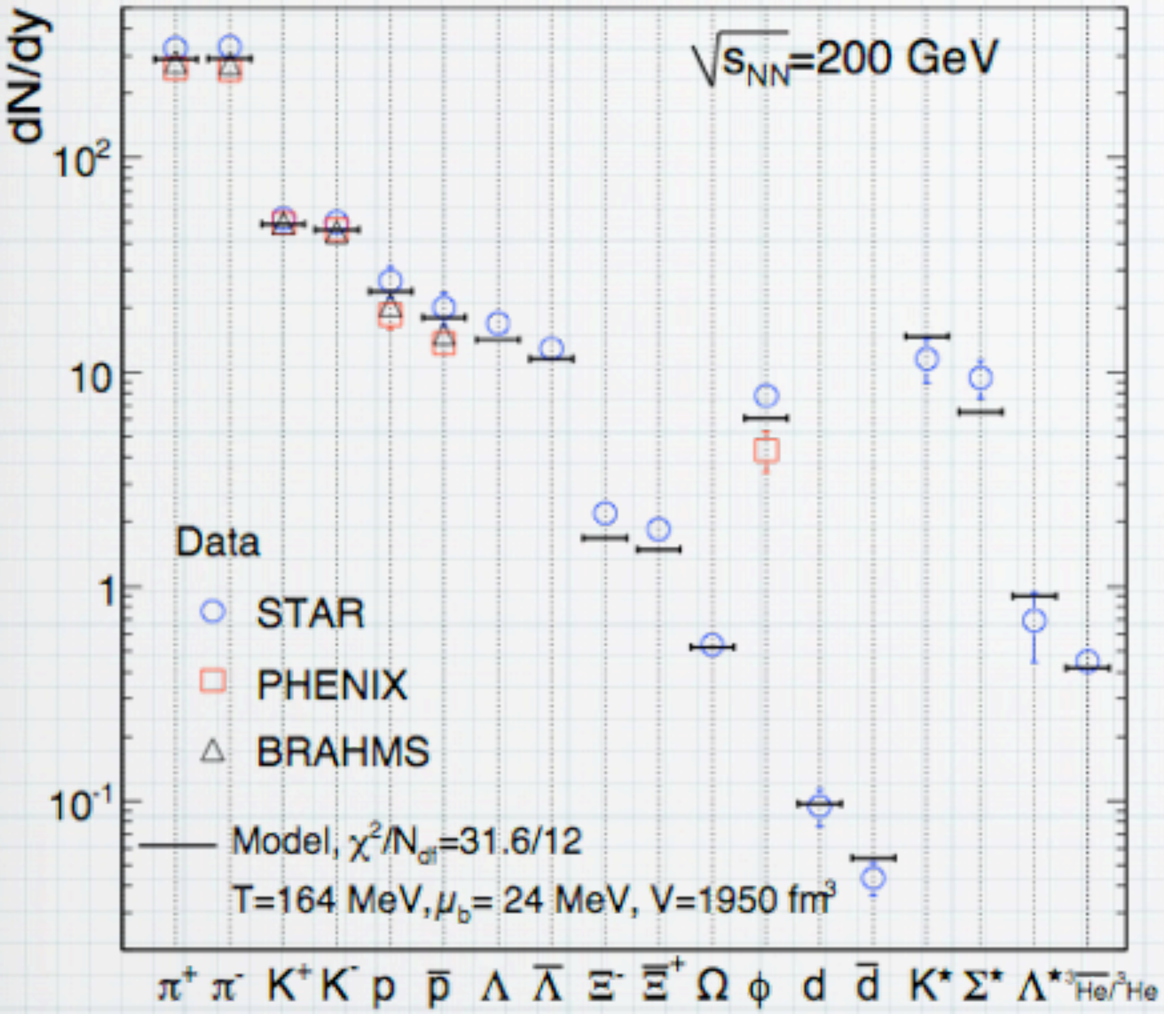
$$\Sigma \rightarrow \Lambda + \gamma$$

* **after chemical freeze-out:** only elastic and quasi-elastic scatterings take place:

$$\pi\pi \rightarrow \rho \rightarrow \pi\pi \qquad p\pi \rightarrow \Delta \rightarrow p\pi \qquad K\pi \rightarrow K^* \rightarrow K\pi$$

$$\bar{N_i} = N_i + \sum_r d_{r \rightarrow i} N_r$$

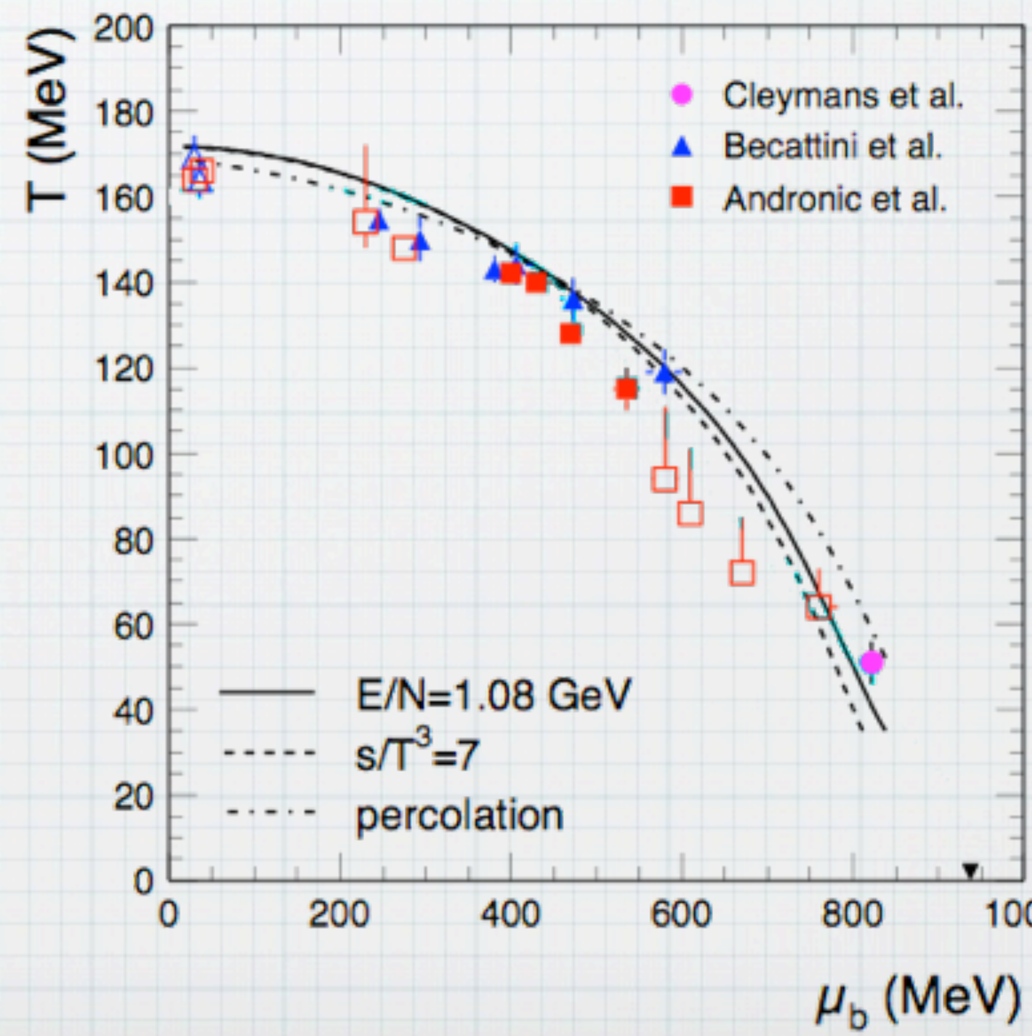
The thermal fits



* Fit is performed minimizing the χ^2

* **Fit to yields**: parameters T, μ_B, V

* **Fit to ratios**: the volume V cancels out



* Changing the collision energy, it is possible to draw the freeze-out line in the T, μ_B plane

Cleymans et al, Becattini et al, Andronic et al.

Caveats

* These results are model-dependent
  * they depend on the particle spectrum which is included in the model
  * possibility of having heavier states with exponential mass spectrum
  * not known experimentally but can be postulated
  * their decay modes are not known

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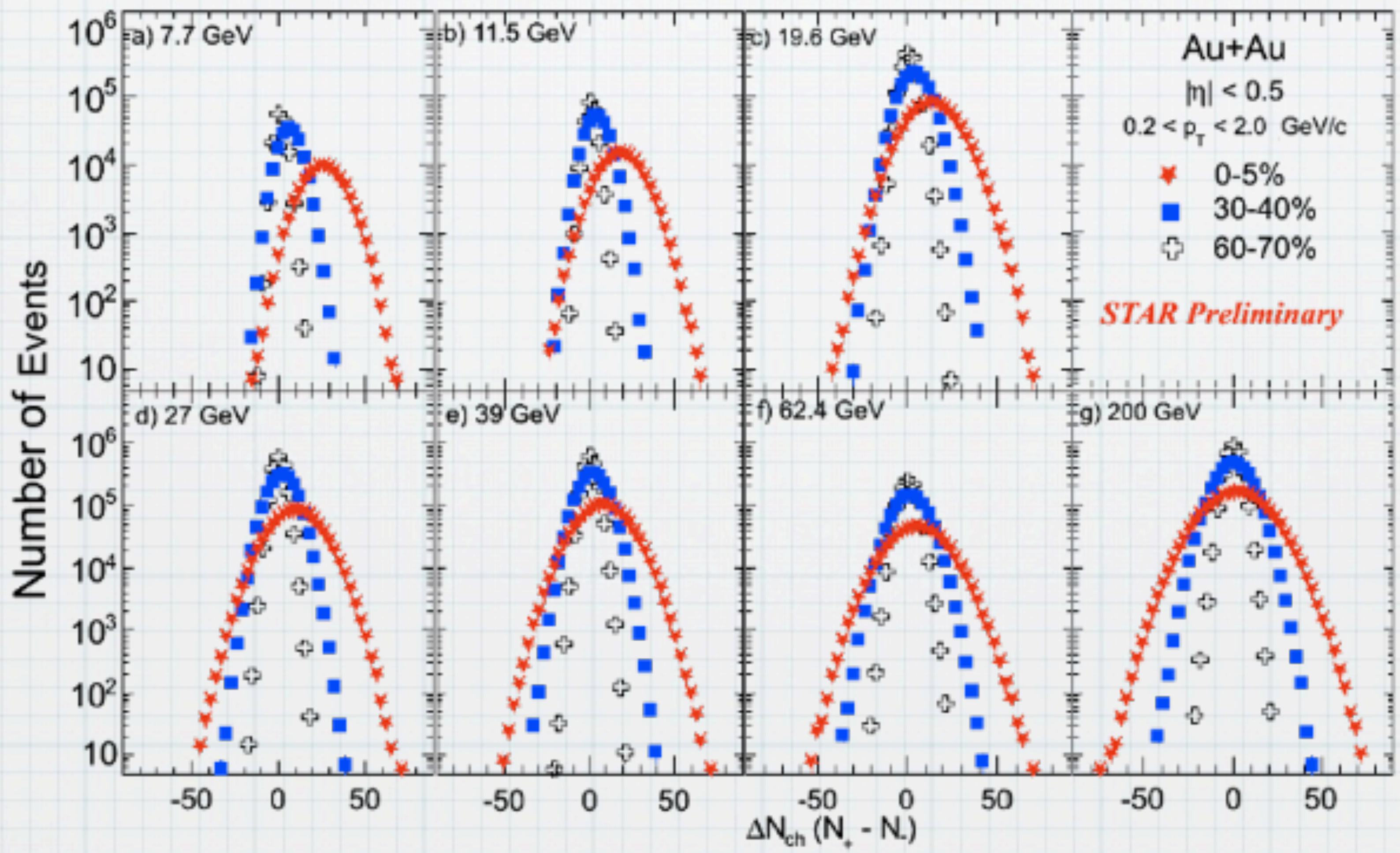
* Purpose: extract freeze-out parameters from first principles
  * direct comparison between experimental measurement and lattice QCD results
  * observable: fluctuations of conserved charges (electric charge, baryon number and strangeness)
  * directly related to moments of multiplicity distribution (measured)
  * lattice QCD looks at conserved charges rather than identified particles

Fluctuations of conserved charges

* Consider the number of electrically charged particles N_Q
* Its average value over the whole ensemble of events is $\langle N_Q \rangle$
* In experiments it is possible to **measure its event-by-event distribution**

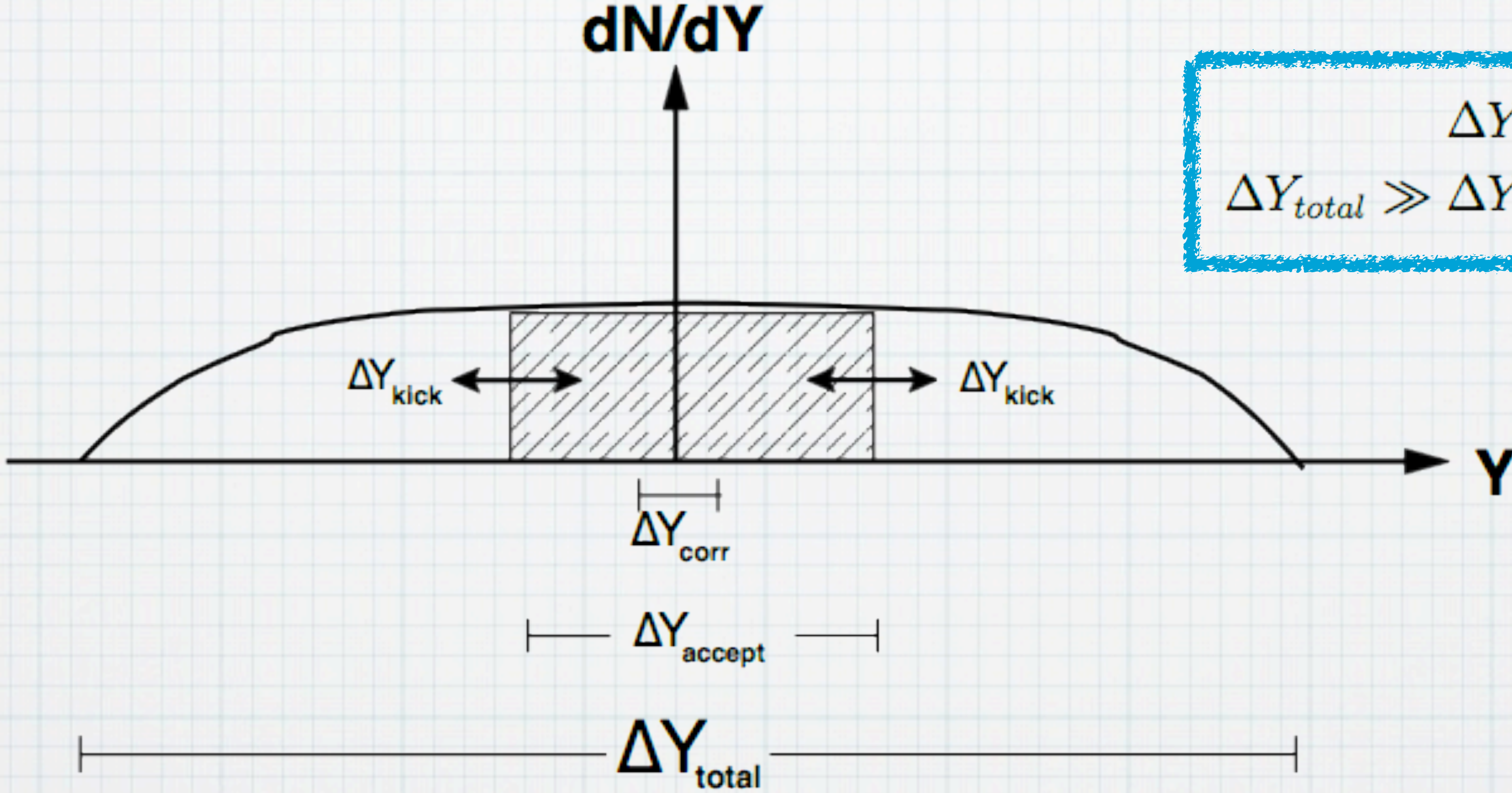


STAR Collab.: 1402.1558

Fluctuations of conserved charges???

* If we look at the entire system, none of the conserved charges will fluctuate

* By studying a sufficiently small subsystem, the fluctuations of conserved quantities become meaningful



$$\Delta Y_{accept} \gg \Delta Y_{corr}$$
$$\Delta Y_{total} \gg \Delta Y_{accept} \gg \Delta Y_{kick}$$

- ΔY_{total}: range for total charge multiplicity distribution
- ΔY_{accept}: interval for the accepted charged particles
- ΔY_{corr}: charge correlation length characteristic to the physics of interest
- ΔY_{kick}: rapidity shift that charges receive during and after hadronization

V. Koch: 0810.2520

Cumulants of multiplicity distribution

* Deviation of N_Q from its mean in a single event: $\delta N_Q = N_Q - \langle N_Q \rangle$

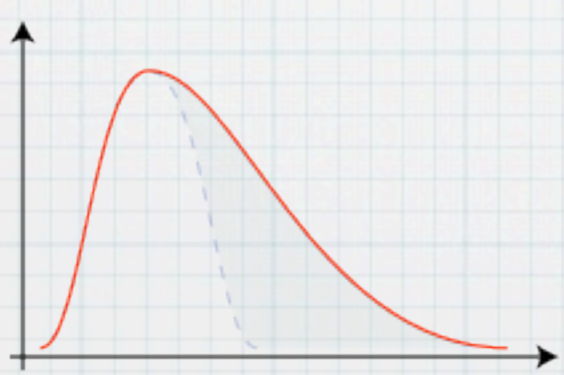
* The cumulants of the event-by-event distribution of N_Q are:

$$K_2 = \langle (\delta N_Q)^2 \rangle \qquad K_3 = \langle (\delta N_Q)^3 \rangle \qquad K_4 = \langle (\delta N_Q)^4 \rangle - 3\langle (\delta N_Q)^2 \rangle^2$$

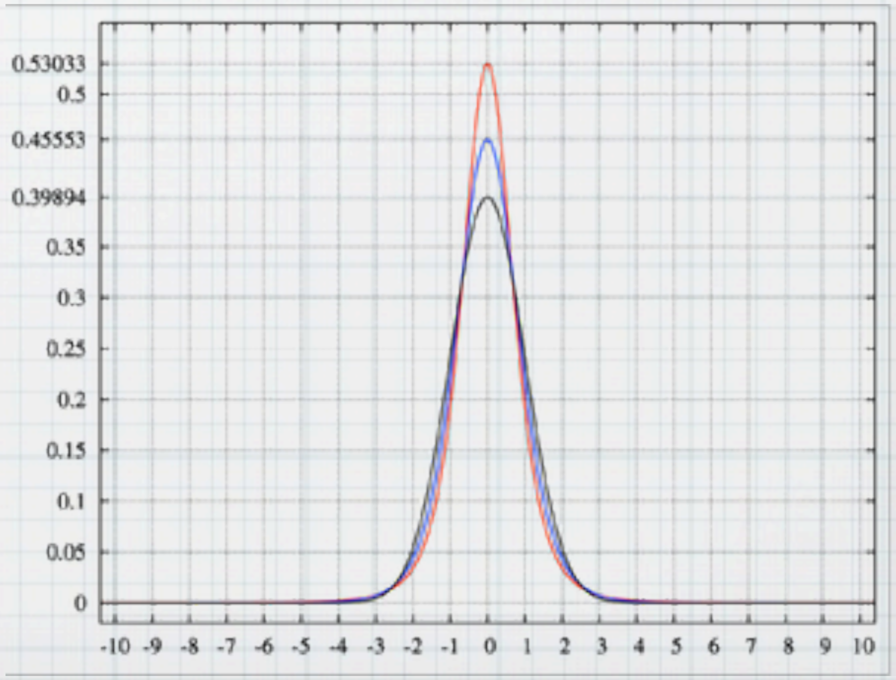
* The cumulants are related to the central moments of the distribution by:

variance: $\sigma^2 = K_2$ Skewness: $S = K_3/(K_2)^{3/2}$ Kurtosis: $\kappa = K_4/(K_2)^2$

Negative Skew



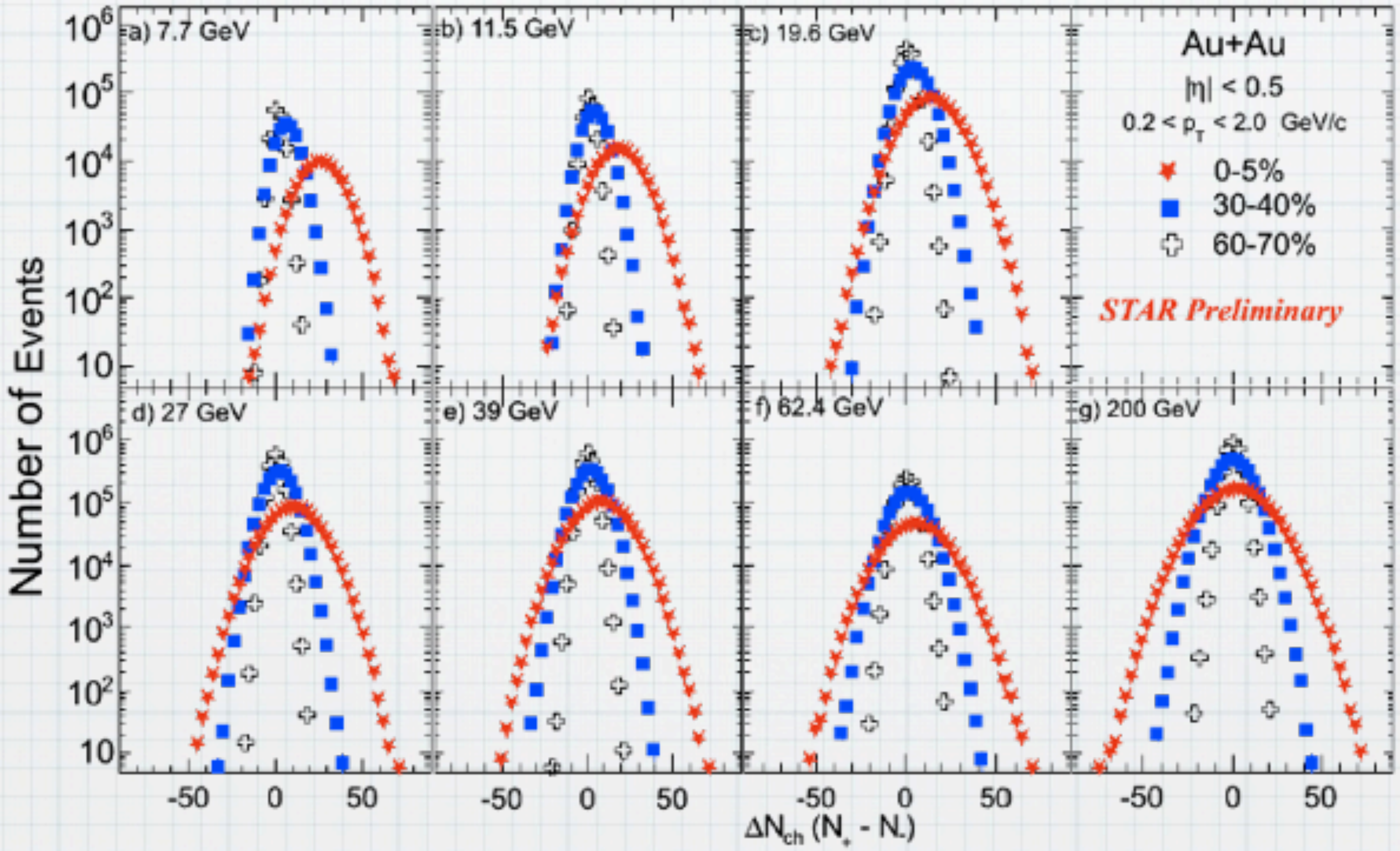
Positive Skew



Experimental measurement

* Volume-independent ratios:

$M/\sigma^2 = K_1/K_2$ $\quad S\sigma = K_3/K_2$ $\quad \kappa\sigma^2 = K_4/K_2$ $\quad S\sigma^3/M = K_3/K_1$



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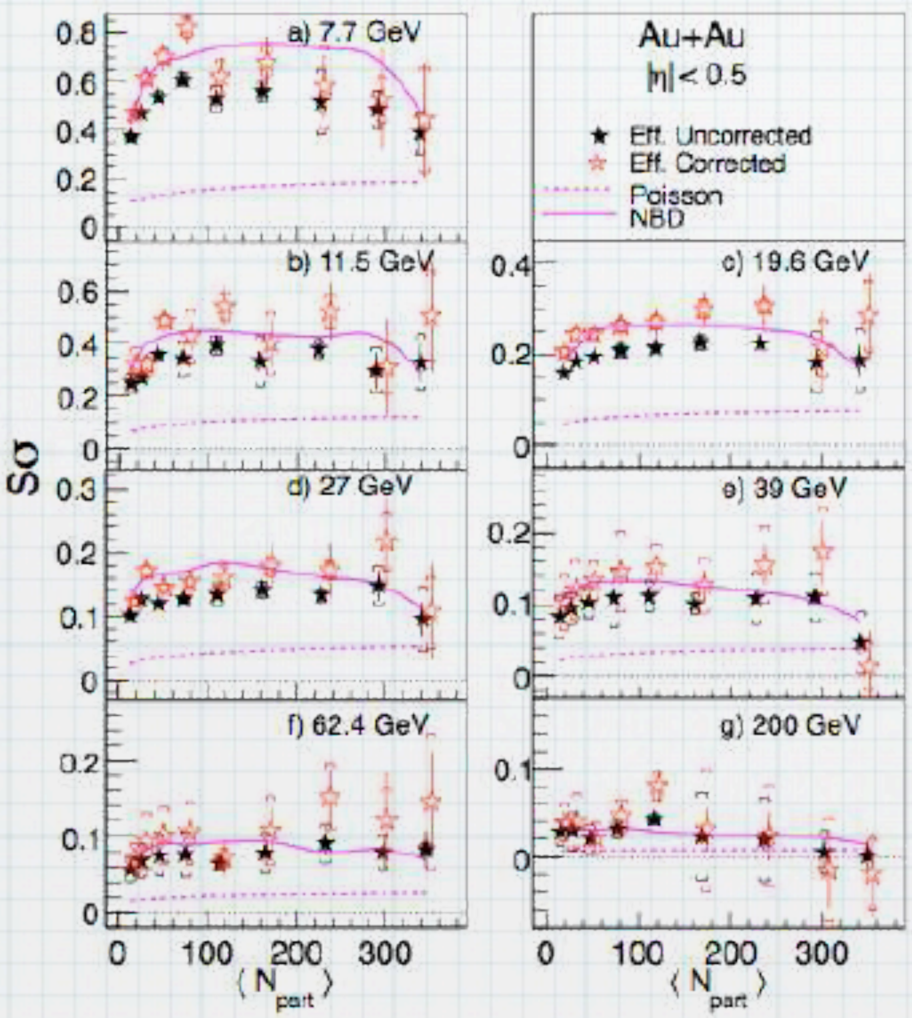
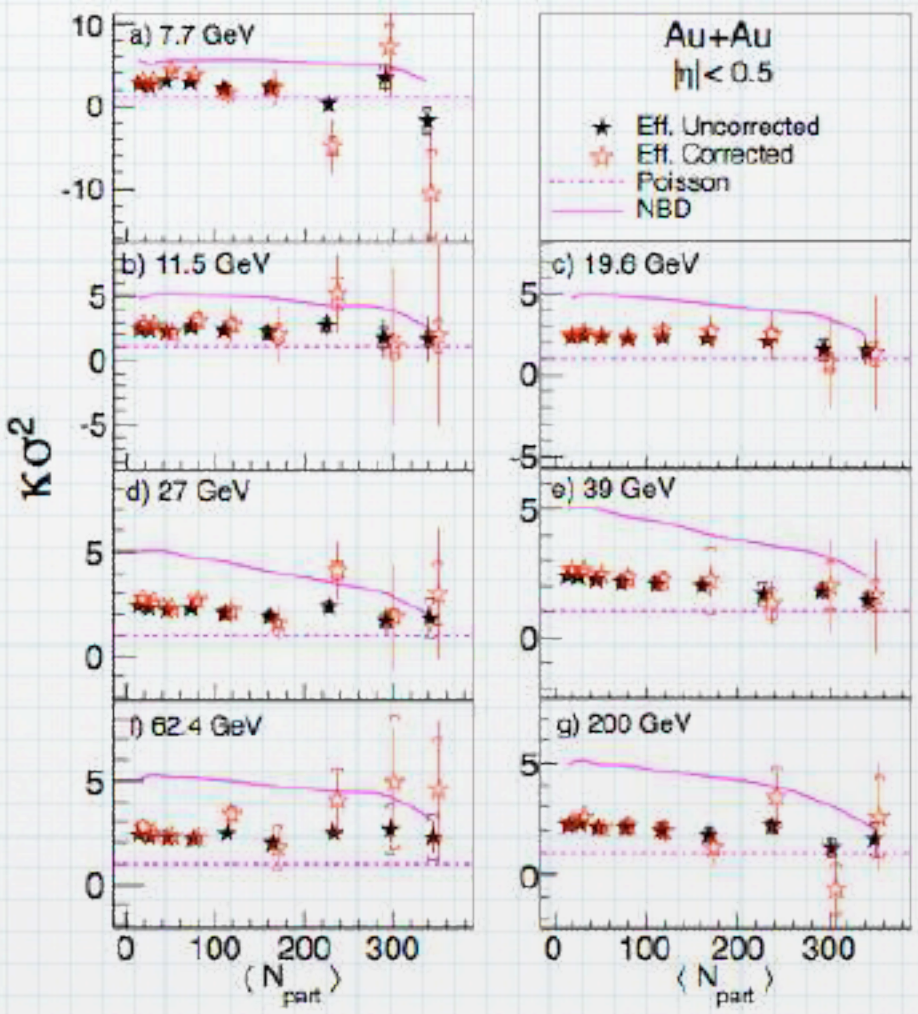




STAR Collab.: 1402.1558

Susceptibilities of conserved charges

* Susceptibilities of conserved charges

$$\chi_{lmn}^{BSQ} = \frac{\partial^{l+m+n} p/T^4}{\partial(\mu_B/T)^l \partial(\mu_S/T)^m \partial(\mu_Q/T)^n}.$$

* **Susceptibilities** of conserved charges are the **cumulants** of their event-by event distribution

$$\text{mean}: \ M = \chi_1 \qquad \text{variance}: \ \sigma^2 = \chi_2$$

$$\text{skewness}: \ S = \chi_3/\chi_2^{3/2} \qquad \text{kurtosis}: \ \kappa = \chi_4/\chi_2^2$$

$$S\sigma = \chi_3/\chi_2 \qquad \kappa\sigma^2 = \chi_4/\chi_2$$

$$M/\sigma^2 = \chi_1/\chi_2 \qquad S\sigma^3/M = \chi_3/\chi_1$$

* Lattice QCD results are functions of **temperature** and **chemical potential**

➡ By comparing lattice results and experimental measurement we can **extract the freeze-out parameters from first principles**

F. Karsch: Centr. Eur. J. Phys. (2012)

Baryometer and thermometer

* Let us look at the Taylor expansion of R^B_{31}

$$R^B_{31}(T,\mu_B) = \frac{\chi_3^B(T,\mu_B)}{\chi_1^B(T,\mu_B)} = \frac{\chi_4^B(T,0) + \chi_{31}^{BQ}(T,0)q_1(T) + \chi_{31}^{BS}(T,0)s_1(T)}{\chi_2^B(T,0) + \chi_{11}^{BQ}(T,0)q_1(T) + \chi_{11}^{BS}(T,0)s_1(T)} + \mathcal{O}(\mu_B^2)$$

* To order μ^2_B it is independent of μ_B: it can be used as a thermometer

* Let us look at the Taylor expansion of R^B_{12}

$$R^B_{12}(T,\mu_B) = \frac{\chi_1^B(T,\mu_B)}{\chi_2^B(T,\mu_B)} = \frac{\chi_2^B(T,0) + \chi_{11}^{BQ}(T,0)q_1(T) + \chi_{11}^{BS}(T,0)s_1(T)}{\chi_2^B(T,0)} \frac{\mu_B}{T} + \mathcal{O}(\mu_B^3)$$

* Once we extract T from R^B_{31}, we can use R^B_{12} to extract μ_B

Caveats

* Effects due to volume variation because of finite centrality bin width
* Finite reconstruction efficiency
* Spallation protons
* Canonical vs Gran Canonical ensemble
* Proton multiplicity distributions vs baryon number fluctuations
* Final-state interactions in the hadronic phase

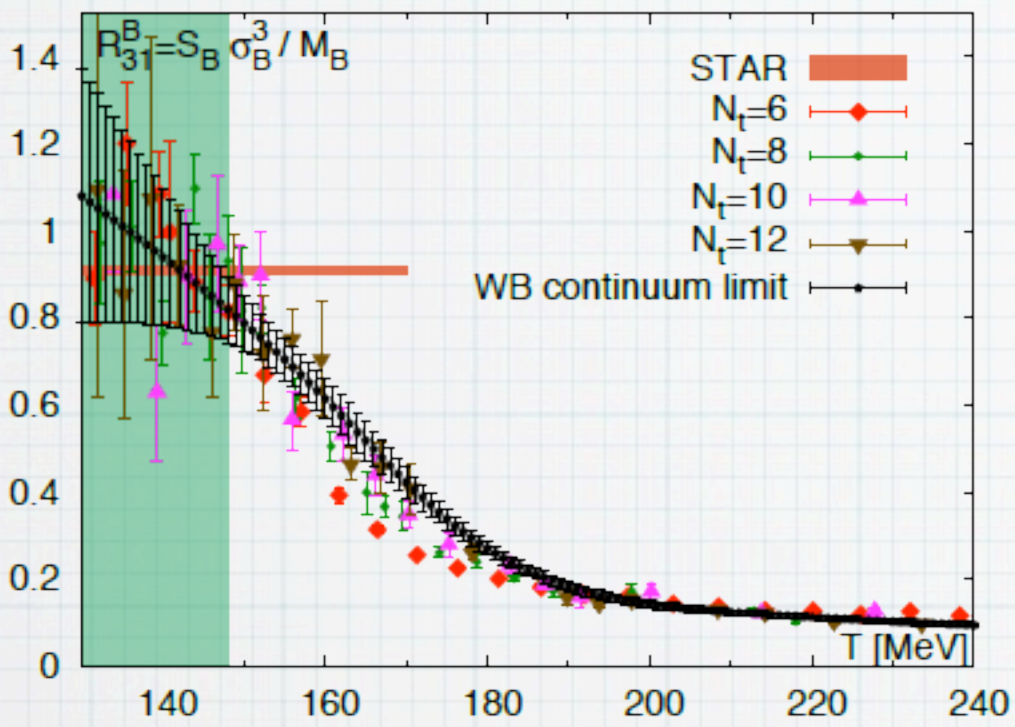
Caveats

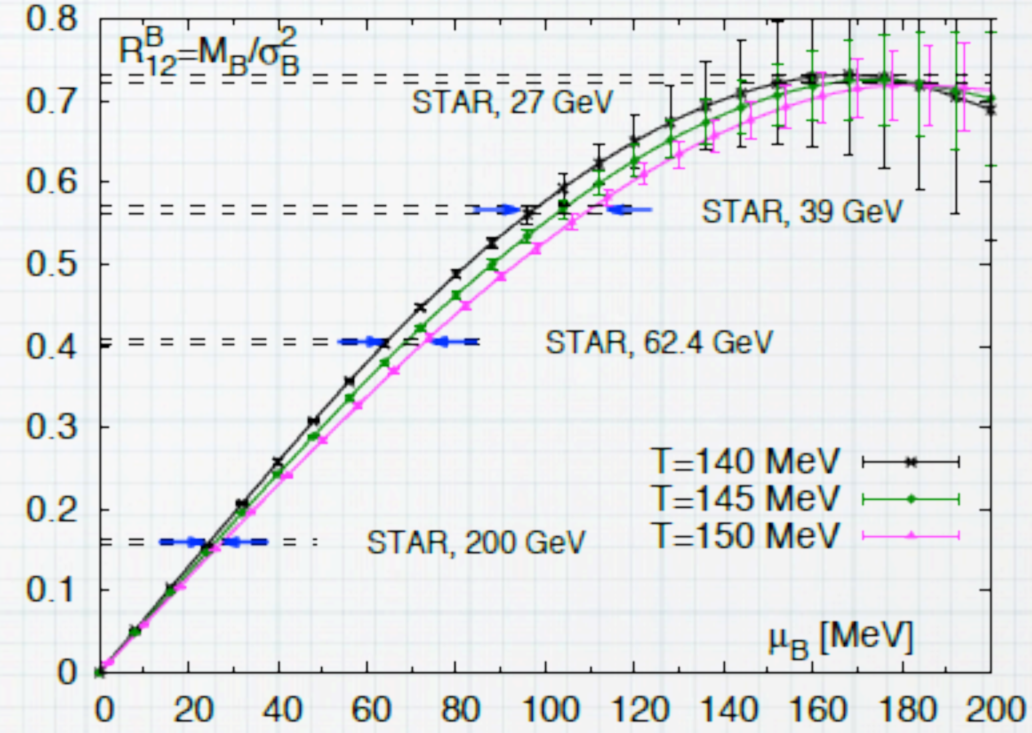
* Effects due to volume variation because of finite centrality bin width
  - ➡ Experimentally corrected by centrality-bin-width correction method
* Finite reconstruction efficiency
  - ➡ Experimentally corrected based on binomial distribution A.Bzdak,V.Koch, PRC (2012)
* Spallation protons
  - ➡ Experimentally removed with proper cuts in p_T
* Canonical vs Gran Canonical ensemble
  - ➡ Experimental cuts in the kinematics and acceptance V. Koch, S. Jeon, PRL (2000)
* Proton multiplicity distributions vs baryon number fluctuations
  - ➡ Numerically very similar once protons are properly treated

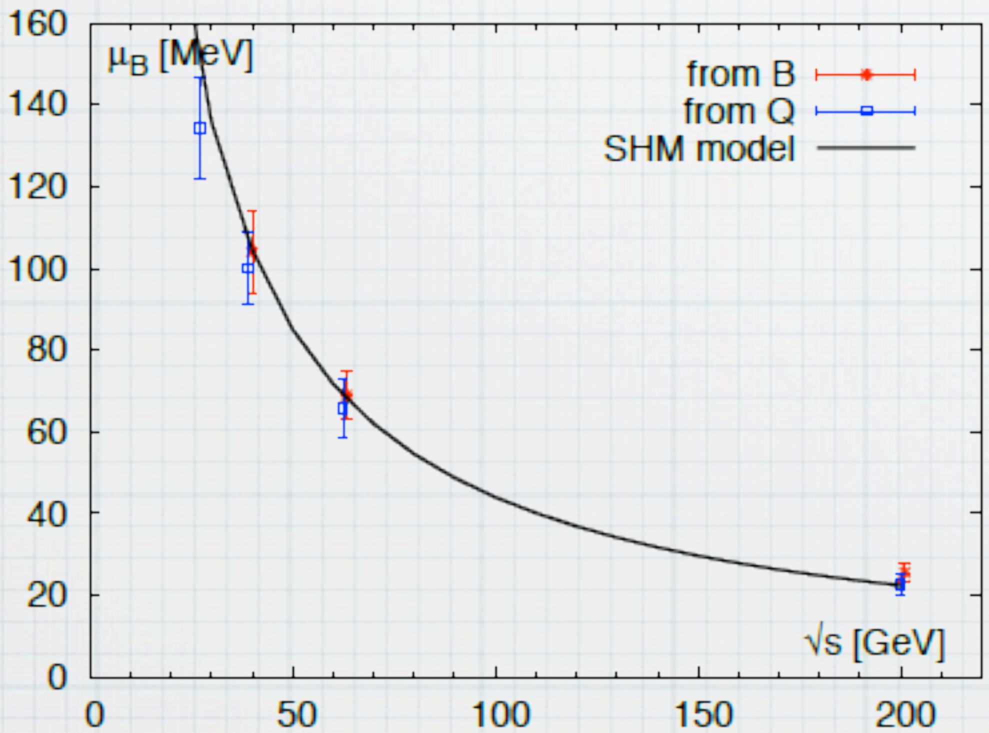
    M. Asakawa and M. Kitazawa, PRC(2012), M. Nahrgang et al., 1402.1238
* Final-state interactions in the hadronic phase

  J.Steinheimer et al., PRL (2013)
  - ➡ Consistency between different charges = fundamental test

Results







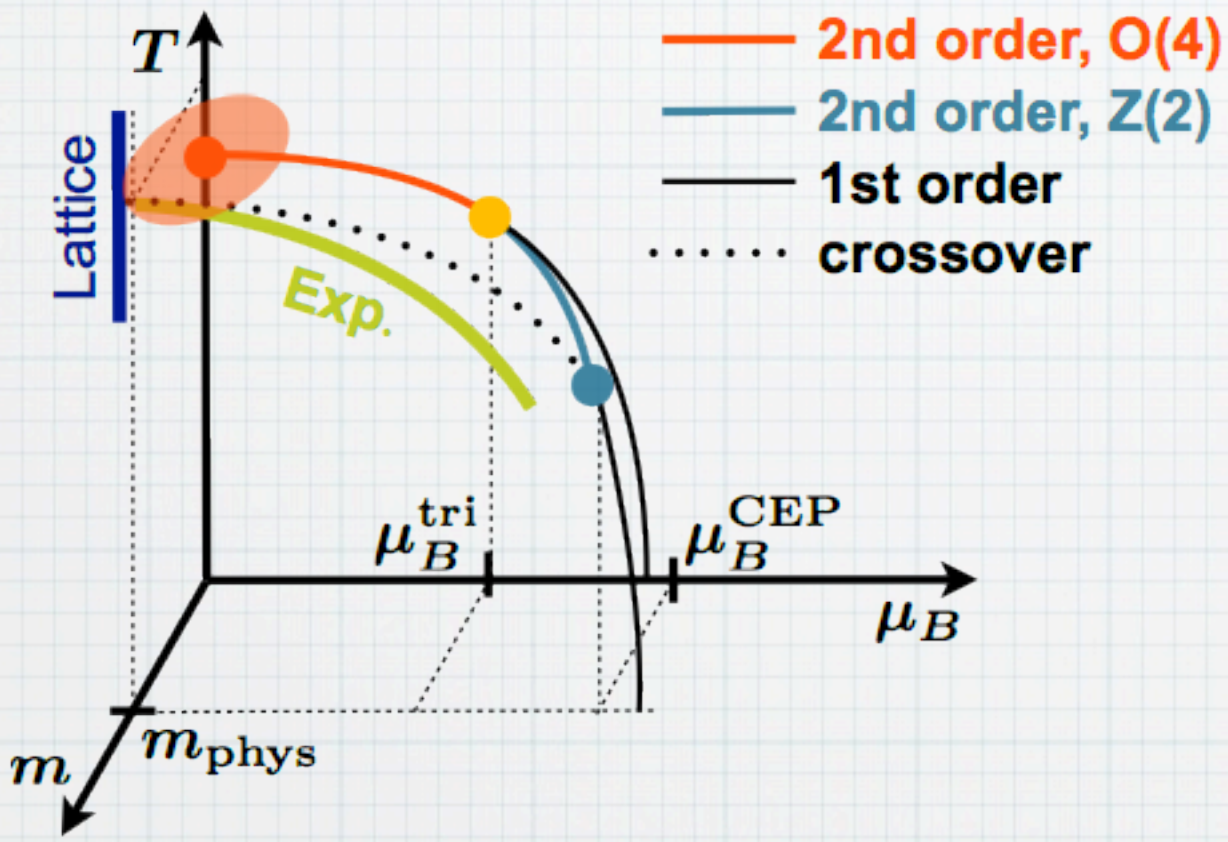
Consistency between different charges and with SHM fits

See QCD Phase diagram session on Wednesday morning

WB Collab.: 1403.4576

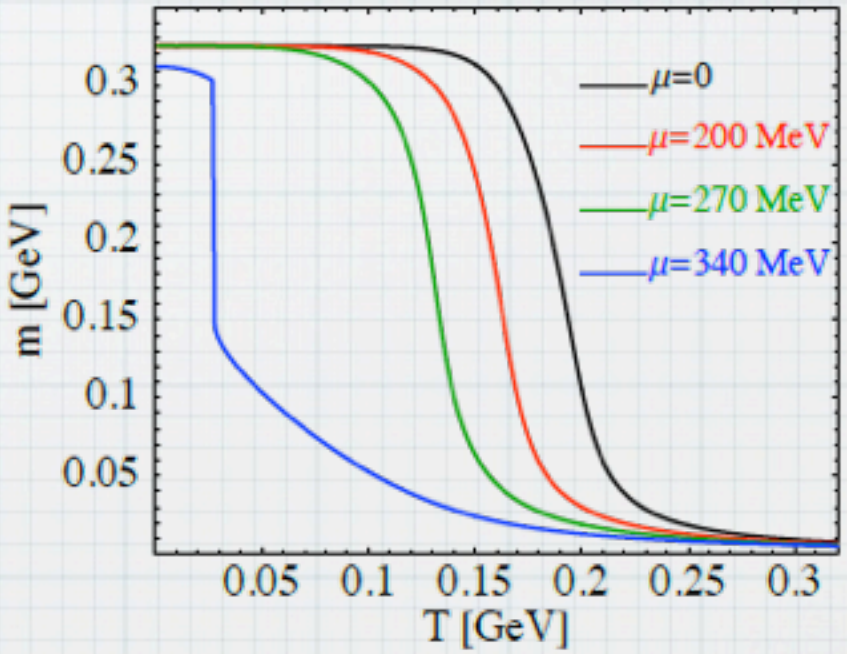
Our world is not ideal:

neither chiral symmetry (m_q=0) nor confinement (m_q=∞) is well defined.



Existence of QCD critical point predicted by models

Z(2) universality class



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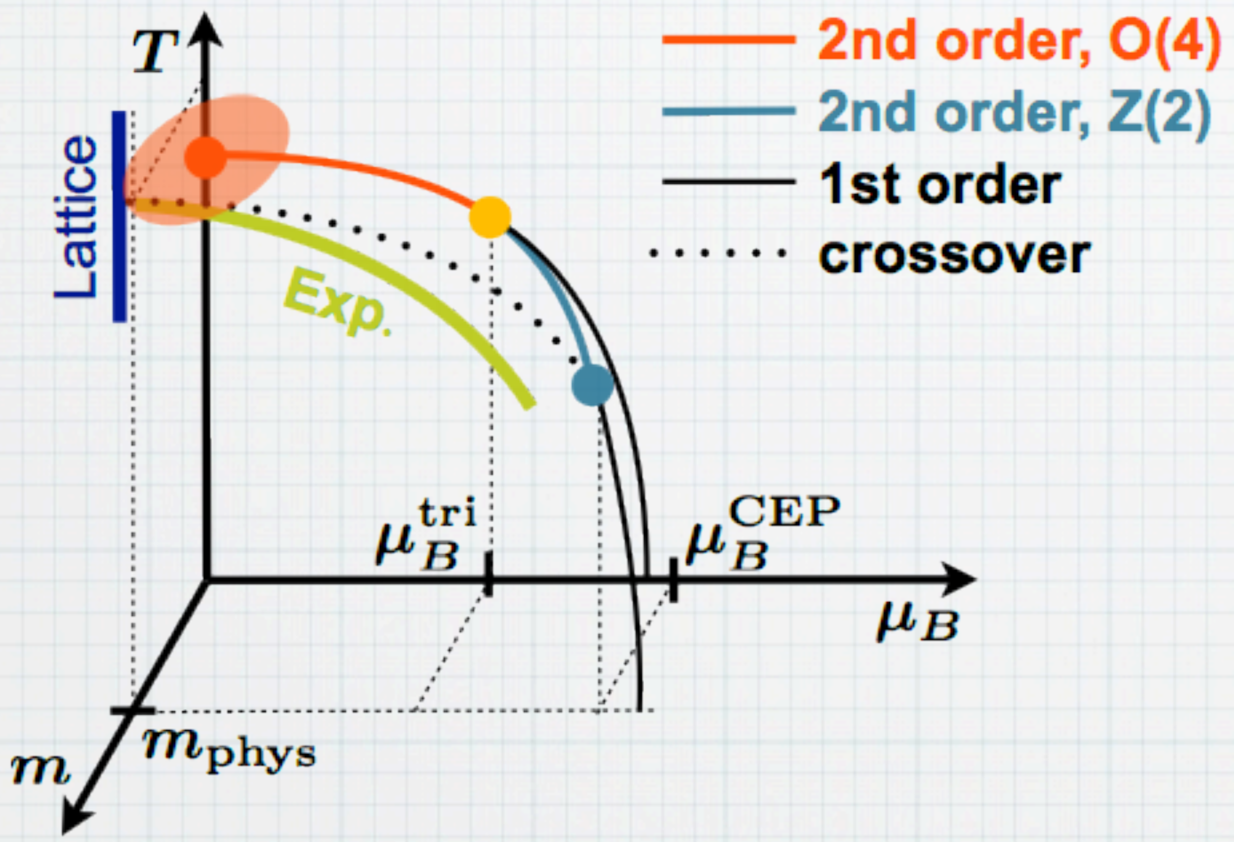

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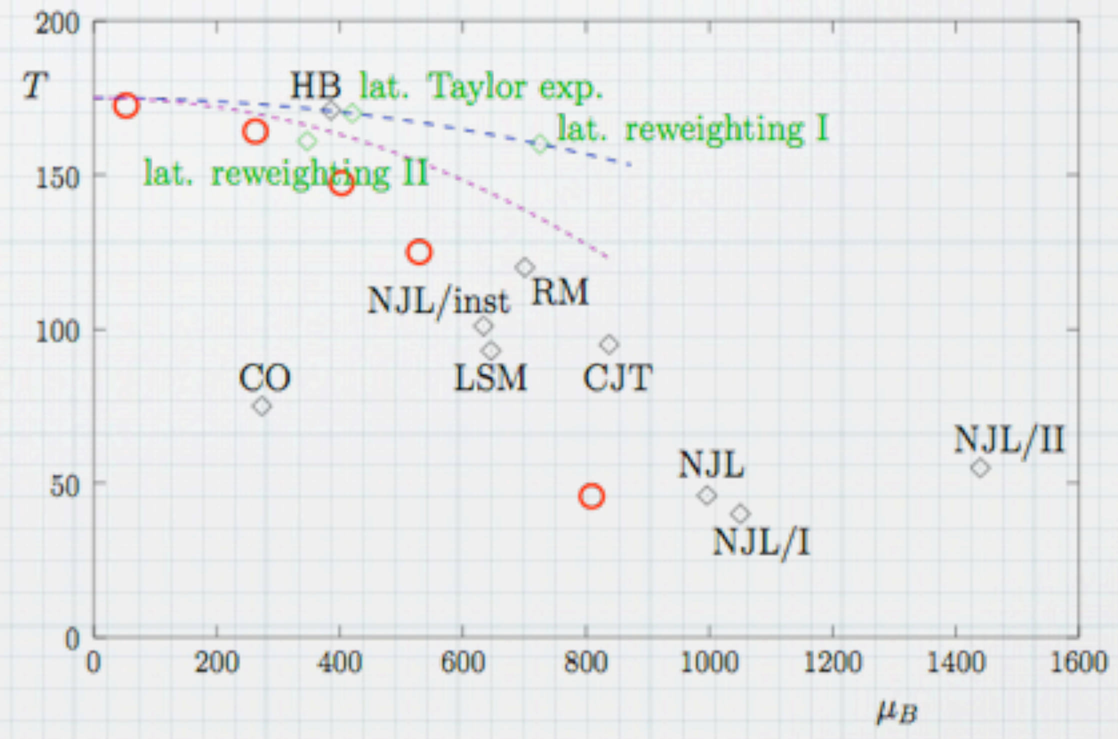
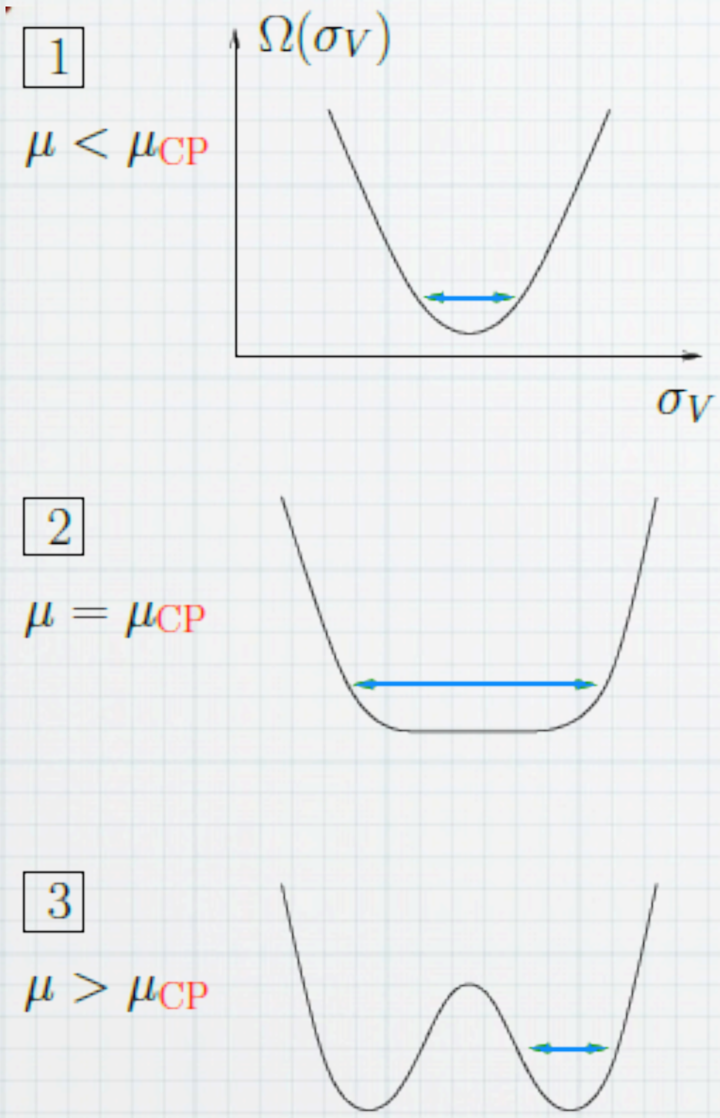


Fluctuations at the critical point

Rajagopal, Shuryak, Stephanov (1998)



Consider the order parameter for the chiral phase transition σ ~< $\bar{\psi}\psi$ >

It has a probability distribution of the form:

$$P[\sigma] \sim \exp\{-\Omega[\sigma]/T\},$$

$$\Omega = \int d^3x \left[\frac{1}{2}(\nabla\sigma)^2 + \frac{m_\sigma^2}{2}\sigma^2 + \frac{\lambda_3}{3}\sigma^3 + \frac{\lambda_4}{4}\sigma^4 + \ldots\right].$$

where: $m_\sigma \equiv \xi^{-1}$

and, near the critical point, ξ->∞:

$$\lambda_3 = \tilde{\lambda}_3\, T\, (T\,\xi)^{-3/2}, \quad \text{and} \;\; \lambda_4 = \tilde{\lambda}_4\, (T\,\xi)^{-1}$$

$\kappa_3 = \langle \sigma_V^3 \rangle = 2VT^{3/2}\,\tilde{\lambda}_3\, \xi^{4.5}$ $\kappa_4 = \langle \sigma_V^4 \rangle = 6VT^2\,[2(\tilde{\lambda}_3)^2 - \tilde{\lambda}_4]\,\xi^7$

M. Stephanov (2009)

correlation length ξ is limited due to critical slowing down, together with the finite time the system has to develop the correlations: ξ<2-3 fm

Berdnikov-Rajagopal

Experimental fluctuations

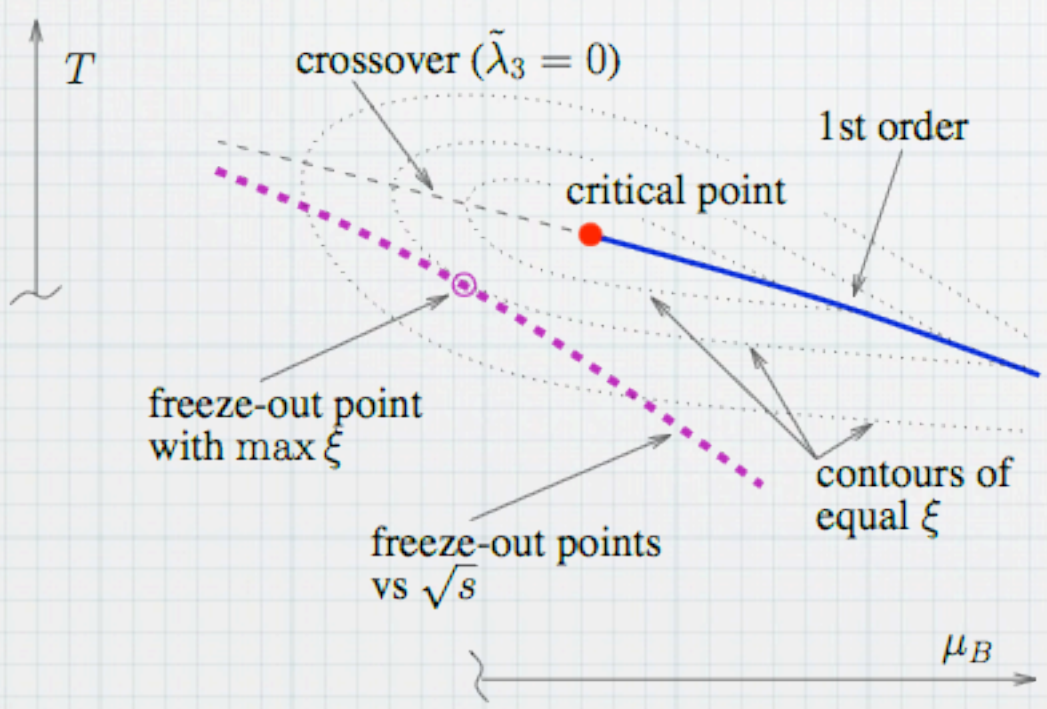
We consider the fluctuation of an observable (e.g. proton multiplicity)

$$\delta N = \sum_{\boldsymbol{p}} \delta n_{\boldsymbol{p}}$$

At the critical point, it receives both a regular and a singular contribution. The latter comes from the coupling to the σ field:

$$\delta n_{\boldsymbol{p}} = \underbrace{\delta n^0_{\boldsymbol{p}}}_{\text{statistical}} + \underbrace{\frac{\partial \bar{n}_{\boldsymbol{p}}}{\partial m} g\, \delta\sigma}_{\text{critical}}$$

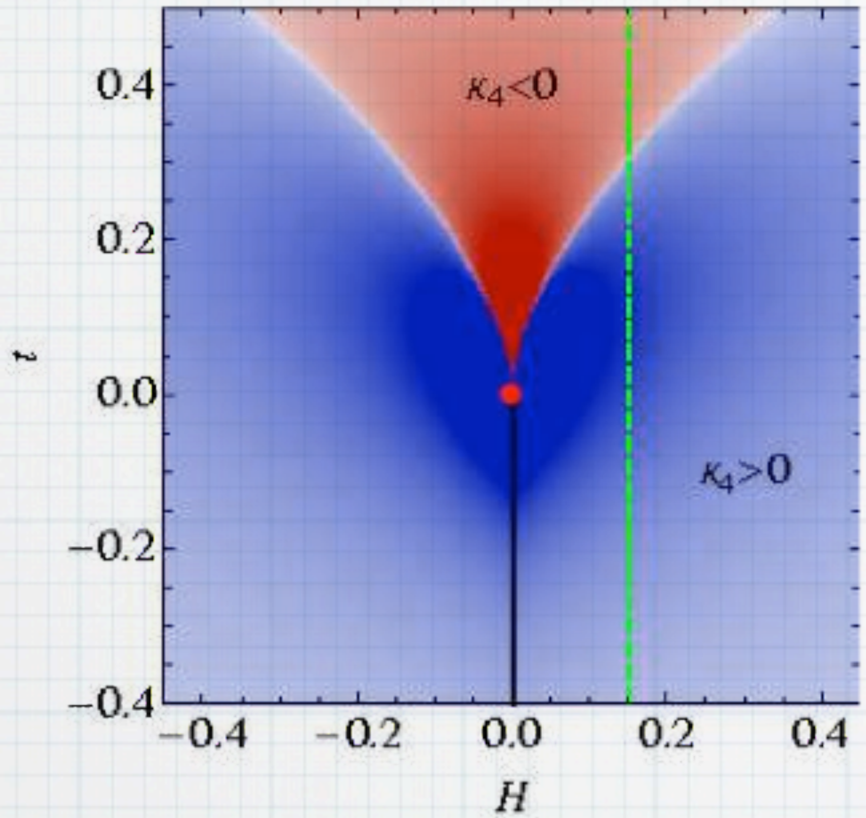
(Poisson)



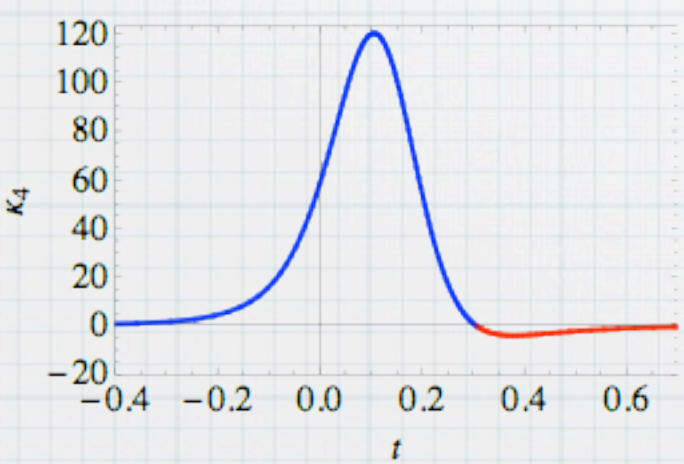
Higher order moments have **stronger dependence on ξ**: they are more sensitive signatures for the critical point

M. Stephanov

Sign of kurtosis



(a)



M. Stephanov (2011)

The 4th order cumulant becomes **negative** when the critical point is **approached from the crossover side**: from Ising model:

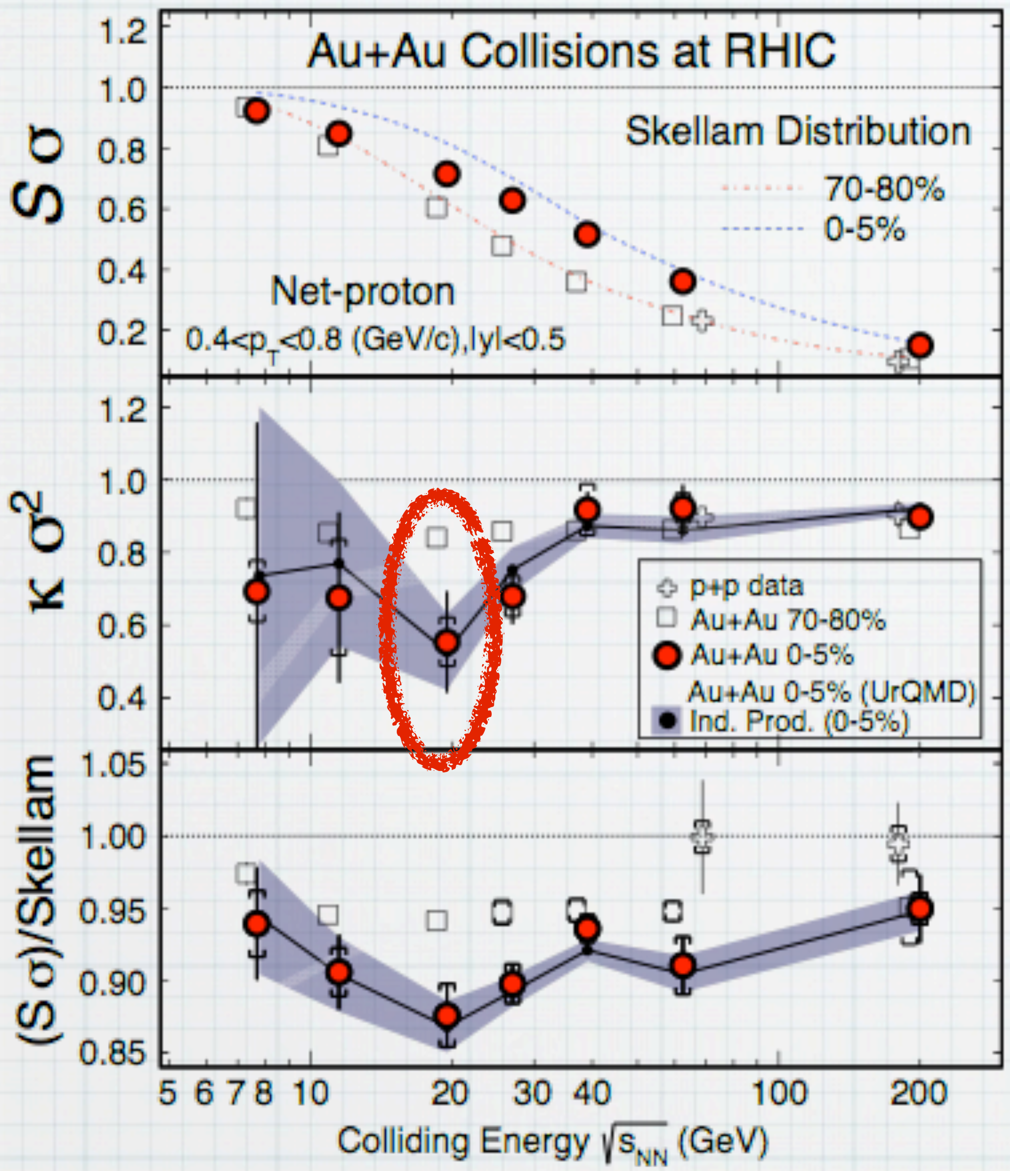
$M = R^{\beta}\theta$, $t = R(1-\theta^2)$, $H = R^{\beta\delta}h(\theta)$

$K_4 = \langle M^4 \rangle$ $(t, H) \to (\mu - \mu_{CP}, T - T_{CP})$

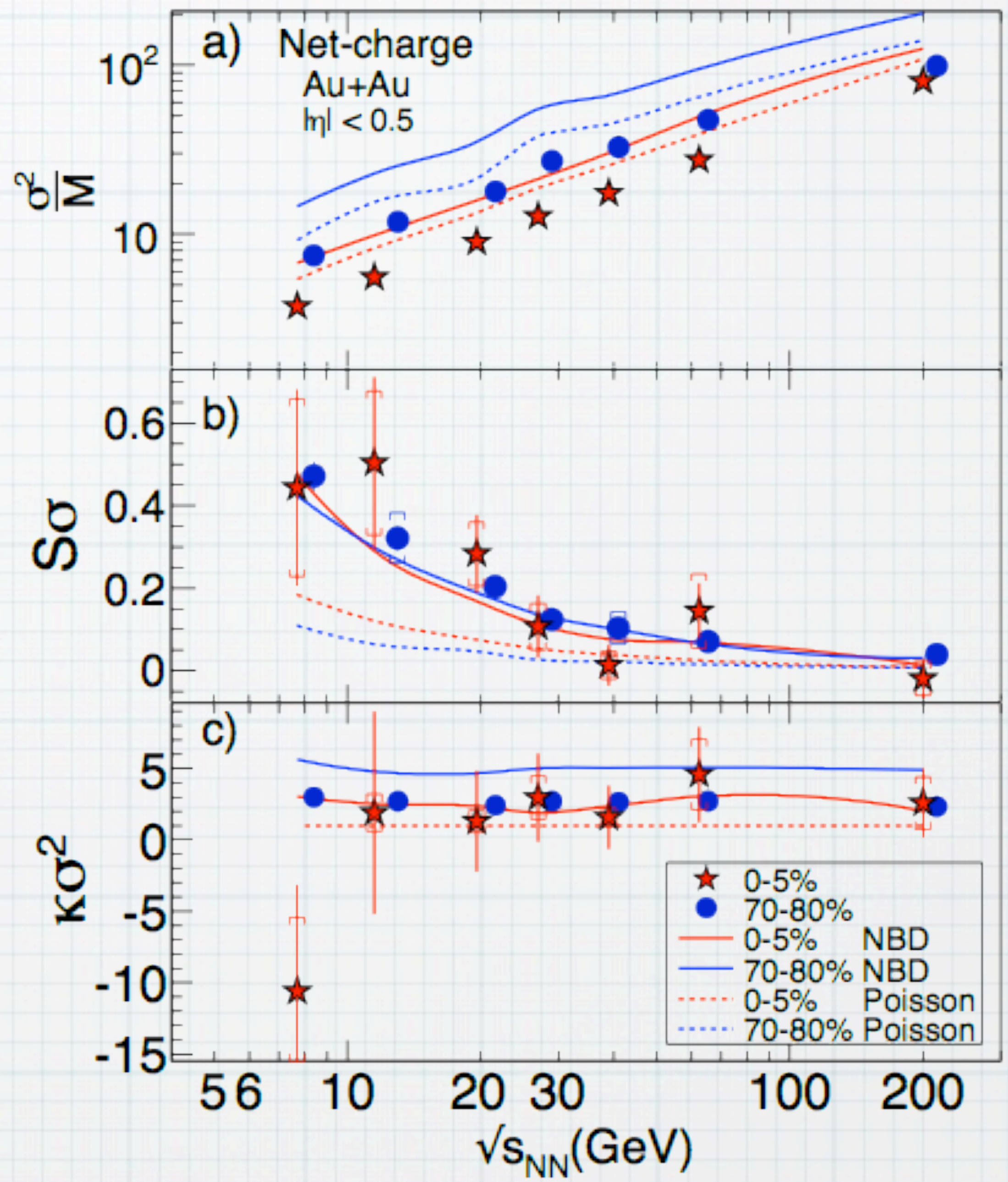
Consequently, the experimental 4th order fluctuation will be **smaller than its Poisson value** (precise value depends on ξ, on how close the freeze-out occurs to the critical point...)

$$\langle(\delta N)^4\rangle = \langle N \rangle + \langle \sigma^4_V \rangle \ldots$$

Experimental results on kurtosis



STAR Collab.: 1309.5681



STAR Collab.: 1402.1558

Kurtosis of Net-protons shows anomalous dip at $\sqrt{s_{NN}}$ = 19 GeV. Not confirmed by kurtosis of Net-charge.

Conclusions

- QCD transition: a smooth crossover at μ_B=0; expected to become first order at large μ_B (critical point)
- Lattice QCD simulations: equilibrium, thermodynamic quantities at small μ_B.
- HRG model: good description below the transition. Fit of hadron yields and ratios -> freeze-out parameters
- Alternative: fluctuations of conserved charges. Determination of freeze-out parameters from first principles
- Fluctuations at the critical point: expected to scale with some power of correlation length
- The kurtosis changes sign in the vicinity of the critical point