

Monte Carlo modelling issues in top physics

[SM@LHC 2013 conference – Albert-Ludwigs-Universität Freiburg]

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→ *mainly discussed in context with AFB.*

- *Colour coherence effect.*
- *Inclusion of top quark decays.*

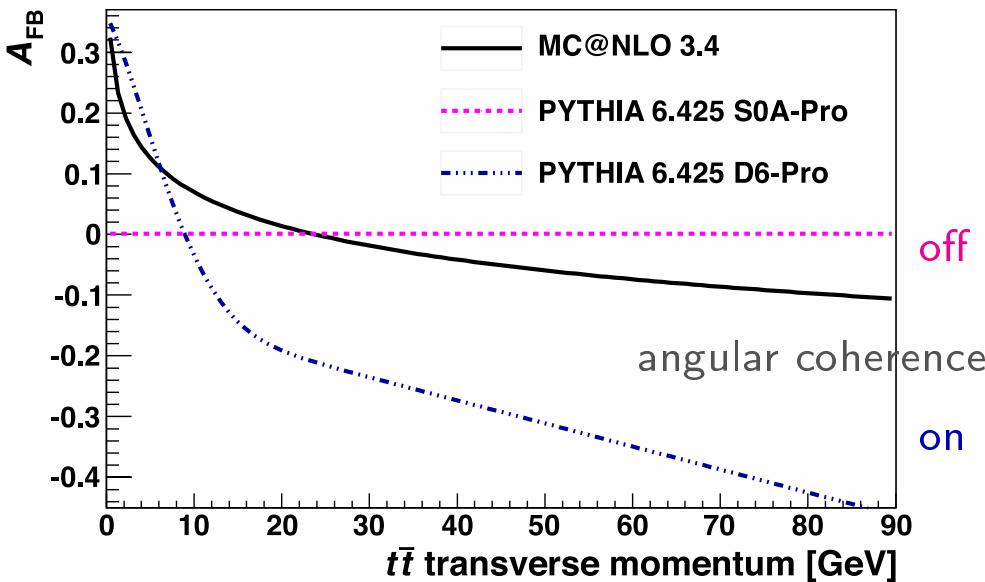
Monte Carlos for top physics

- Standard nowadays is NLO+PS for single top quark, top quark pair and associated production.
 - PowhegBox (and interface to showers). [NASON ET AL.]
 - MC@NLO and aMC@NLO (and interface to showers). [FRIXIONE ET AL.]
 - Sherpa with virtual corrections from BlackHat/GoSam/OpenLoops etc. [KRAUSS ET AL.]
- A number of questions can be asked.
 - Status of validation
 - Level of agreement in differential distributions (Do we understand differences?)
 - Comparison to fixed, higher-order calculations
 - Comparison to merging approaches, MEPS@(N)LO and MENLOPS
 - Assessment of uncertainties
 - Treatment of top quark mass
 - Treatment of top quark decays
 - Treatment of parton showers
 - ...
- Here I will only be able to discuss last few points.

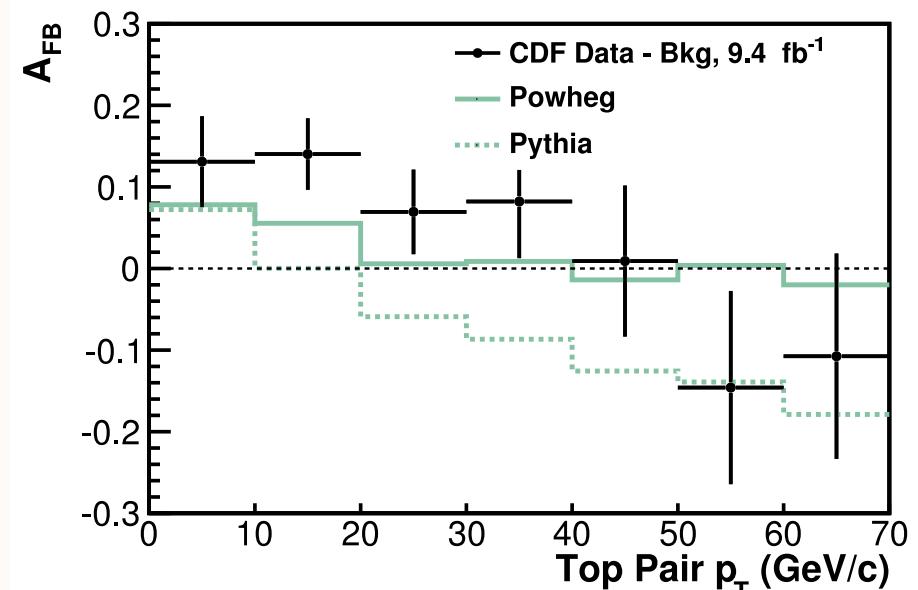
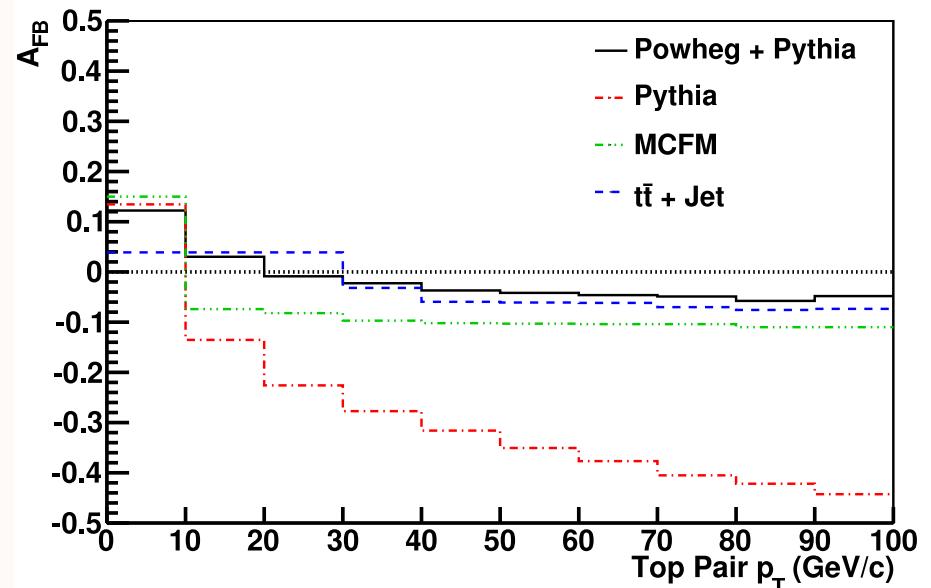
A new benchmark observable for Monte Carlos ?

→ “Observation” first made by DØ.

[PLOT UPPER LEFT FROM DØ – ARXIV:1107.4995]



[PLOTS BELOW FROM CDF – ARXIV:1211.1003]



- coherent shower Monte Carlos (MCs) contain/enhance asymmetry

$$A_{FB}(O) = \frac{\frac{d\sigma}{dO} \Big|_{\Delta y > 0} - \frac{d\sigma}{dO} \Big|_{\Delta y < 0}}{\frac{d\sigma}{dO} \Big|_{\Delta y > 0} + \frac{d\sigma}{dO} \Big|_{\Delta y < 0}}$$

- rapidity difference:
 $\Delta y = y_t - y_{\bar{t}}$

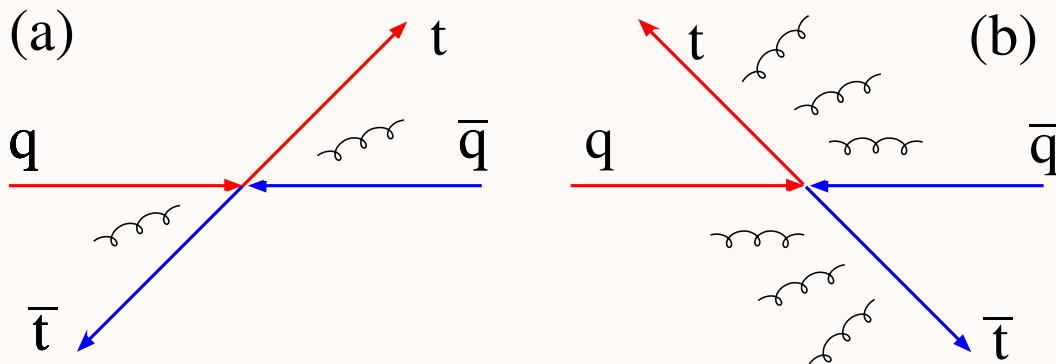
Colour coherence effect

In LO $q\bar{q} \rightarrow t\bar{t}$, there are (IF) colour flows from incoming quark to top quark and vice versa.

“Forward” dipoles – less space space for emission, less likely to radiate.

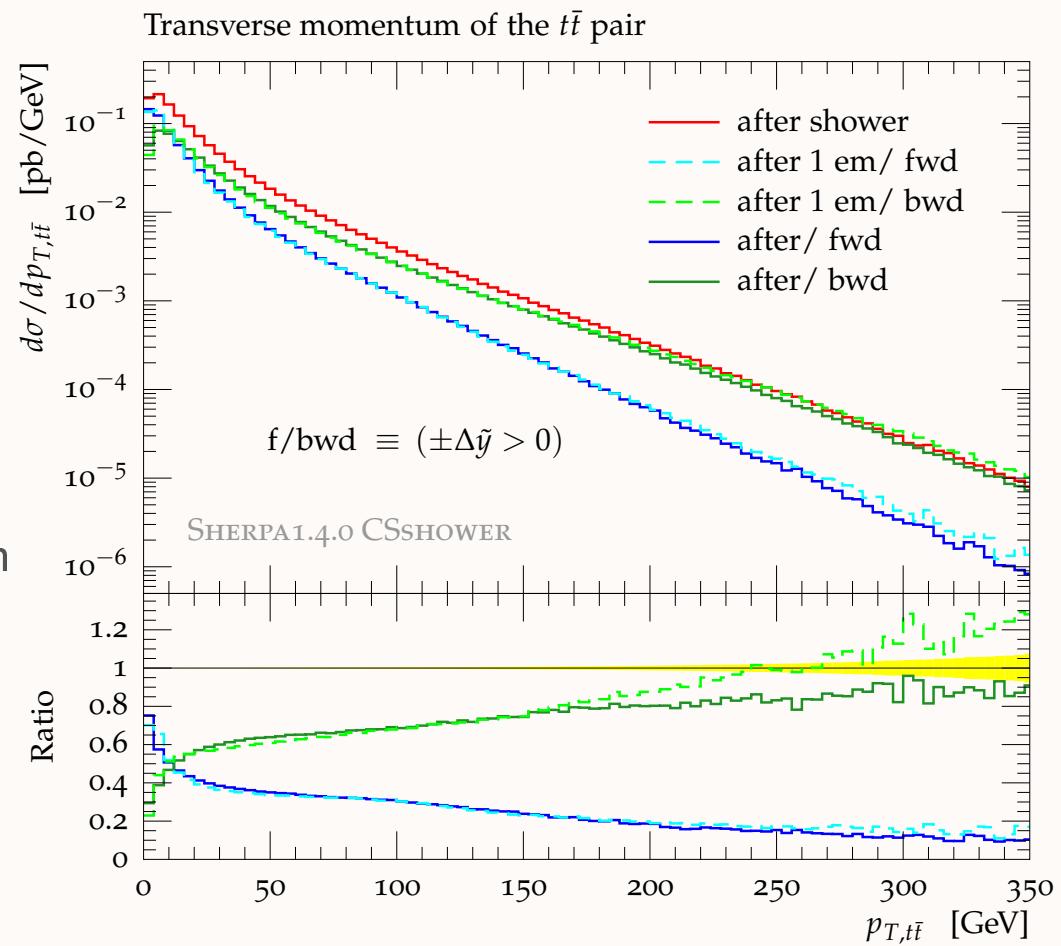
“Backward” dipoles – more violent acceleration of colour, hence more QCD radiation.

- (a) Forward configuration



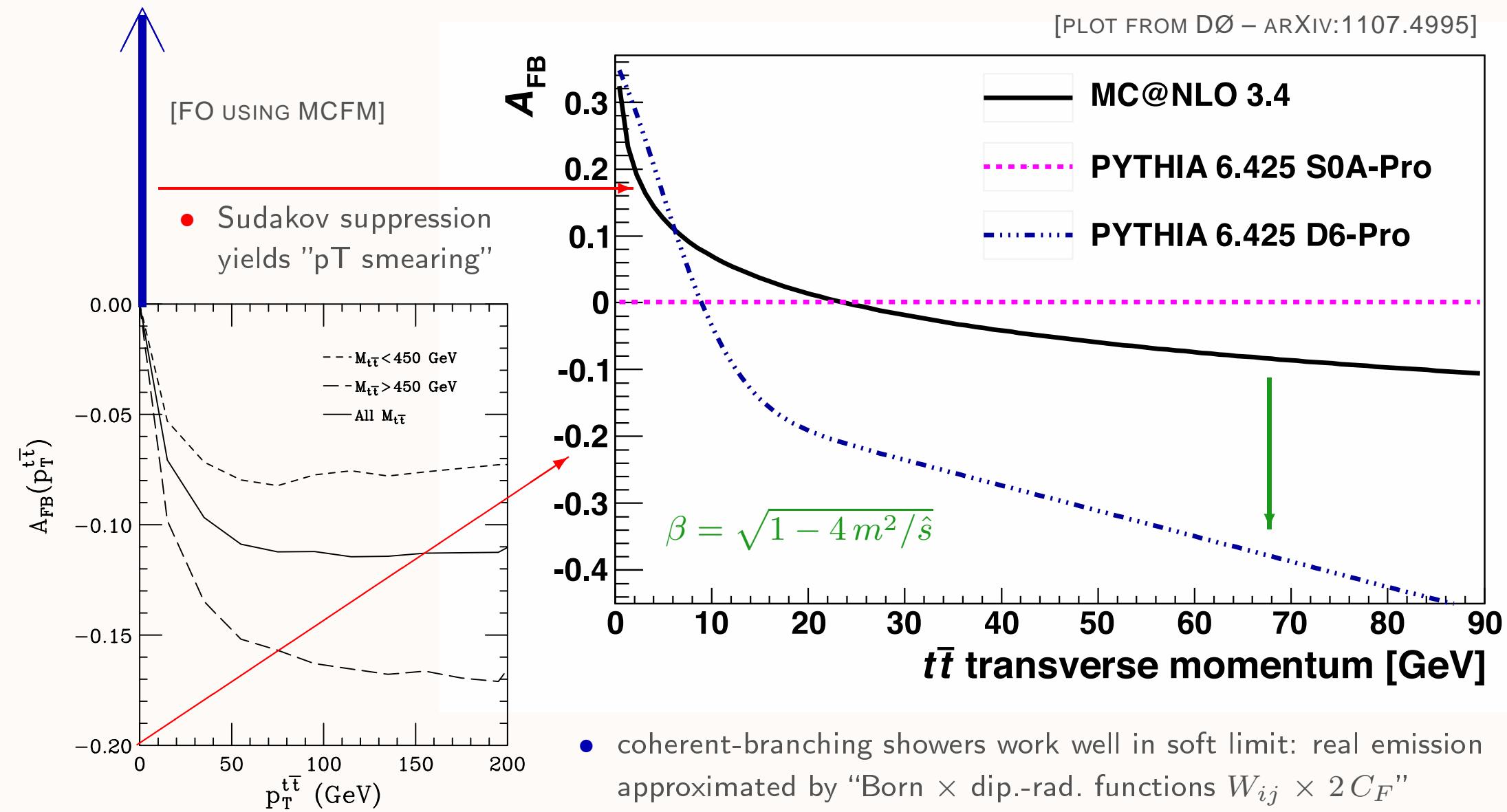
- (b) Backward configuration

- Colour coherence (leads to angular ordering)
 - asymmetric real emission
- Extra emission (confined to cone as large as θ_{ini})
 - more recoil for backward top pairs



Comparison with fixed order

[PLOT FROM DØ – ARXIV:1107.4995]



$$\frac{p_T}{\hat{\sigma}_B} \frac{d\hat{\sigma}_A}{dp_T} = \frac{\alpha_S}{\pi} \frac{(N^2 - 4)}{N} F(\beta, p_T)$$

- coherent-branching showers work well in soft limit: real emission approximated by “Born \times dip.-rad. functions $W_{ij} \times 2C_F$ ”
- LO MCs replace $(N^2 - 4)$ by $(N^2 - 1) \Rightarrow 60\%$ overestimate
- And, $F(\beta, p_T)$ underestimated by $F(\beta, 0) = -4\beta - \beta^3 - \dots$

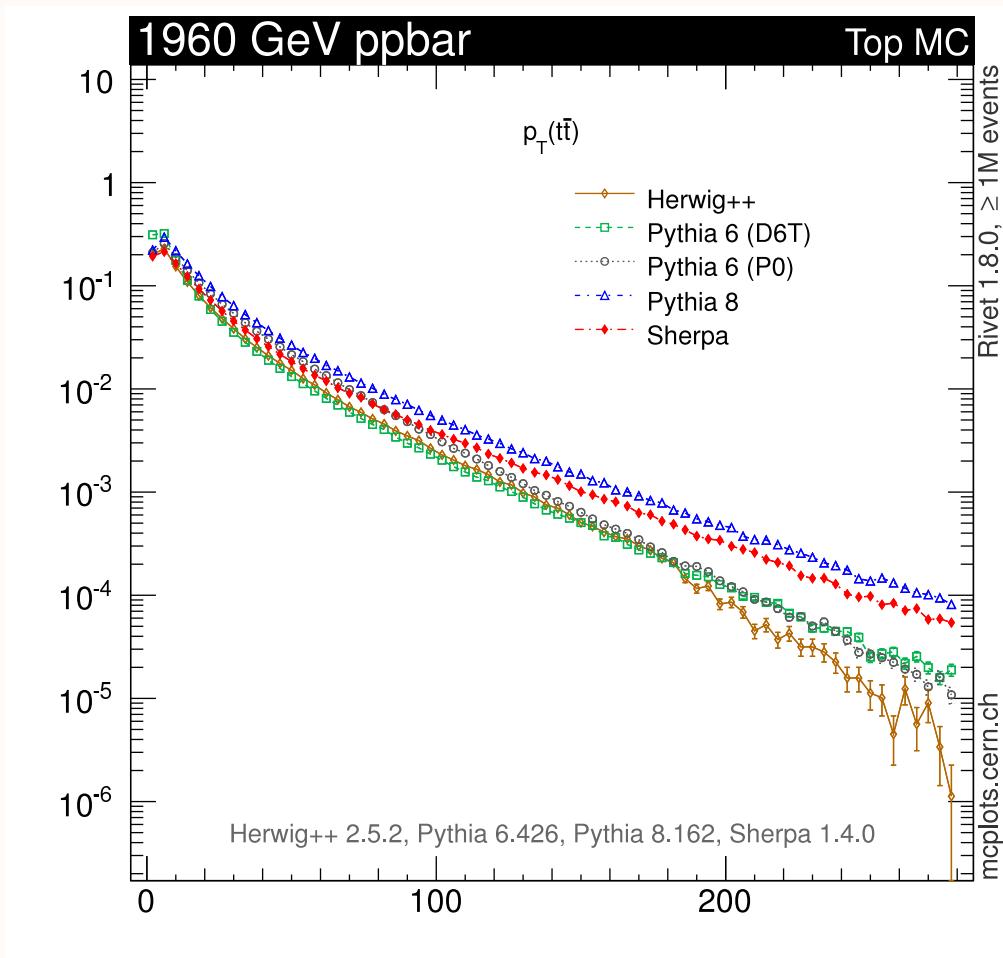
Differential asymmetry produced by LO generators

[SKANDS, WEBBER, WINTER, JHEP 07 (2012) 151 (ARXIV:1205.1466)]

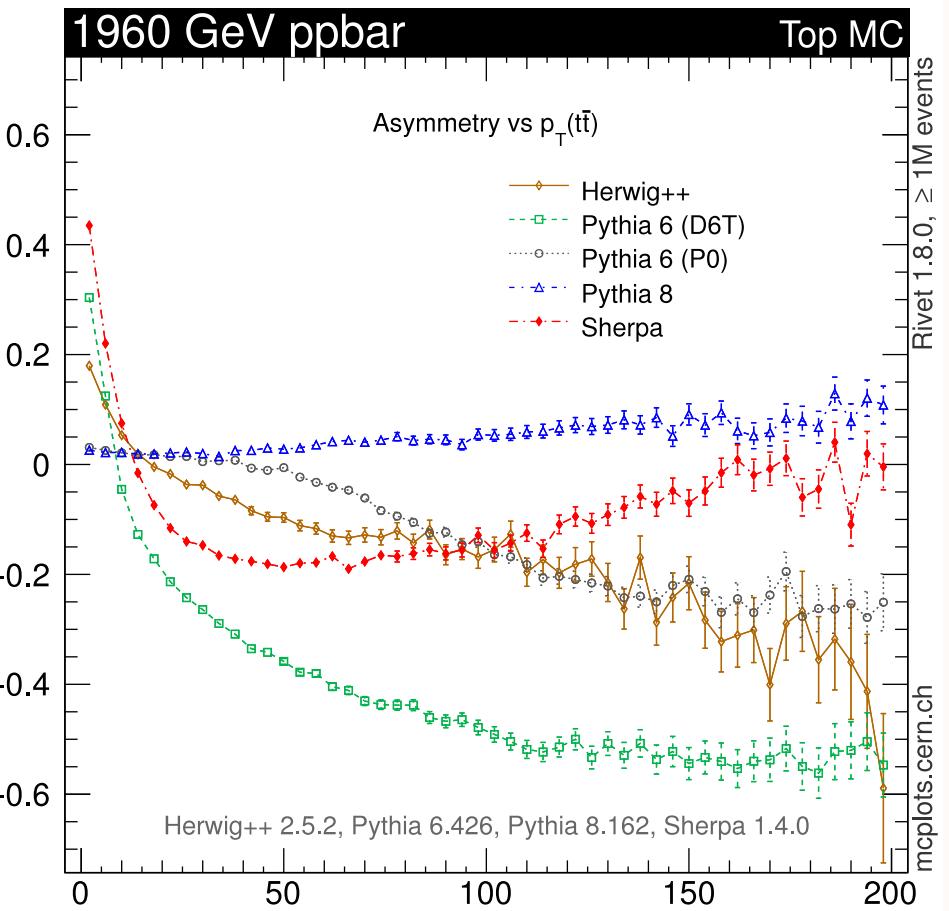
QCD coherence built in for Herwig++ and Sherpa, Pythia 6 has options with varying amounts of coherence.

→ Pythia 8 version used here does not have QCD coherence implemented

LO $q\bar{q}, gg \rightarrow t\bar{t}$ production and showering



Asymmetry as function of the top-pair p_T



- $p_{T,t\bar{t}}$ differential cross section distribution

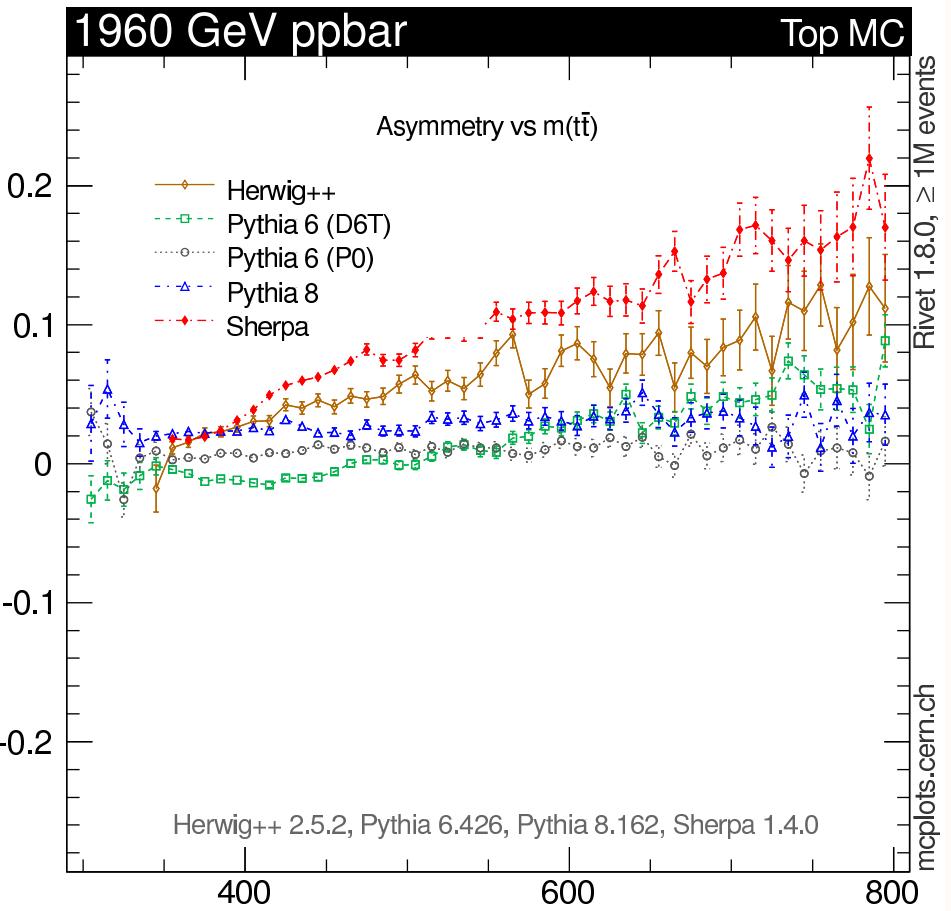
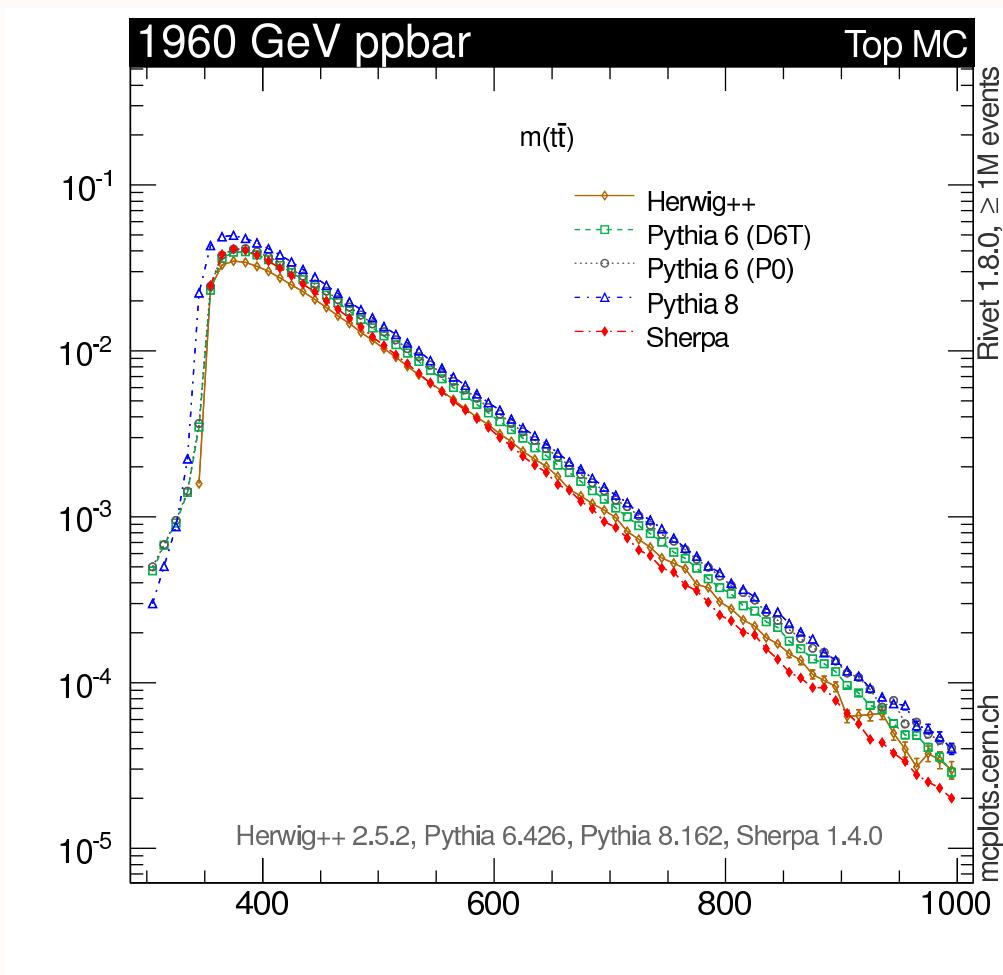
Differential asymmetry produced by LO generators

[SKANDS, WEBBER, WINTER, ARXIV:1205.1466]

→ mass dependence driven by dependence on Δy and $\Delta\phi$, Sudakov region applies over entire mass range

$$m_{t\bar{t}}^2 = m_t^2 + m_{\bar{t}}^2 + 2 E_{T,t} E_{T,\bar{t}} \cosh \Delta y - 2 p_{T,t} p_{T,\bar{t}} \cos \Delta\phi$$

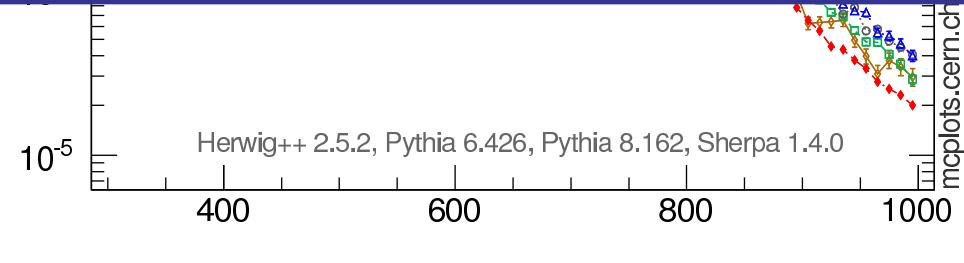
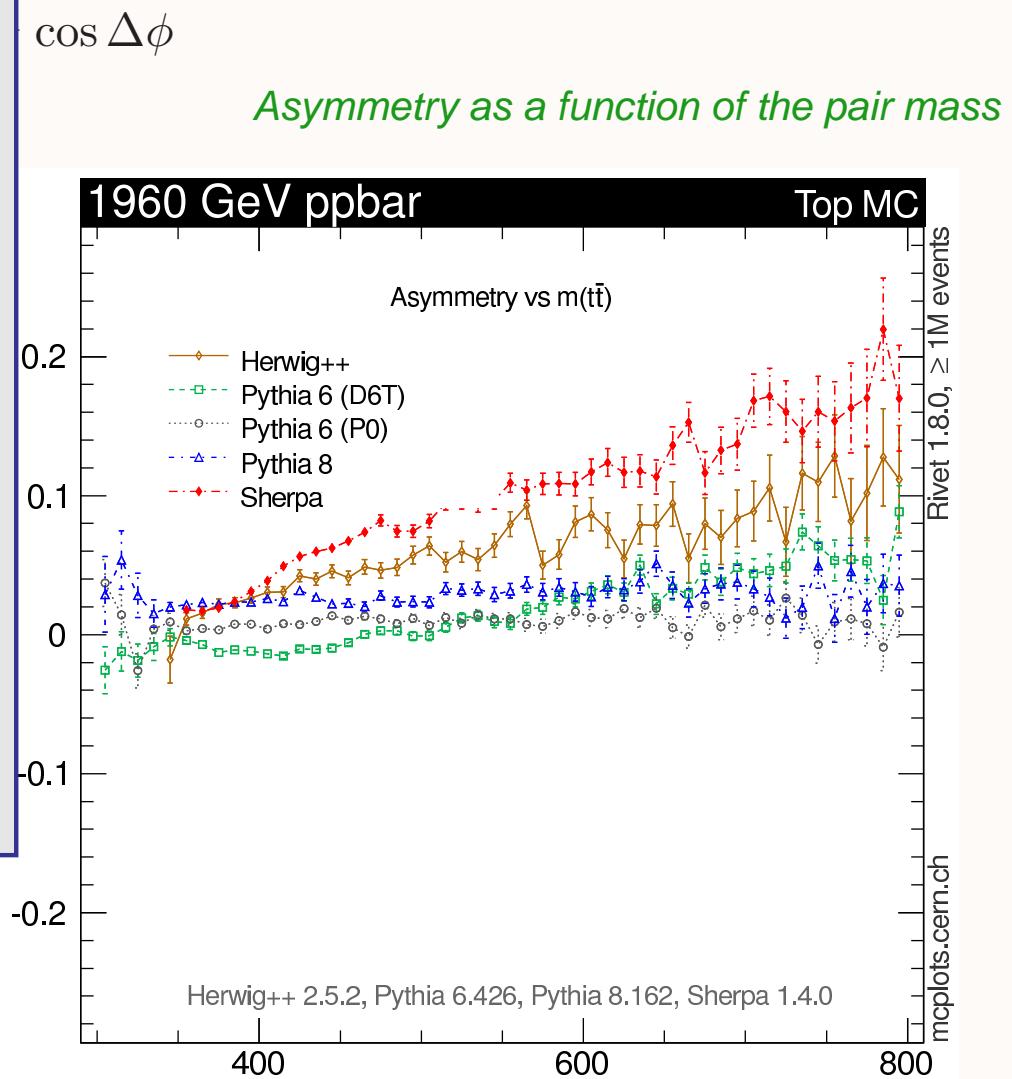
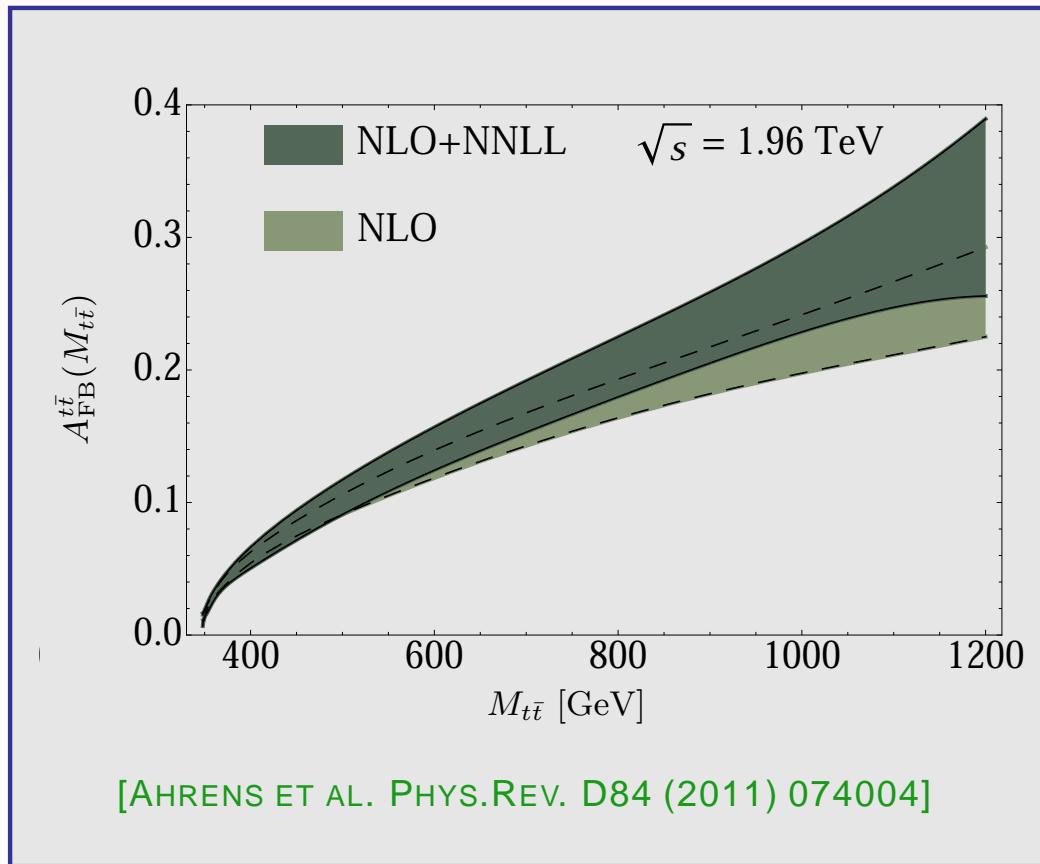
Asymmetry as a function of the pair mass



Differential asymmetry produced by LO generators

[SKANDS, WEBBER, WINTER, ARXIV:1205.1466]

→ mass dependence driven by dependence on Δy and $\Delta\phi$, Sudakov region applies over entire mass range



- $m_{t\bar{t}}$ differential distribution

Inclusive asymmetry produced by LO generators

Model	Version	Inclusive	$m_{t\bar{t}}/\text{GeV}$		$p_{T,t\bar{t}}/\text{GeV}$		• different recoil strategies implemented in the different models
			< 450	> 450	< 50	> 50	
HERWIG++	2.5.2 [def]	3.9	2.7	6.0	5.8	-14.3	
PYTHIA 6	6.426 [def]	-0.1	-0.8	1.2	2.5	-42.5	
PYTHIA 6	6.426 [D6T]	-0.2	-1.1	1.2	3.2	-43.4	
PYTHIA 6	6.426 [P0]	0.8	0.7	1.1	1.8	-8.6	
PYTHIA 8	8.163 [def]	2.5	2.4	2.8	2.4	4.8	
SHERPA	1.4.0 [def]	5.5	3.5	9.2	8.7	-15.4	
SHERPA	1.3.1 [def]	6.3	3.3	12.1	9.6	-15.8	
QCD	LO	6.0	4.1	9.3	7.0	-11.1	

- *If shower kinematics allow for migrations, a net inclusive asymmetry can be generated.*
- showers are unitary, preserving the total inclusive cross section (LO)
- BUT asymmetry is not protected by unitarity
 \Rightarrow migration from $\Delta y > 0$ to $\Delta y < 0$ and v.v.

$$A_{\text{FB}}^{(\text{cut})} = \frac{\sigma^{(\text{cut})}|_{\Delta y > 0} - \sigma^{(\text{cut})}|_{\Delta y < 0}}{\sigma^{(\text{cut})}}$$

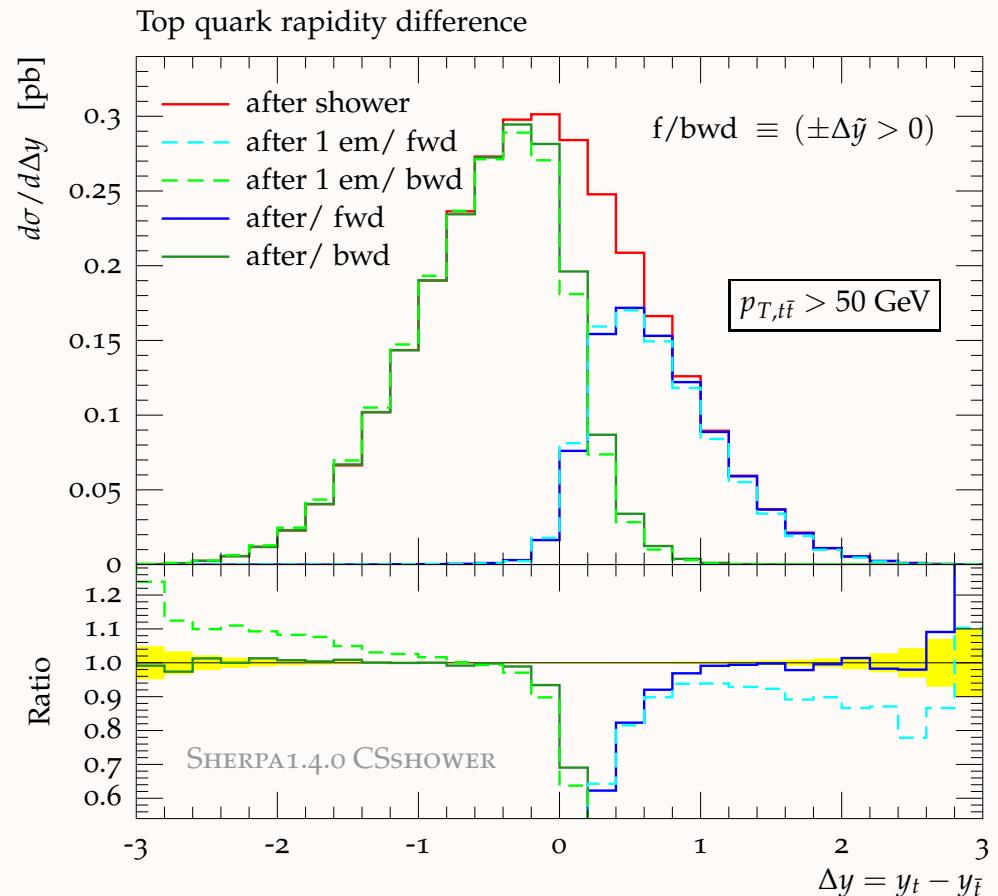
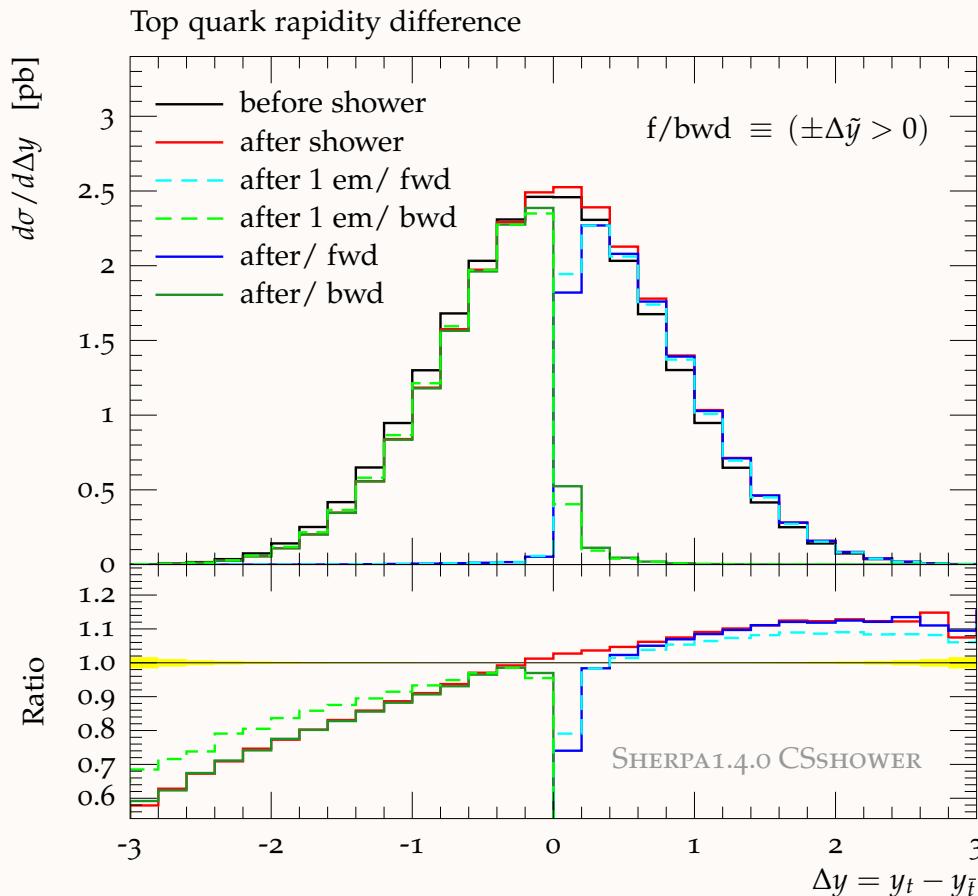
Longitudinal recoil effects – migration

[SKANDS, WEBBER, WINTER, ARXIV:1205.1466]

Simple dipole picture where gluon emission on average stretches LF dipole can give $\Delta y = \Delta \tilde{y} + \epsilon$ ($\epsilon > 0$).

→ migration is small, local, favouring $- \rightarrow +$ direction; largest effect already after 1st emission

radiation imbalance wins over more severe migration



- Δy distribution for various LO $\Delta \tilde{y}$ generation and shower modes

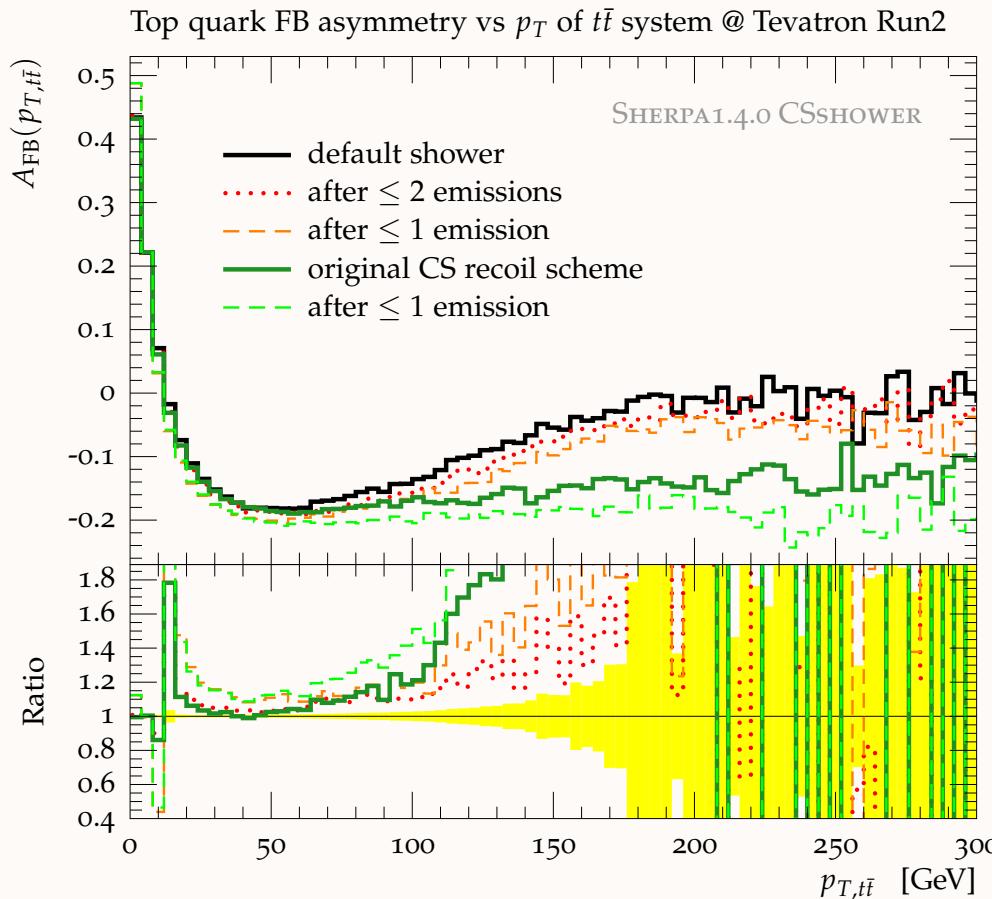
Different recoil-scheme options in Sherpa

[SKANDS, WEBBER, WINTER, ARXIV:1205.1466]

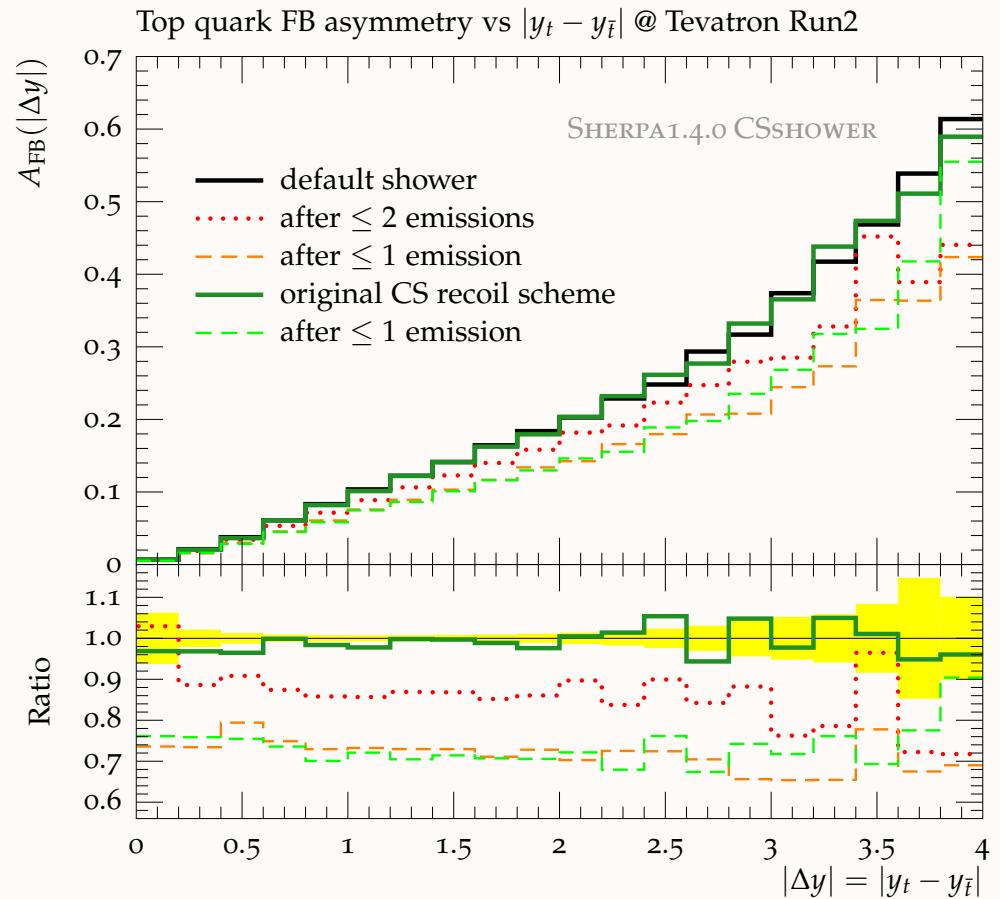
Sherpa's CSshower provides two options for treating the recoils. \Rightarrow Recoil options affect high p_T .

→ original CS scheme treats recoils more locally, IF dipole is decoupled from rest of event

longitudinal recoil treatment is effectively the same



- Asymmetry versus $p_T, t\bar{t}$



- Asymmetry versus $|\Delta y|$

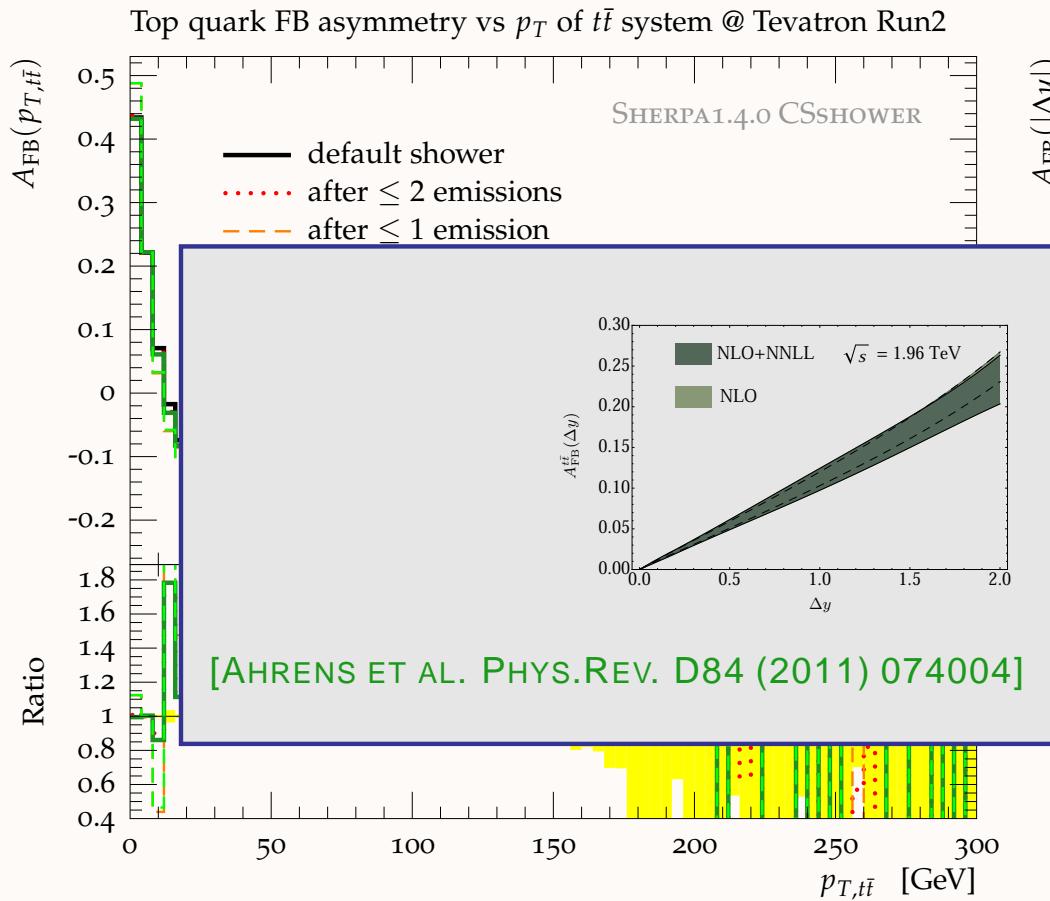
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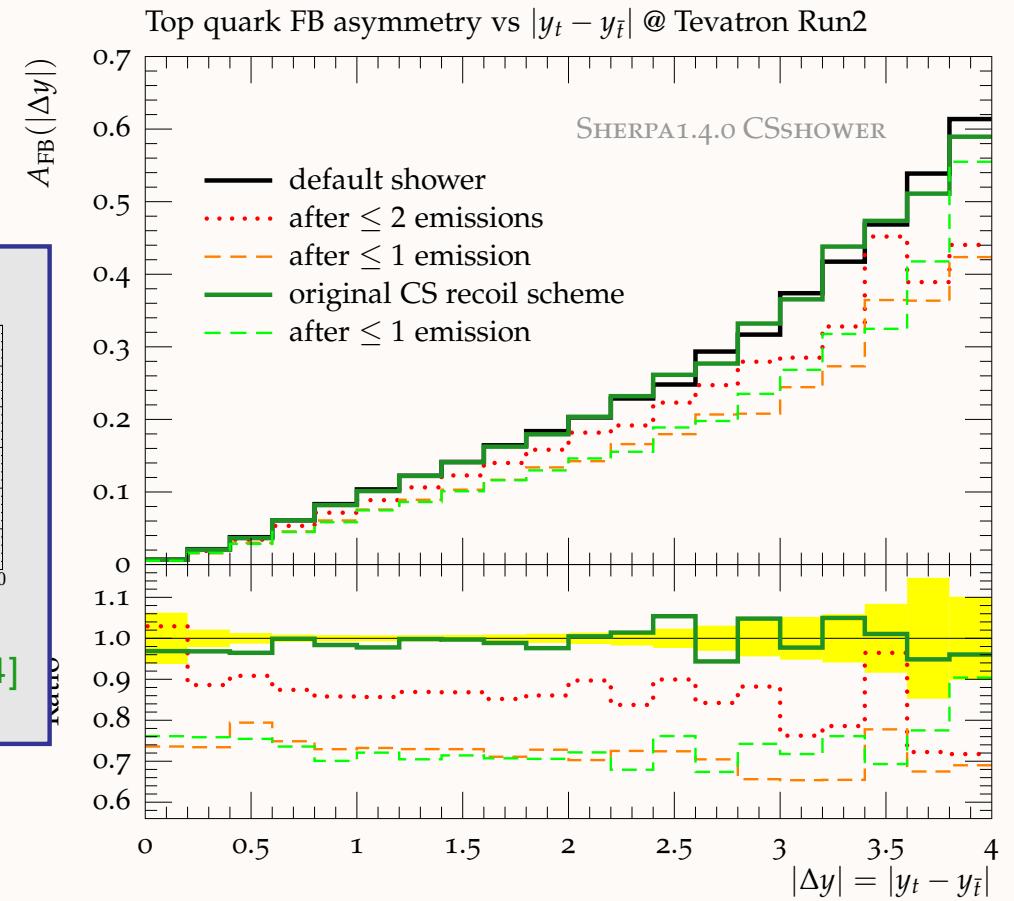
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- Asymmetry versus $p_{T,t\bar{t}}$

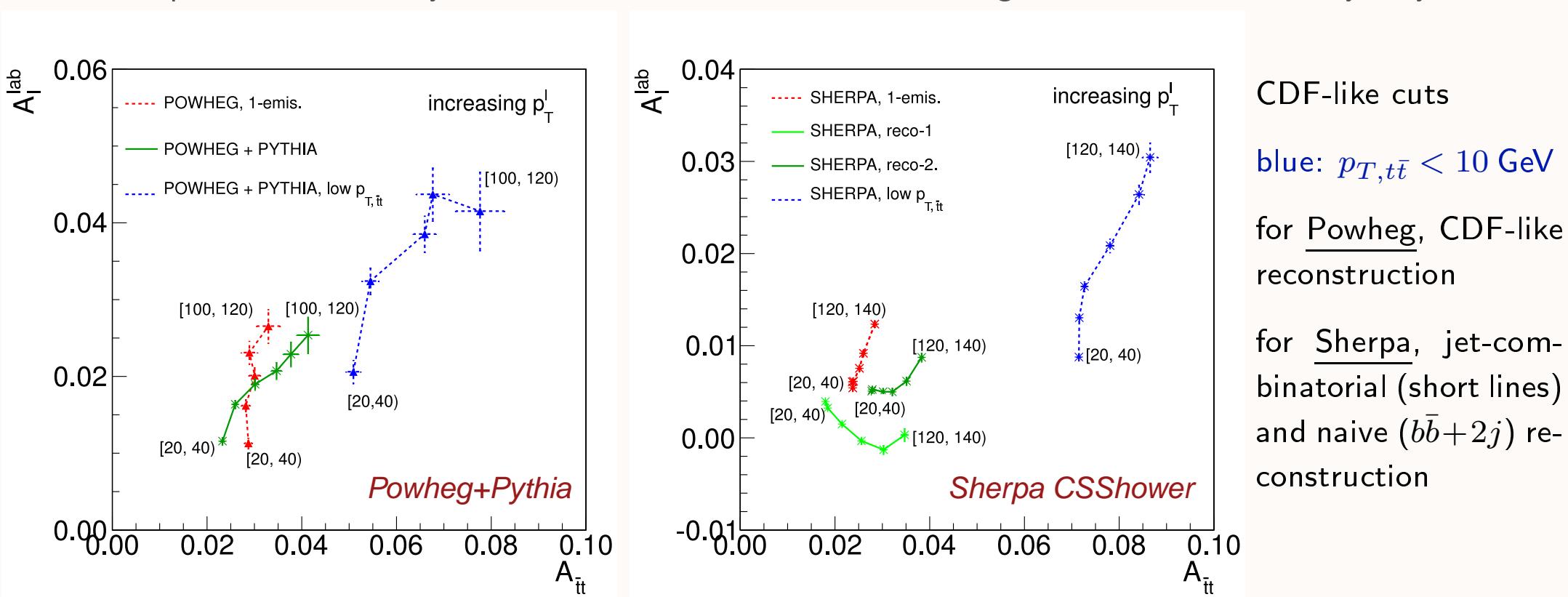


- Asymmetry versus $| \Delta y |$

Asymmetry from the lepton perspective

[FALKOWSKI, MANGANO, MARTIN, PEREZ, WINTER, ARXIV:1212.4003]

- consider top quark $A_{t\bar{t}}$ and lepton-based A_ℓ at the same time (as a function of $p_{T,\ell}$)
- gain extra handles in your analysis: exploit correlation of asymmetries, make use of “no or less” need for top quark reconstruction \Rightarrow Correlation is robust, but modelling of $p_{T,t\bar{t}}$, hence $A_{t\bar{t}}(p_{T,t\bar{t}})$ turns out to be crucial. Just colour coherence alone as in Sherpa’s shower is not sufficient.
- BSM predictions are very different but should come with warning label: LO+PS accuracy only.



→ Important for such a study, the reliable treatment of top quark decays.

Top quark decays included

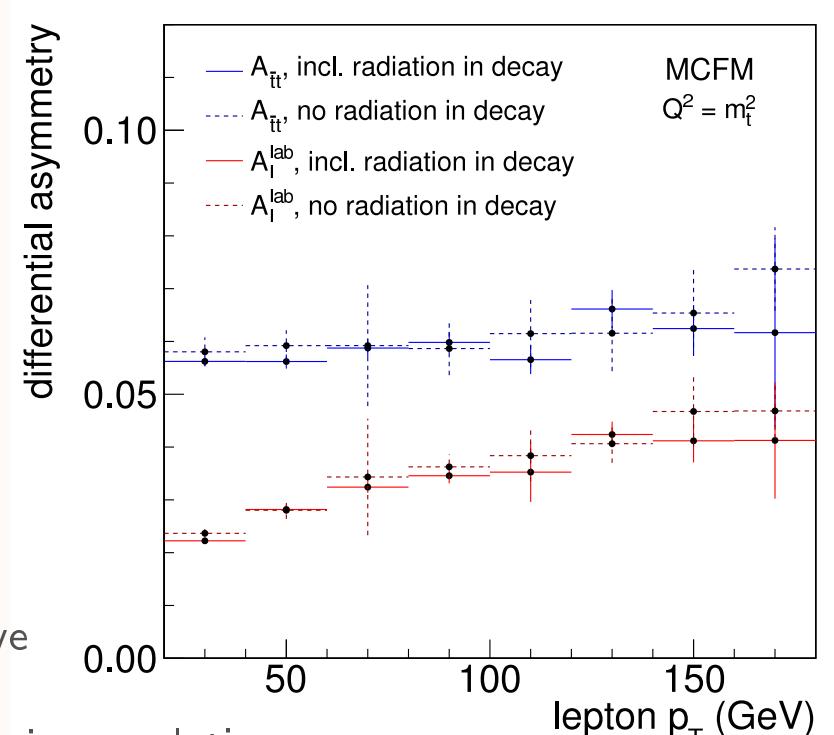
Effect of radiation in decays

- using **MCFM** [CAMPBELL, ELLIS, ARXIV:1204.1513]
- top production in pole approximation (allows for NLO)
- all spin correlations kept in tree, real and virtual MEs
- in our study, corrections on $A_{t\bar{t}}$, A_ℓ turn out to be small
- how does it compare to full $WWb\bar{b}$ computation
[DENNER ET AL, ARXIV:1012.3975, ARXIV:1204.1513]
[BEVILACQUA ET AL, ARXIV:1012.4230]
effects Γ_t/m_t or $\alpha_s \Gamma_t/m_t$ suppressed if sufficiently inclusive
- LO **Sherpa showering** includes top quark off-shellness and spin correlations
- NLO+PS, here **Powheg**, generate un-decayed top quark pairs and introduce the decay spin correlation by reweighting (approximate spin correlation mechanism)

$$\frac{d\sigma(\bar{e}, \nu, b; e, \bar{\nu}, \bar{b}, [p])}{d\Phi_{\text{dec}} d\Phi_{t\bar{t}[p]}} = \frac{d\sigma^{\text{NLO+PS}}(t; \bar{t}, [p])}{d\Phi_{t\bar{t}[p]}} \times \underbrace{\frac{\frac{d\sigma^{\text{Tree}}(\bar{e}, \nu, b; e, \bar{\nu}, \bar{b}, [p])}{d\Phi_{\text{dec}} d\Phi_{t\bar{t}[p]}}}{\frac{\sigma^{\text{Tree}}(t, \bar{t}, [p])}{d\Phi_{t\bar{t}[p]}}}}_{\text{Integrates to 1 in } d\Phi_{\text{dec}}}$$

[FRIXIONE ET AL, HEP-PH/0702198]

→ Full validation between approaches should be done to assess differences.

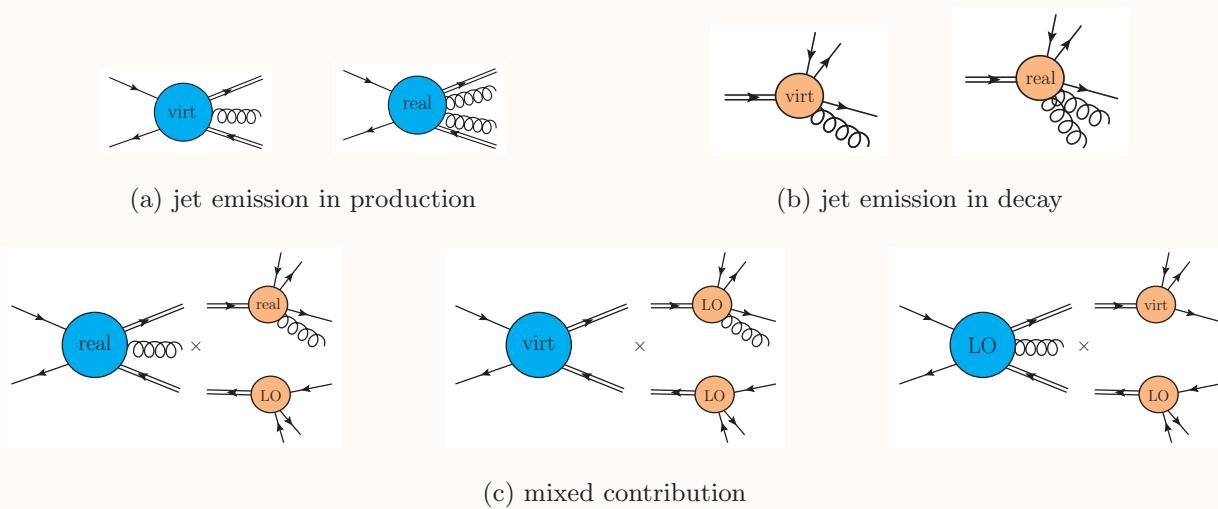


no cuts comparison

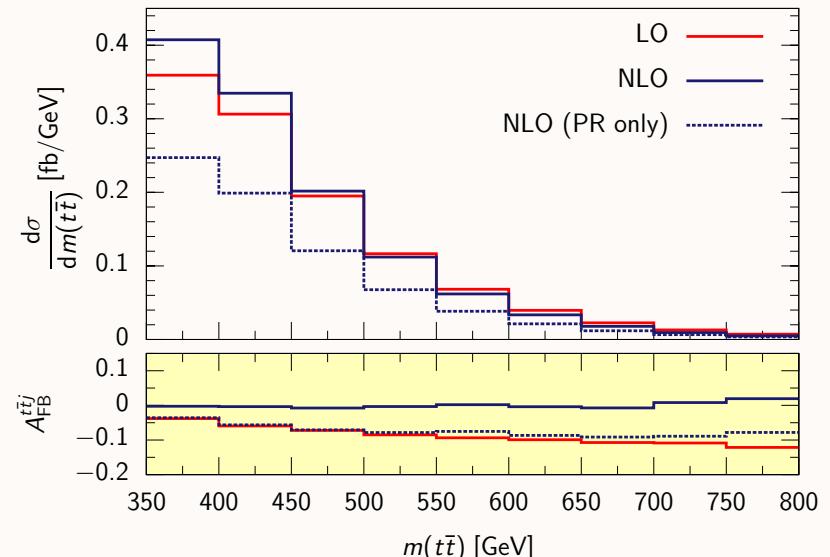
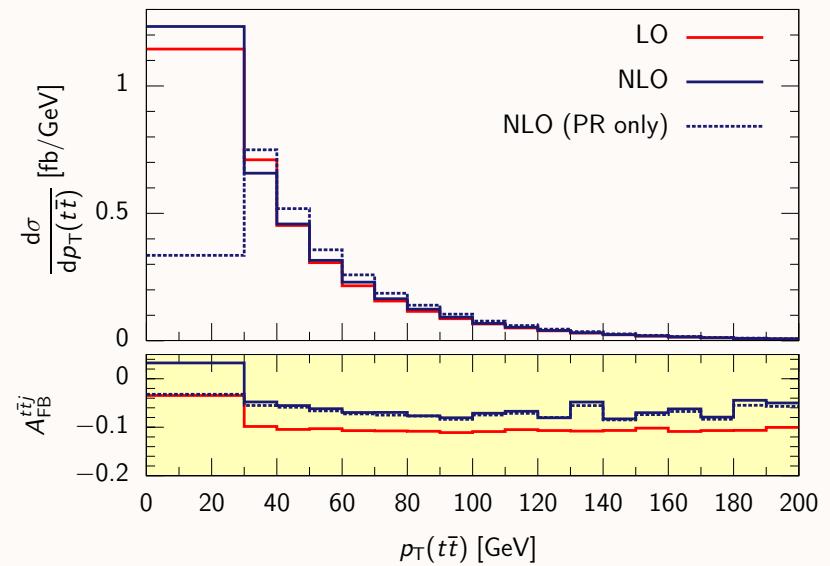
One more jet included (at NLO)

Effect of radiation in decays

- prediction based on the narrow width approximation with the NLO QCD corrections to both production and decay
- in general more complicated pattern of corrections and larger effects due to radiation in decay
- for example, the negative forward–backward asymmetry given at LO receives fairly sizeable positive corrections



[MELNIKOV, SCHARF, SCHULZE, ARXIV:1111.4991]

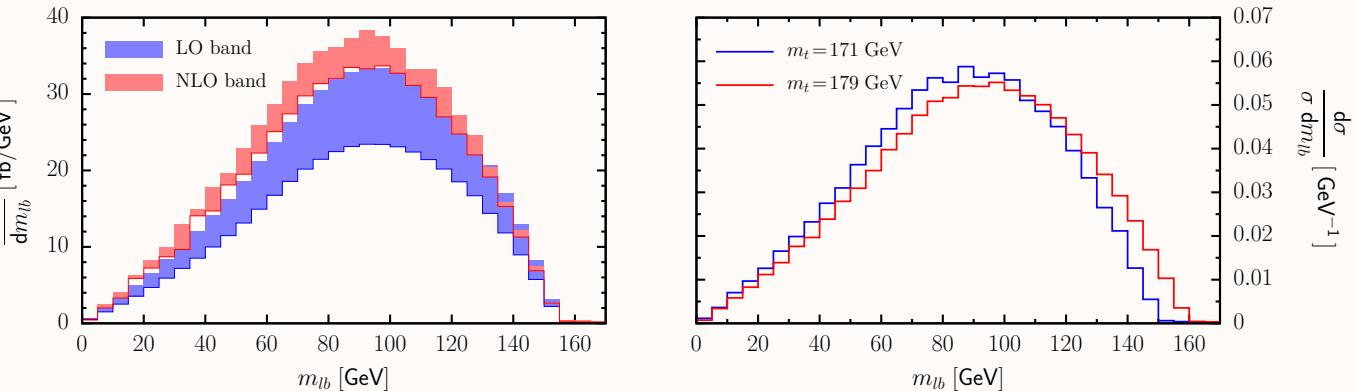


→ How do we match NLO calculations as such to parton showers.

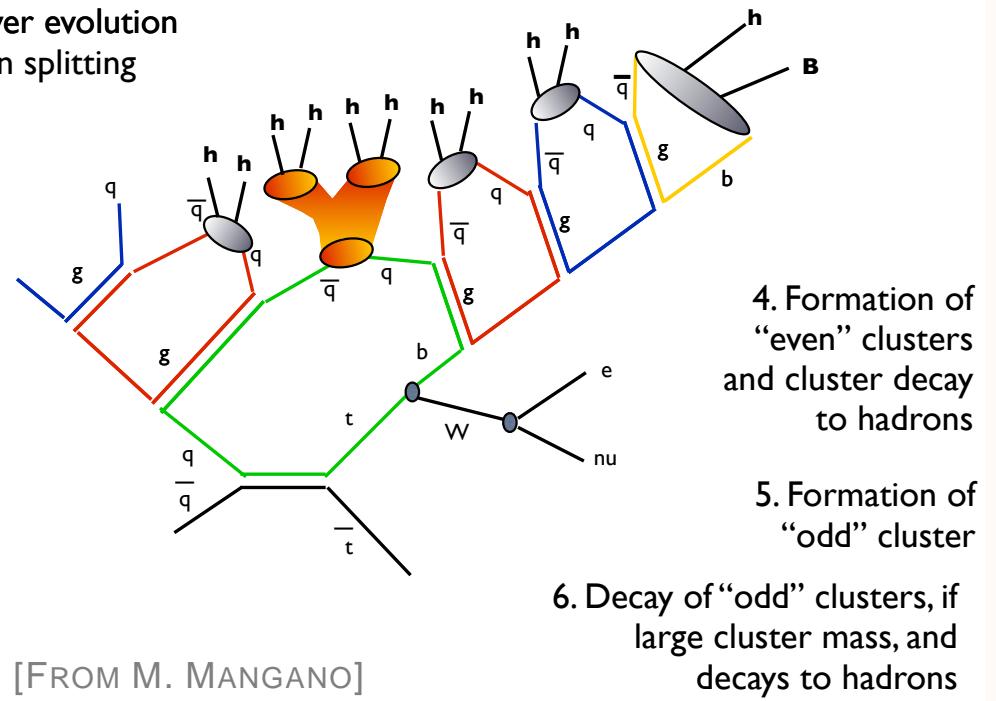
Quick note: top quark mass determination

[FROM BISWAS, MELNIKOV, SCHULZE, ARXIV:1006.0910]

- accurate and reliable measurement of top quark mass is important
- precisely measured in experiments but we have trouble to relate this mass parameter to a top quark mass we can control
- several proposals to extract it from kinematic distributions describable with pQCD (renormalization schemes can be switched)
→ S. Moch's talk
- this still requires object identification, i.e. non-perturbative corrections have to be estimated
- preliminary studies by M. Mangano indeed point to $\mathcal{O}(1)$ GeV effects



1. Hard Process
2. Shower evolution
3. Gluon splitting



[FROM M. MANGANO]

Summary & Implications

- Monte Carlo LO event generators can produce significant (differential and inclusive) asymmetries where none were previously expected.
 - One needs to be aware of that for the interpretation of experimental data.
 - Monte Carlo estimates of corrections to asymmetries could be affected by this.
 - Model-dependent corrections!?
- Asymmetries in Monte Carlos arise from valid physics built into generators with coherent parton or dipole showering.
- While not quantitatively correct in every detail, important features are captured by coherent showering approximation.
 - unequal Sudakov factors for forward and backward top production
 - migration of recoiling tops between hemispheres (Use (N)LO for A_{FB} to optimize recoils?)
 - from modelling point of view we are safe since standard is NLO+PS
 - but shall we neglect contributions from secondary emissions?
- Inclusion of top quark decays complicates the picture, more checks and work needed.