



SPIN DETERMINATION WITH INVISIBLE PARTICLES

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OUTLINE

- Introduction
- Review of Wigner d functions, spin and the Feynman amplitude
- Kinematical method
- The angular distributions
- Summary and conclusion

DARK MATTER AT A COLLIDER

- Substantial astronomical evidence points towards the existence of a nonluminous form of matter
- We anticipate dark matter production at a collider
- What are the properties of dark matter?
- We focus on the International Linear Collider (ILC)

SPIN

- Given total initial angular momentum j and initial net spin m along the z axis and final net spin m' along the z' axis then we define the Wigner d function as

$$d_{m,m'}^j(\theta) = \langle j, m', \theta | j, m \rangle = \langle j, m' | \exp(iJ_y\theta) | j, m \rangle$$

where the reaction takes place in the x - z plane, and θ is the angle between the z and z' axes in the CM frame

- So for a particular spin combination m, m'

$$\mathcal{M}_{m,m'} = C_{m,m'} d_{m,m'}^j(\theta)$$

- The cross section is then

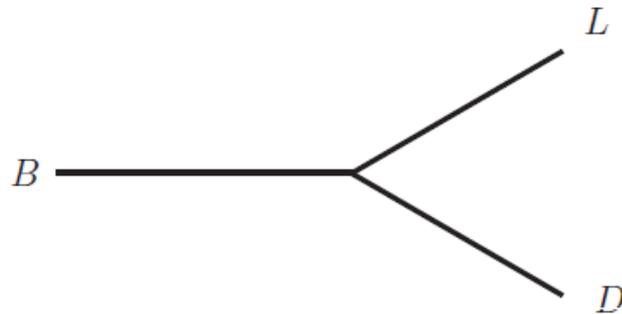
$$\sigma \propto \sum_{m,m'} |\mathcal{M}_{m,m'}|^2 = \sum_{m,m'} |C_{m,m'}|^2 (d_{m,m'}^j(\theta))^2$$

SPIN DIRECTLY FROM THE VERTEX

- In principle, we want to proceed as directly as possible to the dark matter vertex
- We need to calculate the angle of L with respect to B in the B 's CM frame:

$$\cos \theta_{LB} = \frac{\vec{p}_B \cdot \vec{p}_{LBCMS}}{|\vec{p}_B| |\vec{p}_{LBCMS}|}$$

- From this, we can reconstruct which Wigner d functions are responsible and thus which spins are involved

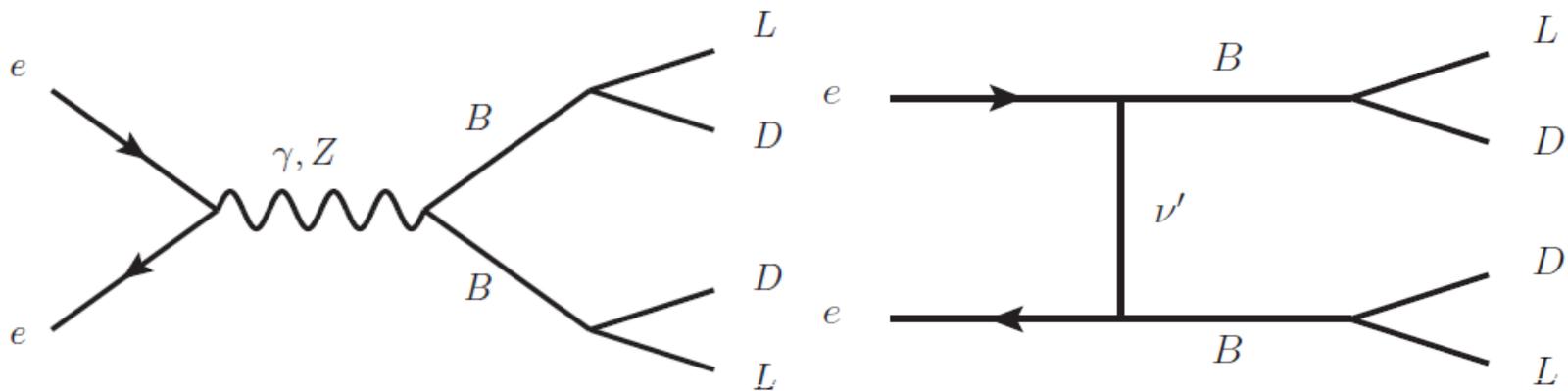


WHY THE ILC

- The ILC is nice because
 - The collision CM frame is the LAB frame
 - We know \sqrt{s}
 - Have the ability to polarize the incoming beams

SOME CONSIDERATIONS

- We assume the particles B , L and D are all on their mass shells
- The energy and momentum of L are completely known
- The masses of the B and D particles are known
 - Antler Diagram Topology
 - T. Han, I.-W. Kim and J. Song, Phys. Lett. B **693**, 575 (2010) [arXiv:0906.5009 [hep-ph]].
 - For example, in the MSSM, $e^+e^- \rightarrow \tilde{\mu}^+\tilde{\mu}^- \rightarrow \mu^+\mu^-\chi_1^0\chi_1^0$

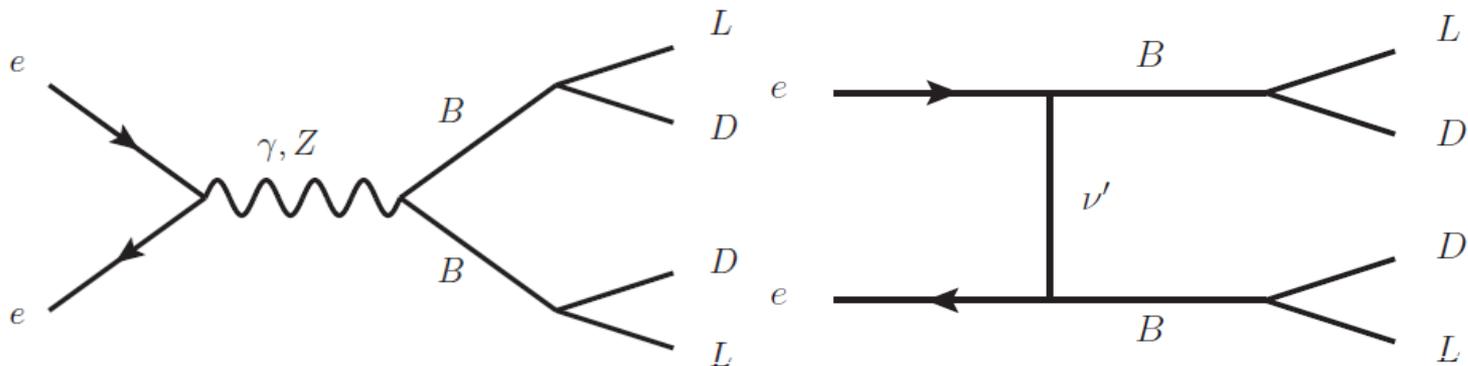


KINEMATICS

- The angle of L with respect to B :

$$\cos \theta_{LB} = \frac{\vec{p}_B \cdot \vec{p}_{LBCMS}}{|\vec{p}_B| |\vec{p}_{LBCMS}|}$$

- We found that we can calculate the magnitude of $\cos \theta_{LB}$ and partially reconstruct the sign
- To be published soon

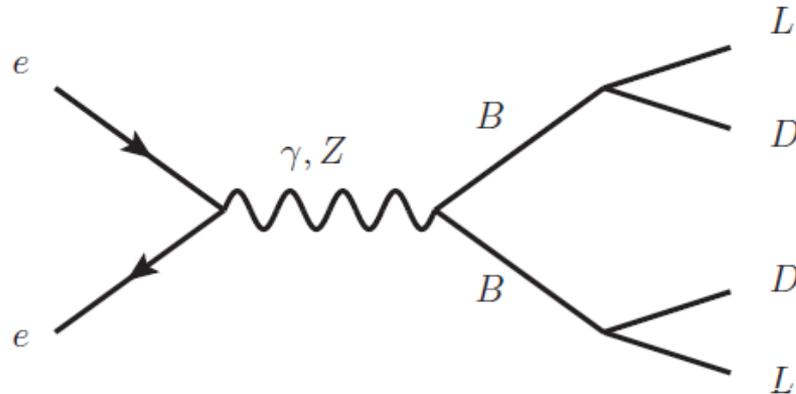


NOTES ABOUT NOTATION

- All of the processes we will consider here proceed via

$$e^+e^- \rightarrow \gamma/Z \rightarrow B^+B^- \rightarrow \mu^+\mu^- DD$$

- Something like $sB \rightarrow \mu fD$ is simply this process in which we analyze the specific case where B is a spin-0 particle (sB) and D is a spin-1/2 particle (fD)



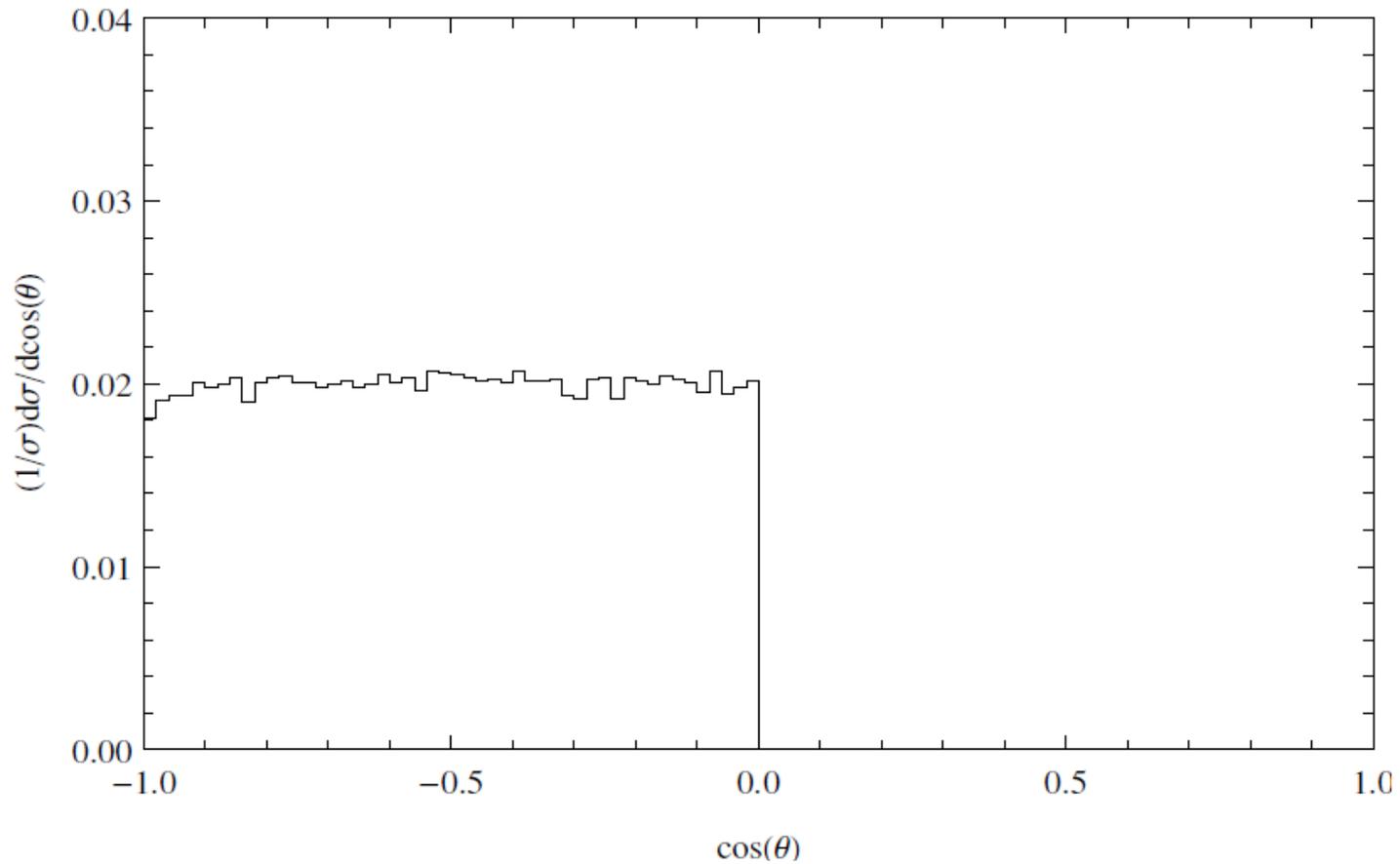
ANGULAR DISTRIBUTIONS ($sB \rightarrow \mu f D$)

- For a scalar particle, there is no spin correlation
- The Wigner d function is

$$d_{0,0}^0(\theta) = 1$$

- We expect a flat line

PLOT OF $sB \rightarrow \mu f D$



ANGULAR DISTRIBUTIONS ($fB \rightarrow \mu sD$)

- The Wigner d functions associated with spin $1/2$:

$$d_{m,m'}^{1/2} = \begin{pmatrix} \cos \theta/2 & \sin \theta/2 \\ -\sin \theta/2 & \cos \theta/2 \end{pmatrix}$$

- And, of course

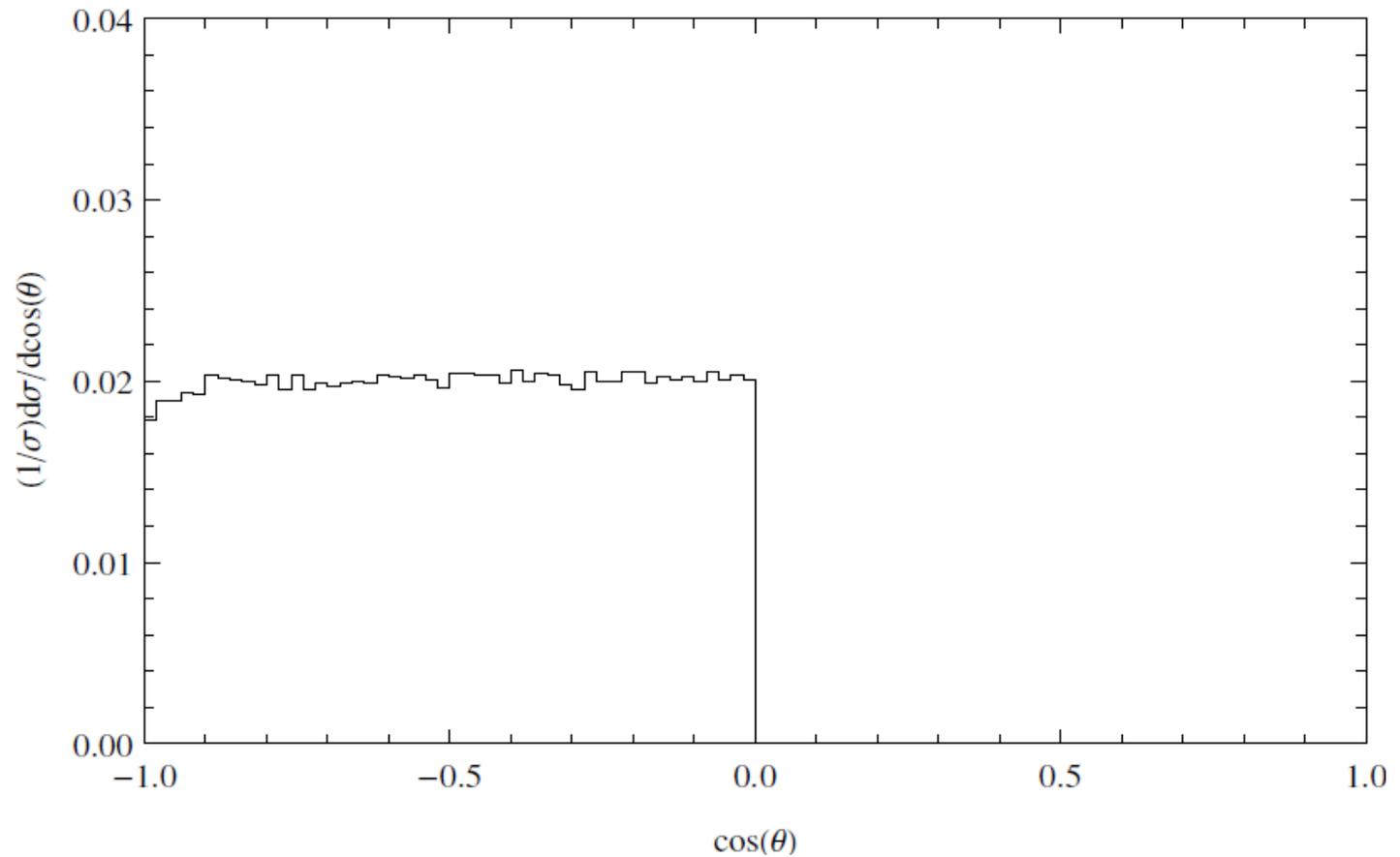
$$\sin^2 \frac{\theta}{2} = \frac{1 - \cos \theta}{2} \qquad \cos^2 \frac{\theta}{2} = \frac{1 + \cos \theta}{2}$$

- So as a linear combination we should have

$$\sum_{m,m'} |\mathcal{M}_{m,m'}|^2 = A \sin^2 \theta/2 + B \cos^2 \theta/2 = \frac{A}{2} (1 - \cos \theta) + \frac{B}{2} (1 + \cos \theta)$$

- We find that $A = B$

PLOT OF $fB \rightarrow \mu s D$



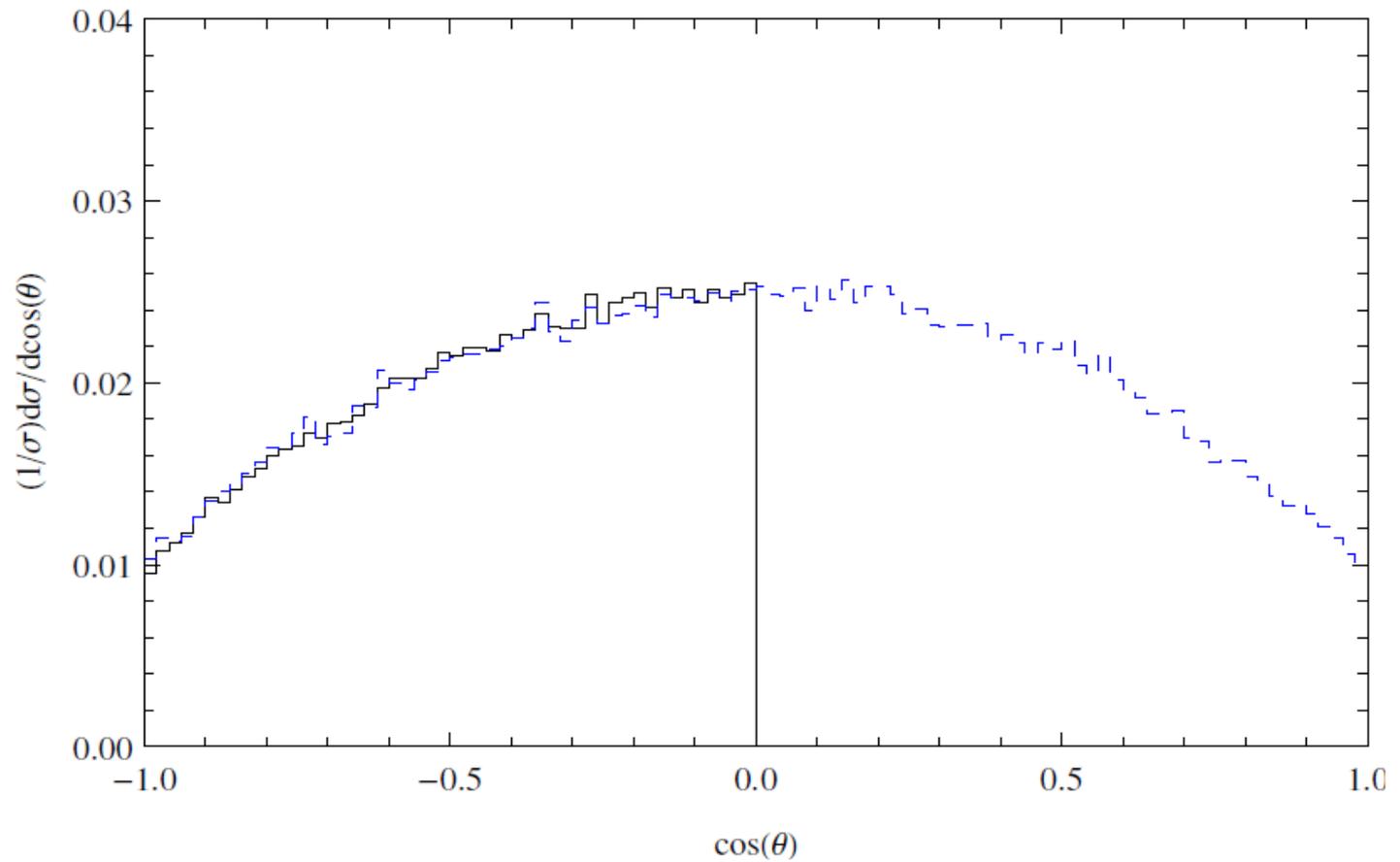
ANGULAR DISTRIBUTIONS ($\nu B \rightarrow \mu f D$)

- The Wigner d functions corresponding to the spin configurations are

$$d_{m,m'}^1(\theta) = \begin{pmatrix} \frac{1}{2}(1 + \cos \theta) & \frac{1}{\sqrt{2}} \sin \theta & \frac{1}{2}(1 - \cos \theta) \\ -\frac{1}{\sqrt{2}} \sin \theta & \cos \theta & \frac{1}{\sqrt{2}} \sin \theta \\ \frac{1}{2}(1 - \cos \theta) & -\frac{1}{\sqrt{2}} \sin \theta & \frac{1}{2}(1 + \cos \theta) \end{pmatrix}$$

- There is an unequal contribution from the longitudinal term
- There is a dominance of $\sin^2 \theta$

PLOT OF $\nu B \rightarrow \mu f D$



SUMMARY AND CONCLUSION

- The ILC is an opportune place to make measurements on the fundamental properties of dark matter such as its mass and spin
- Our method of kinematics is model independent
- From it we can determine the magnitude of $\cos \theta_{LB}$ and partially reconstruct its sign from only visible momenta
- We can apply our method in other situations in which there is “missing” energy and momentum in the final state
 - Drell Yan process at the LHC
 - C. -W. Chiang, N. D. Christensen, G. -J. Ding and T. Han, Phys. Rev. D85, 015023 (2012)[arXiv:1107.5830 [hep-ph]].
- So far, our angular distributions produced from our kinematics seem to fit expectations

Thank You

FURTHER ANALYSIS OF ($\nu B \rightarrow \mu FD$)

■ Spin contributions

