

Sterile neutrino analysis of reactor neutrino oscillation

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Introduction

- In spite of the confirmation of three flavors of neutrinos, we do not concretely exclude the existence of sterile neutrinos.
- Singlets generally come out of many BSM.
(Remnants of GUT, seesaw mechanism, etc...)
- CMB, LSS and BBN prefer additional relativistic degree of freedom.
- Although sterile neutrinos do not interact with the electroweak gauge bosons, they can mix with three active neutrinos, leading to the oscillation between active and sterile neutrinos.

Evidence for sterile neutrino ?

- There exist **some anomalies** which cannot be explained within 3-flavor framework and might point towards the existence of additional ν 's.
- **LSND** : evidence for $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_e$ with $\Delta m^2 \sim 1 \text{ eV}^2$.
- **MiniBooNE** : a yet unexplained event excess in the low-E region of ν_e & anti- ν_e event spectra.
 - No significant excess is found at higher E_{ν} .
 - Interpreting MiniBooNE in terms of oscillations, there are parameter regions consistent with LSND result.

- **“Gallium anomaly”** : Gallium radioactive source experiments SAGE and GALLEX have obtained an event rate which is somewhat lower than expected → indication of ν_e disappearance, possibly due to $\nu_e \rightarrow \nu_s$
- **“Reactor anomaly”** : recent re-evaluation of the neutrino flux emitted by reactor led to increased fluxes compared with previous one. Inconsistent with measurements from short-baseline reactor experiments.

Four neutrino mixing

- **Our aim** : examine whether $\nu_s \leftrightarrow \nu_a$ is plausible, especially by interpreting the results released from Daya Bay and RENO.
- For simplicity, we introduce a sterile neutrino into the game.
- Since the setups of both expts. are optimized for Δm_{31}^2 , the search for $\nu_s \leftrightarrow \nu_a$ is possible **only if Δm_{41}^2 is not far different from Δm_{31}^2 .**
- Thus, our study does not cover the sterile neutrino introduced to reconcile the anomalies observed.

4-neutrino oscillations

- Unitary transformation from $(\nu_1, \nu_2, \nu_3, \nu_4)$ to $(\nu_e, \nu_\mu, \nu_\tau, \nu_s)$ parameterized by **6 angles and 3 Dirac-type Phases.**
- $$U_F = R_{34}(\theta_{34}) R_{24}(\theta_{24}, \delta_2) R_{14}(\theta_{14}) \cdot \\ R_{23}(\theta_{23}) R_{13}(\theta_{13}, \delta_1) R_{12}(\theta_{12}, \delta_3)$$

($R_{ij}(\theta_{ij})$: rotation of i-j block by θ_{ij})

- More explicitly except for phases,

$$\begin{aligned}
 \tilde{U}_F &= \begin{pmatrix} c_{14} & 0 & 0 & s_{14} \\ -s_{14}s_{24} & c_{24} & 0 & c_{14}s_{24} \\ -c_{24}s_{14}s_{34} & -s_{24}s_{34} & c_{34} & c_{14}c_{24}s_{34} \\ -c_{24}c_{34}s_{14} & -s_{24}c_{34} & -s_{34} & c_{14}c_{24}c_{34} \end{pmatrix} \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & 0 \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & 0 \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \\
 &= \begin{pmatrix} c_{14}U_{e1} & c_{14}U_{e2} & c_{14}U_{e3} & s_{14} \\ \cdots & \cdots & \cdots & c_{14}s_{24} \\ \cdots & \cdots & \cdots & c_{14}c_{24}s_{34} \\ \cdots & \cdots & \cdots & c_{14}c_{24}c_{34} \end{pmatrix},
 \end{aligned}$$

Survival probability of $\bar{\nu}_e$

$$\begin{aligned}
 P_{\text{th}}(\bar{\nu}_e \rightarrow \bar{\nu}_e) &= \left| \sum_{j=1}^4 |\tilde{U}_{ej}|^2 \exp i \frac{\Delta m_{j1}^2 L}{2E_\nu} \right|^2 \\
 &= 1 - \sum_{i < j} 4 |\tilde{U}_{ei}|^2 |\tilde{U}_{ej}|^2 \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E_\nu} \right)
 \end{aligned}$$

- Assuming $\Delta m_{41}^2 \sim \Delta m_{42}^2 \sim \Delta m_{43}^2 > \Delta m_{31}^2$
 $|\tilde{U}_{e3}|$ and $|\tilde{U}_{e4}| \ll 1$

$$\begin{aligned}
 P_{\text{th}}(\bar{\nu}_e \rightarrow \bar{\nu}_e) &= 1 - c_{14}^4 c_{13}^4 \sin^2 2\theta_{12} \sin^2 \left(1.27 \Delta m_{21}^2 \frac{L}{E} \right) \\
 &\quad - c_{14}^4 \sin^2 2\theta_{13} \sin^2 \left(1.27 \Delta m_{31}^2 \frac{L}{E} \right) \\
 &\quad - \sin^2 2\theta_{14} \sin^2 \left(1.27 \Delta m_{41}^2 \frac{L}{E} \right).
 \end{aligned}$$

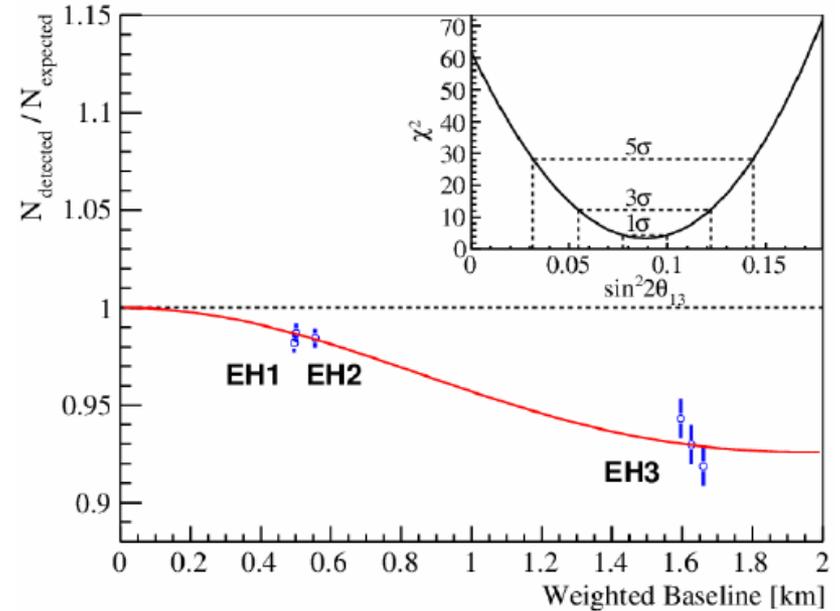
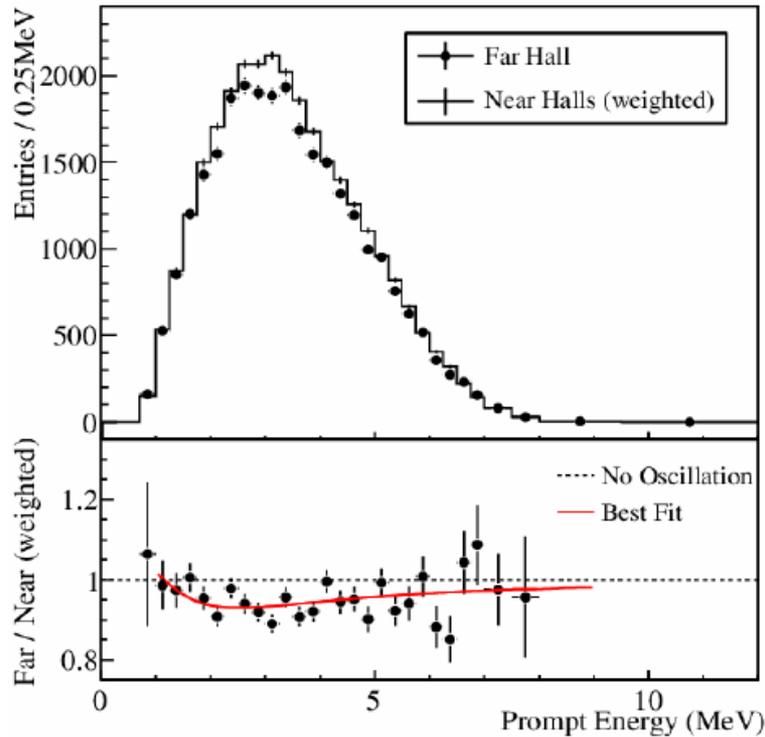
θ_{13} at Daya Bay

(Xianonan Li, NOW2012)

The Daya Bay Experiment

- 6 reactor cores, 17.4 GW_{th} total power
- Relative measurement
 - 2 near sites, 1 far site
- Multiple detector modules
- Good cosmic shielding
 - 250 m.w.e @ Daya Bay near
 - 860 m.w.e @ far site





$\bullet R = 0.944 \pm 0.007 \text{ (stat)} \pm 0.003 \text{ (syst)}$
 $\bullet \sin^2 \theta_{13} = 0.089 \pm 0.010 \text{ (stat)} \pm 0.005 \text{ (sys)}$

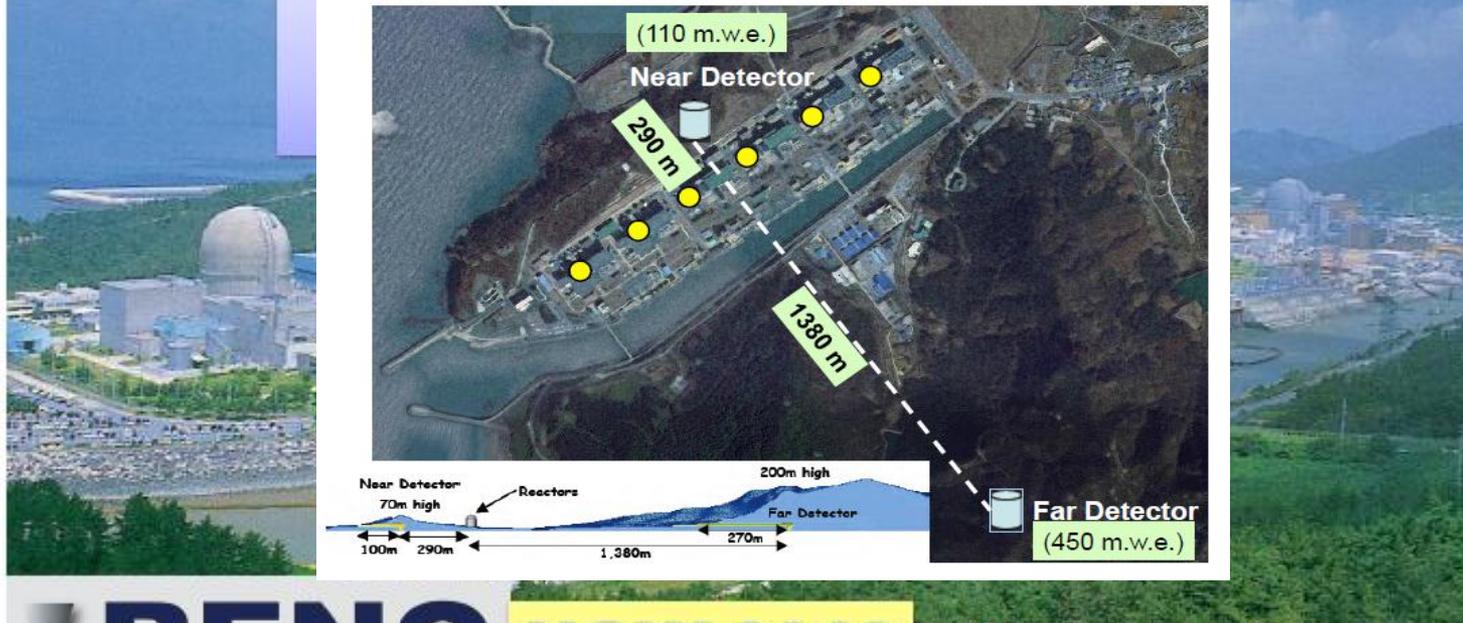
\bullet With 2.5x more statistics, an improved measurement to θ_{13}

θ_{13} at RENO

(Seon-Hee Seo, NOW2012)

Observation of Reactor Antineutrino Disappearance at RENO & Future Plan

RENO Experimental Setup



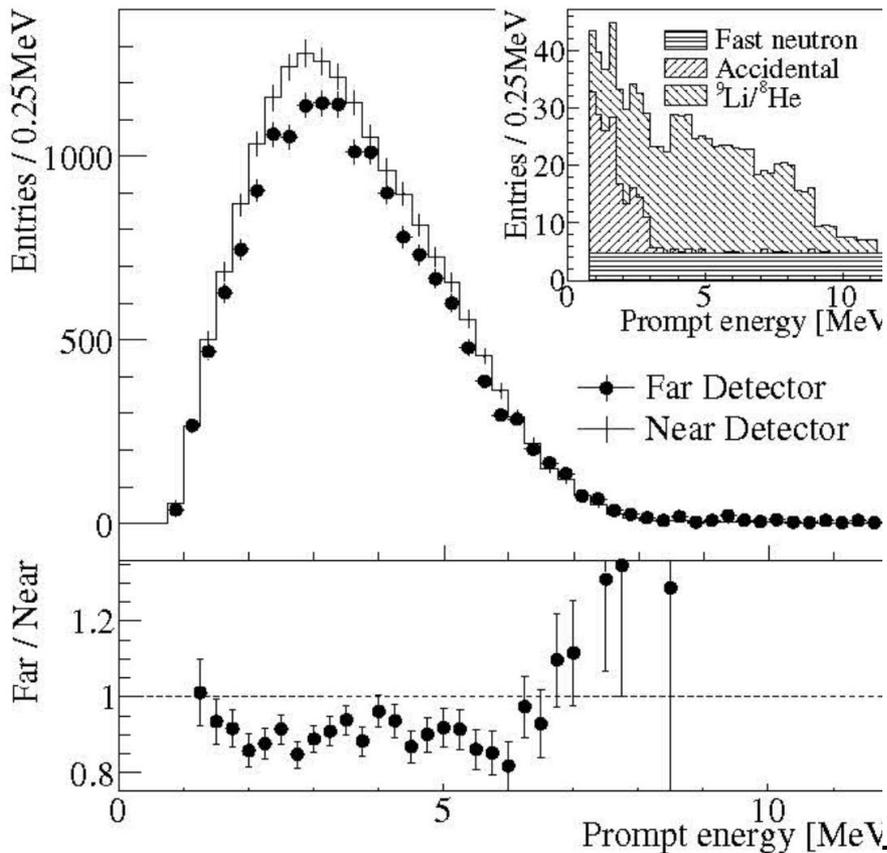
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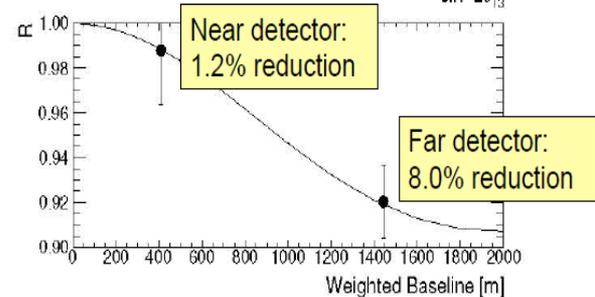
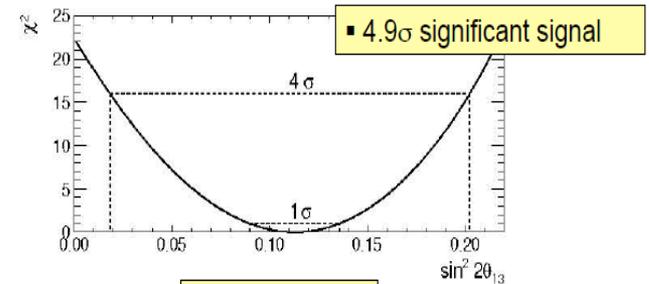
Reactor Antineutrino Disappearance

$$R = \frac{\Phi_{\text{observed}}^{\text{Far}}}{\Phi_{\text{expected}}^{\text{Far}}} = 0.920 \pm 0.009(\text{stat}) \pm 0.014(\text{syst})$$



Definitive Measurement of θ_{13}

$$\sin^2 2\theta_{13} = 0.113 \pm 0.013(\text{stat.}) \pm 0.019(\text{syst.})$$



Rate-only analysis

- Averaging the survival probability over E_ν :

$$\langle P \rangle = \frac{\int P_s(E) \sigma_{\text{tot}}(E) \phi(E) dE}{\int \sigma_{\text{tot}}(E) \phi(E) dE}$$

$\sigma_{\text{tot}}(E)$: total cross section of inverse beta decay

(Voget PRD60)

$\Phi(E)$: neutrino flux distribution from reactors (Huber PRC84)

$P_s(E)$: survival prob. at each detector

- Baselines

Daya Bay : EH1=494m, EH2=554m (near)
EH3=1628m (far)

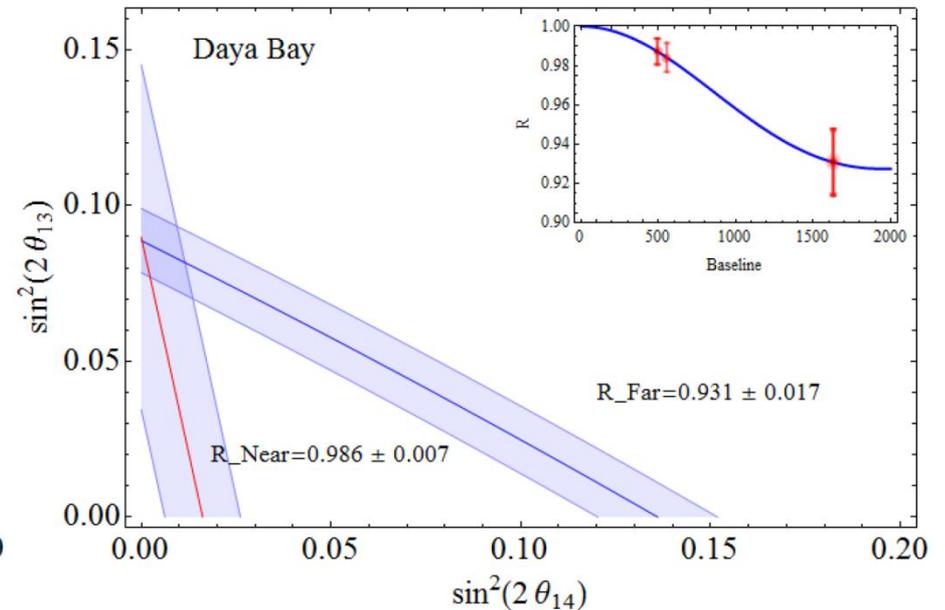
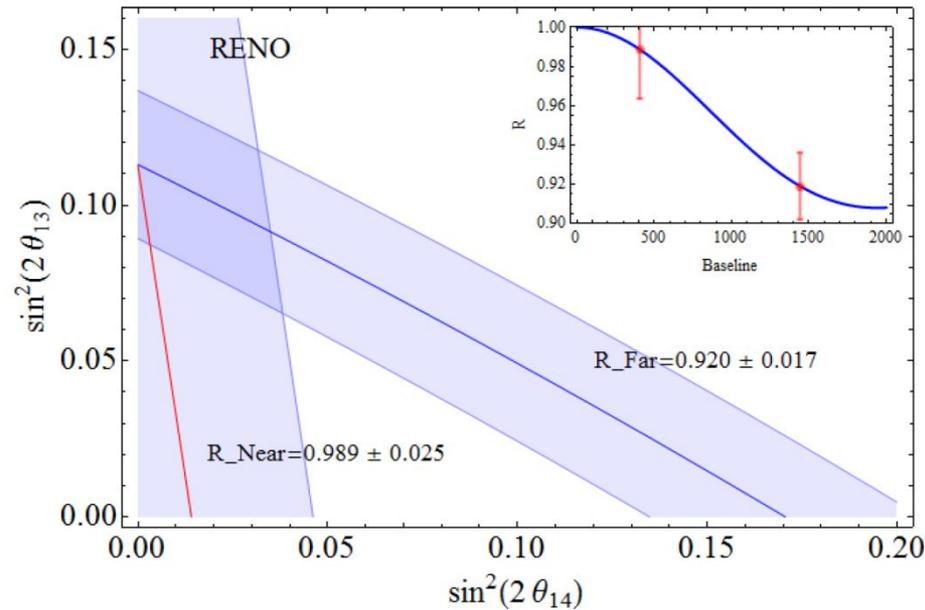
RENO : (near)={660,445,302,340,520,746}m
(far)={1560,1460,1400,1380,1410,1480}m

➔ flux-weighted averages :

(near)=407.3m (far)=1443m

- Taking into account detection efficiency at each detector.

Averaged $N_{\text{obs}}/N_{\text{exp}}$ in terms of 4-neutrino oscillation



- RENO :

$$\langle P_{\text{near}} \rangle = 0.989 \pm 0.025$$

$$\langle P_{\text{far}} \rangle = 0.920 \pm 0.017$$

- Daya Bay :

$$\langle P_{\text{near}} \rangle = 0.986 \pm 0.007$$

$$\langle P_{\text{far}} \rangle = 0.931 \pm 0.017$$

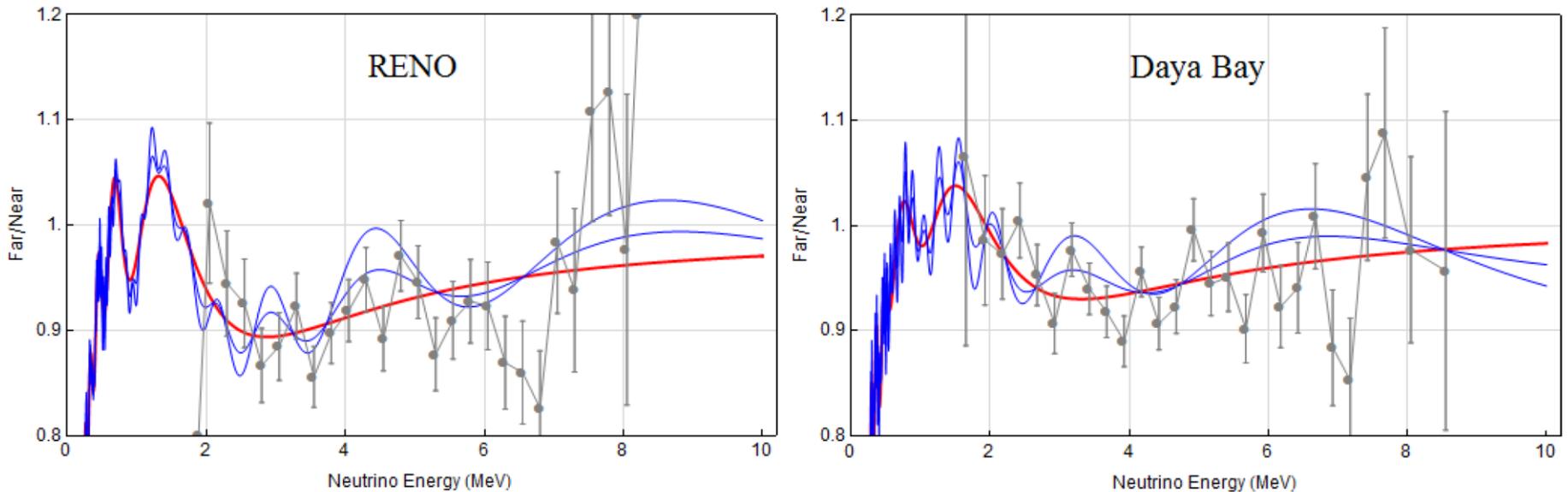
Spectral-shape analysis

- $R \equiv \frac{\# \text{ of events at FD}}{\# \text{ of events at ND}}$

→ data obtained per 0.25MeV bin as E varies from 1.8MeV to 8.5 MeV

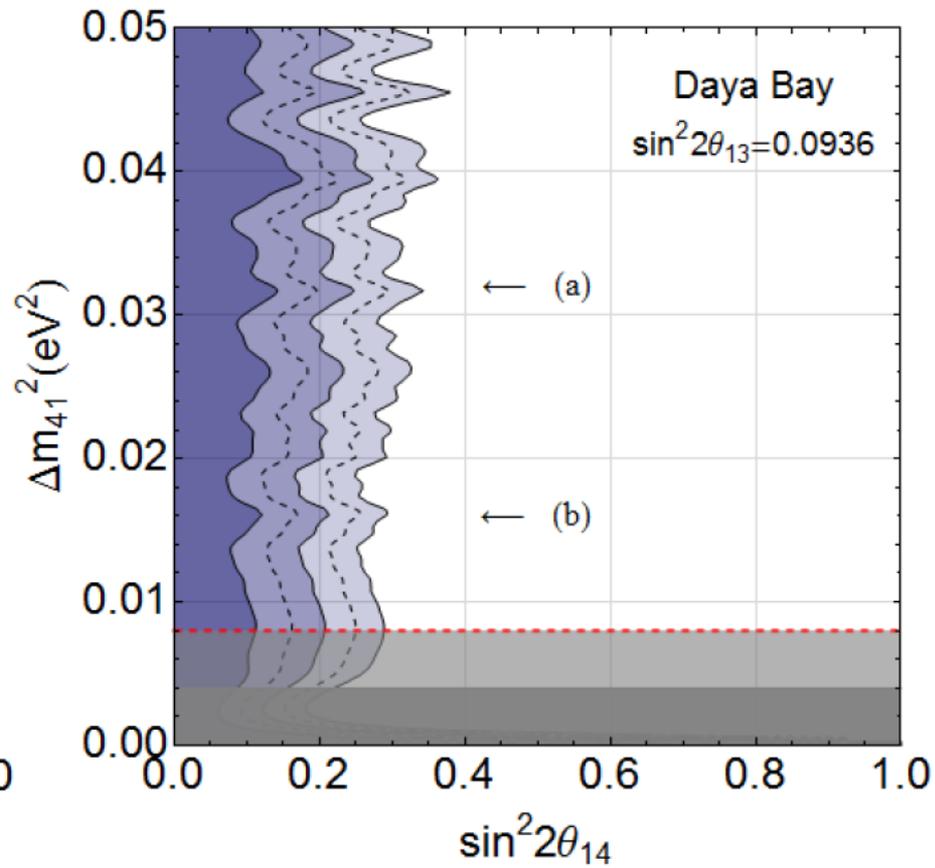
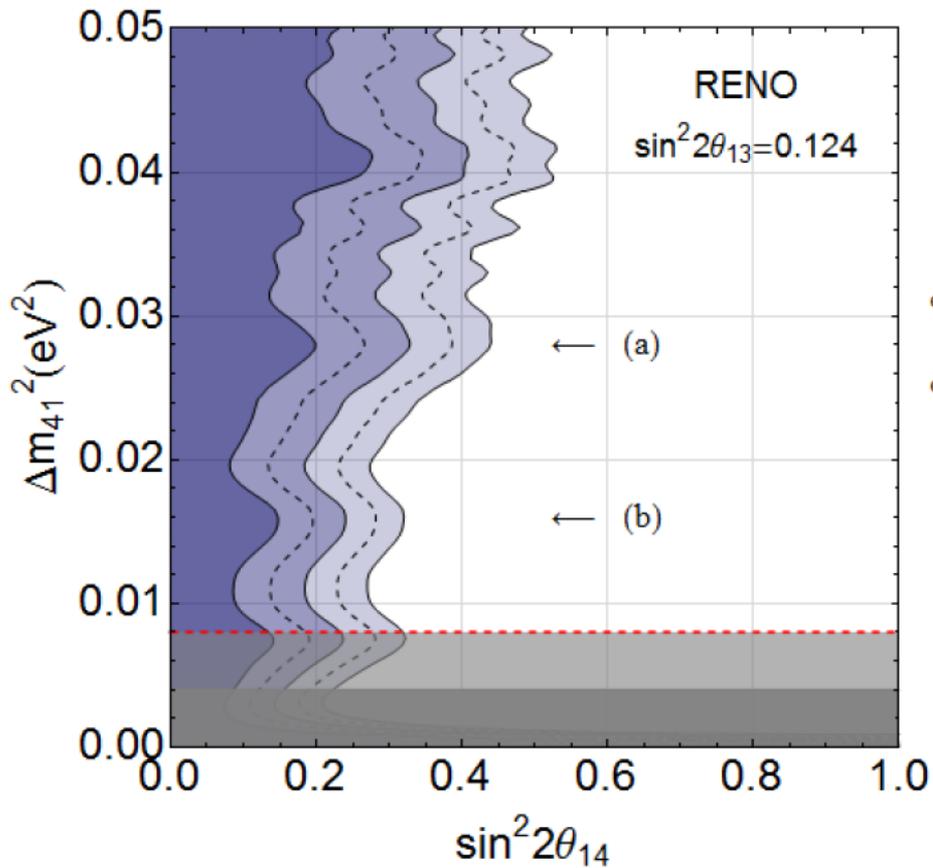
We plot the curve for $R = \frac{P_{th}(L_{far})}{P_{th}(L_{near})}$

Far-to-near ratio in terms of 4-neutrino oscillation

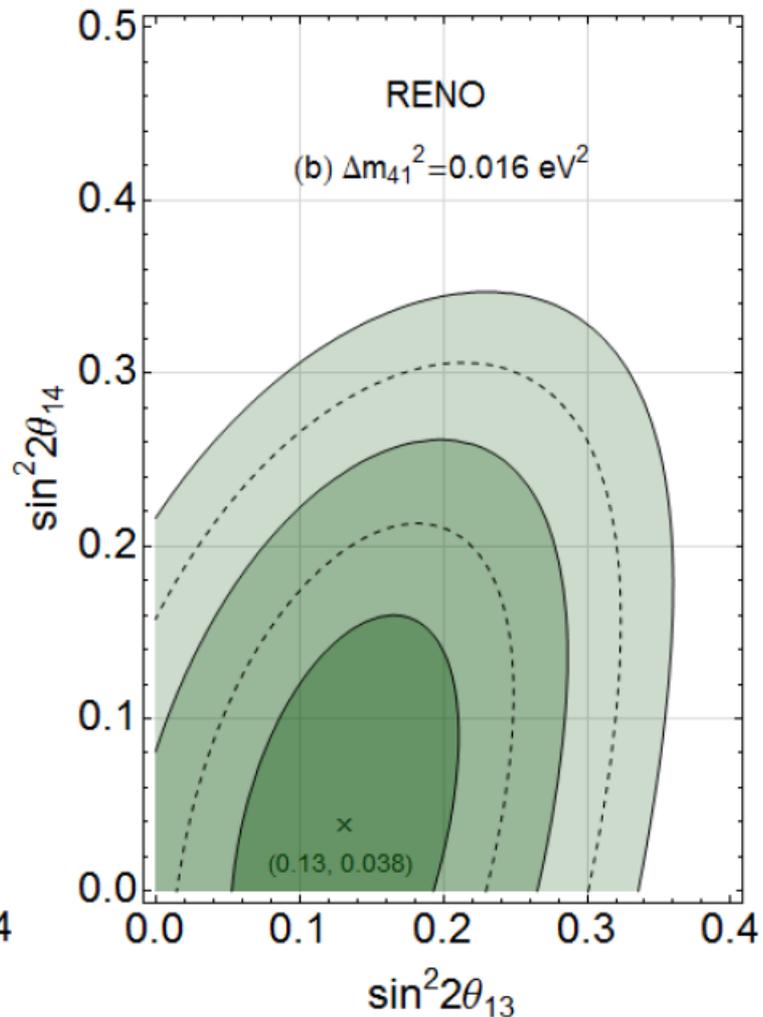
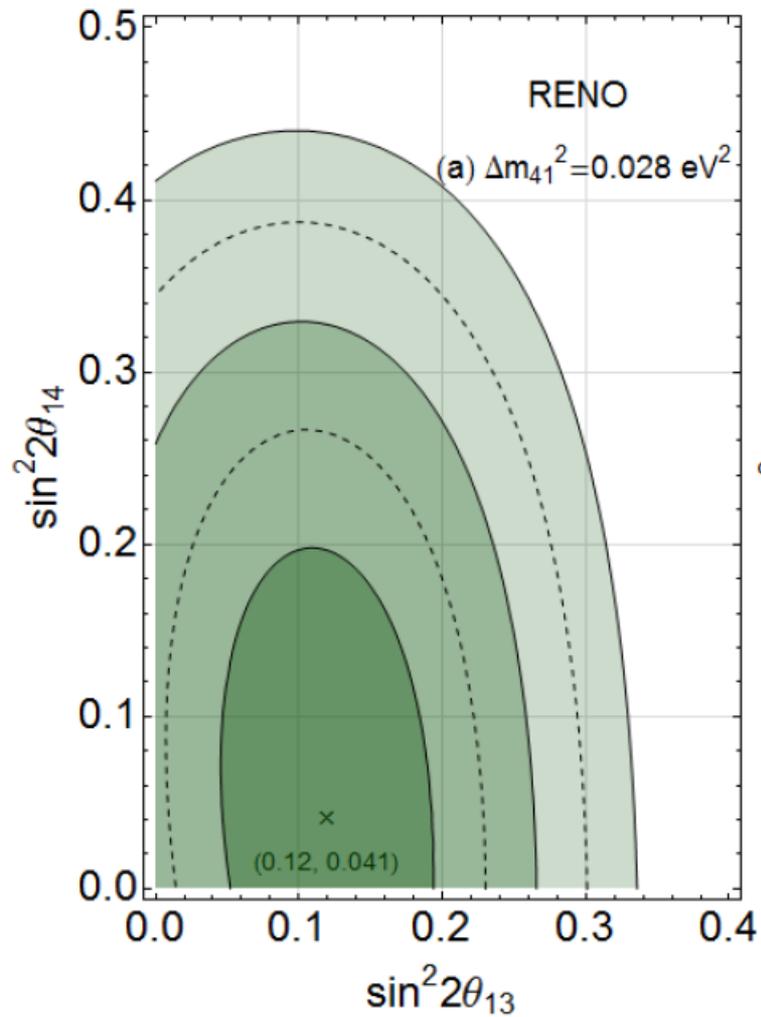


- **Red curves** : $\theta_{14} = 0$ & corresponding to the best fit points in the rate-only analysis
- **Blue curves** : $\sin^2 2\theta_{14} = 0.05, 0.1$
 $\Delta m_{41}^2 = 0.0185, 0.008$ eV²

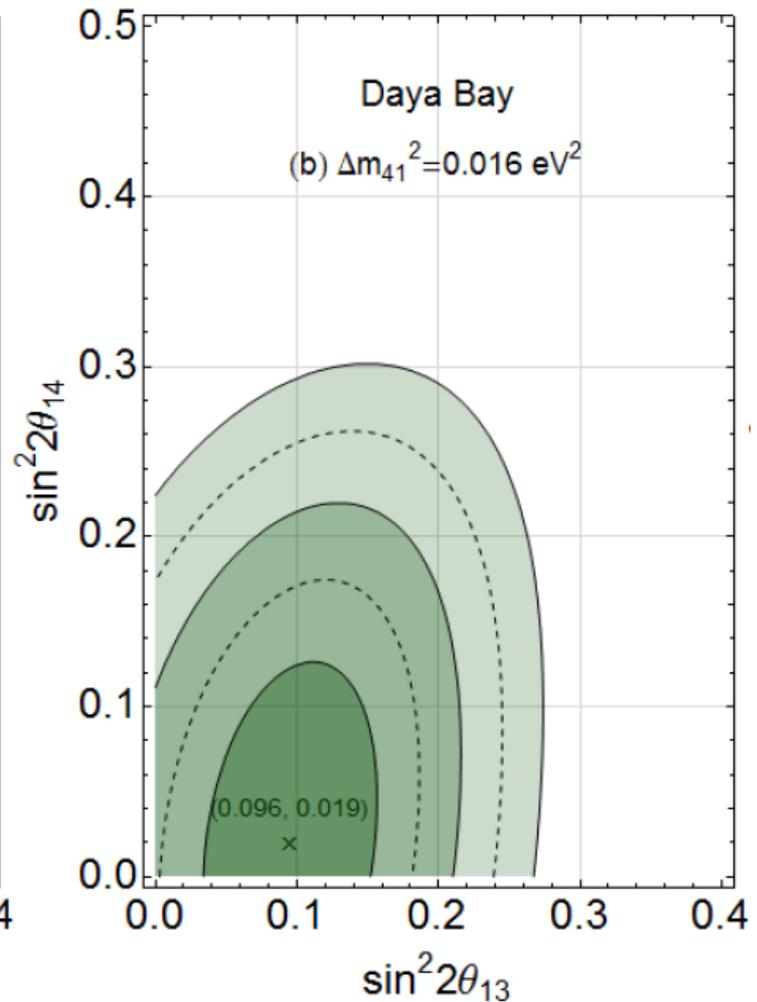
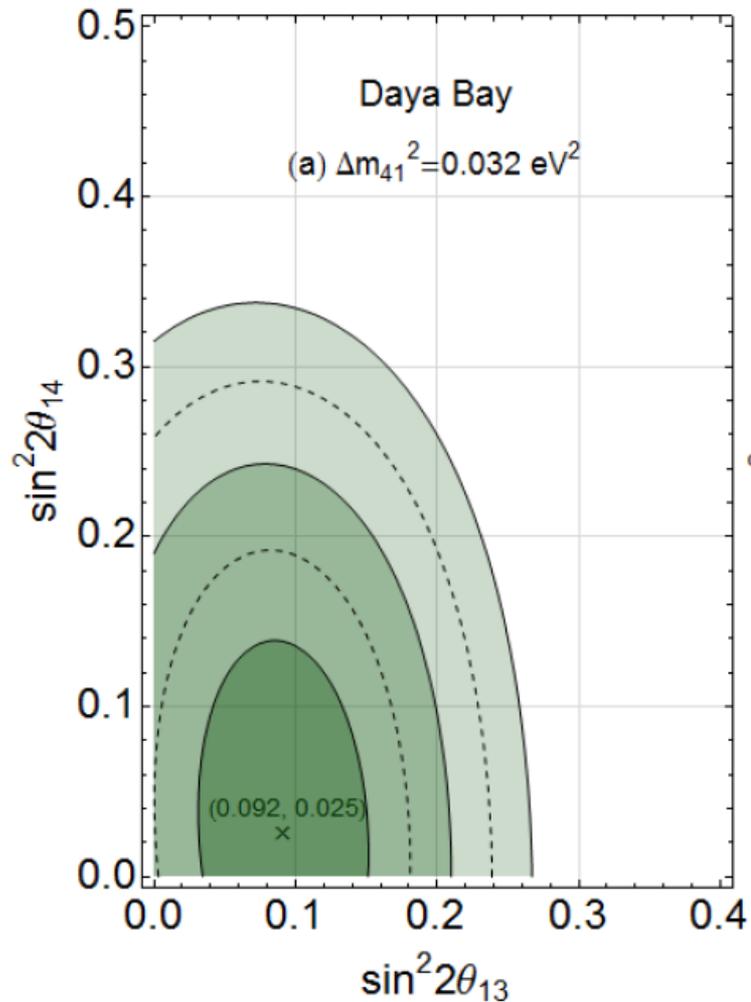
Exclusion curves (RENO)



Allowed regions in $(\theta_{13}, \theta_{14})$ (RENO)



Allowed regions in $(\theta_{13}, \theta_{14})$ (Daya Bay)



Summary

- We have examined whether oscillation between a sterile neutrino and active neutrinos is plausible at Daya Bay & RENO experiments.
- Under a few assumptions, our analysis shows that $\nu_s \leftrightarrow \nu_a$ is plausible and gives even better fit compared with 3-flavor framework analysis.

