

Anomalous Gauge Couplings using Effective Theories

Celine Degrande
University of Illinois at Urbana-Champaign
Snowmass-electroweak

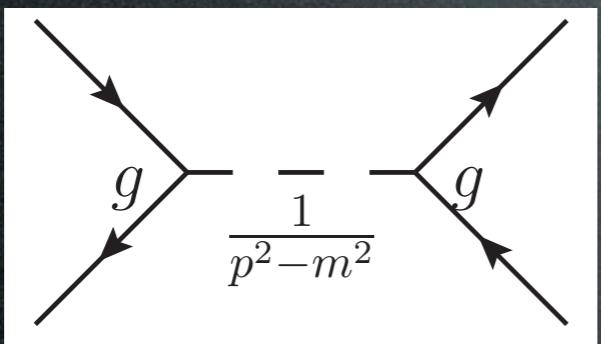
Plan

- Introduction to EFT
- EFT for $pp(ee) \rightarrow WW$
- The constraints
- The effects
- EFT for $pp \rightarrow WWW$, $pp \rightarrow WWjj$
- Concluding remarks

Model Independent searches

Resonances

Assumption : One particle exchanged

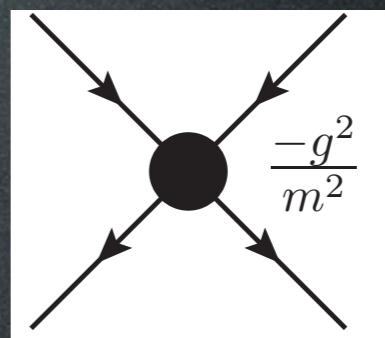


Example : Z'

Effects : Peak in the invariant mass distribution

Effective theory

Assumption : The new physics is heavy



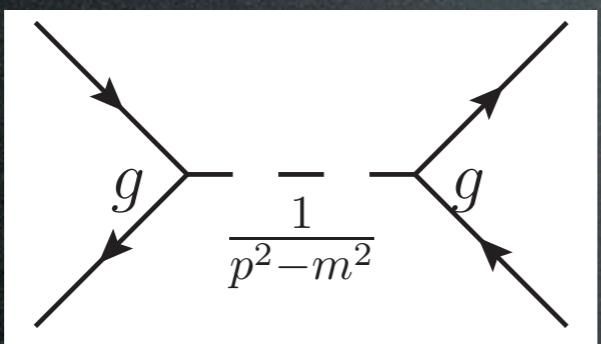
Example : Fermi th.

Effects :
Normalisation : m^2/Λ^2
Shape : s/Λ^2

Model Independent searches

Resonances

Assumption : One particle exchanged

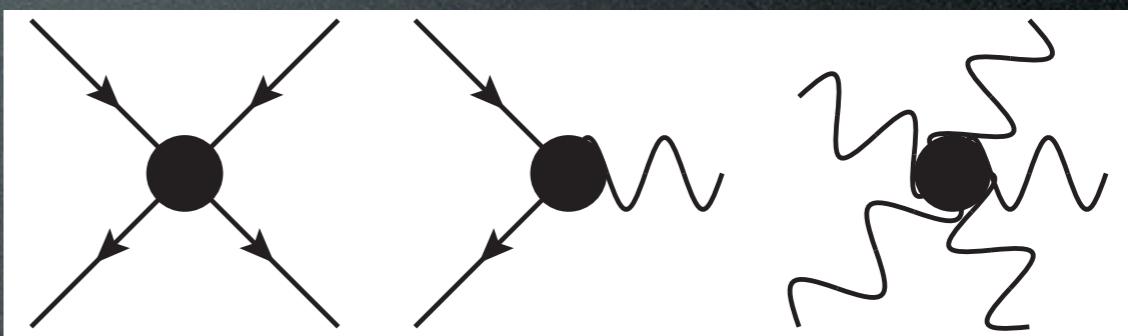


Example : Z'

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Effective theory

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Example : Fermi th.

Effects :
Normalisation : m^2/Λ^2
Shape : s/Λ^2

Effective Lagrangian

- From the SM fields only
- Invariant under the SM symmetries
- Dimension d of the new operator is >4
- New operators are suppressed by $1/\Lambda^{d-4}$
- Keep only the first order (lowest dimension)

$$\mathcal{O}_{WWW} = \text{Tr}[W_{\mu\nu} W^{\nu\rho} W_{\rho}^{\mu}]$$

$$\mathcal{L} = \mathcal{L}^{SM} + \sum \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

Th. error

Effective field theories

$$\mathcal{L} = \mathcal{L}^{SM} + \sum \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

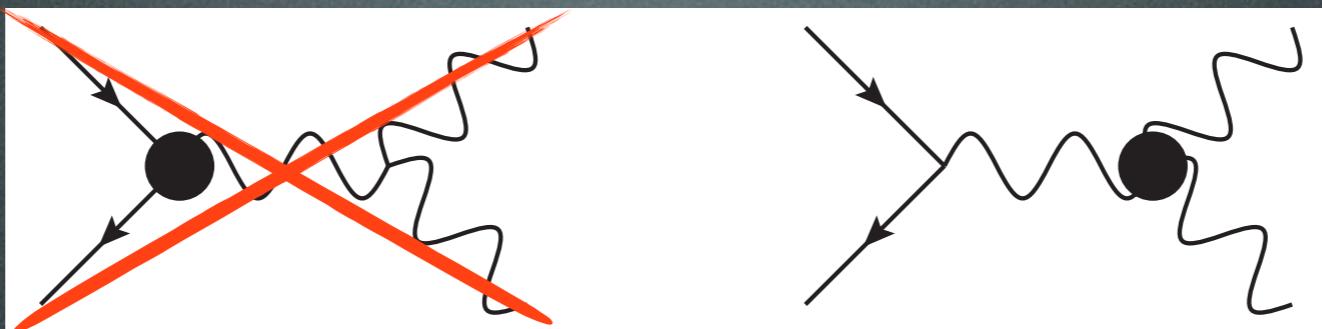
- SM symmetries (B&L) : 59 dimension-six operators (one flavor)
- Only few operators/process and different effects
- More predictive than anomalous couplings
- Unitarity is satisfied (no form factors)
- More than one vertex/operator (high multiplicities)
- Loop computation

C.D. et al, arXiv:1205.4231

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WW(WZ/WA) production



CP even operators

$$\mathcal{O}_{WWW} = \text{Tr}[W_{\mu\nu} W^{\nu\rho} W_\rho^\mu]$$

$$\mathcal{O}_W = (D_\mu \Phi)^\dagger W^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_B = (D_\mu \Phi)^\dagger B^{\mu\nu} (D_\nu \Phi)$$

TGC's and weak boson masses are affected by different operators at the tree-level in this basis

CP odd operators

$$\mathcal{O}_{\tilde{W}WW} = \text{Tr}[\tilde{W}_{\mu\nu} W^{\nu\rho} W_\rho^\mu]$$

$$\mathcal{O}_{\tilde{W}} = (D_\mu \Phi)^\dagger \tilde{W}^{\mu\nu} (D_\nu \Phi)$$

Only 5 operators!

WW production

$$\begin{aligned}\mathcal{L} = ig_{WWV} \left(& g_1^V (W_{\mu\nu}^+ W^{-\mu} - W^{+\mu} W_{\mu\nu}^-) V^\nu + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda_V}{M_W^2} W_\mu^{\nu+} W_\nu^{-\rho} V_\rho^\mu \right. \\ & + ig_4^V W_\mu^+ W_\nu^- (\partial^\mu V^\nu + \partial^\nu V^\mu) \\ & \left. - ig_5^V \epsilon^{\mu\nu\rho\sigma} (W_\mu^+ \partial_\rho W_\nu^- - \partial_\rho W_\mu^+ W_\nu^-) V_\sigma + \tilde{\kappa}_V W_\mu^+ W_\nu^- \tilde{V}^{\mu\nu} + \frac{\tilde{\lambda}_V}{m_W^2} W_\mu^{\nu+} W_\nu^{-\rho} \tilde{V}_\rho^\mu \right)\end{aligned}$$

$$g_{WW\gamma} = -e \quad g_{WWZ} = -e \cot \theta_W$$

EM gauge invariance implies : $g_1^\gamma = 1$ $g_4^\gamma = g_5^\gamma = 0$

11(5+6) parameters

WW production

$$\mathcal{L} = ig_{WWV} \left(g_1^V (W_{\mu\nu}^+ W^{-\mu} - W^{+\mu} W_{\mu\nu}^-) V^\nu + \kappa_V W_\mu^+ W_\nu^- V^\mu + \frac{\lambda_V}{M_W^2} W_\mu^{\nu+} W_\nu^{-\rho} V_\rho^\mu \right. \\ \left. + ig_4^V W_\mu^+ W_\nu^- (\partial^\mu V^\nu + \partial^\nu V^\mu) \right. \\ \left. - ig_5^V \epsilon^{\mu\nu\rho\sigma} (W_\mu^+ \partial_\rho W_\sigma^- - \partial_\rho W_\mu^+ W_\sigma^-) V_\sigma + \tilde{\kappa}_V W_\mu^+ W_\nu^- \tilde{V}^{\mu\nu} + \frac{\tilde{\lambda}_V}{m_W^2} W_\mu^{\nu+} W_\nu^{-\rho} \tilde{V}_\rho^\mu \right)$$

Why not adding derivatives

$$g_{WW\gamma} = -e \quad g_{WWZ} = -e \cot \theta_W$$

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CP even Operators

$$\begin{aligned}\mathcal{O}_{WWW} &= \text{Tr}[W_{\mu\nu} W^{\nu\rho} W_\rho^\mu] \\ \mathcal{O}_W &= (D_\mu \Phi)^\dagger W^{\mu\nu} (D_\nu \Phi) \\ \mathcal{O}_B &= (D_\mu \Phi)^\dagger B^{\mu\nu} (D_\nu \Phi)\end{aligned}$$

CP odd operators

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$$\begin{aligned}g_1^Z &= 1 + c_W \frac{m_Z^2}{2\Lambda^2} \\ \kappa_\gamma &= 1 + (c_W + c_B) \frac{m_W^2}{2\Lambda^2} \\ \kappa_Z &= 1 + (c_W - c_B \tan^2 \theta_W) \frac{m_W^2}{2\Lambda^2} \\ \lambda_\gamma &= \lambda_Z = c_{WWW} \frac{3g^2 m_W^2}{2\Lambda^2} \\ g_4^V &= g_5^V = 0 \\ \tilde{\kappa}_\gamma &= c_{\tilde{W}} \frac{m_W^2}{2\Lambda^2} \\ \tilde{\kappa}_Z &= -c_{\tilde{W}} \tan^2 \theta_W \frac{m_W^2}{2\Lambda^2} \\ \tilde{\lambda}_\gamma &= \tilde{\lambda}_Z = c_{\tilde{W}WW} \frac{3g^2 m_W^2}{2\Lambda^2}\end{aligned}$$

constants

WW Production

$$\mathcal{L} = ig_{WWV} \left(g_1^V (W_{\mu\nu}^+ W^{-\mu} - W^{+\mu} W_{\mu\nu}^-) V^\nu + \kappa_V W_\mu^+ W_\nu^- V^{\mu\nu} + \frac{\lambda_V}{M_W^2} W_\mu^{\nu+} W_\nu^{-\rho} V_\rho^\mu + ig_4^V W_\mu^+ W_\nu^- (\partial^\mu V^\nu + \partial^\nu V^\mu) \right. \\ \left. - ig_5^V \epsilon^{\mu\nu\rho\sigma} (W_\mu^+ \partial_\rho W_\nu^- - \partial_\rho W_\mu^+ W_\nu^-) V_\sigma + \tilde{\kappa}_V W_\mu^+ W_\nu^- \tilde{V}^{\mu\nu} + \frac{\tilde{\lambda}_V}{m_W^2} W_\mu^{\nu+} W_\nu^{-\rho} \tilde{V}_\rho^\mu \right)$$

CP even Operators

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CP odd operators

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$$\mathcal{O}_{\tilde{W}} = (D_\mu \Phi)^\dagger \tilde{W}^{\mu\nu} (D_\nu \Phi)$$

$\Delta X = X - 1$

$g_1^Z = 1 + c_W \frac{m_Z^2}{2\Lambda^2}$

$\Delta g_1^Z = \Delta \kappa_Z + \tan^2 \theta_W \Delta \kappa_\gamma$

$\kappa_Z = 1 + (c_W \frac{m_Z^2}{2\Lambda^2}) \frac{3g^2 m_W^2}{2\Lambda^2}$

$\lambda_\gamma = \lambda_Z = c_{WWW} \frac{3g^2 m_W^2}{2\Lambda^2}$

$g_4^V = g_5^V = 0$

$\tilde{\kappa}_\gamma = \frac{m_W^2}{2\Lambda^2}$

$\tilde{\kappa}_Z = \tilde{\kappa}_Z + \tan^2 \theta_W \tilde{\kappa}_\gamma$

$\tilde{\lambda}_\gamma = \tilde{\lambda}_Z = c_{\tilde{W}WW} \frac{3g^2 m_W^2}{2\Lambda^2}$

constants

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PDG constraints

$$g_1^Z = 0.984^{+0.022}_{-0.019}$$

$$\kappa_\gamma = 0.979^{+0.044}_{-0.045}$$

$$\lambda_\gamma = 0.028^{+0.020}_{-0.021}$$

$$\tilde{\kappa}_Z = 0.12^{+0.06}_{-0.04}$$

$$\tilde{\lambda}_Z = 0.09 \pm 0.07$$

At 68% C.L.

$$c_{WWW}/\Lambda^2 \in [-11.9, 1.94] \text{TeV}^{-2}$$

$$c_W/\Lambda^2 \in [8.42, 1.44] \text{TeV}^{-2}$$

$$c_B/\Lambda^2 \in [-7.9, 14.9] \text{TeV}^{-2}$$

$$c_{\tilde{W}WW}/\Lambda^2 \in [-185.3, -82.4] \text{TeV}^{-2}$$

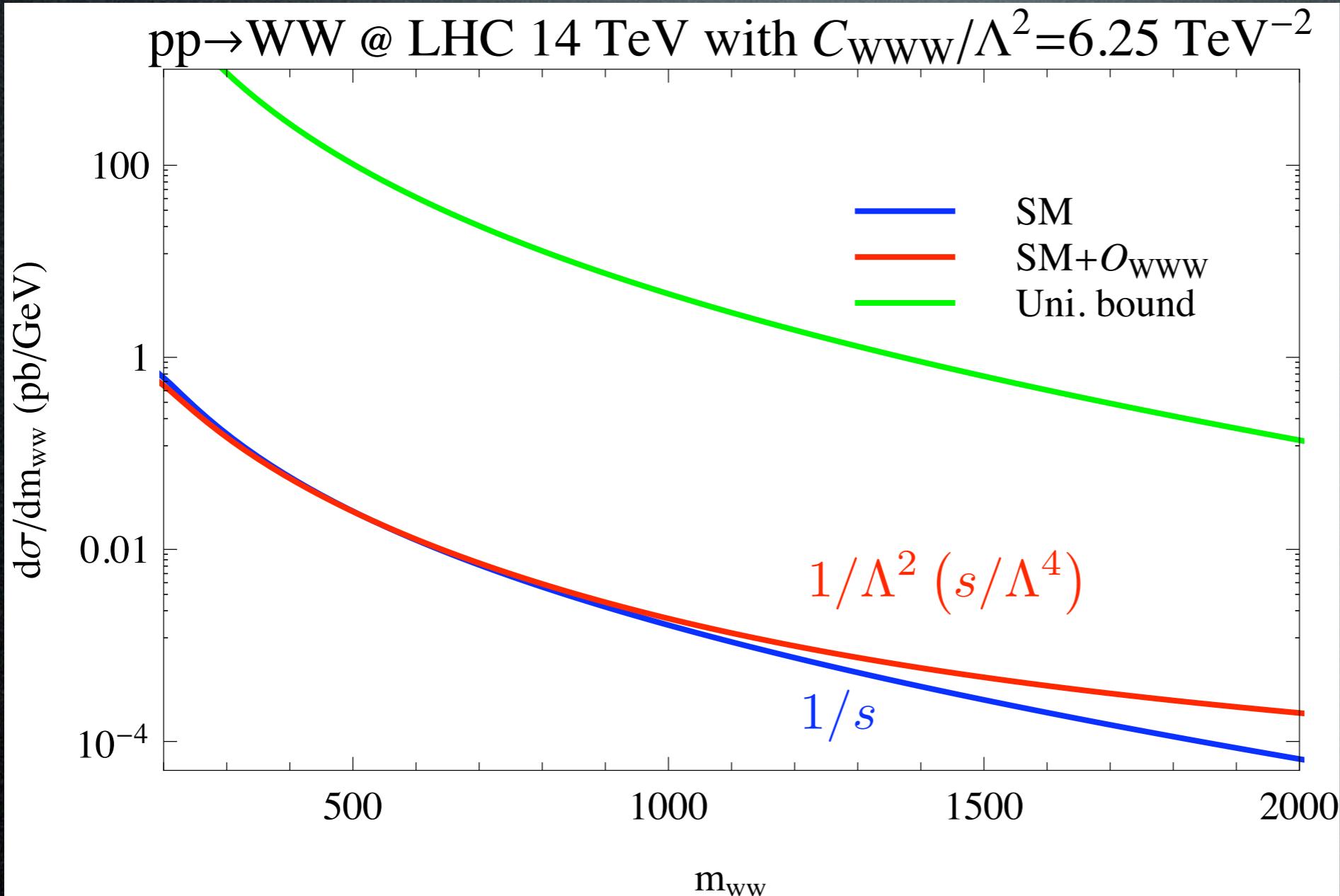
$$c_{\tilde{W}}/\Lambda^2 \in [-39.3, -4.9] \text{TeV}^{-2}$$

- Only LEP combination
- Tevatron measurements are more precise but use form factors/other relations

Plan

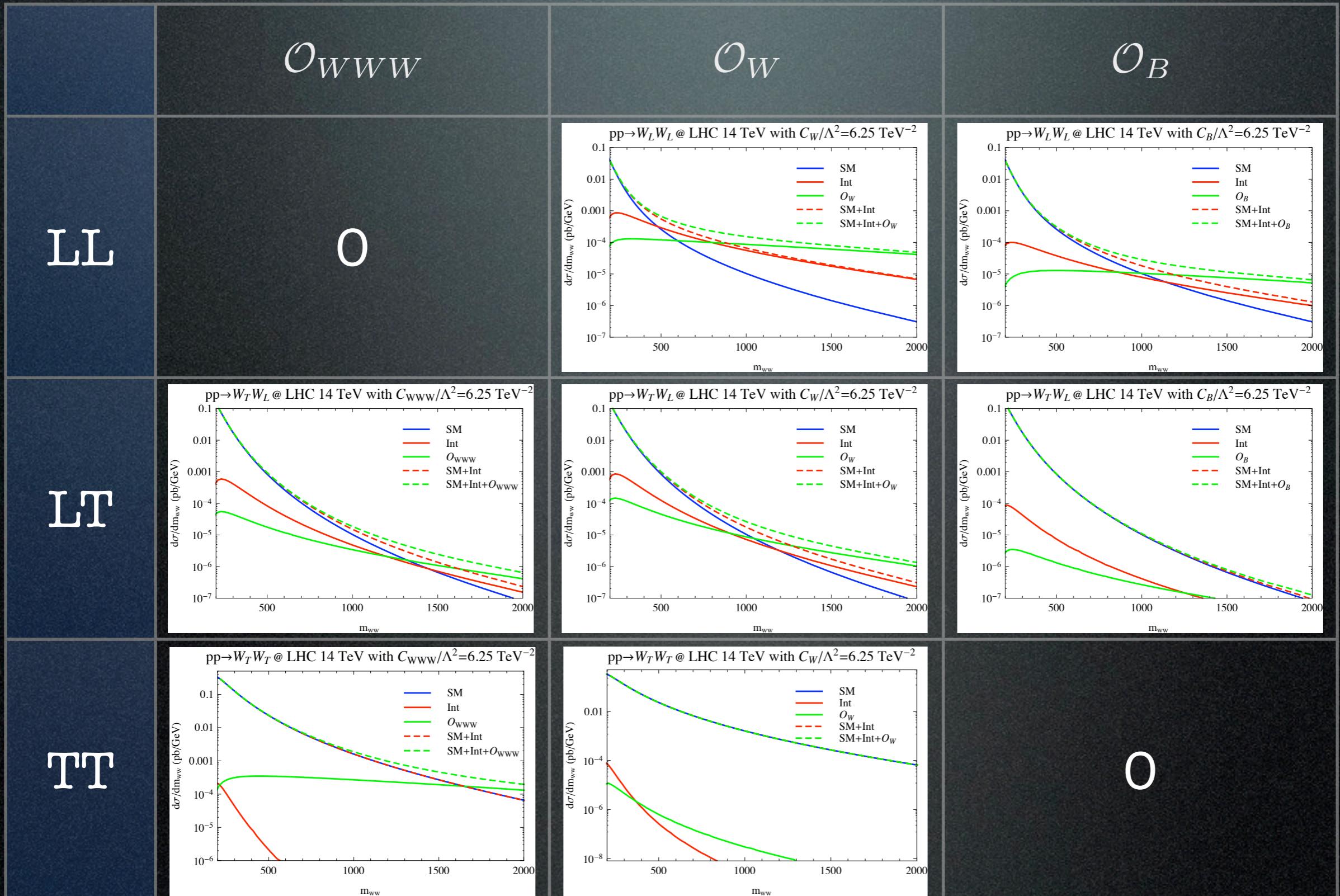
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Unitarity bound

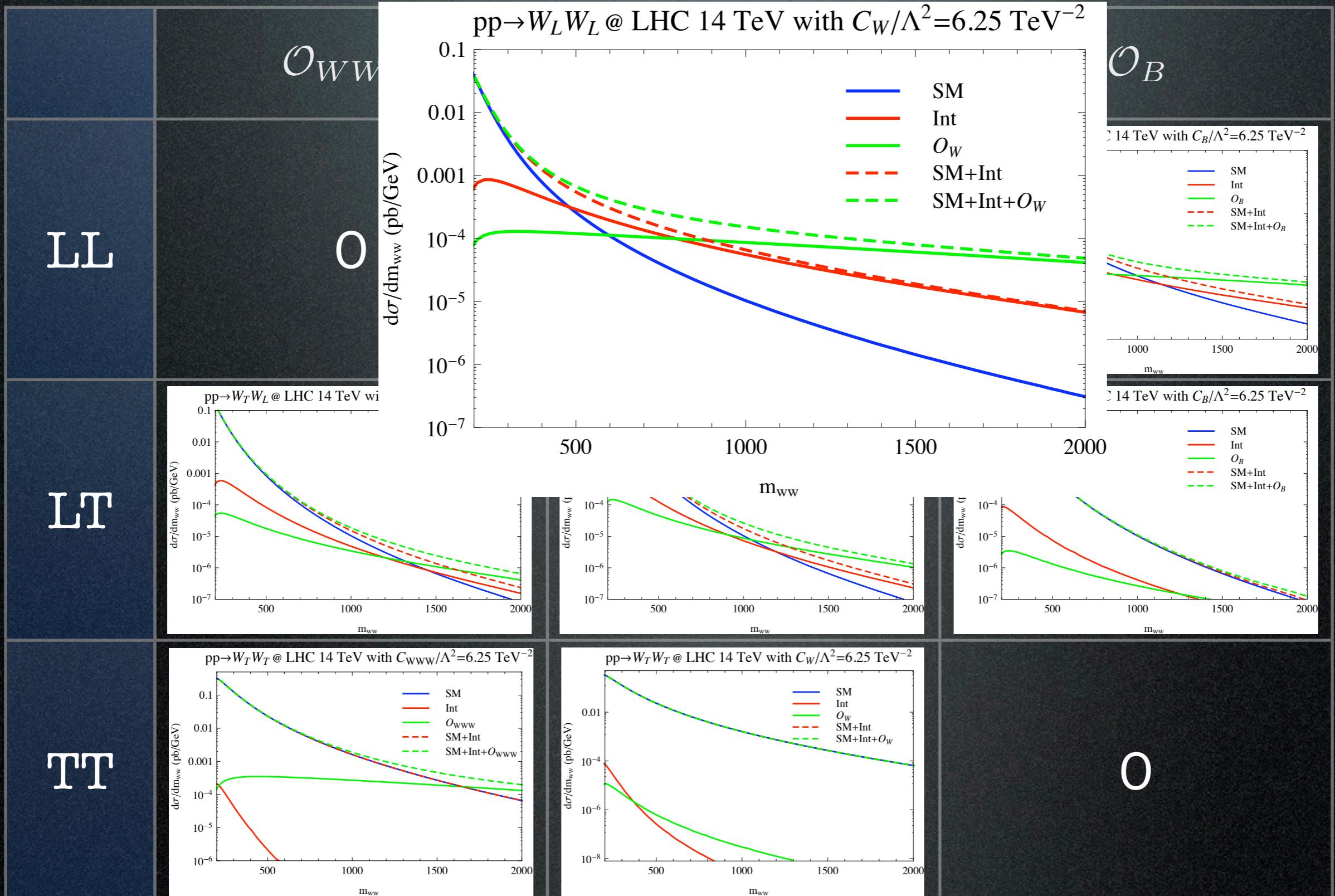


Form factors are not needed!

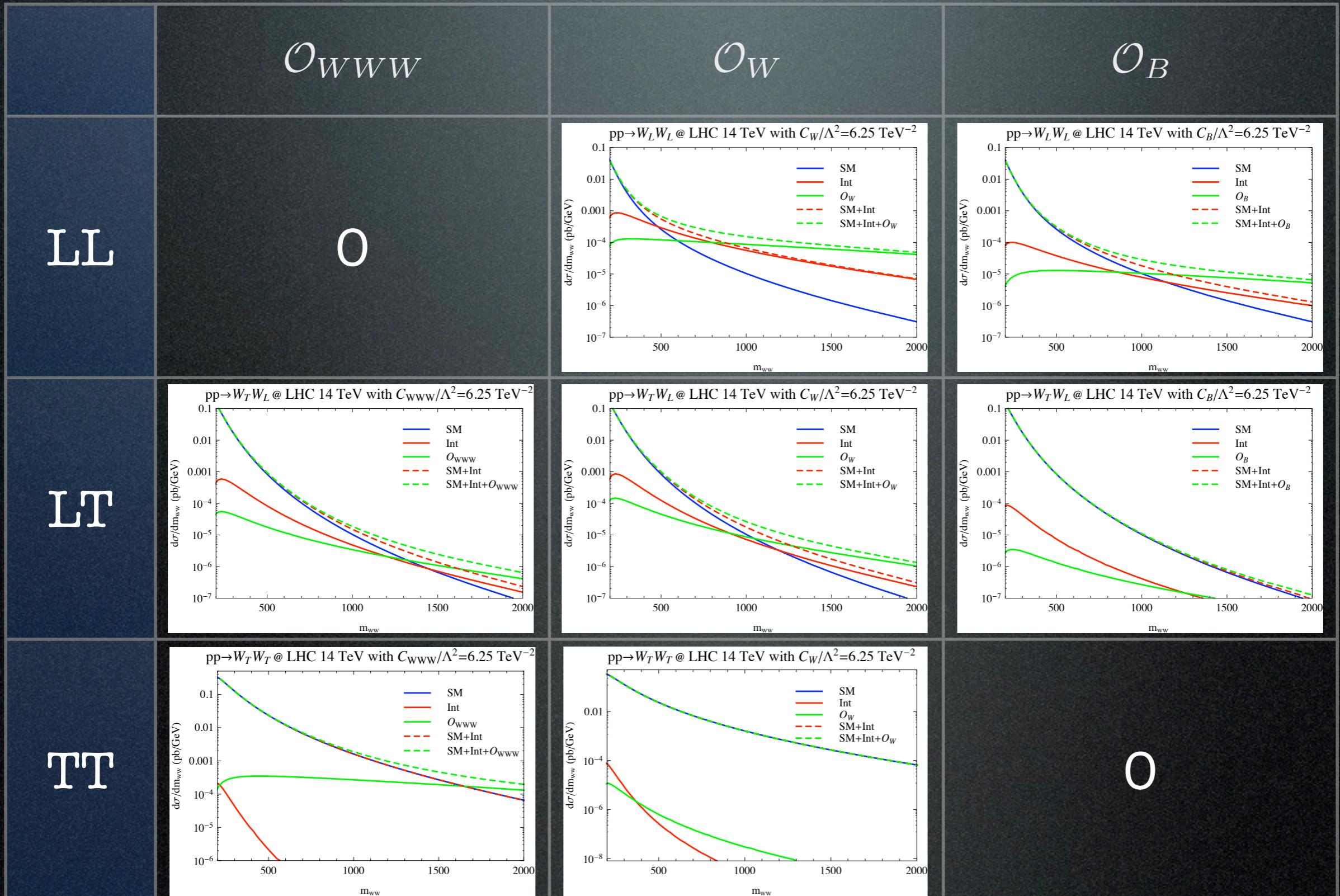
Invariant mass and polarisations



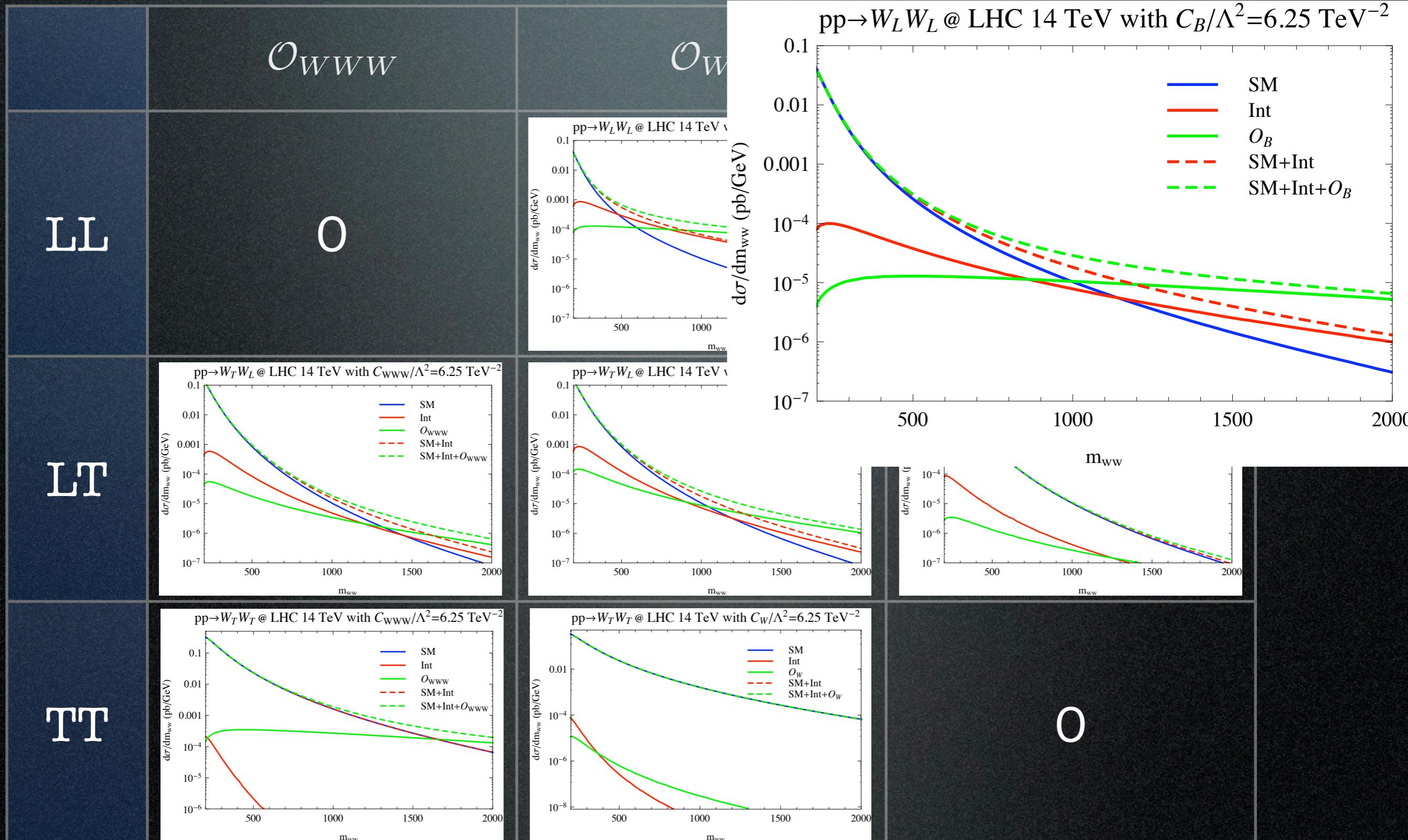
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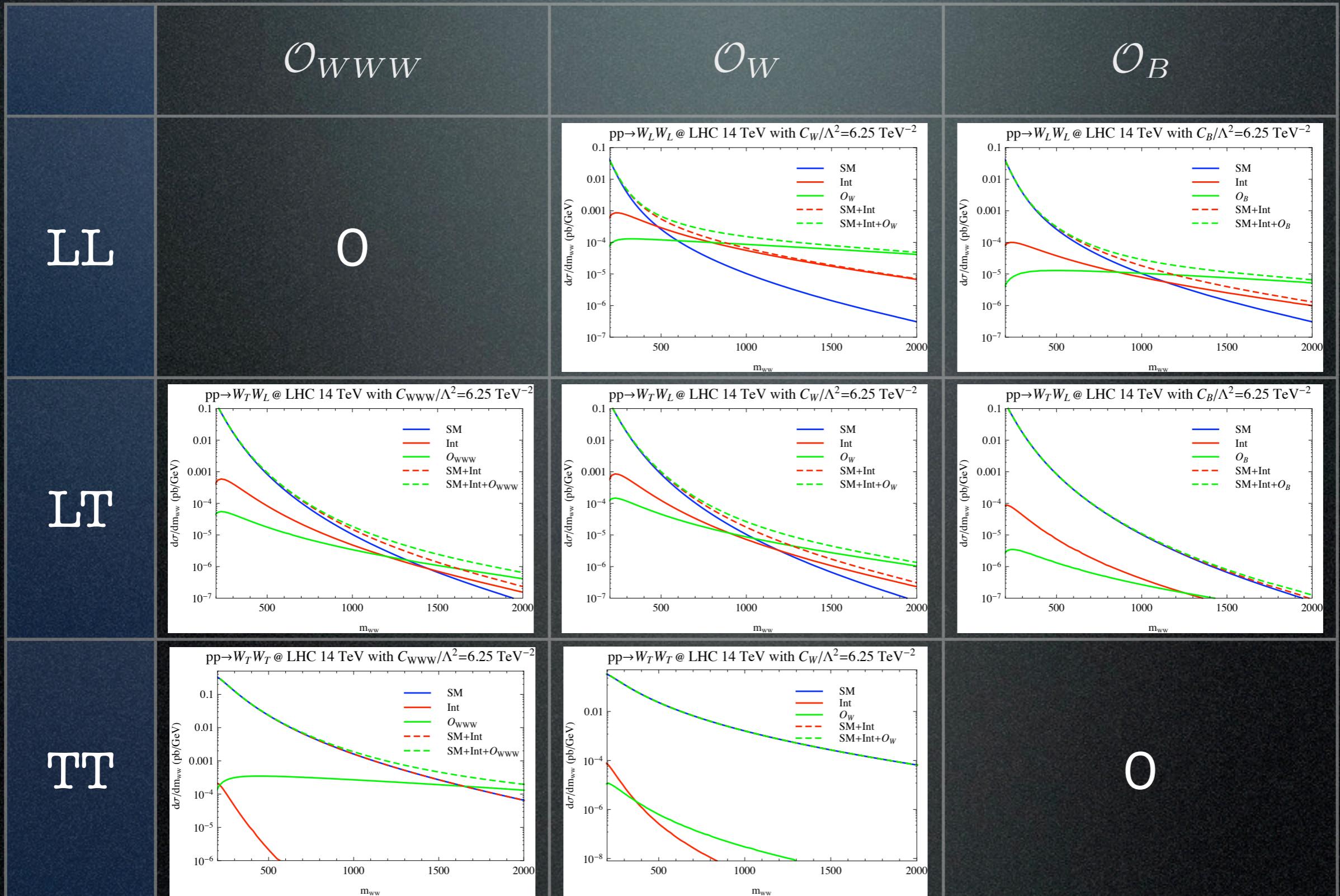
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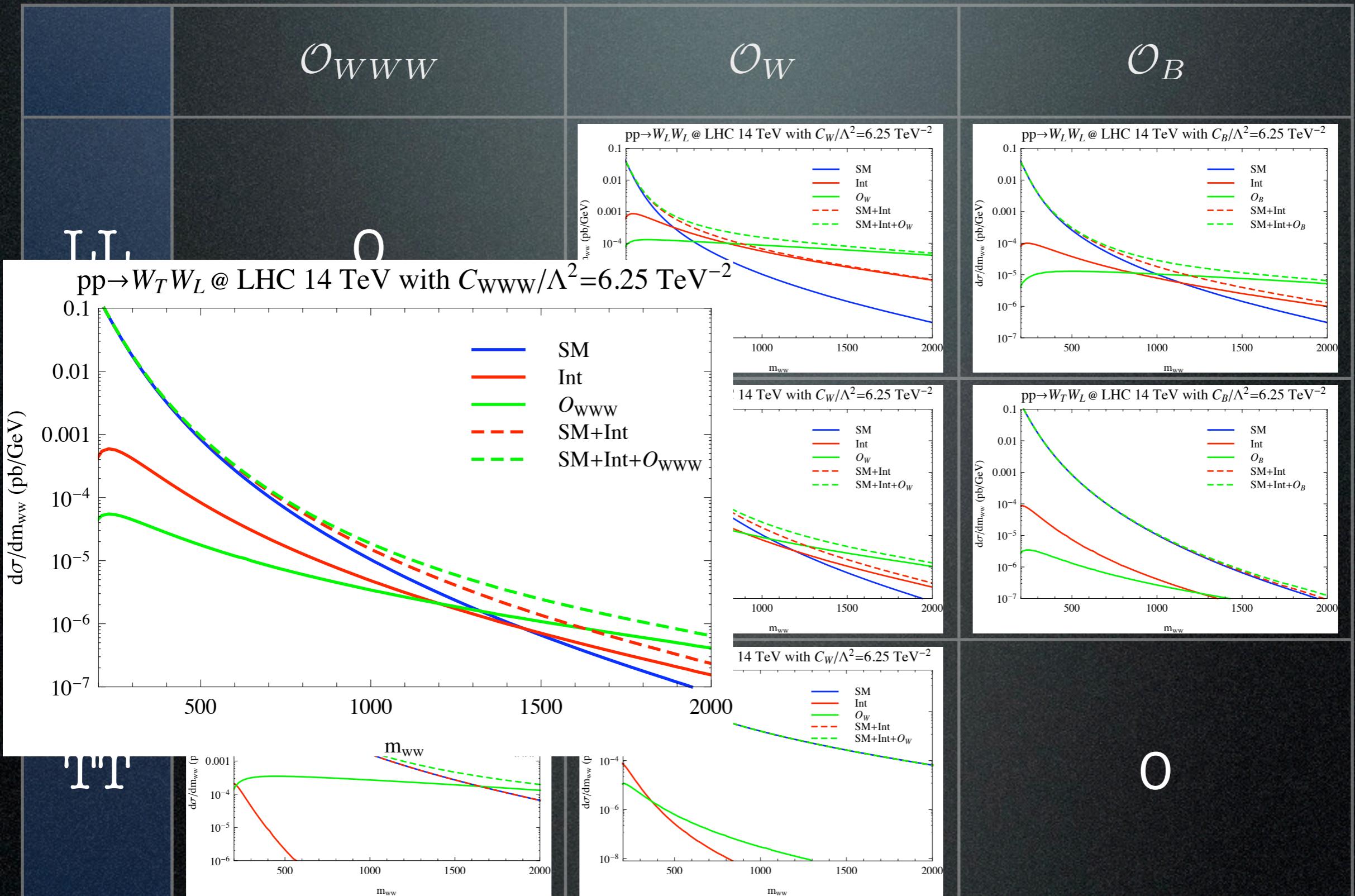
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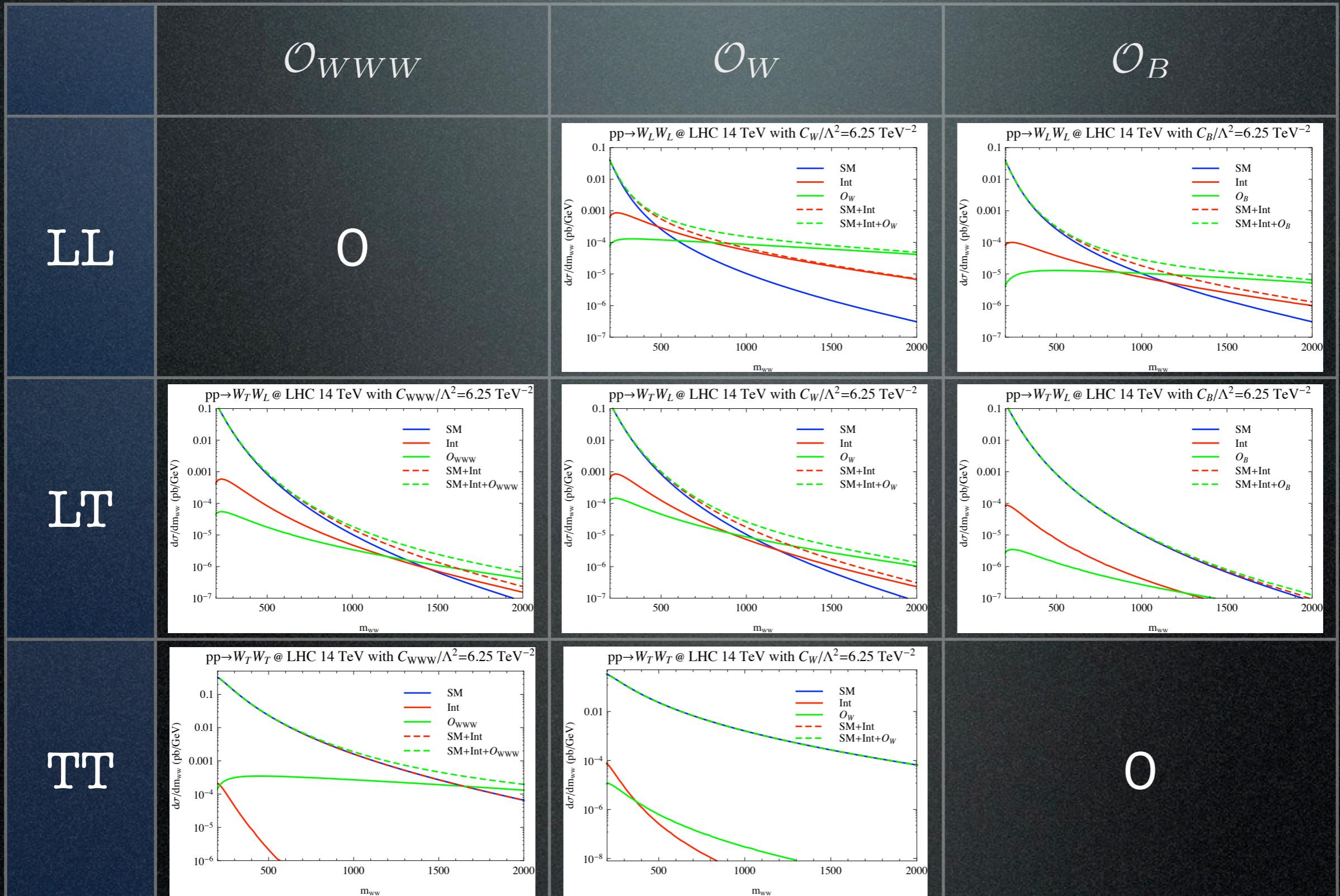
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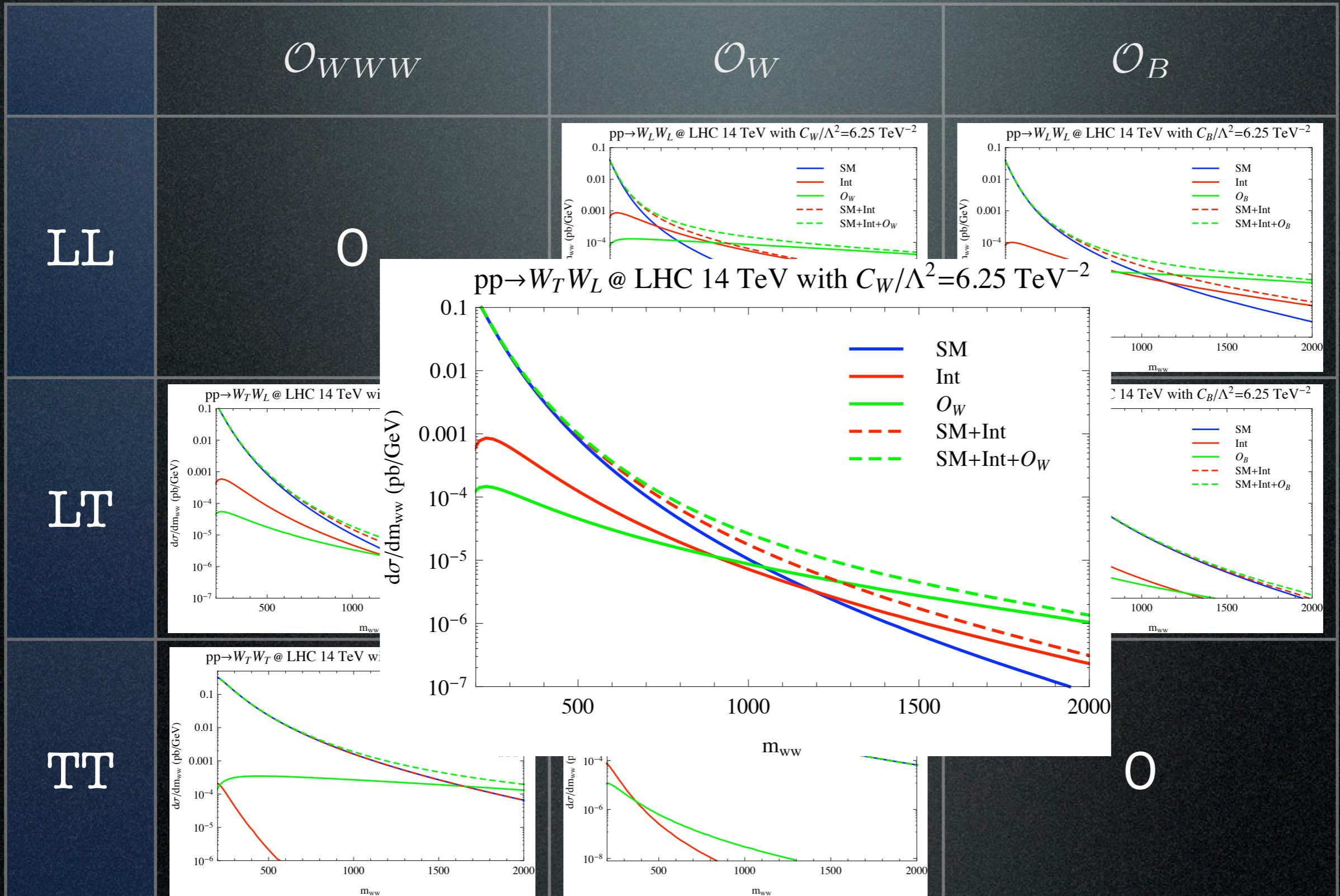
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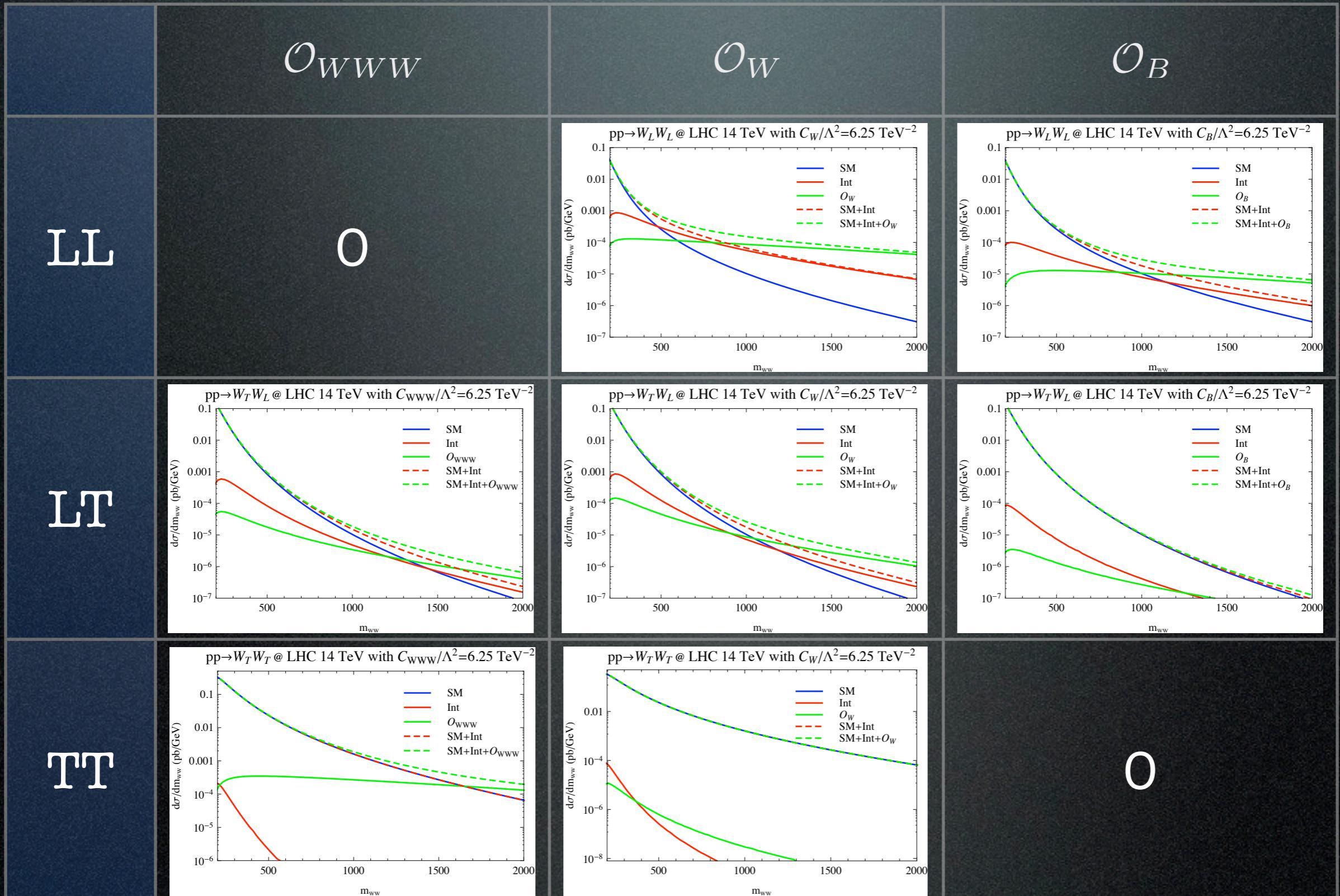
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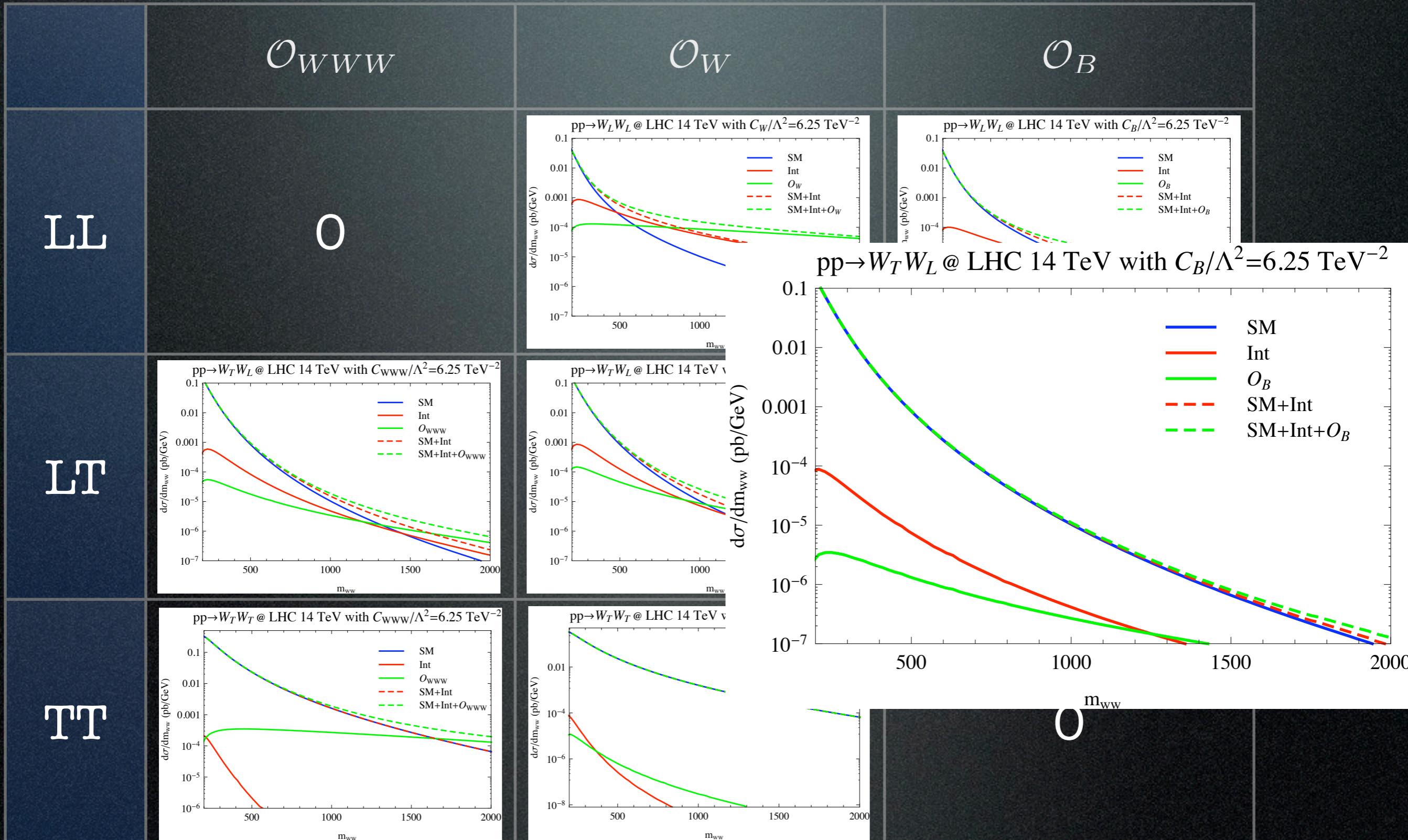
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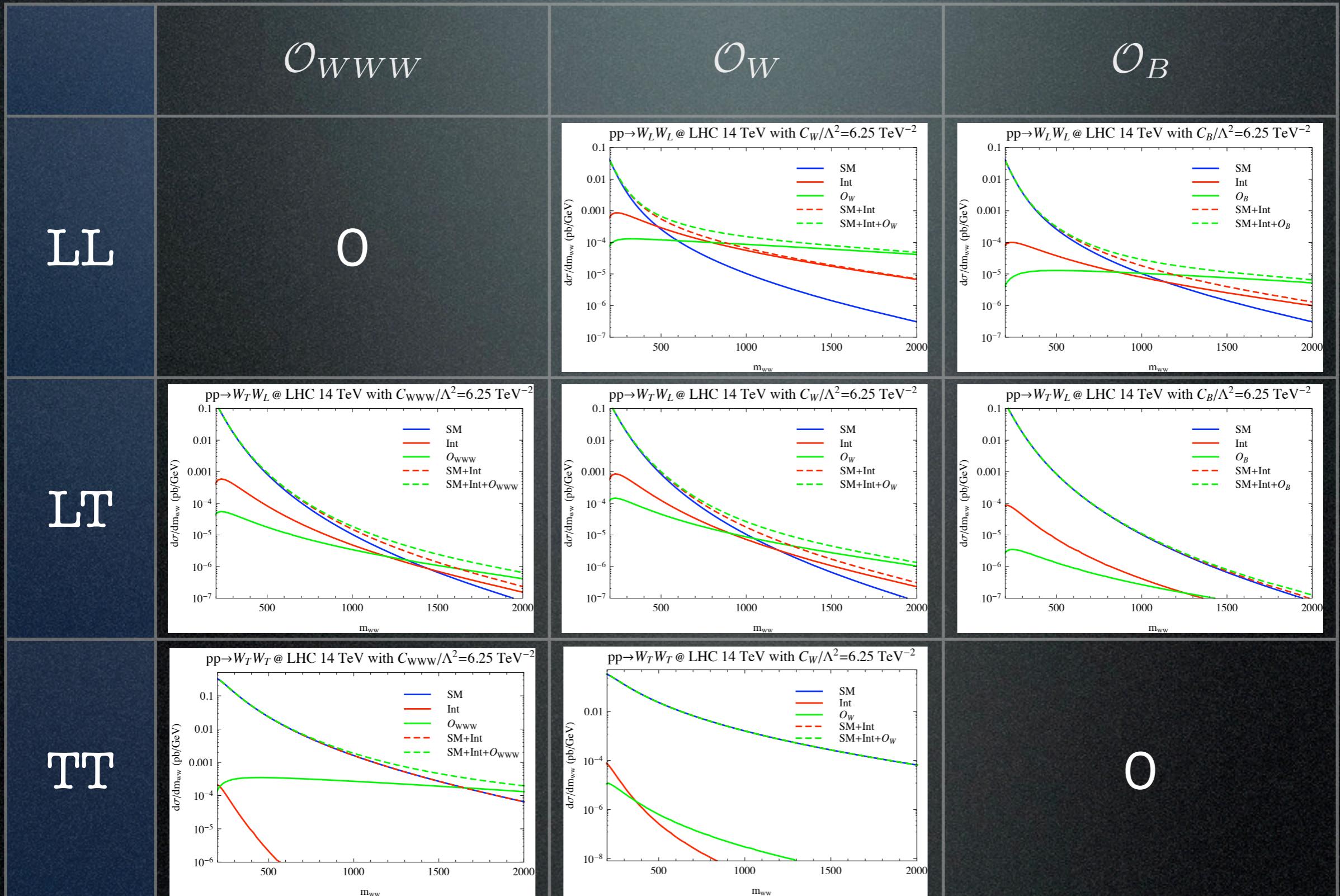
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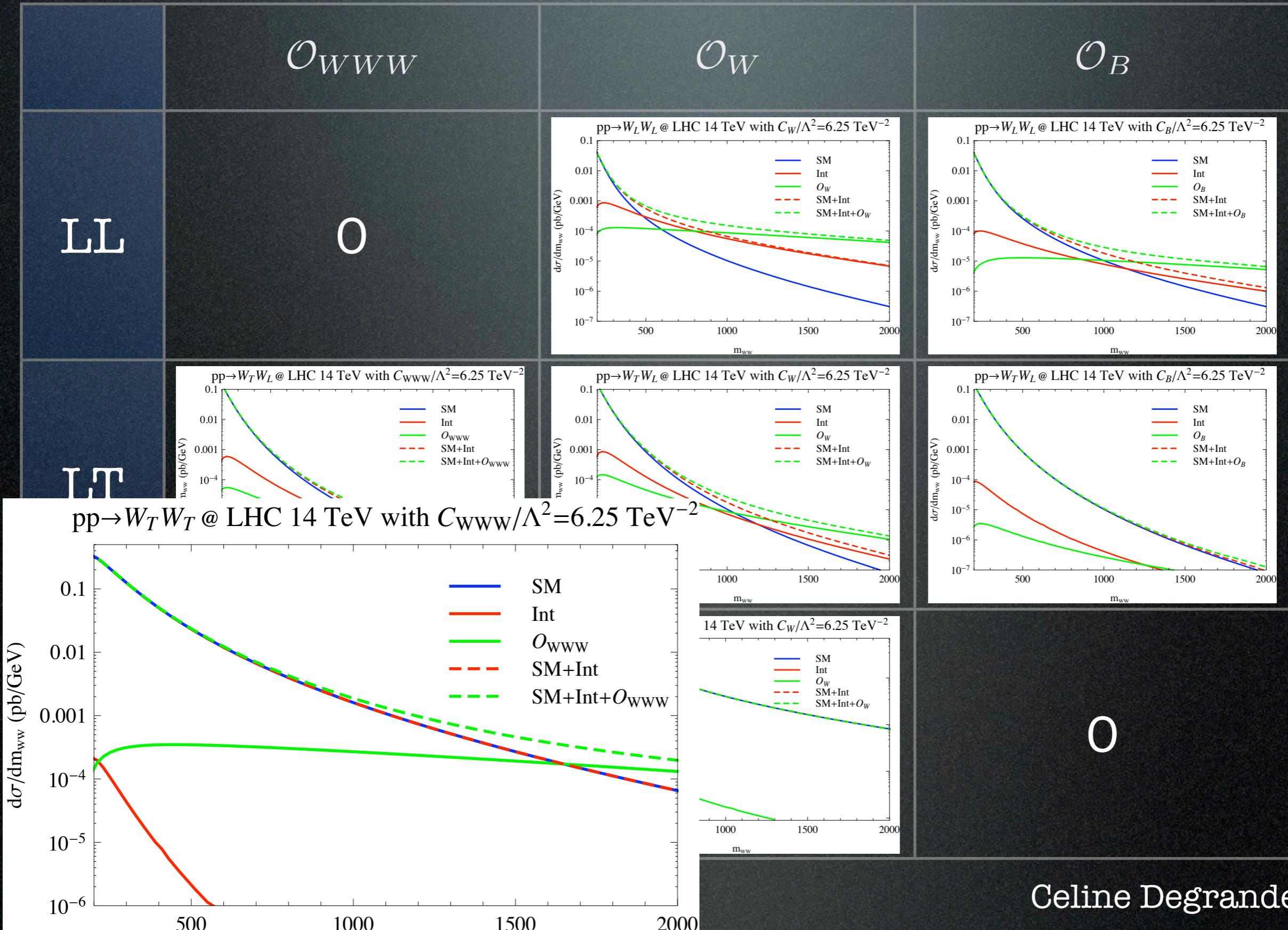
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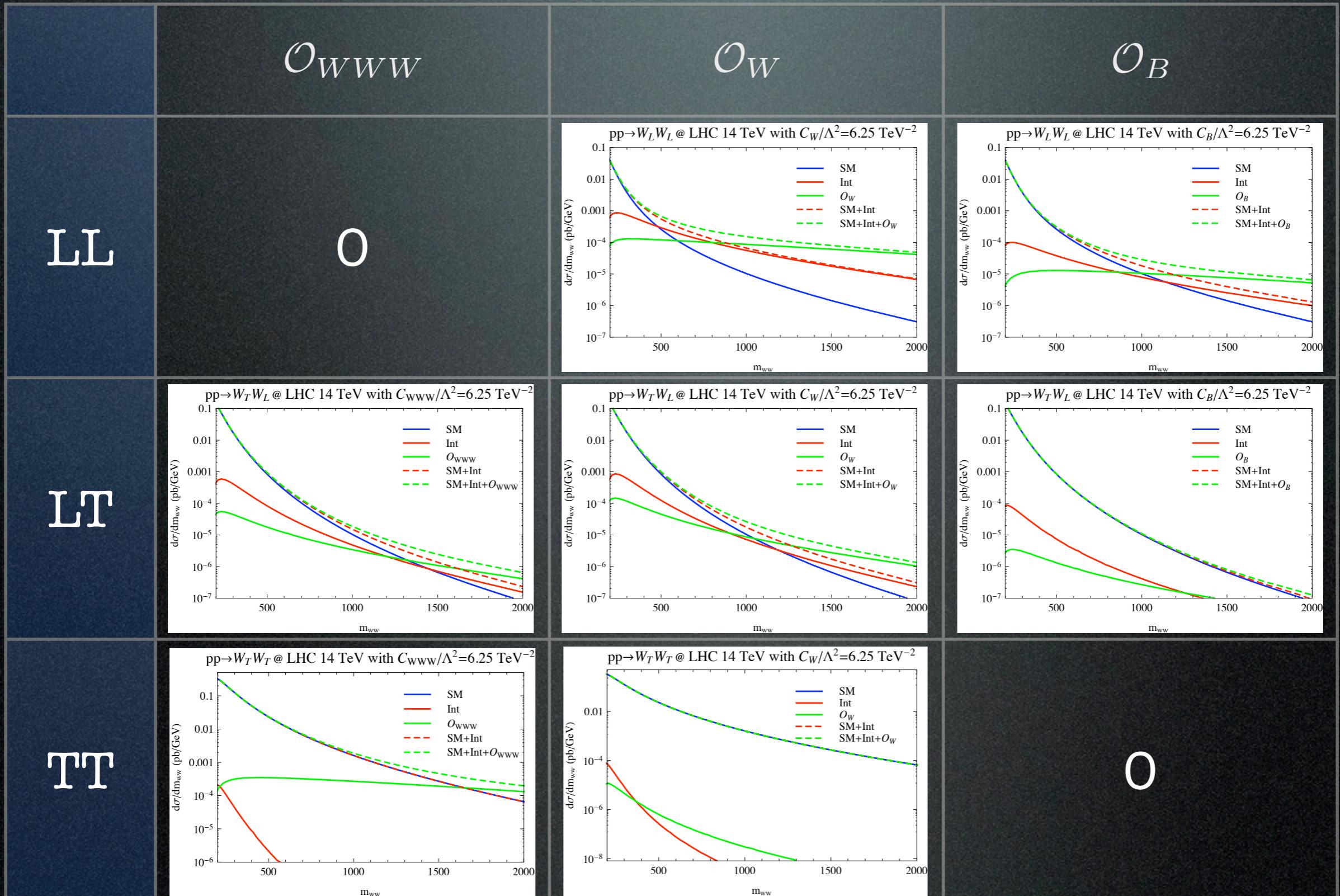
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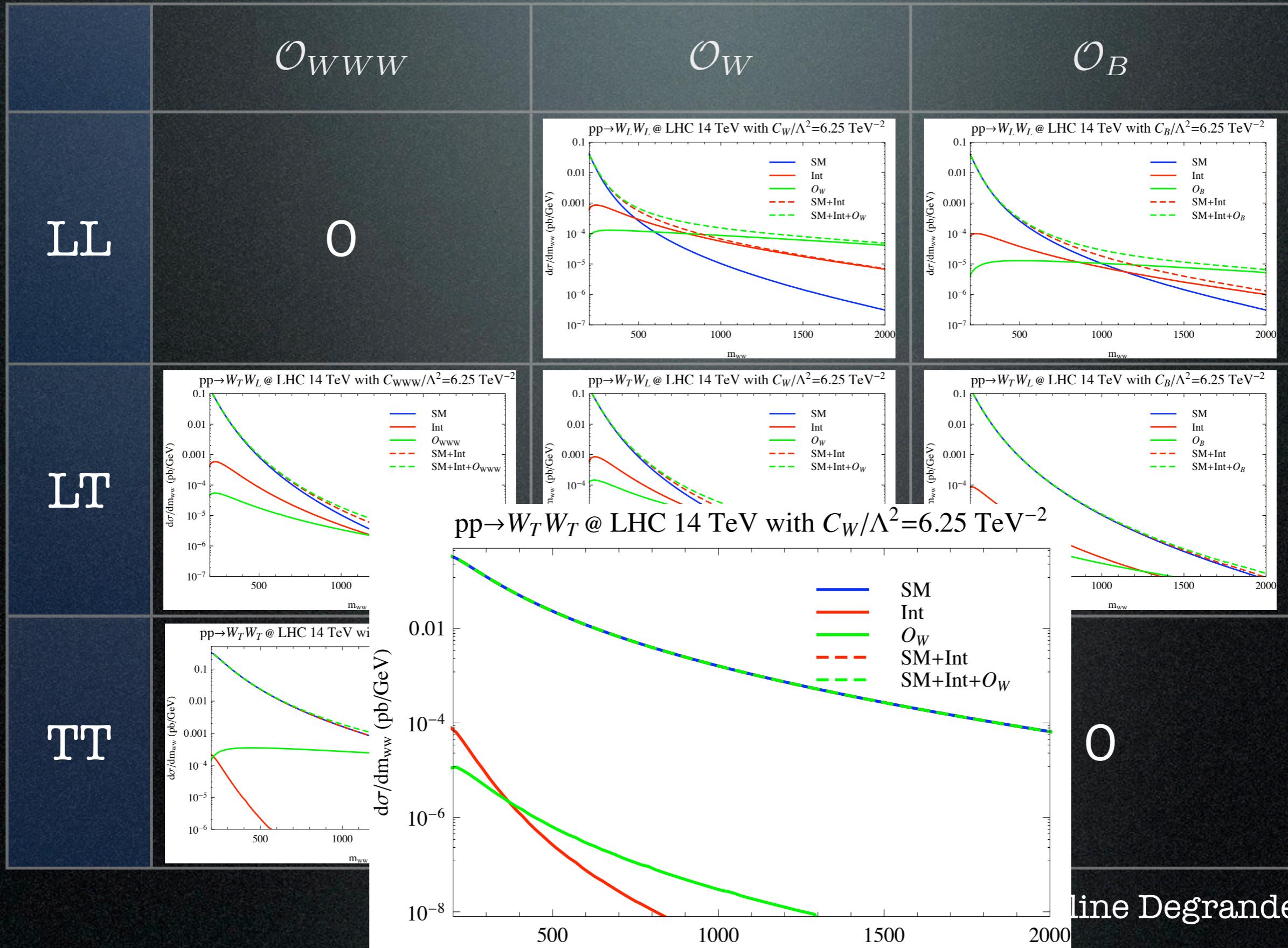
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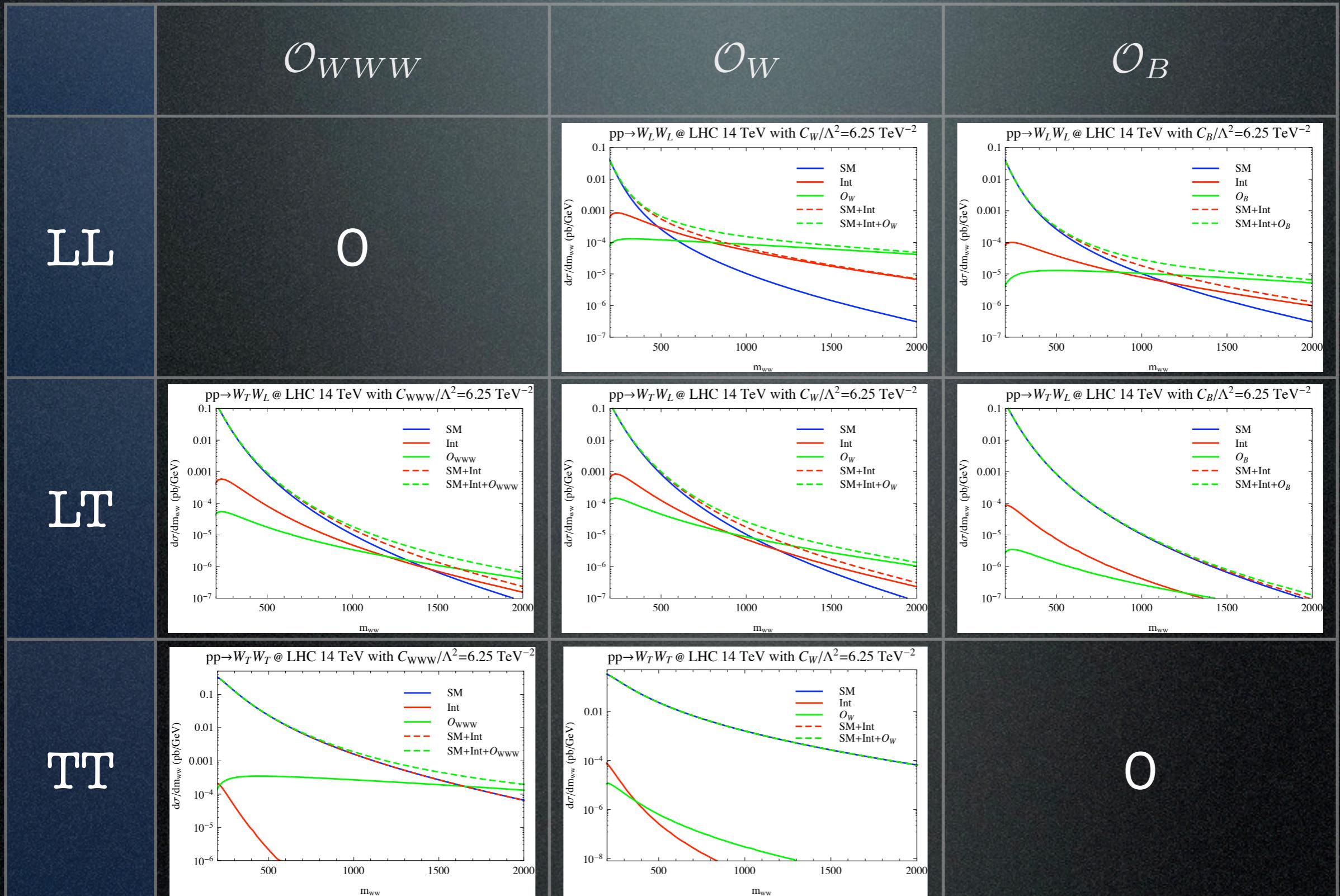
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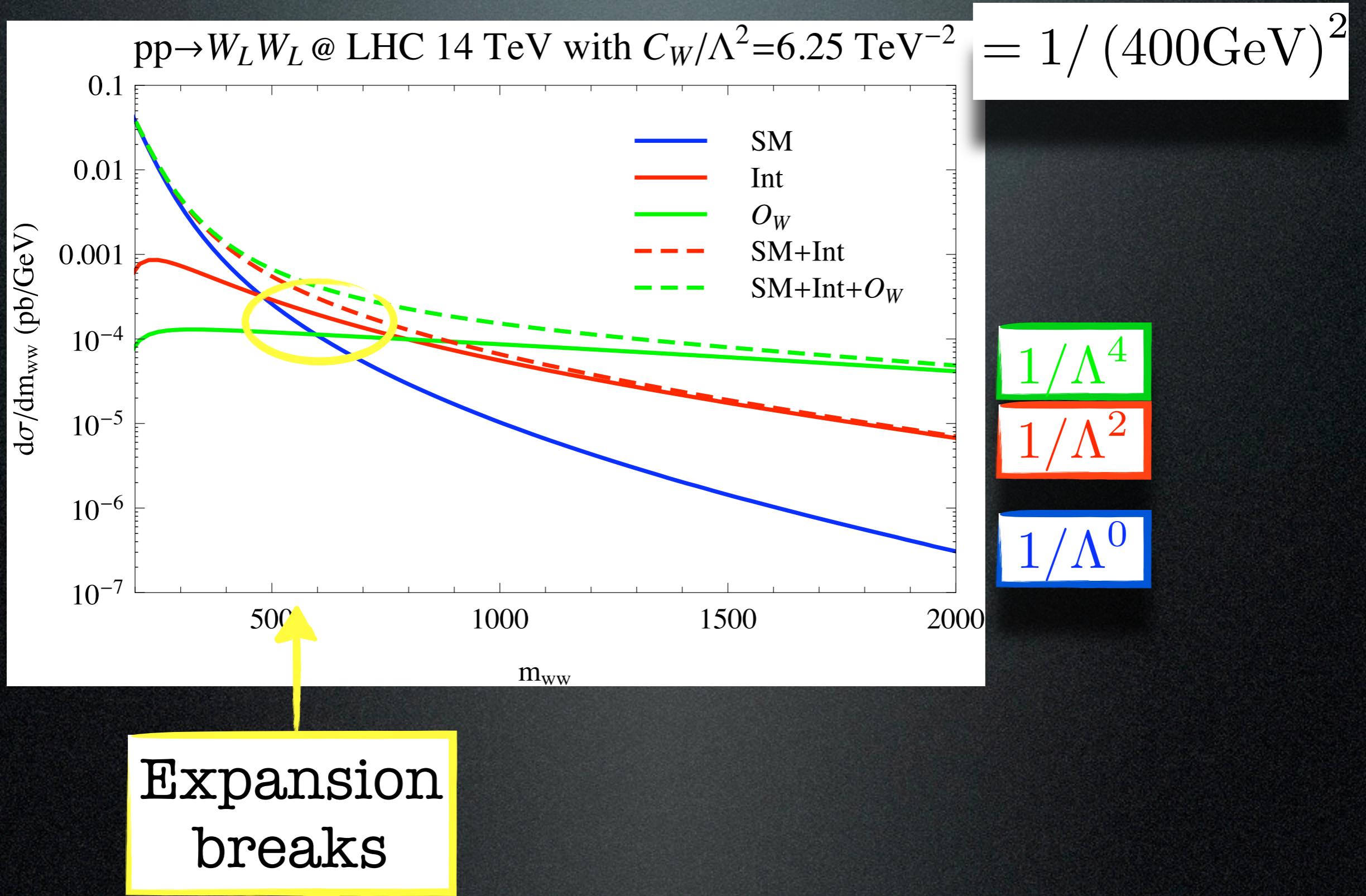
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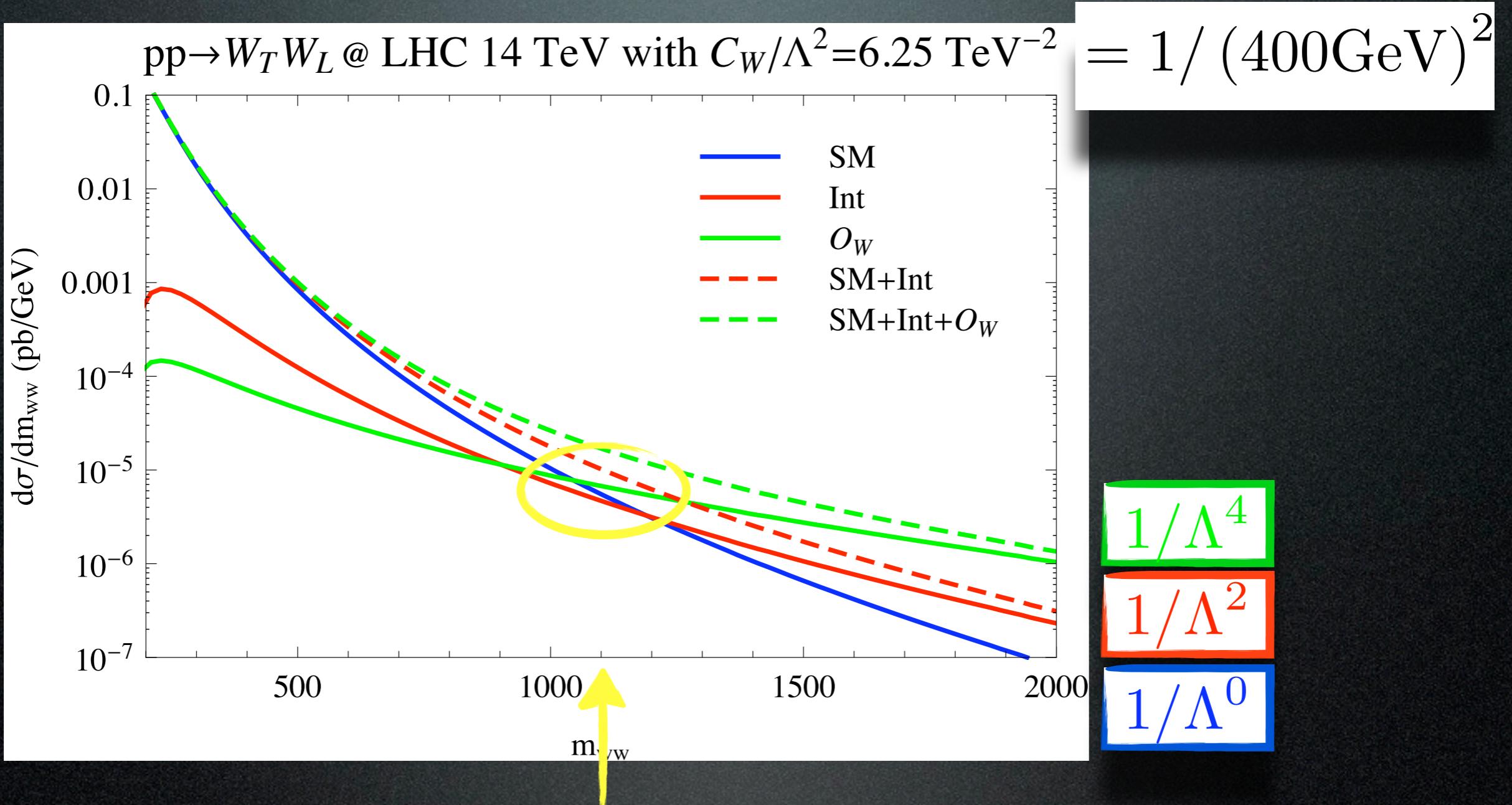
Invariant mass and polarisations

	\mathcal{O}_{WWW}	\mathcal{O}_W	\mathcal{O}_B	SM
LL	0	$1/(s)$	$1/(s)$	$1/s$
LT	$1/s(1)$	$1/s(1)$	$1/s(1)$	$1/s^2$
TT	$1/s(s)$	$1/s^2(1/s)$	0	$1/s$

Expansion and error

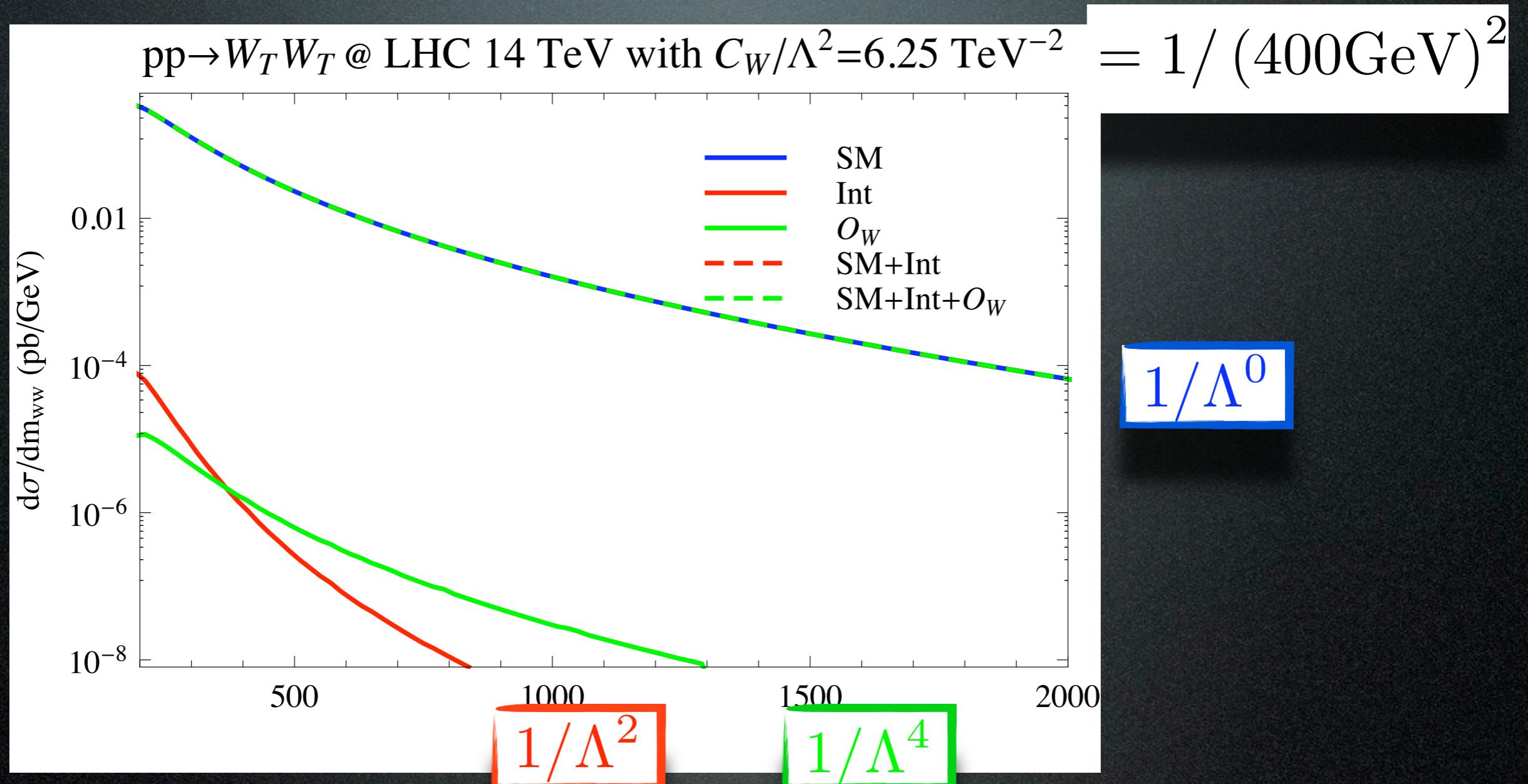


Expansion and error



Expansion
breaks

Expansion and error



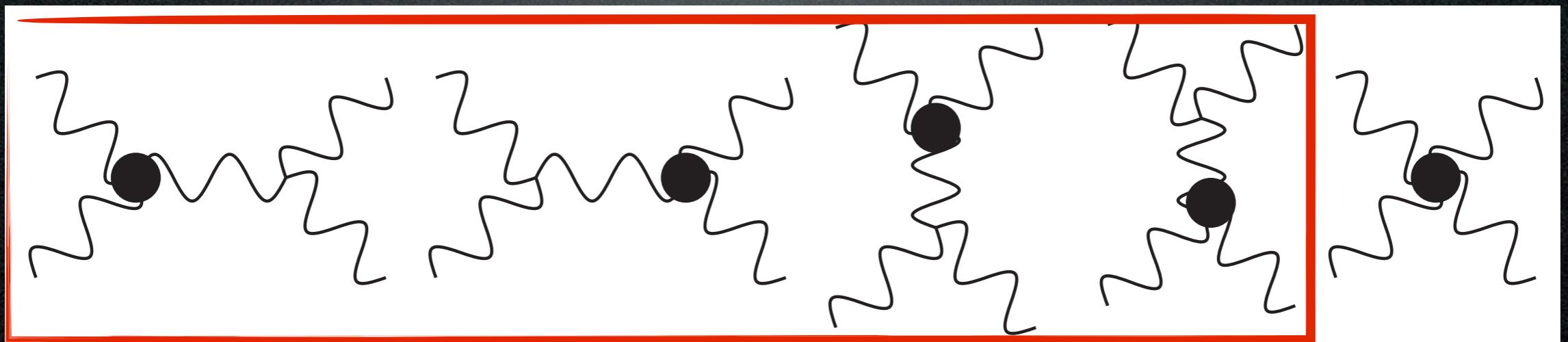
NP is suppressed : Bad estimate of the scale

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$pp>WWW$ or $pp>WWjj$

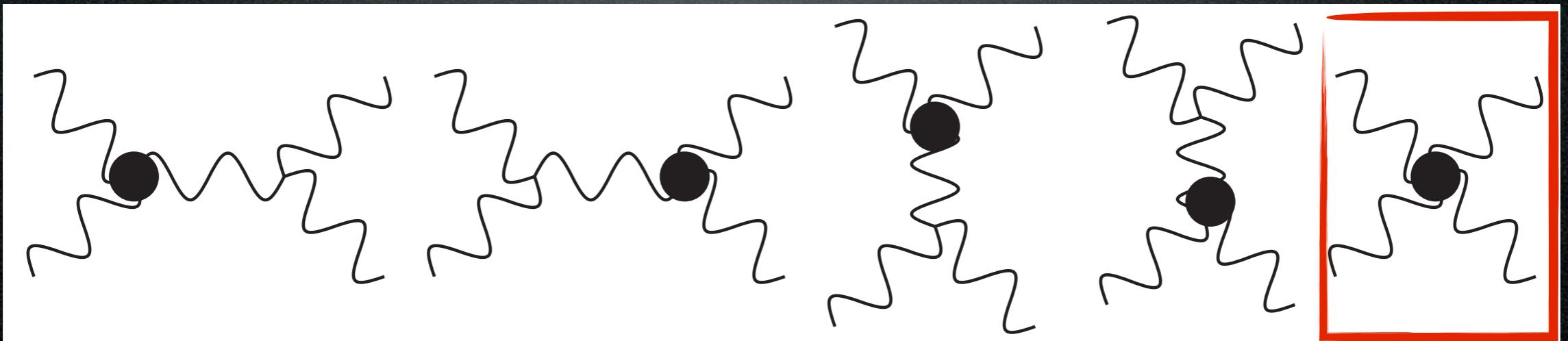
- Same operators for 4W amplitudes
- gauge invariance requires 3 and 4 legs vertices to be related



TGC's alone are not *gauge*
invariant

$pp>WWW$ or $pp>WWjj$

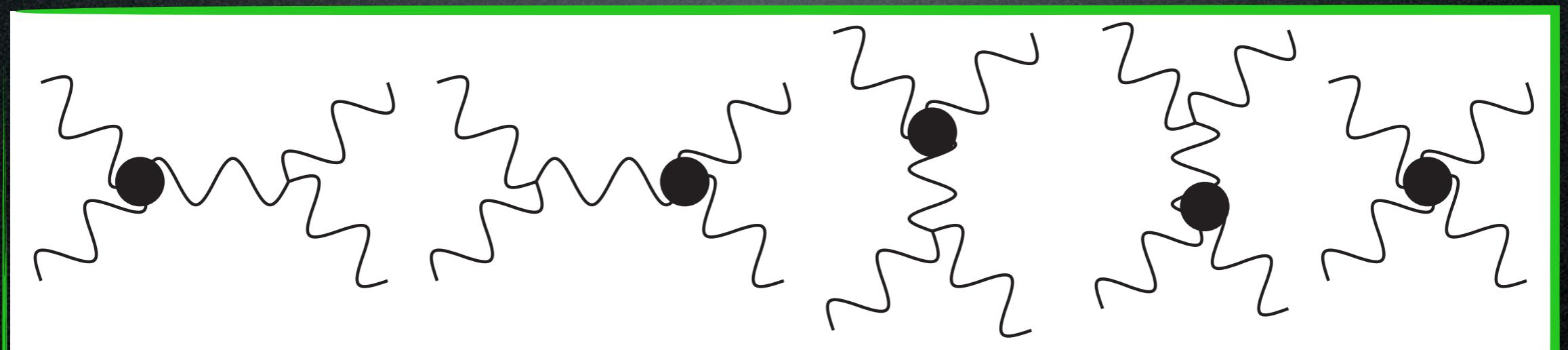
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QGC's alone are
not gauge
invariant

$pp>WWW$ or $pp>WWjj$

- Same operators for 4W amplitudes
- gauge invariance requires 3 and 4 legs vertices to be related



TGC's and QGC's from the dimension-six
operators are gauge invariant

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Concluding remarks

- EFT are a good way to search for heavy new physics
 - More predictive (guidance)
 - Satisfy unitarity
 - Take care of gauge invariance for any process
 - Allow loop computation
- EFT are available in MadGraph (<https://cp3.irmp.ucl.ac.be/projects/madgraph/wiki/Models/EWdim6>)

Back-up

Dim-6 versus dim-8

- Smaller effects or larger errors

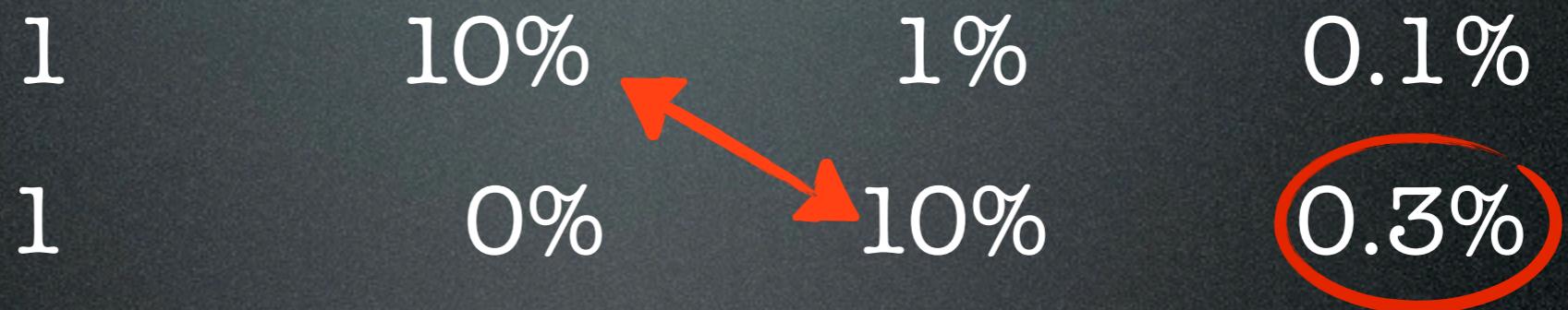
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1	10%	1%	0.1%
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Dim-6 versus dim-8

- Smaller effects or larger errors

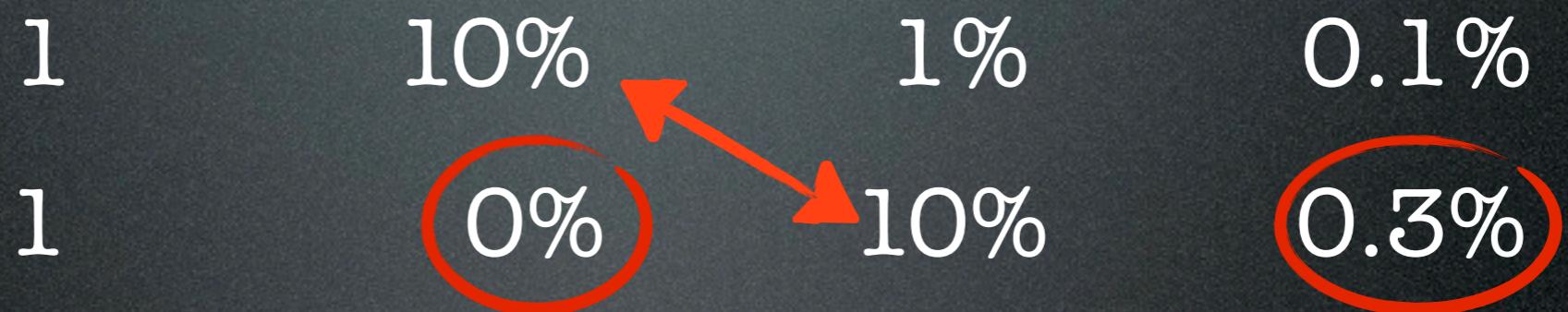
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- Extra assumptions

Dim-6 versus dim-8

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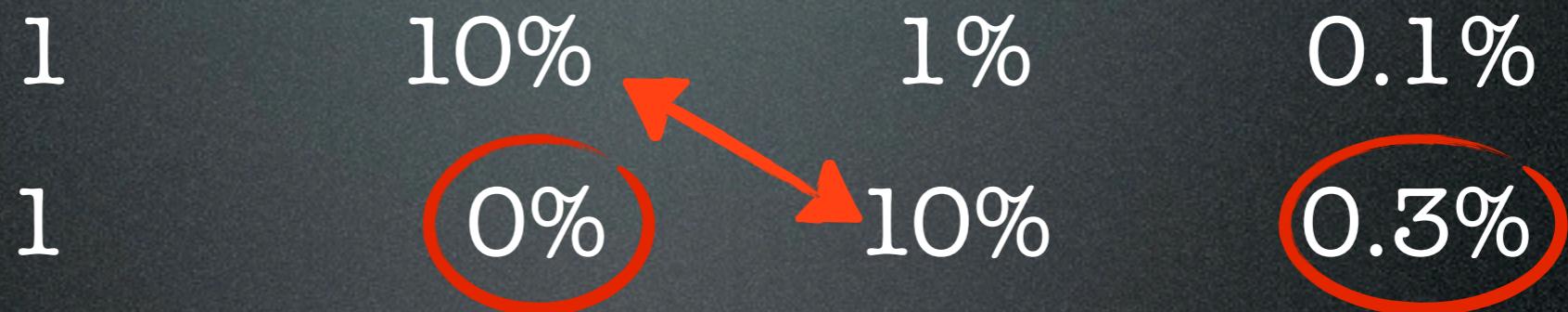


- Extra assumptions
- More parameters/less guidance

Dim-6 versus dim-8

- Smaller effects or larger errors

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- Extra assumptions
- More parameters/less guidance
- Can affect a new observable