

# What do we learn from the recent cosmic-ray positron measurements?

K. Blum\*, B. Katz\*, E. Waxman  
Weizmann Institute

\*currently at IAS, Princeton

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arXiv:1305.1324

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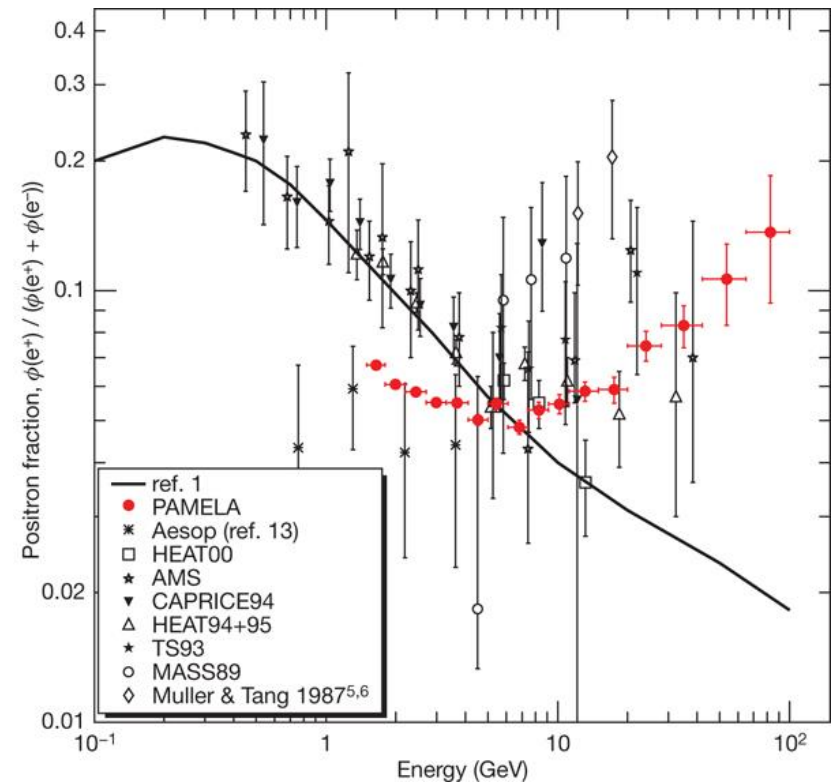
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### Letter

*Nature* **458**, 607–609 (2 April 2009) | doi:10.1038/nature07942; Received 28 February 2009

## An anomalous positron abundance in cosmic rays with energies 1.5–100 GeV

...Our results clearly show an increase in the positron abundance at high energy that cannot be understood by standard models describing the secondary production of cosmic-rays.



# The positron “anomaly”

- A common argument:

$$\begin{aligned} \frac{n_+(E)}{n_e(E)} &= \frac{\dot{n}_+ \tau_{+, \text{conf.}} f_{+, \text{rad}}}{\dot{n}_e \tau_{e, \text{conf.}} f_{e, \text{rad}}} \propto \frac{n_p \tau_{+, \text{conf.}} f_{+, \text{rad}}}{\dot{n}_e \tau_{e, \text{conf.}} f_{e, \text{rad}}} = \frac{\dot{n}_p \tau_{p, \text{conf.}} \tau_{+, \text{conf.}} f_{+, \text{rad}}}{\dot{n}_e \tau_{e, \text{conf.}} f_{e, \text{rad}}} \\ &= \frac{\dot{n}_p}{\dot{n}_e} \frac{\tau_{+, \text{conf.}}}{\tau_{e, \text{conf.}}} \frac{f_{+, \text{rad}}}{f_{e, \text{rad}}} \tau_{p, \text{conf.}} \propto \tau_{p, \text{conf.}} \propto E^{-0.5} \end{aligned}$$

- Unsubstantiated assumptions:

$$\dot{n}_p \propto \dot{n}_e, \quad \frac{\tau_{+, \text{conf.}}}{\tau_{e, \text{conf.}}} = 1, \quad \frac{f_{+, \text{rad}}}{f_{e, \text{rad}}} = 1$$

- The “anomaly” implies that (some of) these assumptions, which are not based on theory or observations, are not valid.

# What is the $e^+$ excess claim based on?

- On assumptions not supported by data/theory
- \* primary  $e^-$  &  $p$  produced with the same spectrum,  
and  $e^-$  and  $e^+$  suffer same  $f_{\text{rad}}$   
 $\rightarrow e^+/e^- \sim \Sigma_{\text{sec}} \sim \varepsilon^{-0.5}$

Or

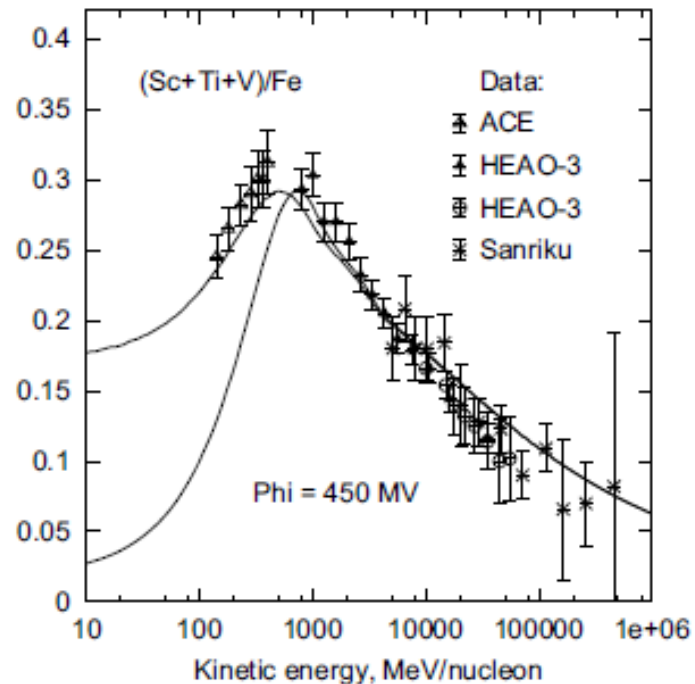
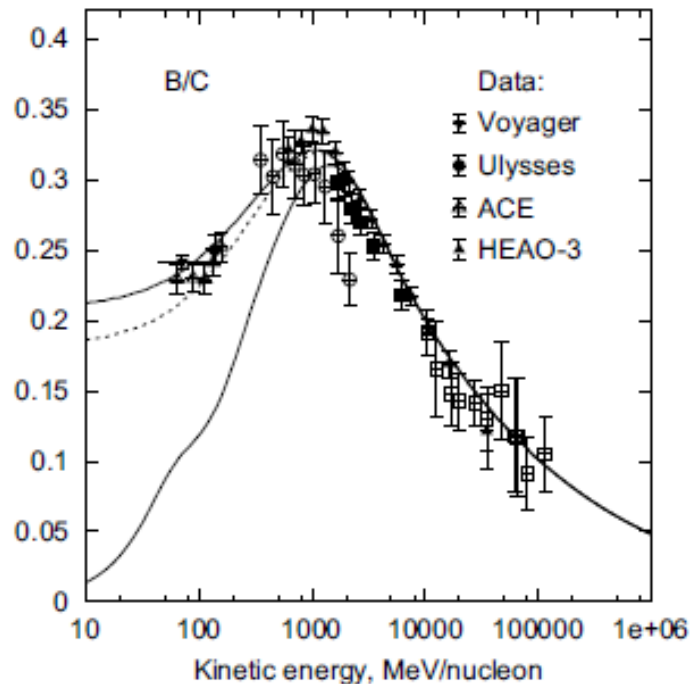
- \* detailed assumptions RE CR propagation,  
e.g. isotropic diffusion,  $D \sim \varepsilon^\delta$ ,  
within an  $\varepsilon$ -independent box  
 $\rightarrow f_{\text{rad}} \sim \varepsilon^{(\delta-1)/2}$
- If PAMELA/AMS correct, these assumptions are wrong

# What we really know

- For all secondary nuclei (e.g.  $C \rightarrow B, Be; \dots$ ):

$$n_{\text{sec}}(\varepsilon) = n_{\text{prim.}} \times \sigma_{\text{eff.}} \times n_{p,\text{ISM}} ct_{\text{conf.}} = n_{\text{prim.}} \times \sigma_{\text{eff.}} \times \frac{\Sigma_{\text{sec}}(\varepsilon/Z)}{m_p}$$

$$\Sigma_{\text{sec}} \approx 8.7 \left( \frac{\varepsilon}{10Z \text{ GeV}} \right)^{-0.5} \text{ g/cm}^2 ; \quad \frac{n_{i,\text{sec}}}{n_{j,\text{sec}}} = \frac{\tilde{\sigma}_i}{\tilde{\sigma}_j}$$



# Why does it work?

- We have no basic principles model for CR propagation.
- In general

$$\frac{dn_{i,\text{sec},0}}{dE} = \sum_{j \neq i} \int d^3x \int dE' \frac{dn_j(x, E')}{dE'} n(x) \frac{d\sigma_{j \rightarrow i}(E', E)}{dE} cG(x, E/Z)$$

- If: The CR composition ( $n_j/n_p$ ) is independent of  $x$ ,  
The CR spectrum is independent of  $x$   
(or:  $E/Z$  the same for prim. & sec.),

Then:

$$\frac{dn_{i,\text{sec},0}}{dE} = n_{p,0} \frac{d\sigma_i}{dE} \frac{\Sigma_{\text{sec}}}{m_p}, \quad \text{where}$$

$$\frac{d\sigma_i}{dE} \equiv \sum_{j \neq i} \frac{n_j}{n_p} \int dE' n_j^{-1} \frac{dn_j(E')}{dE'} \frac{d\sigma_{j \rightarrow i}(E', E)}{dE}$$

$$t_{\text{sec}}(E/Z) \equiv \int d^3x \frac{n_p}{n_{p,0}} \frac{n(x)}{n_0} G(x, E/Z), \quad \Sigma_{\text{sec}}(E/Z) \equiv n_0 m_p c t_{\text{sec}}(E/Z)$$

# What we really know: application to $e^+$

- Positron secondaries:

$$n_{i,\text{sec}}(E) = \left[ \sum_{j \neq i} n_j \frac{\sigma_{j \rightarrow i}}{m_p} - \frac{n_i \sigma_i}{m_p} \right] \Sigma_{\text{sec}}(E/Z)$$

$$\frac{n_{i,\text{sec}}}{n_{j,\text{sec}}} = \frac{\tilde{\sigma}_i}{\tilde{\sigma}_j} \Rightarrow n_+ = n_p \tilde{\sigma}_+ \frac{\Sigma_{\text{sec}}(E/Z)}{m_p} \times f_{+,\text{rad}}$$

Known from Lab

Measured from CR sec.

Suppression due to synchrotron  
And inverse Compton energy loss

- A robust, model independent, upper limit:

$$n_+ \leq n_p \tilde{\sigma}_+ \Sigma_{\text{sec}}(E/Z) / m_p$$

- Estimate  $f_{\text{rad}}$  at 20 GeV from  $^{10}\text{Be}$  ( $10^7\text{yr}$ ):  $\sim 0.3$   
(CMB and starlight  $\sim 1\text{eV}/\text{cm}^3$ ,  $e^+$  lifetime  $\sim 10^7\text{yr}$ ).

# Some cautionary notes

- $t_{\text{sec}}$  is not the residence time

$$t_{\text{res}}(E/Z) \equiv \int d^3x \frac{\dot{n}(x)}{\langle \dot{n} \rangle} \int \frac{d^3x'}{V} G(x, E/Z; x')$$

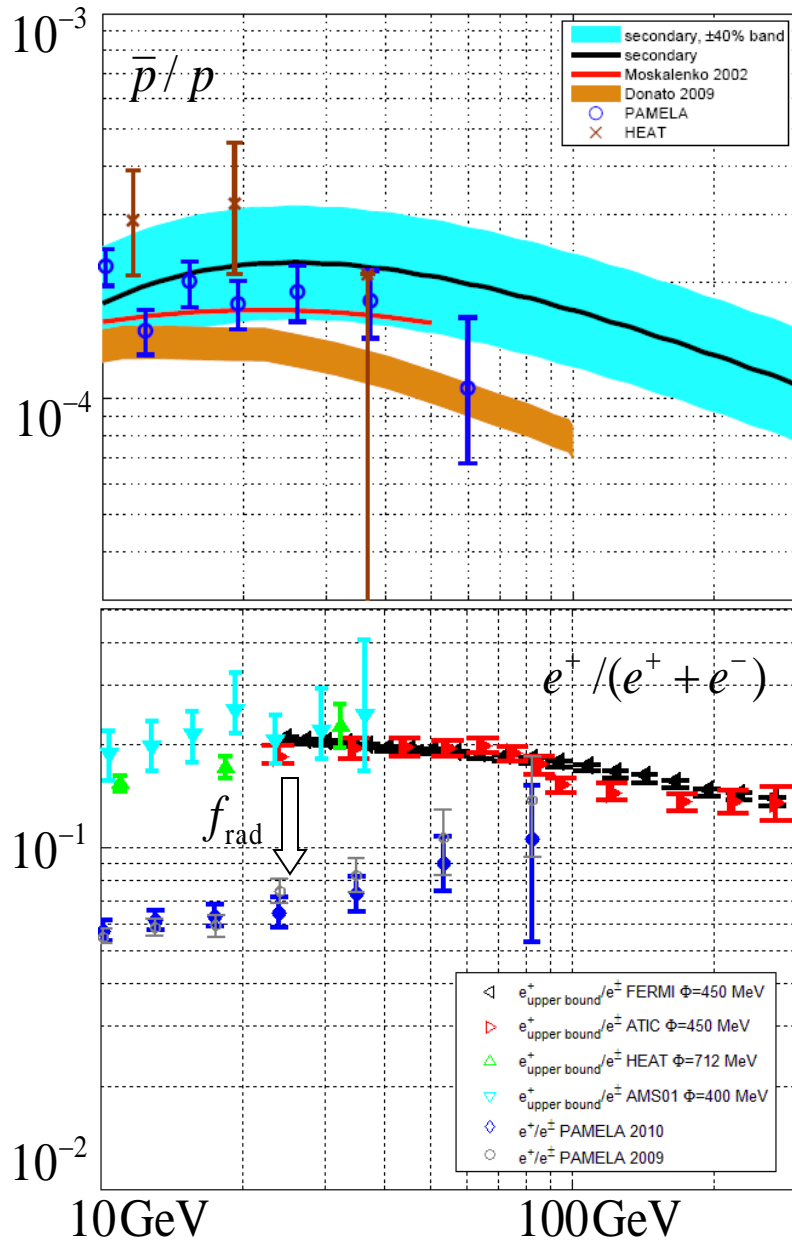
- If  $G(x, x')$  is independent of  $x'$  then

$$t_{\text{res}} = \frac{\dot{n}_0}{\langle \dot{n} \rangle} t_{\text{sec}}, \quad \text{and} \quad n_{i,0}(E) = \dot{n}_0 t_{\text{sec}} = \langle \dot{n} \rangle t_{\text{res}}$$

- For particles with life time  $\tau$ ,  
 $f_{\text{sec}}$  depends on model assumptions. It may, e.g.,  
attain the values  $\tau/t_{\text{res}}$  or  $V(t_{\text{prop}} < \tau)/V_G$  under  
different model assumptions.



# PAMELA 09-10



- For all secondaries (e.g. anti-p)

$$n_i(\varepsilon) = \left[ \sum_{j \neq i} n_j \frac{\sigma_{j \rightarrow i}}{m_p} - \frac{n_i \sigma_i}{m_p} \right] \Sigma_{\text{sec}}(\varepsilon/Z)$$

$$\Sigma_{\text{sec}} \approx 8.7 \left( \frac{\varepsilon}{10Z \text{ GeV}} \right)^{-0.5} \text{ g/cm}^2$$

- Radiative  $e^+$  losses- depend on propagation in Galaxy (poorly understood)

$$n_+(\varepsilon) = n_p \tilde{\sigma}_+ \frac{\Sigma_{\text{sec}}(\varepsilon/Z)}{m_p} f_{\text{rad}}$$

- \* At  $\sim 20 \text{ GeV}$ :  $f_{\text{rad}} \sim 0.3 \sim f_{10\text{Be}}$

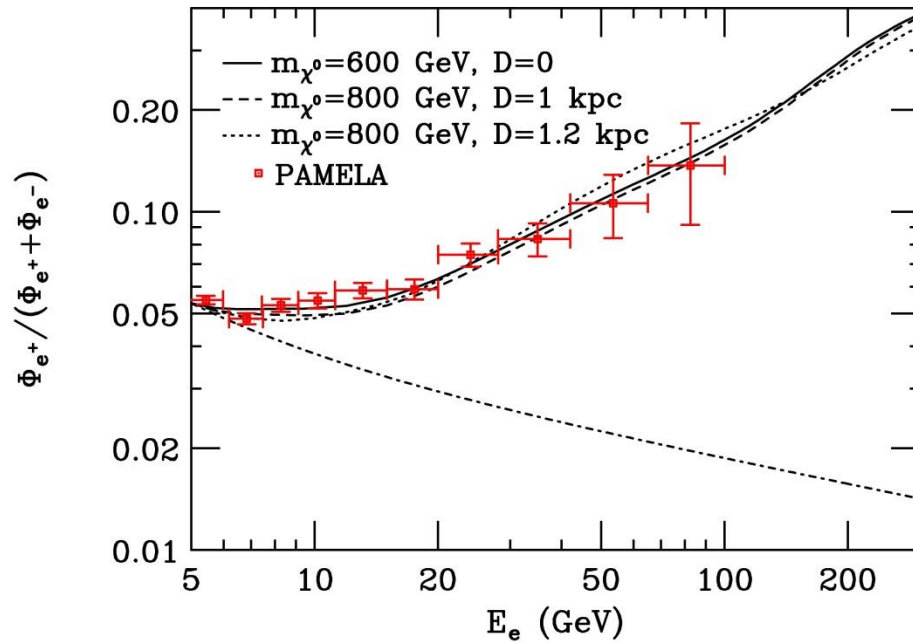
→  $e^+$  consistent with 2<sup>nd</sup>ary origin

- \* Above  $20 \text{ GeV}$ :

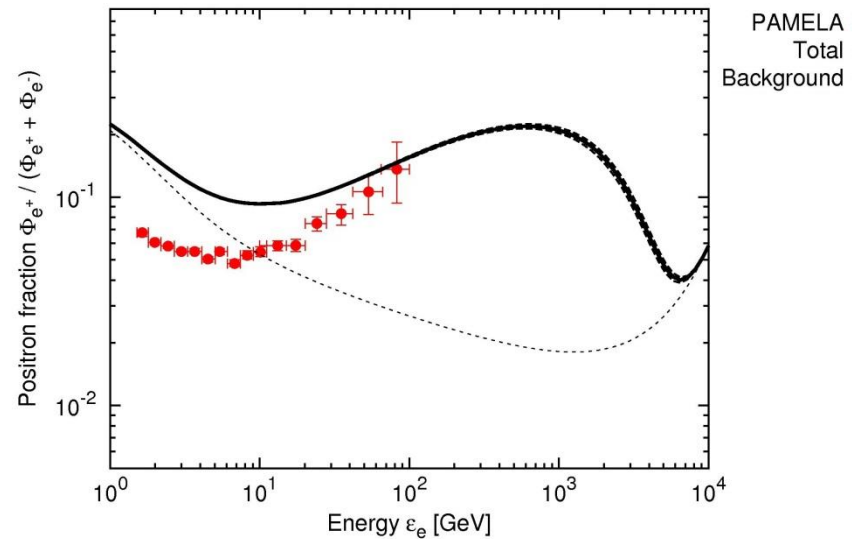
If PAMELA correct

→ energy independent  $f_{\text{rad}}(\varepsilon)$

# Primary e<sup>+</sup> sources

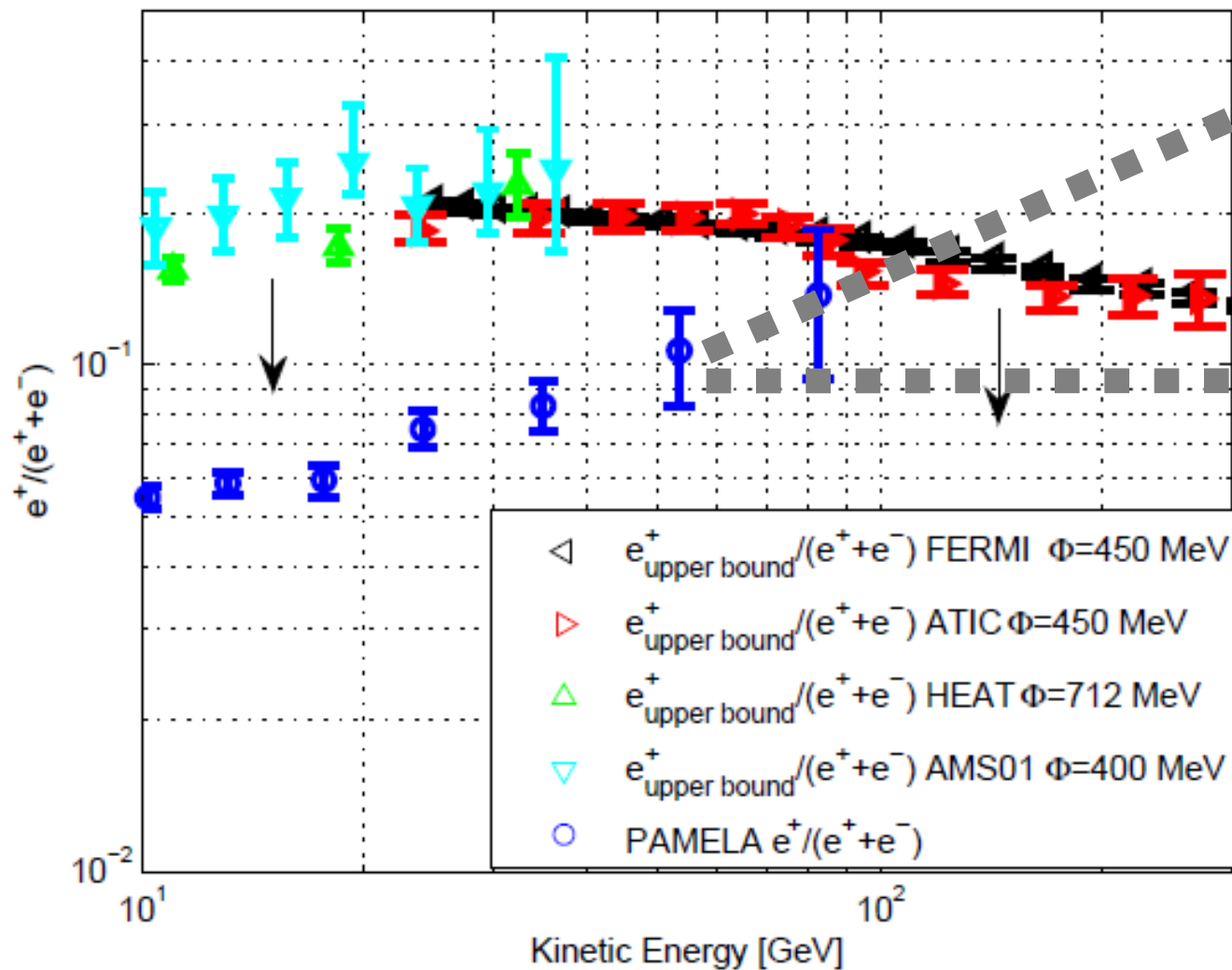


DM annihilation [Hooper et al. 09]



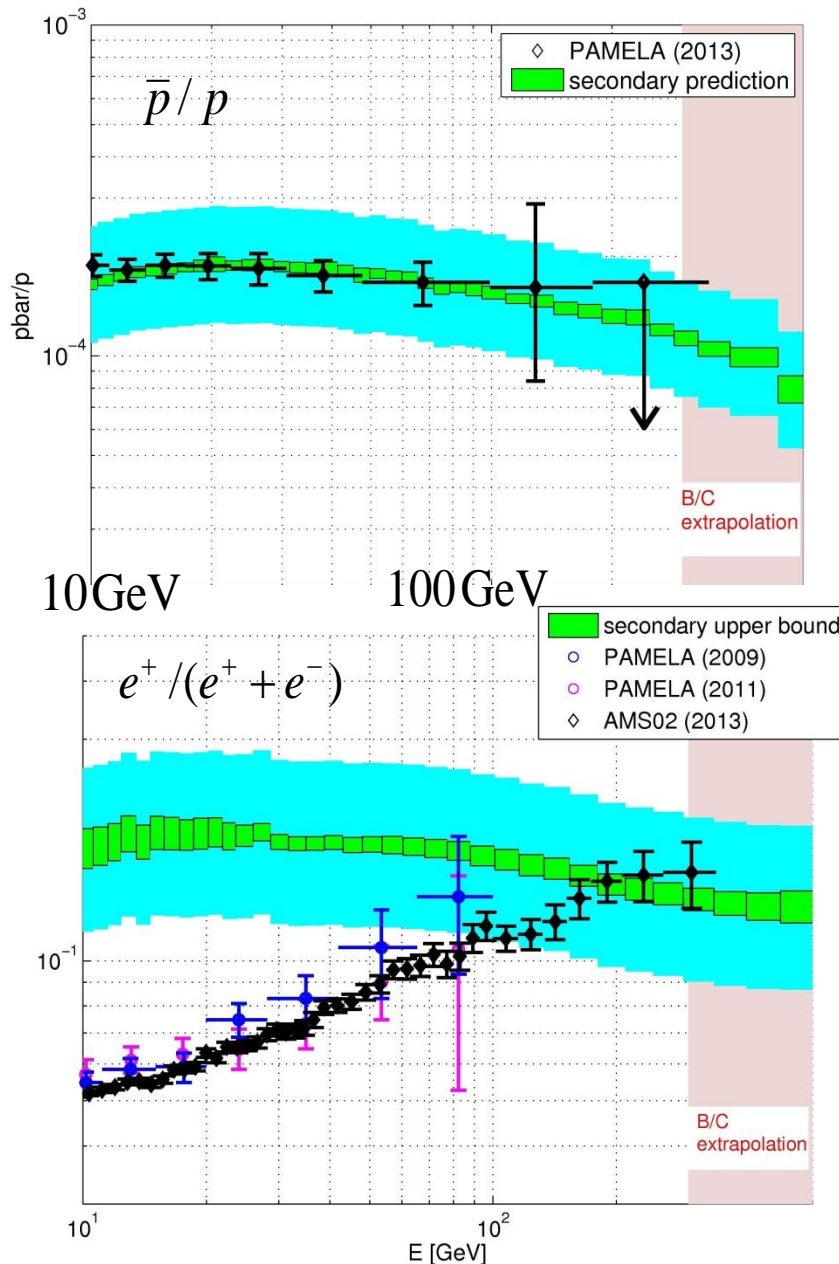
Pulsars [Kashiyama et al. 11]

New primary sources



Secondary origin

# PAMELA & AMS 2013



- Anti-p consistent with prediction.
- $e^+$  flux saturates at the secondary upper bound.
- The ABSOLUTE  $e^+$  flux matches the secondary bound.
- In all primary  $e^+$  models (DM, pulsars...) there is no intrinsic scale that would explain why the ABSOLUTE observed flux lies near the data-driven secondary bound.

# Conclusions

- Anti-p and  $e^+$  measurements are in excellent agreement with secondary production model predictions.
- The saturation of the ABSOLUTE  $e^+$  flux at the predicted secondary upper bound is a strong indication for a secondary origin, and is not expected in primary production models.
- Upcoming measurements at yet higher energy will further test the validity of the model.
- The main constraints that may be derived from the  $e^+$  measurements are on models of Galactic CR propagation.
- In particular, the measurements constrain  $f_{\text{rad}}(e^+)$ .

# Implications of $f_{\text{sec}}(e^+) \sim 0.3$ - an example

Under some model assumptions

$$f_{\text{sec}}(e^+) \sim t_{\text{cool}}/t_{\text{res}}$$

which would imply

$$t_{\text{res}}(E/Z=10\text{GeV}) > 30\text{Myr}, \quad t_{\text{res}}(E/Z=200\text{GeV}) < 1\text{Myr}$$

and using  $\Sigma_{\text{sec}}$

$$\langle n_{\text{ISM}} \rangle (E/Z=10\text{GeV}) < 0.2\text{g/cc},$$

$$\langle n_{\text{ISM}} \rangle (E/Z=200\text{GeV}) > 0.6\text{g/cc},$$

which in turn implies that the CR halo scale height decreases with  $E$ .

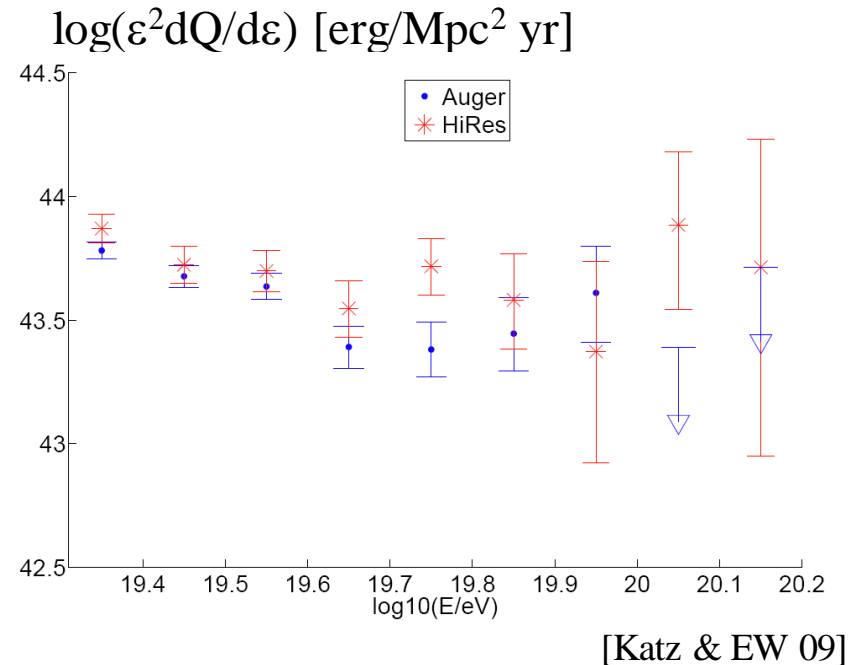
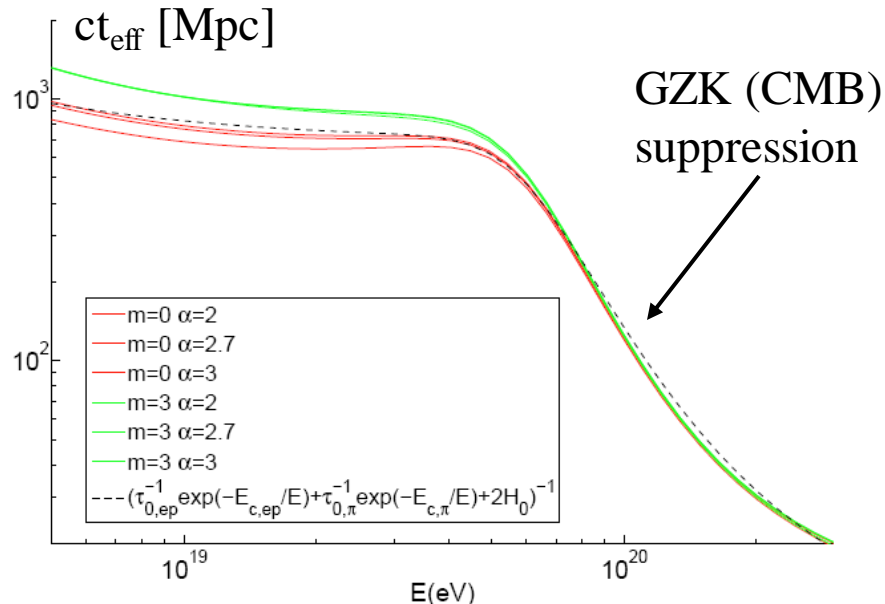
*A note on the IceCube detection*

# UHE: Flux & Generation Spectrum

- $\varepsilon^2(dN/d\varepsilon)_{\text{Observed}} = \varepsilon^2(dQ/d\varepsilon) t_{\text{eff.}}$

$$(t_{\text{eff.}} : p + \gamma_{\text{CMB}} \rightarrow N + \pi)$$

Assume:  $p, dQ/d\varepsilon \sim (1+z)^m \varepsilon^{-\alpha}$



- $>10^{19.3} \text{ eV}$ : consistent with protons,  $\varepsilon^2(dQ/d\varepsilon) = 0.5(+0.2) \times 10^{44} \text{ erg/Mpc}^3 \text{ yr} + \text{GZK}$



# HE $\nu$ : UHECR bound

- $p + \gamma \rightarrow N + \pi$

$$\pi^0 \rightarrow 2\gamma; \quad \pi^+ \rightarrow e^+ + \nu_e + \nu_\mu + \bar{\nu}_\mu$$

→ Identify UHECR sources

Study BH accretion/acceleration physics

- For all known sources,  $\tau_{\gamma p} \leq 1$ :

$$\varepsilon_\nu^2 \frac{dj_\nu}{d\varepsilon_\nu} \leq \Phi_{\text{WB}} \equiv 10^{-8} \zeta \left( \frac{\varepsilon^2 dQ/d\varepsilon}{10^{44} \text{ erg/Mpc}^3 \text{ yr}} \right) \frac{\text{GeV}}{\text{cm}^2 \text{ s sr}} \quad [\text{EW \& Bahcall 99; Bahcall \& EW 01}]$$

$$\zeta = 1, 5 \quad \text{for} \quad f(z) = 1, (1+z)^3$$

- If X-G p's:  $\varepsilon_\nu^2 \frac{dj_\nu}{d\varepsilon_\nu} (10^{19} \text{ eV}) = \Phi_{\text{WB}}$

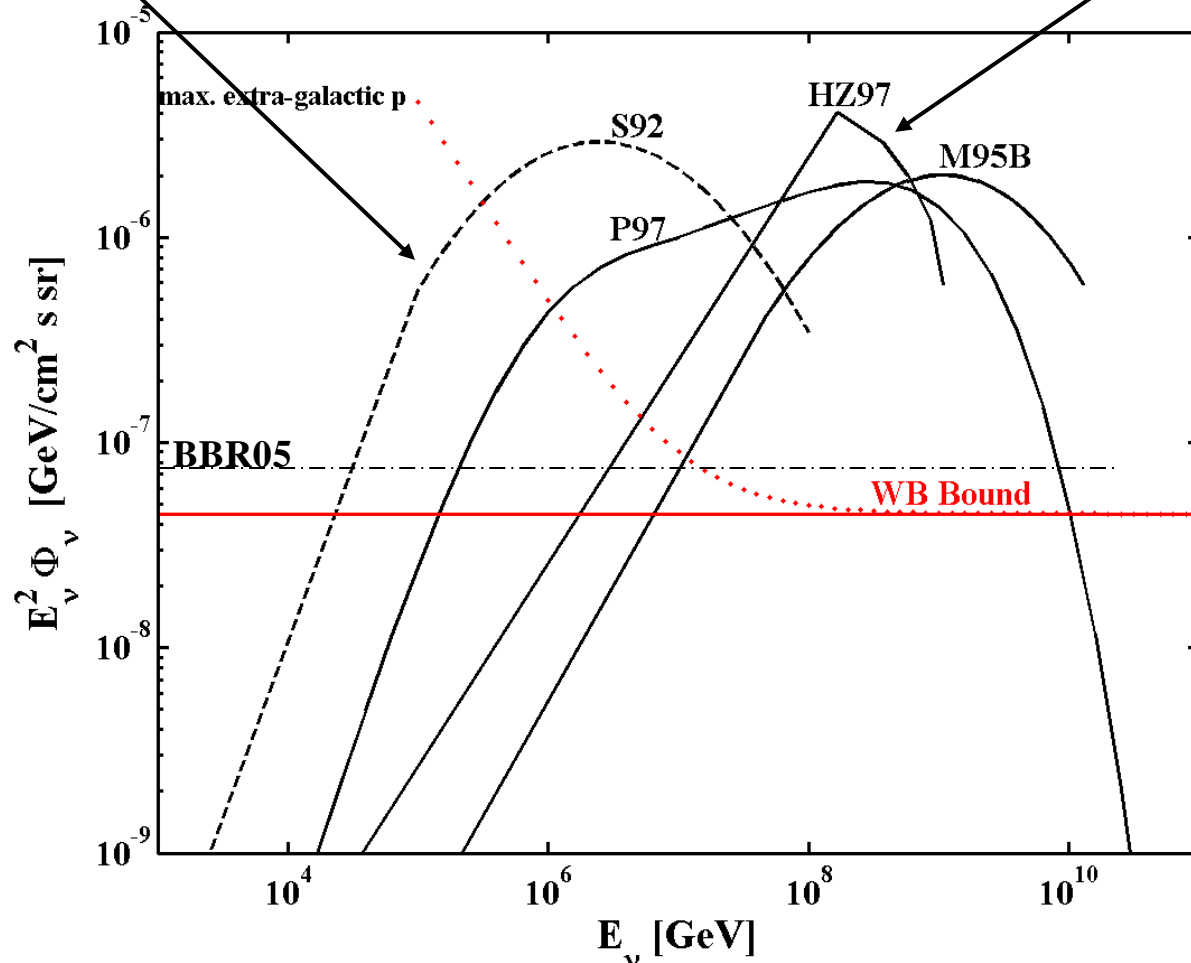
[Berezinsky & Zatsepin 69]

→ Identify primaries, determine  $f(z)$

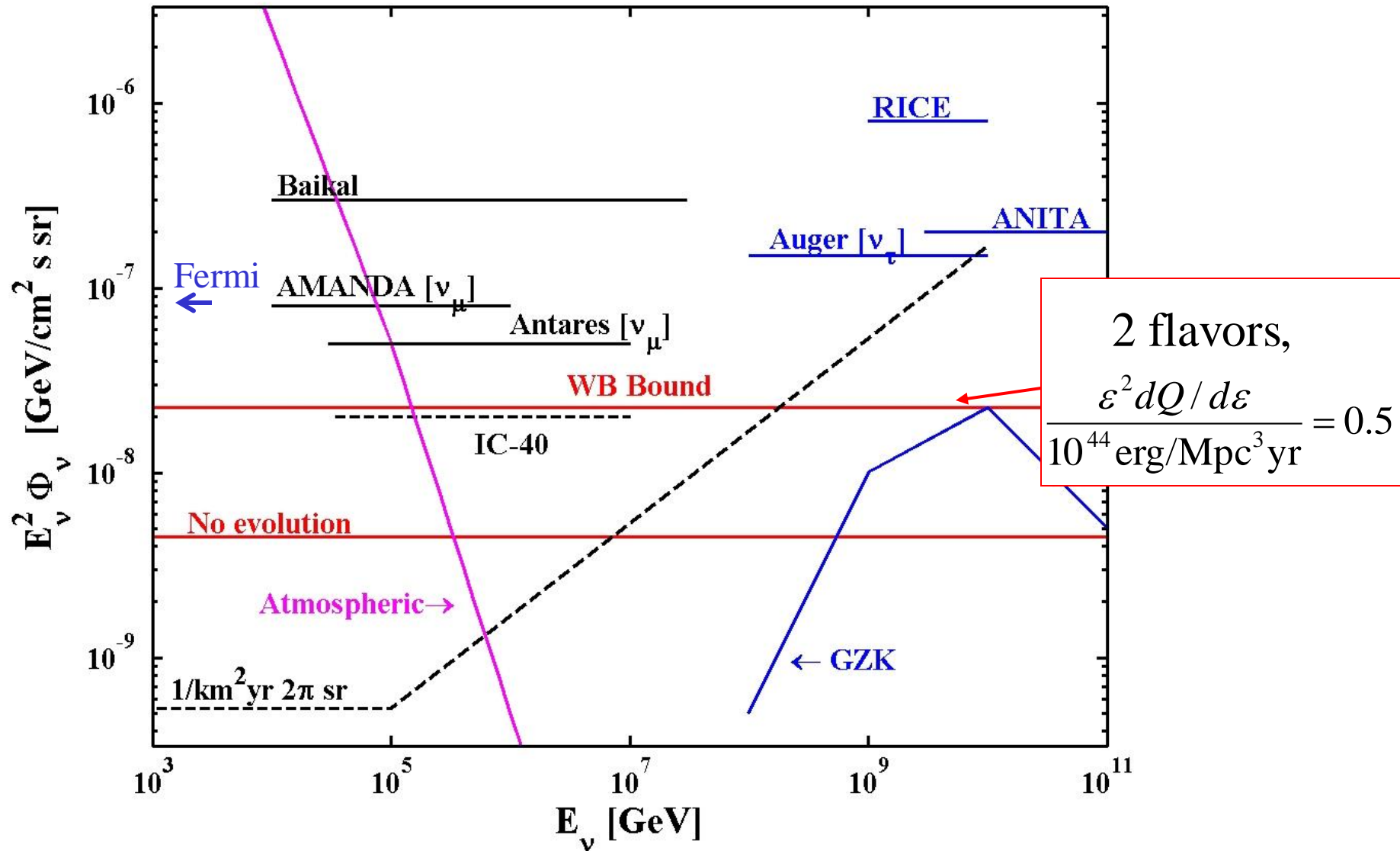
# Bound implications: I. AGN $\nu$ models

“Hidden” ( $\nu$  only)  
sources

Violating UHECR  
bound



# Bound implications: $\nu$ experiments

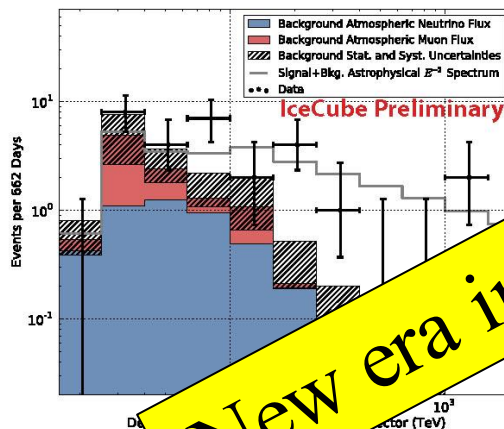


# IceCube (preliminary) detection

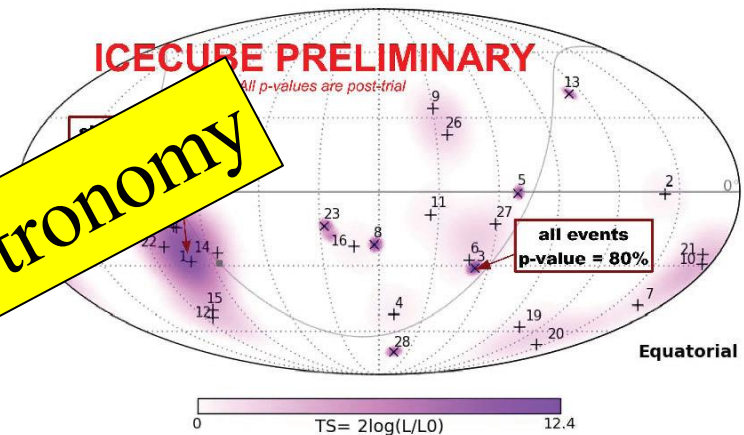
- 28 events, compared to 12 expected, above 50 TeV;  $\sim 4\sigma$  (cutoff at 2 PeV?)
- $1/E^2$  spectrum,  $4 \times 10^{-8} \text{ GeV/cm}^2 \text{ s sr}$
- Consistent with  $\nu_e:\nu_\mu:\nu_\tau=1:1:1$
- Consistent with isotropy

Energy Spectrum

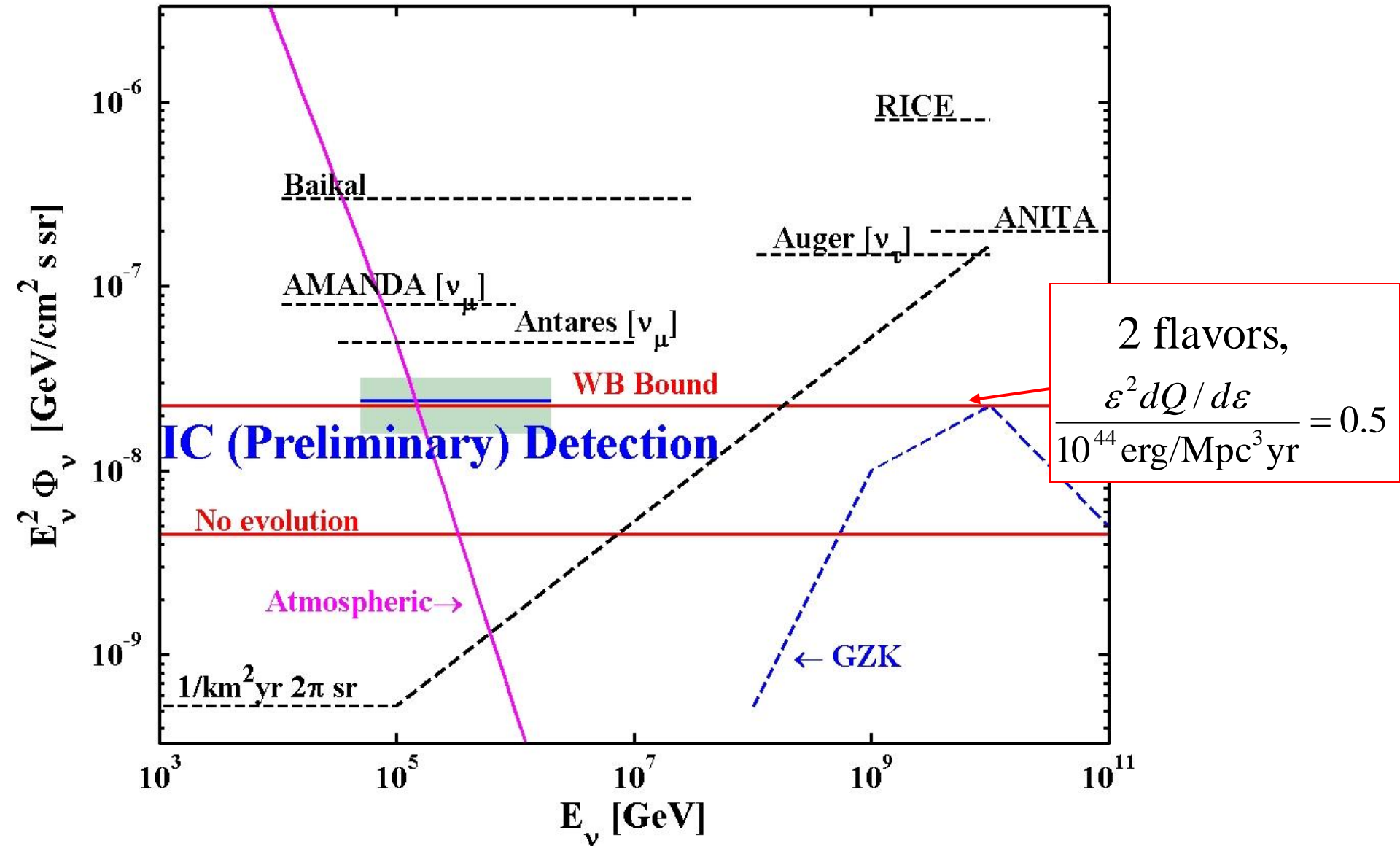
- Harder than any expected atmospheric background
- Merges well into expected backgrounds at low energies
- Potential cutoff at  $1.6^{+1.5}_{-0.4} \text{ PeV}$



Skymap: No Significant Clustering



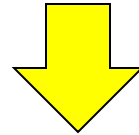
# IceCube (preliminary) detection



# IceCube's detection: Some implications

The coincidence of  $50\text{TeV} < E < 2\text{PeV}$   $\nu$  flux, spectrum (& flavor) with the WB bound is unlikely a chance coincidence.

- Unlikely Galactic:  $\varepsilon^2 \Phi_\gamma \sim 10^{-7} (E_{0.1\text{TeV}})^{-0.7} \text{GeV/cm}^2 \text{s sr}$  [Fermi]  
 $\rightarrow \sim 10^{-9} (E_{0.1\text{PeV}})^{-0.7} \text{GeV/cm}^2 \text{s sr}$



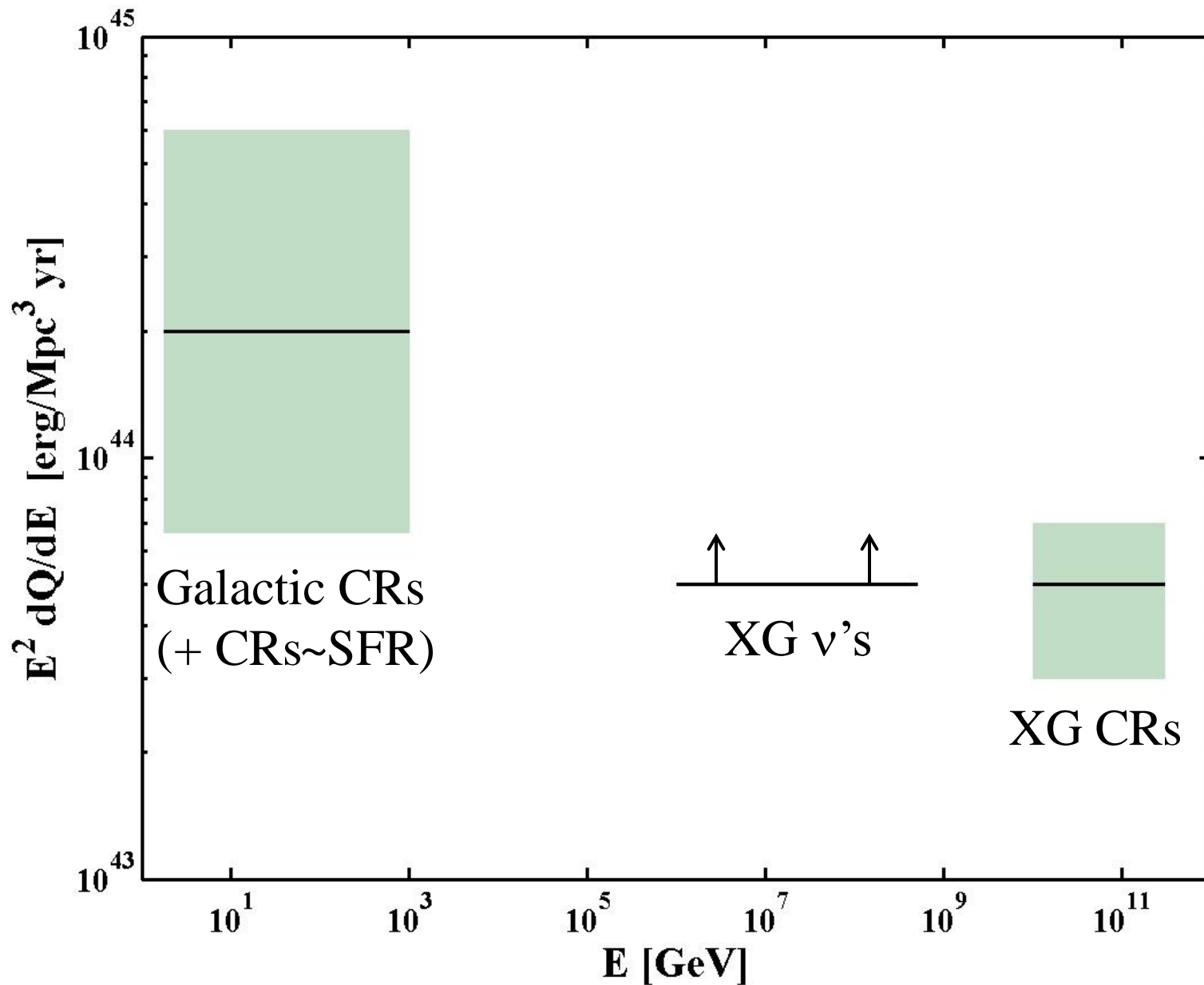
- XG distribution of sources,  
 $p, \varepsilon^2 (dQ/d\varepsilon)_{\text{PeV-EeV}} \sim \varepsilon^2 (dQ/d\varepsilon)_{>10\text{EeV}}, \tau_{\gamma p(pp)} > \sim 1$

Or:

$$\varepsilon^2 (dQ/d\varepsilon)_{\text{PeV-EeV}} \gg \varepsilon^2 (dQ/d\varepsilon)_{>10\text{EeV}}, \tau_{\gamma p(pp)} \ll 1$$

& Coincidence (over a wide energy range)

# The cosmic ray generation spectrum



# $\nu$ 's from Star Bursts

- Starburst galaxies
  - {Star formation rate, density, B}  $\sim 10^3 \times$  Milky way.  
Most stars formed in  $z > 1.5$  star bursts.
- CR e's lose all energy to synchrotron radiation,
- $\epsilon < 10 \text{ PeV}$  p's likely lose all energy to  $\pi$  production,  
at higher  $\epsilon$  may escape.

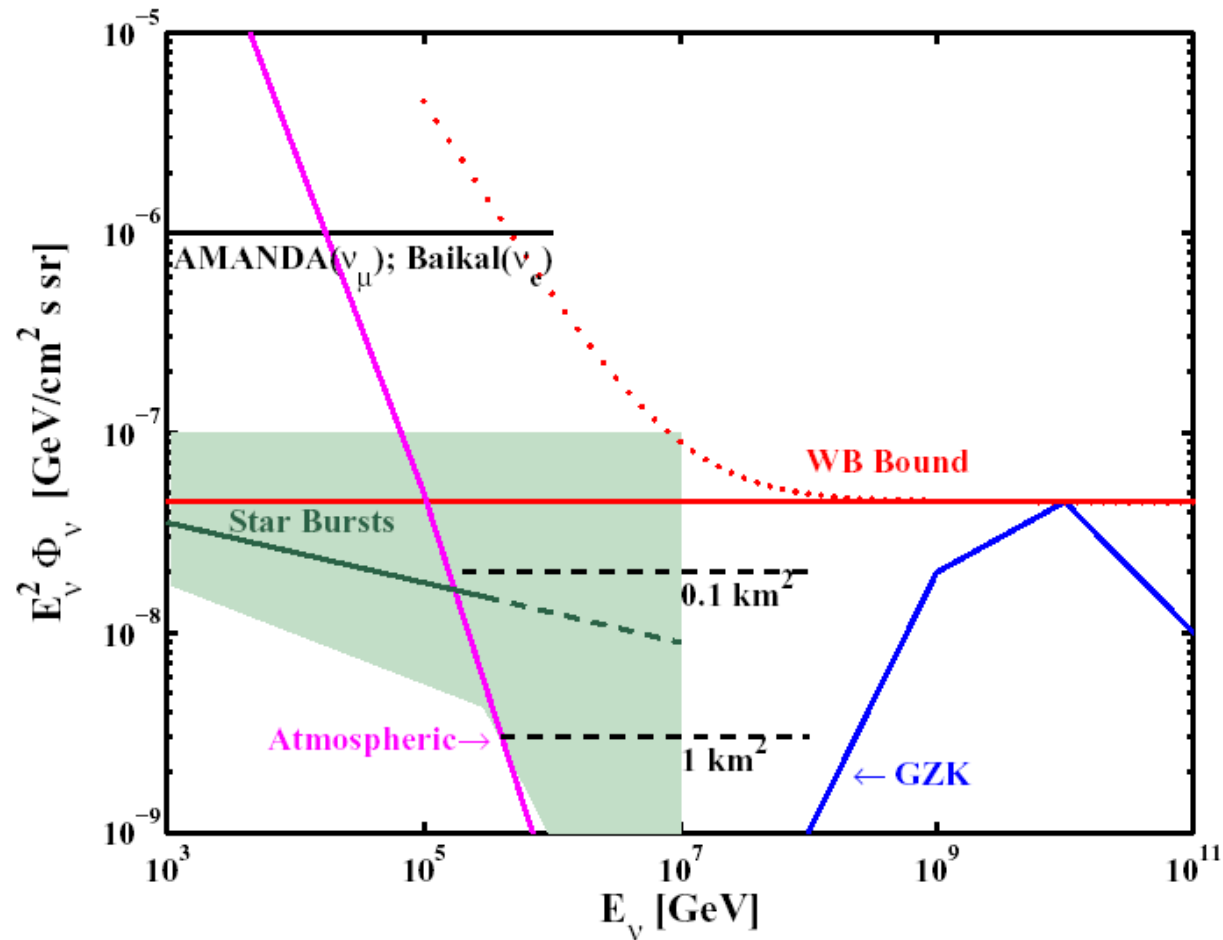
[Quataert et al. 06]

[Loeb & EW 06]

$$\epsilon^2(dQ/d\epsilon) \sim 10^{44} \text{ erg/Mpc}^3 \text{ yr} \rightarrow \Phi_\nu \sim \Phi_{\text{WB}}$$



# Starburst galaxies: predicted $\nu$ emission



[Loeb & EW 06]

# IceCube's detection: Some implications

- The identity of the CR source(s) is still debated.
- Open Q's RE candidate source(s) physics [accreting BHs].
- $\varepsilon^2(dQ/d\varepsilon) \sim 10^{44} \text{erg/Mpc}^3 \text{yr}$  at all energies ( $10\text{--}10^{10} \text{ GeV}$ ).  
Suggests: CRs of all E produced in galaxies @ a rate  $\sim \text{SFR}$ ,  
by transients releasing  $\sim 10^{50.5 \pm 1.5} \text{erg}$ .
- IceCube's detection: new era in  $\nu$  astro.  
Next: spectrum, flavor,  $>1 \text{PeV}$ , GZK,  
EM association- Bright transients are the prime targets.
- Coordinated wide field EM transient monitoring- crucial.
- EM Association may resolve outstanding puzzles:
  - Identify CR (UHE & G-CR) sources,
  - Resolve open “cosmic-accelerator” physics Q's  
(related to BH-jet systems, particle acc., rad. mechanisms),
  - Constrain  $\nu$  physics, LI, WEP.