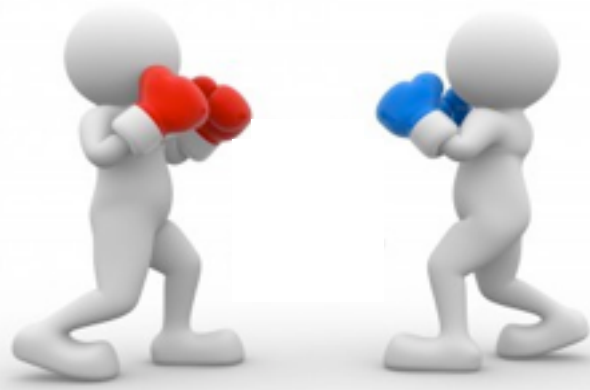


# LHC vs. Dark Matter

Roni Harnik, Fermilab



**Can LHC make a *model independent* statement about DM?**

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Yes!

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DM has not been found at LHC yet.



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**Well, duh... anything else?**

**Can LHC make a *model independent* statement about DM?**

Yes!

DM has not been found at LHC yet.

**Well, duh... anything else?**

No.

Beyond that any statement is valid within a model.

# Models

- \* Plan: a laundry list of DM models.
- \* As an organizing principle, take the “Tait axis”:



Lets take a tour  
of the axis.

# Models

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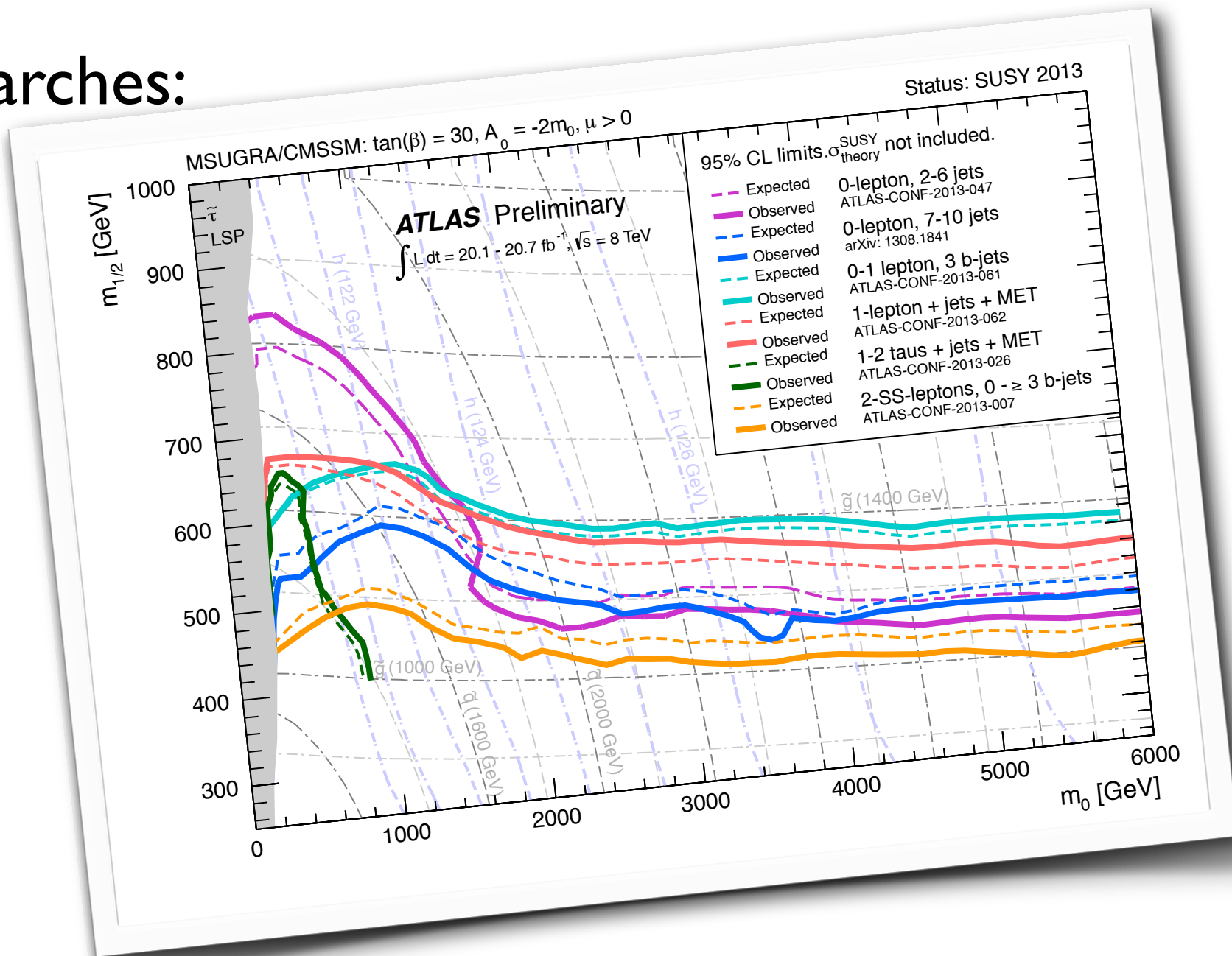
Lets take a tour  
of the axis.



# Complete Models

# SUSY Searches

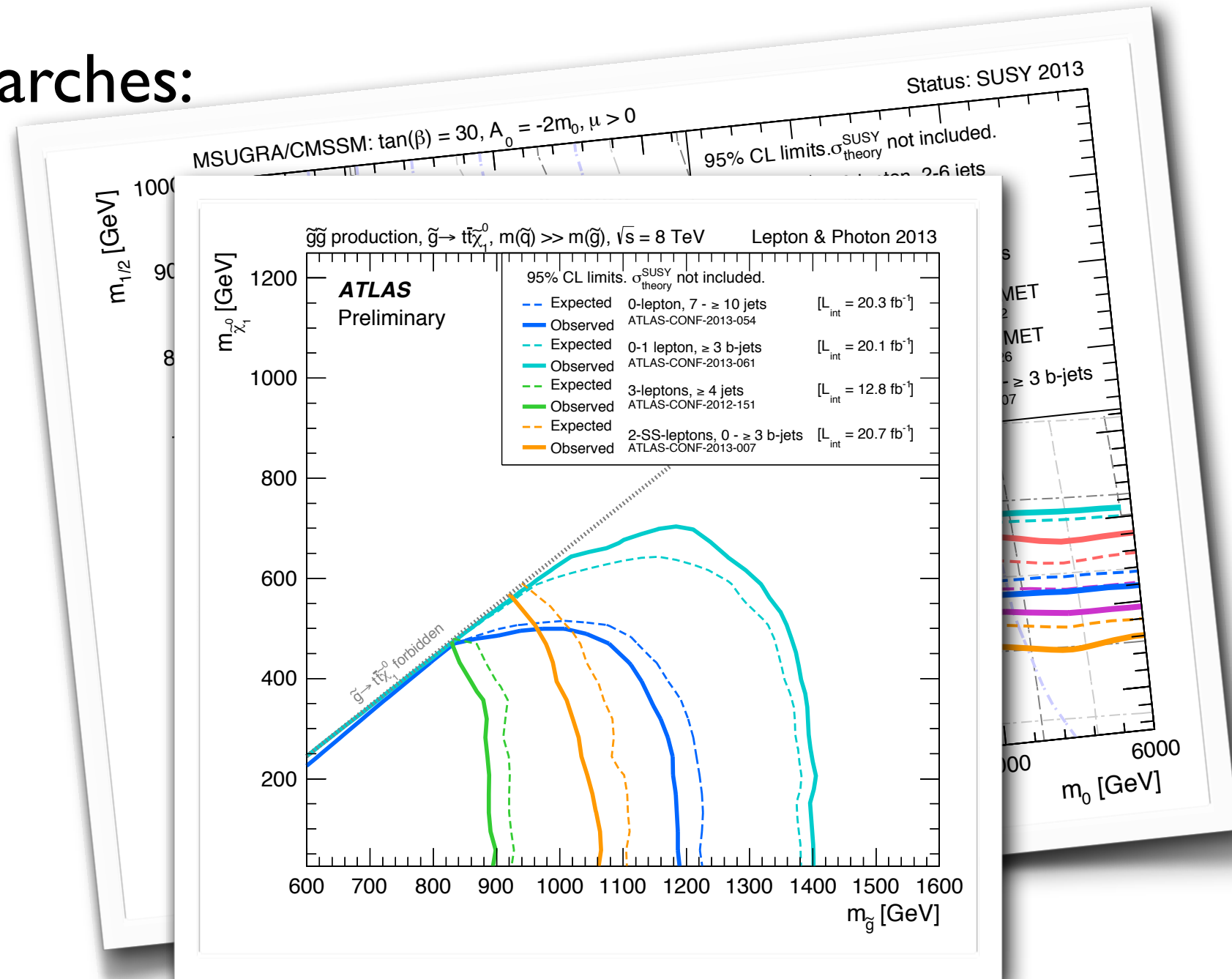
\* Many many searches:



TeV SUSY is not looking so great.  
 A statement about naturalness. not DM so much.

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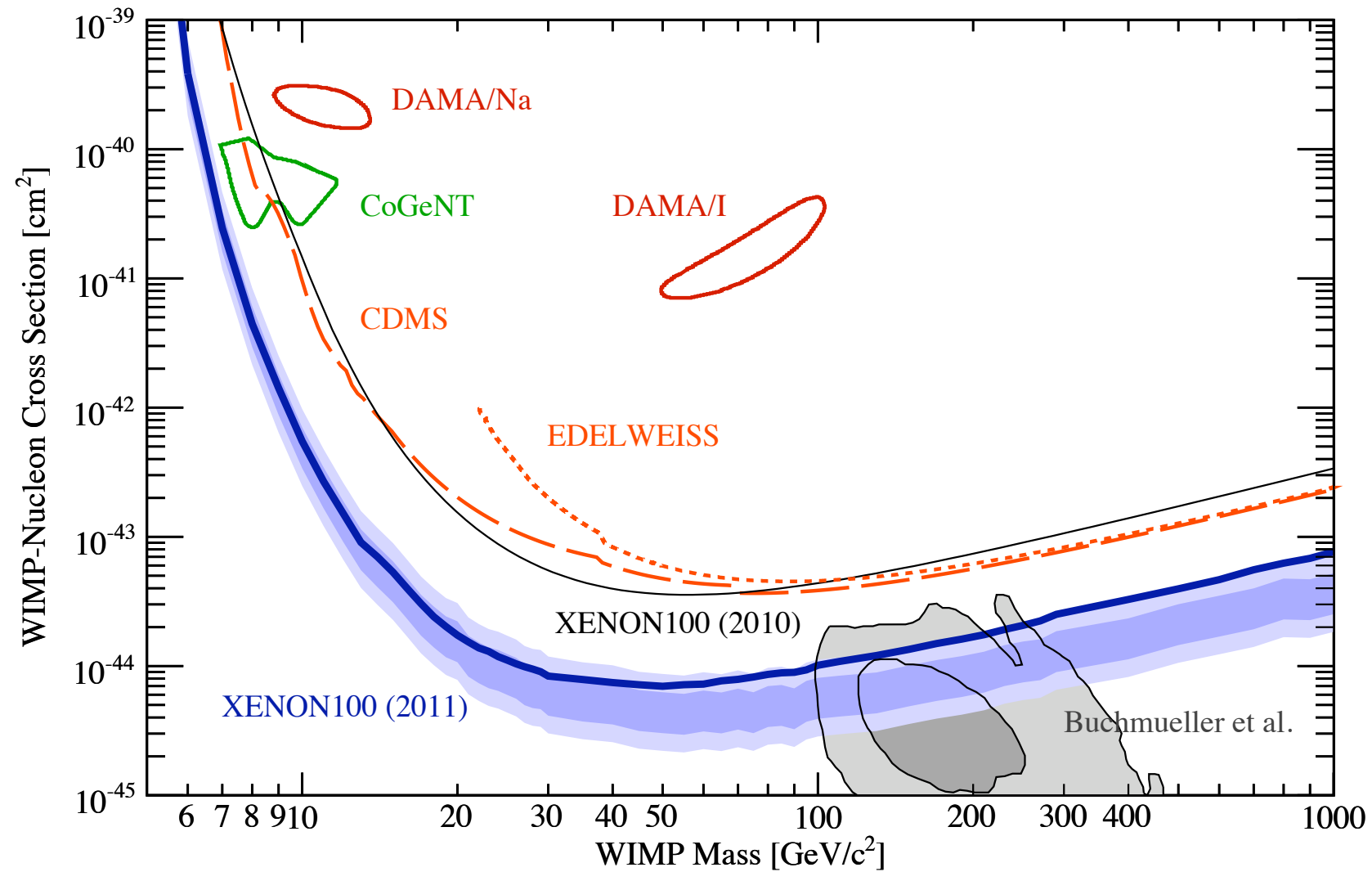


TeV SUSY is not looking so great.  
A statement about naturalness. not DM so much.

**What do these models look like  
on a DM parameter space?**

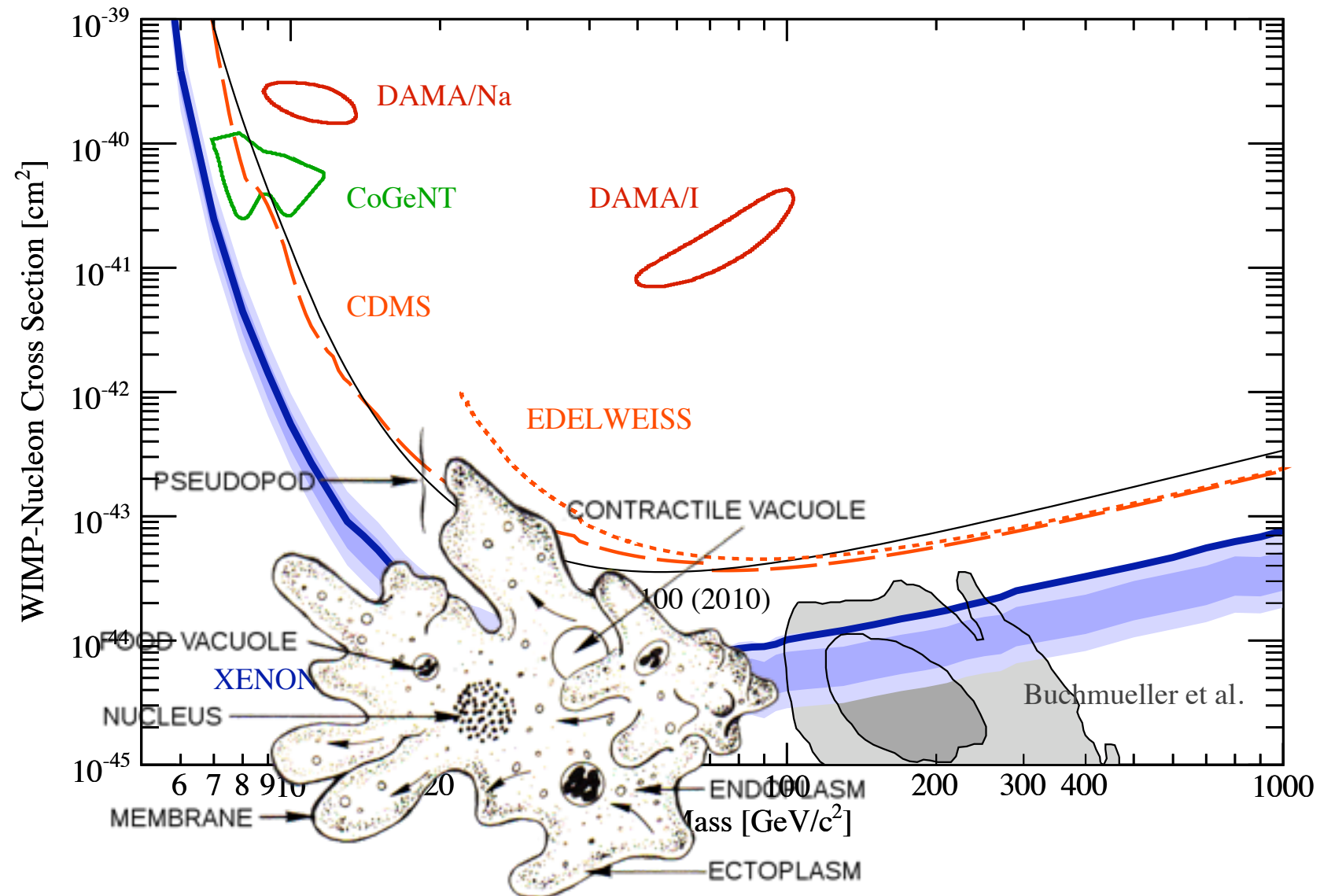
# Over-Specific?

- \* DM constraints are often presented with amoeba-shaped parameter spaces.



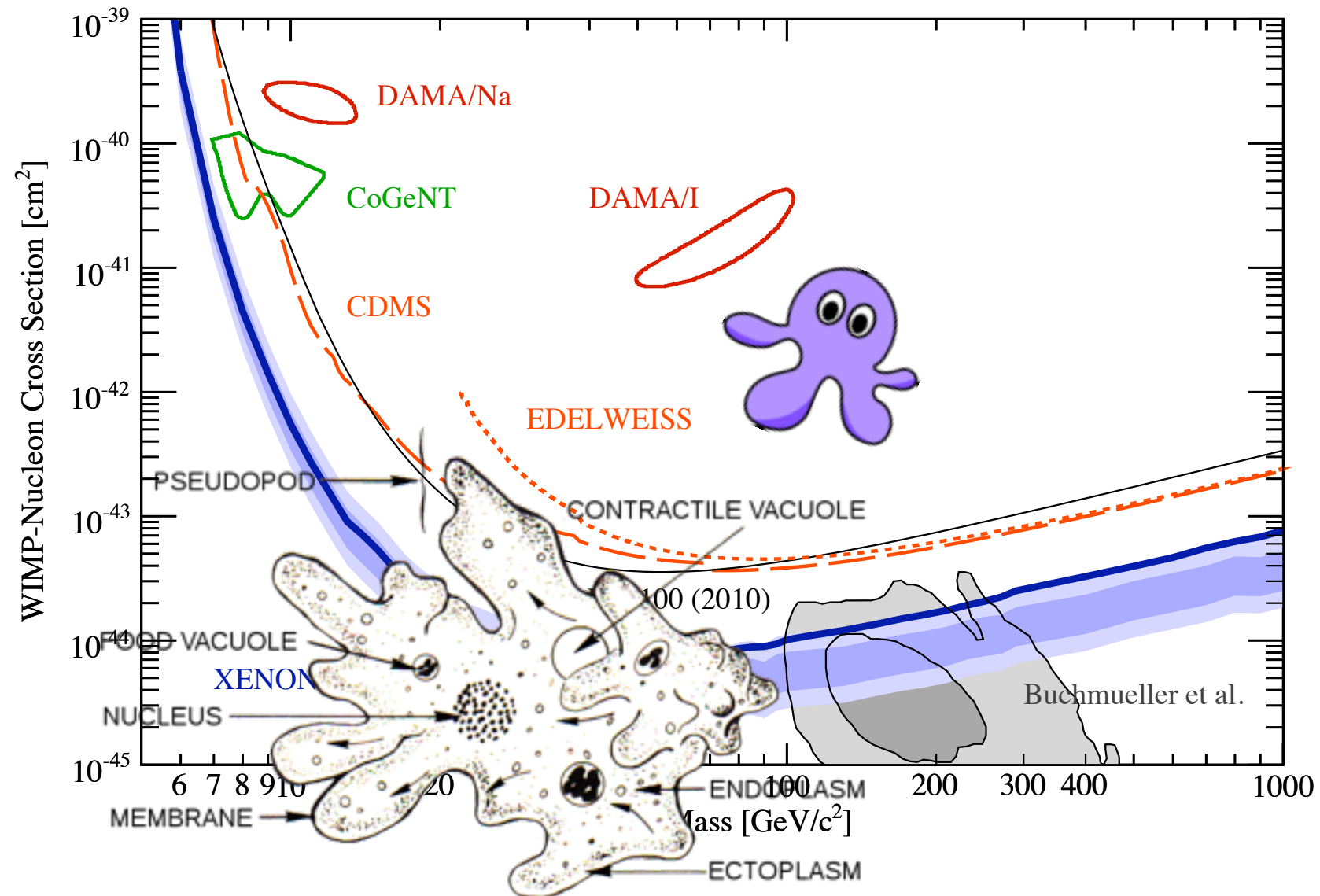
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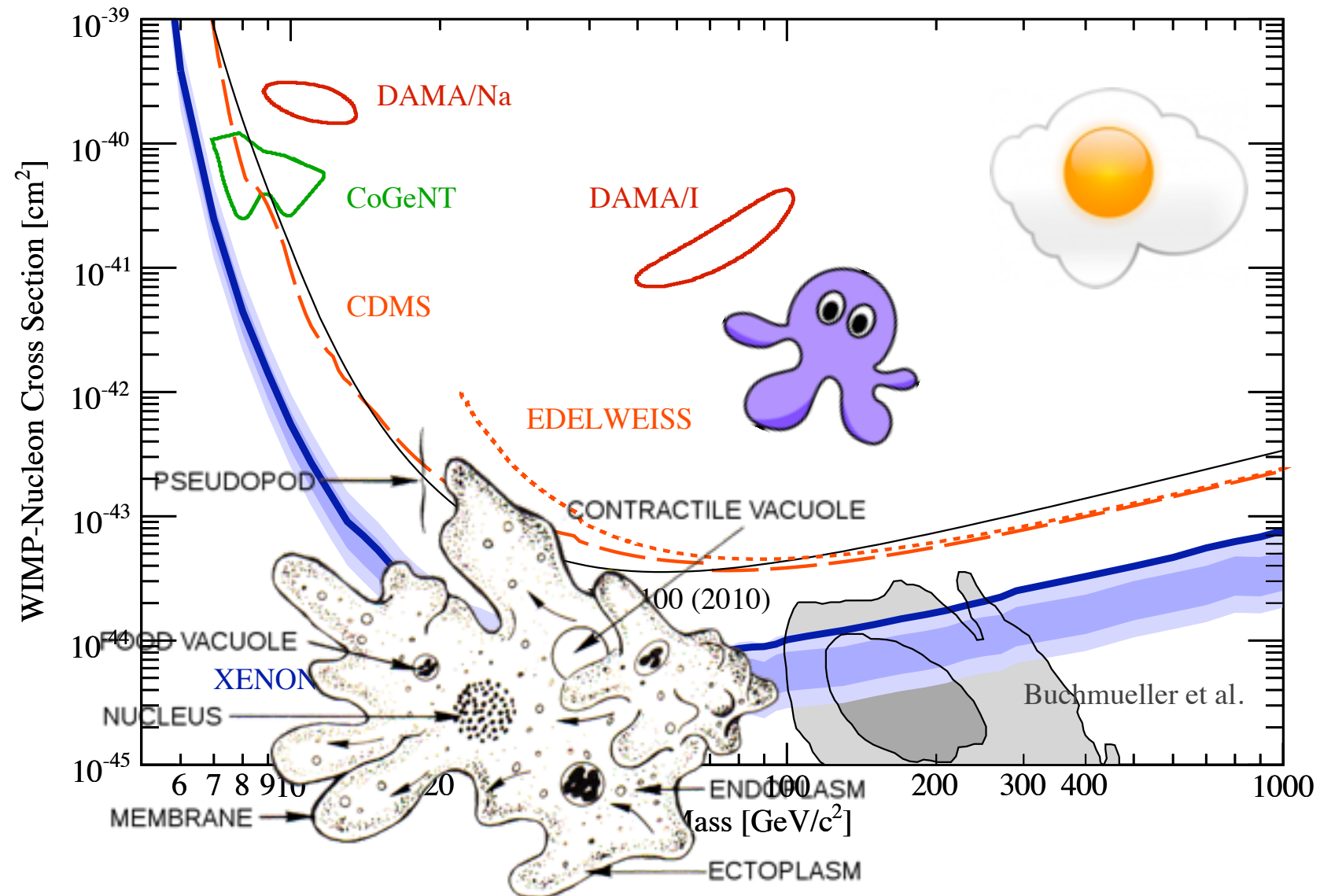
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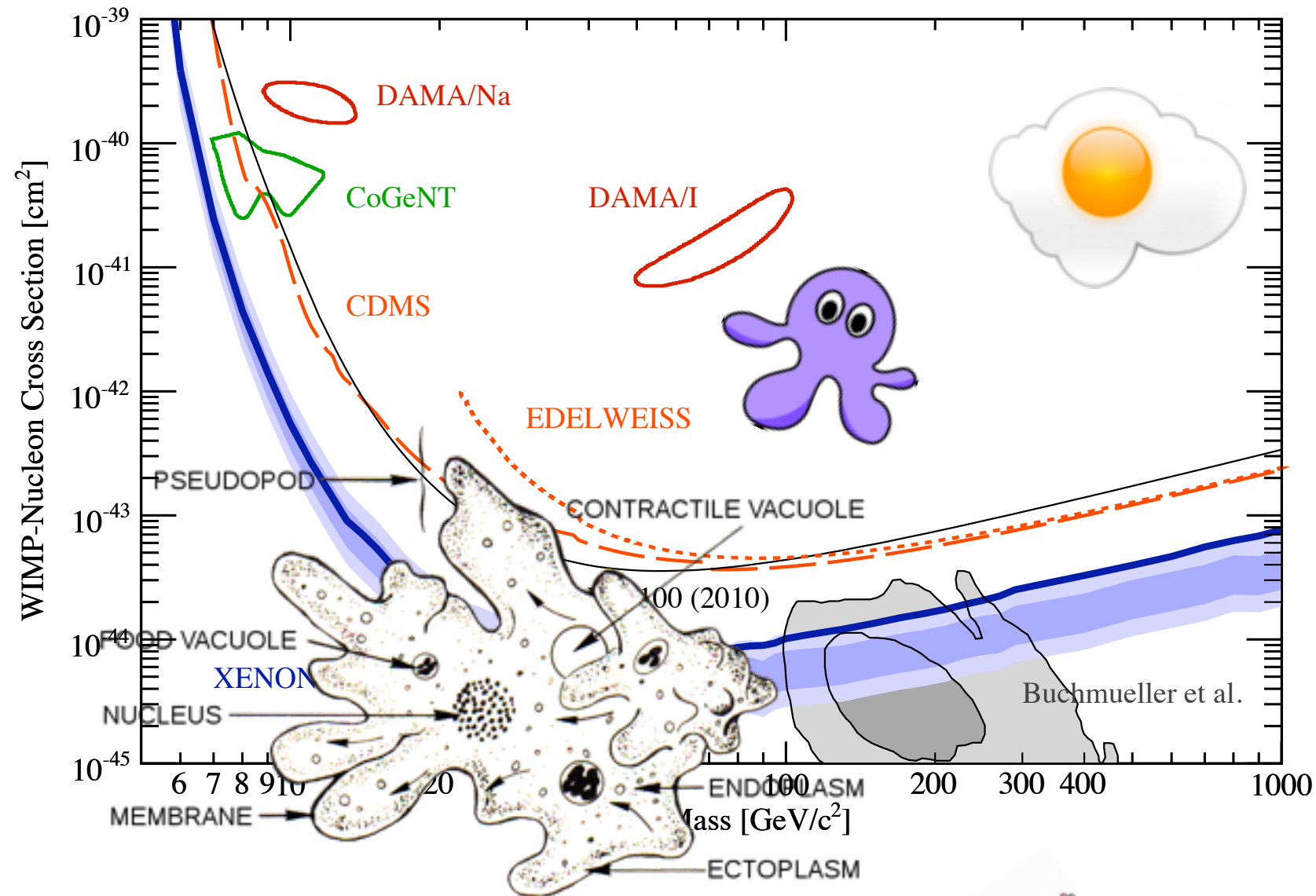
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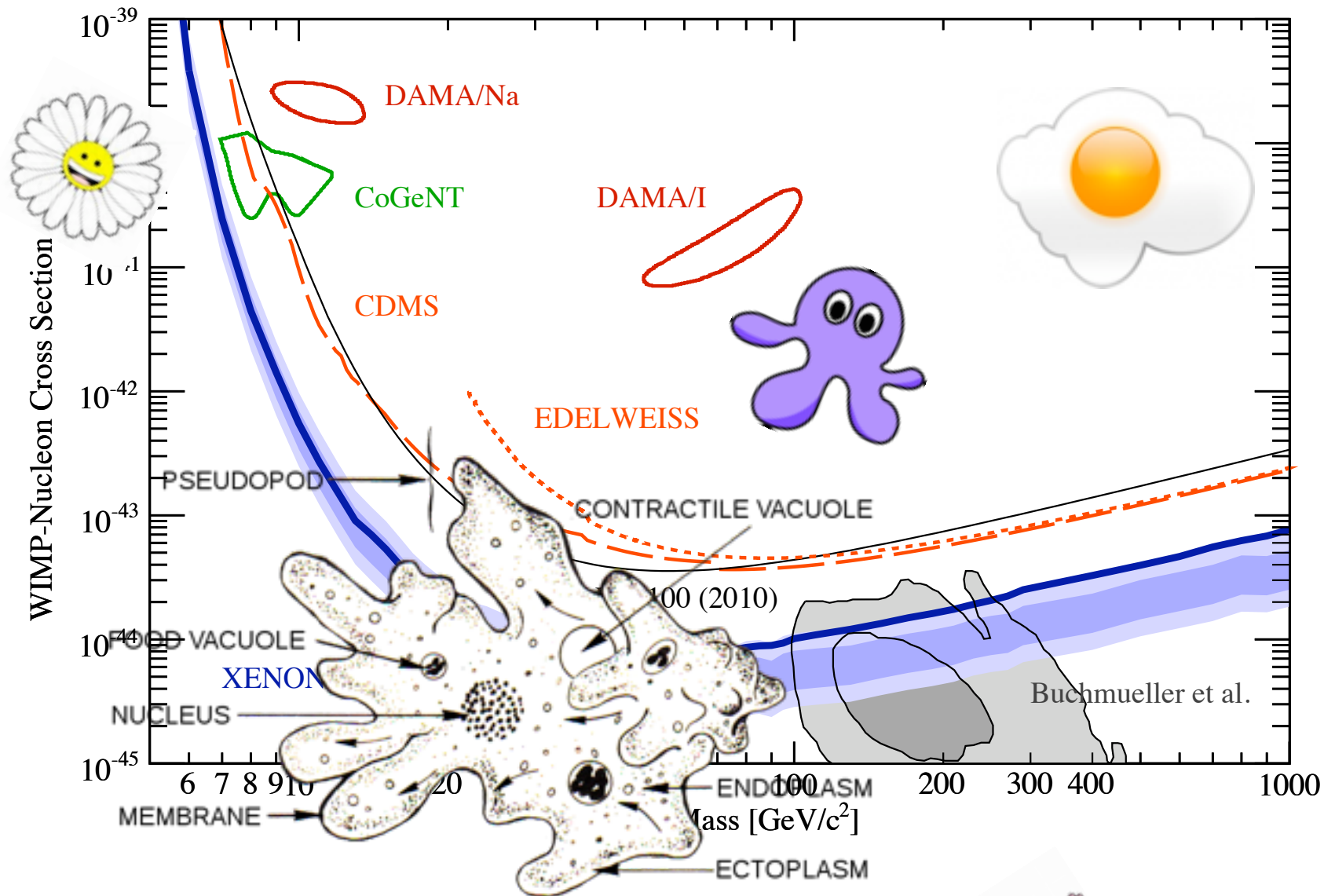
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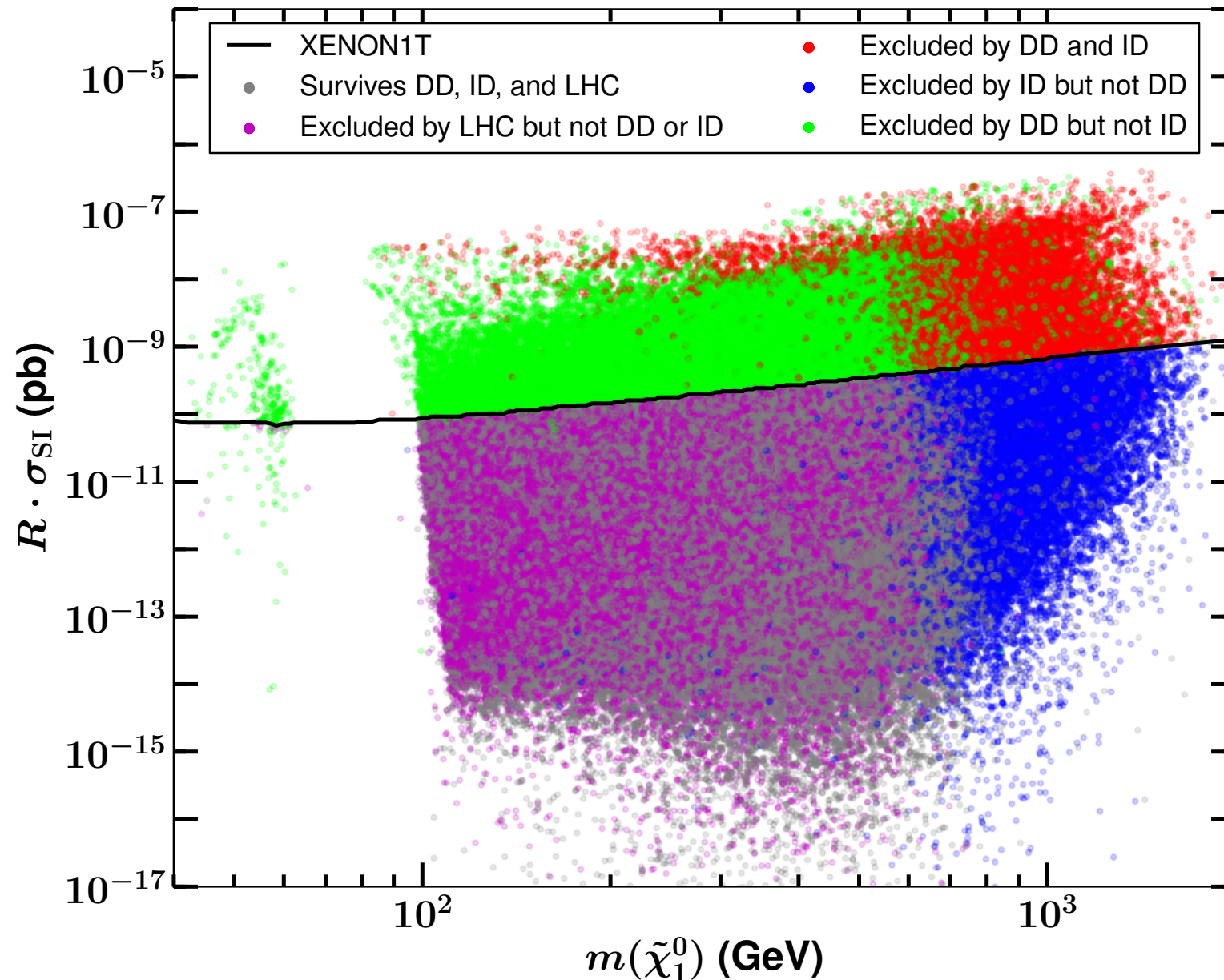
- \* DM constraints are often presented with amoeba-shaped parameter spaces.





# SUSY

Note:



pMSSM. Cotta & co. 1305.6921

\* The plane is full.

\* LHC-accessible and LHC-proof models co-exist everywhere.

\* A specific model will often have a knob to turn...

# DM EFT's

## “Fermilab Clan”:

Bai, Fox, RH - 1005.3797

Fox, RH, Kopp, Tsai - 1103.0240

Fox, RH, Kopp, Tsai - 1109.4389

Fox, RH, Primulando, Yu - 1203.1662

## “Irvine Clan”:

Goodman, Ibe, Rajaraman, Shepherd, Tait and Haibo Yu - 1005.1286

Goodman, Ibe, Rajaraman, Shepherd, Tait and Haibo Yu - 1008.1783

Fortin and Tait - 1103.3289

Rajaraman, Shepherd, Tait and Wijangco - 1108.1196

Shepherd and Goodman - 1111.2359

# Direct Detection - EFT

- \* The EFT that described DM interaction in direct detection experiments:

$$\mathcal{O}_V = \frac{(\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu q)}{\Lambda^2},$$

SI, vector exchange

$$\mathcal{O}_A = \frac{(\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu\gamma_5q)}{\Lambda^2},$$

SD, axial-vector exchange

$$\mathcal{O}_t = \frac{(\bar{\chi}P_Rq)(\bar{q}P_L\chi)}{\Lambda^2} + (L \leftrightarrow R),$$

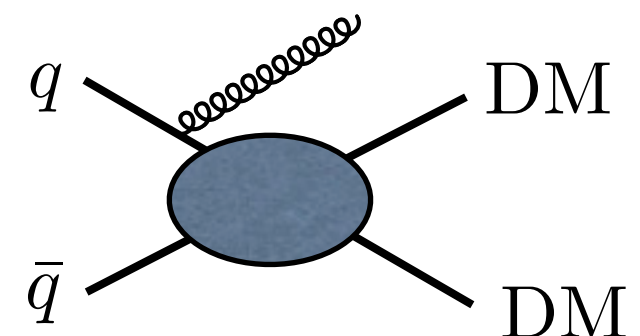
SI (or SD), t-channel  
“squark exchange”

$$\mathcal{O}_g = \alpha_s \frac{(\bar{\chi}\chi)(G_{\mu\nu}^a G^{a\mu\nu})}{\Lambda^3}$$

SI gluon operator

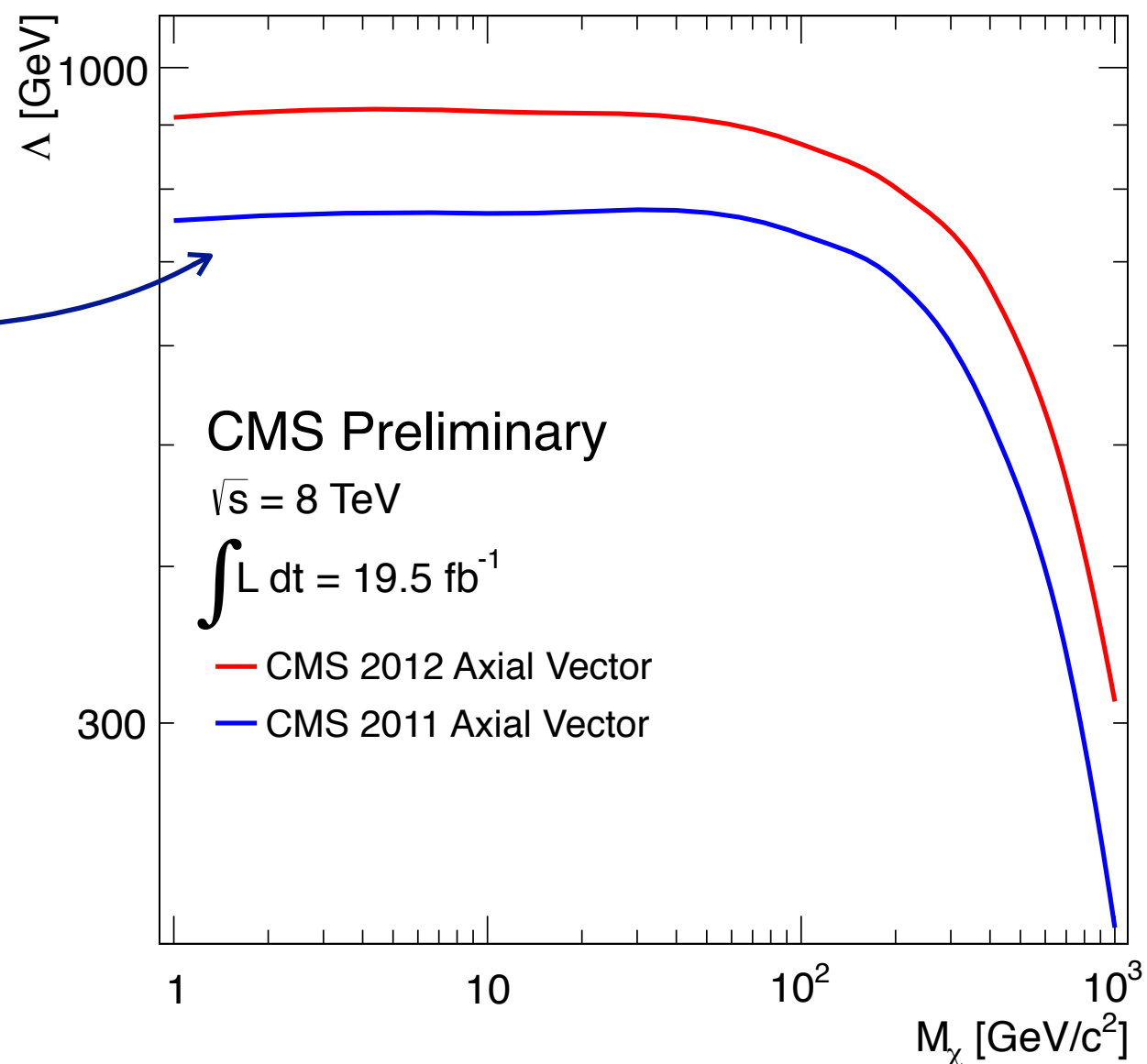
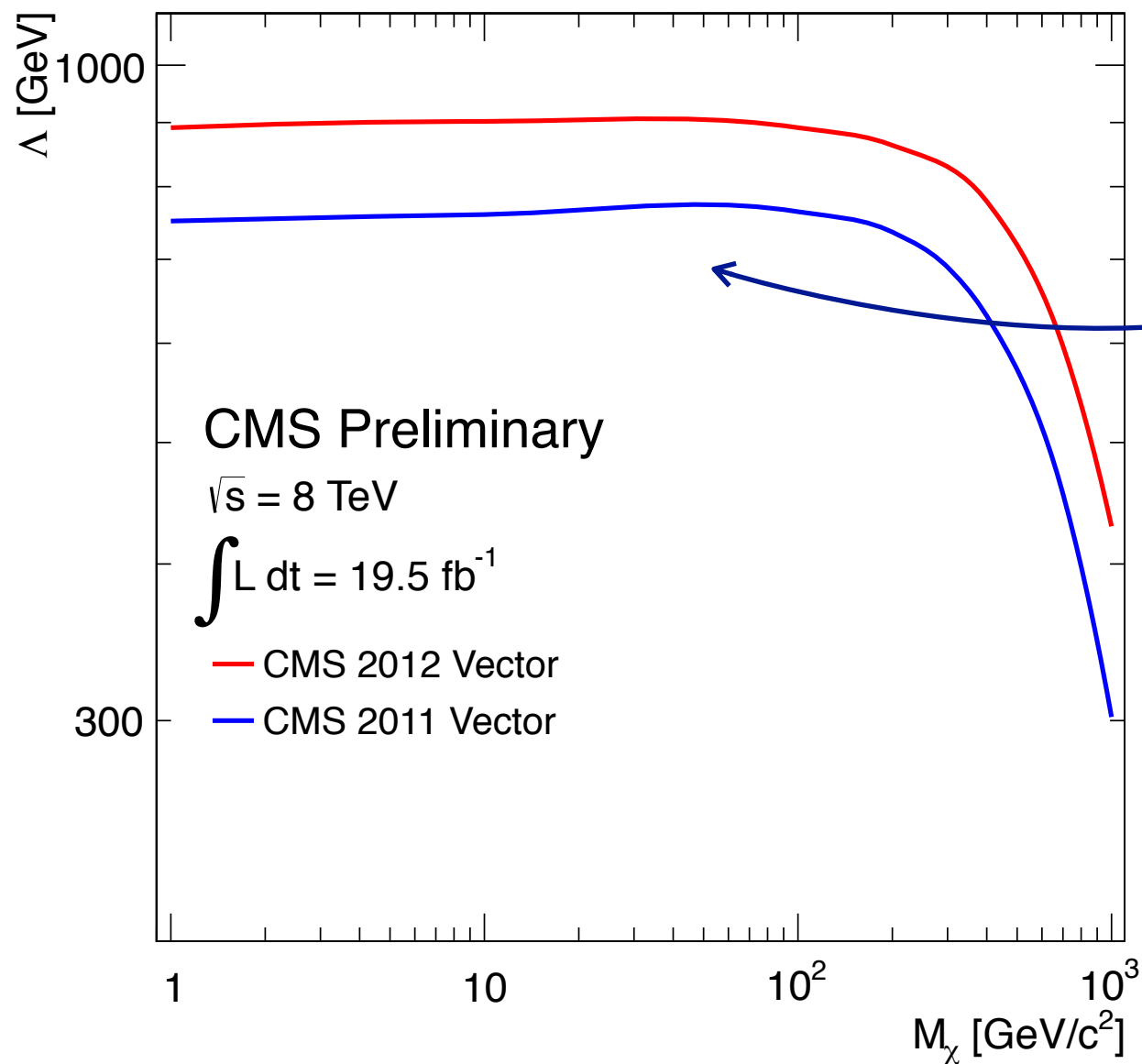
Look for mono-anything:

(mono-jet, mono-photon, mono-Z,  
mono-W, mono-di-jet...?)



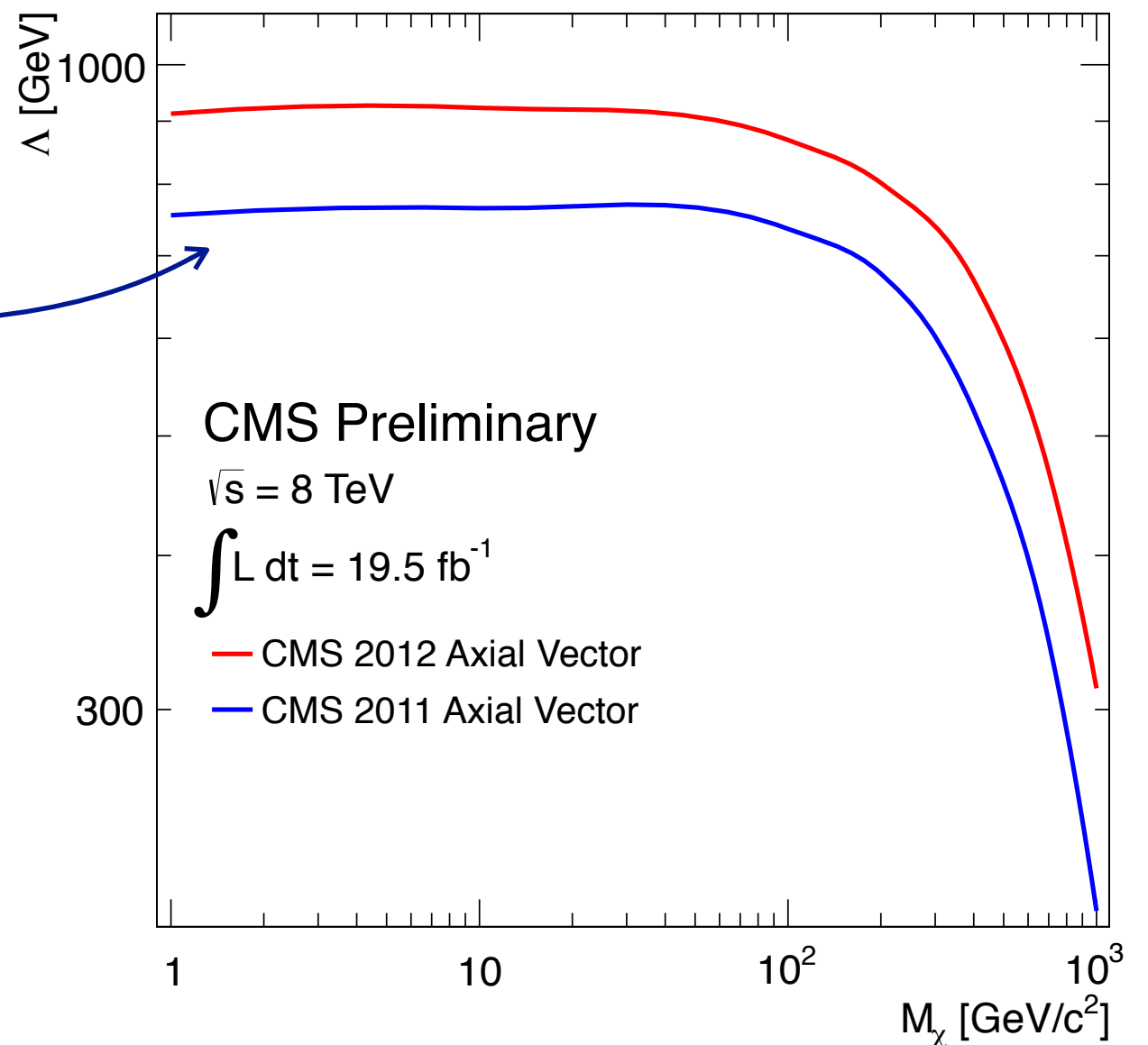
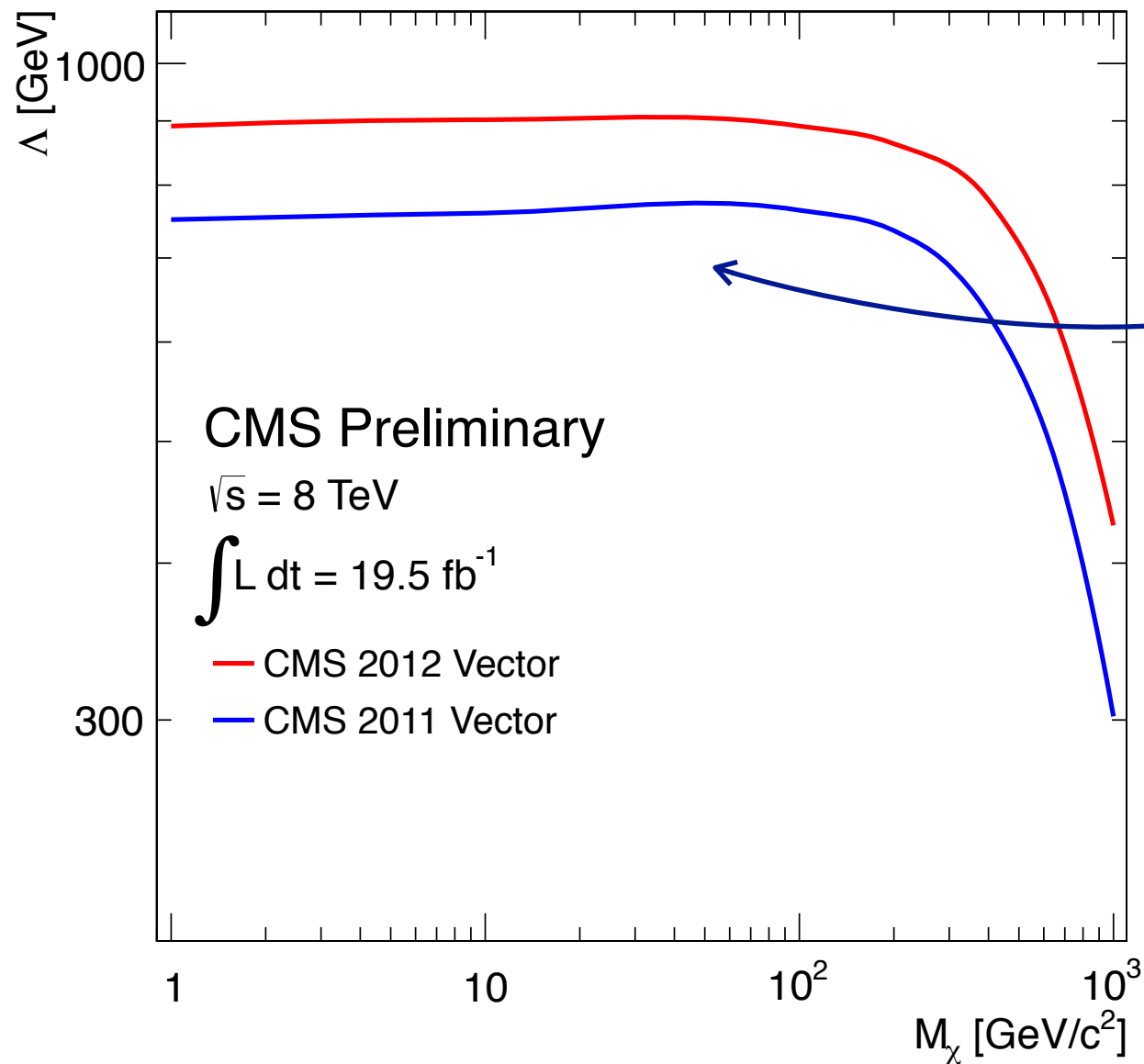
# Limits on $\Lambda \equiv \frac{M}{\sqrt{g_\chi g_1}}$ :

*same limit for SI and SD*



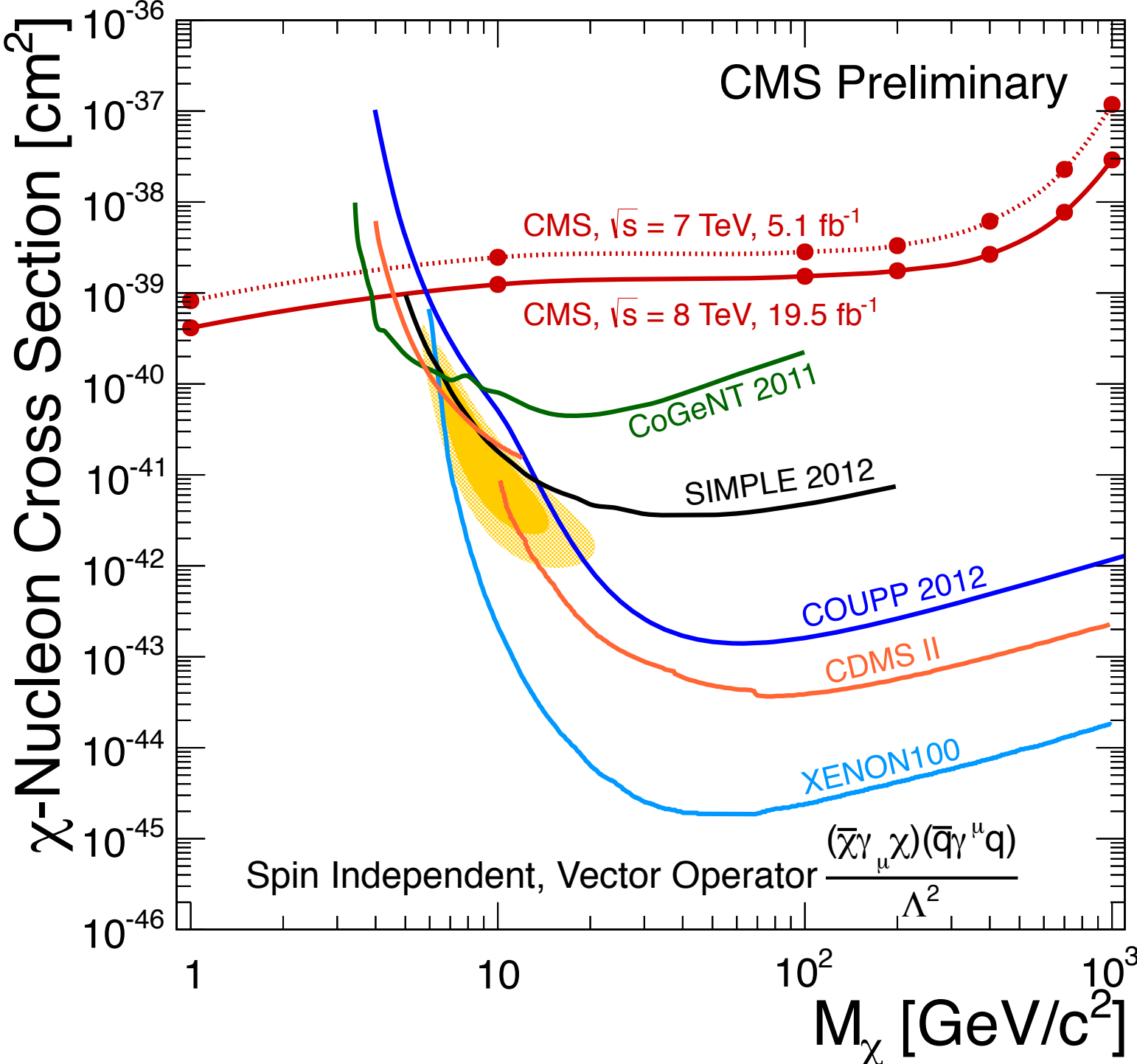
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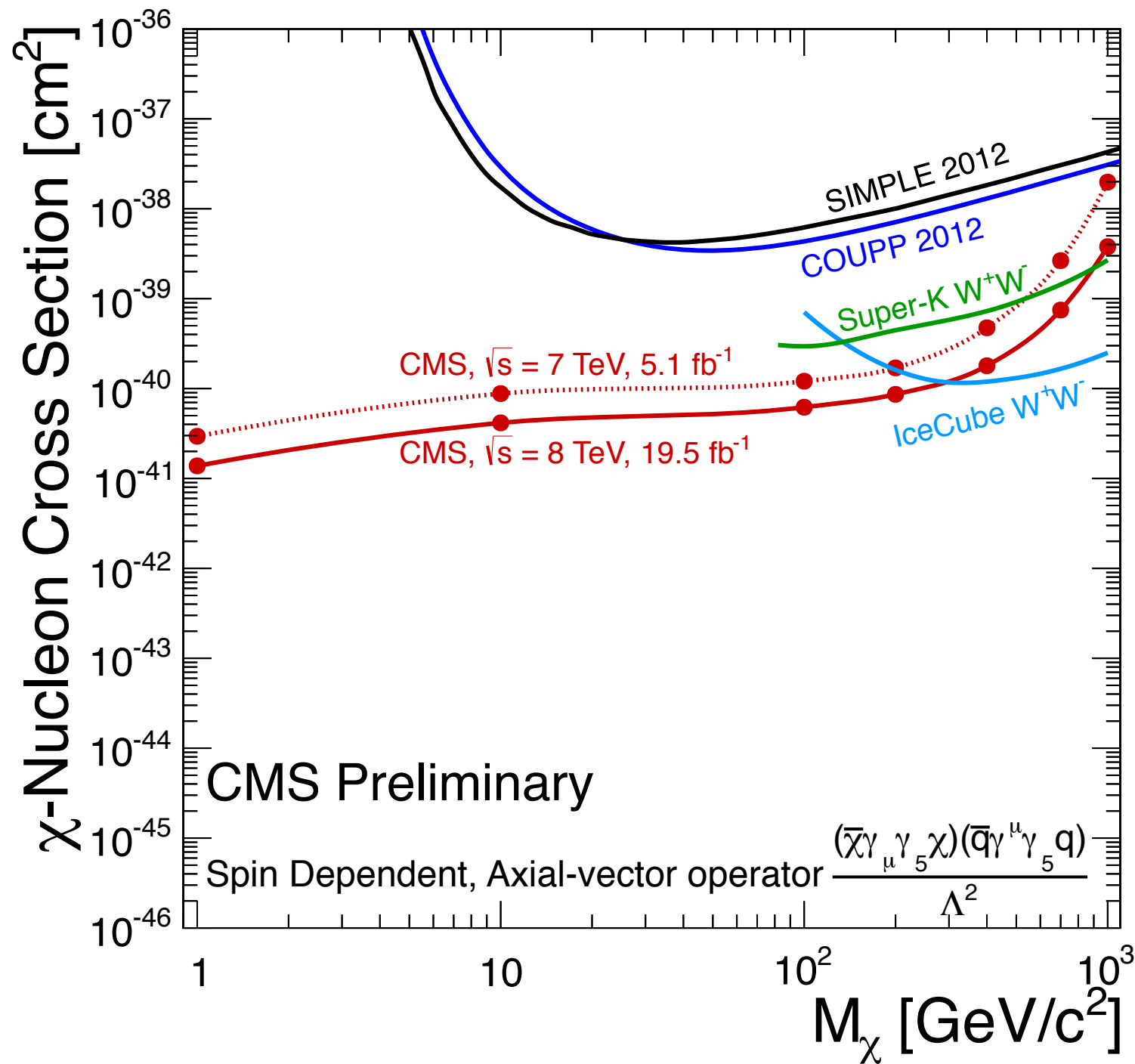
The limit is flat at low mass.

# SI Limit



**Strong limit  
at low mass**

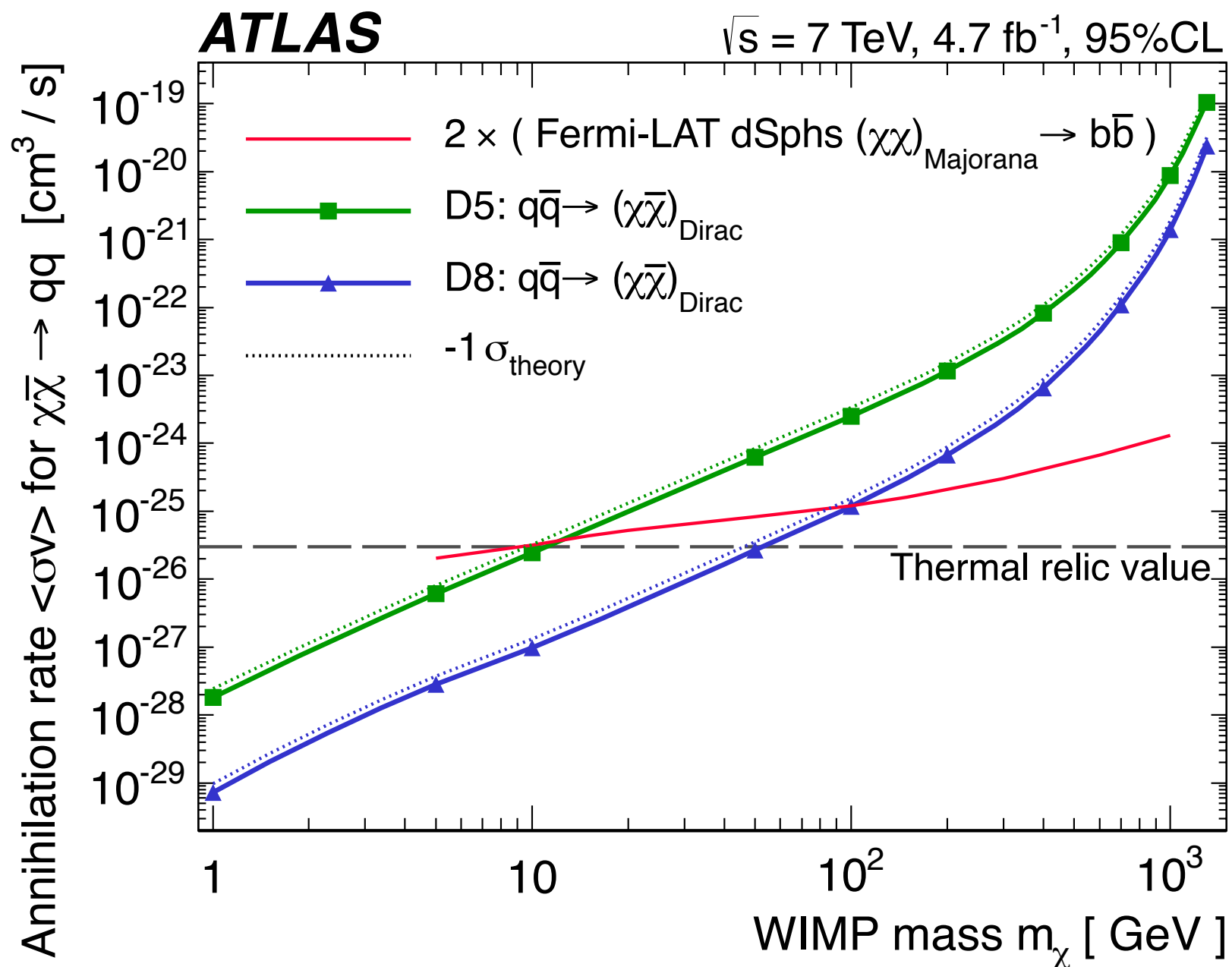
# SD Limit



**Strong spin  
dependent  
limit.**

# Annihilation

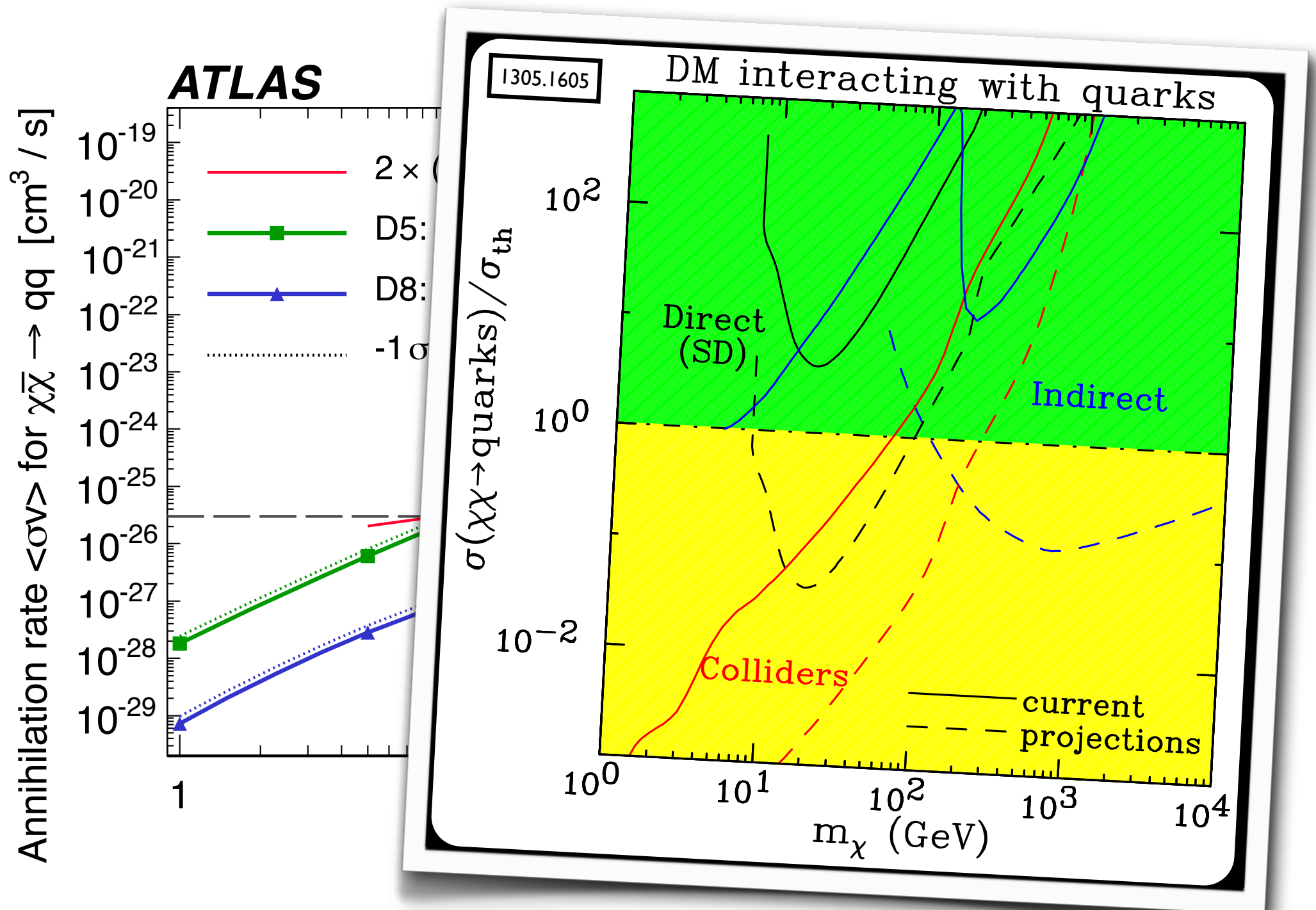
- \* A minimal light thermal relic is constrained:





# Annihilation

- \* A minimal light thermal relic is constrained:



**Is the EFT a valid description at LHC?**

Two possibilities:

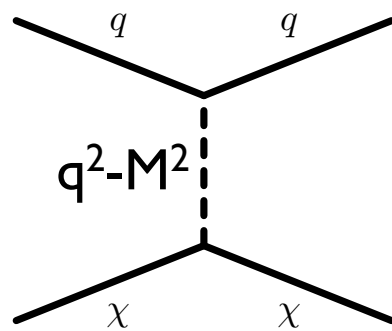
# **Is the EFT a valid description at LHC?**

Two possibilities:

- 1) EFT is valid at LHC.
- 2) It's not.

# EFT vs Light Mediator

- \* The EFT is valid for direct detection ( $q \sim 100$  MeV):

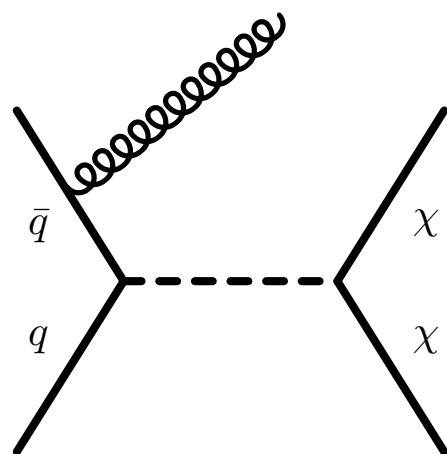


$$\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\mu = \frac{m_{\chi} m_N}{m_N + m_{\chi}}$$

$$\Lambda \equiv \frac{M}{\sqrt{g_q g_{\chi}}}$$

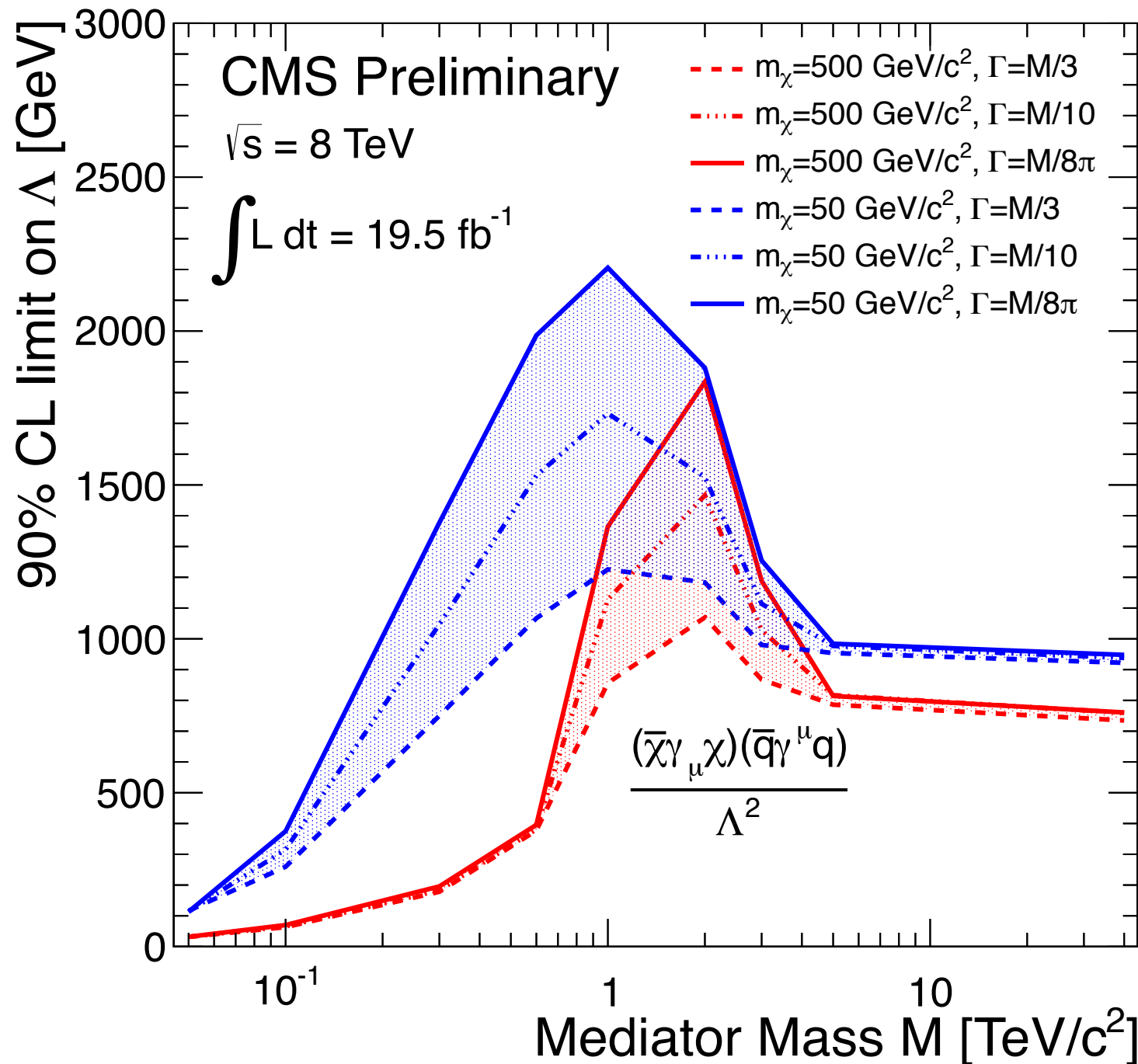
- \* At a collider consider two extreme limits:



$$\sigma_{1j} \sim \begin{cases} \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2} & M \ll \sqrt{s_*} \\ \alpha_s g_{\chi}^2 g_q^2 \frac{p_T^2}{M^4} & M > \sqrt{s_*} \end{cases}$$

# Light Mediator

\* The limit become better before it gets worse:

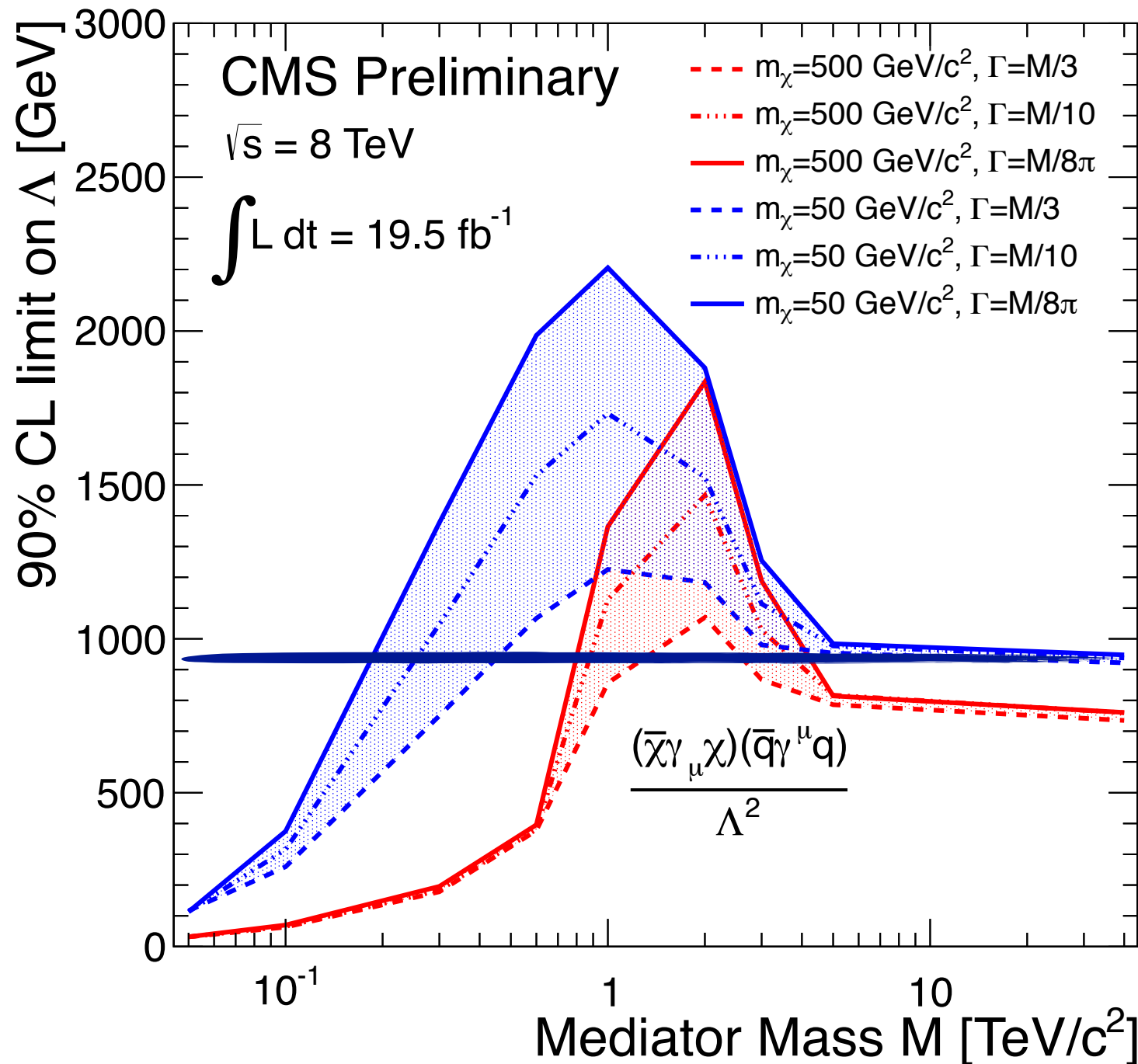


EFT limits are **conservative** so long as the mediator is above  $\sim$  twice DM mass.

They are not valid for light mediator.

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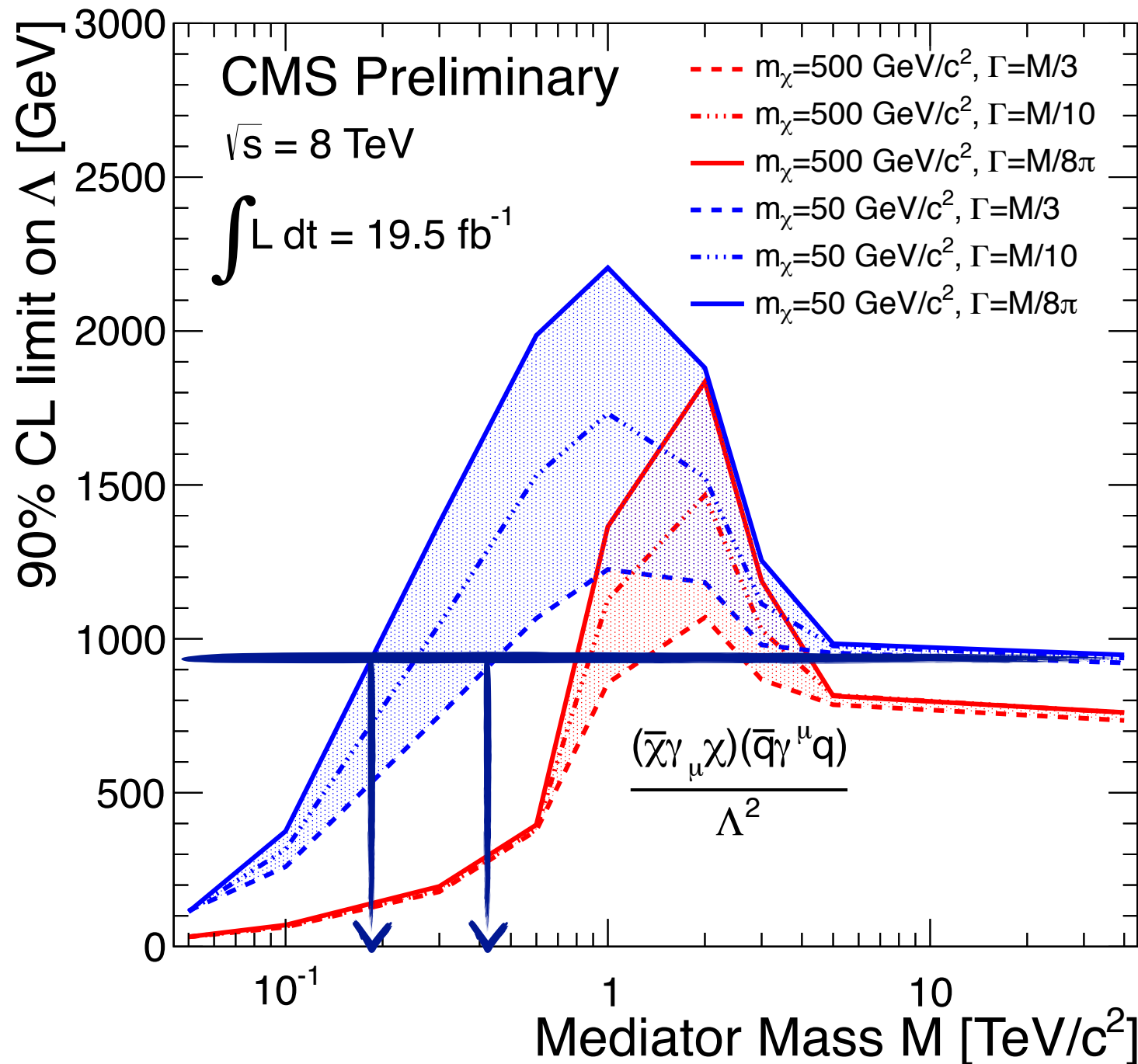


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EFT limits are **conservative** so long as the mediator is above  $\sim$  twice DM mass.

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# Simplified Model: Higgs

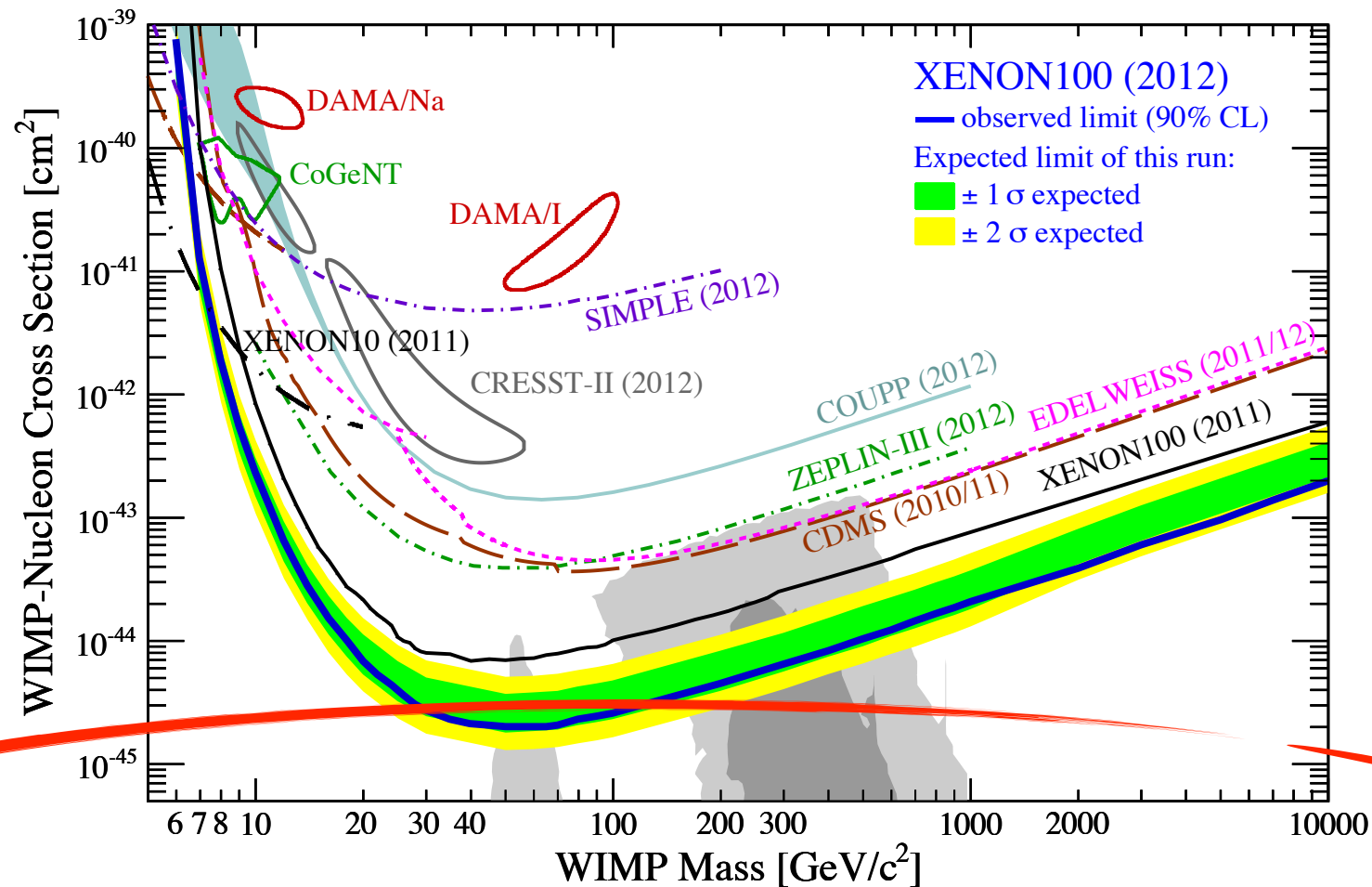
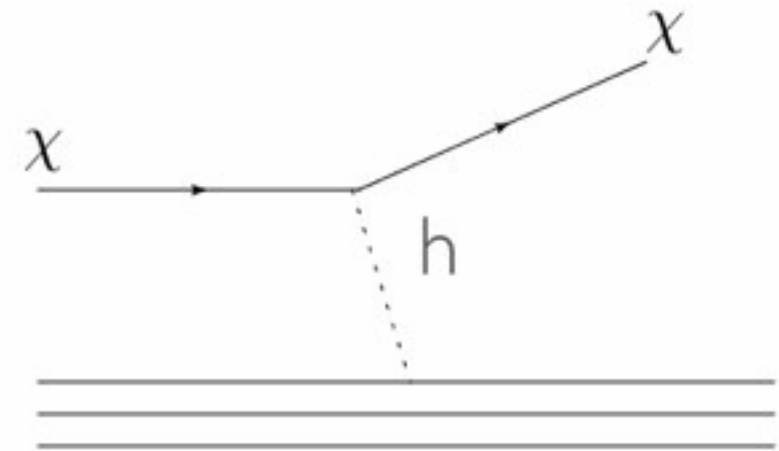
For other examples see talk  
By DiFranzo this afternoon.



# Higgs Mediation

DM exists  
+  
HIGGS exists

?

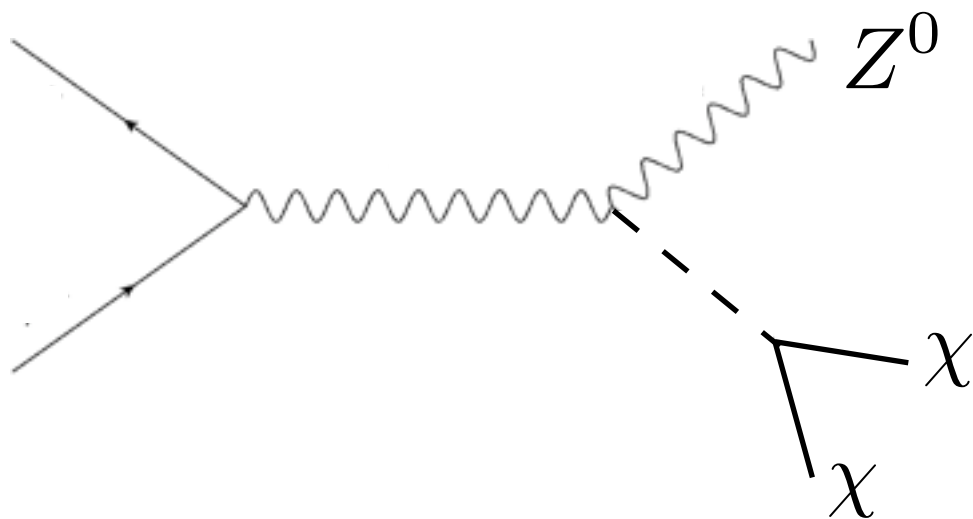


Higgs region  
is being  
PROBED!

# Higgs mediation

\* Direct detection will not cover light DM. But...

DM is light  
+  
couples to HIGGS = Invisible h decay!



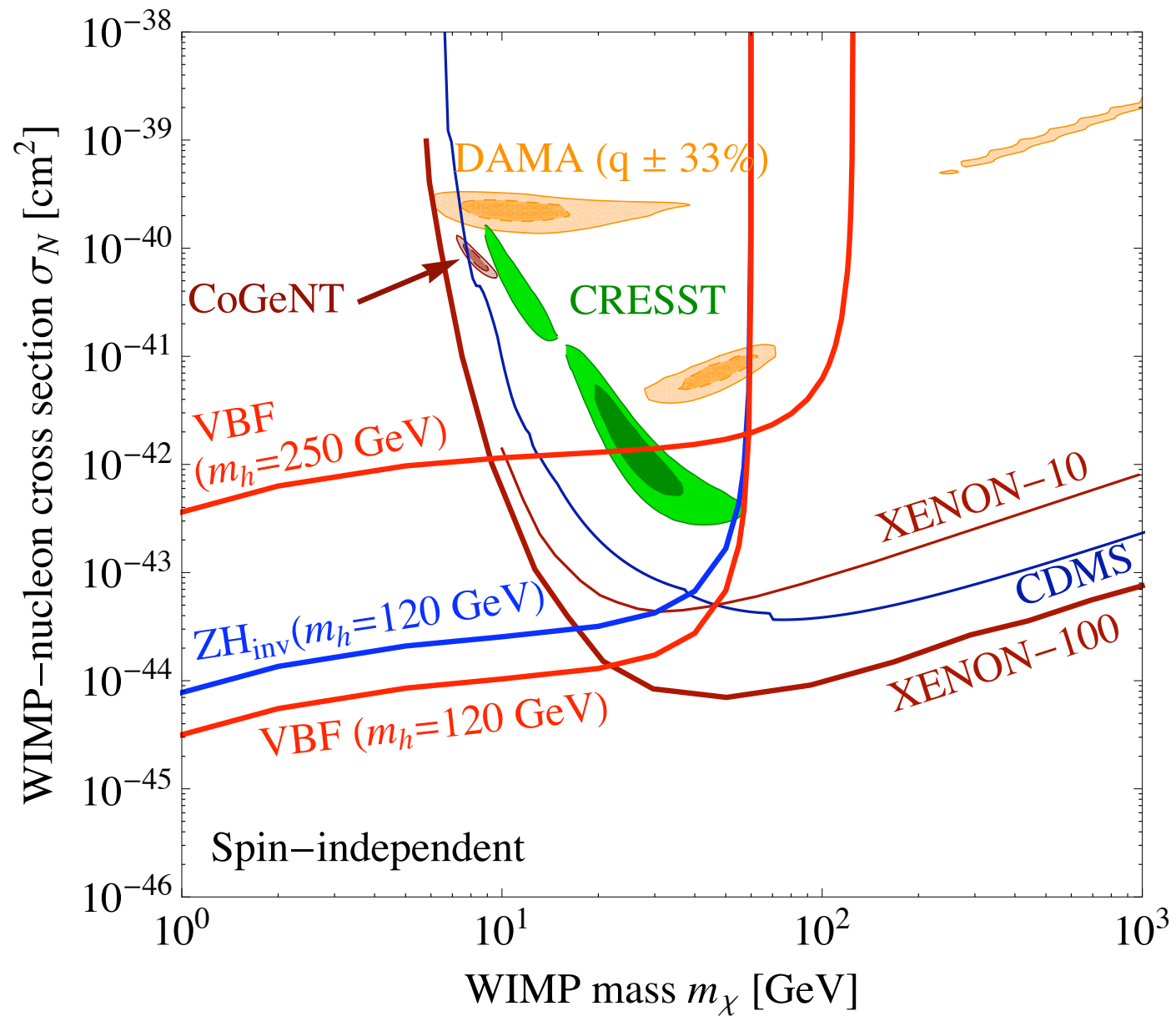
current ATLAS limit:  
 $BR(h \rightarrow inv) < 65\%$

(David's Talk)

(Higgs is narrow!  
a strong limit on  
h-DM coupling)

# Higgs Mediation

ATLAS 30 fb<sup>-1</sup> upper bound ( projected )



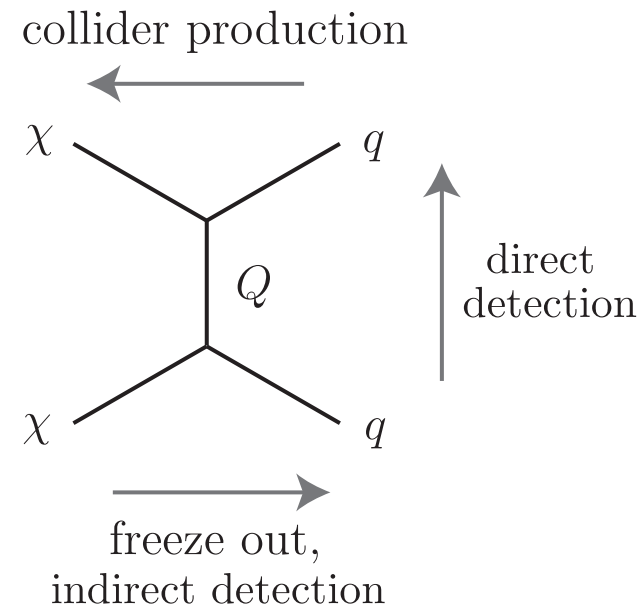
**Complementary.**

**Higgs window will be fully covered!**

# Other Simplified Models

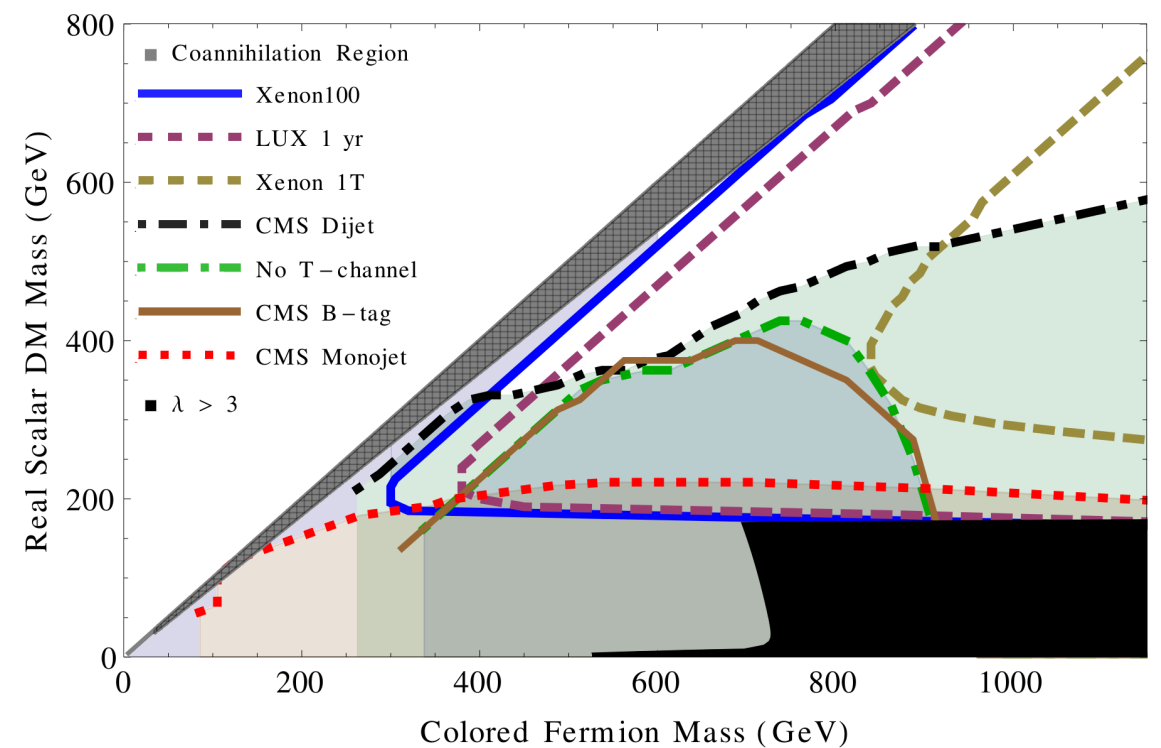
## \* “Squark mediator”:

Chang et al. 1307.8120  
 Bai and Berger 1308.0612  
 DiFranzo et al 1308.2679



## \* Small parameter space:

(e.g. can fix coupling from relic density)



# Conclusion

- \* **The LHC is also looking for DM.**
- \* The precise mapping of limits is model dependent.
- \* LHC searches are often **complementary**:
  - o Low mass.
  - o Spin dependent.
  - o Invisible Higgs.







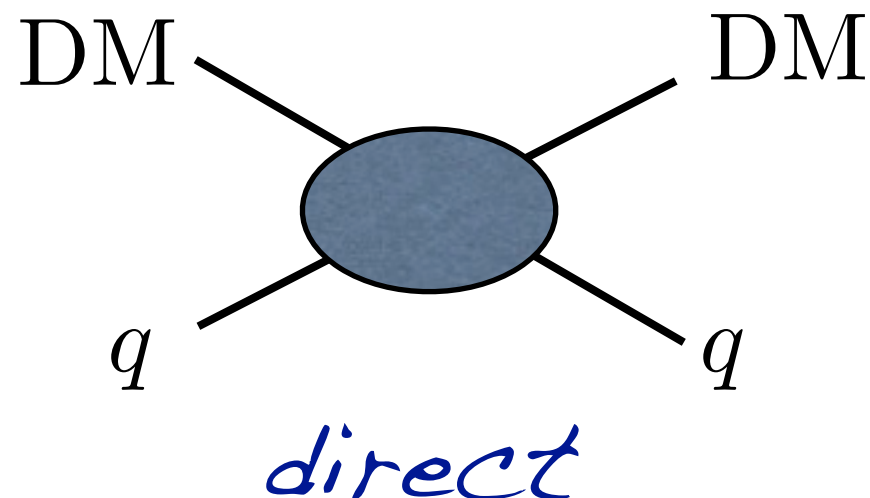
**Sales Pitch.**



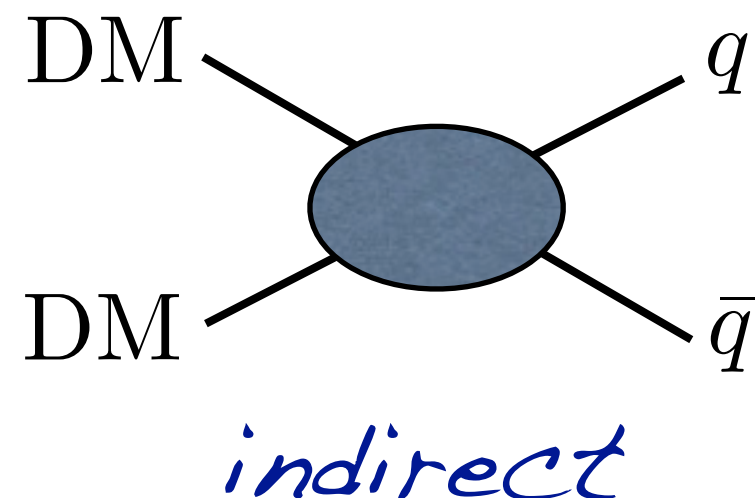
# Probes of DM Interactions

- \* “WIMP coincidence” hints that DM has an interaction w/ matter. picobarn-ish cross sections!?
- \* We hope to probe dark matter in several ways:

*DM-nucleus scattering*



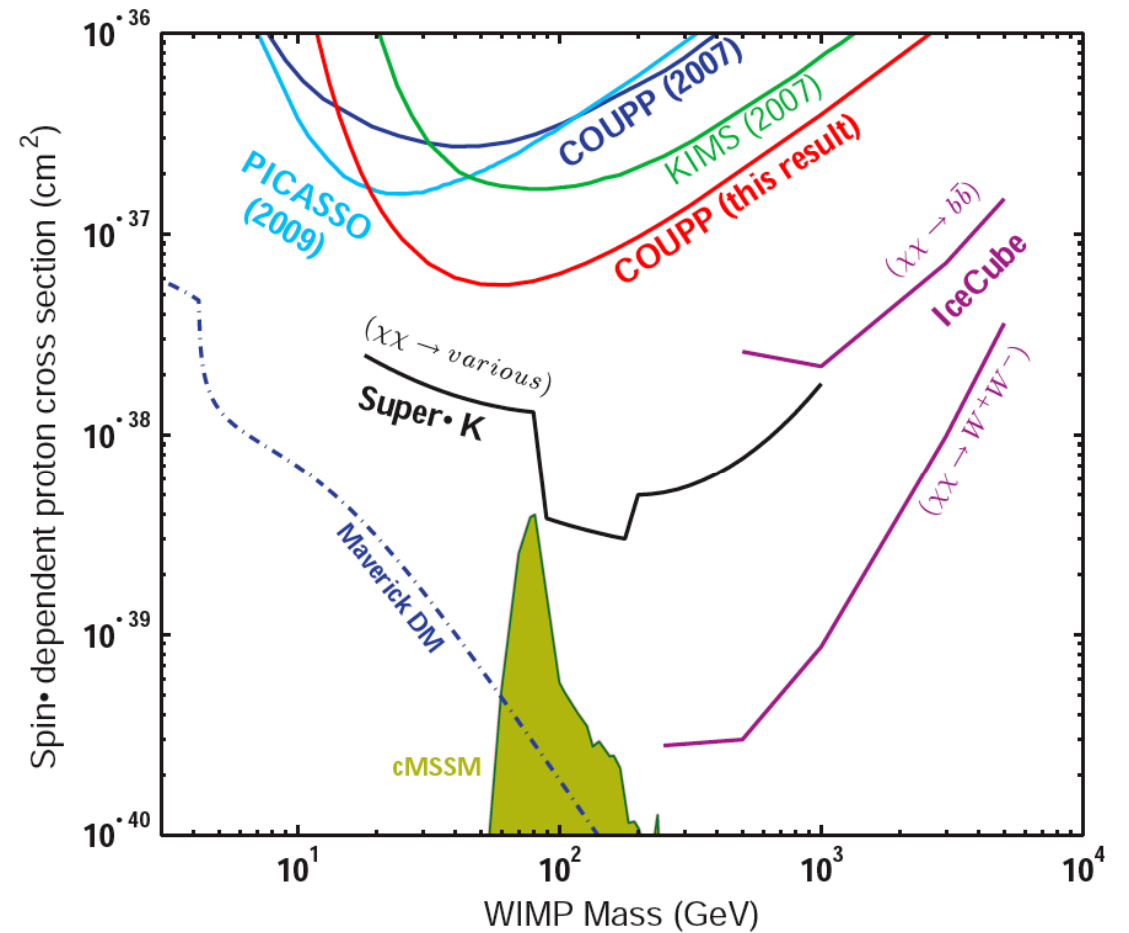
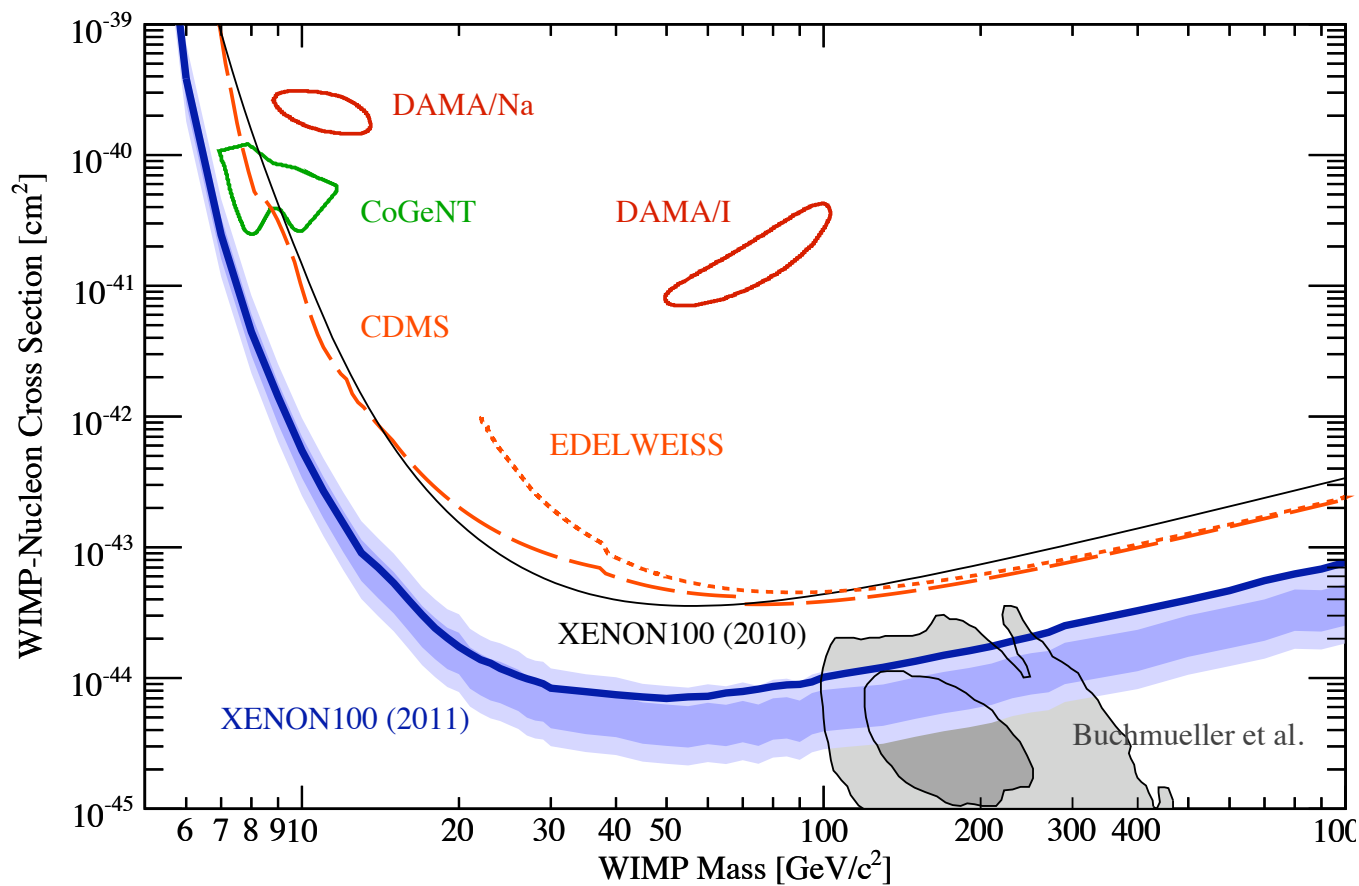
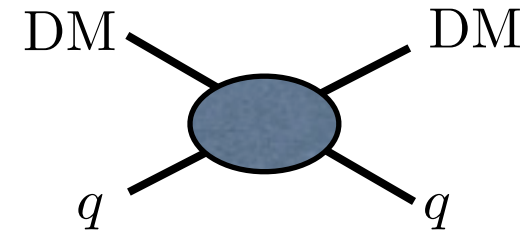
*DM annihilation*



Focus on direct detection in this talk.  
(a similar game can be played for indirect)

# Direct detection

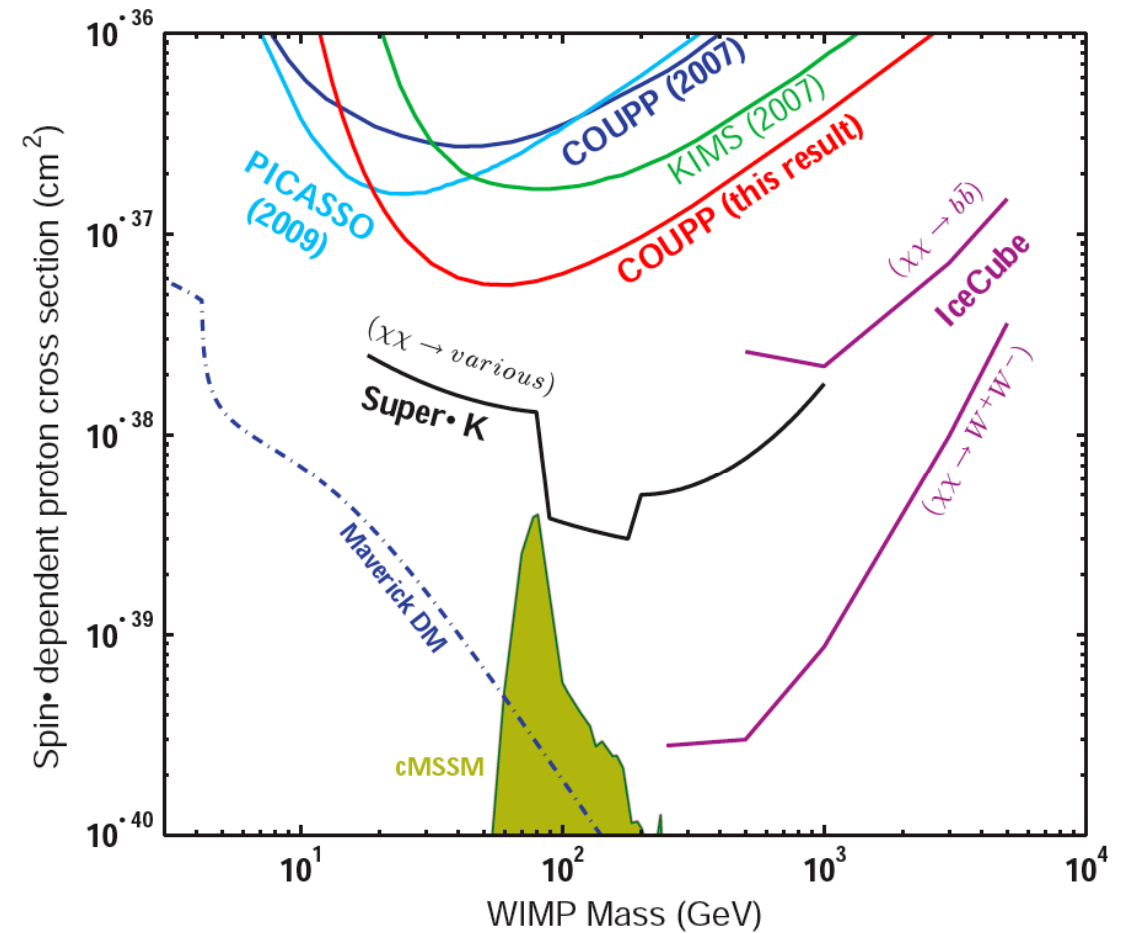
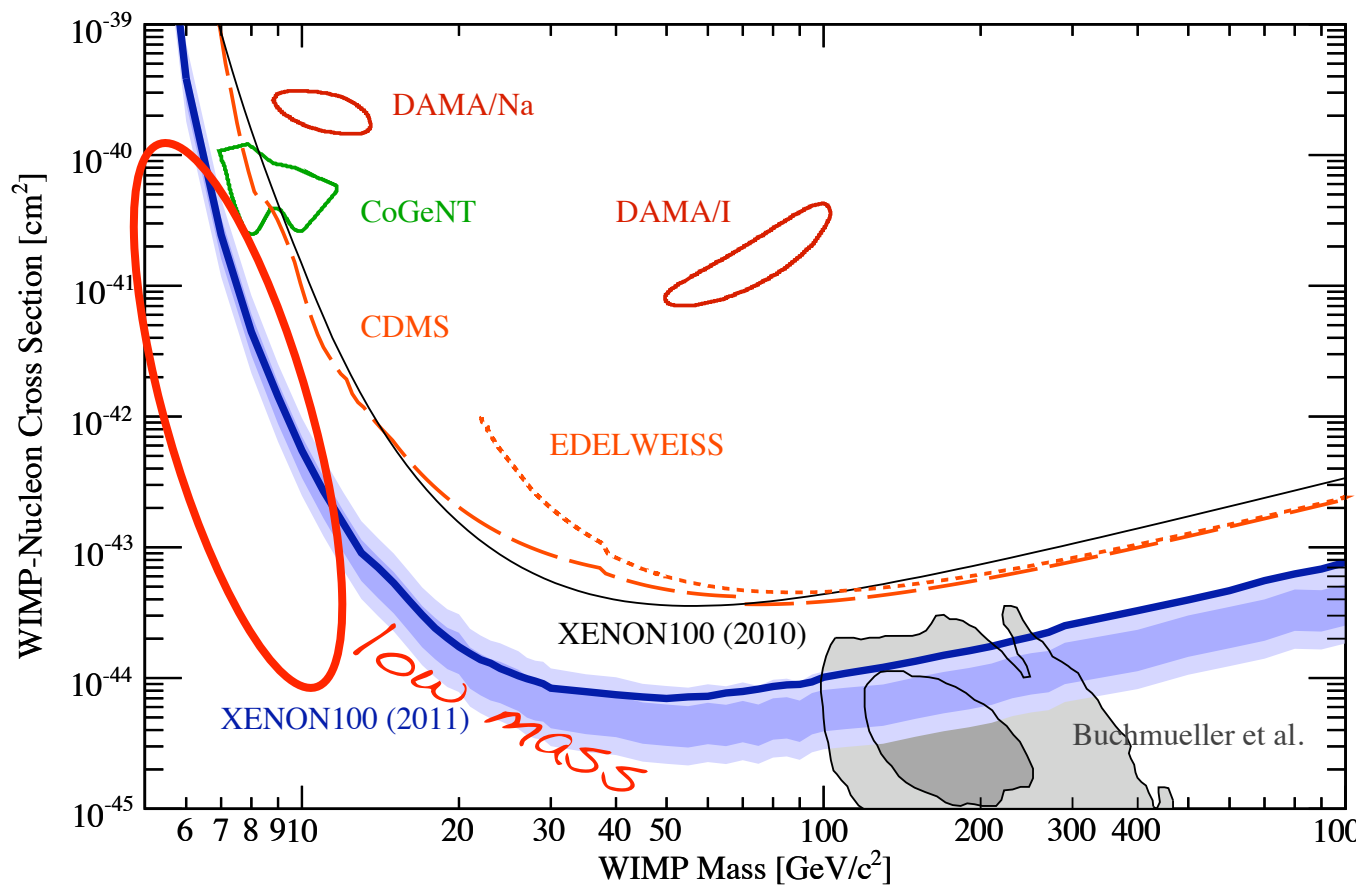
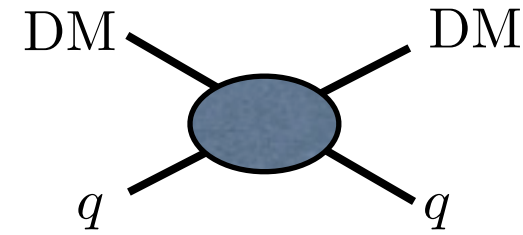
- \* Direct detection places limits on
- \* Heroic effort with remarkable results.
- \* DD has some weaknesses.



What do colliders tell us about this parameter space?

# Direct detection

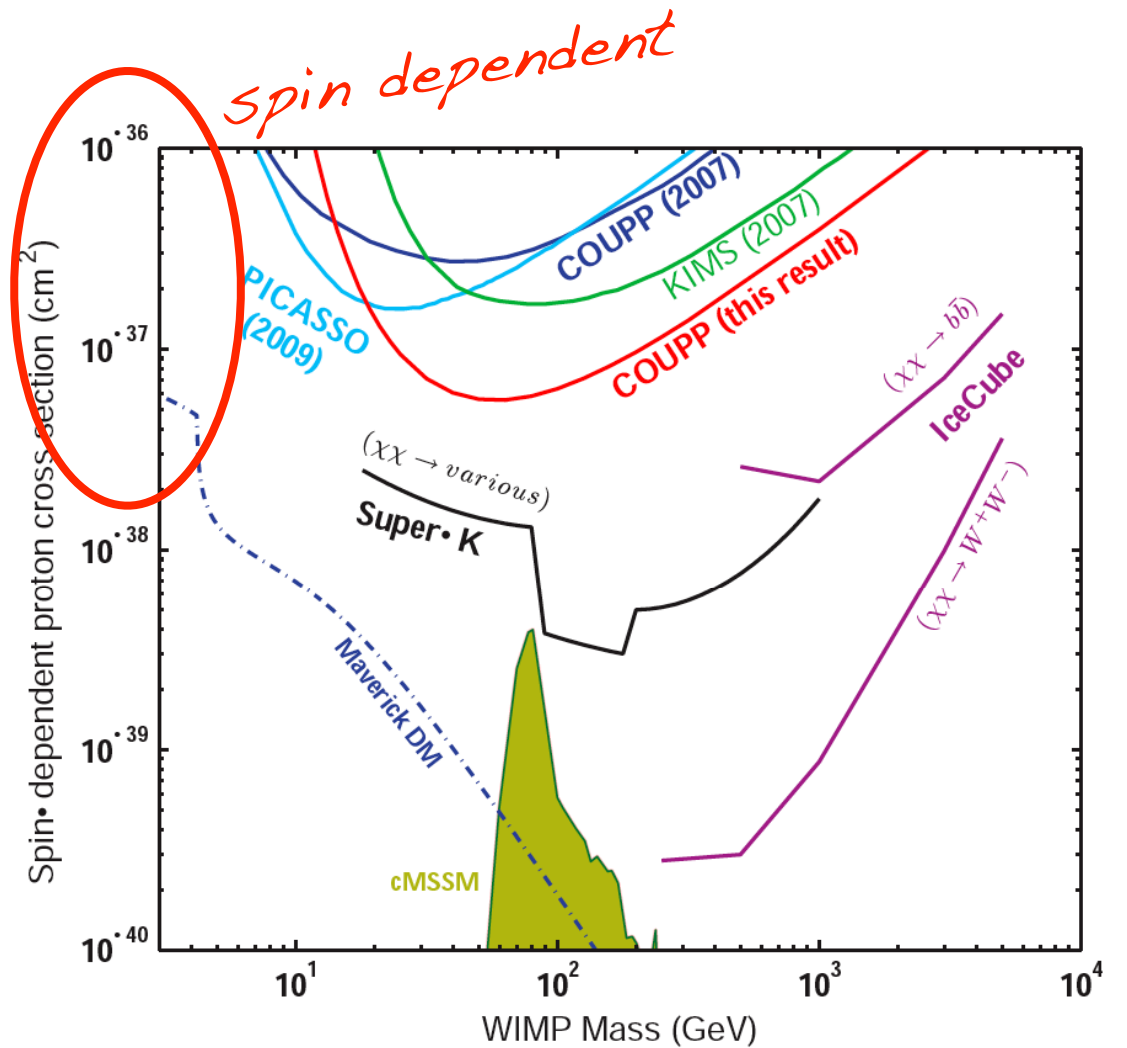
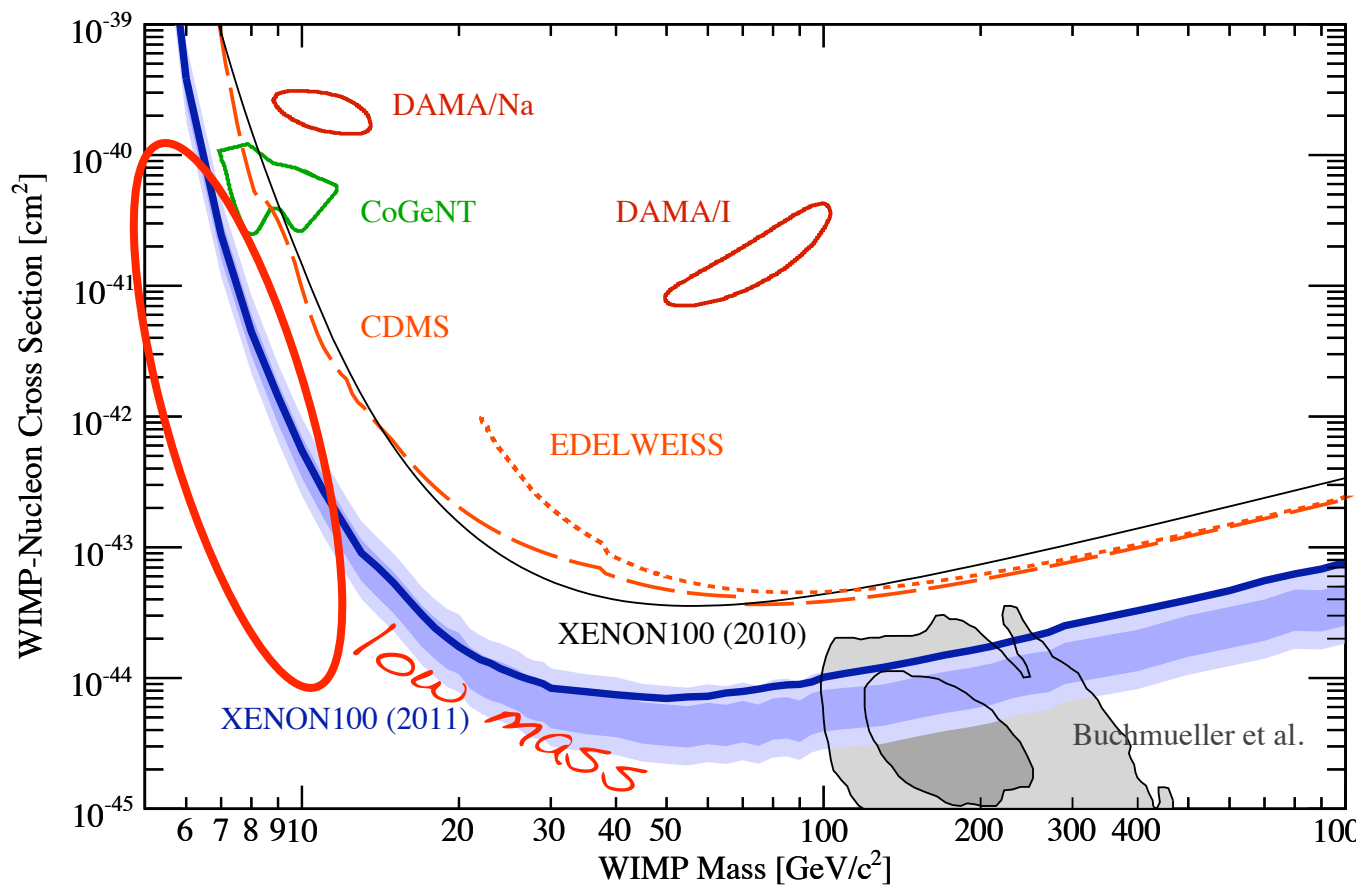
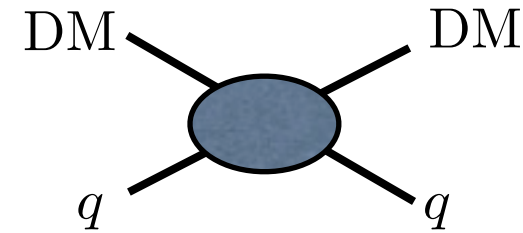
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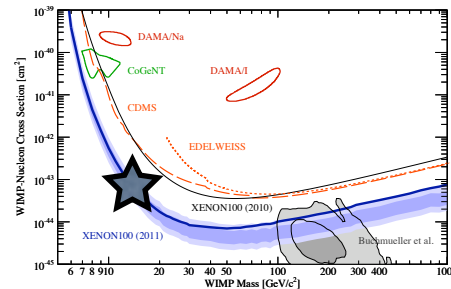


What do colliders tell us about this parameter space?

# A Simple Point

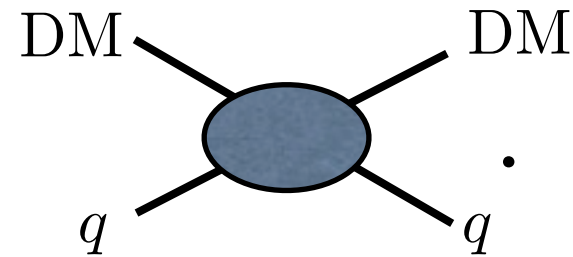
\* In order to get a particular DM-nucleon cross

section,

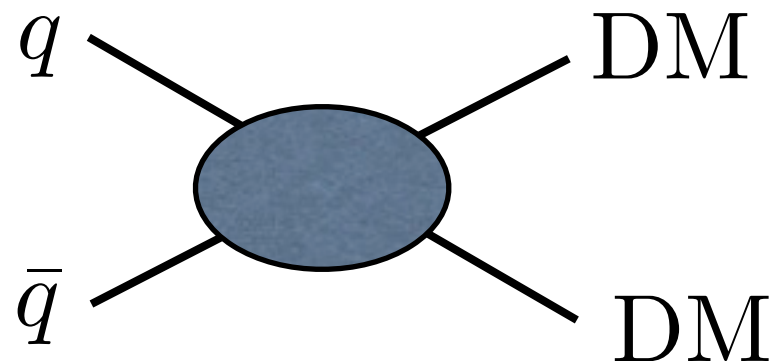


, we assume the existence of

a DM-hadron interaction,



\* The same interaction can lead to DM production at a hadron machine.

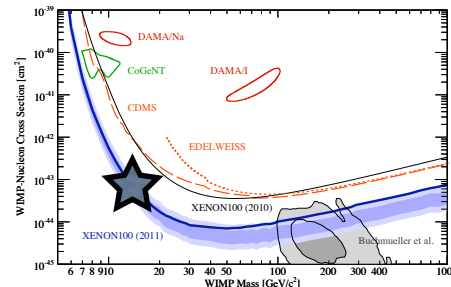


$p\bar{p} \rightarrow \text{nothing}$

# A Simple Point

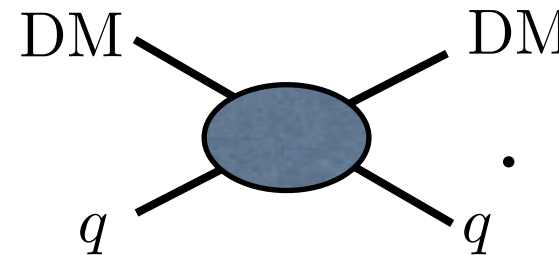
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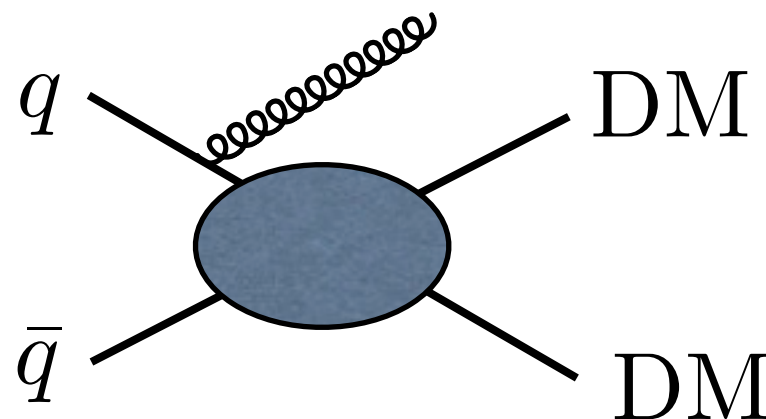


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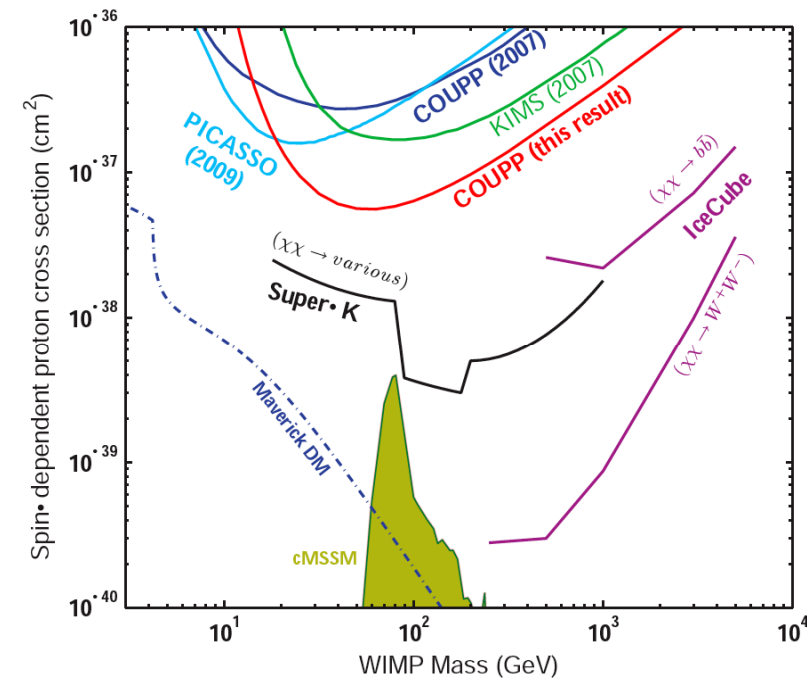
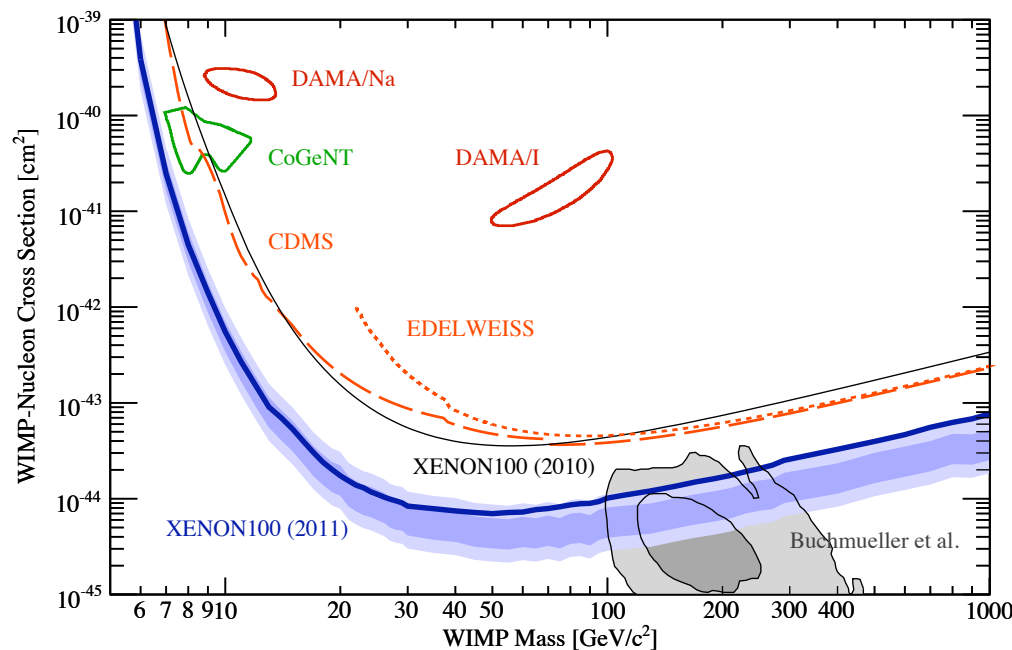
\* The same interaction can lead to DM production at a hadron machine.



$$p\bar{p} \rightarrow j + \cancel{E}_T$$

# A Simple Point

- \* **Mono-jet searches can place limits on the direct detection plane.**

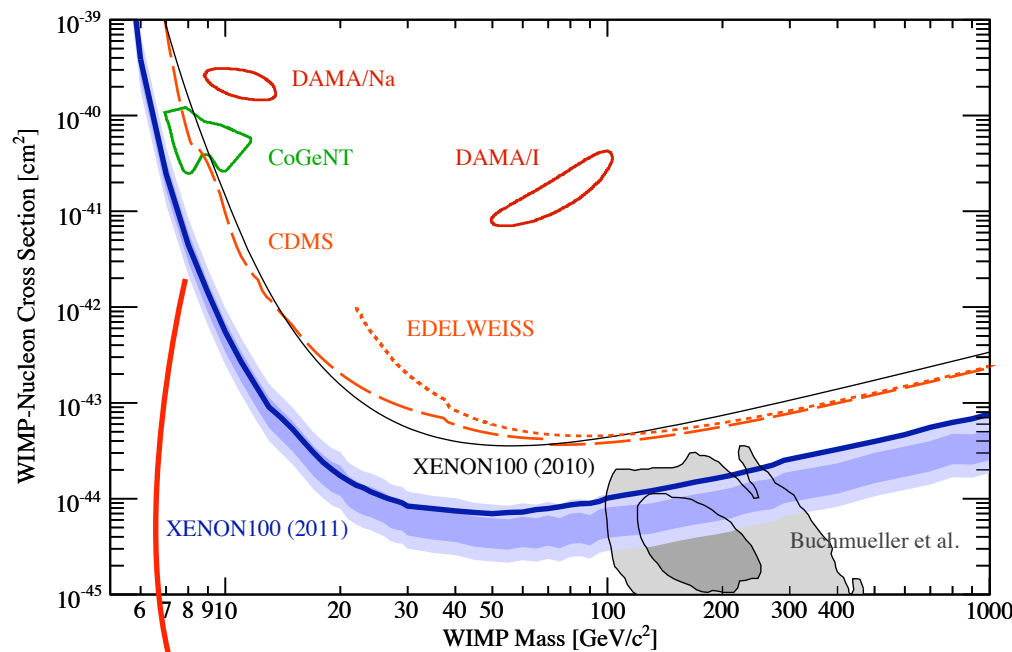


- \* These are **conservative** limits.  
In a specific model there may be other ways to produce DM, e.g. through cascades from heavy colored states.

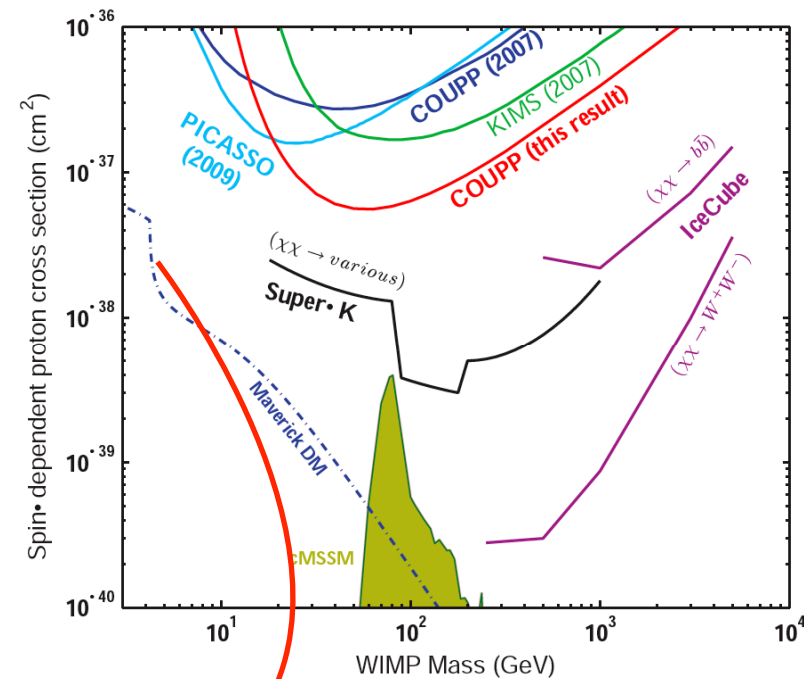
*But mono-jet are certainly good to set bounds.*

# A Simple Point

\* **Mono-jet searches can place limits on the plane.**



*The collider does not have a low energy threshold*



*The collider does not pay a price for spin dependence*



# Direct Detection - EFT

- \* The EFT that describes DM interaction in direct detection experiments:

$$\mathcal{O}_V = \frac{(\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu q)}{\Lambda^2},$$

SI, vector exchange

$$\mathcal{O}_A = \frac{(\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu\gamma_5q)}{\Lambda^2},$$

SD, axial-vector exchange

$$\mathcal{O}_t = \frac{(\bar{\chi}P_Rq)(\bar{q}P_L\chi)}{\Lambda^2} + (L \leftrightarrow R),$$

SI (or SD), t-channel  
“squark exchange”

$$\mathcal{O}_g = \alpha_s \frac{(\bar{\chi}\chi)(G_{\mu\nu}^a G^{a\mu\nu})}{\Lambda^3}$$

SI gluon operator

---

**Two possibilities:**

# Direct Detection - EFT

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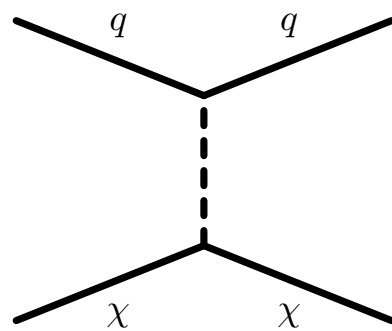
---

**Two possibilities:**

- 1) EFT is valid at LHC.
- 2) It's not.

# EFT vs Light Mediator

- \* The EFT is valid for direct detection ( $q \sim 100$  MeV):

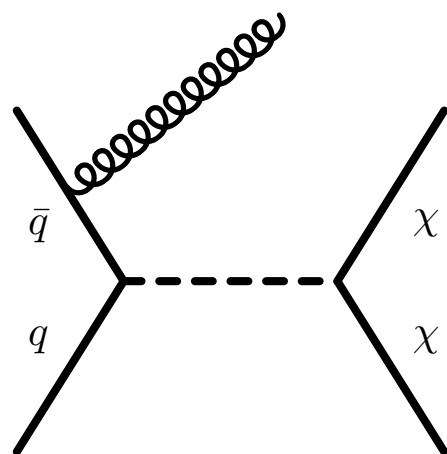


$$\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$$

$$\mu = \frac{m_{\chi} m_N}{m_N + m_{\chi}}$$

$$\Lambda \equiv \frac{M}{\sqrt{g_q g_{\chi}}}$$

- \* At a collider consider two extreme limits:



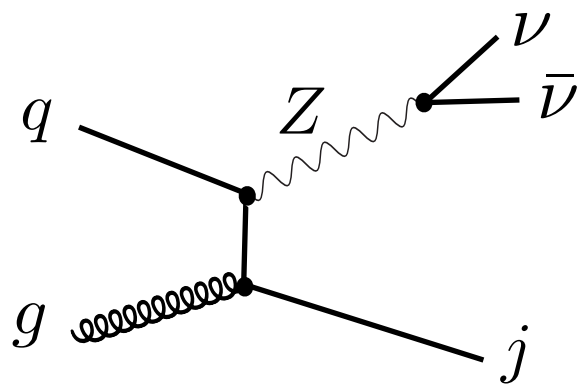
$$\sigma_{1j} \sim \begin{cases} \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2} & M \ll \sqrt{s_*} \\ \alpha_s g_{\chi}^2 g_q^2 \frac{p_T^2}{M^4} & M > \sqrt{s_*} \end{cases}$$

**Handles for S vs B  
&  
Inclusive Jets plus MET**

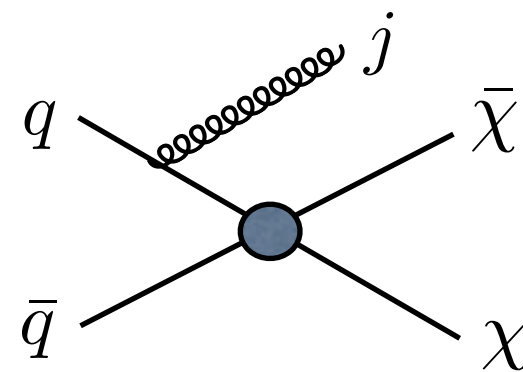
# Mono-Jet

- \* Assume the EFT is valid at the LHC.
- \* Consider contact operator involving u or d.
- \* The signal spectrum is harder than backgrounds.

dominant background:  
Z plus jet ( $qg$  initial state).



dominant signal:  
 $q\bar{q}$  initial state.

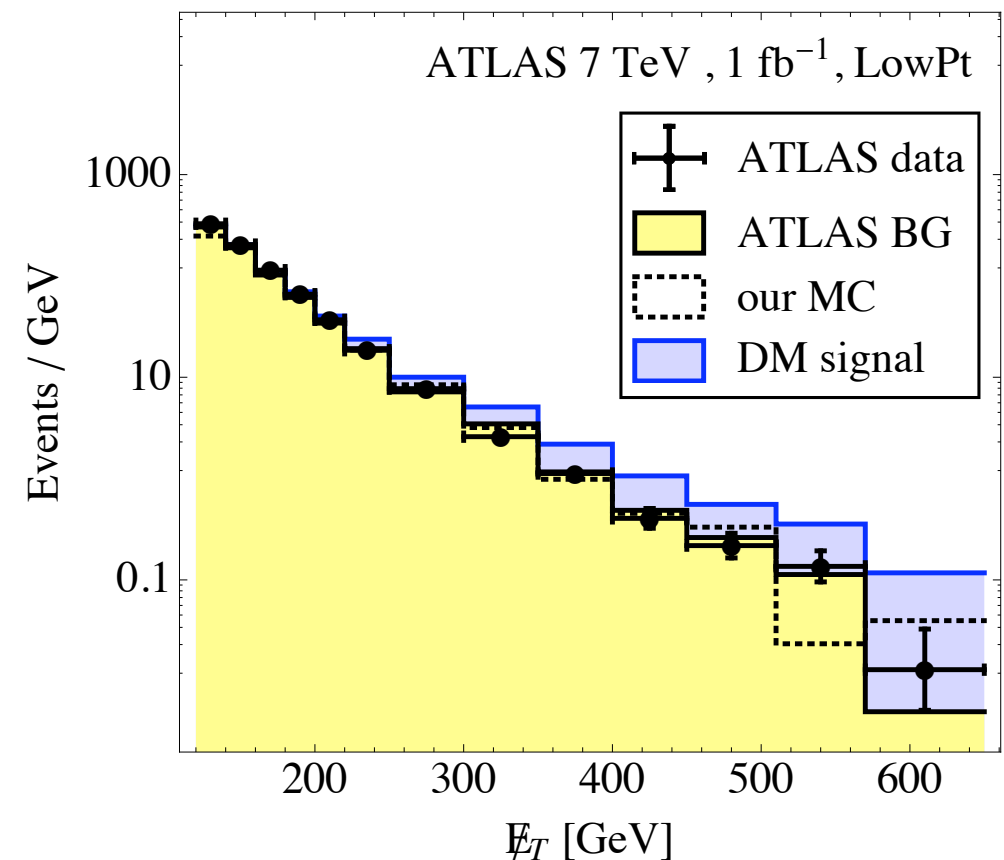


$Z$  is typically emitted forward,  
with low  $p_T$ .

DM system emitted isotropically.

# Mono-Jet

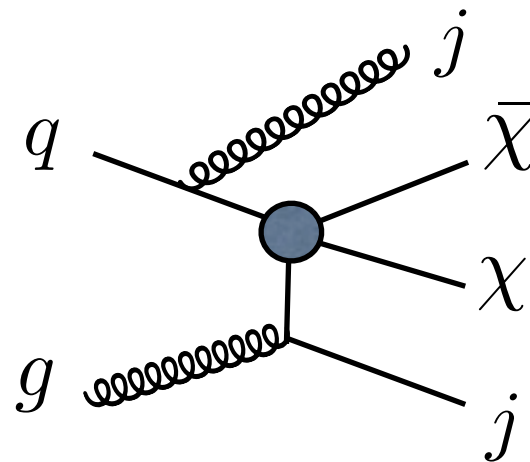
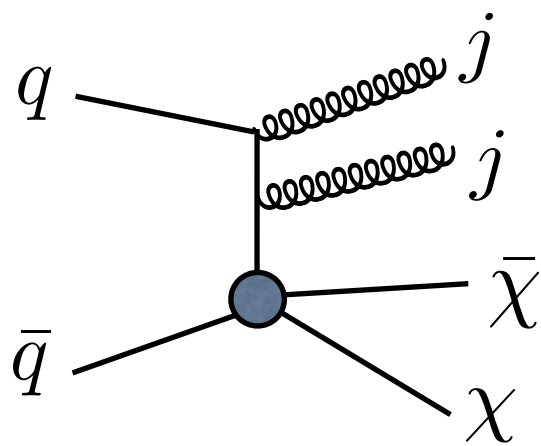
- \* Signal is **harder** and **more central** (unless DM couples to sea quarks).



- \* Add angular info to mono-jet analyses? (could improve ADD limits as well)

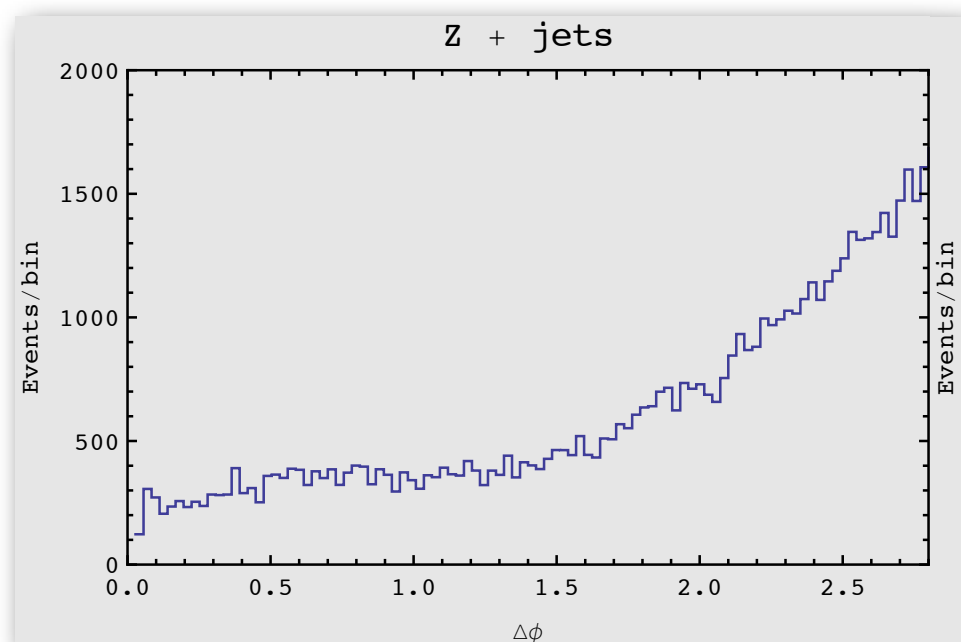
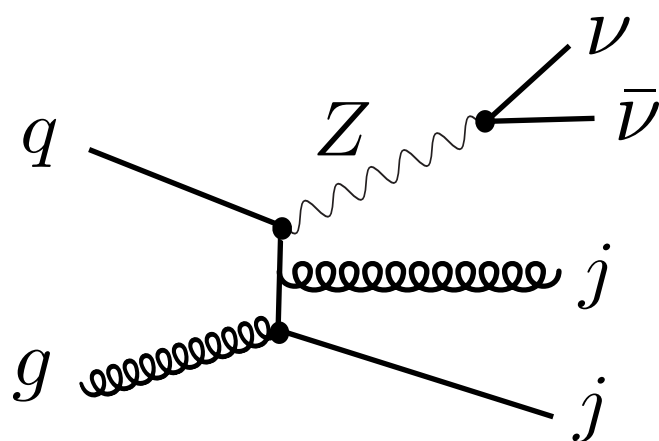
# More Jets

- \* Applying a veto on the second jet reduces signal efficiency, and increases theory uncertainty.
- \* Indeed, most recent CMS (and upcoming ATLAS) mono-jet searches allow for a **hard second jet**, so-long as its not back to back with first jet.
- \* Inclusive jets plus MET searches for SUSY exist. Can we use them as searches for dark matter?

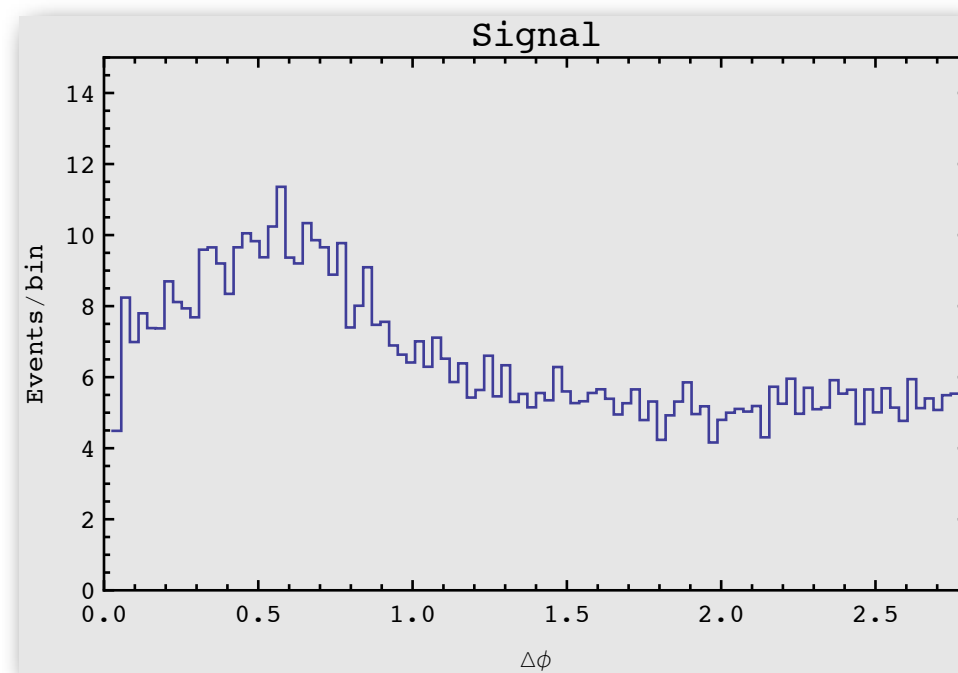
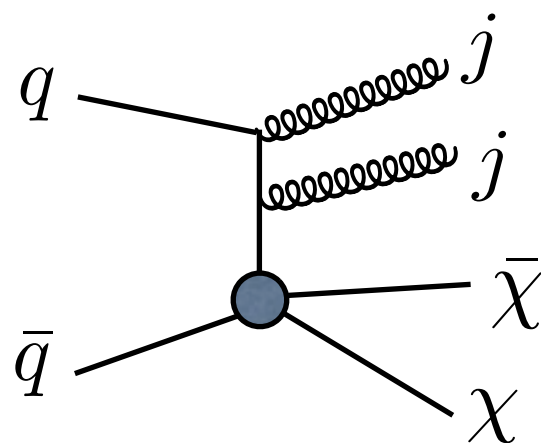


# More Jets

- \* Signal and Background have different dominant initial states, different angular distributions:



*Azimuthal  
angle  
b/w jets*





# Razor

- \* We would like an analysis that is also sensitive to the angular distribution of jets and to MET.
- \* CMS's Razor analysis limits SUSY by inspecting a 2D distribution of two kinematic variables.
- \* The Razor variables follow simple exponential distributions - a data driven analysis.

# Razor

\* Consider events

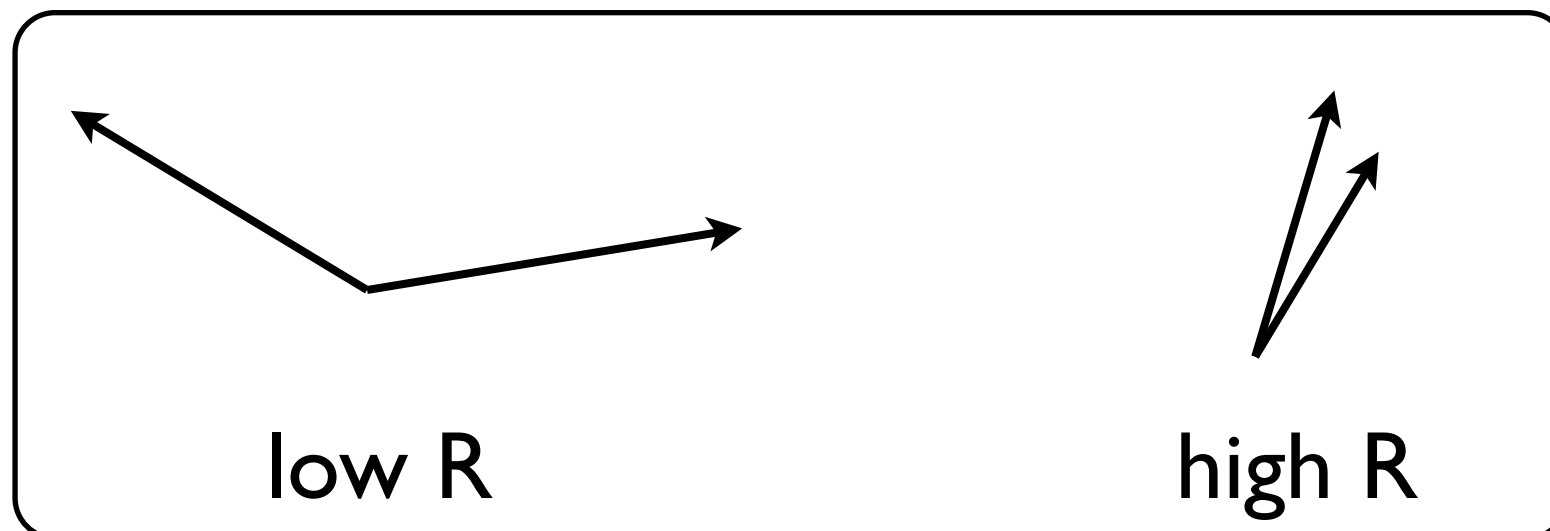
$$M_R = \sqrt{(E_{j_1} + E_{j_2})^2 - (p_z^{j_1} + p_z^{j_2})^2}$$

$$M_R^T = \sqrt{\frac{\cancel{E}_T(p_T^{j_1} + p_T^{j_2}) - \vec{\cancel{E}}_T \cdot (\vec{p}_T^{j_1} + \vec{p}_T^{j_2})}{2}}$$

$$R = \frac{M_R^T}{M_R}$$

*Estimators for  
characteristic  
scale in the event*

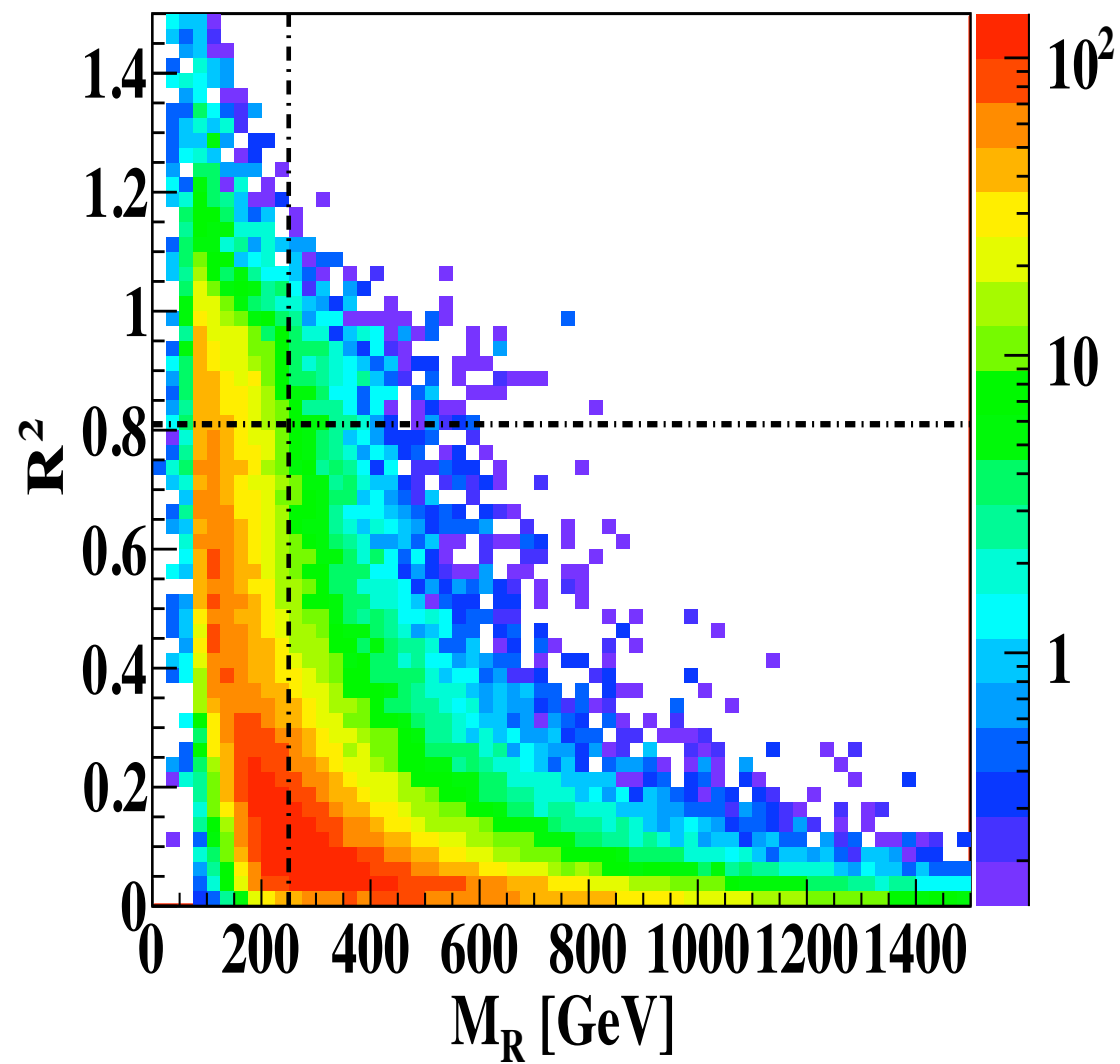
*roughly: R=ratio of  
MET to visible.*



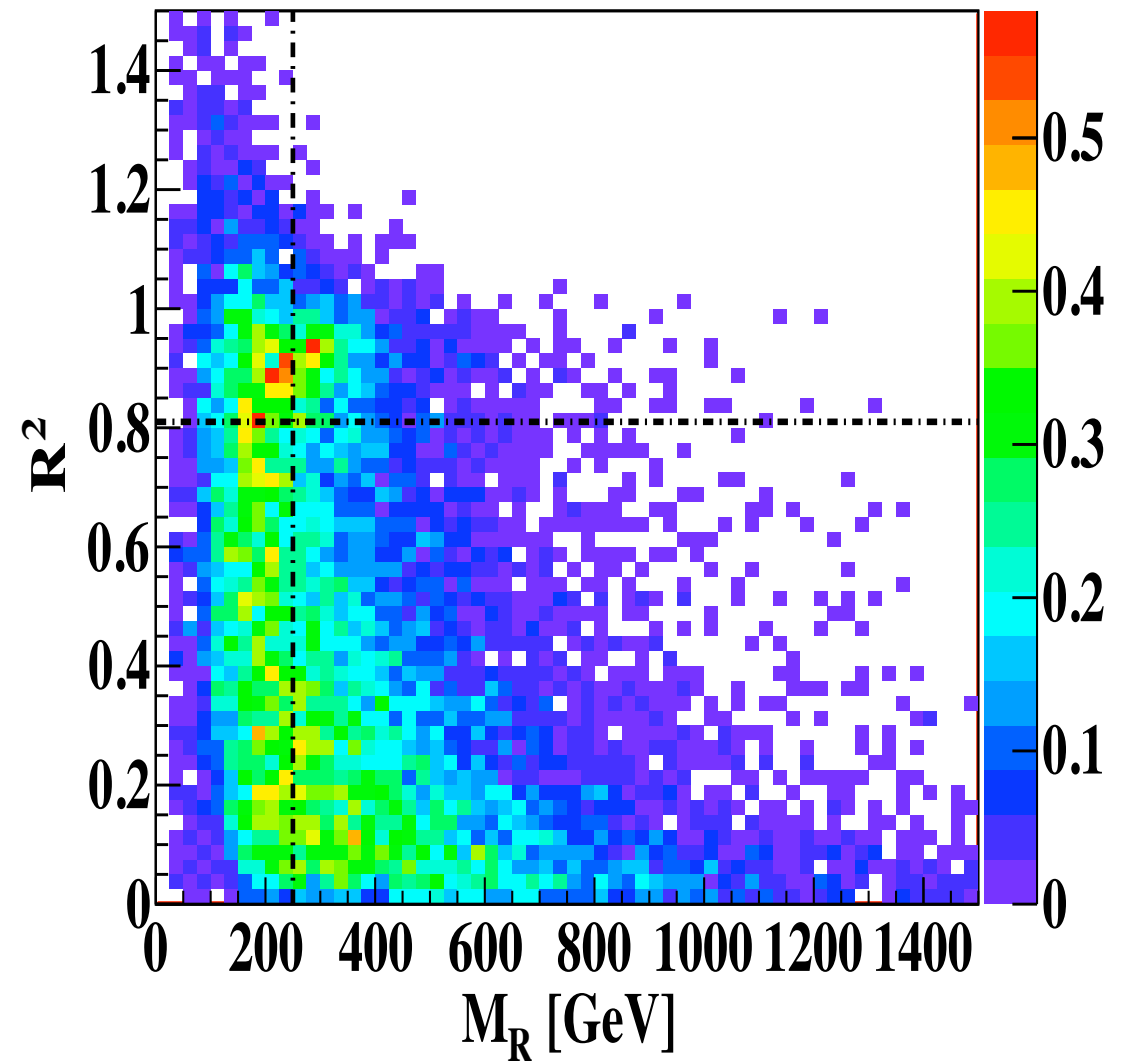
# Razor

\* Background and Signal shapes in  $M_R$  vs  $R$ :

Z+jets



Signal. 100 GeV DM.

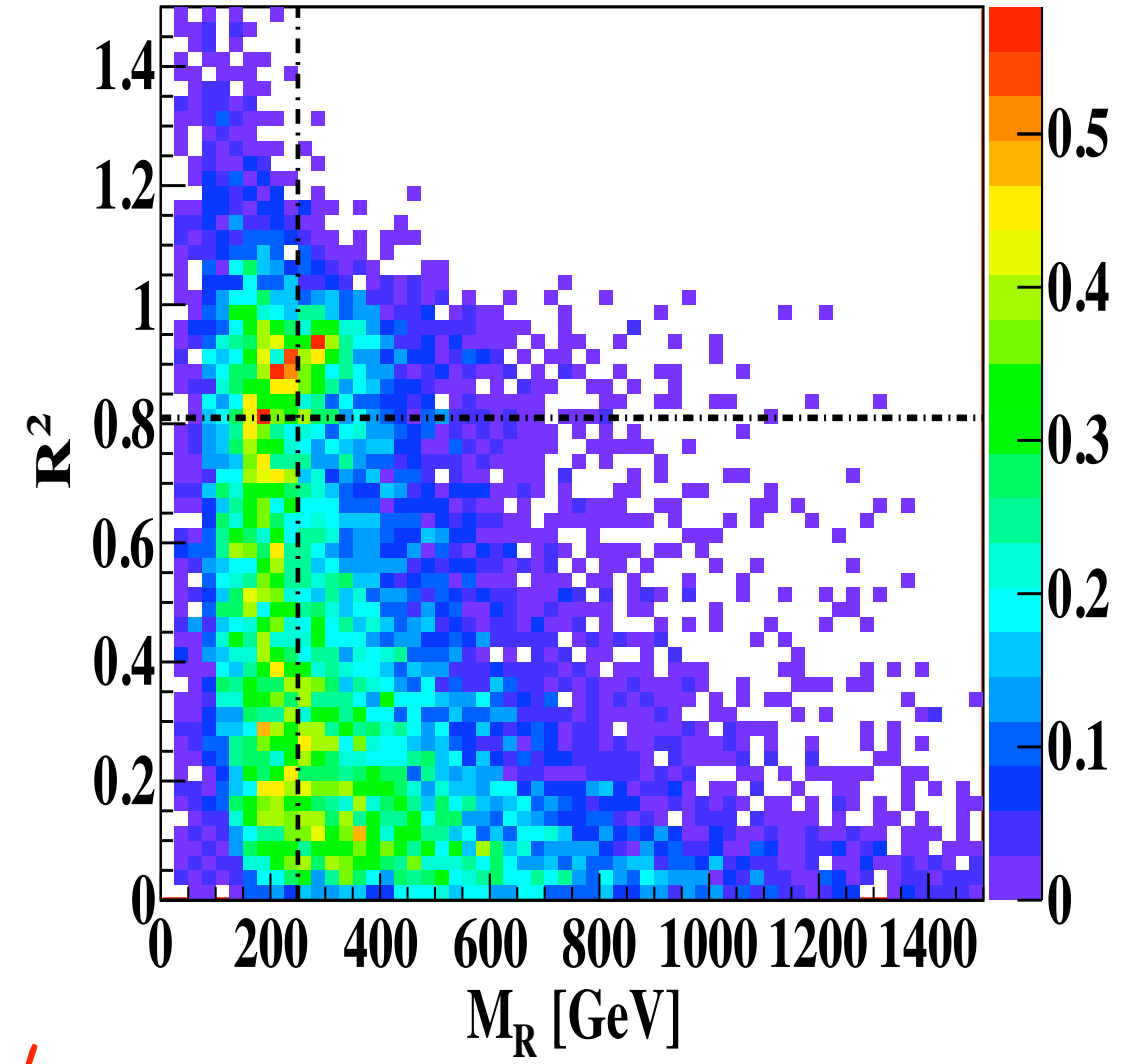
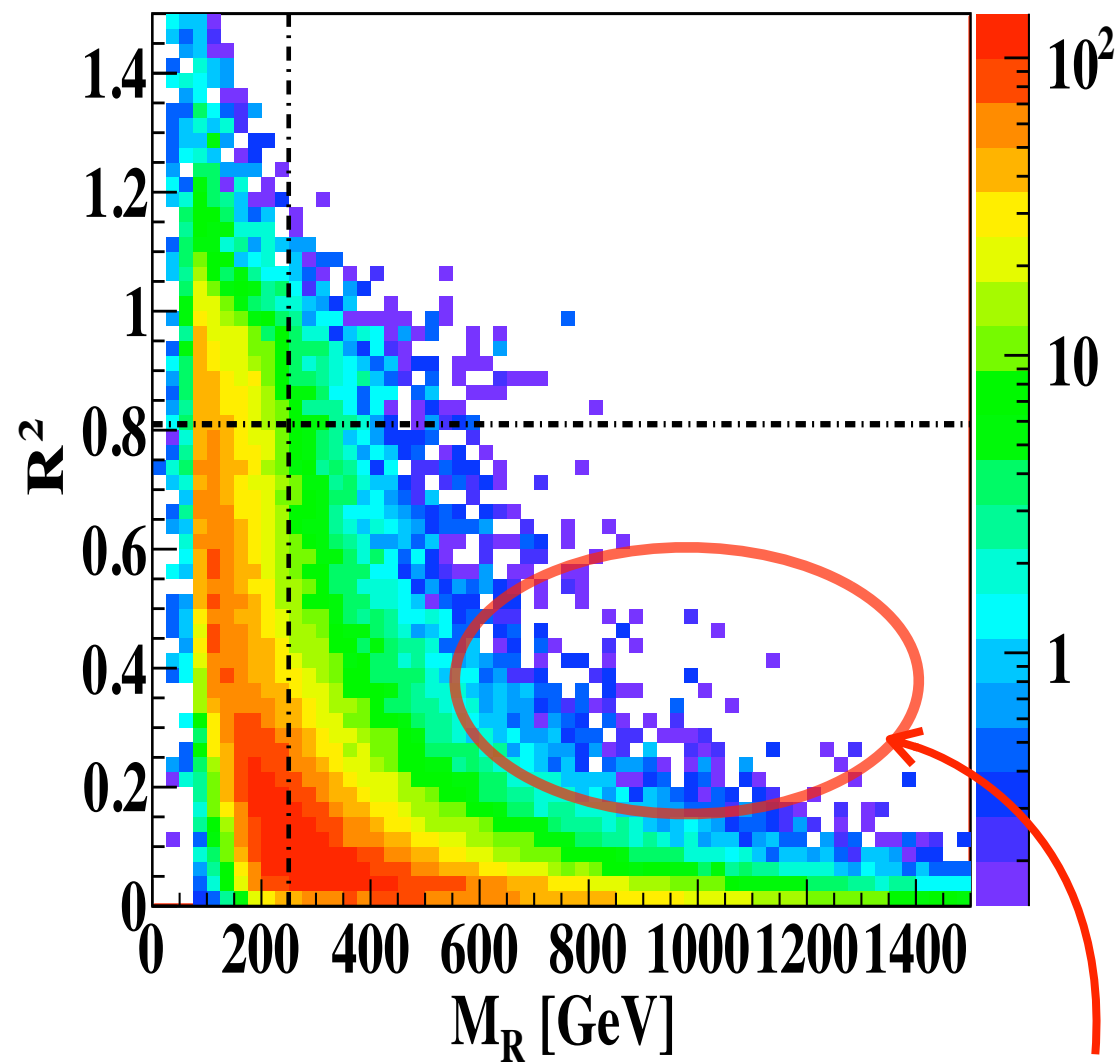


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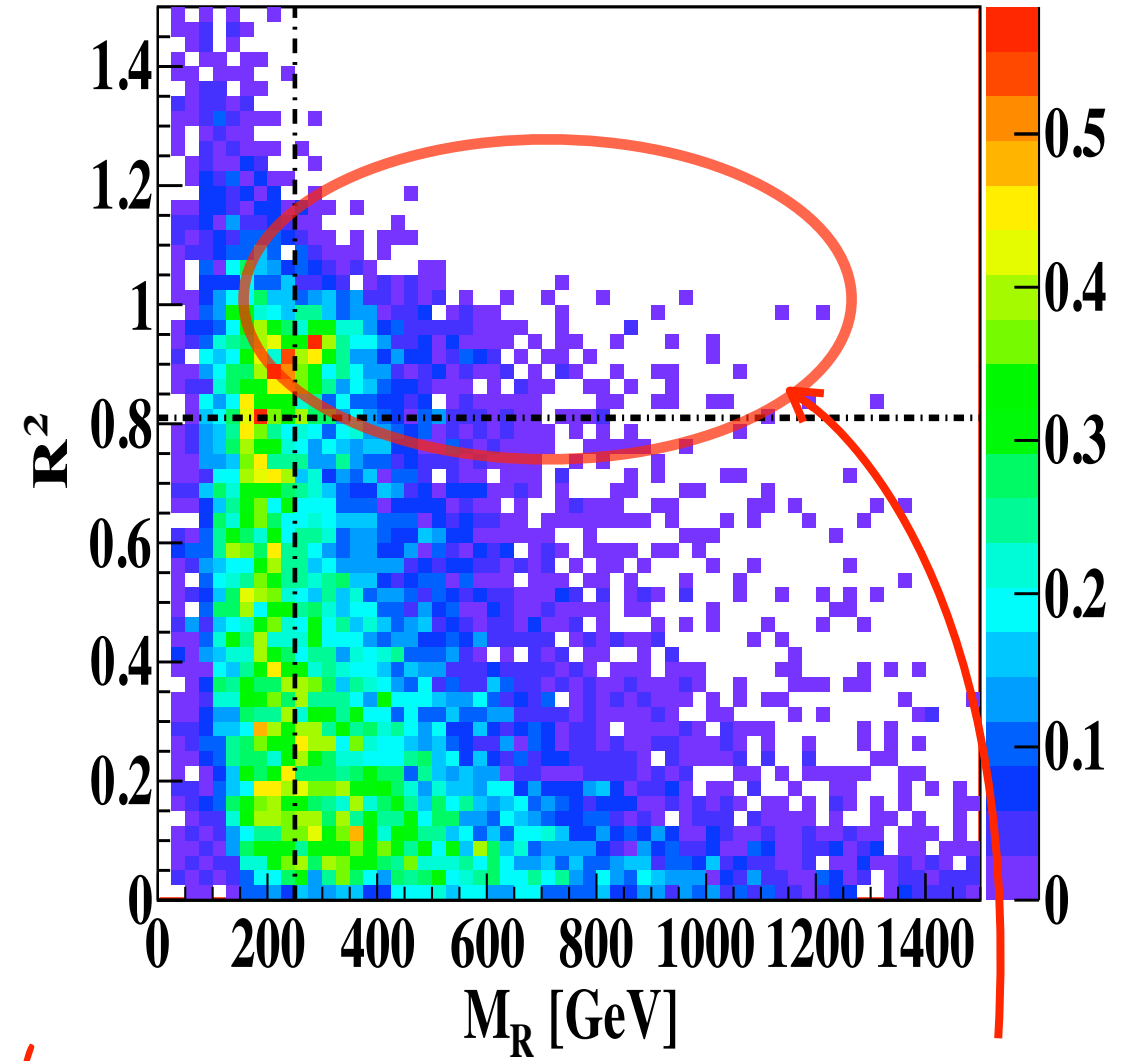
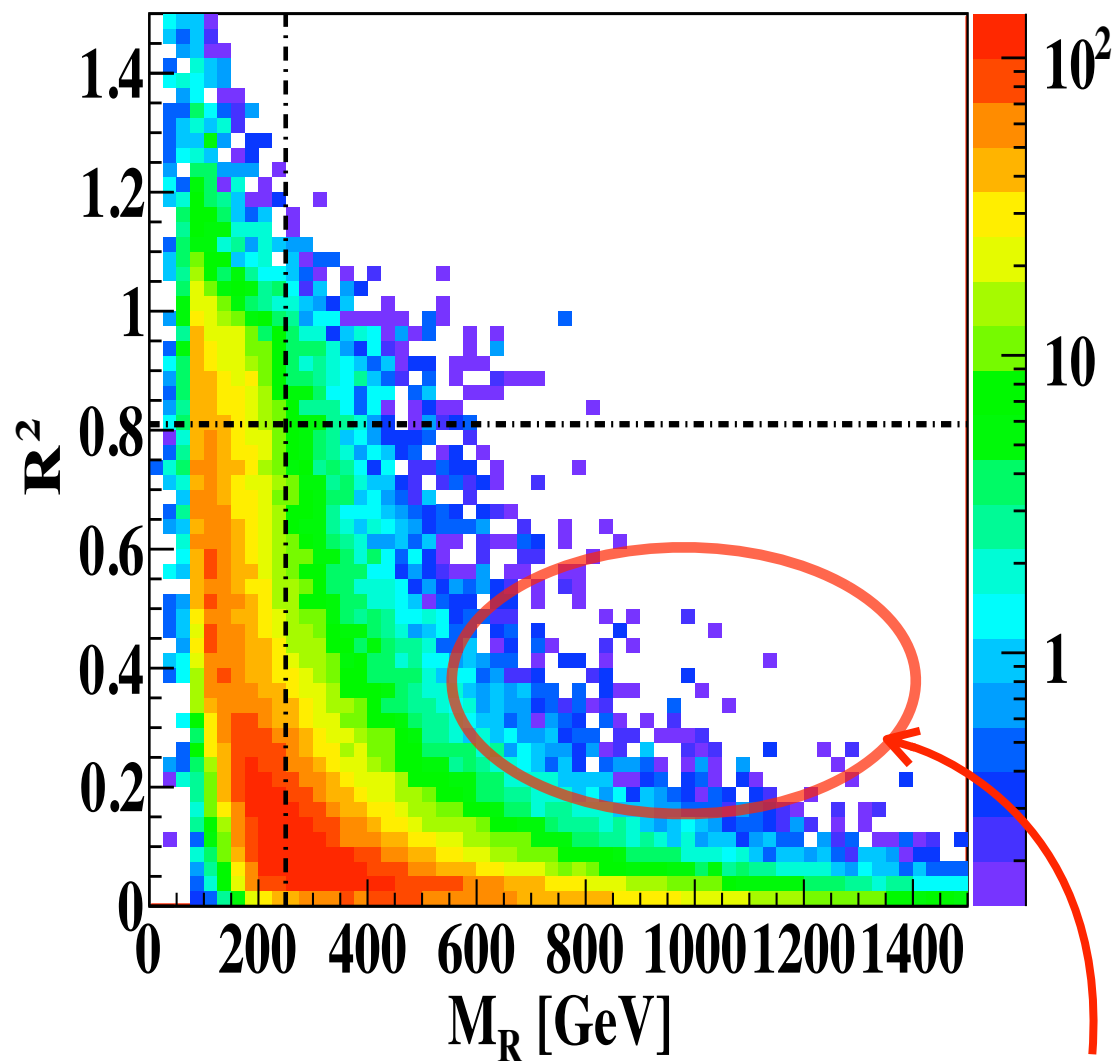
*SUSY search*

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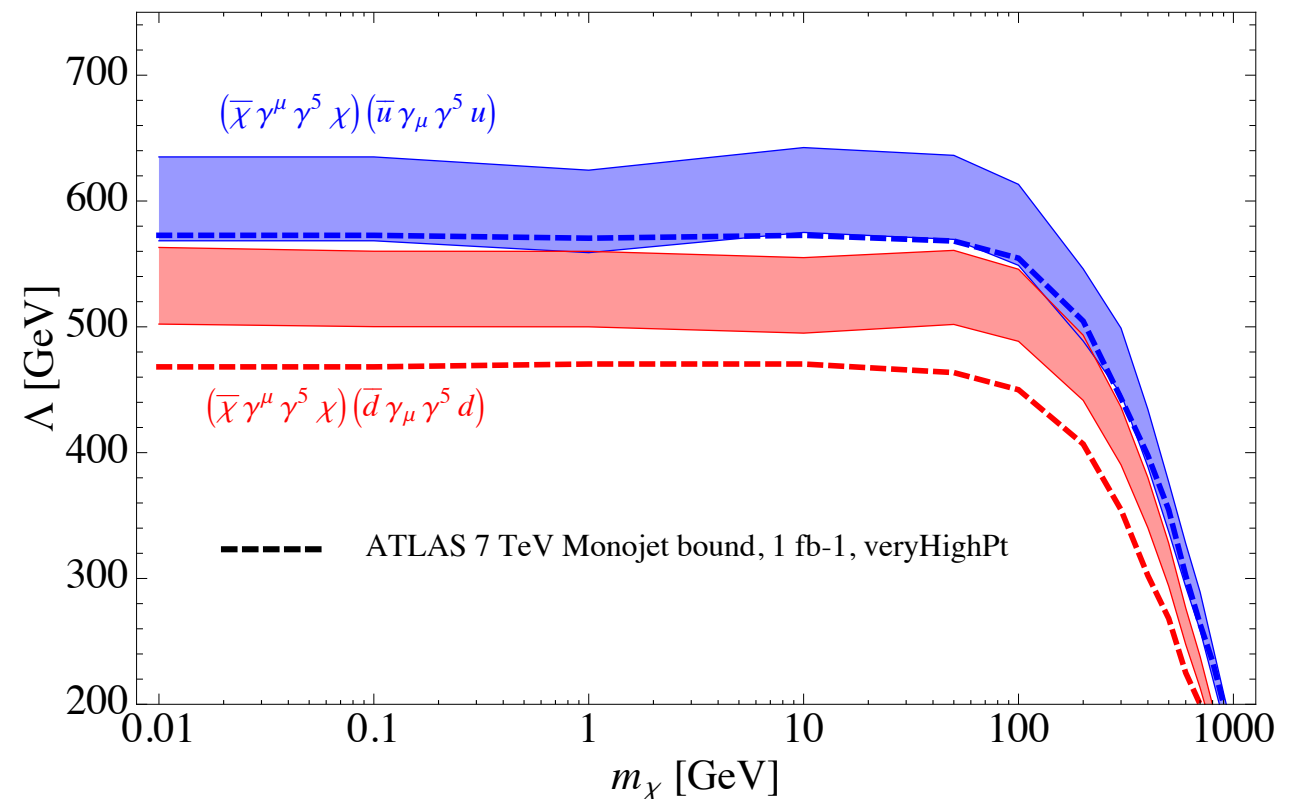
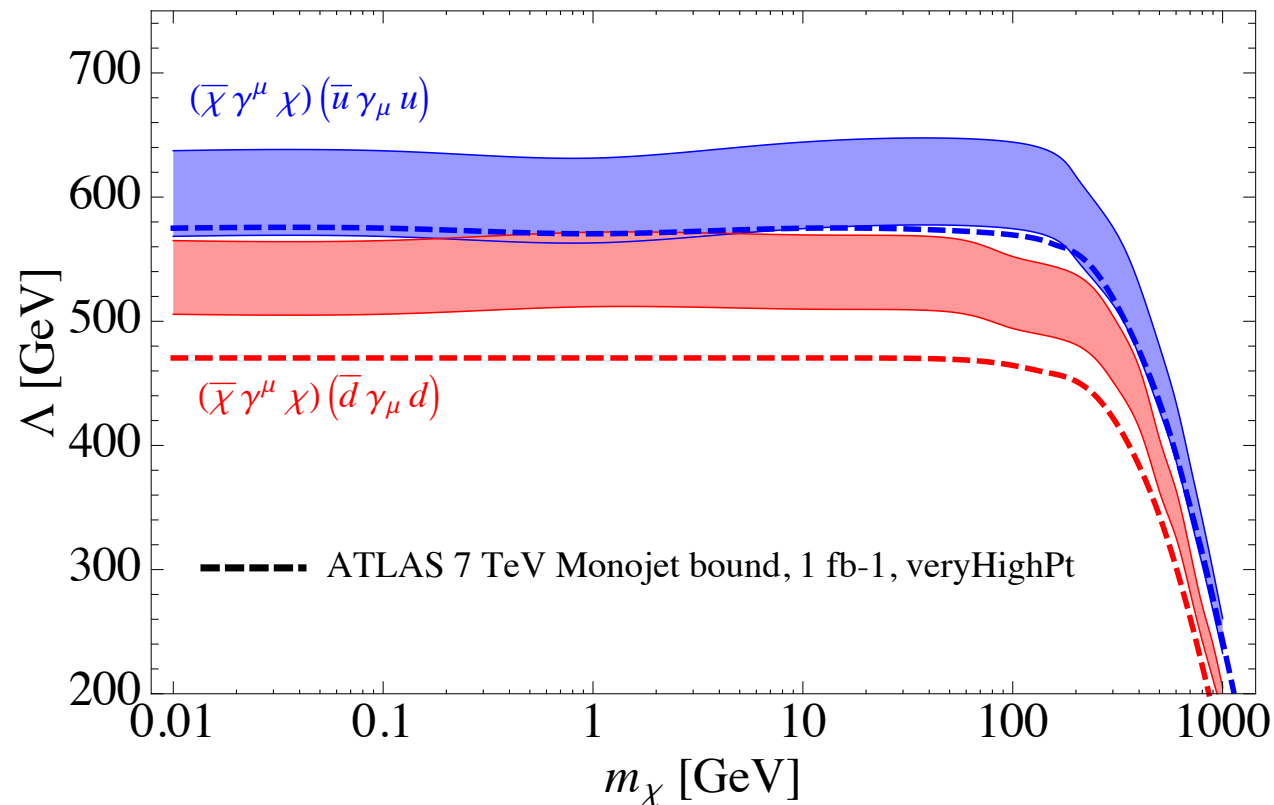


*SUSY search*

*our search*

# Razor Limits

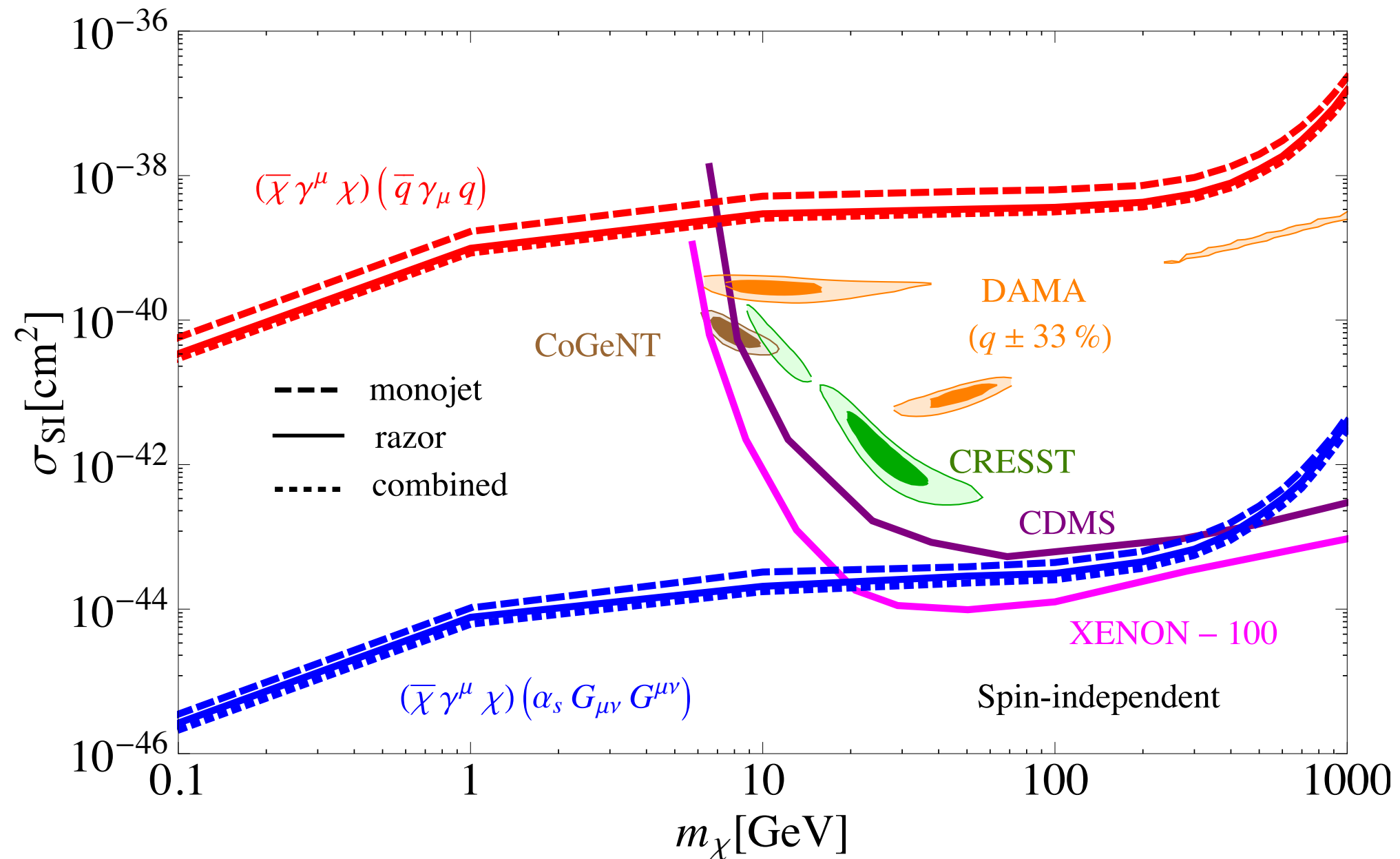
\* With 800 fb<sup>-1</sup>:



\* The band represents two assumption about uncertainties:  
 (1) statistics only, and (2) systematic=statistic.  
 truth is probably closer to (1).

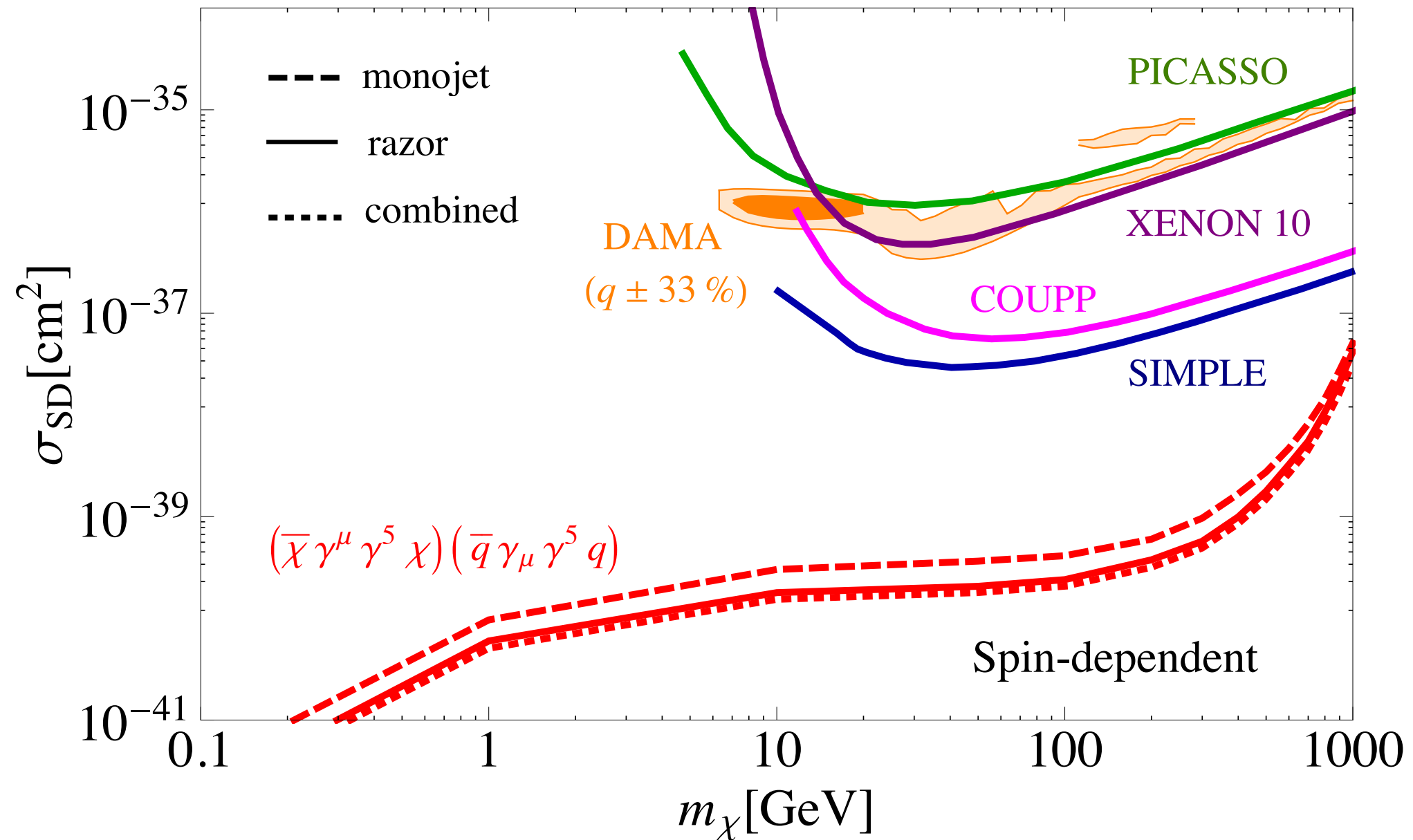
# Razor Limits

\* Spin-Independent (with  $800 \text{ fb}^{-1}$ ):



# Razor Limits

\* Spin-Dependent:





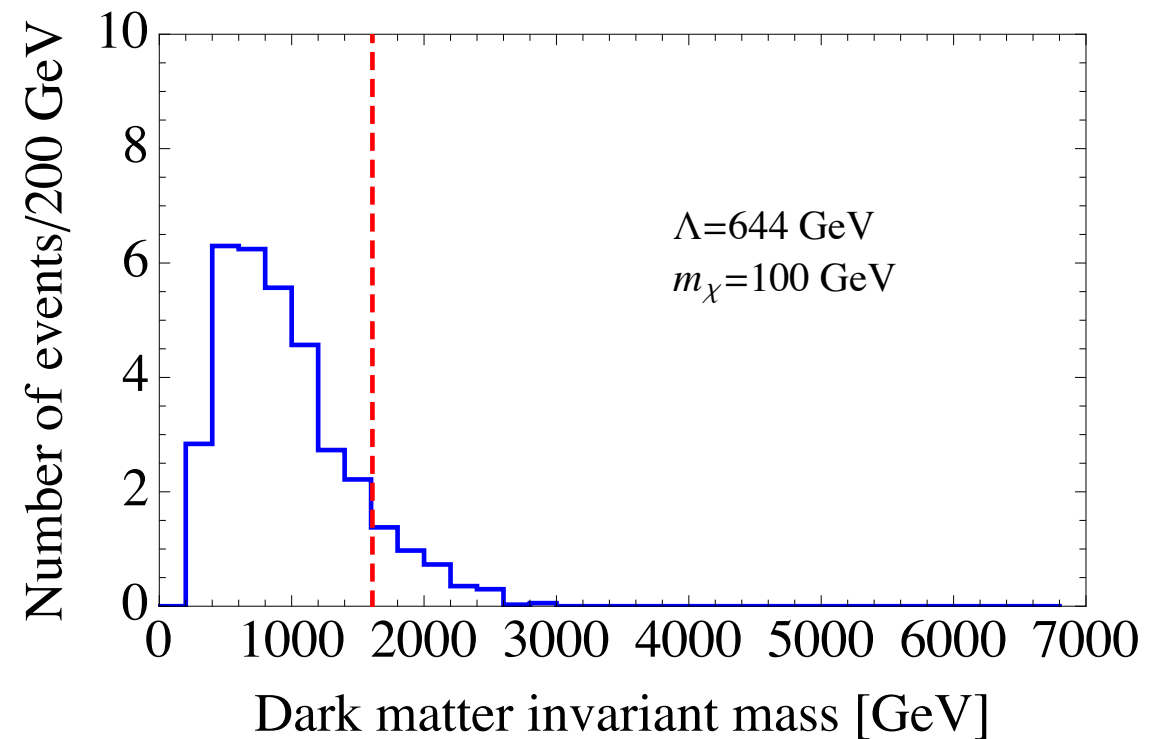
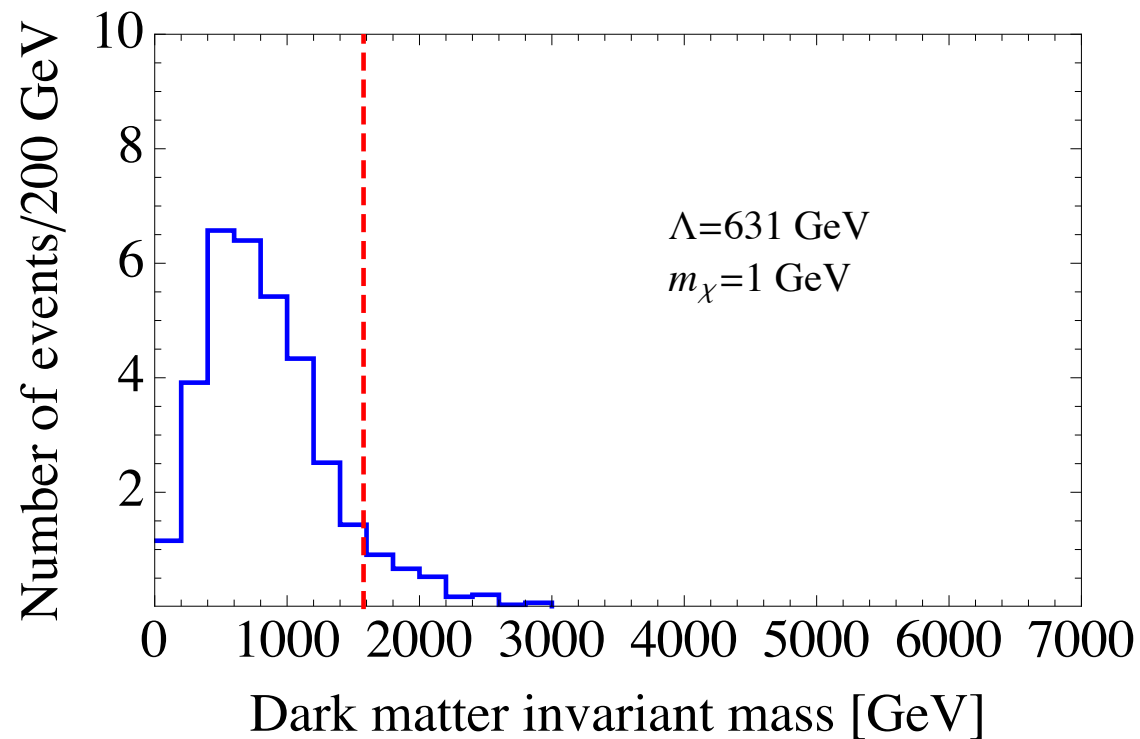
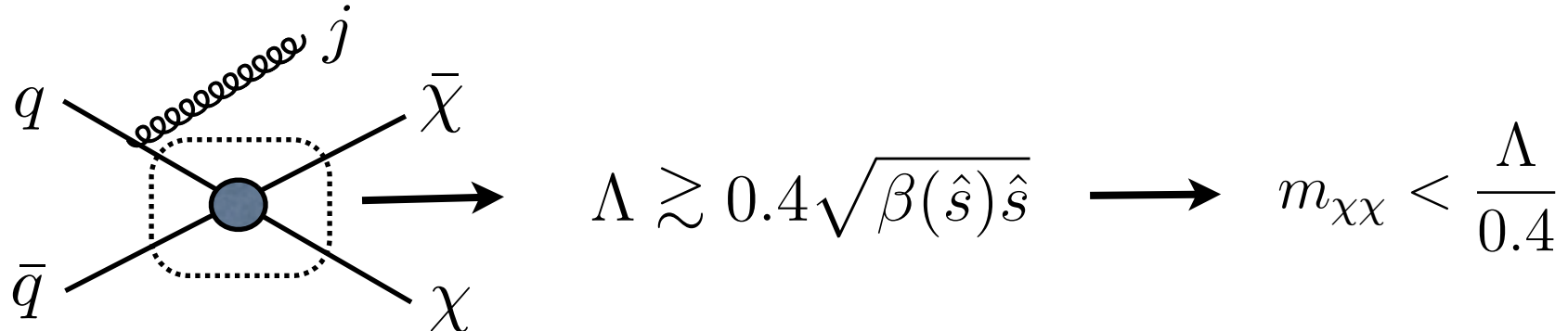
# EFT

Is it always valid?  
“Model independent”?  
Certainly not.

Is it useful?  
Certainly.

# Unitarity

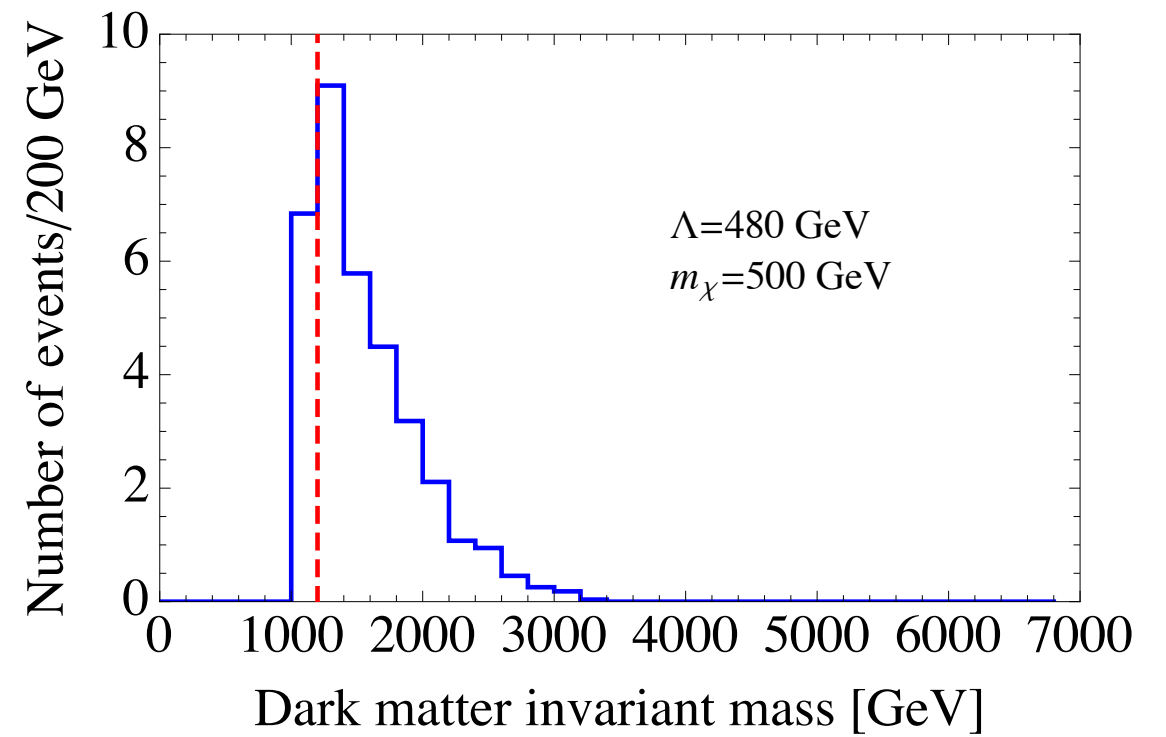
- \* Vechi and Shoemaker point to a unitarity limit:



The EFT may be accurate to  $\sim 10\%$  for light DM.

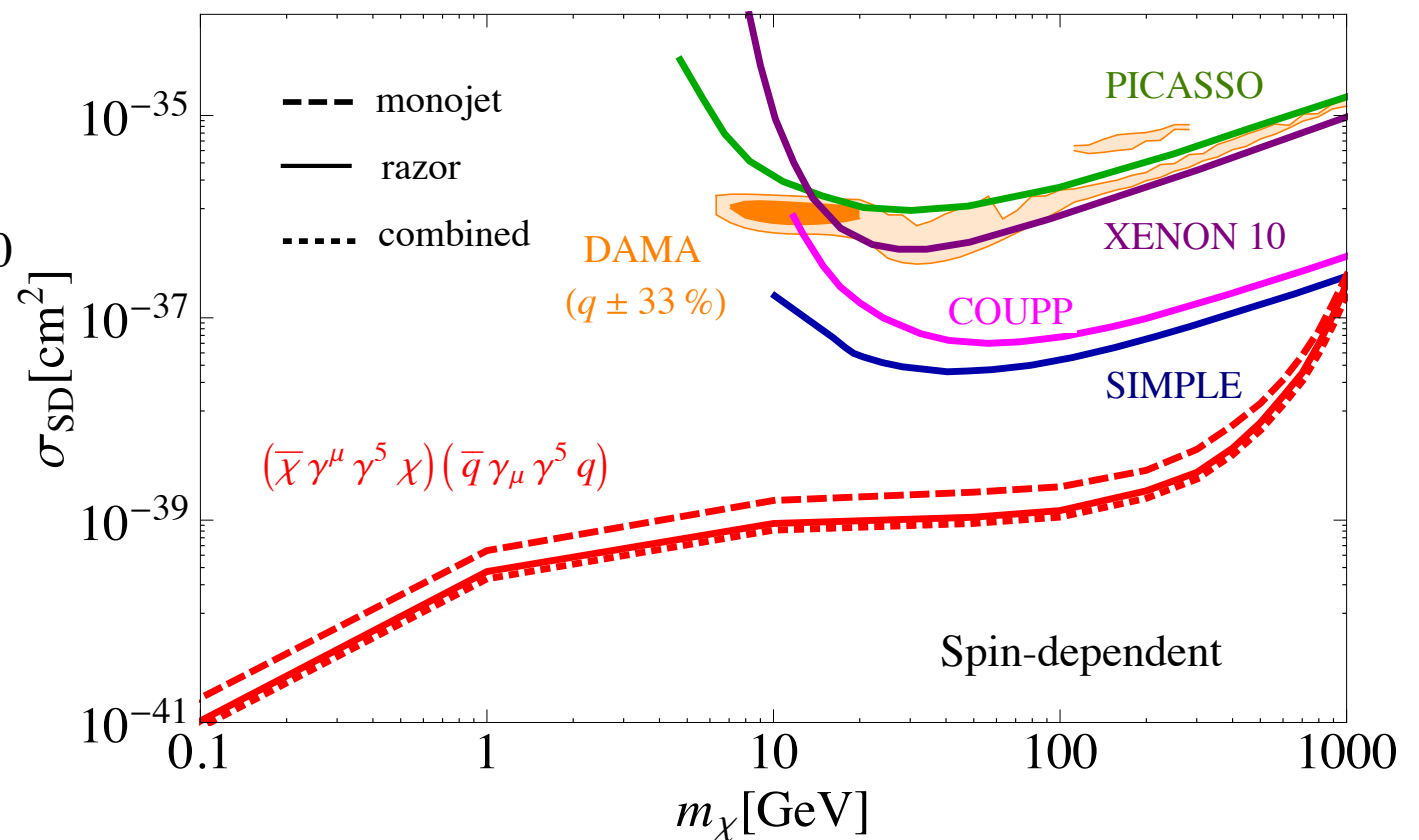
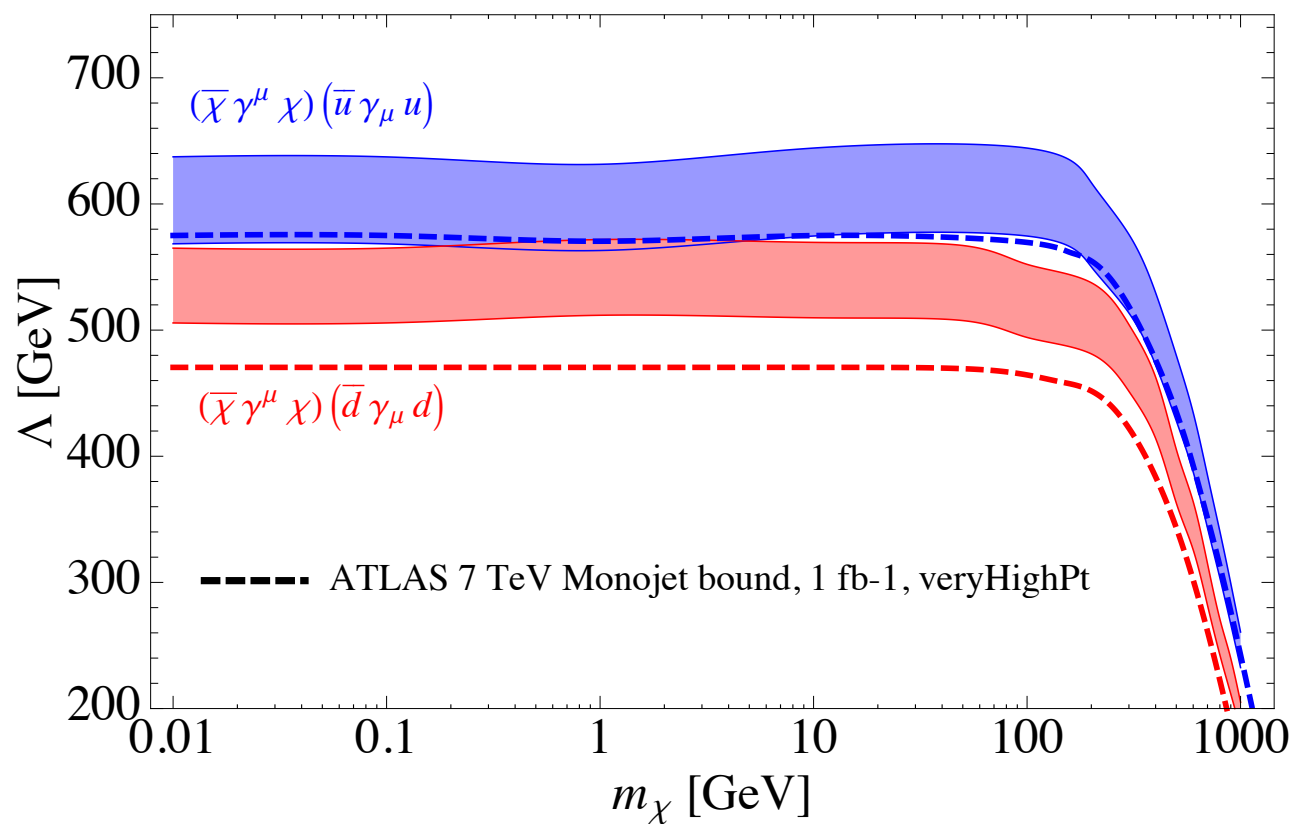
# Unitarity

- \* For heavy dark matter the EFT is not valid:



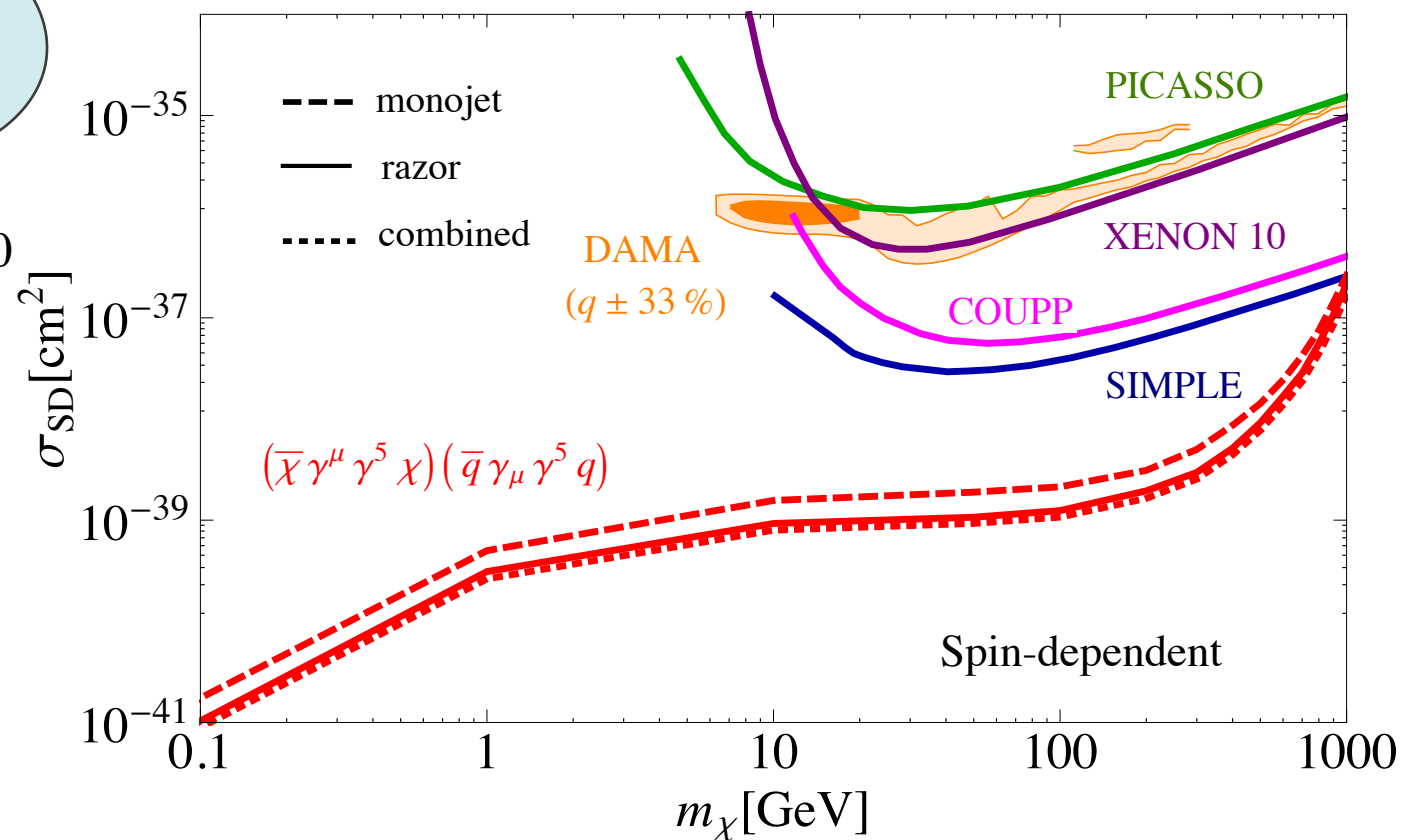
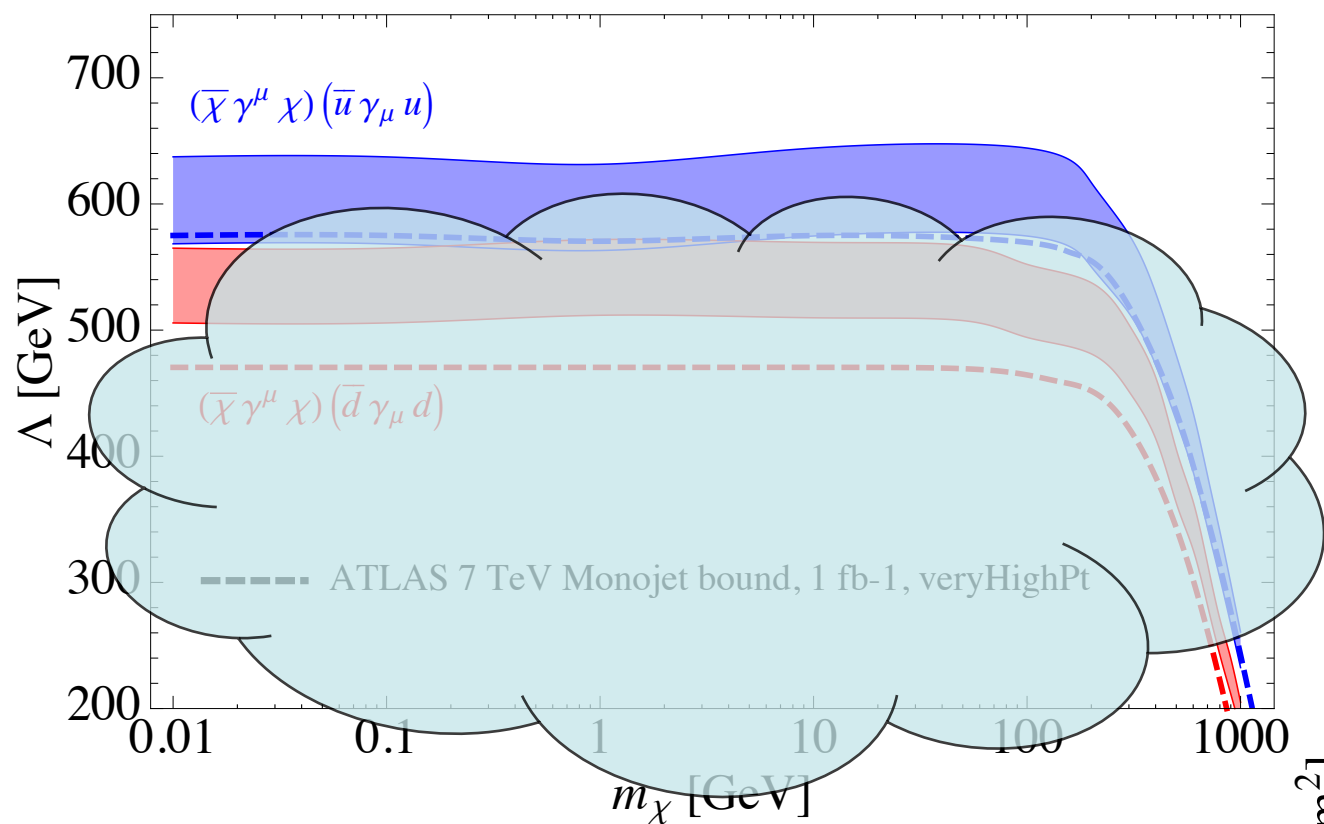
# Unitarity

\* Strictly speaking, the EFT analysis is not valid at small  $\Lambda$ .



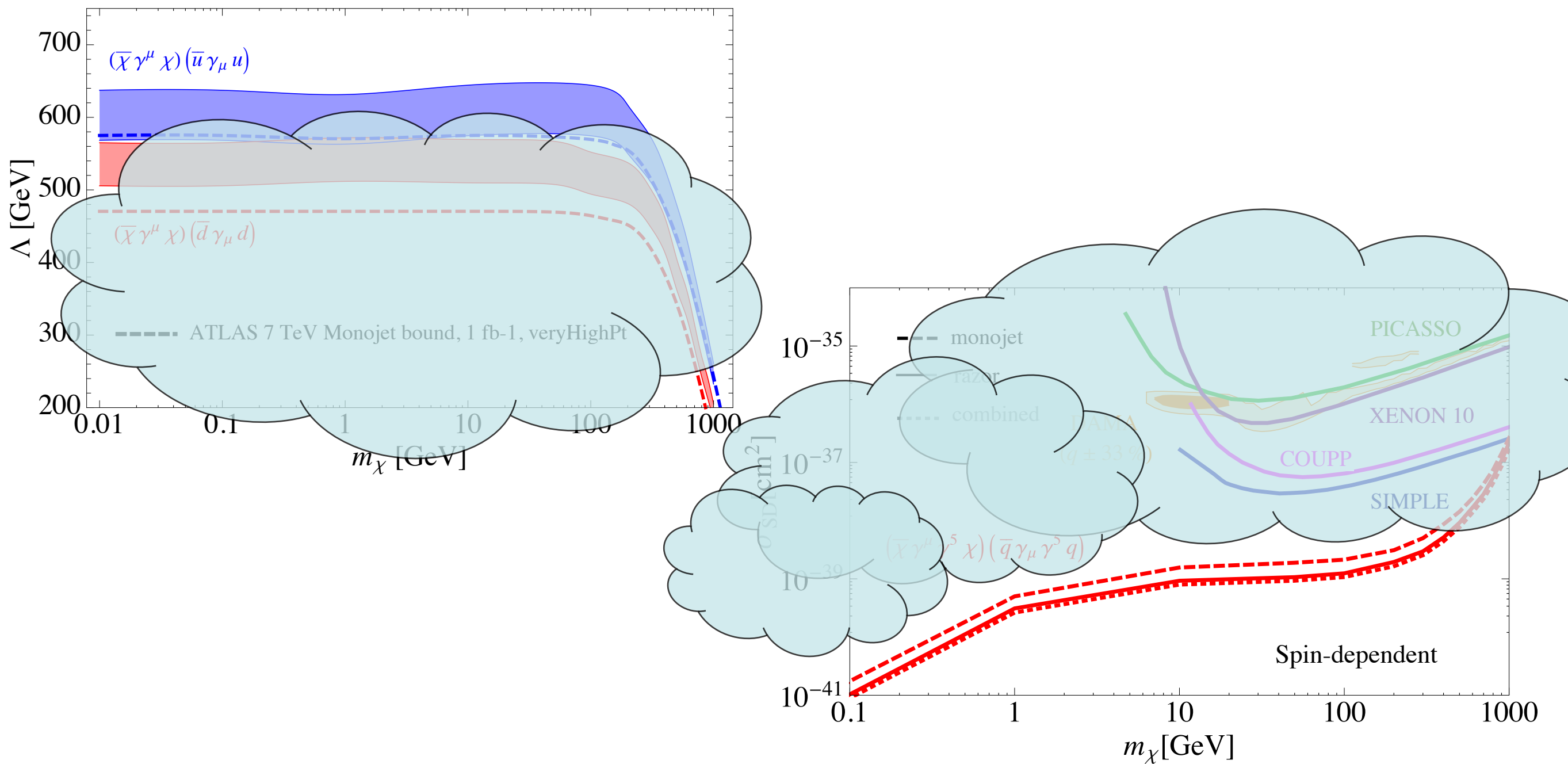
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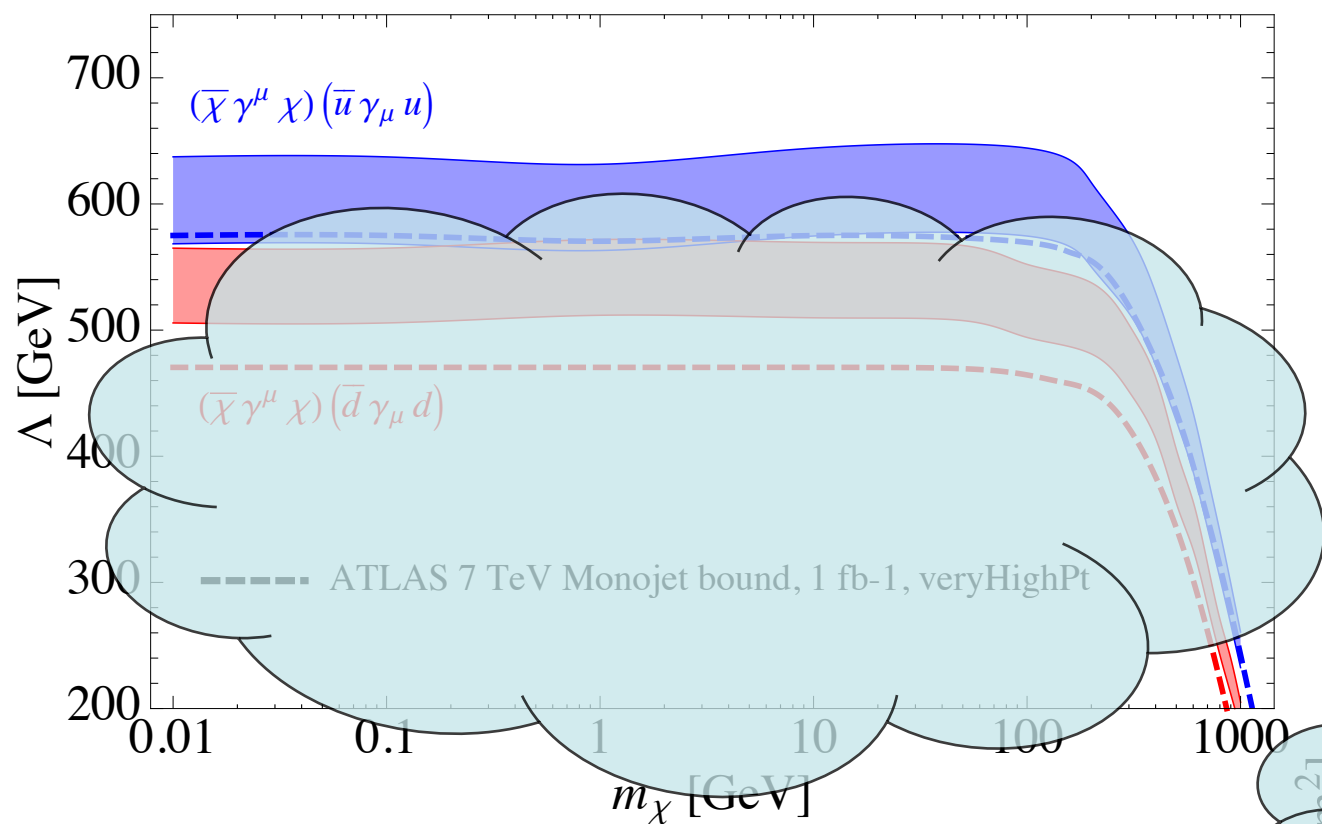
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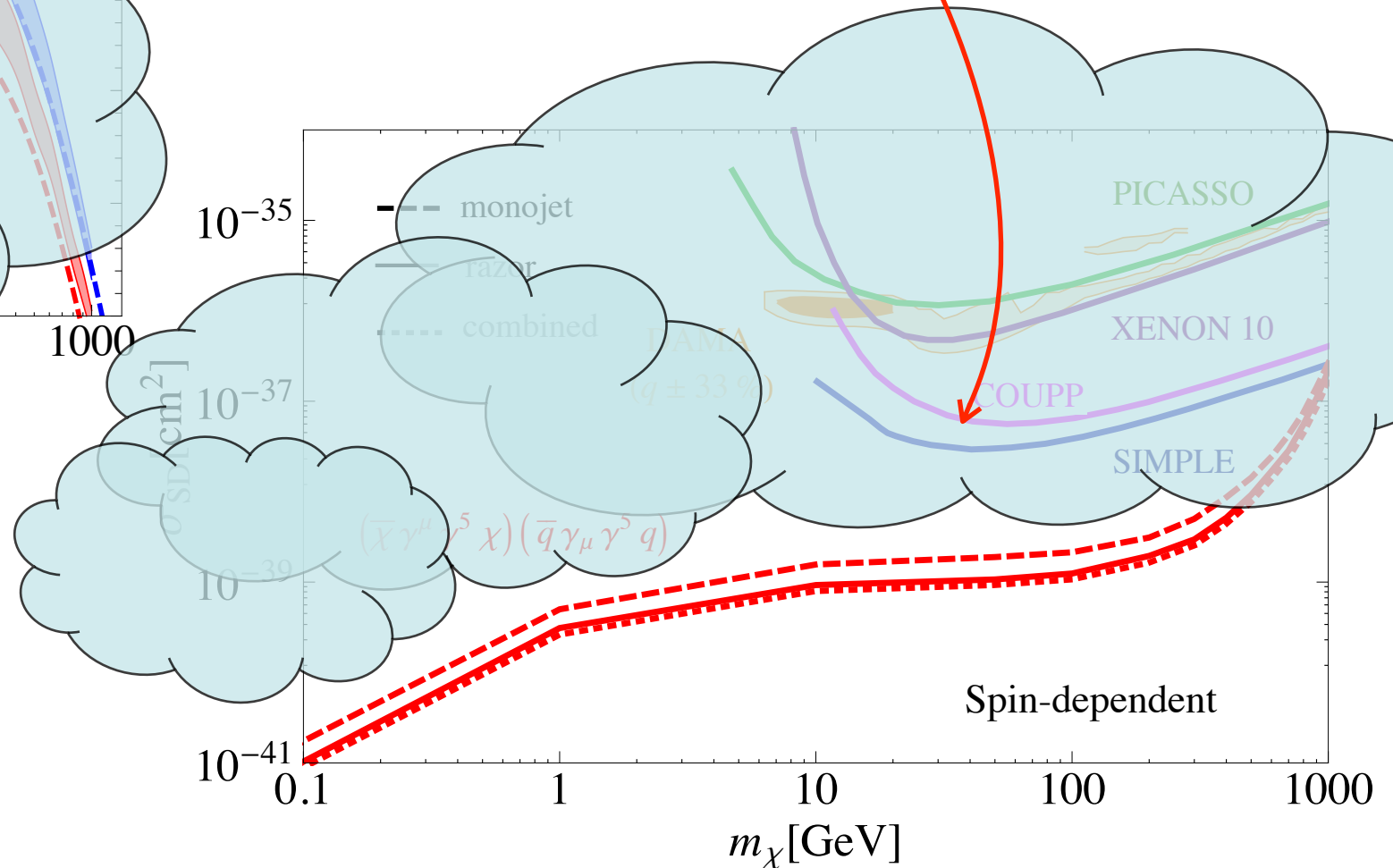


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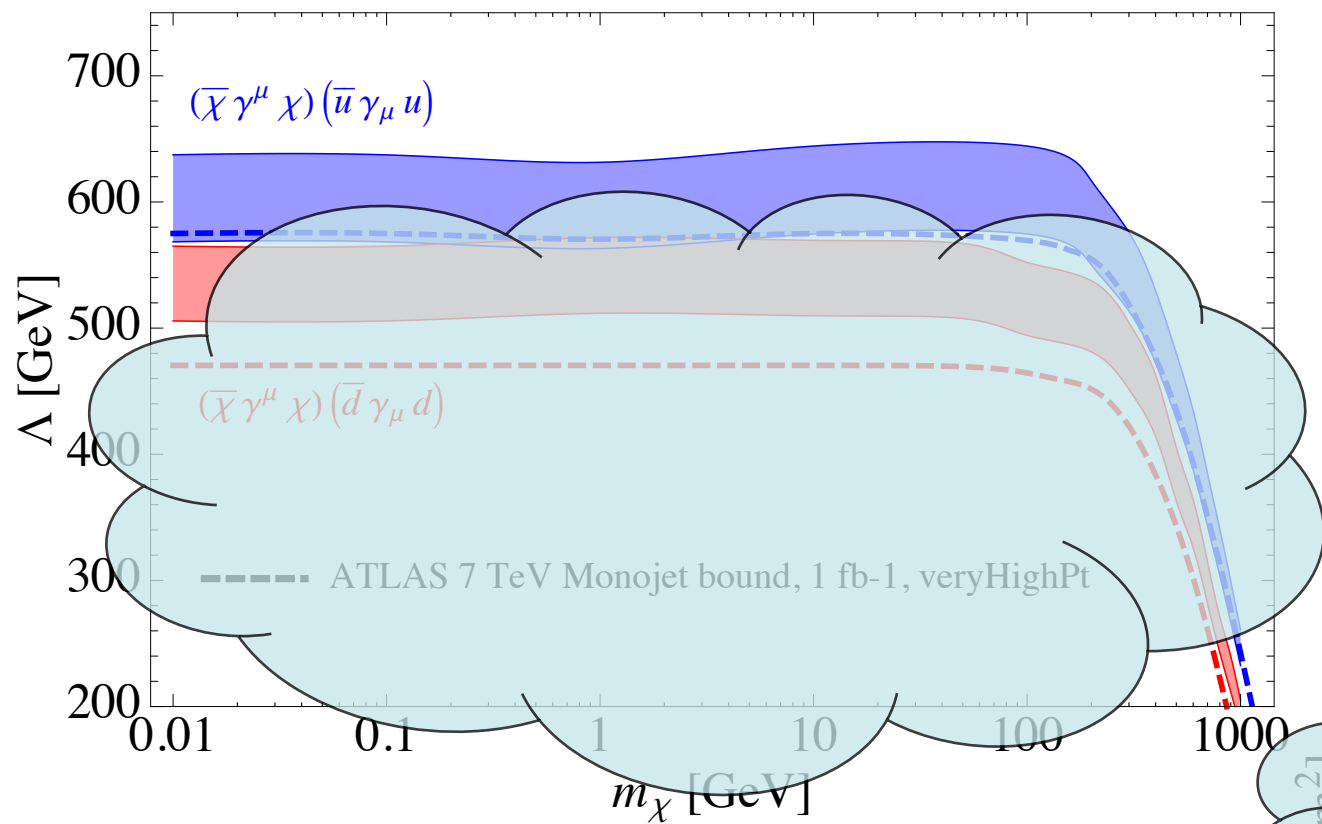


Does this mean a theory that lives here is alive?

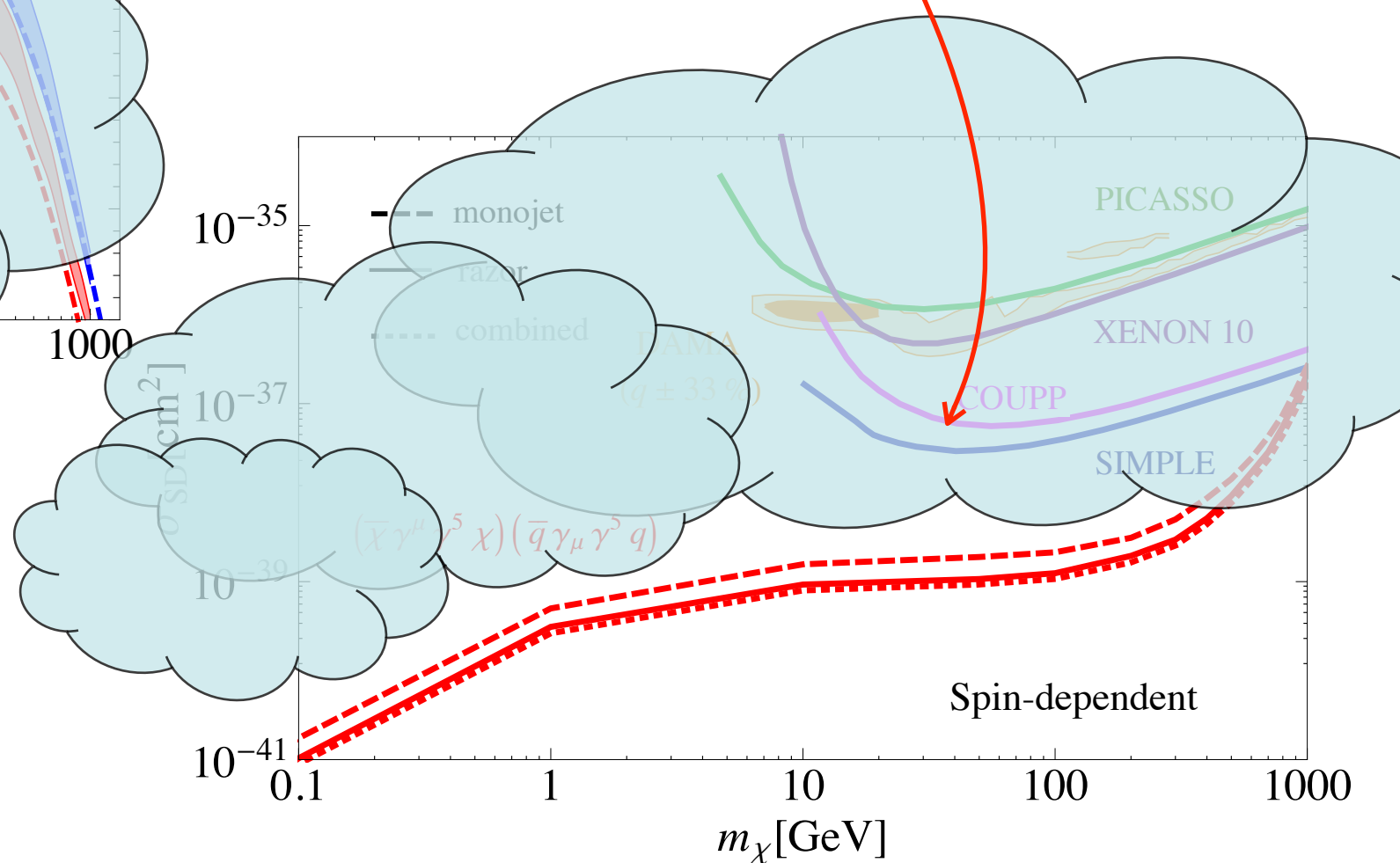


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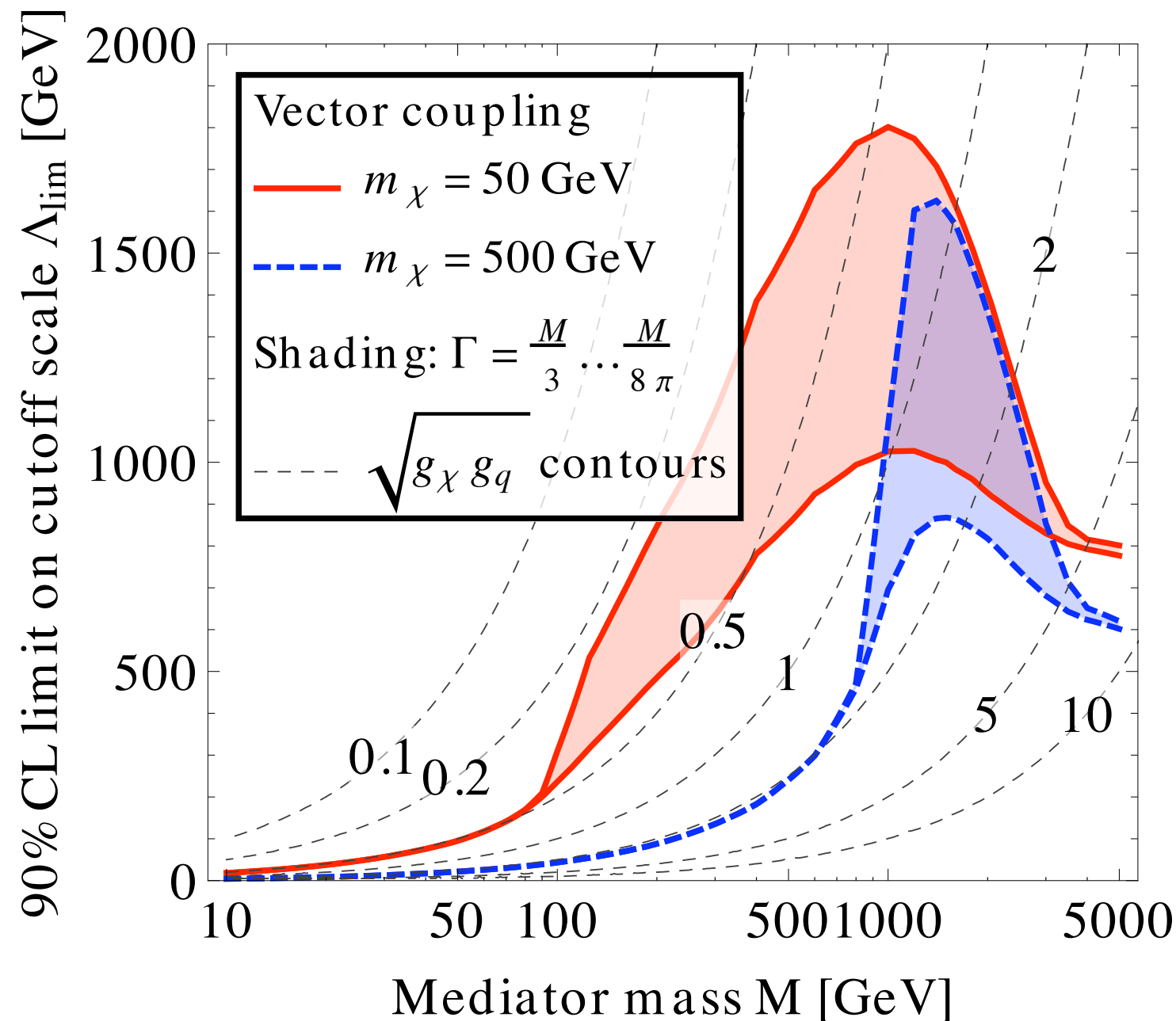


No.  
In fact it may be even more dead.



# Light Mediator

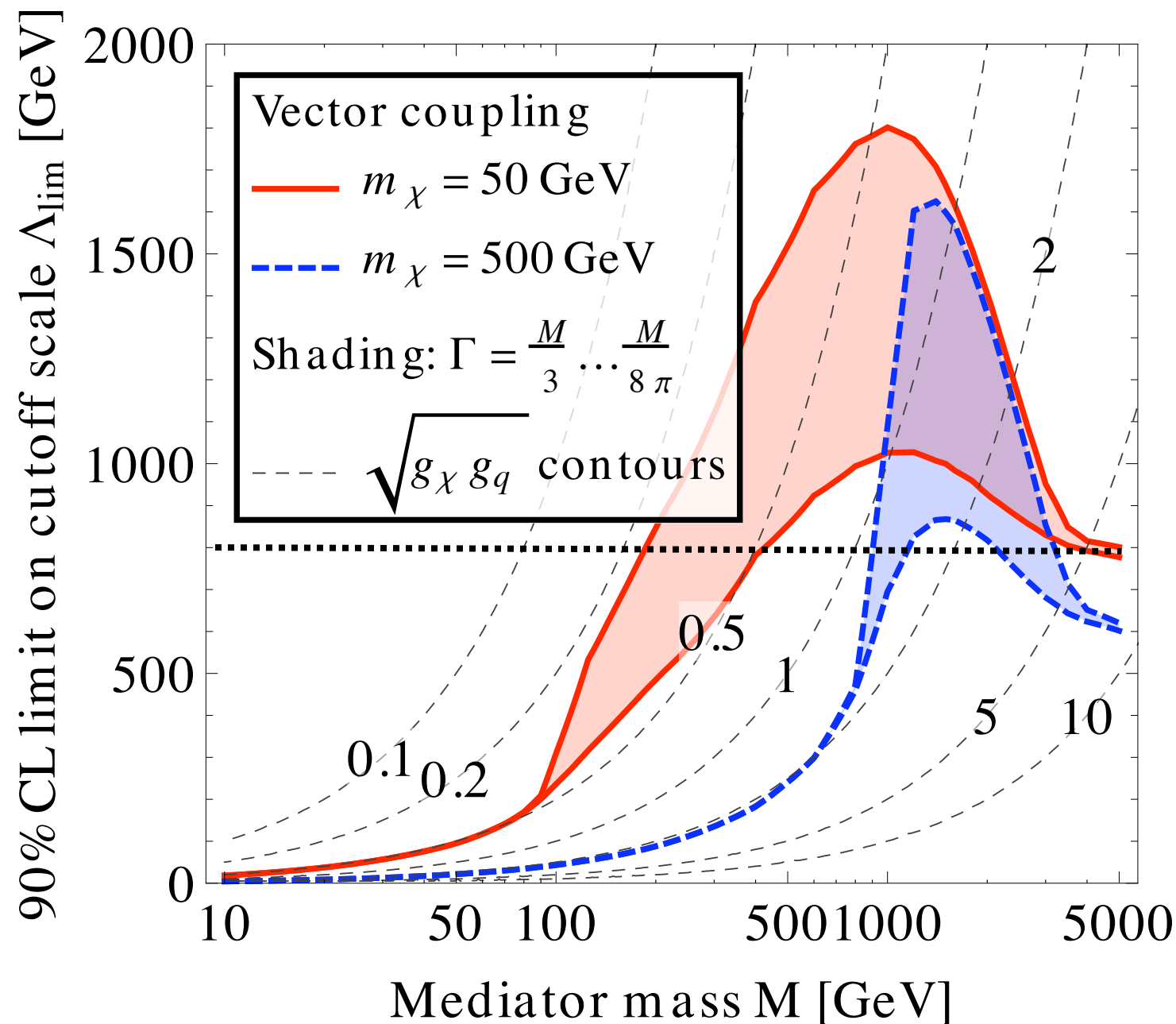
\* The limit become better before it gets worse:



EFT limits are **conservative** so long as the mediator is above a few hundred GeV (and the mediator decays to DM).

# Light Mediator

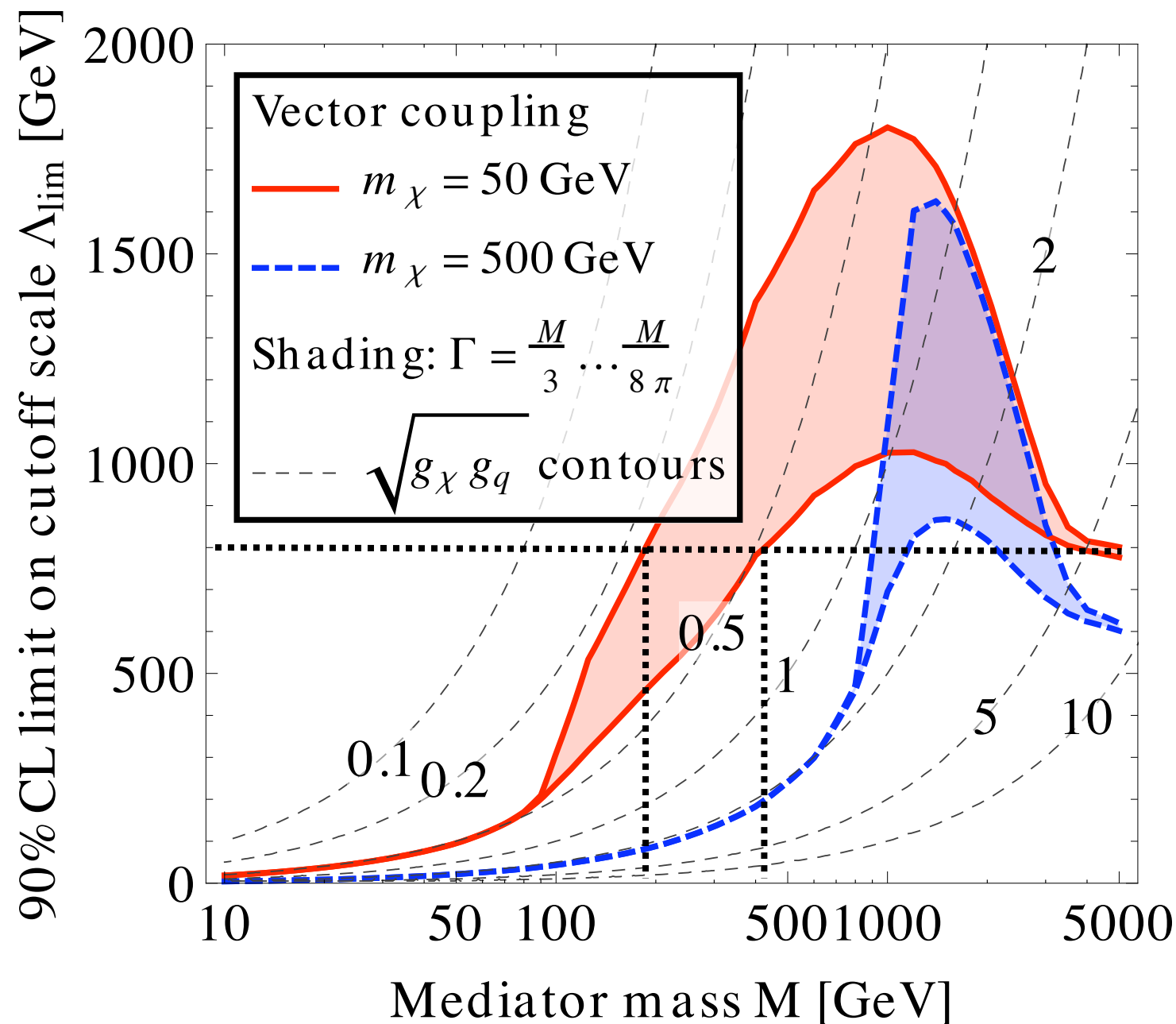
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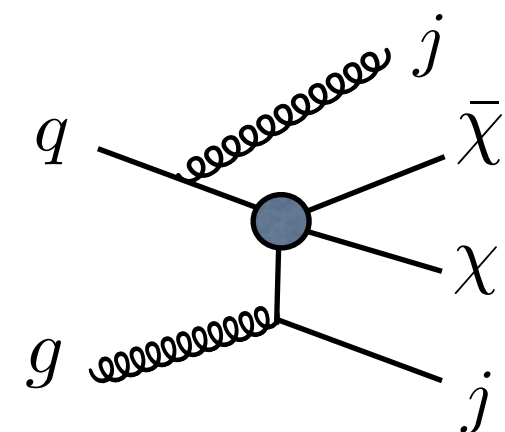
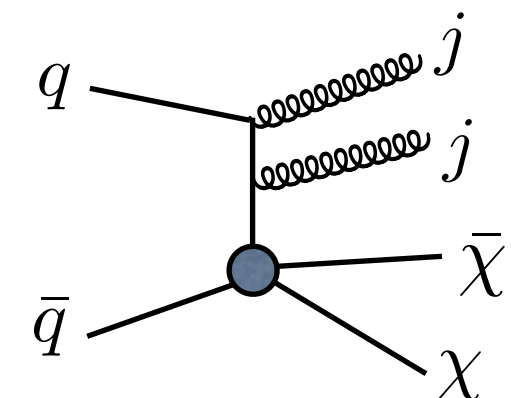
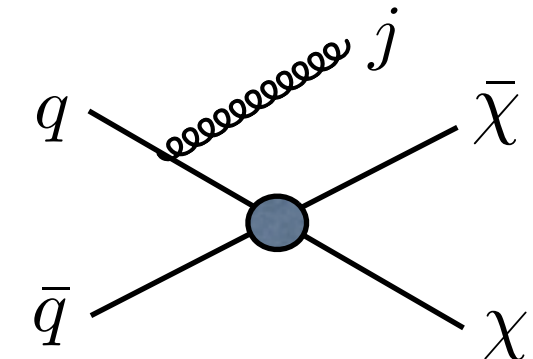


# Questions

- \* How can we improve these bounds?
- \* Are there other LHC searches that are relevant?
- \* What theoretical assumptions go into the bounds?
- \* How can collider bounds be evaded?

# Outline

- \* Setup - operators and mediators
- \* Mono-Jets & Mono-photons.
- \* Multi-jets plus MET.
  - o Handles on S vs B.
  - o Razor analysis.
- \* Note on validity of EFT.
- \* DM Higgs interplay?

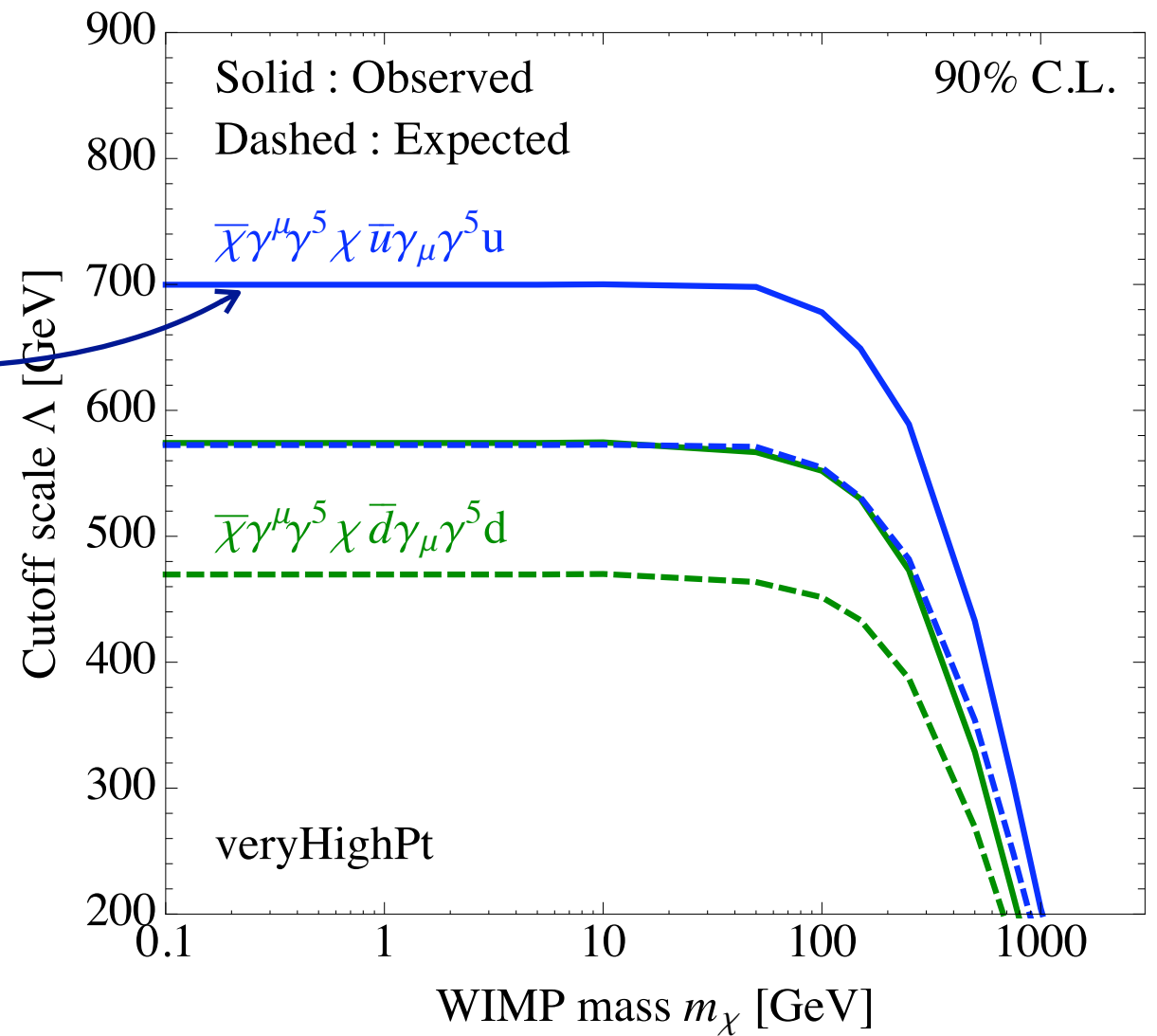
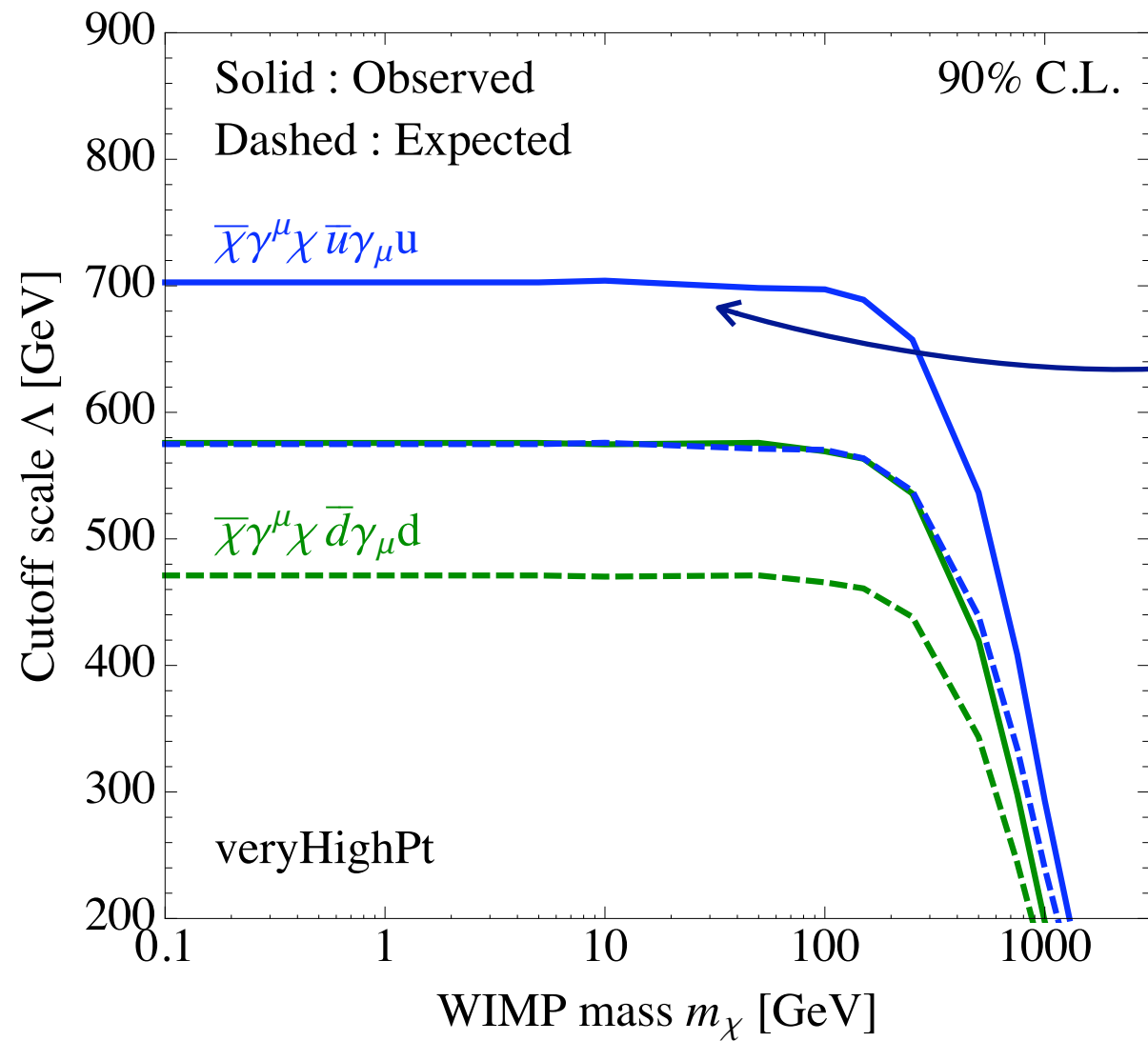


**Mono-Jet**



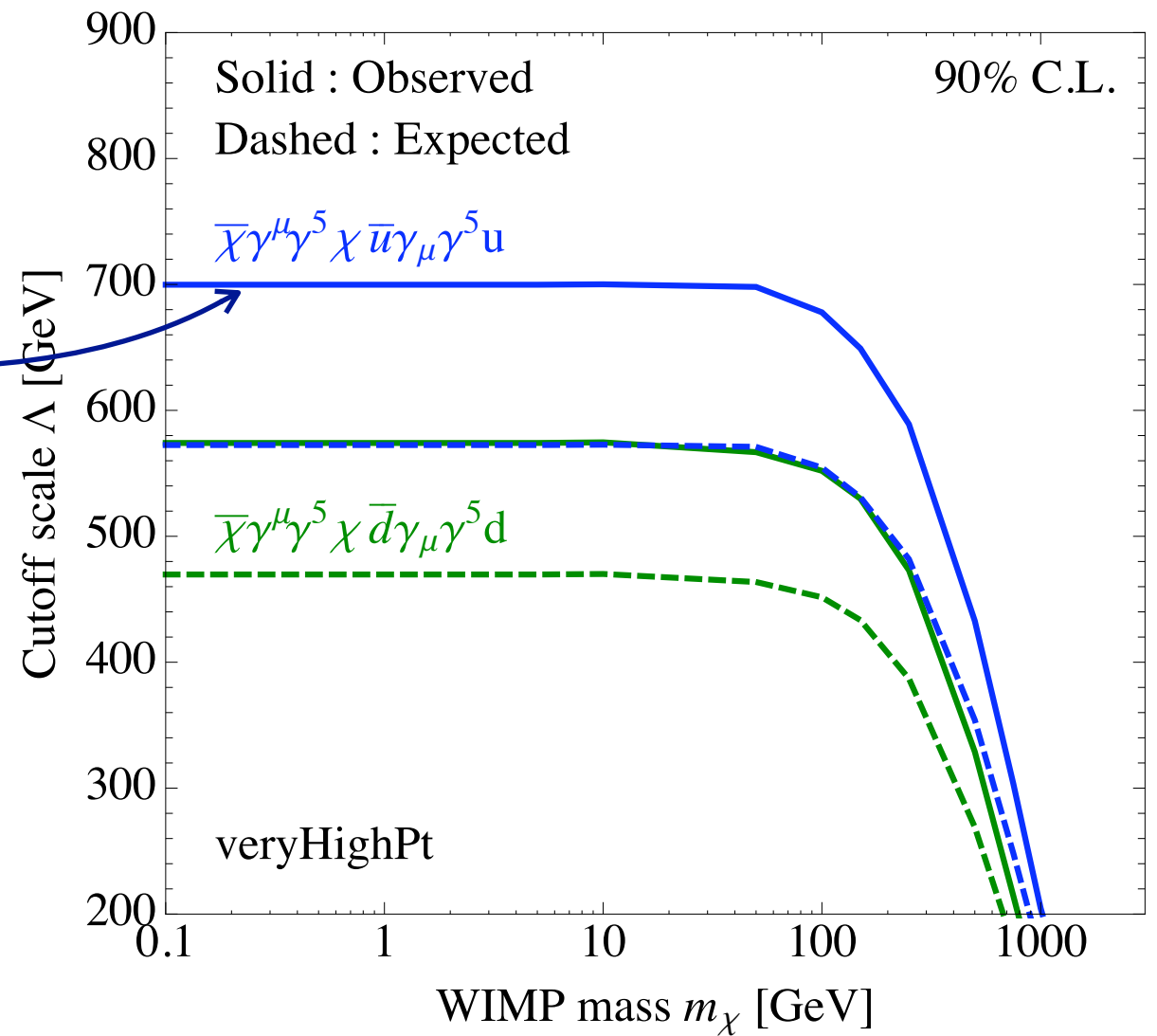
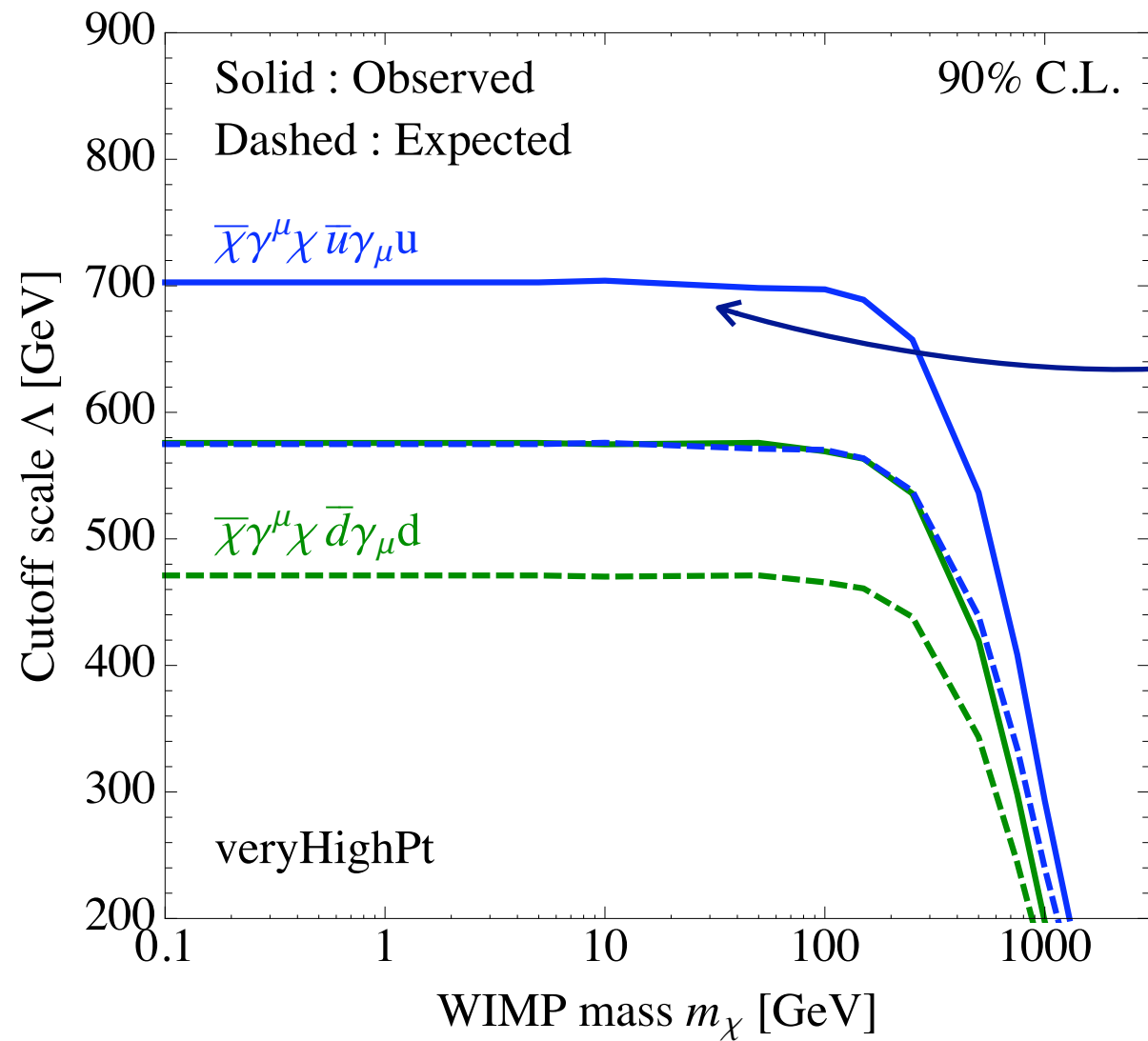
# Limits on $\Lambda \equiv \frac{M}{\sqrt{g_\chi g_1}}$ :

*same limit for SI and SD*



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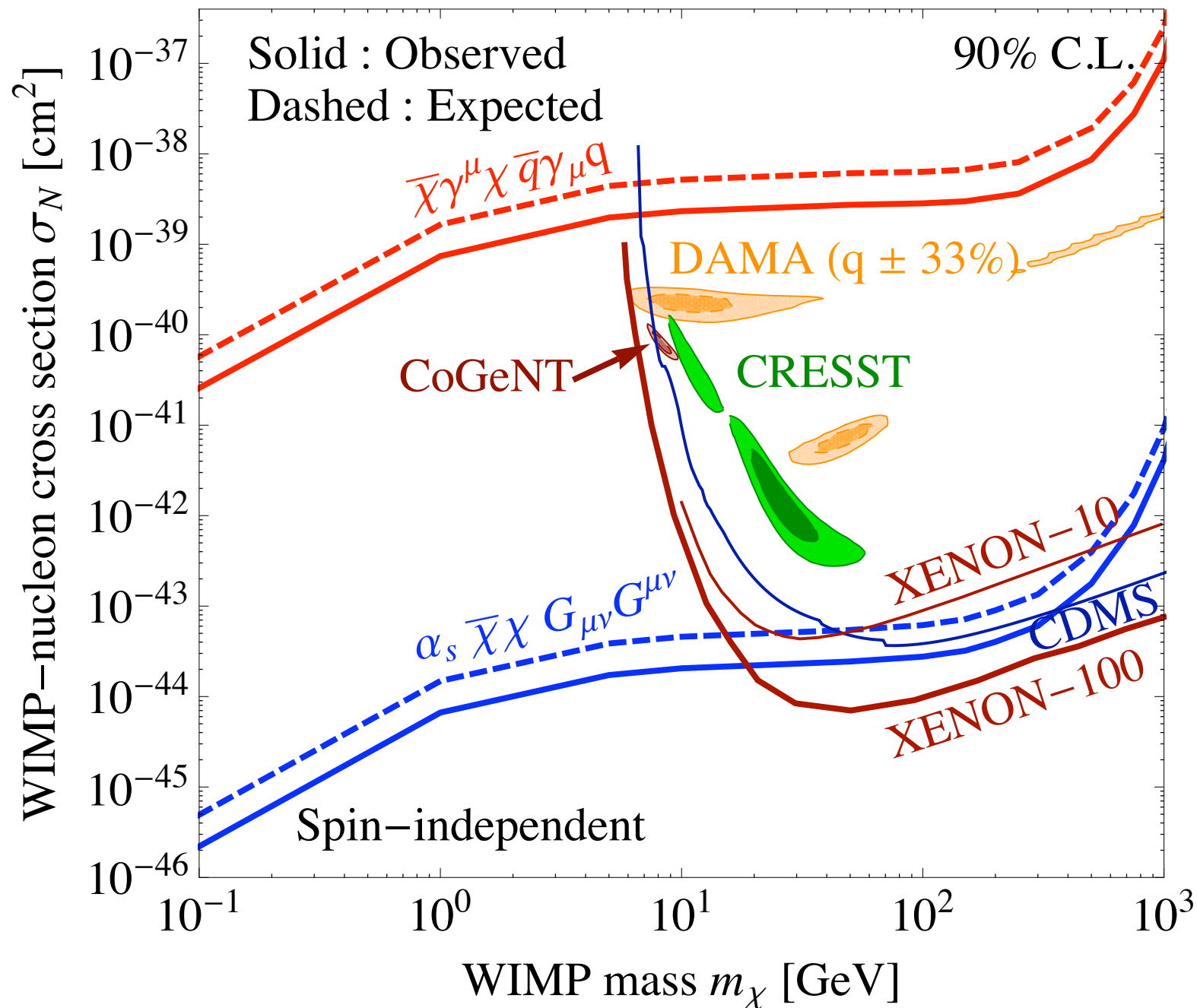
*The limit is flat up to ~200 GeV.  
Goes all the way to zero.*

# SI Limit

$$\sigma_1^{Nq} = \frac{\mu^2}{\pi\Lambda^4} B_{Nq}^2,$$

$$\sigma_2^{Nq} = \frac{\mu^2}{\pi\Lambda^4} f_{Nq}^2,$$

ATLAS 7TeV,  $1\text{fb}^{-1}$  VeryHighPt



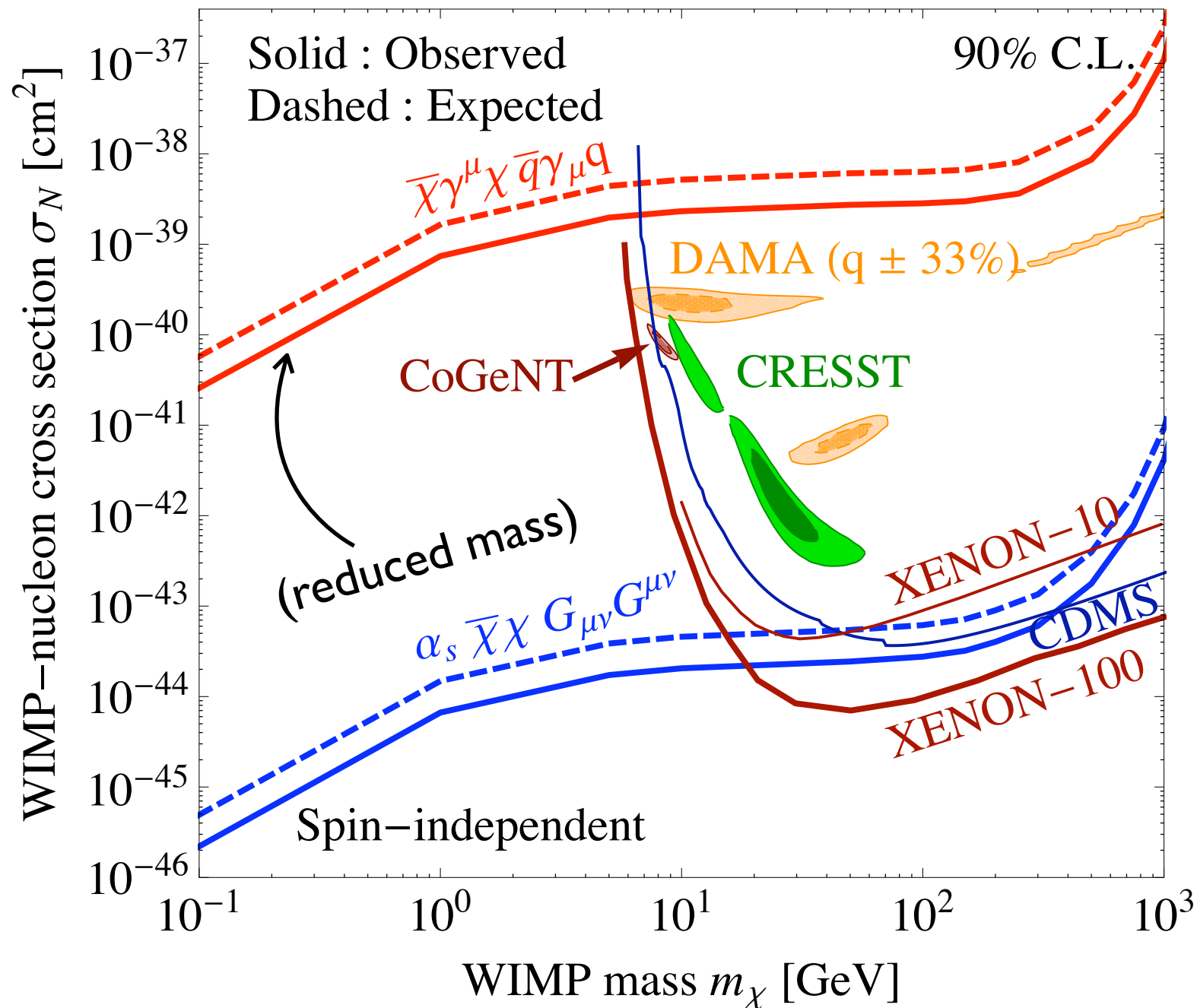
**Best limit  
at low mass**

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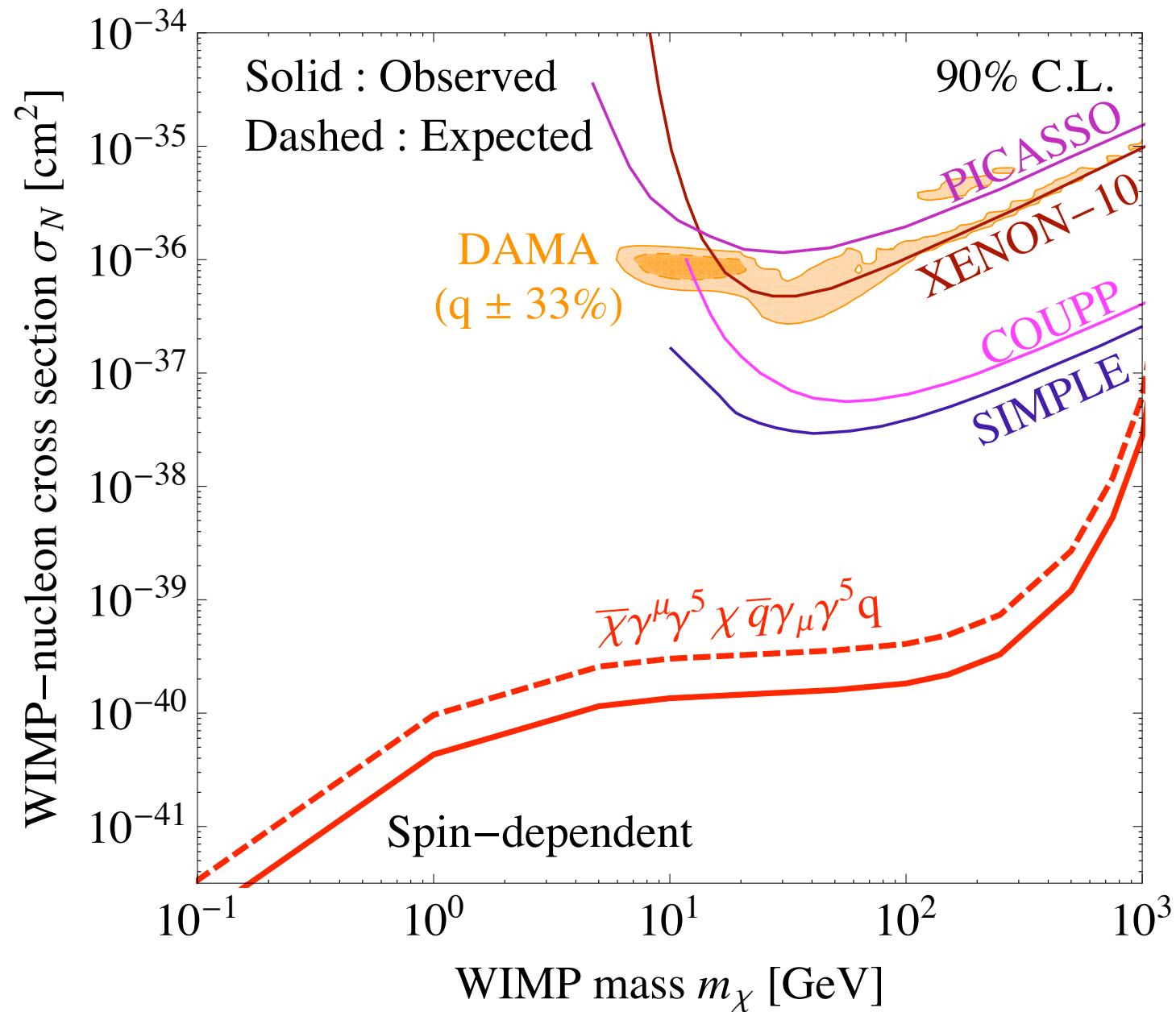
ATLAS 7TeV,  $1\text{fb}^{-1}$  VeryHighPt



**Best limit  
at low mass**

# SD Limit

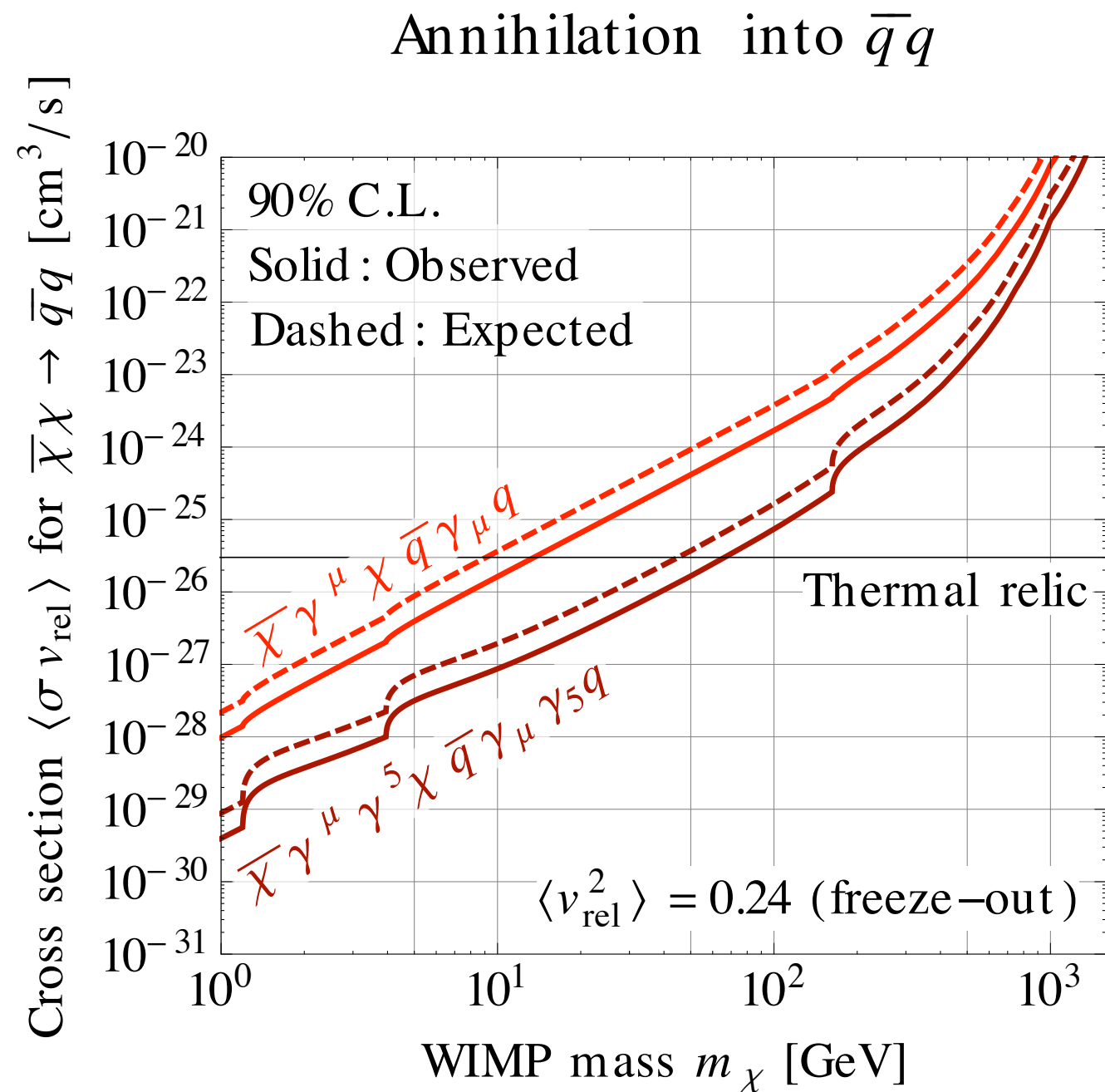
ATLAS 7TeV,  $1\text{fb}^{-1}$  VeryHighPt



**Best spin  
dependent  
limit.**

# Annihilation

- \* A minimal light thermal relic is ruled out:



# Can we evade bounds?

\* Lets fix  $\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$  and lower  $M$ .

The couplings must be decreased to compensate.

\* Then for very small  $M$  we get to the regime where

$$\sigma_{1j} \sim \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2}$$

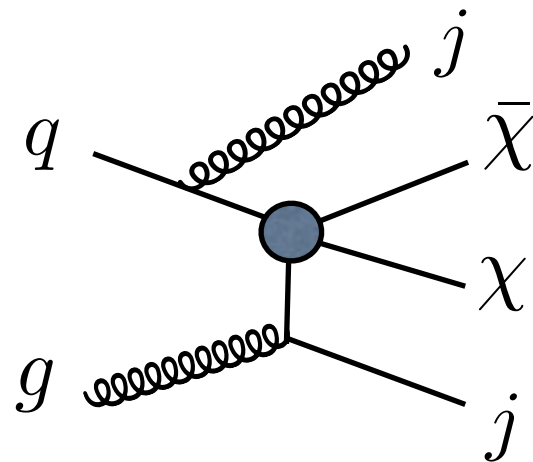
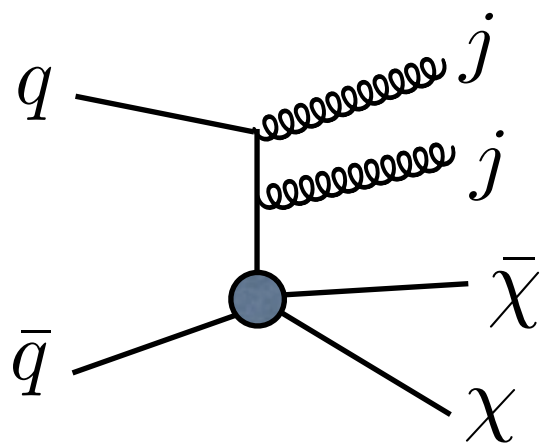
\* The cross section drops as  $M^4$ .

\* **Theories with light mediators always evade the collider bound.**

(**But...**more on the intermediate regime later)

# Jets plus MET

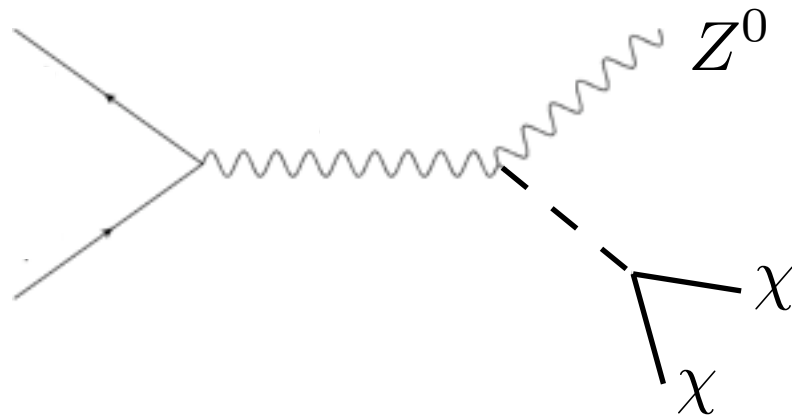
(with Fox, Primulando and Yu, 1203.1662 )





# Higgs Portal DM

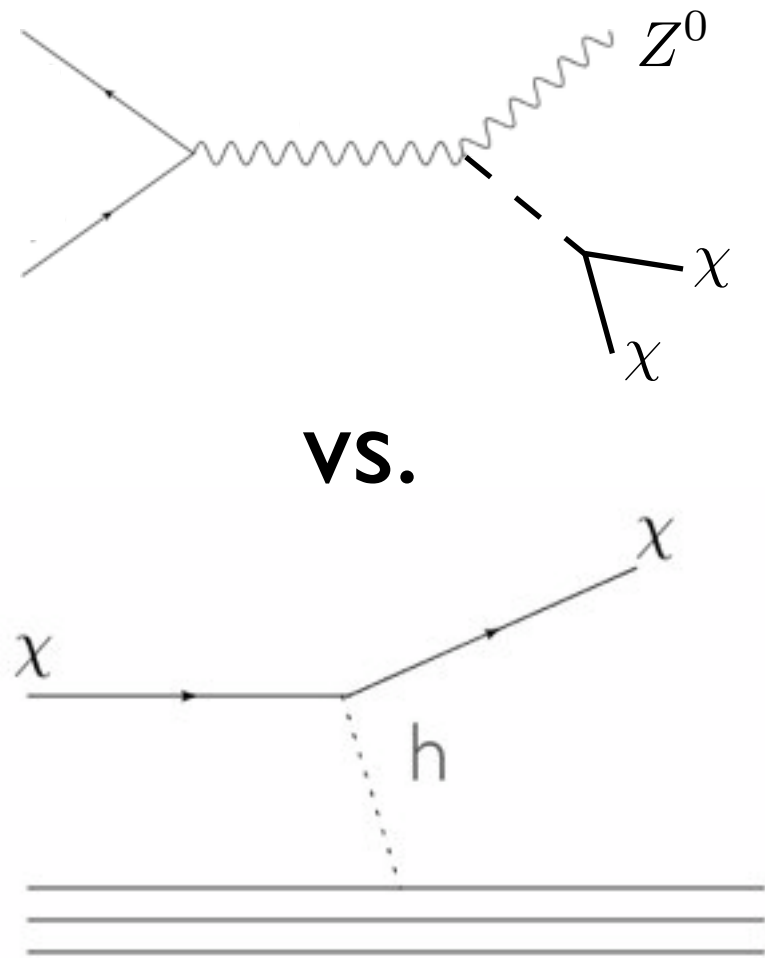
- \* For specific models, we can probe the identity of the mediator with other mono-somthings.
- \* In many models DM couples via the **Higgs**. **Mono-Z** (and **VBF**) may be sensitive to this.



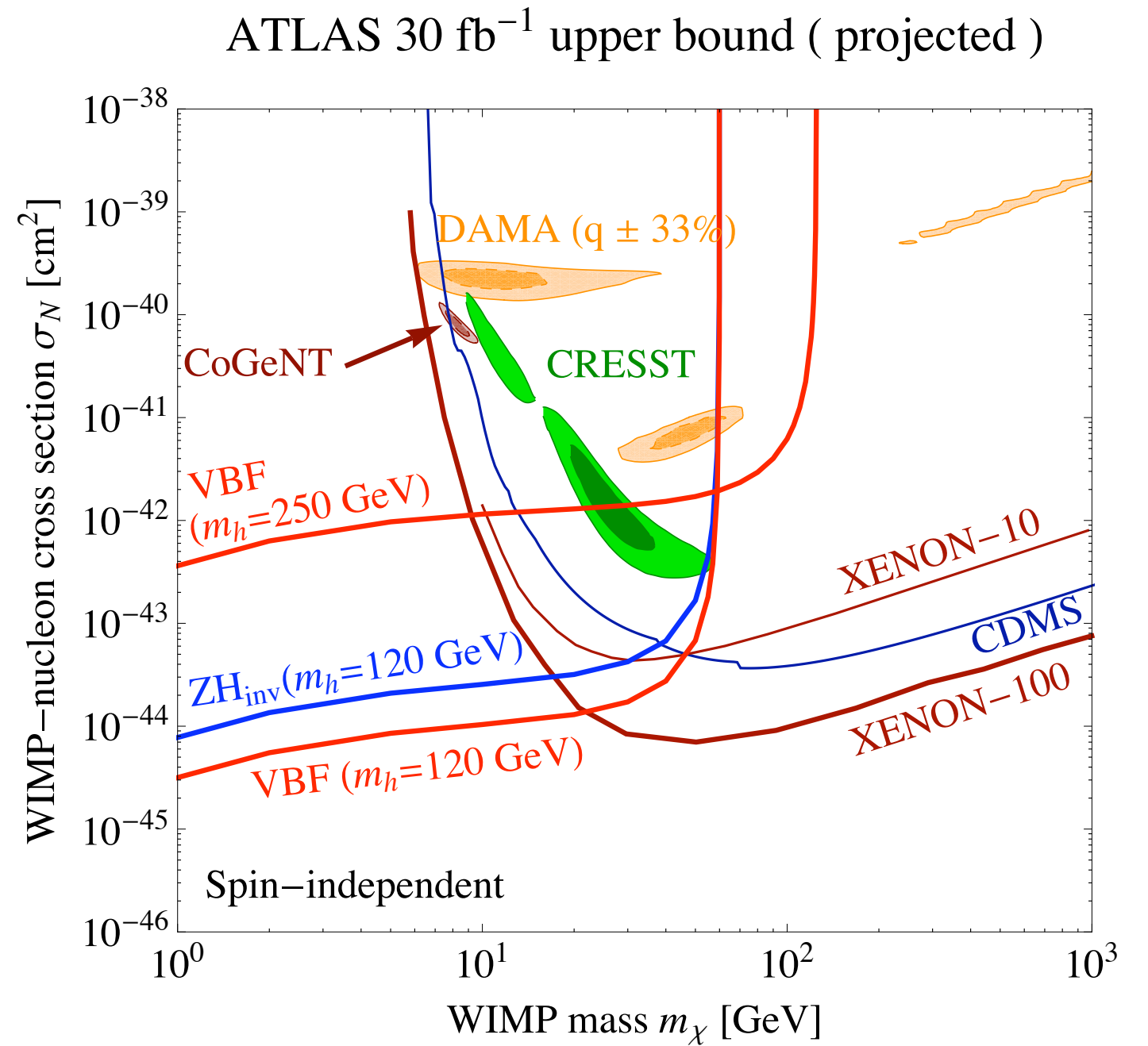
Invisible Higgs searches can be interpreted as “direct detection” experiments!

*A Characteristic Higgs Channel can confirm Higgs mediation!*

# Higgs Mediator



**Direct detection is parametrically smaller!**



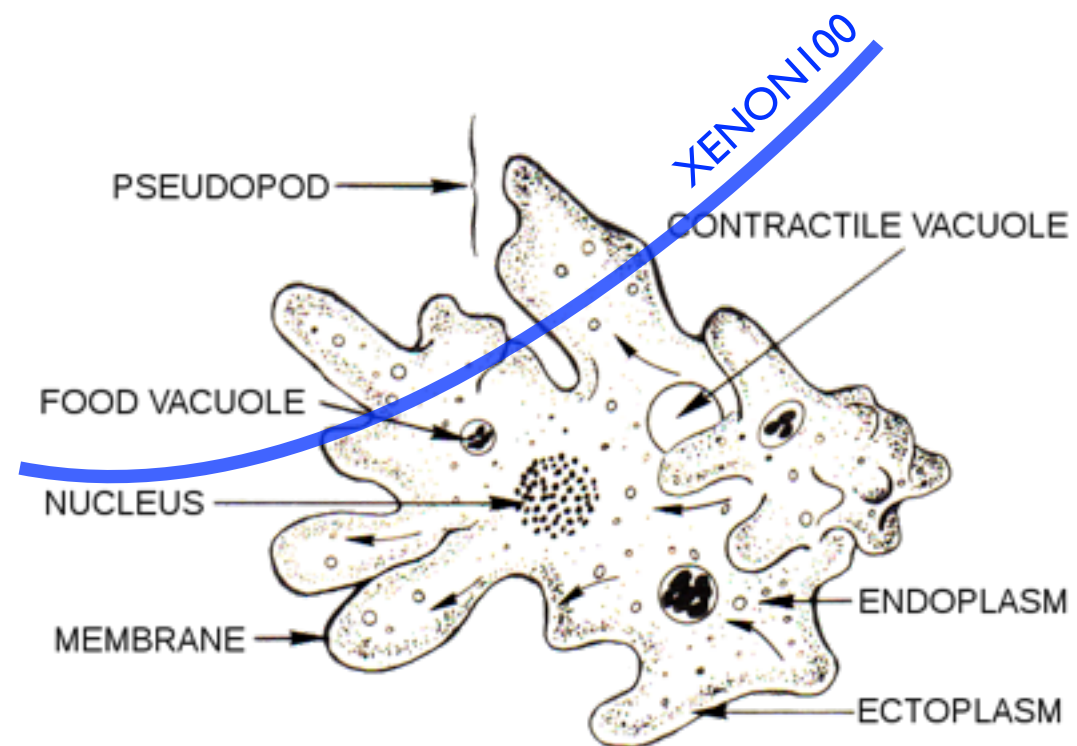
Fox, RH, Kopp and Tsai

# To Conclude:

**Colliders are placing competitive and complementary bounds to direct and to indirect detection:**

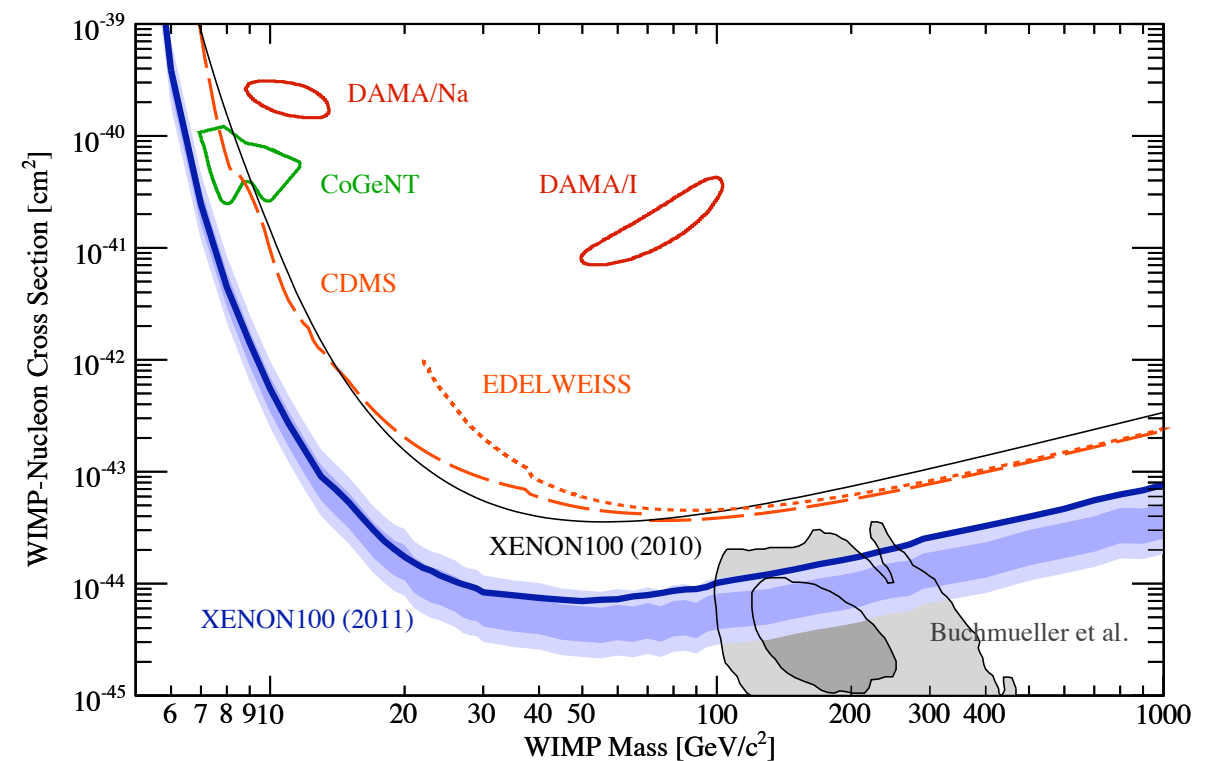
- \* The ~~Tevatron~~<sup>LHC</sup> is the world record holder for light dark matter and for spin dependent.
- \* Dedicated CDF, CMS, and ATLAS **mono-jet** studies are out (or underway). CMS **mono-photon** too.
- \* Inclusive Jets plus MET studies may have additional discriminating power (Razor).
- \* **Higgs** and **DM** play nicely together.

# Deleted Scenes



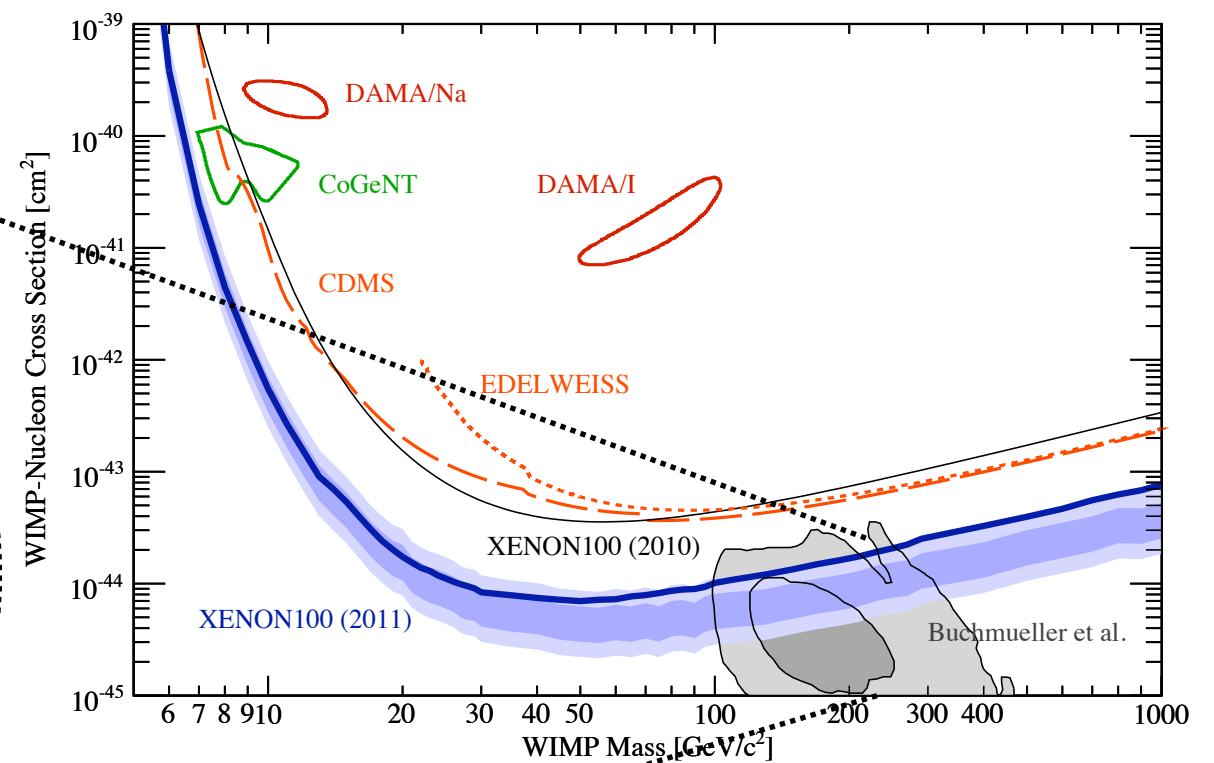
# Collider Connections?

- \* DM experiments and colliders are often said to be related *in a specific framework* (SUSY).



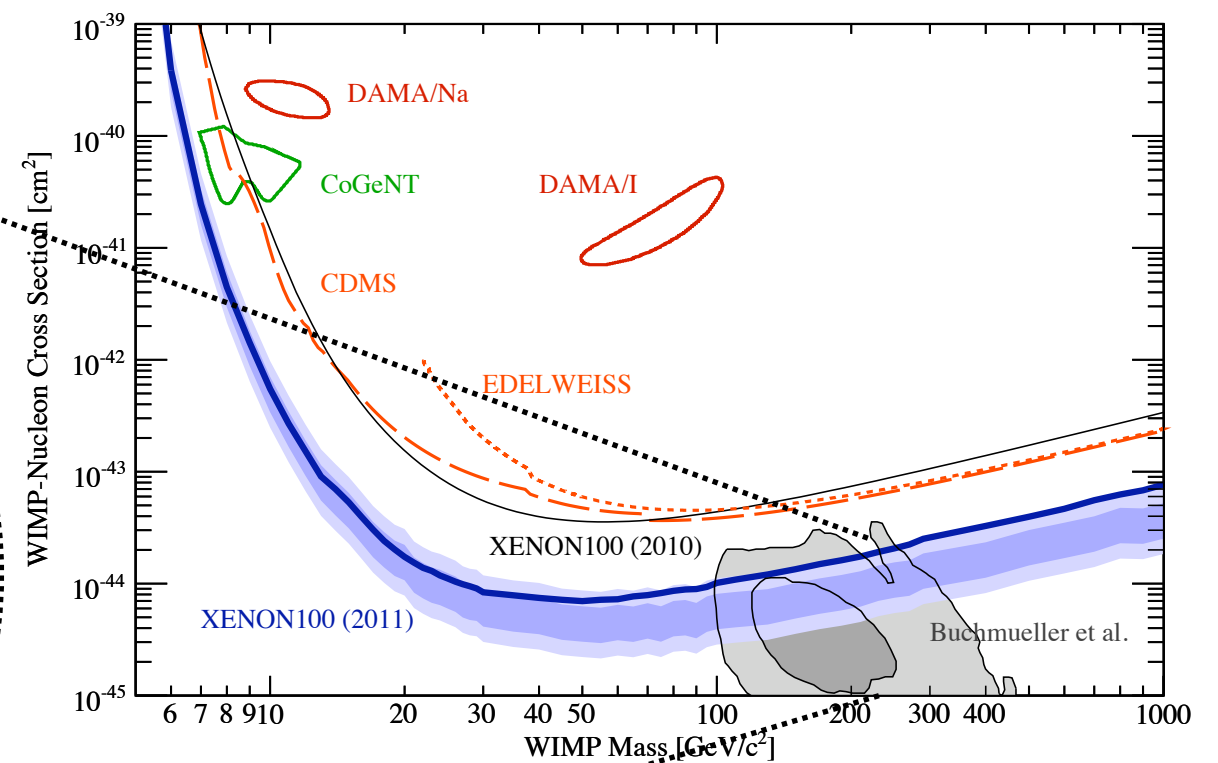
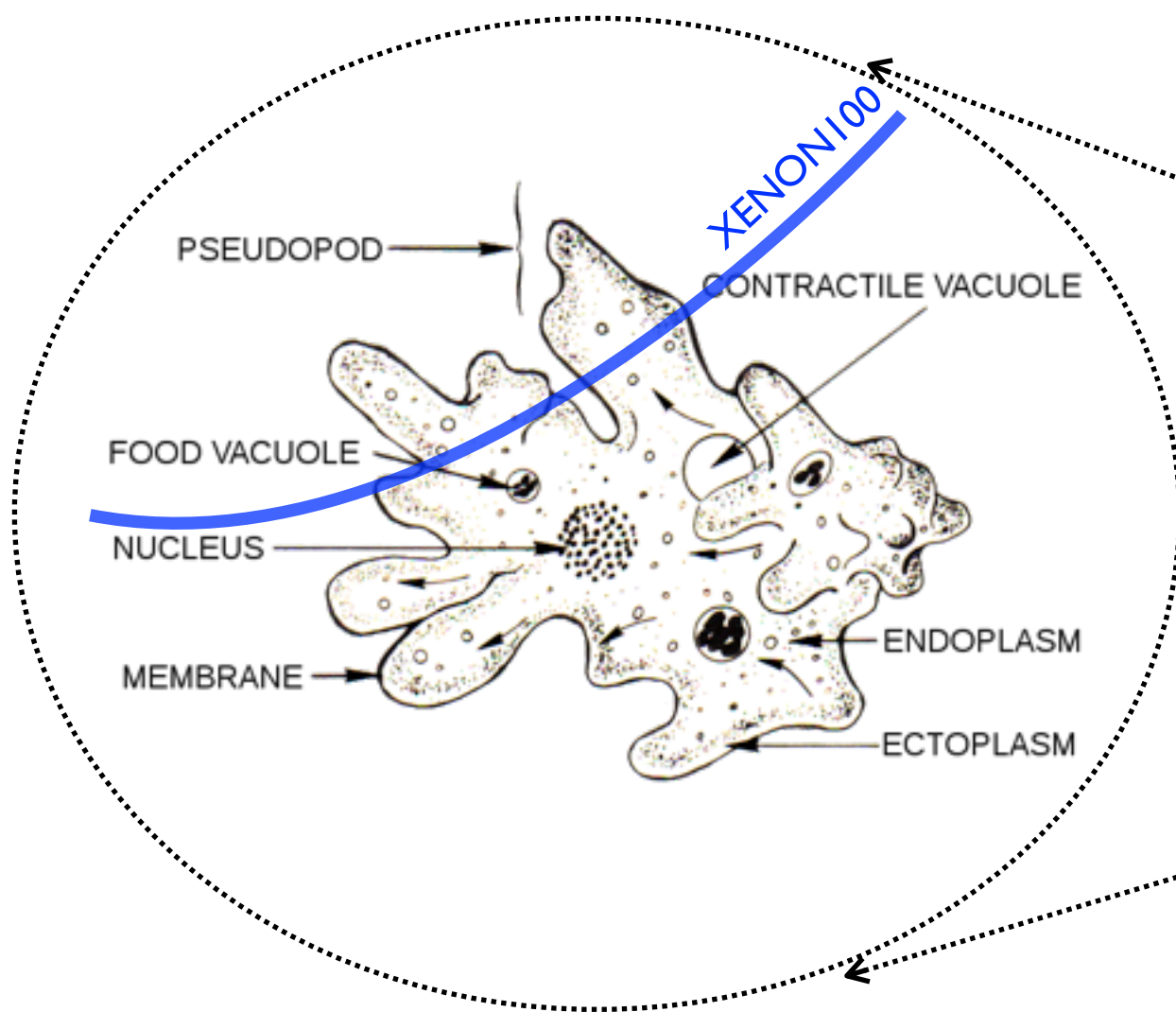
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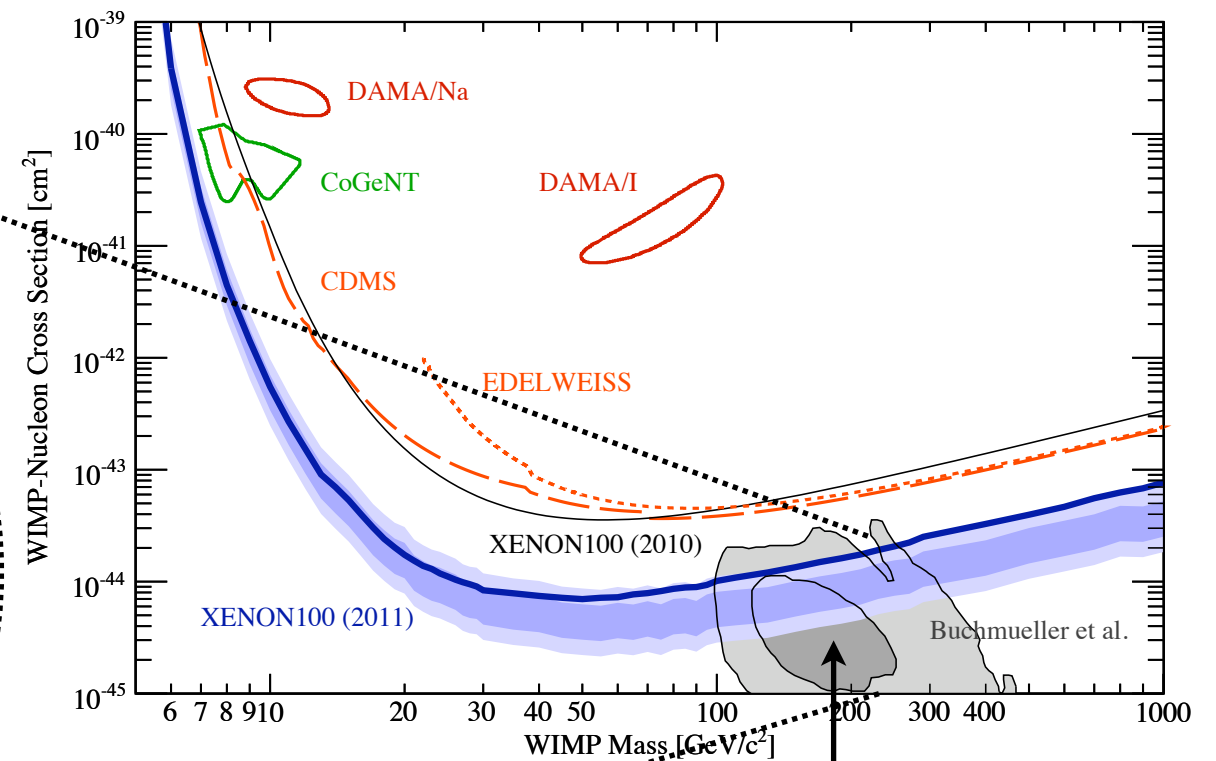
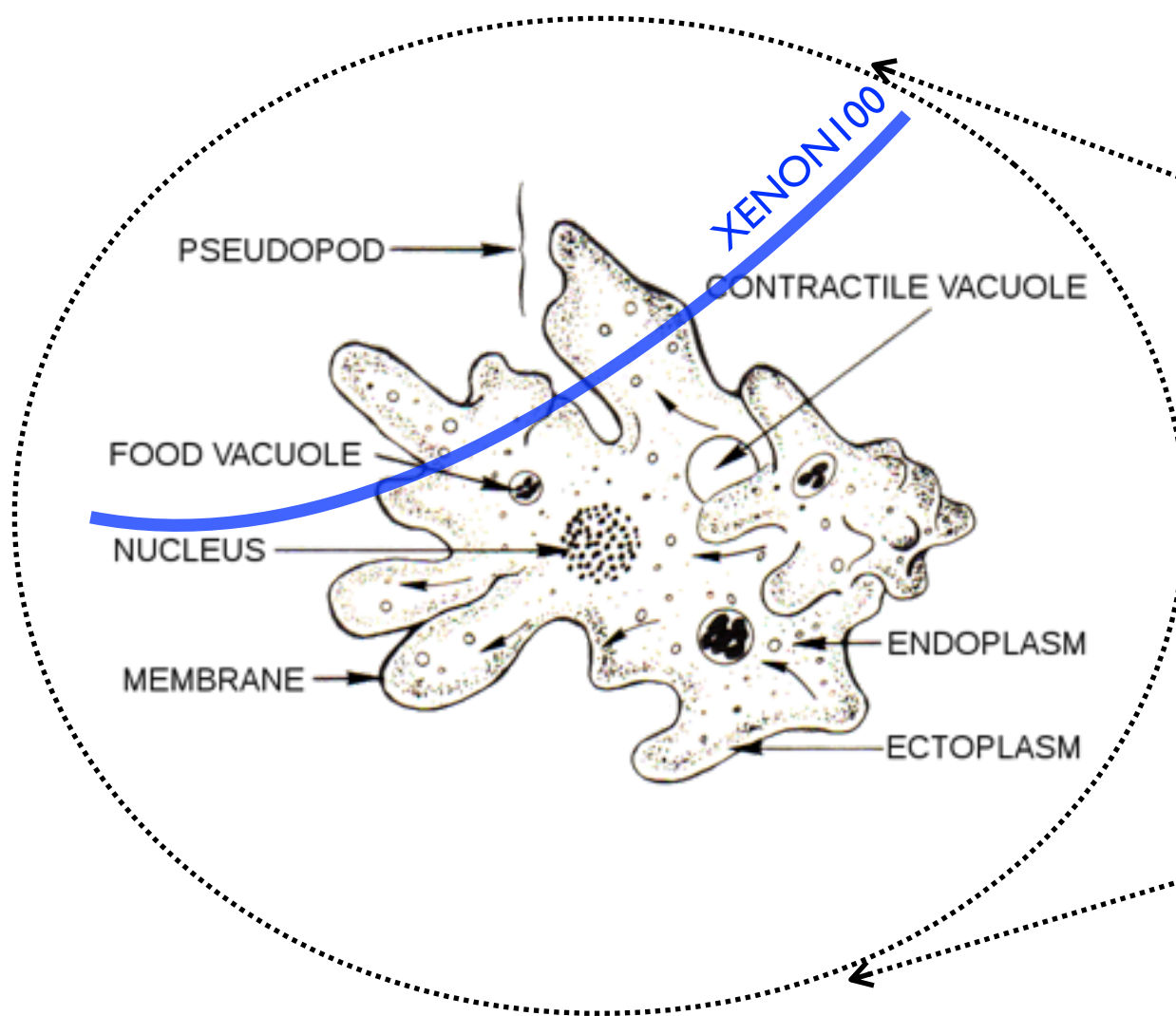
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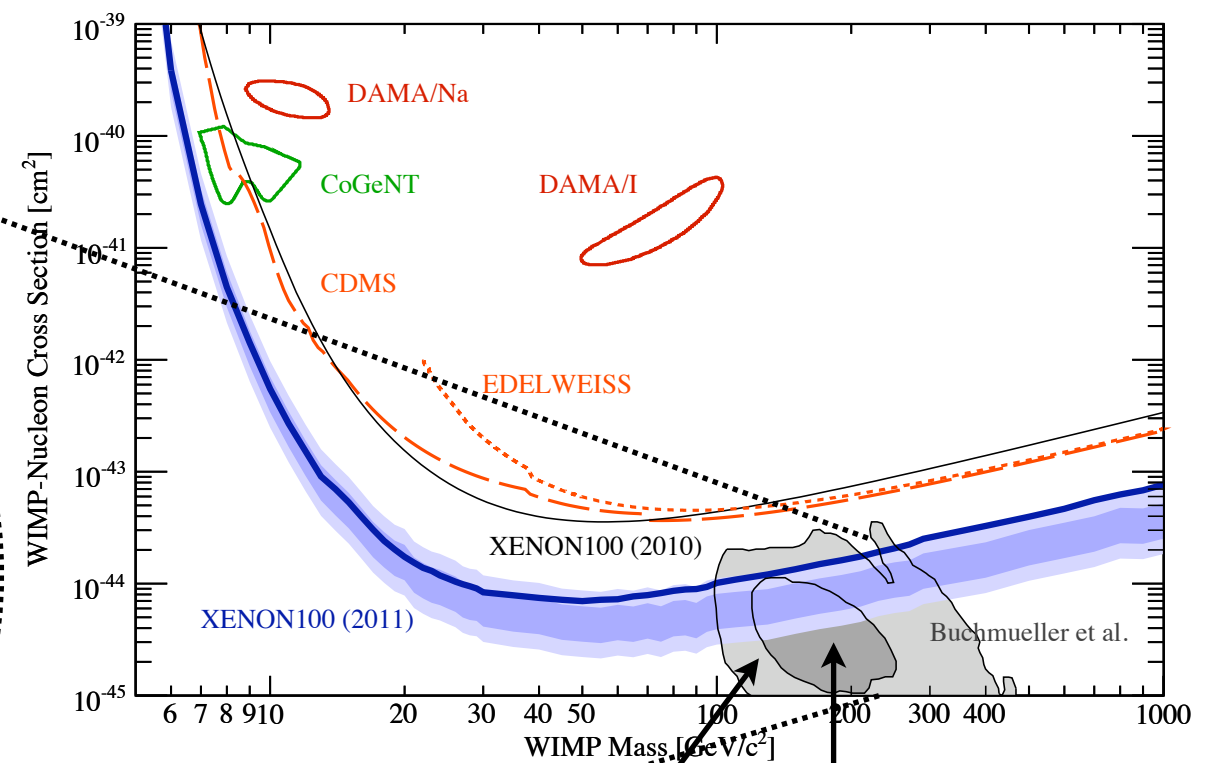
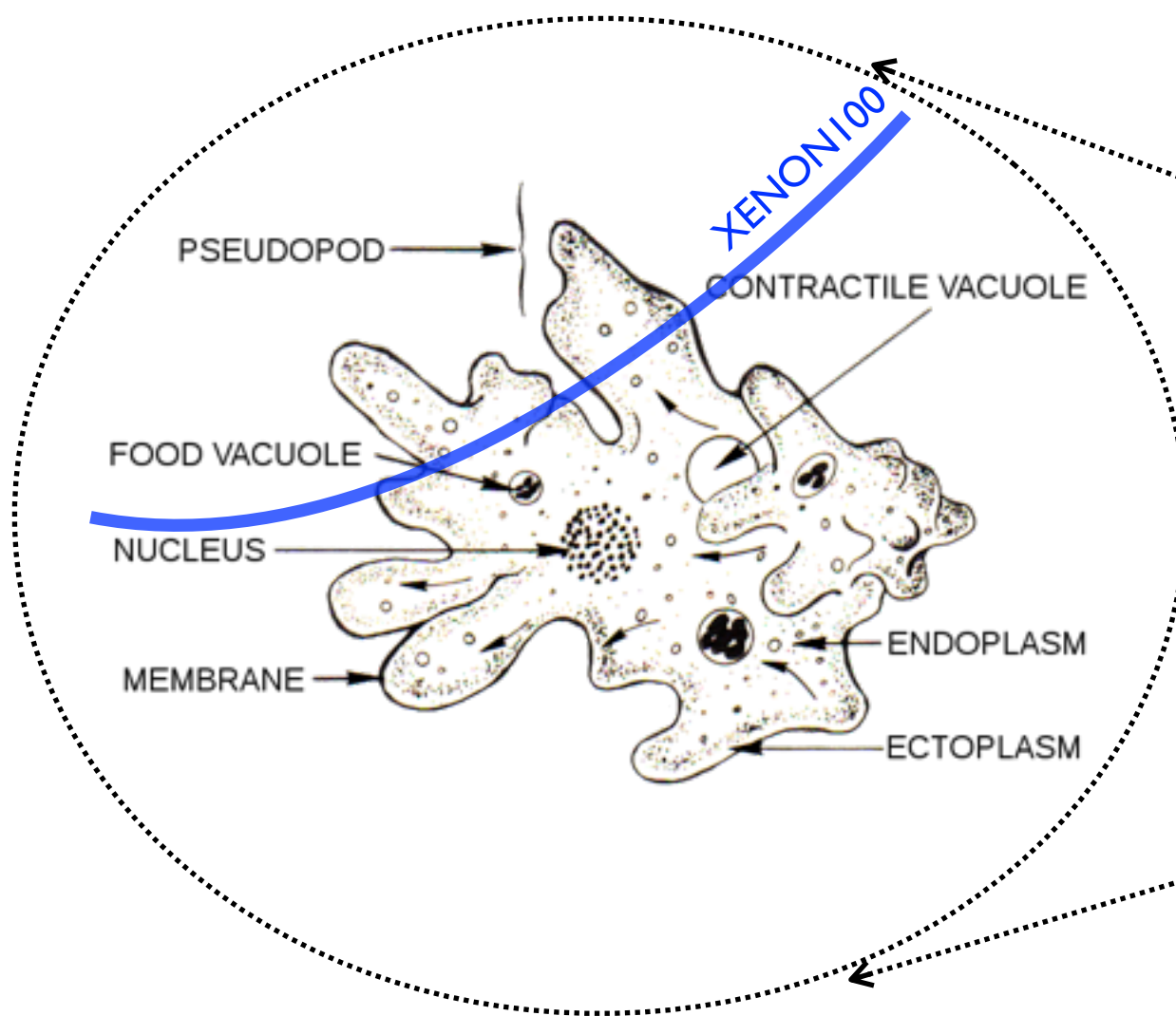


tri-leptons+  
jets + MET



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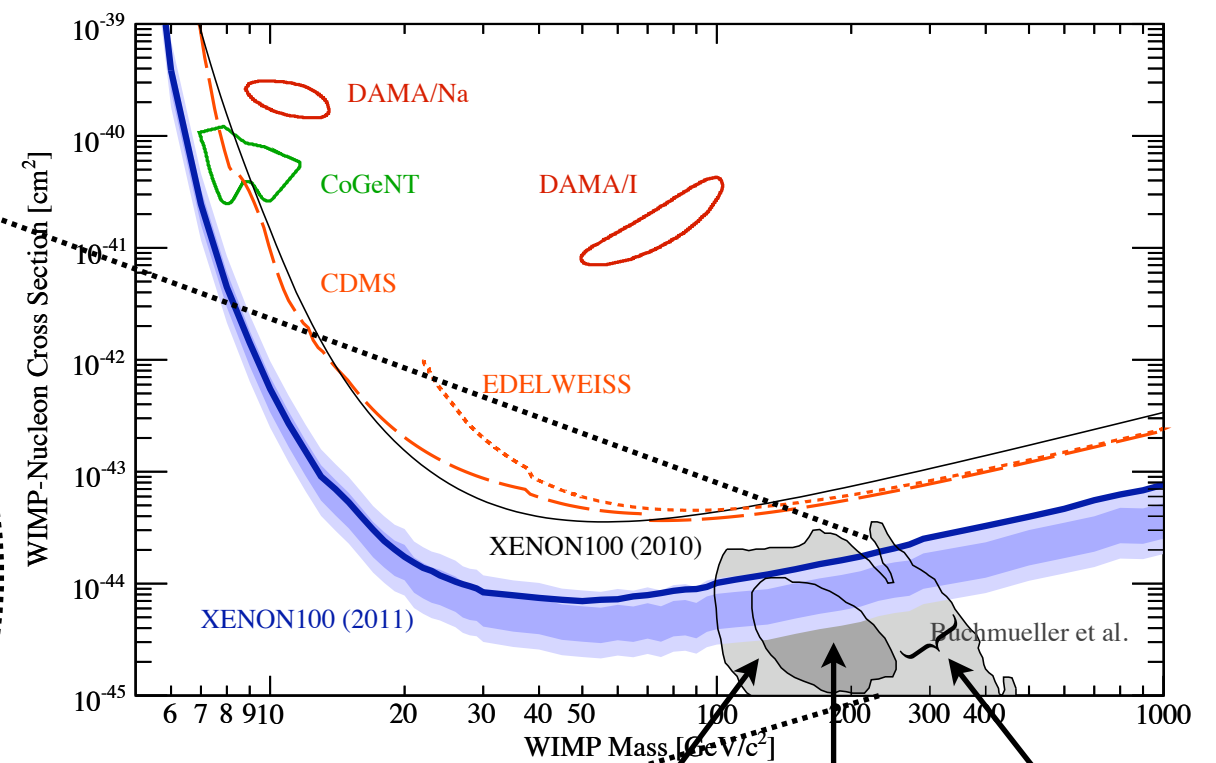
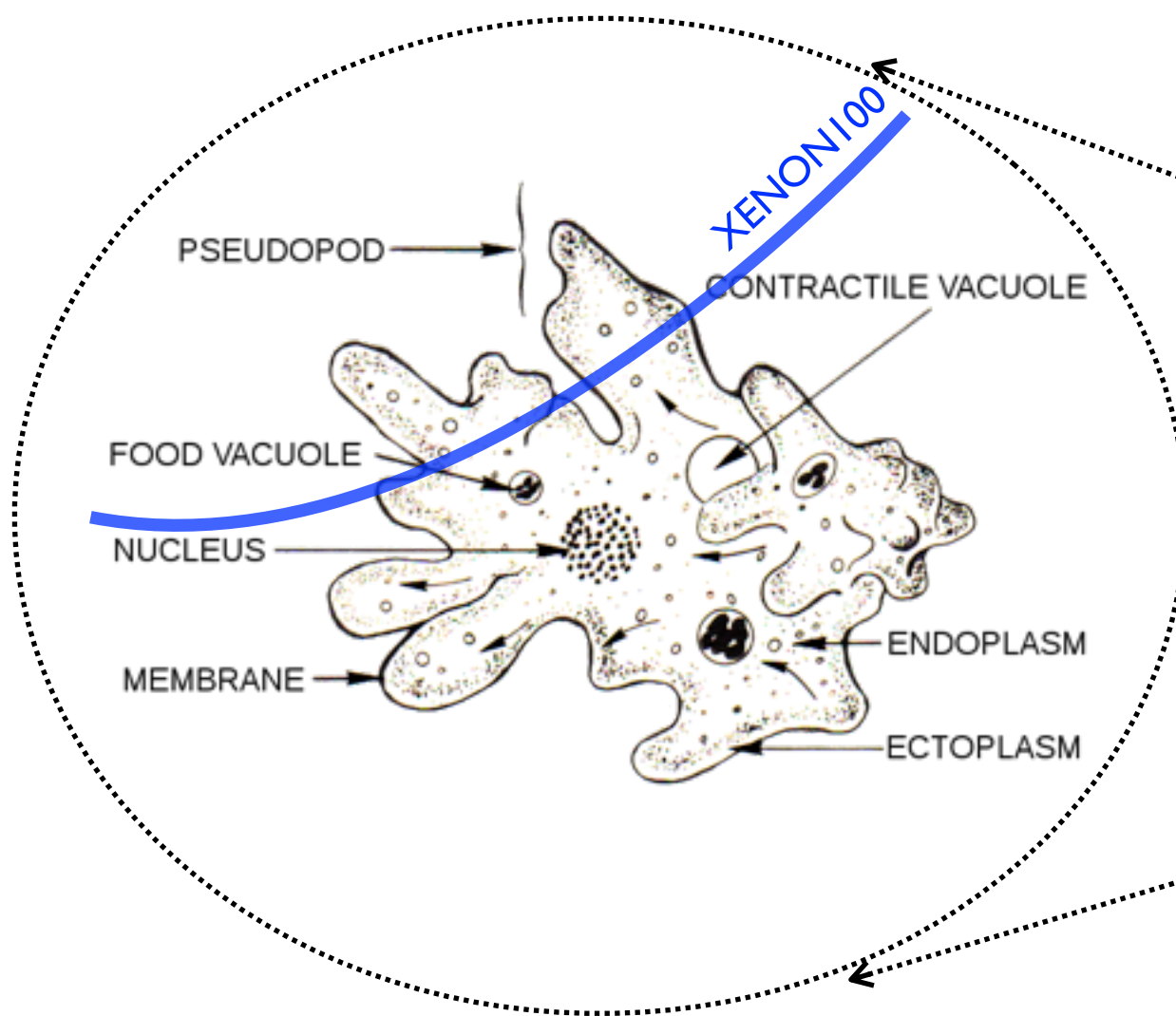


jets + MET

tri-leptons+ jets + MET

# Collider Connections?

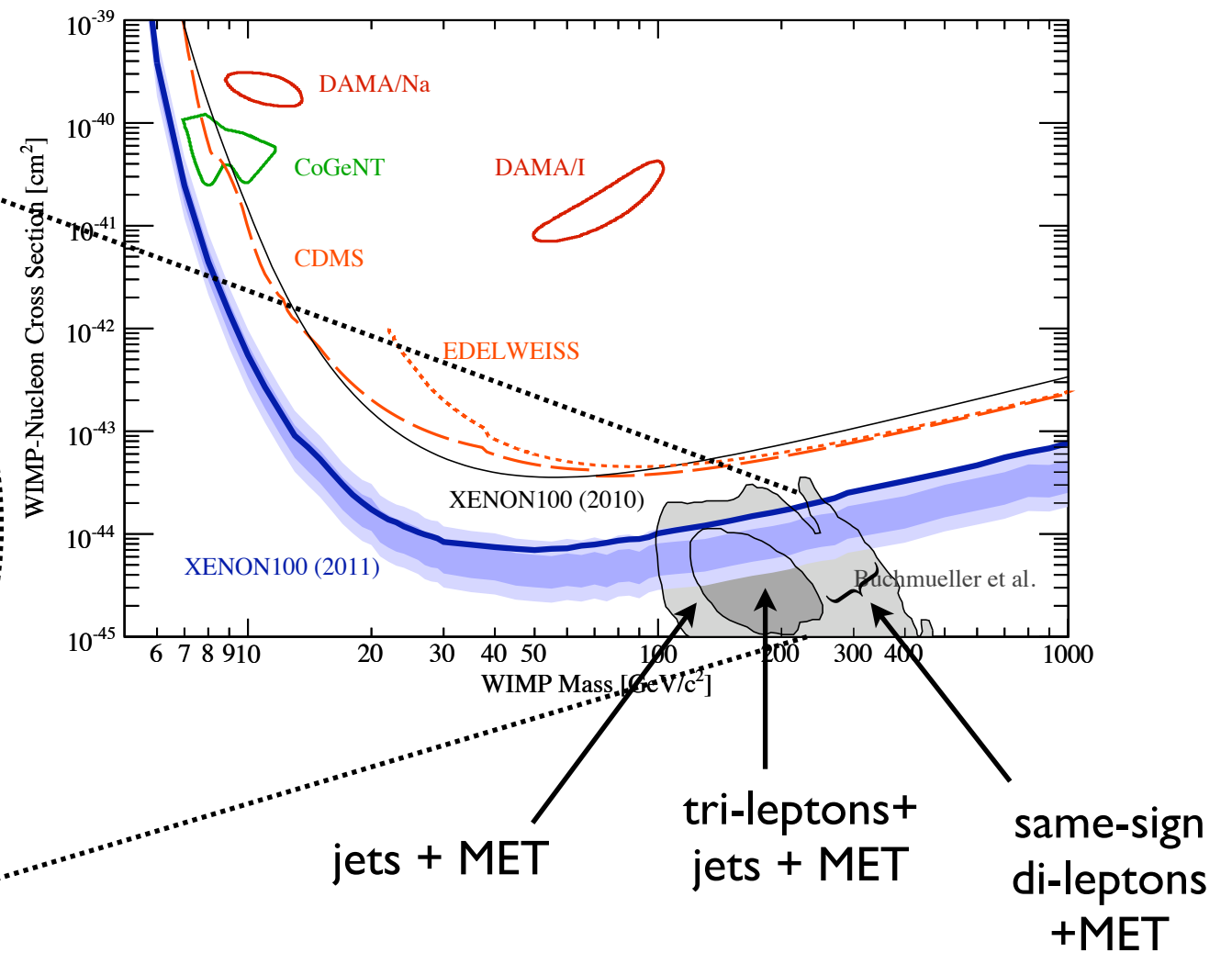
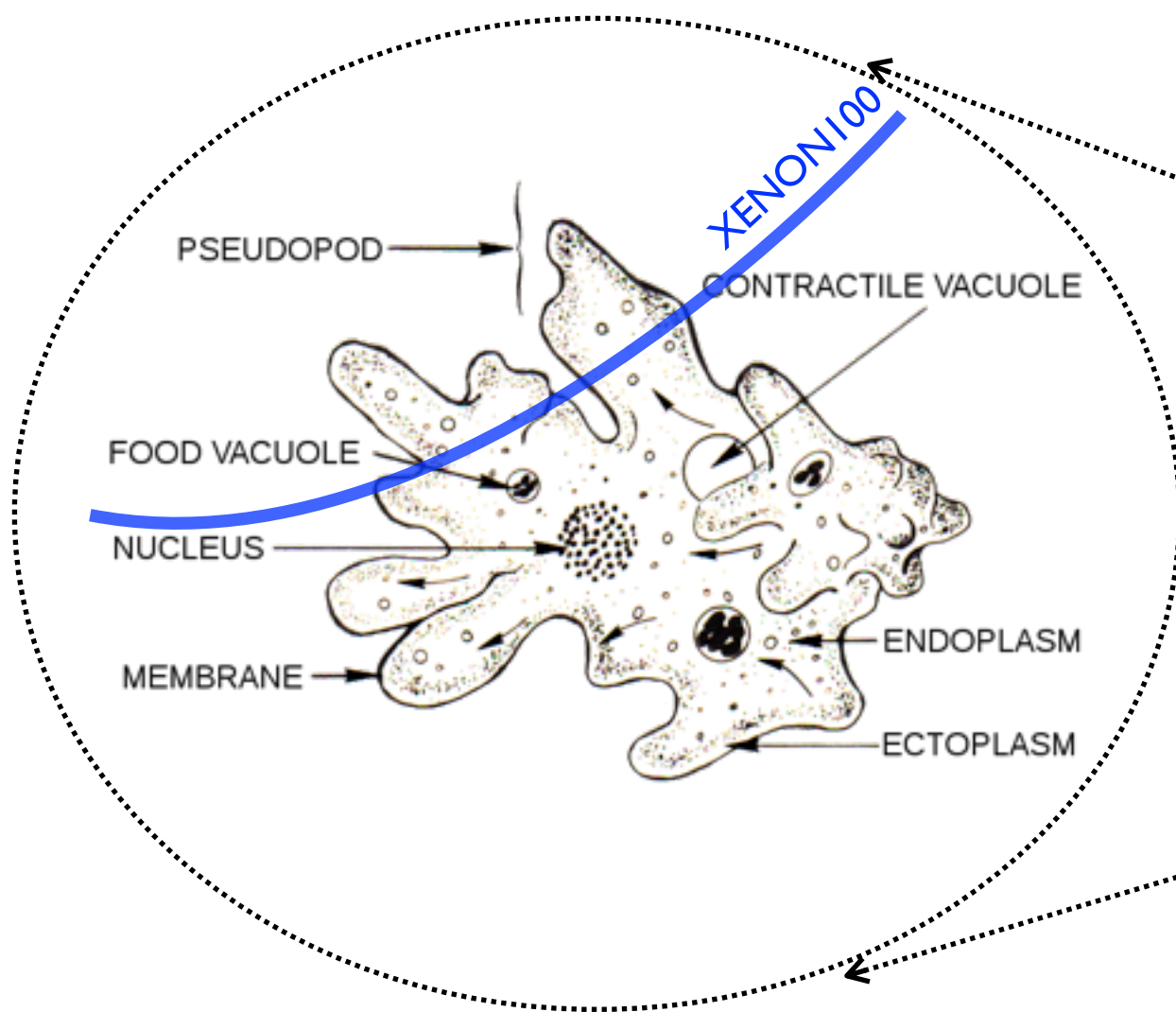
- \* DM experiments and colliders are often said to be related *in a specific framework* (SUSY).



jets + MET  
tri-leptons+ jets + MET  
same-sign di-leptons +MET

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- \* DM experiments and colliders are often said to be related *in a specific framework* (SUSY).



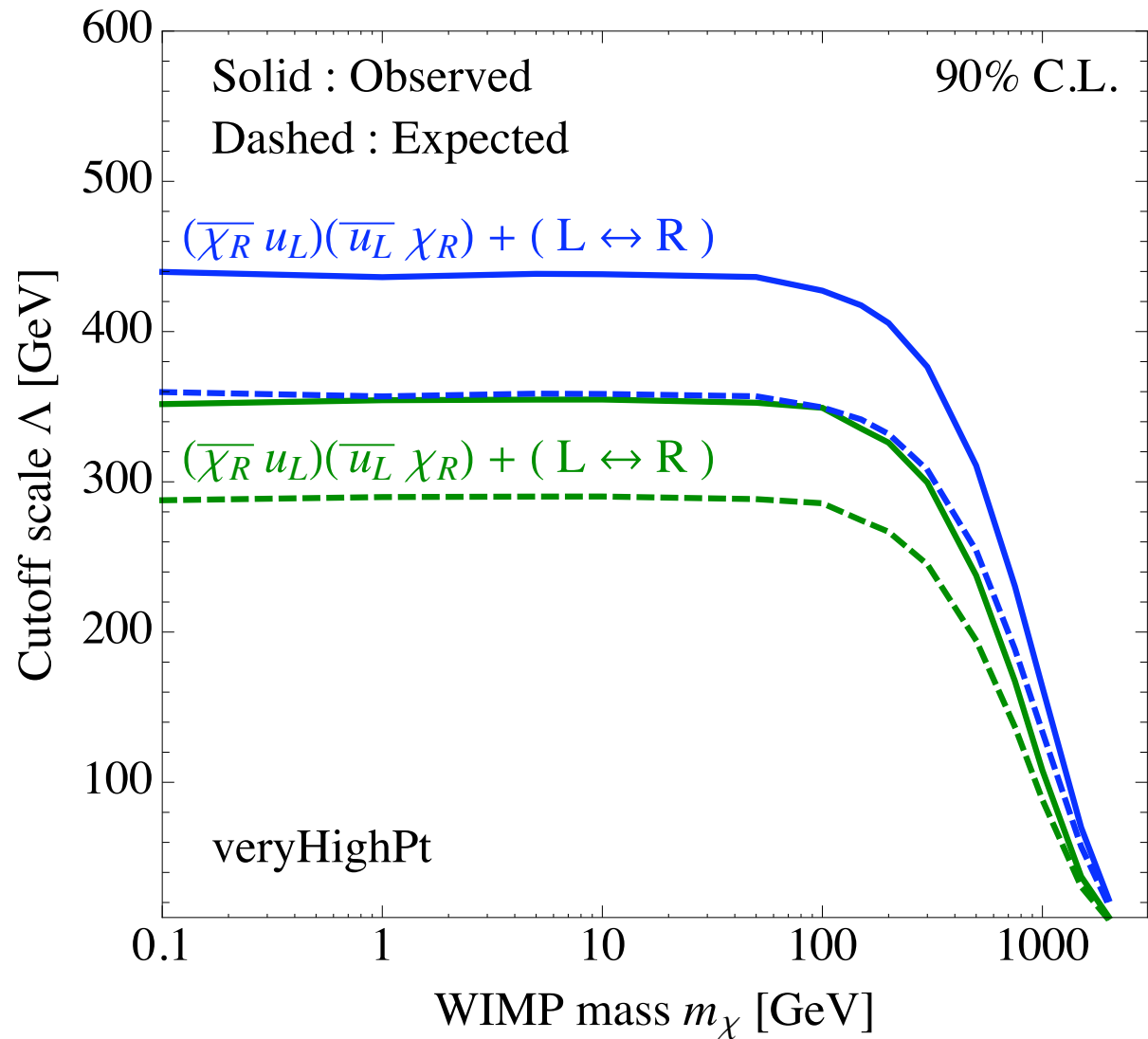
“XENON100 is starting to probe the MSSM’s pseudopod, LHC killed the Membrane, but the ectoplasm is still safe.” [submitted to nature]

# Games: Higgs searches & DM

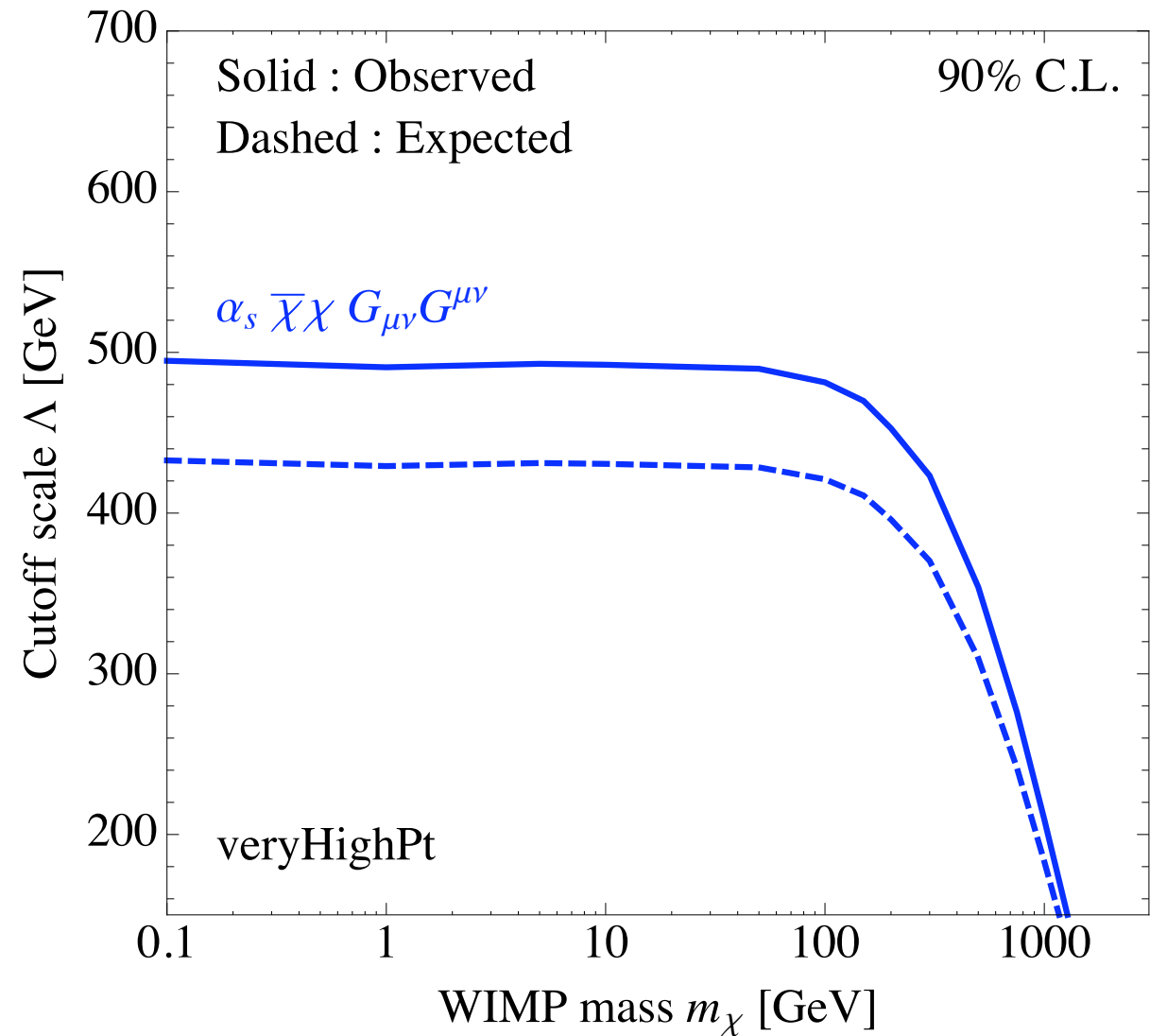
- \* *Assume* the Higgs hint is real w/ SM production.
- \* The fact that it was seen in di-photon with the rate that it has, places limits on competing modes, e.g. Higgs to invisible.
- \* Places **upper** limit on higgs mediated direct detection.
- \* *Assume* a Higgs mass that is already excluded for SM.
- \* *Assume* the reason it was excluded is an invisible branching fraction.
- \* This places a lower limit on the invisible BR.  
Places a **lower** limit on higgs mediated direct detection.

# Other Operators:

ATLAS 7 TeV , 1 fb<sup>-1</sup>



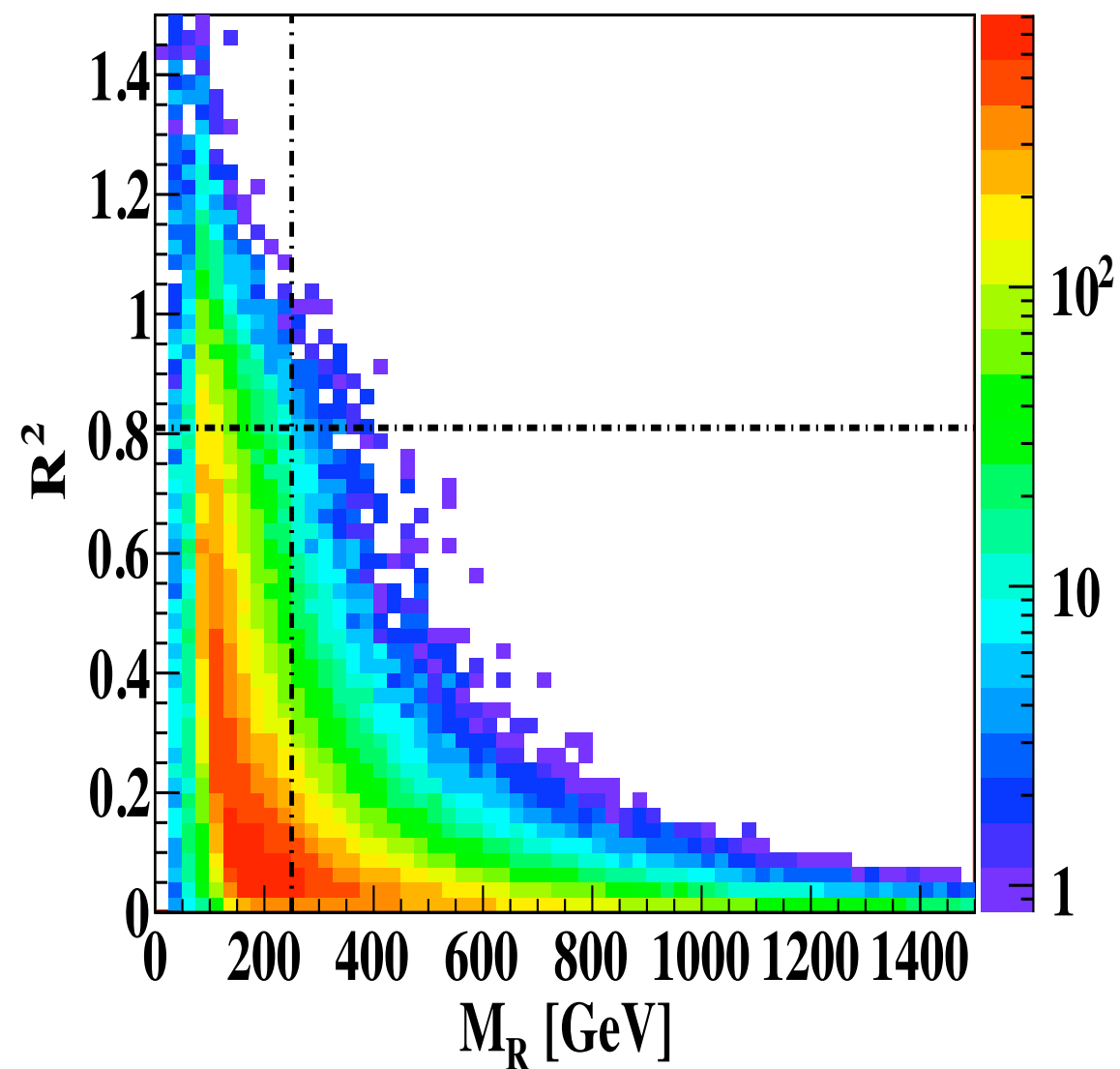
ATLAS 7 TeV , 1 fb<sup>-1</sup>



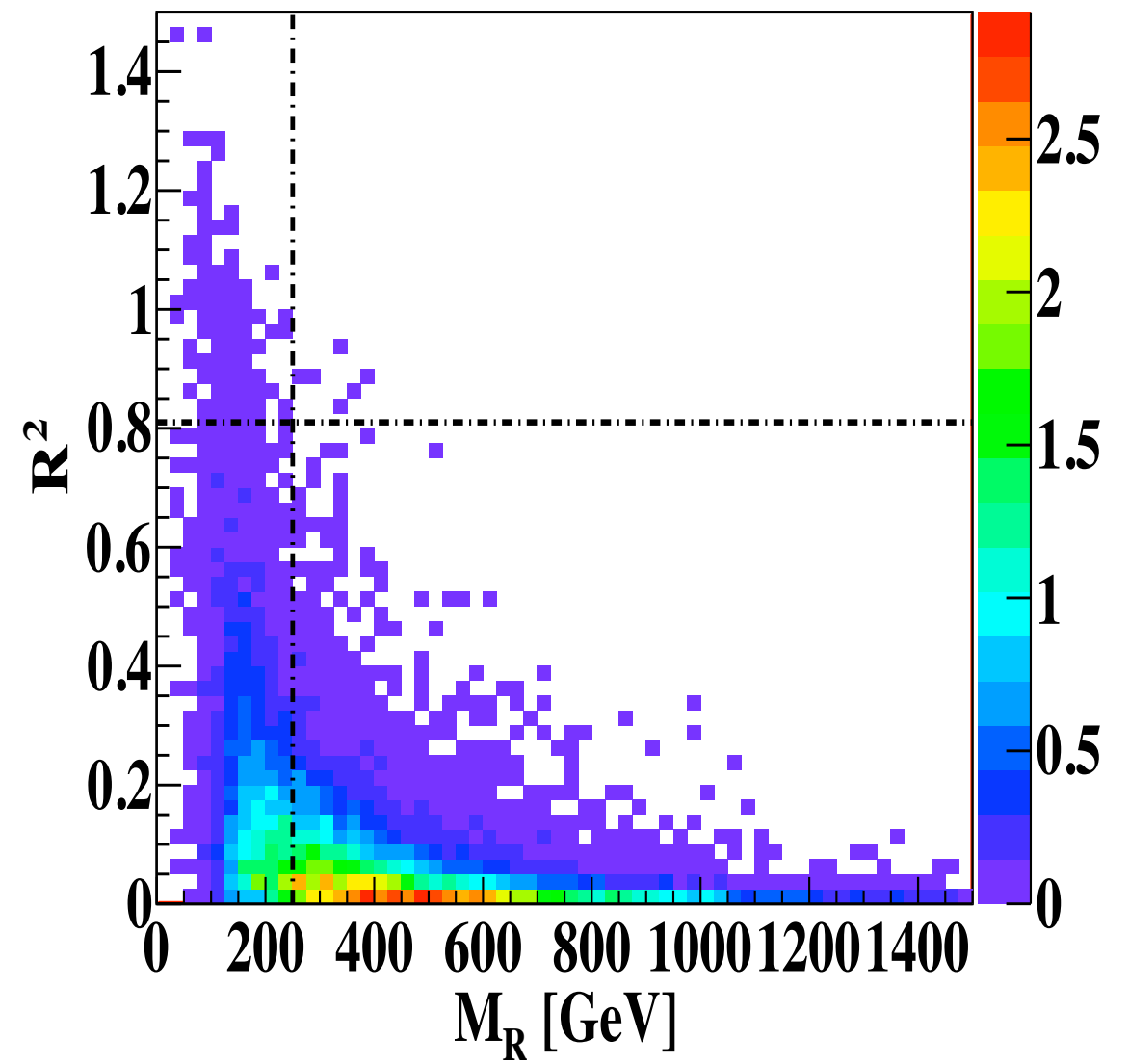
*The limit is flat up to ~200 GeV.  
Goes all the way to zero.*

# Razor Shapes

\* Other Backgrounds:



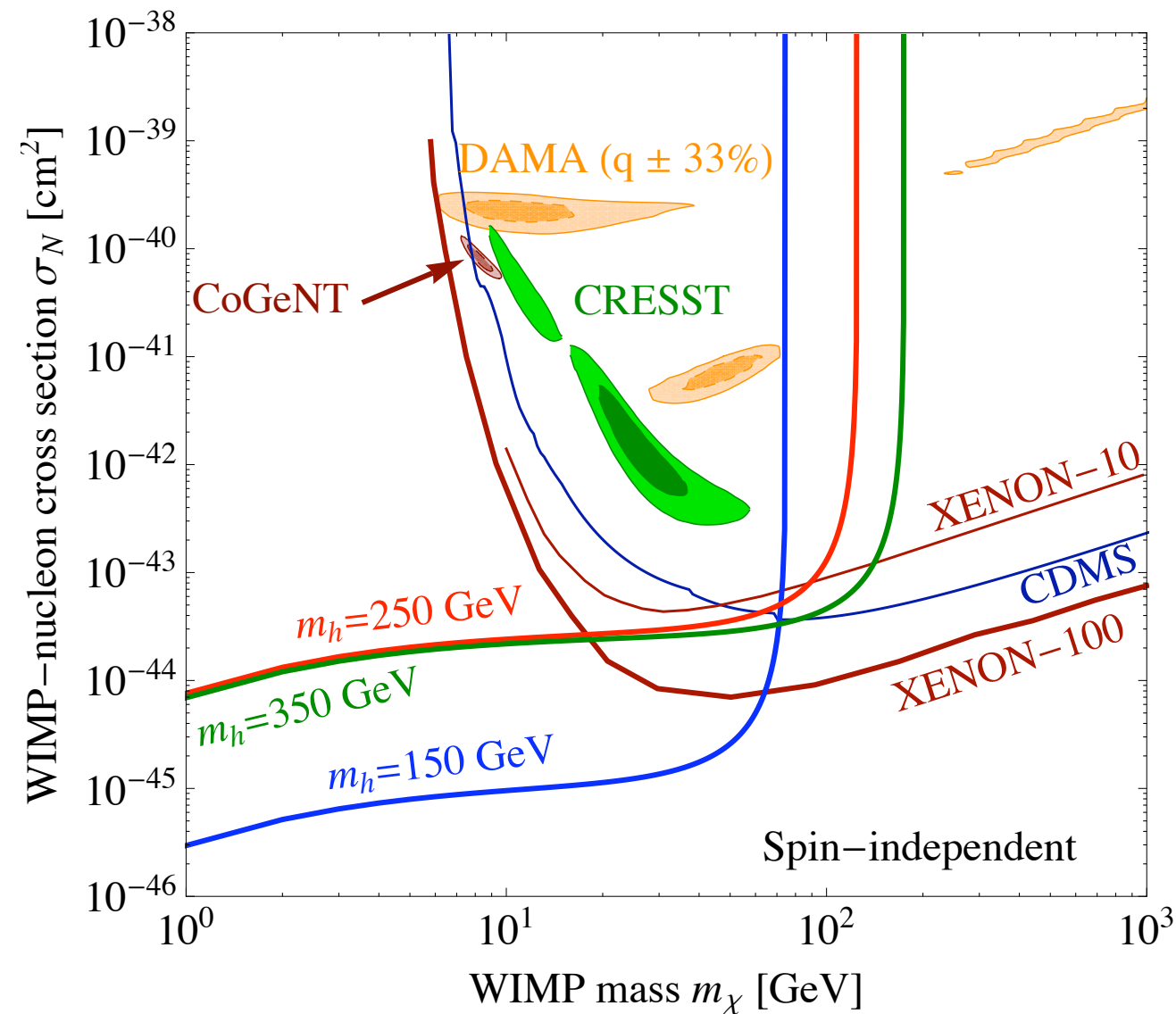
(b)  $W$ +jets.



(c)  $t\bar{t}$ .

# Current Higgs limits vs DM

CMS Higgs combined lower bound



*Also, if a light SM Higgs is discovered, an upper limit on DD can be extracted.*



# CDF: jet + MET $(1 \text{ fb}^{-1})$

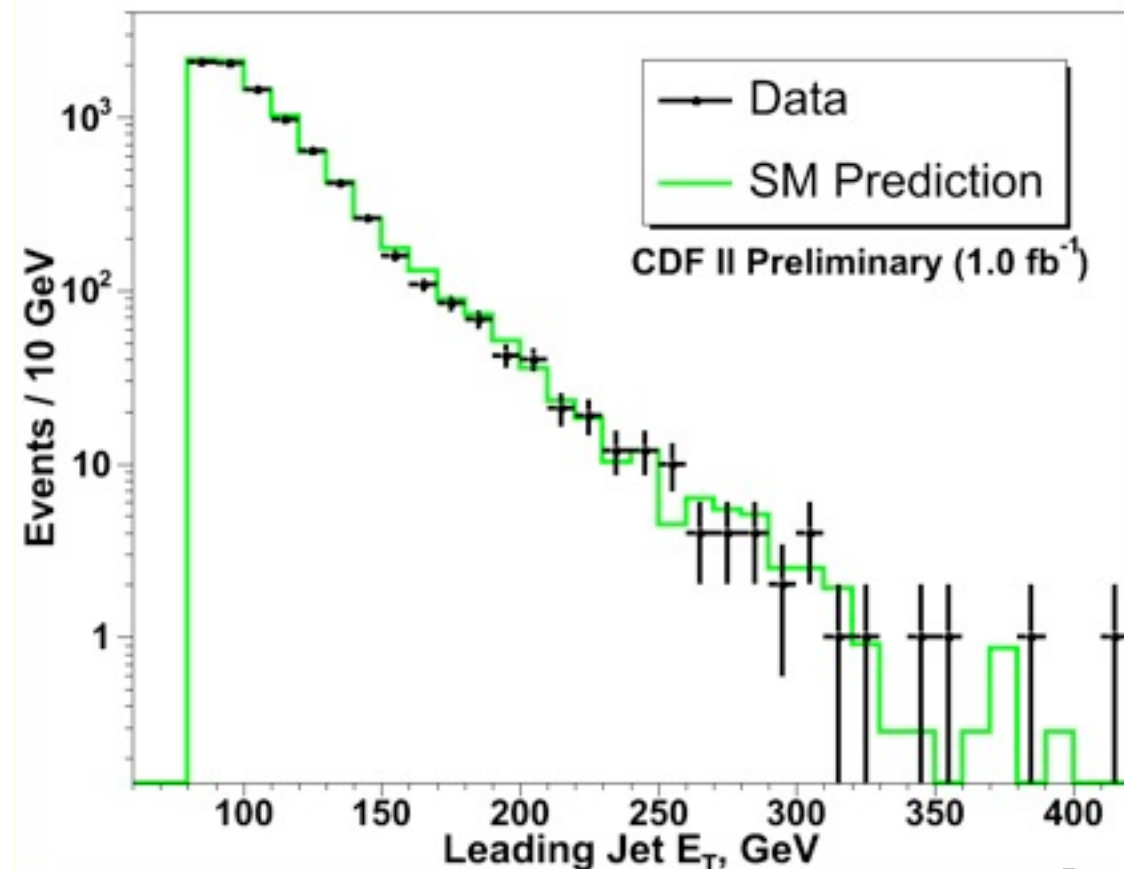
counting experiment:

$$\cancel{E}_T > 80 \text{ GeV}$$

$$p_T(j1) > 80 \text{ GeV}$$

$$p_T(j2) < 30 \text{ GeV}$$

$$p_T(j3) < 20 \text{ GeV}$$



Background	Number of Events
Z $\rightarrow$ nu nu	3203 +/- 137
W $\rightarrow$ tau nu	2010 +/- 69
W $\rightarrow$ mu nu	1570 +/- 54
W $\rightarrow$ e nu	824 +/- 28
Z $\rightarrow$ ll	87 +/- 3
QCD	708 +/- 146
Gamma plus Jet	209 +/- 41
Non-Collision	52 +/- 52
Total Predicted	8663 +/- 332
Data Observed	8449

**Observed: 8449 events**



# ATLAS Analysis

## \* ATLAS's 1fb analysis employs 3 sets of cuts

LowPT Selection requires  $\cancel{E}_T > 120$  GeV, one jet with  $p_T(j_1) > 120$  GeV,  $|\eta(j_1)| < 2$ , and events are vetoed if they contain a second jet with  $p_T(j_2) > 30$  GeV and  $|\eta(j_2)| < 4.5$ .

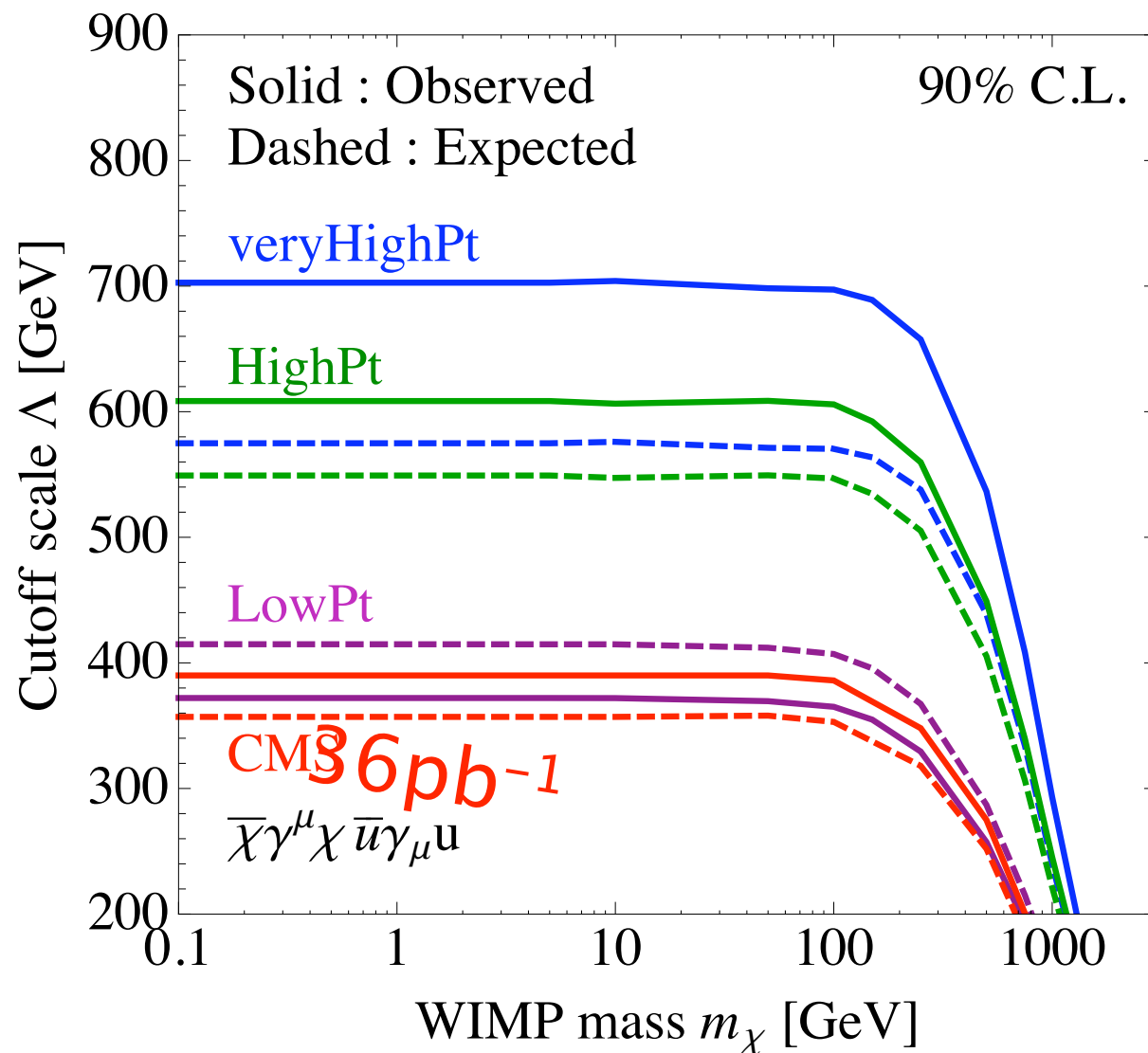
HighPT Selection requires  $\cancel{E}_T > 220$  GeV, one jet with  $p_T(j_1) > 250$  GeV,  $|\eta(j_1)| < 2$ , and events are vetoed if there is a second jet with  $|\eta(j_2)| < 4.5$  and with either  $p_T(j_2) > 60$  GeV or  $\Delta\phi(j_2, \cancel{E}_T) < 0.5$ . Any further jets with  $|\eta(j_2)| < 4.5$  must have  $p_T(j_3) < 30$  GeV.

veryHighPT Selection requires  $\cancel{E}_T > 300$  GeV, one jet with  $p_T(j_1) > 350$  GeV,  $|\eta(j_1)| < 2$ , and events are vetoed if there is a second jet with  $|\eta(j_2)| < 4.5$  and with either  $p_T(j_2) > 60$  GeV or  $\Delta\phi(j_2, \cancel{E}_T) < 0.5$ . Any further jets with  $|\eta(j_2)| < 4.5$  must have  $p_T(j_3) < 30$  GeV.

	ATLAS LowPT	ATLAS HighPT	ATLAS veryHighPT
	$1.0 \text{ fb}^{-1}$	$1.0 \text{ fb}^{-1}$	$1.0 \text{ fb}^{-1}$
Expected	$15100 \pm 700$	$1010 \pm 75$	$193 \pm 25$
Observed	15740	965	167

# Limits on $\Lambda \equiv \frac{M}{\sqrt{g_\chi g_1}}$ :

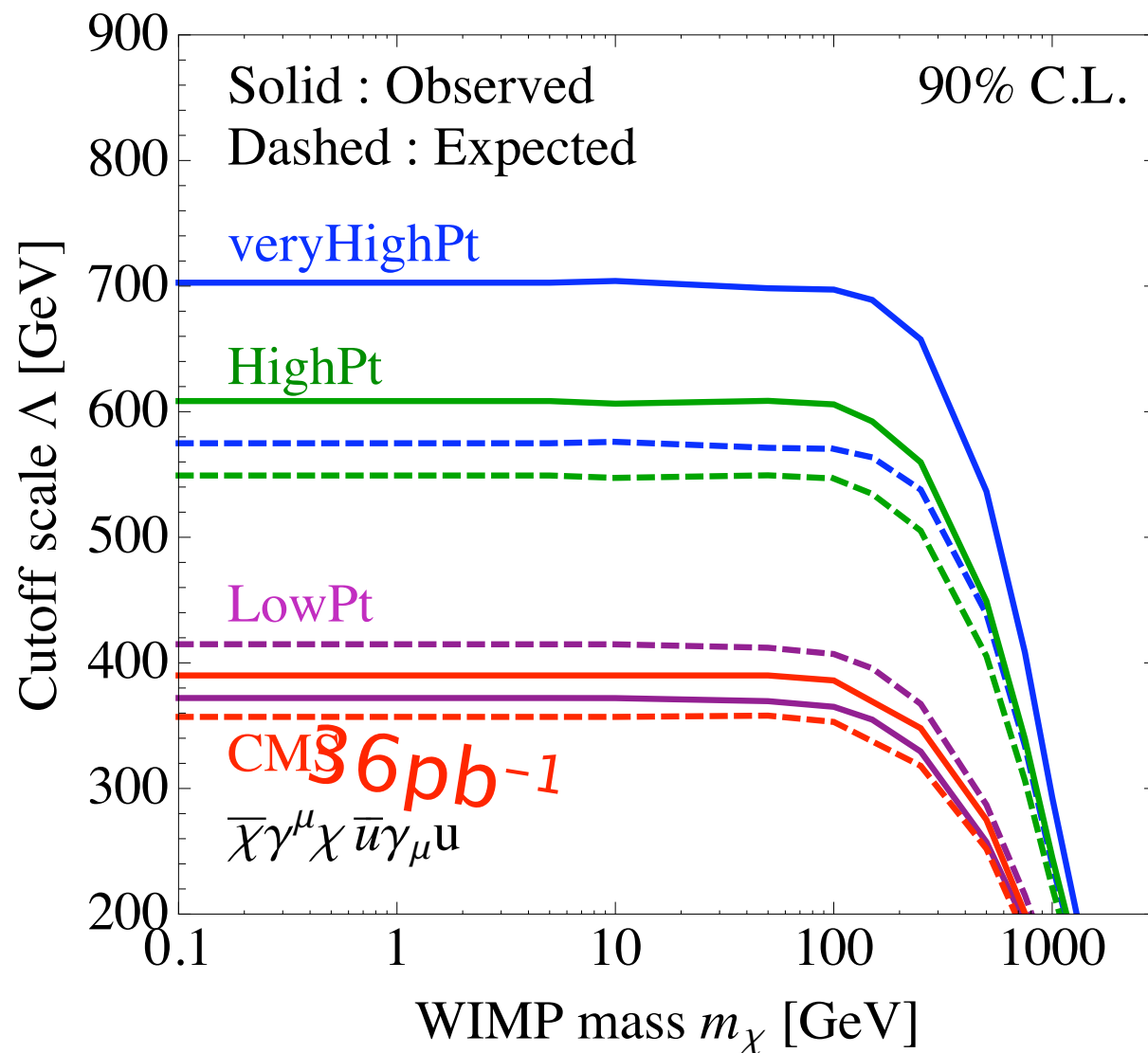
\* Set 90% CL limits:  $\chi^2 \equiv \frac{[\Delta_N - N_{\text{DM}}(m_\chi, \Lambda)]^2}{N_{\text{DM}}(m_\chi, \Lambda) + N_{\text{SM}} + \sigma_{\text{SM}}^2} = 2.71.$



$$\Delta_N = \begin{cases} 0 & \text{expected bound} \\ N_{\text{obs}} - N_{\text{SM}} & \text{observed bound} \end{cases}$$

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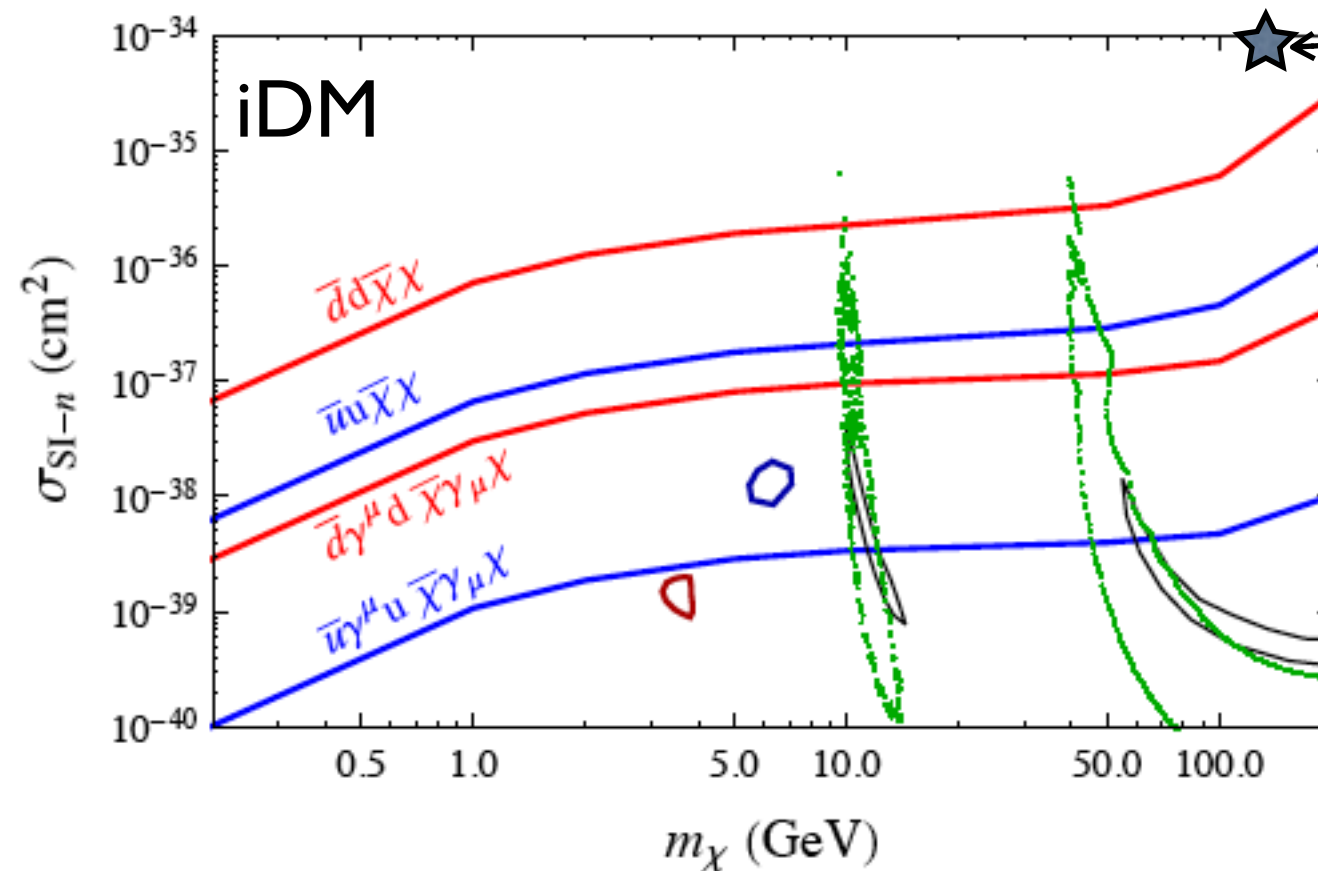


$$\Delta_N = \begin{cases} 0 & \text{expected bound} \\ N_{\text{obs}} - N_{\text{SM}} & \text{observed bound} \end{cases}$$

Harder is better.  
in the future:  
populate the tail  
and keep cutting harder

# iDM, MDDDM, ...

- \* There are other scenarios in which DD is suppressed, but colliders don't care:



impure thoughts...  
can survive with light  
mediator.

# LEP mono-photon

w/ Fox, Kopp and Tsai  
arXiv:1103.0240

# LEP

- \* Directly constrain DM coupling to electrons.
- \* **But**, in many models quark and lepton coupling are related (consider 2 benchmarks).
- \* LEP is a clean environment. Ability to measure missing mass.
- \* Places non-trivial limits also on indirect searches in lepton channels (e.g. the Hooperon).

# Operators

\* Same story w/ leptons (assume universality)

$$\mathcal{O}_V = \frac{(\bar{\chi}\gamma_\mu\chi)(\bar{\ell}\gamma^\mu\ell)}{\Lambda^2}, \quad (\text{vector, } s\text{-channel})$$

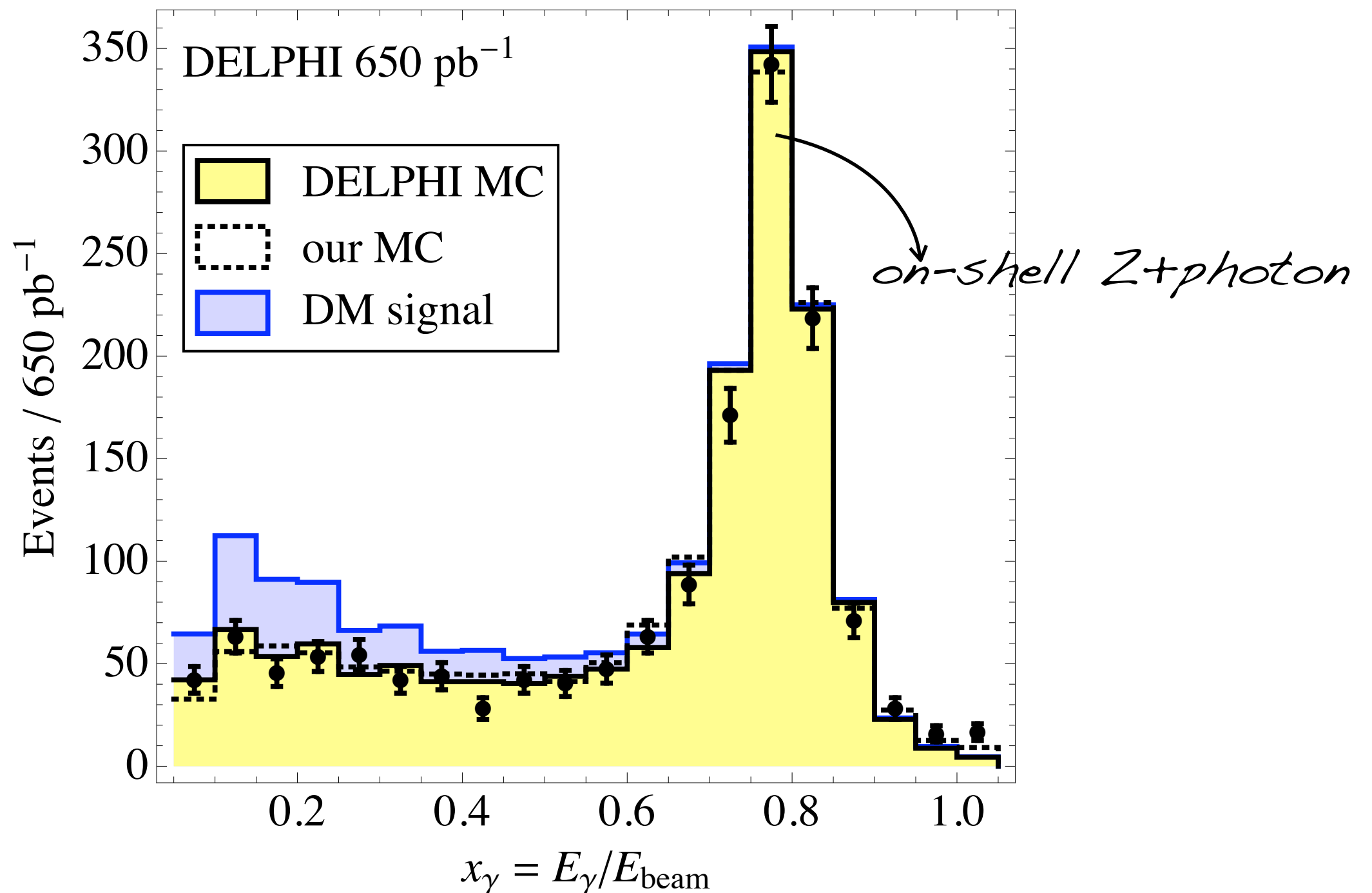
$$\mathcal{O}_S = \frac{(\bar{\chi}\chi)(\bar{\ell}\ell)}{\Lambda^2}, \quad (\text{scalar, } s\text{-channel})$$

$$\mathcal{O}_A = \frac{(\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{\ell}\gamma^\mu\gamma_5\ell)}{\Lambda^2}, \quad (\text{axial vector, } s\text{-channel})$$

$$\mathcal{O}_t = \frac{(\bar{\chi}\ell)(\bar{\ell}\chi)}{\Lambda^2}, \quad (\text{scalar, } t\text{-channel})$$

# Mono-photon

- \* Use spectrum shape to reject background peak.



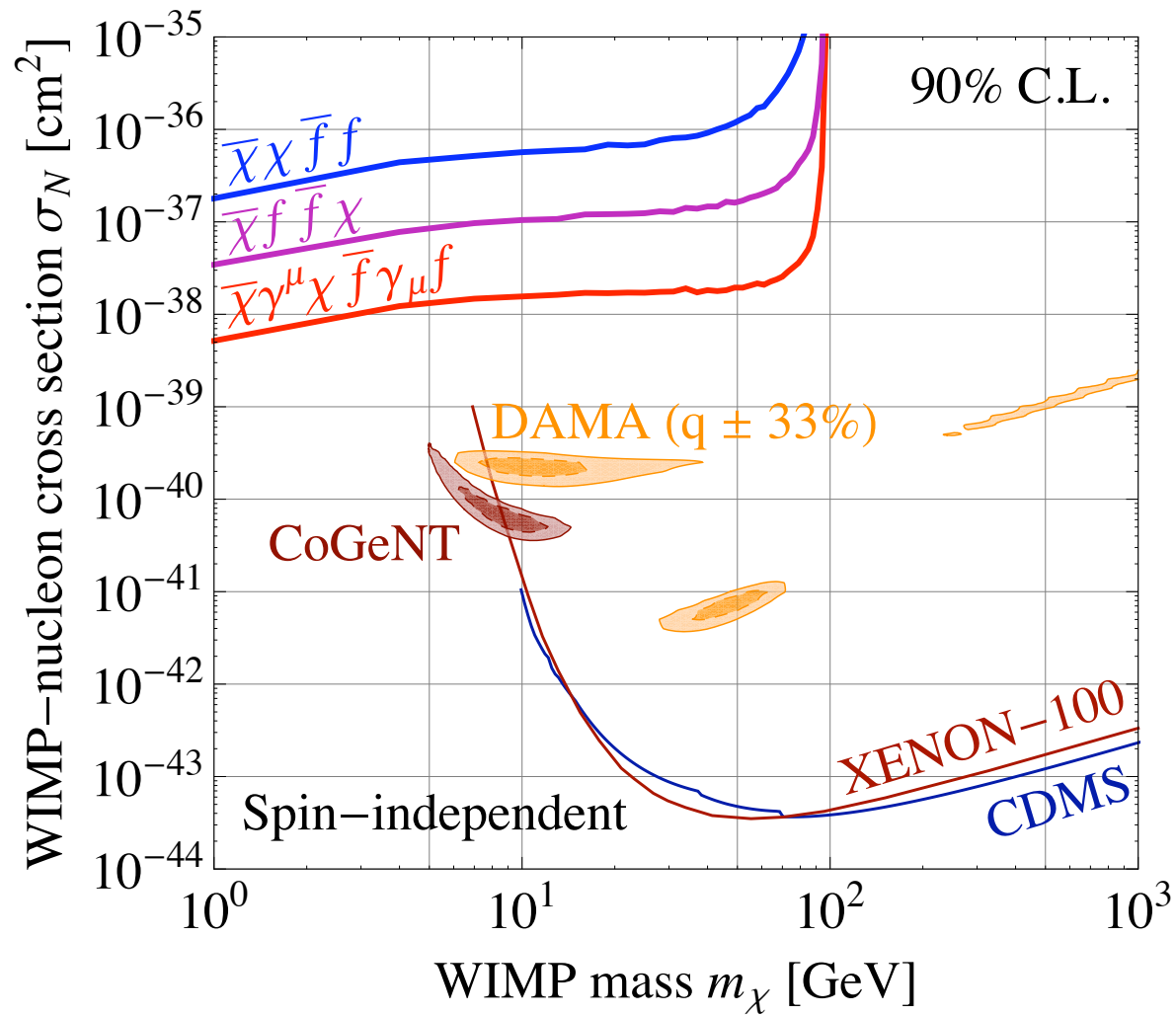


# Model Dependence

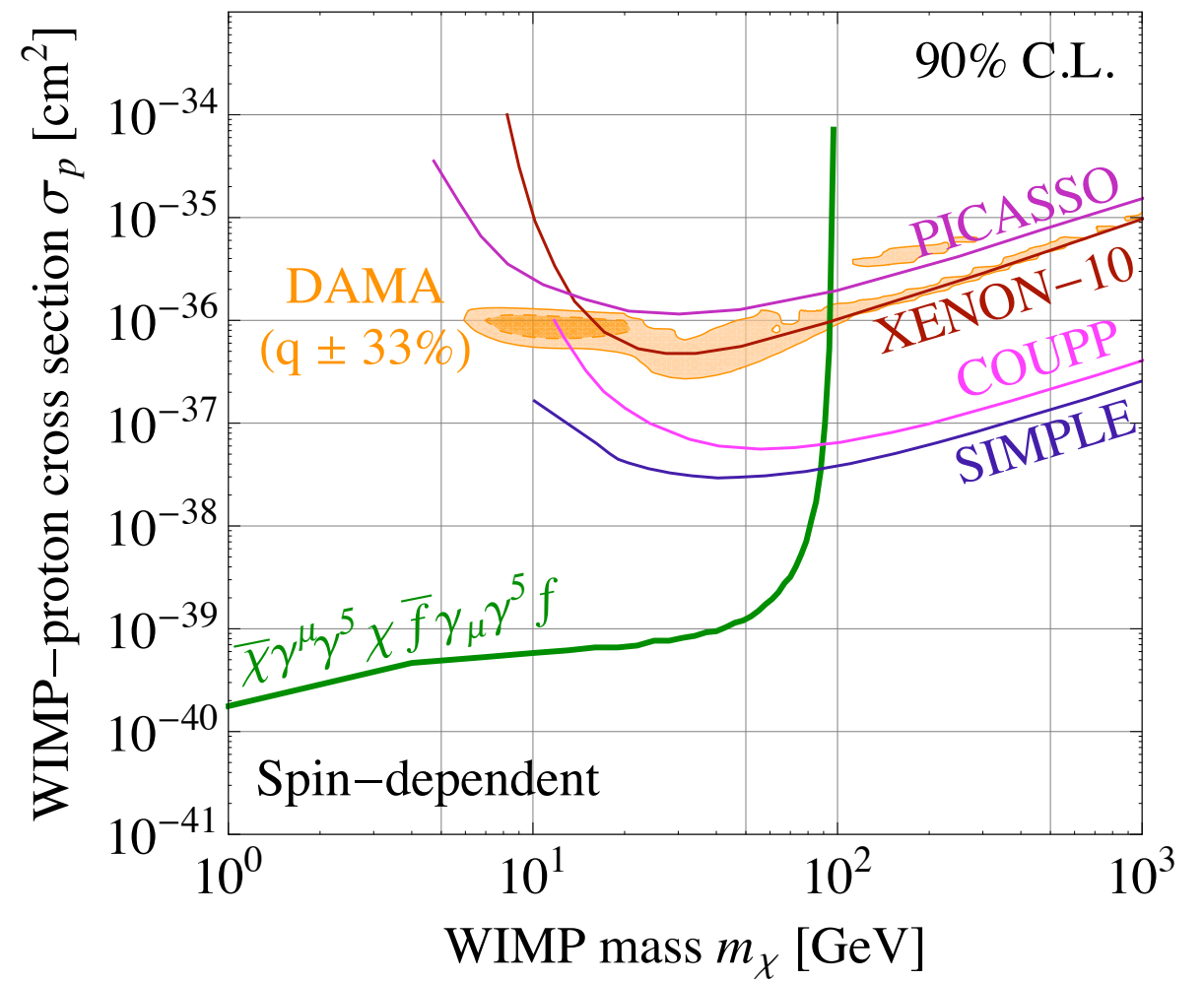
- \* We limit lepton couplings.
- \* But how does DM couple to quarks?
- \* Consider 2 extreme cases:
  - o Couplings to **quarks are same as leptons.**
  - o Couplings to quarks are **zero.**
- \* *Any other case can be derived from these two.*

# DD Limits

Equal couplings to all SM fermions

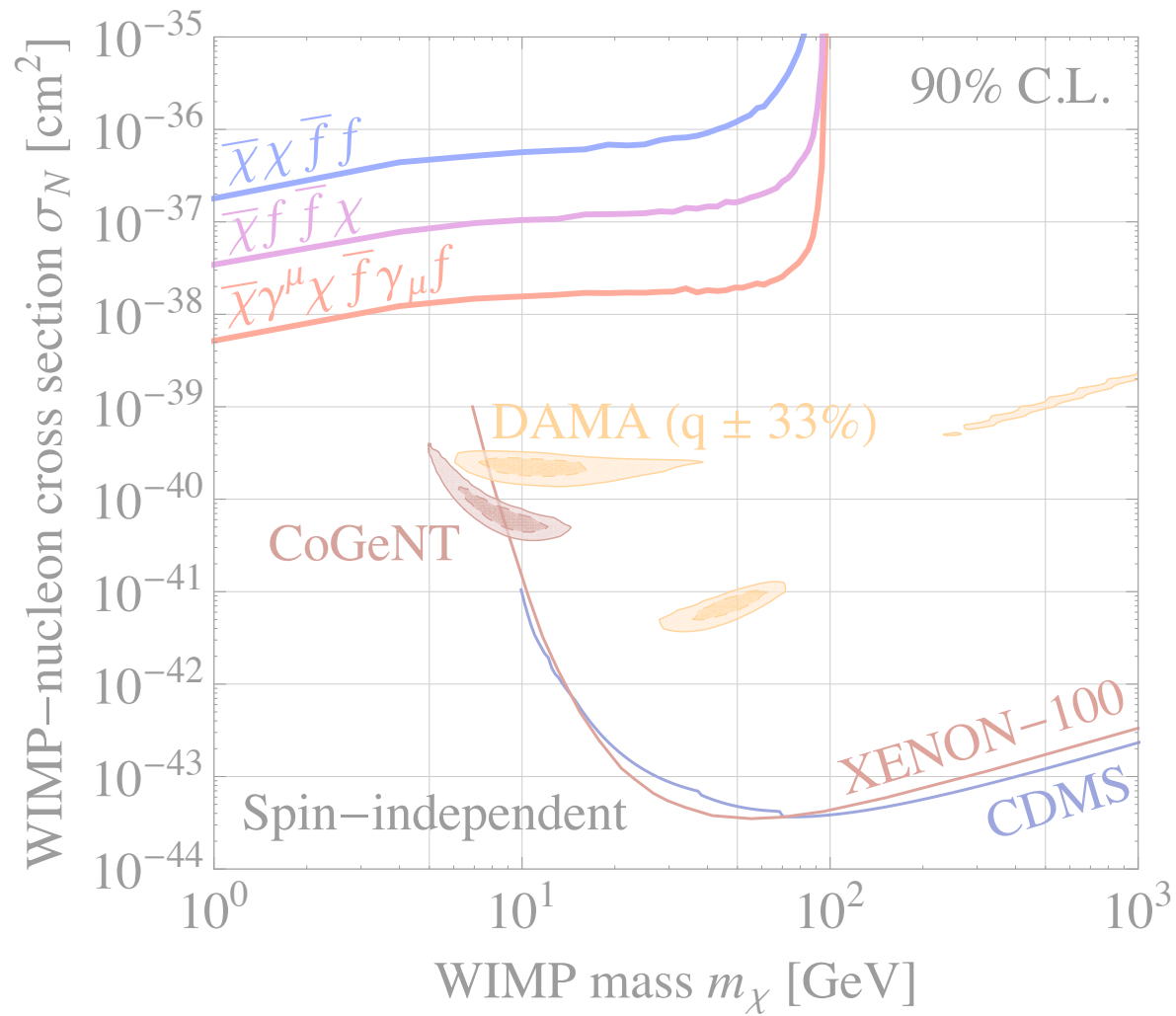


Equal couplings to all SM fermions

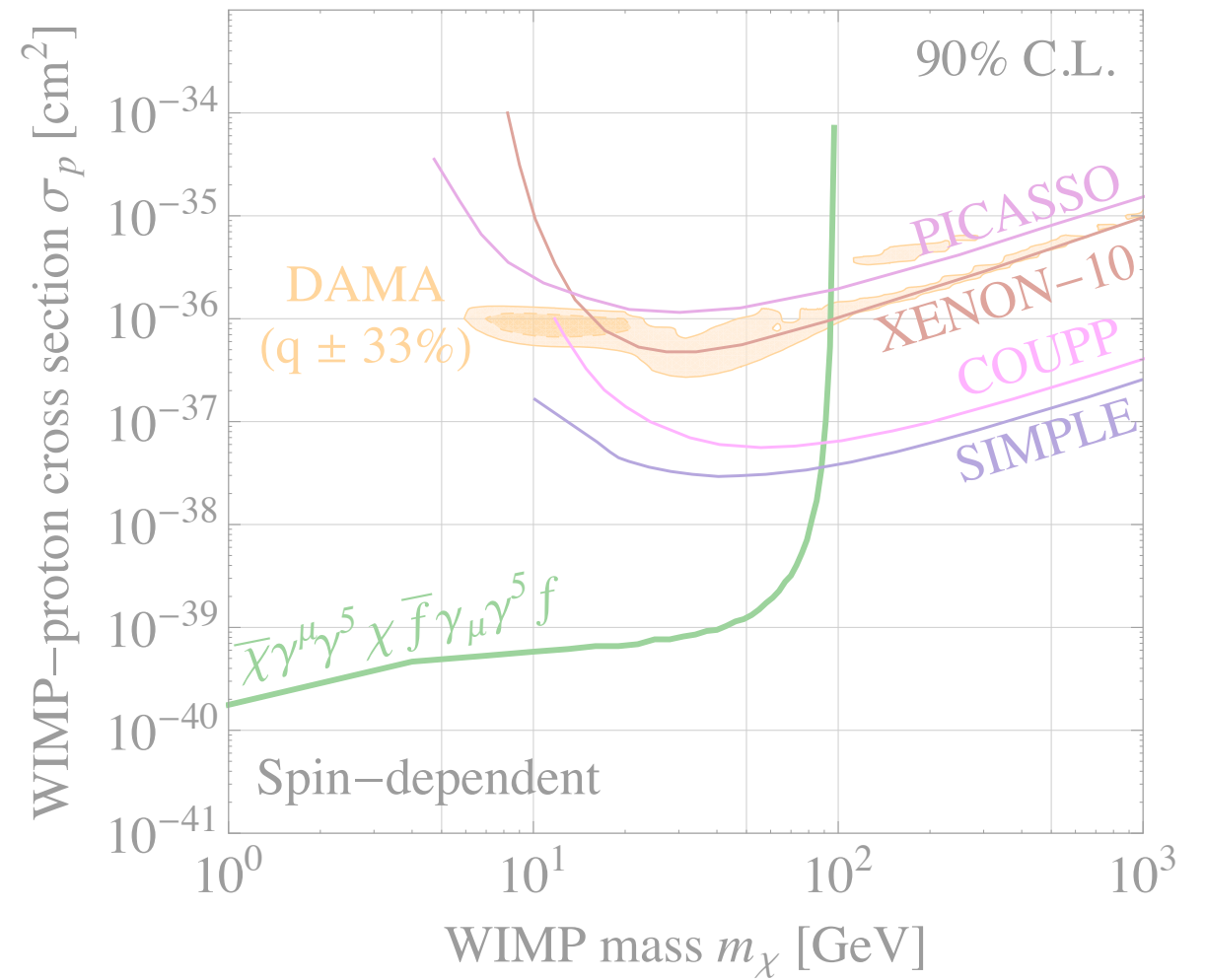


# DD Limits

Equal couplings to all SM fermions

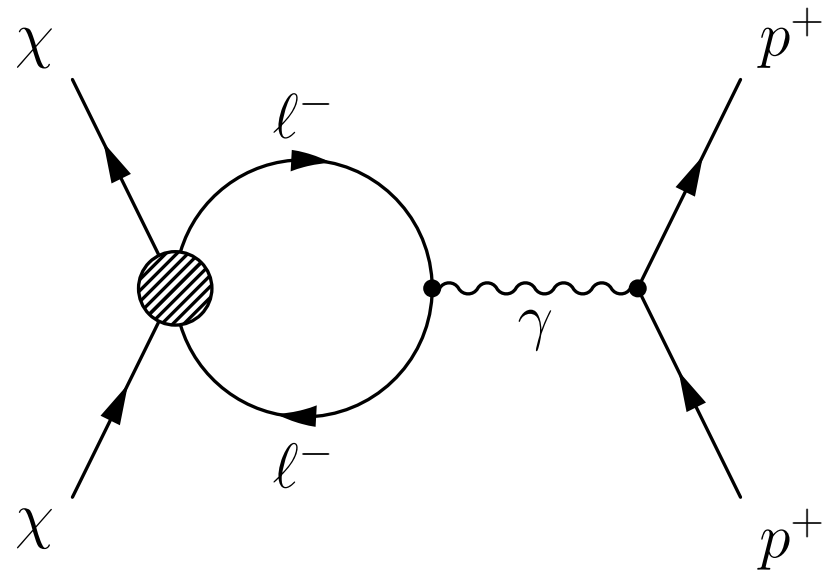


Equal couplings to all SM fermions



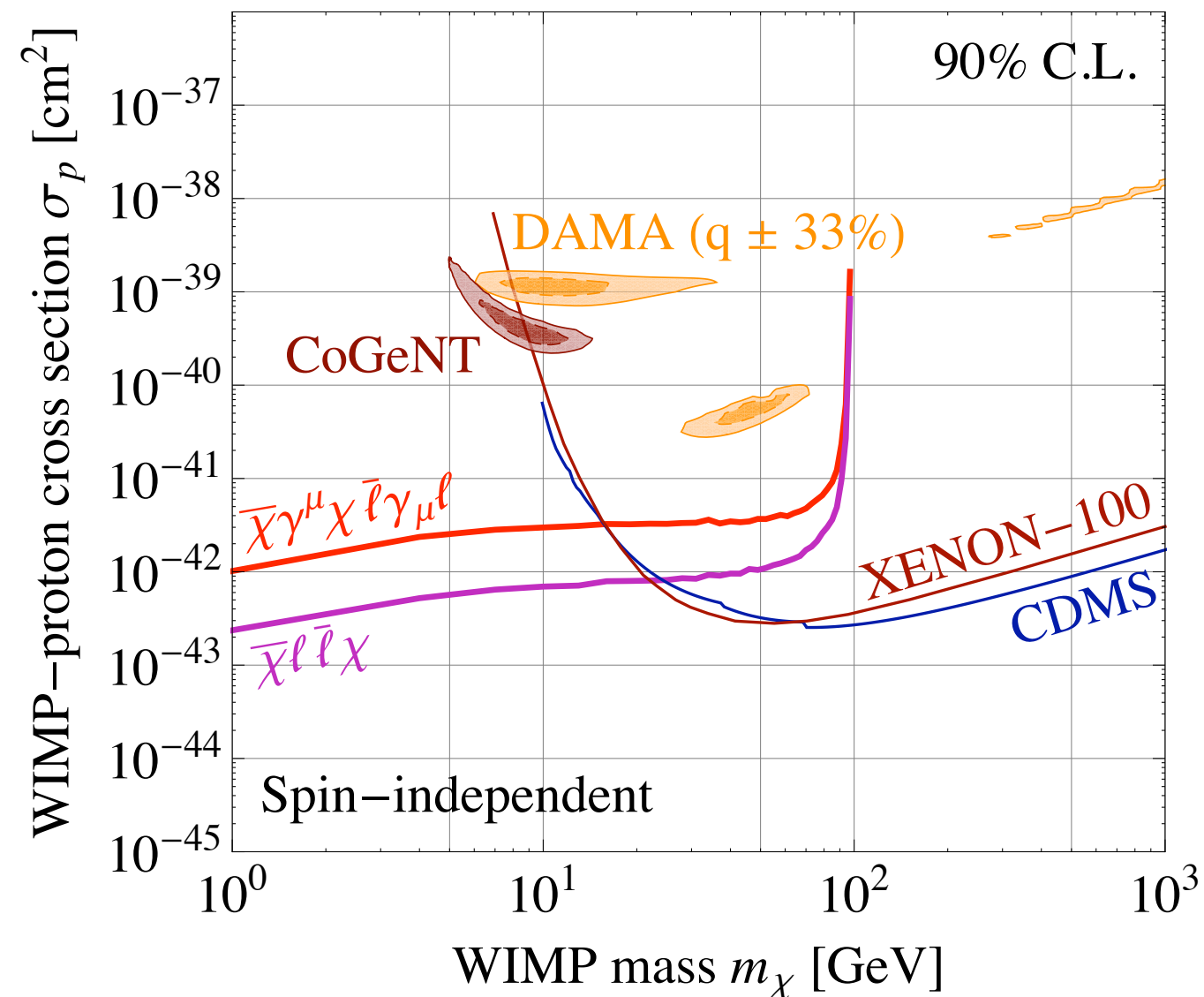
# Leptophilic DM

\* Consider zero couplings to quarks.



Direct detection  
pays a big price.  
Collider limits are strong.

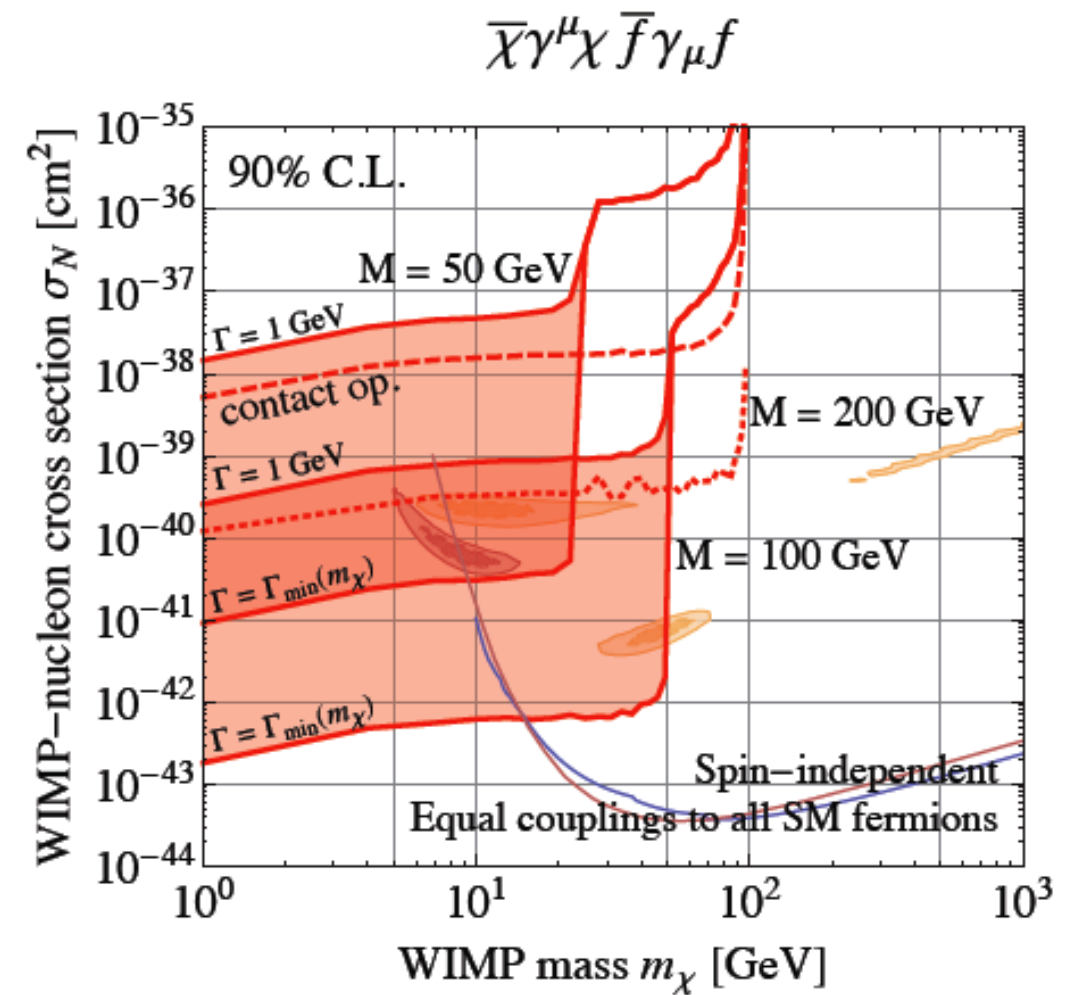
Couplings to leptons only



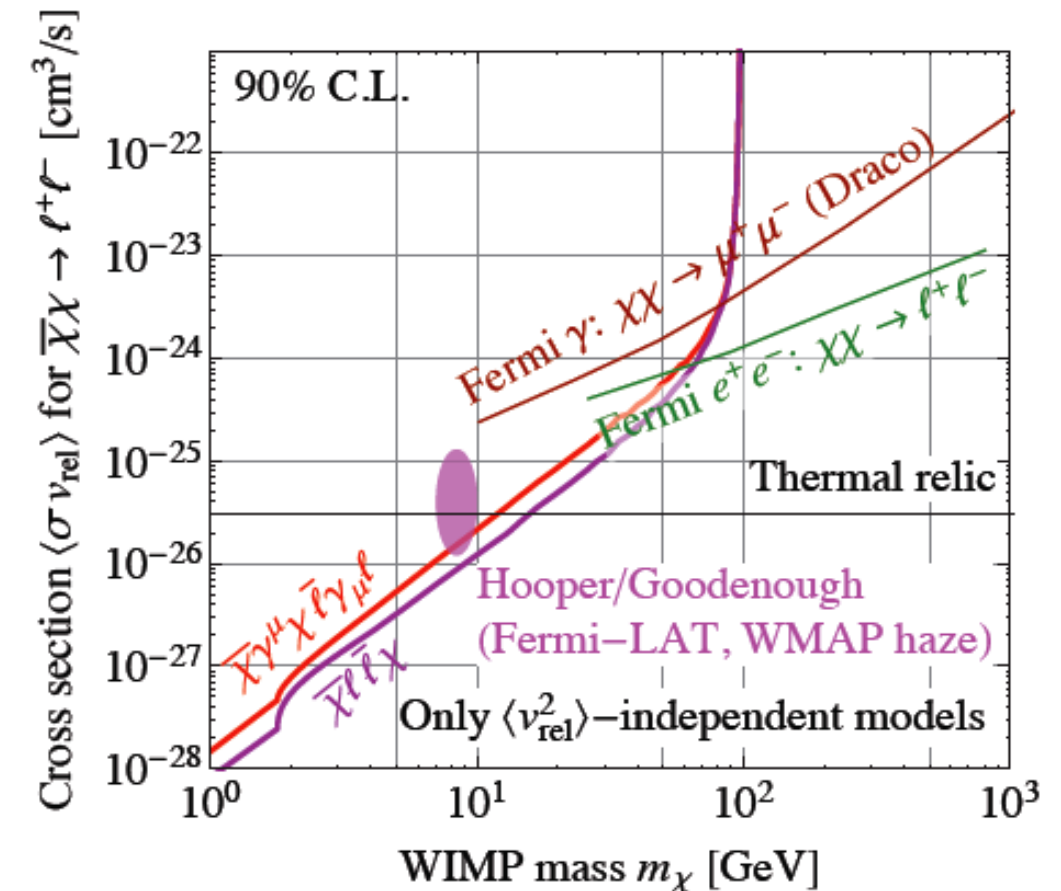
# Many more..

\* Light mediators:

\* Indirect detection:

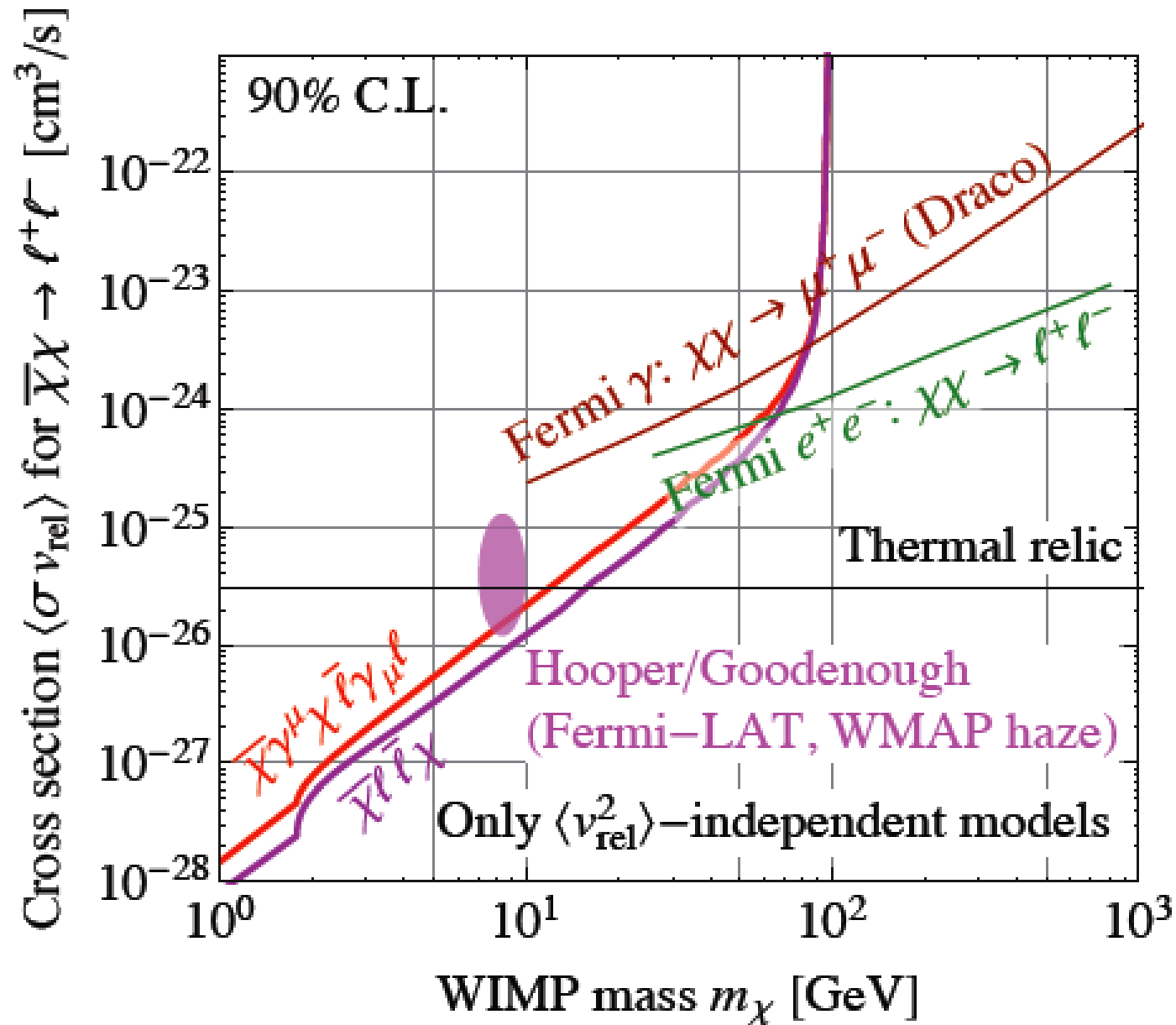


Equal coupling to all charged leptons



# Indirect Detection

Equal coupling to all charged leptons



Tension with the “Hooperon”. Light thermal relic ruled out.

# Light Mediators

\* Lets fix  $\sigma_{\text{DD}} \sim g_{\chi}^2 g_q^2 \frac{\mu^2}{M^4}$  and lower  $M$ .

The couplings must be decreased to compensate.

\* Then for very small  $M$  we get to the regime where

$$\sigma_{1j} \sim \alpha_s g_{\chi}^2 g_q^2 \frac{1}{p_T^2}$$

\* The cross section drops as  $M^4$ .

\* **But what happens in the intermediate regime?**

# Dark Matter @ Colliders

Roni Harnik, Fermilab

Bai, Fox, RH - 1005.3797

Fox, RH, Kopp, Tsai - 1103.0240

Fox, RH, Kopp, Tsai - 1109.4389

Fox, RH, Primulando, Yu - 1203.1662

Very related work by the “Irvine Clan”:

Goodman, Ibe, Rajaraman, Shepherd, Tait and Haibo Yu - 1005.1286

Goodman, Ibe, Rajaraman, Shepherd, Tait and Haibo Yu - 1008.1783

Fortin and Tait - 1103.3289

Rajaraman, Shepherd, Tait and Wijangco - 1108.1196

Shepherd and Goodman - 1111.2359



# A Search For Dark Matter in the Monojet + Missing Transverse Energy Signature in $6.7 \text{ fb}^{-1}$

S.Z. Shalhout<sup>1</sup>, T. Schwarz<sup>2</sup>, R. Erbacher<sup>1</sup>, J. Conway<sup>1</sup>, P. Fox<sup>2</sup>, R. Harnik<sup>2</sup>, Y. Bai<sup>2</sup>  
UC Davis<sup>1</sup> Fermilab<sup>2</sup>

A neural net with our name on it ?! :-0

