



Revealing Deaths of Massive Stars with High-Energy Neutrinos

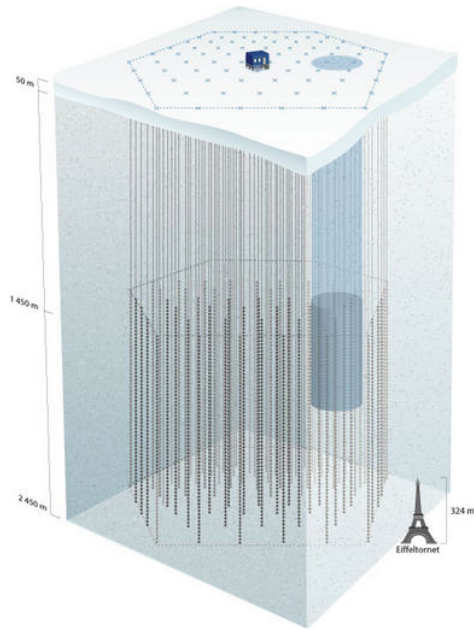
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**Realtime Astroparticle Physics
February 2013**

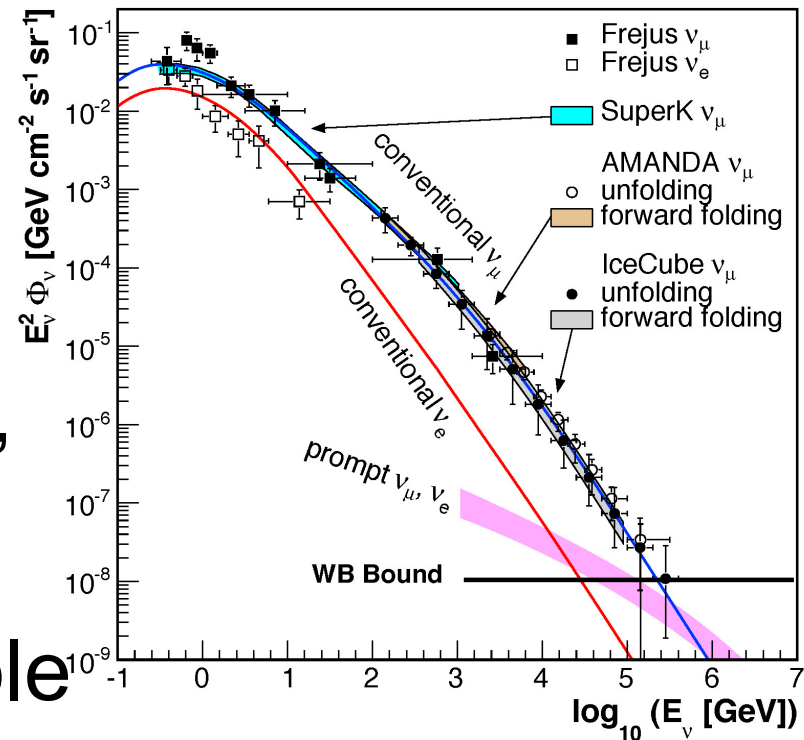


High-Energy Neutrinos



Now we have IceCube
Interesting era!

>PeV: cosmic-ray origin
connection w. γ rays
<PeV: atm. $\nu \rightarrow$ “transients”
more compact
 γ rays may be invisible



Outline

Astrophysical scenarios for HE neutrinos from explosive phenomena such as GRBs and SNe

- Origin of extragalactic cosmic rays
- Physical mechanisms, GRB-SN connection etc.

1. Gamma-Ray Bursts
2. Supernovae

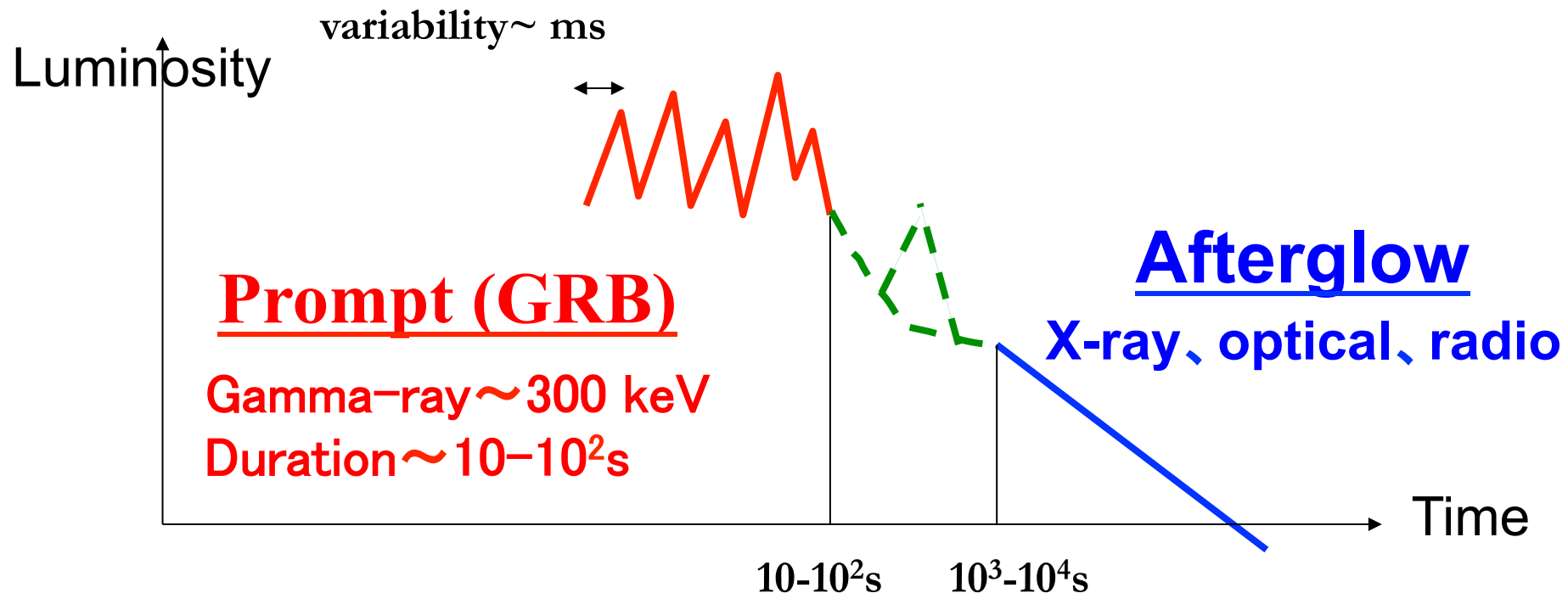


1. Gamma-Ray Bursts

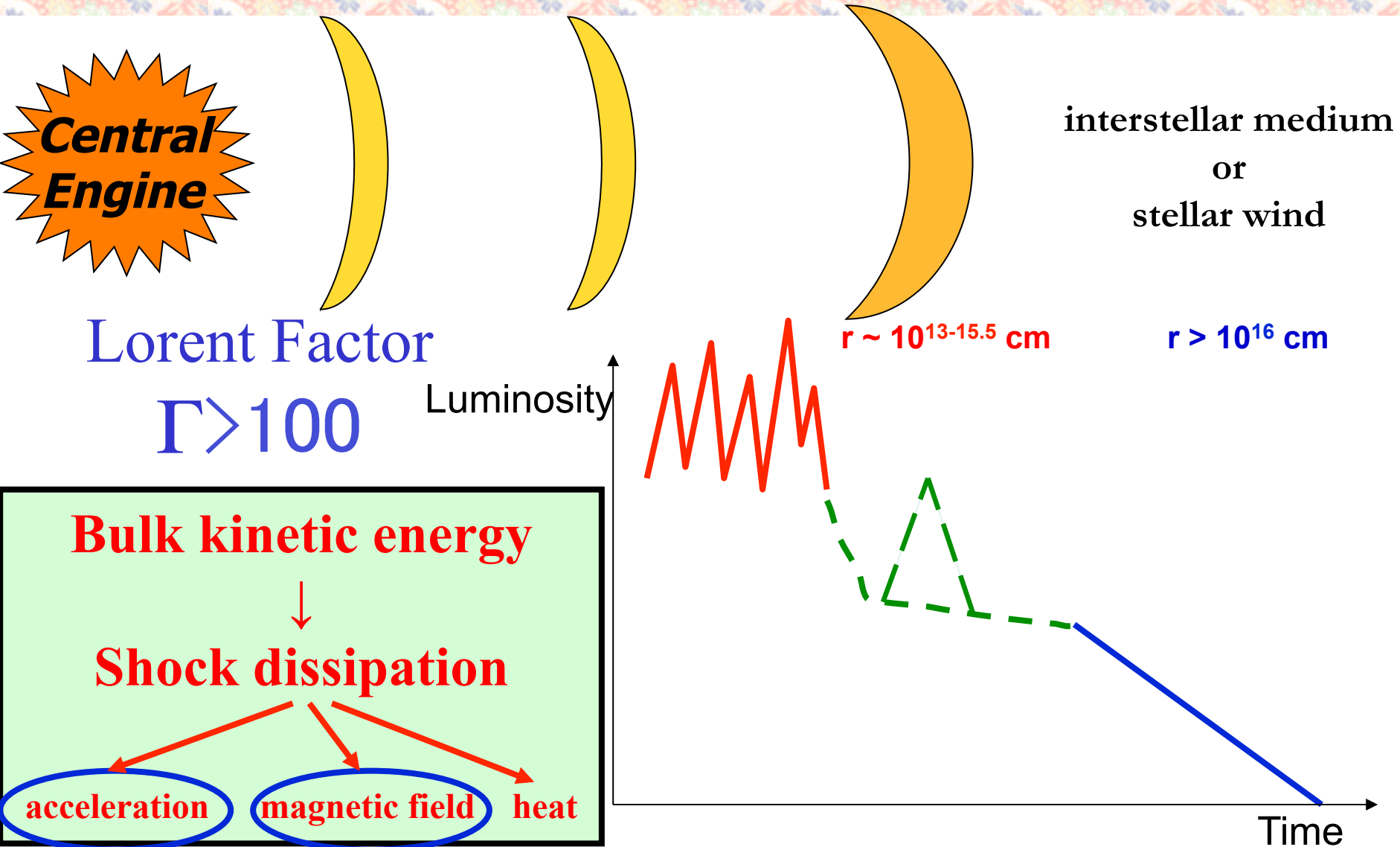


(Long) Gamma-Ray Bursts

- The most violent phenomena in the universe ($L_\gamma \sim 10^{51-52}$ ergs s $^{-1}$)
- Cosmological events ($z \sim 1-3$), ~ 1000 per year
- Relativistic jet ($\Gamma \sim 100-1000$; $\theta_{\text{jet}} \sim 0.1$ rad)
- Related to death of massive stars (association with supernovae)



“Classical” Internal-External Shock Model (Baryonic Jet Model)



Ultra-High-Energy Cosmic Rays?

Fermi shock acceleration mechanism

-> not only electrons but protons are accelerated

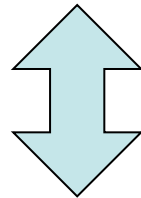
$$\varepsilon_p < e r B \sim 3 \times 10^{20} \text{ eV } r_{14} B_4 \text{ (Waxman 1995)}$$

If UHECR energy output \sim GRB radiation energy

$$E_{\text{HECR}}^{\text{iso}} \sim E_{\gamma}^{\text{iso}} \sim 10^{53} \text{ erg}$$

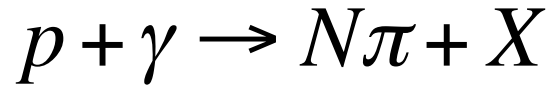
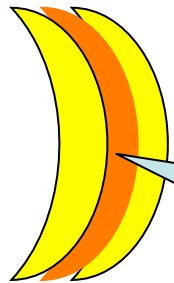
with local GRB rate density: $\sim 1 \text{ Gpc}^{-3} \text{ yr}^{-1}$

(e.g., Wanderman & Piran 2010)



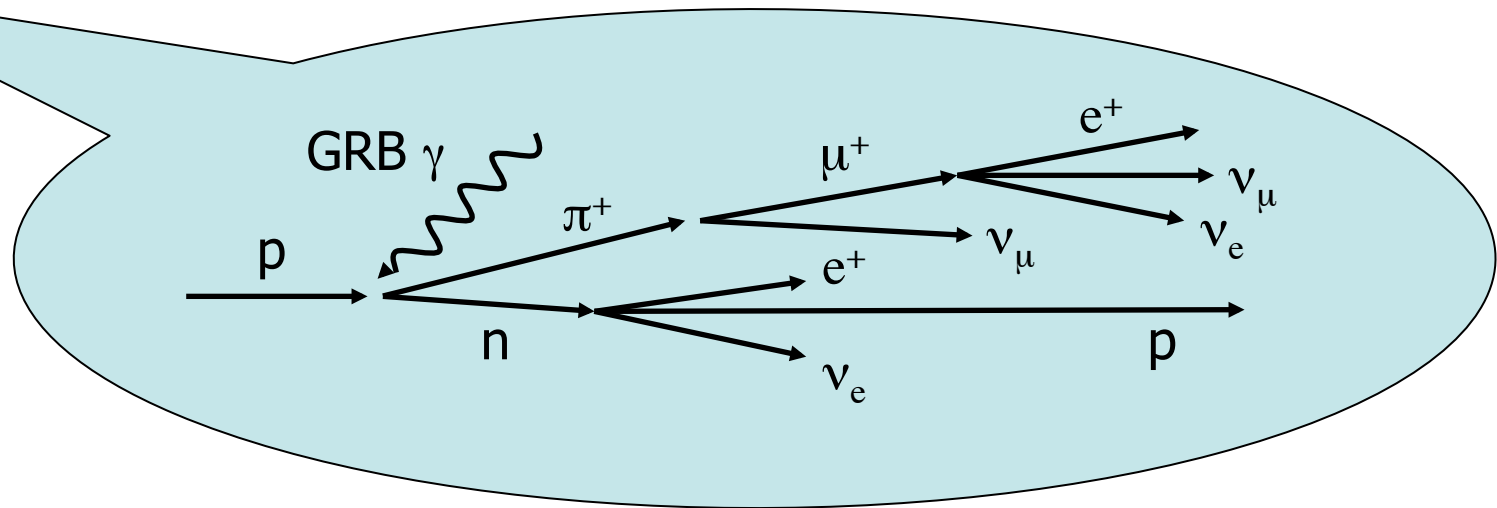
UHECR budget (from obs.): $Q_{\text{HECR}} \sim 10^{44} \text{ erg/Mpc}^3/\text{yr}$

Neutrino Production in the Source



$$\sigma_{p\gamma} \sim \text{a few} \times 10^{-28} \text{ cm}^2$$

**baryonic resonances,
direct production,
multi-pion production etc.**



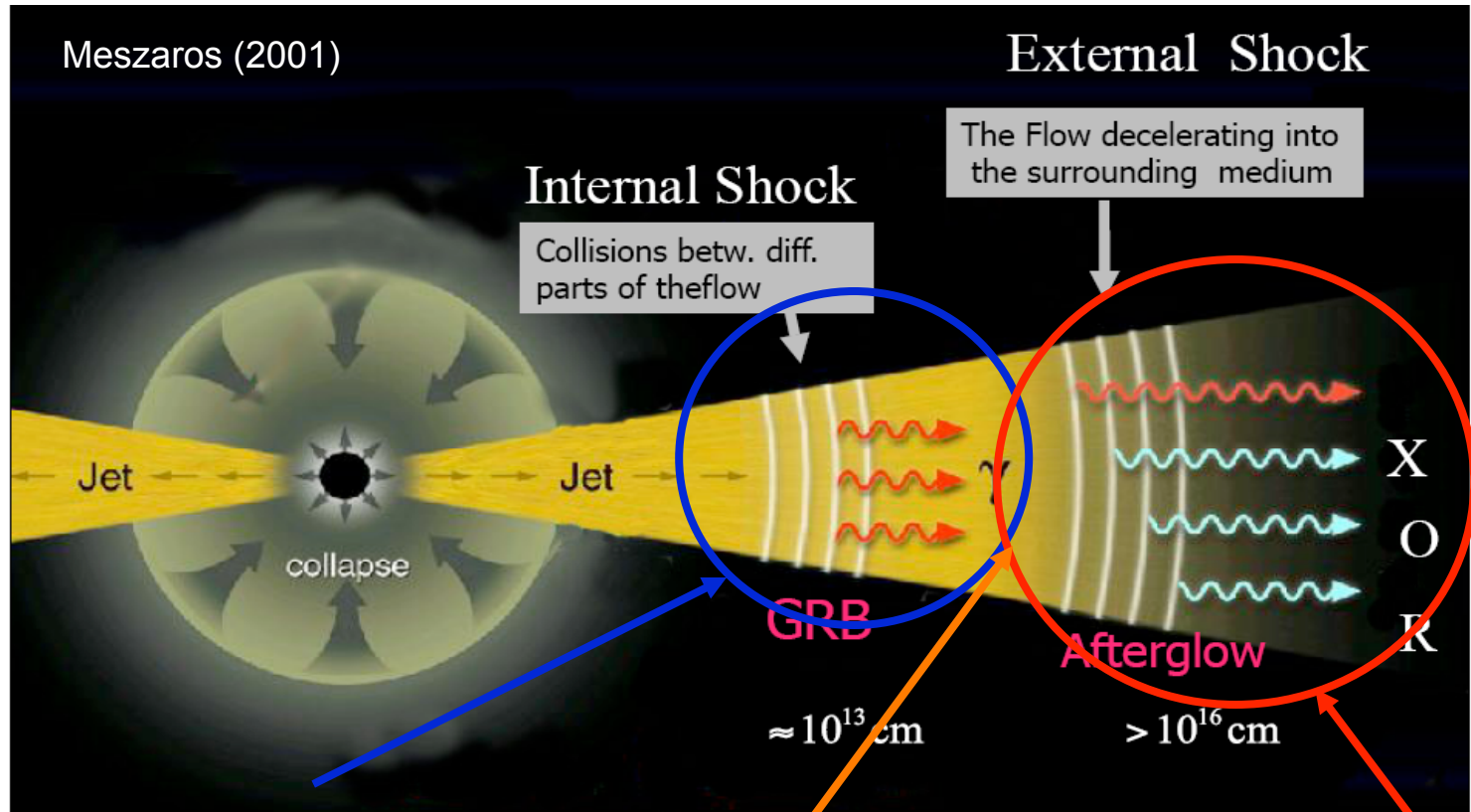
at Δ -resonance ($\epsilon_p \epsilon_\gamma \sim 0.2 \Gamma^2 \text{ GeV}^2$)

$$\epsilon_\nu^b \sim 0.05 \epsilon_p^b \sim 0.01 \text{ GeV}^2 \Gamma^2 / \epsilon_{\gamma, pk} \sim \mathbf{1 \text{ PeV}}$$
 (if $\epsilon_{\gamma, pk} \sim 1 \text{ MeV}$)

Meson production efficiency (large astrophysical uncertainty)

$$\mathbf{f_{p\gamma}} \sim 0.2 n_\gamma \sigma_{p\gamma} (r/\Gamma) \propto \mathbf{r^{-1} \Gamma^{-2}} \propto \mathbf{\Gamma^{-4} \delta t^{-1}}$$
 (if $r \sim \Gamma^2 \delta t$)

CR Acceleration in “Classical” Pictures



Inner jet (prompt emission)

$r \sim 10^{12}-10^{16}$ cm $B \sim 10^{2-6}$ G

PeV ν , GeV-TeV γ

Waxman & Bahcall 97 PRL

Inner jet (flares)

$r \sim 10^{14}-10^{16}$ cm $B \sim 10^{2-4}$ G

PeV-EeV ν , GeV-TeV γ

KM & Nagataki 06 PRL

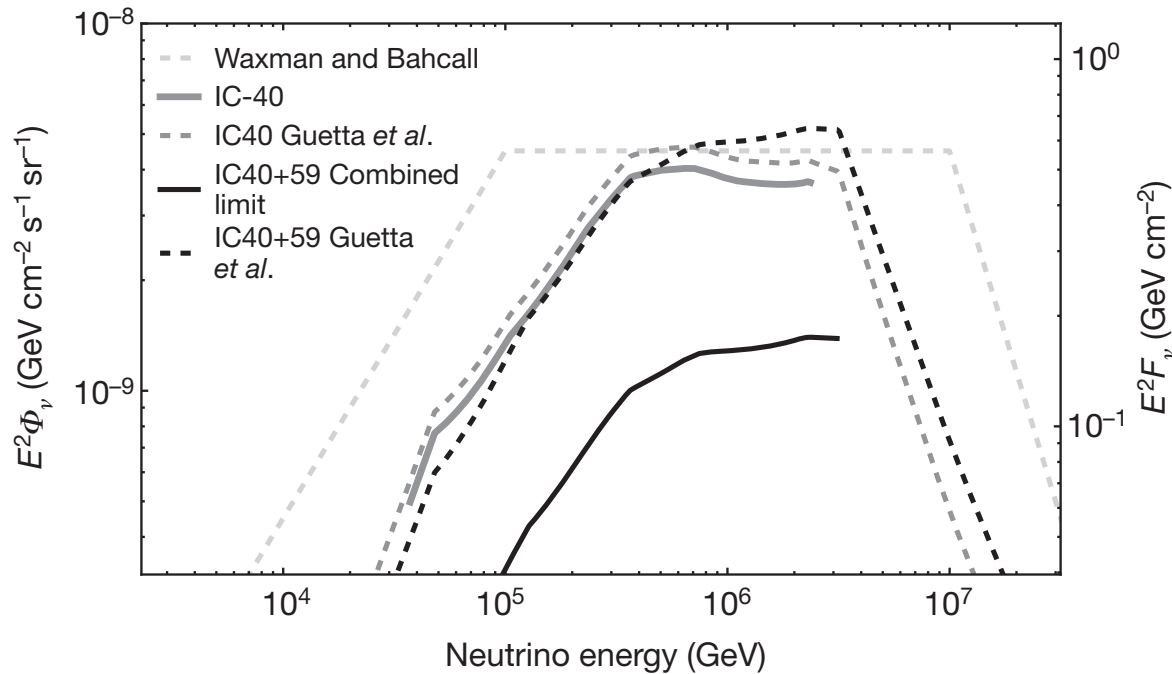
External shock (afterglow)

$r \sim 10^{16}-10^{17}$ cm $B \sim 0.1-100$ G

EeV ν , GeV-TeV γ

e.g., Waxman & Bahcall 00,
Dermer 02, KM 07

Recent IceCube Limits on Prompt ν Emission



IceCube collaboration 12 Nature

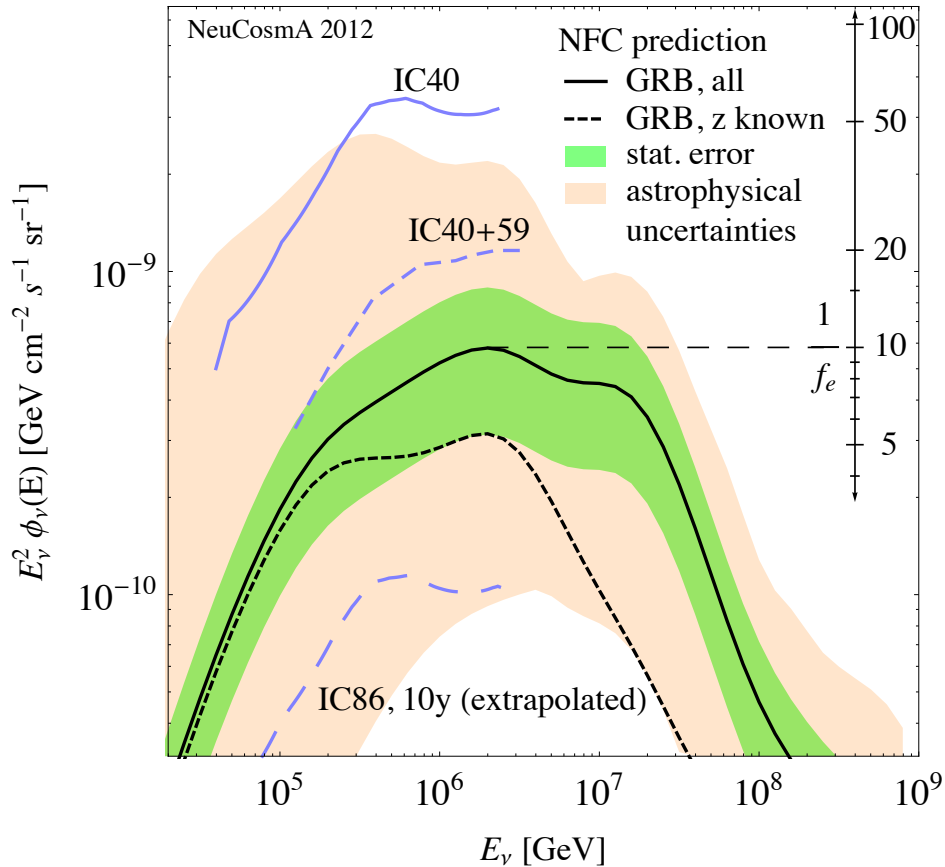
But theoretical fluxes should be lower than IceCube-Guetta *et al.*

- $f_{p\gamma}$ is energy-dependent, π -cooling $\rightarrow \sim 4 \downarrow$** (Li 11 PRD, Hummer *et al.* 12 PRL)
- $(\epsilon_\gamma^2 \phi_\gamma \text{ at } \epsilon_{\gamma,pk}) \neq (\int d\epsilon_\gamma \epsilon_\gamma \phi_\gamma) \rightarrow \sim 3-6 \downarrow$** (Hummer *et al.* 12 PRL, He *et al.* 12 ApJ)
- details (multi- π , ν mixing etc.) \rightarrow ex., multi- $\pi \sim 2-3 \uparrow$** (KM & Nagataki 06 PRD)

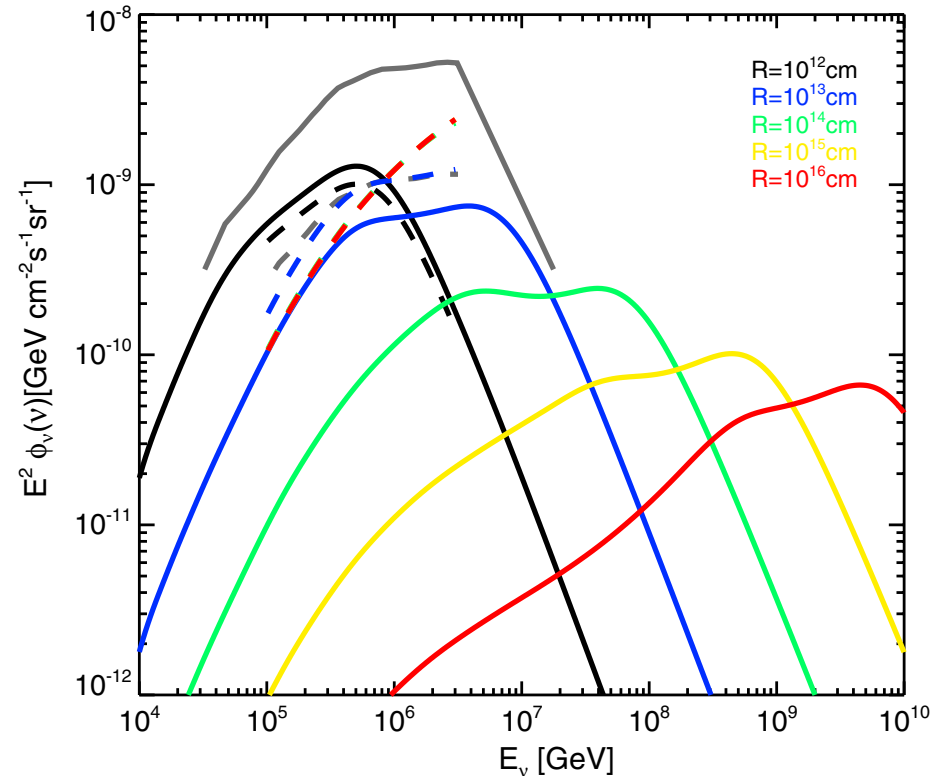
- In addition, there is “astrophysical” model-uncertainty in calculating $f_{p\gamma}$

Recent IceCube Limits on Prompt ν Emission

Hummer, Baerwald & Winter 12 PRL



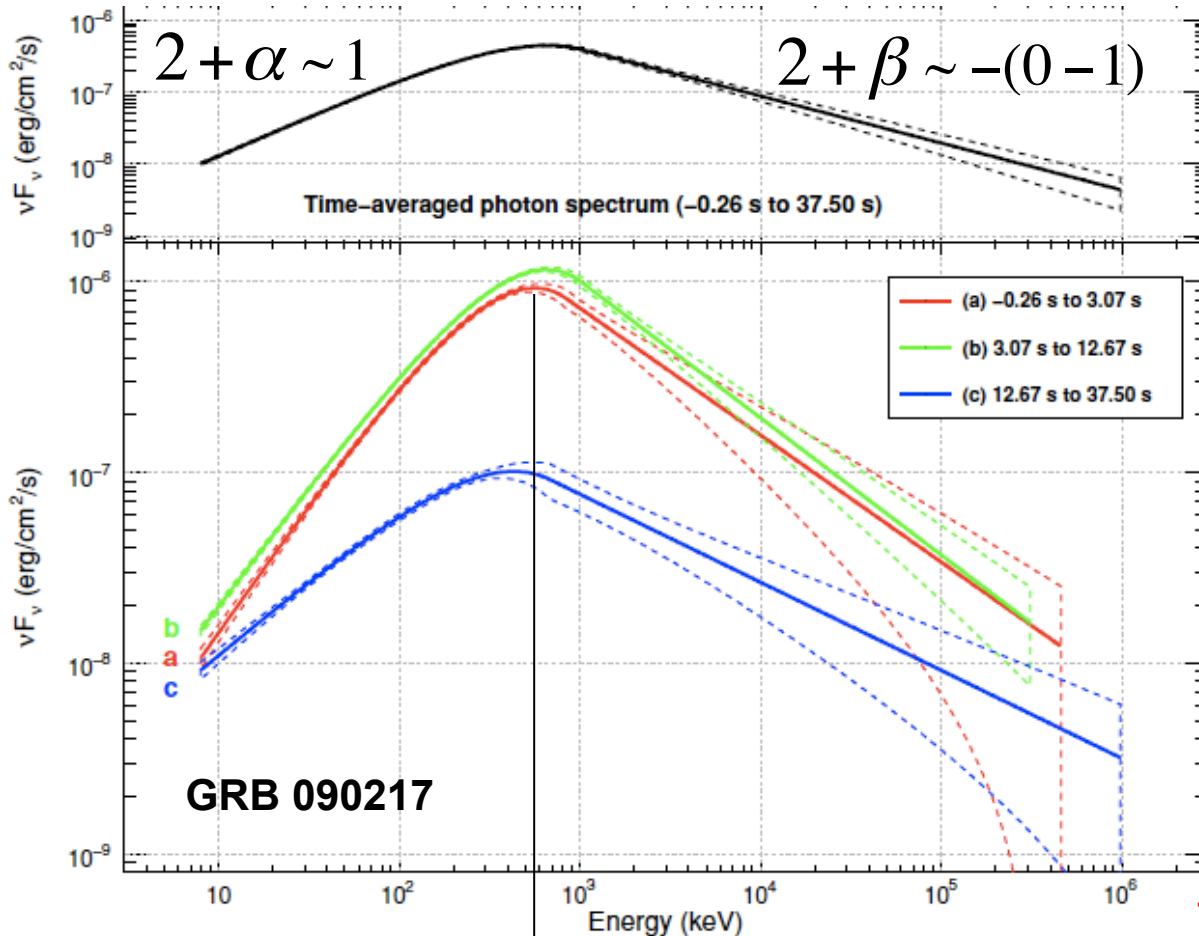
He, Liu, Wang, Nagataki, KM & Dai 12 ApJ



~10 yr observations by IceCube can cover reasonable parameter space required for the GRB-UHECR scenario

Problems in Internal-Shock-Synchrotron Scenario

Fermi collaboration 10 ApJ



$$\mathcal{E}_{\gamma, \text{pk}} \sim 0.1 - 1 \text{ MeV}$$

Band function
 ~ broken power-law

$$F_{\nu} \propto \nu^{\beta+1}$$

“synchrotron scenario”

$$\mathcal{E}_{\gamma, \text{pk}} = \Gamma \hbar \gamma_{ei}^2 \frac{eB}{m_e c}$$

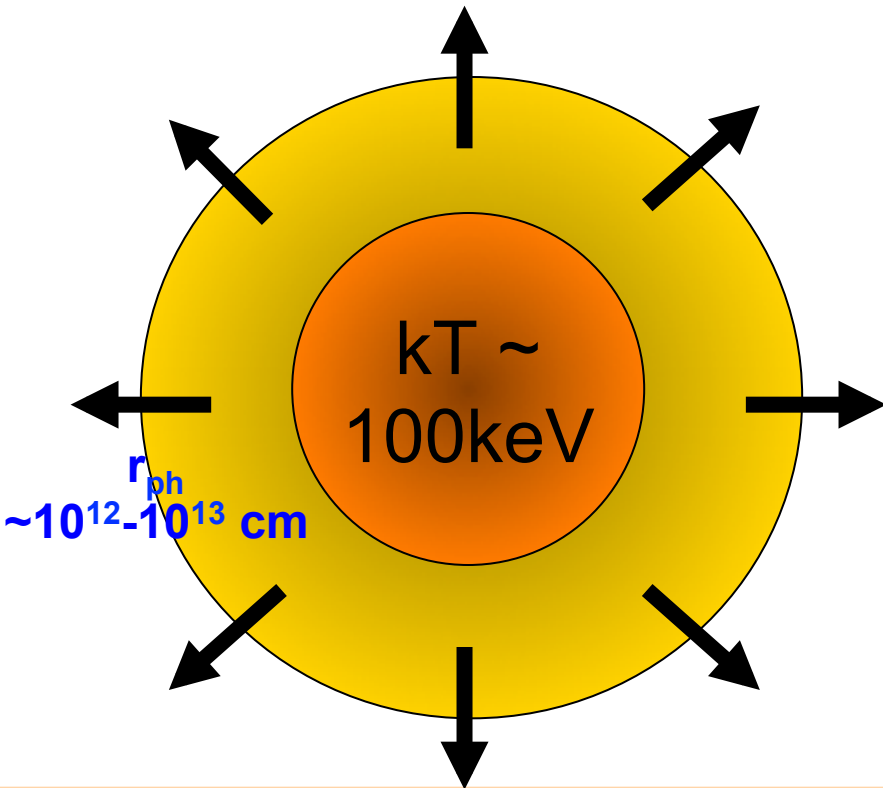
$2 + \alpha \sim 0.5$ (fast-cooling)

Other problems

- high radiative efficiency
- empirical relations ($\mathcal{E}_{\gamma, \text{pk}} - L_{\gamma}$)
- theoretical issues

Dissipative Photosphere Scenario

e.g., Thompson 1994, Meszaros & Rees 2000, Rees & Meszaros 2005,
Peer et al. 2006, Giannios 2006, Ioka, KM et al. 2007, Beloborodov 2010

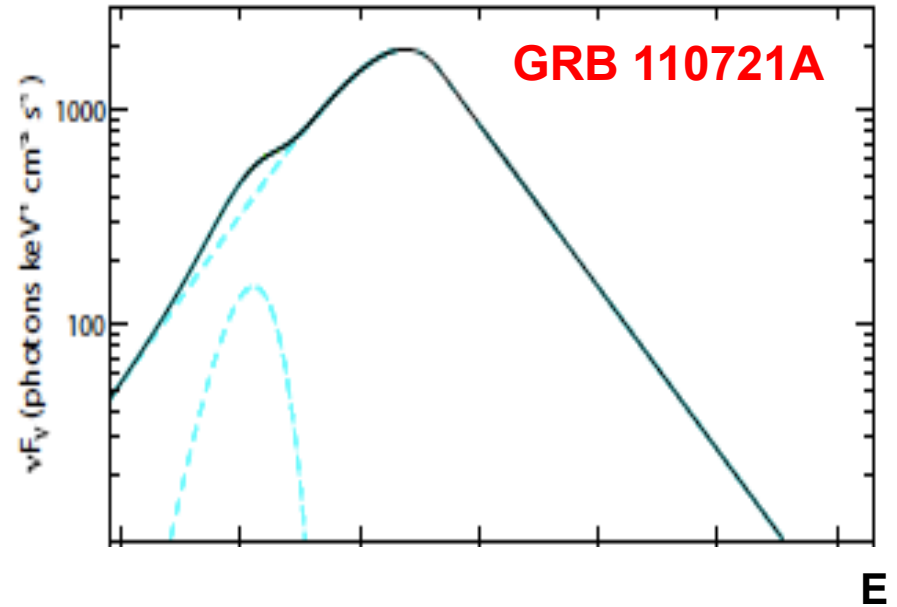
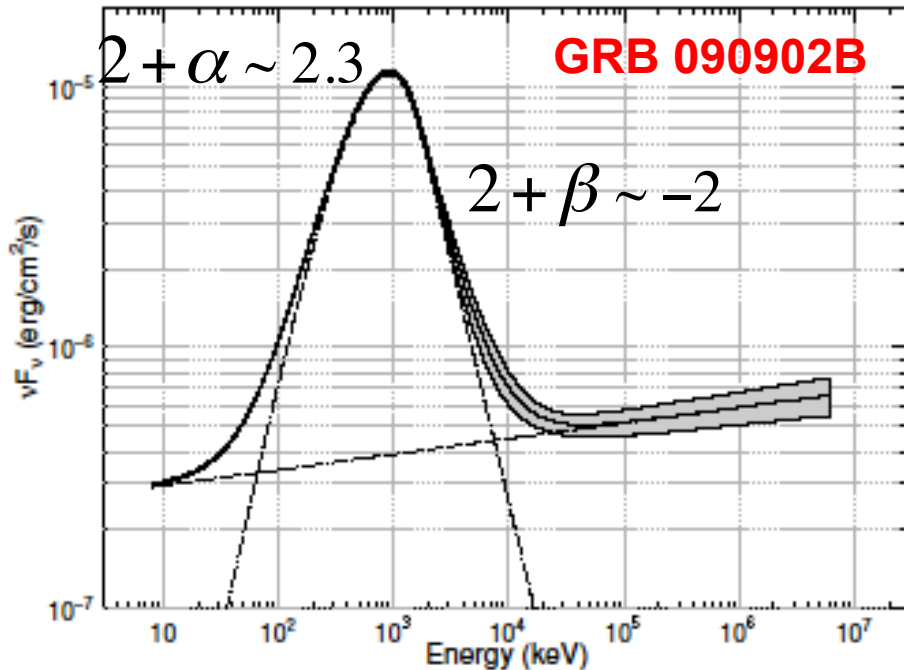


Emissions from $\tau_T \sim 1-10$ “dissipative photosphere”

- internal shocks
- interaction with star or wind
- recollimation shocks
- magnetic reconnection
- collisions with neutrons

- Re-conversion of kinetic energy to radiation energy
- High radiative efficiency & stabilization of $\epsilon_{\gamma, \text{pk}}$

Observational Hints



“modified” black-body emission

or

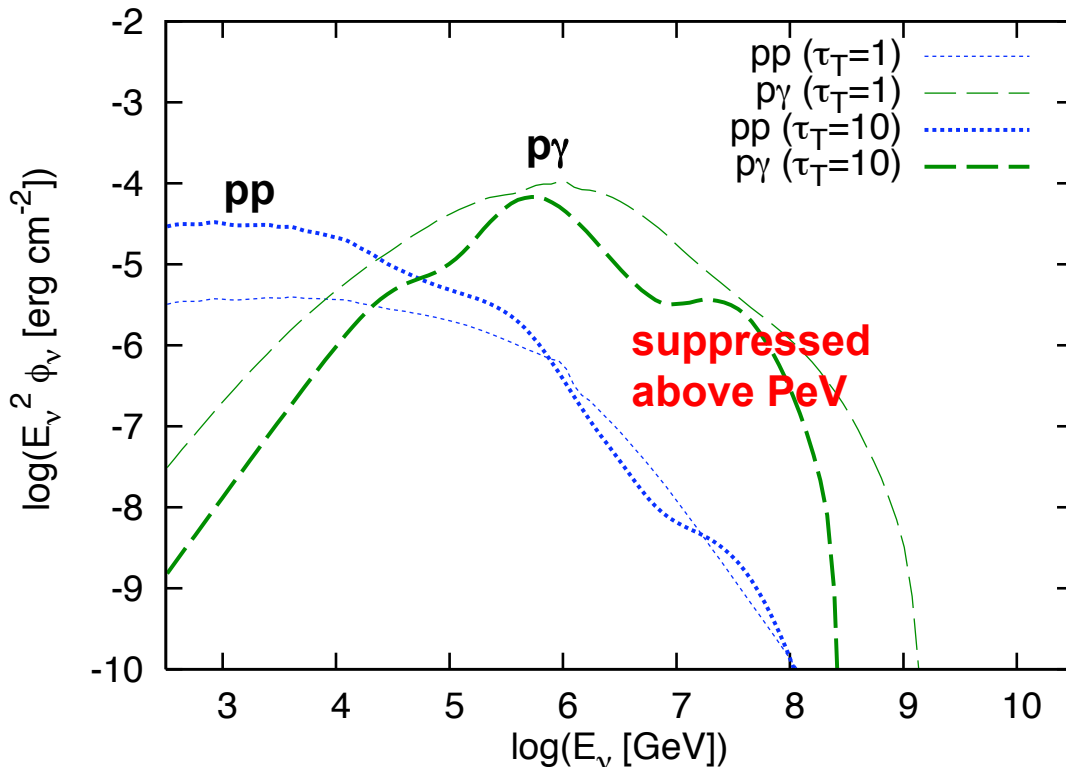
thermal + nonthermal emission

Observed spectra can be reproduced by theories

Photospheric Neutrinos

- Dissipative photosphere** (e.g., Rees & Meszaros 05 ApJ)

$$\tau_T = n_e \sigma_T (r/\Gamma) \sim 1-10 \Leftrightarrow f_{pp} = (\kappa_{pp} \sigma_{pp} / \sigma_T) \tau_T \sim 0.05-0.5$$



$$\tau_T = 1-10 \quad (r \sim 10^{12}-10^{12.5} \text{ cm})$$

$$\Gamma = 10^{2.5}, \quad U_e = U_B$$

$$E_{\text{CR}}^{\text{iso}} \sim E_\gamma^{\text{iso}} \sim 10^{53.5} \text{ erg}$$

→ # of $\mu\text{s} \sim 1-2$ for GRB @z=0.1

KM, PRD(R), 78, 101302 (2008)
 cf. Wang & Dai, ApJL, 691, L67 (2009)

Detection of pp neutrinos strongly supports dissipative photospheres

Quasi-Thermal Neutrinos?

We have assumed ε_p^{-2} spectrum w. $E_{\text{CR}}^{\text{iso}} \sim E_{\gamma}^{\text{iso}}$
But highly uncertain....

- $p\gamma$ is unimportant (pp is more important) if steeper
- non-thermal component may be absent

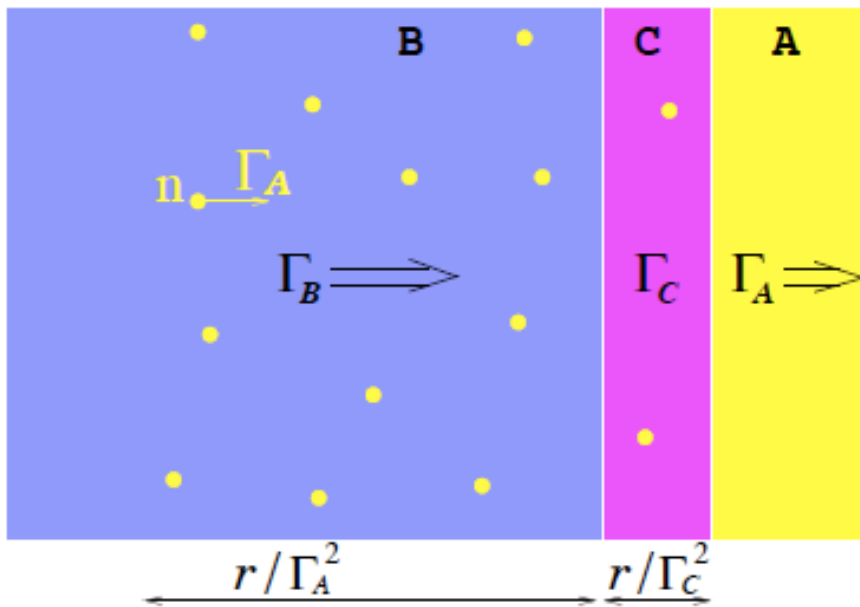
Q. Can we expect “thermal” neutrinos?

A. Yes! (Paczynski & Xu 1994, Bahcall & Meszaros 2000, Meszaros & Rees 2000)

inevitable when **neutrons** are loaded in jets

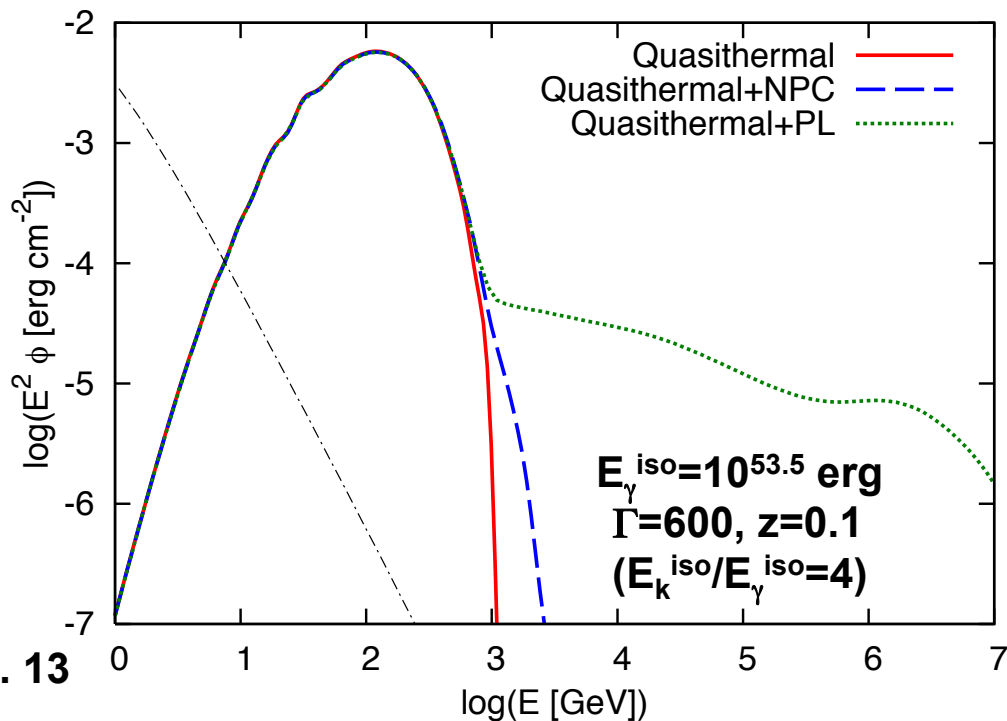
(But uncertain if not, since thermal protons would not cause inelastic collisions for radiation-mediated shocks)

Collisional Dissipation by Neutrons



Decoupled neutrons are swept up by the faster flow w. $\Gamma_B \gg \Gamma_A = \Gamma_n$
inelastic pn collision

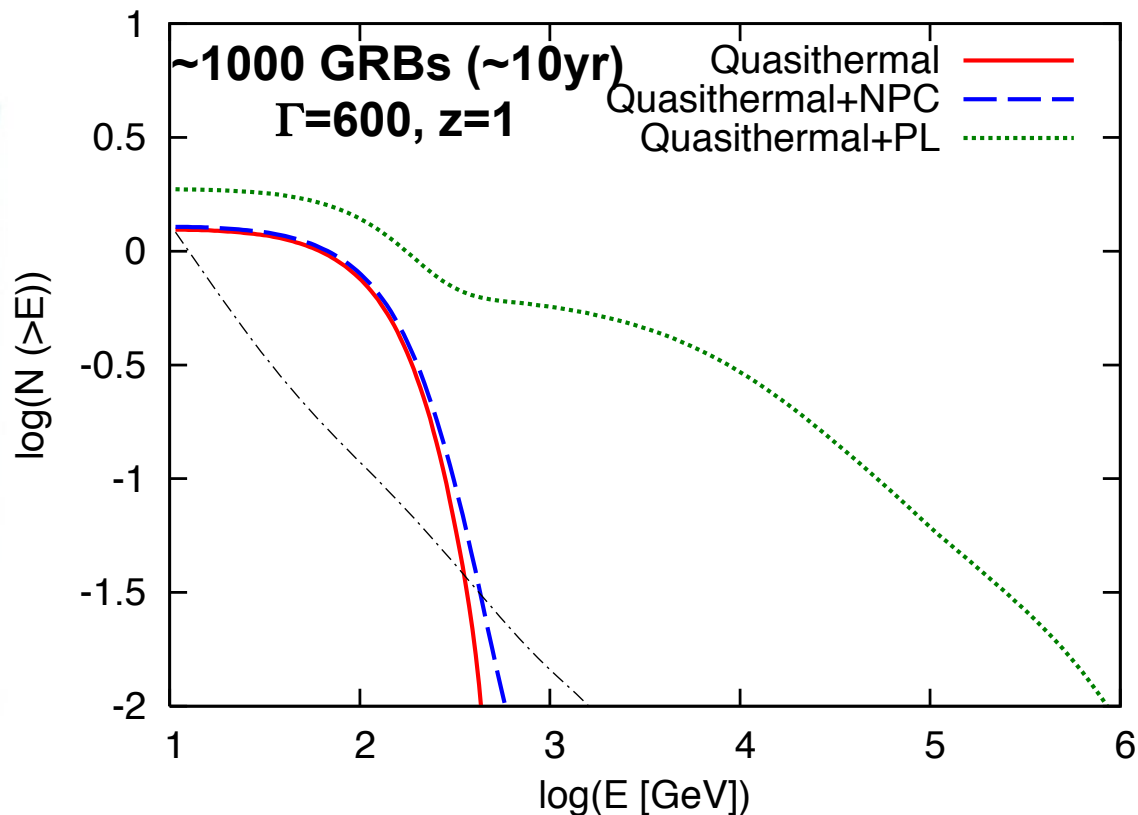
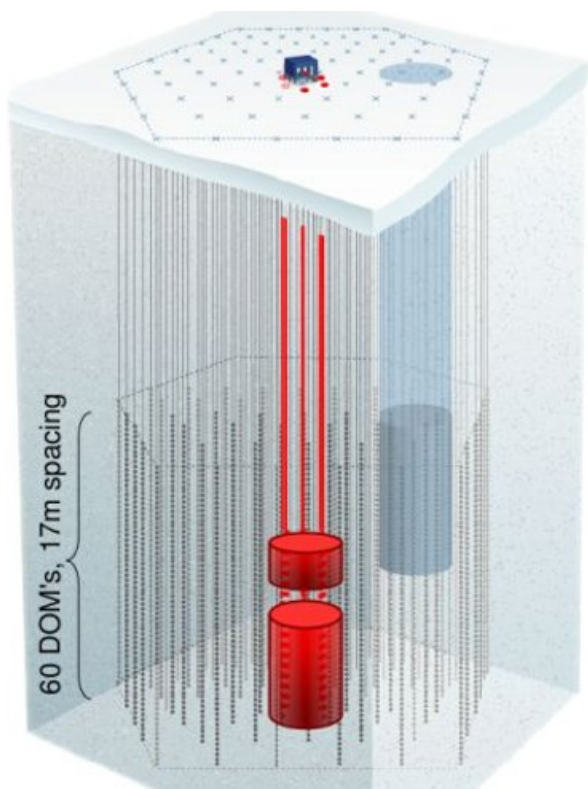
↓
 γ rays \Leftrightarrow prompt emission
 + **neutrinos!**



- “Quasi-thermal”
 $\epsilon_\nu \sim 100 \text{ GeV } (\Gamma/500) (\Gamma_{\text{rel}}/2)$
- γ s & ν s come from ps
 $\epsilon_\nu^2 \phi_\nu \sim \epsilon_\gamma^2 \phi_\gamma$ (normalizable)

Prospects for DeepCore+IceCube

- Including DeepCore is essential at **10-100 GeV**
- Reducing atmospheric ν background is essential
→ select only bright GRBs w. $> 10^{-6}$ erg cm $^{-2}$



Summary: Testable Cases for GRB Neutrinos

vs in GRB-UHECR hypothesis

- Prompt ν s: **PeV**
 - ~ 10 yr to cover parameter space in classical scenario
 - maybe difficult in magnetic scenarios
- Afterglow cases are allowed (\rightarrow Askaryan Radio Array)

Prompt ν s in dissipative photosphere scenario

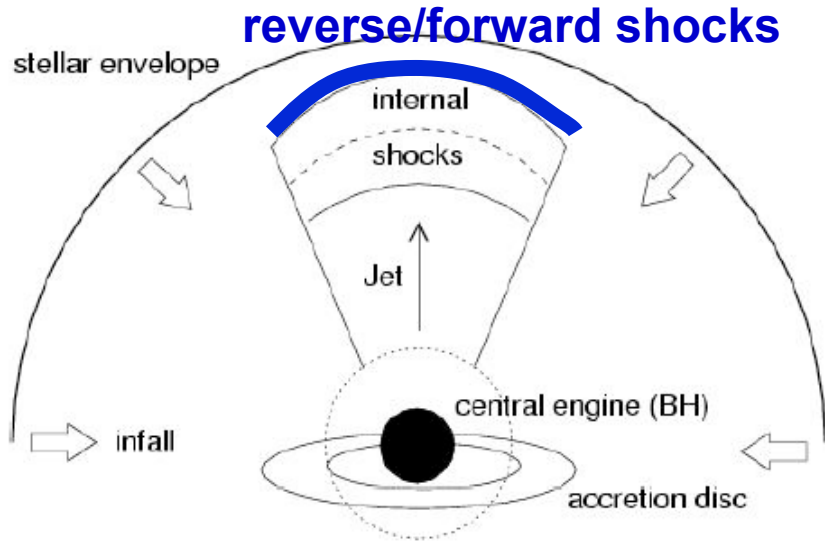
- Different predictions (**GeV-TeV** from pp, suppression at $> \text{PeV}$)
 - Consistent w. non-observations by IceCube
 - **Better** things
 - pp(or pn)/p γ efficiencies are fixed by $\tau_T \sim 1-10$
 - inevitable quasi-thermal ν s in pn collisional dissipation
- \rightarrow **searches w. DeepCore, improving ang. resolution for cascades**



2. Supernovae



TeV Neutrinos from Choked Jets



Razzaque et al. 2005

- **Jet penetration?**
GRBs=successful jets
failed GRBs=SNe w. choked jets

(Meszaros & Waxman 01 PRL)

- **Slow jets embedded in SNe?**
Some SNe may be driven by a slow jet

(Razzaque et al. 04 PRL,
Ando & Beacom 05 PRL, Horiuchi & Ando 08 PRD)

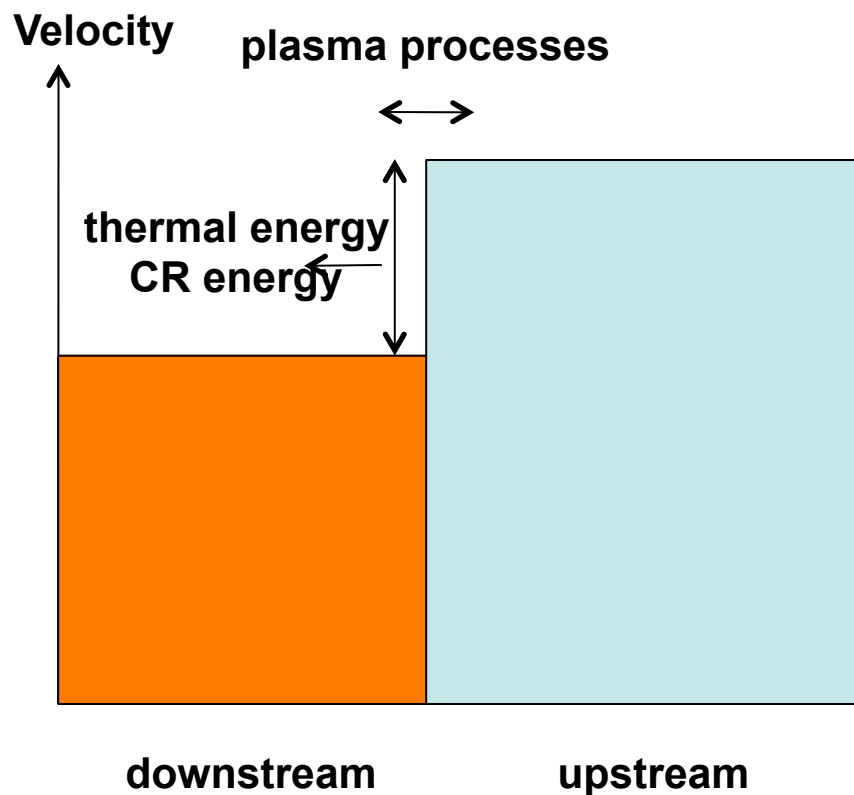
If CRs carry $E_{\text{CR}}^{\text{iso}} \sim 0.5 \times 10^{53}$ erg (GRB)
→ # of $\mu\text{s} \sim 30$ for SN@10Mpc

Neutrinos from such SNe are interesting if detected
(neutrino tomography, neutrino mixing etc.)

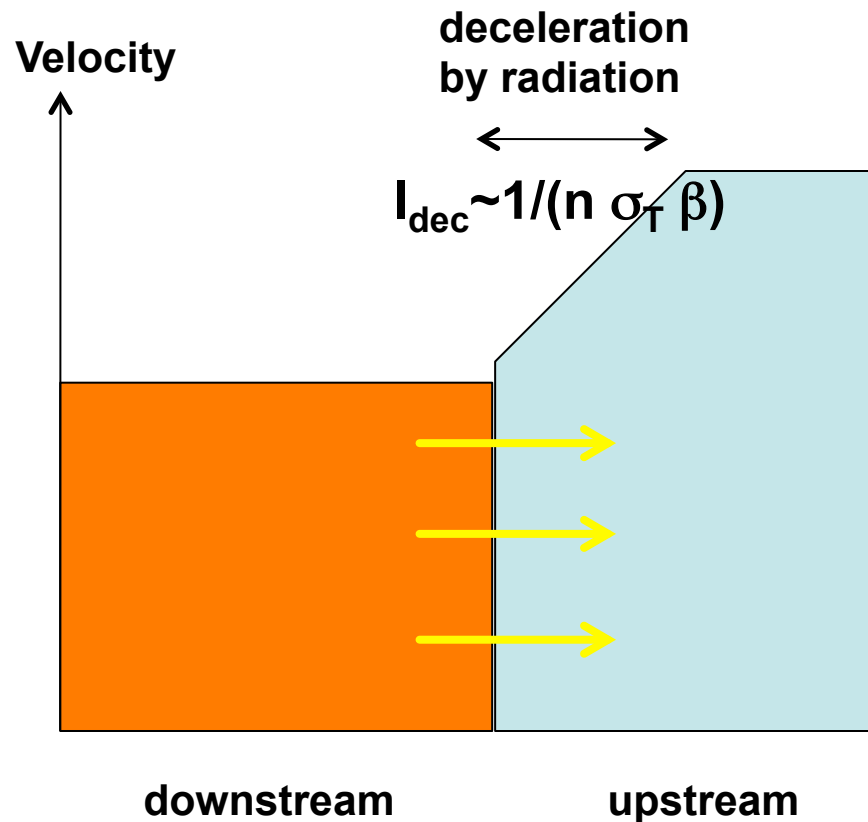
But can we expect CR acceleration deep inside stars?

Limitation of Shock Acceleration

Collisionless shock



Radiation-mediated shock (ex. Weaver 76 ApJ, Katz+ 10 ApJ, Nakar & Sari 12 ApJ)



Shock Breakout & Collisionless Shocks

- Necessary condition for collisionless shocks

$$l < \sim l_{\text{dec}} \sim (1/n \sigma_T \beta) \Leftrightarrow \tau_T < \sim 1/\beta$$

(not sufficient condition: ex. steep density profile)

(Waxman & Loeb 01 PRL, KM et al. 11 PRD, Katz, Sapid & Waxman 11)

- Shock breakout: $t_{\text{diff}} \sim t_{\text{dyn}} \Leftrightarrow \tau_T \sim 1/\beta$

$$t_{\text{diff}} \sim l^2/\kappa \quad (\kappa \sim (c/n \sigma_T))$$

$$t_{\text{dyn}} \sim l/\beta c$$

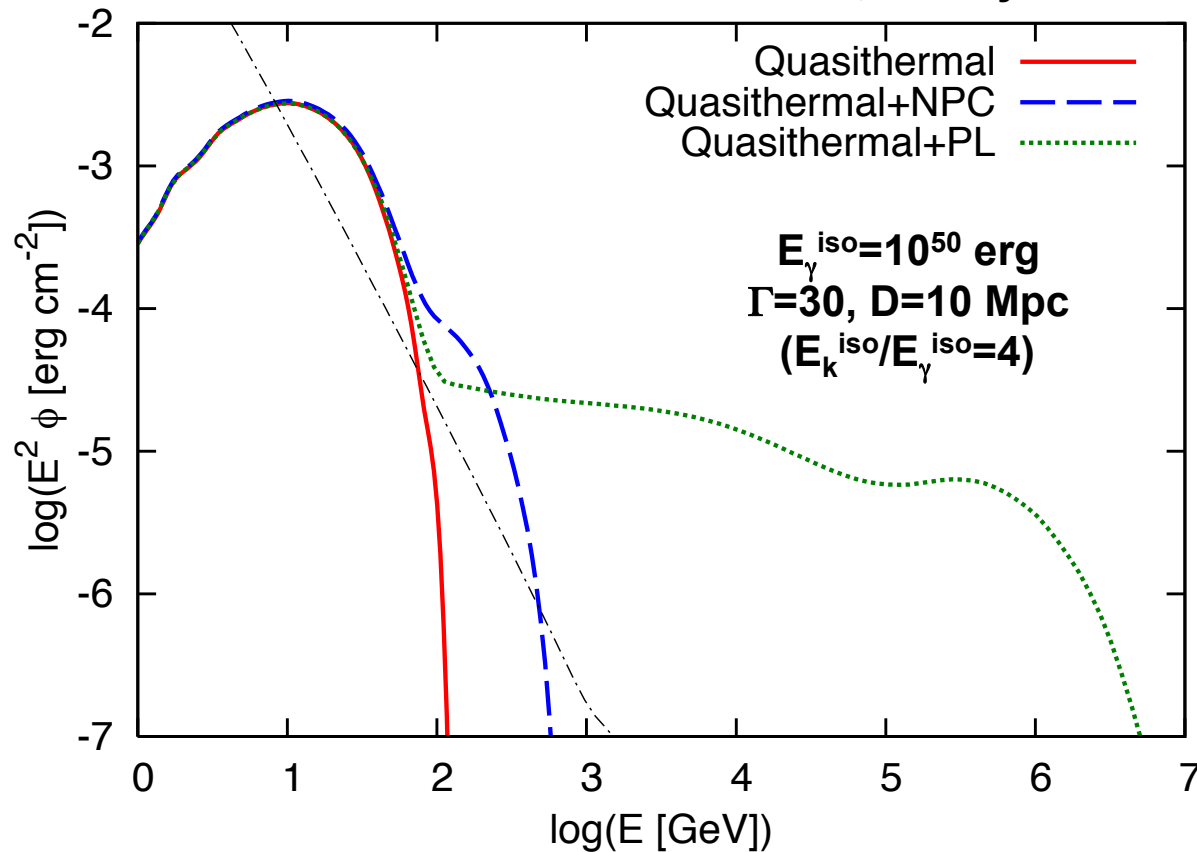
$$\text{wind CSM} \rightarrow r_{\text{bo}} \sim l_{\text{bo}} \sim (1/n \sigma_T \beta) \quad (\text{unless ultra-relativistic})$$

Ex. int./rev. shock at $r=10^9$ cm in choked jets ($L_k=10^{48}$ erg/s, $\Gamma=10$)

$\rightarrow \tau_T \sim 10^3$, CR acc. is difficult (see also Levinson & Bromberg 08 PRL)

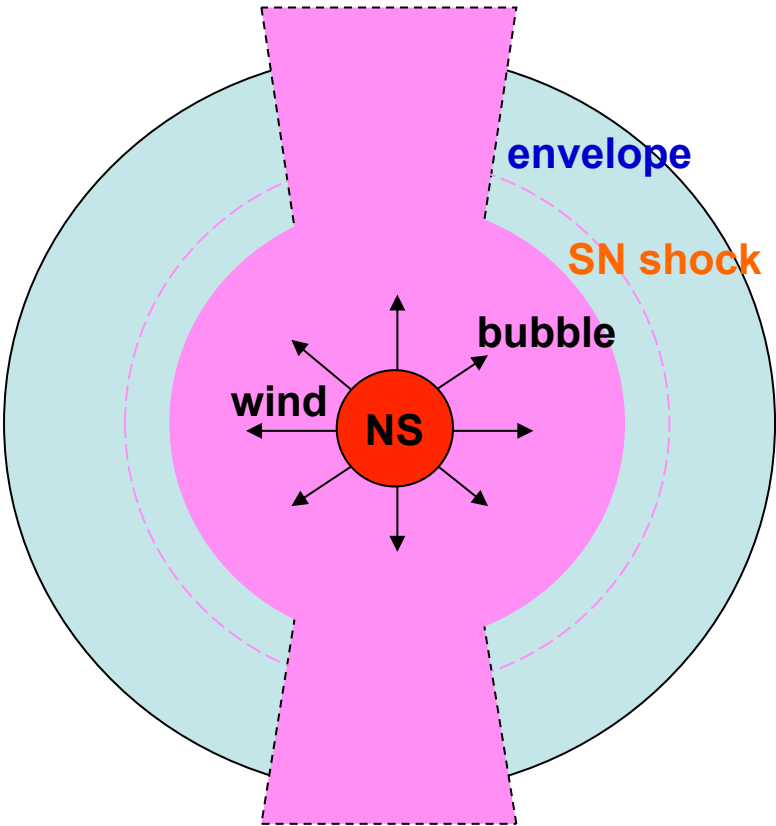
Possibility: Quasi-Thermal Neutrinos

KM, Kashiyama & Meszaros 13



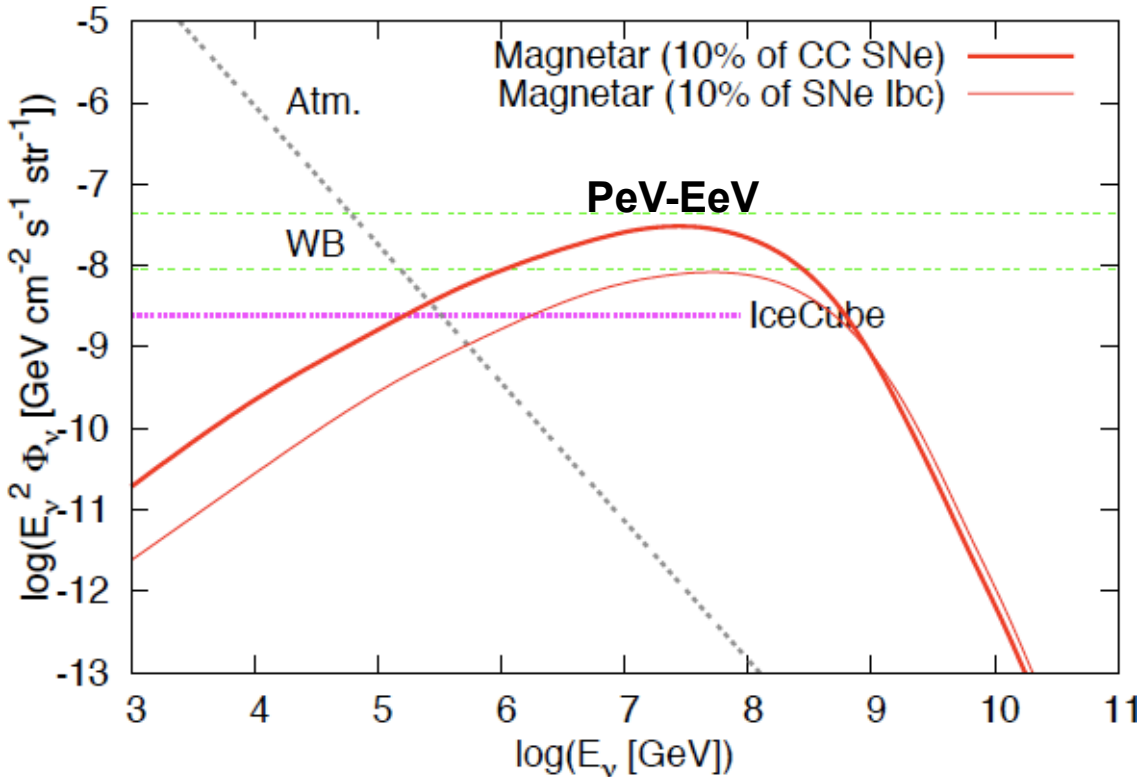
- pn collisions in LL GRBs → possible detections up to $< \sim 3$ Mpc

Possibility: Non-Shock Acceleration?



- **Magnetar-UHECR hypothesis**
UHECR acc. may occur in wind in ~hrs after the birth (Arons 03 ApJ)
- ↓
- **Accelerated CRs should interact with stellar material and rad. field**
 - pp/p γ reactions
 - **vs should be produced**
 - **Escape of UHECRs ?**
ex. puncturing envelope by jets

Neutrinos from Fast-Rotating Magnetars



KM, Meszaros, & Zhang, PRD, 79, 103001 (2009)

Time scale ~ day
soft-hard-soft time-evolution

Detectable for $D < 5$ Mpc
Probe of the magnetar birth
(birth rate ~ 0.02 - 0.05 yr $^{-1}$)

Magnetar-UHECR hypothesis can be tested by IceCube

Possibility: Post-Shock-Breakout?

Expect formation of collisionless shocks & CRs

$$\text{pp cooling: } t_{\text{pp}} = 1/(n \kappa_{\text{pp}} \sigma_{\text{pp}} c)$$

$$\text{dynamical: } t_{\text{dyn}} = l/\beta c$$

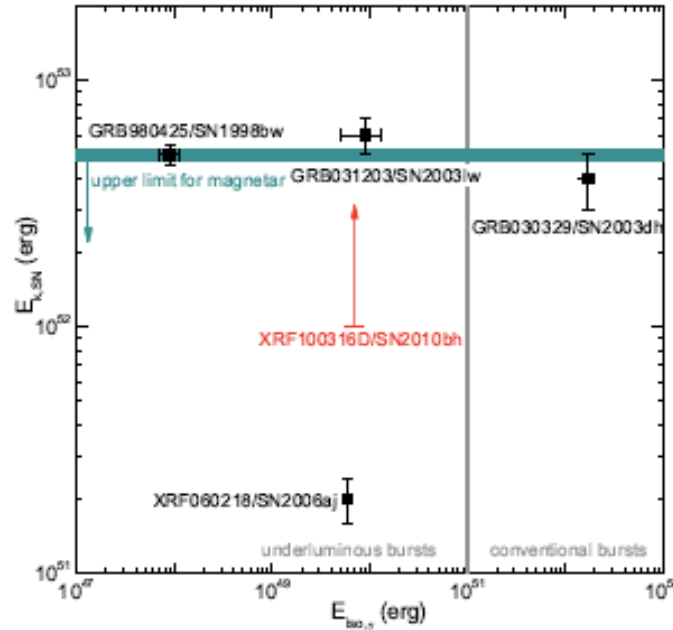
$$\rightarrow f_{\text{pp}} = (l/\beta) n \kappa_{\text{pp}} \sigma_{\text{pp}}$$

$$f_{\text{pp}}(r_{\text{bo}}) \sim \beta^{-2} (\kappa_{\text{pp}} \sigma_{\text{pp}}/\sigma_{\text{T}}) \sim 0.03 \beta^{-2}$$

$\beta \sim 1 \Leftrightarrow$ trans-relativistic SNe
($p\gamma$ efficiency ~ 1 : dominant)

$\beta \sim 0.01-0.03 \Leftrightarrow$ typical SN velocity
pp efficiency ~ 1

Trans-Relativistic SNe (Low-Luminosity GRBs)



from Fan et al. 10

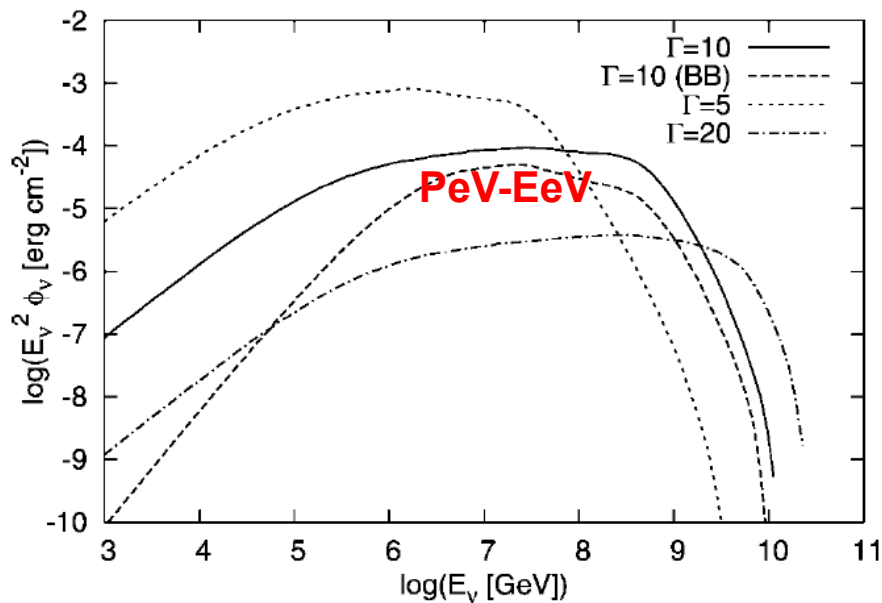
Nearby GRBs (ex. 060218@140Mpc, 980425@40Mpc) may form another class

- **much dimmer** ($E_{LL\gamma}^{\text{iso}} \sim 10^{50}$ erg $\Leftrightarrow E_{\text{GRB}\gamma}^{\text{iso}} \sim 10^{53}$ erg/s)
- **more frequent** ($\rho_{\text{LL}} \sim 10^{2-3}$ Gpc $^{-3}$ yr $^{-1}$ $\Leftrightarrow \rho_{\text{GRB}} \sim 0.05-1$ Gpc $^{-3}$ yr $^{-1}$)
- **relativistic ejecta** (the other GRB-SNe + 2009bb) (Soderberg+ 10 Nature)
- more baryon-rich? (e.g., Zhang & Yan 11 ApJ), relevant for UHECRs? (KM et al. 06 ApJ)

Two Competing Scenarios

- Inner jet dissipation (similar to GRBs)

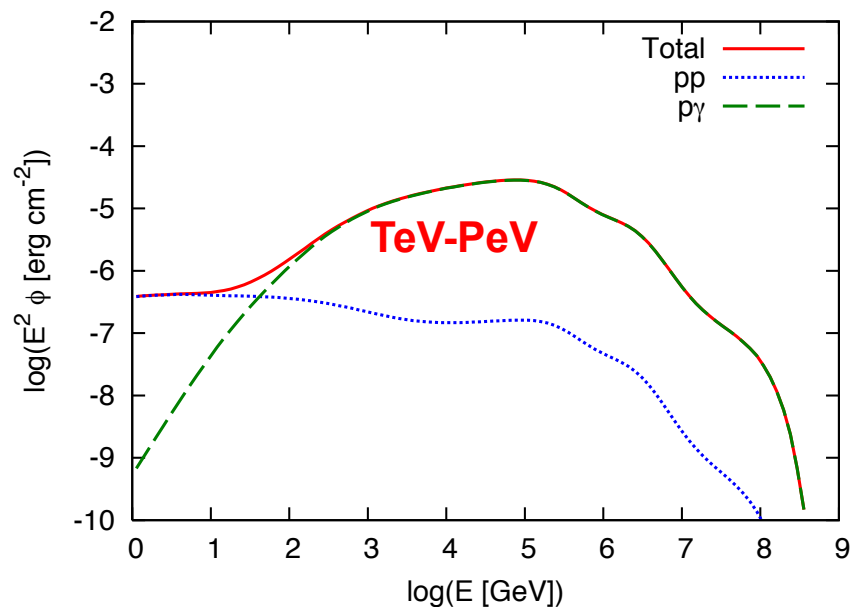
(Toma et al. 07 ApJ, Fan et al. 10 ApJL)



KM et al. 06 ApJL

- Shock breakout from optically-thick wind

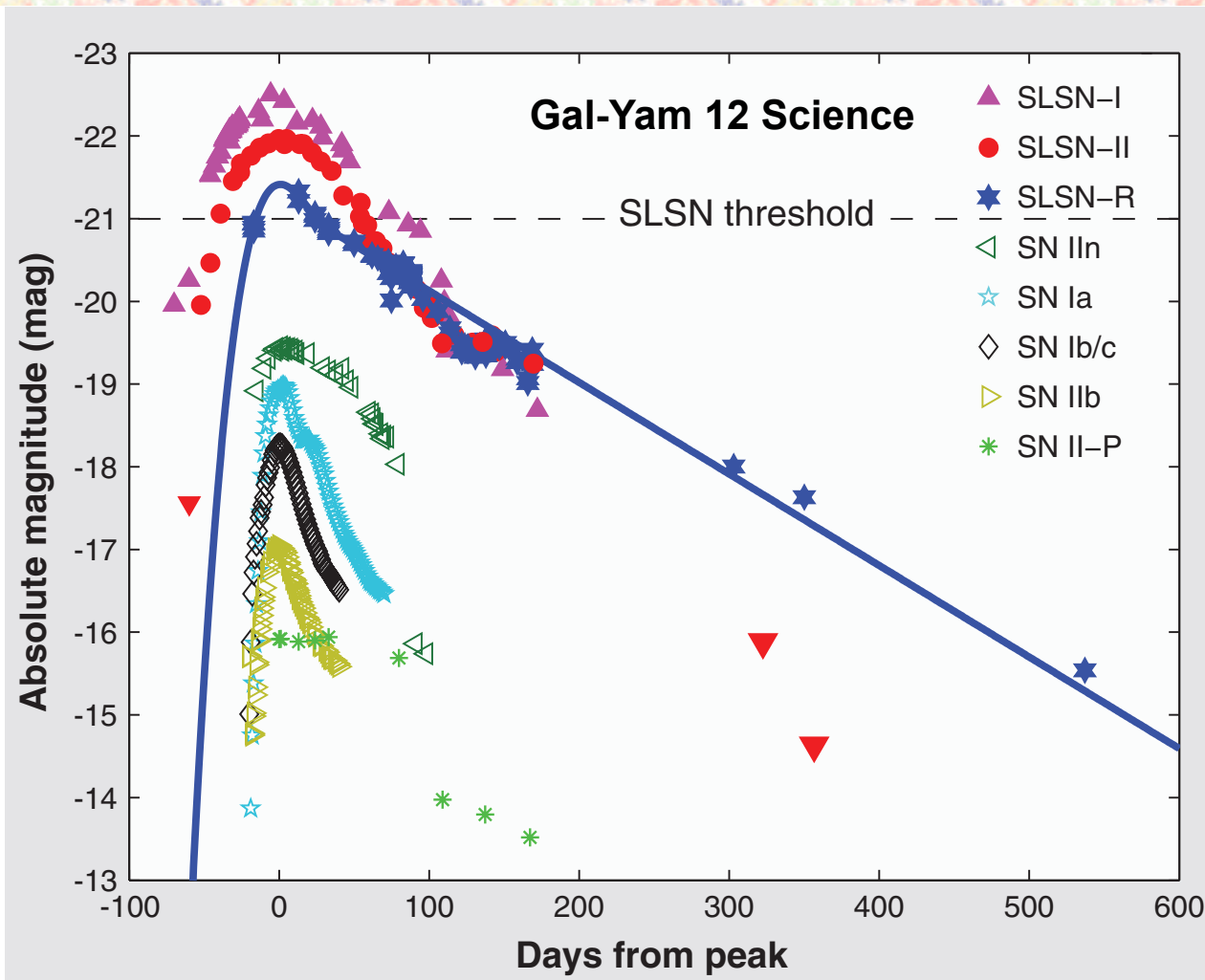
(Waxman et al. 07 ApJ, Nakar & Sari 12 ApJ)



Kashiyama et al. 12

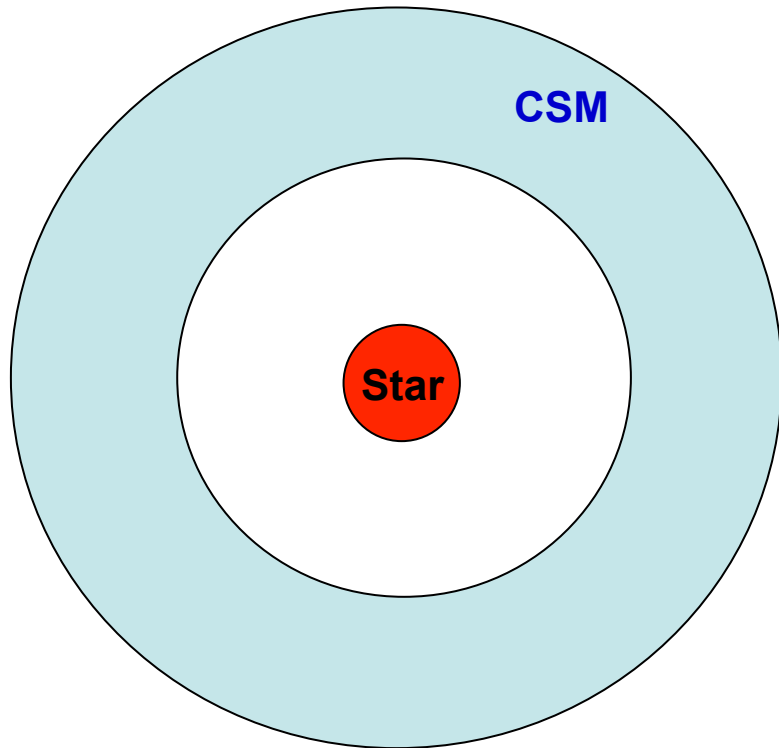
The signal is detectable for nearby SNe at $D < 10$ Mpc

SNe IIn & Super-Luminous SNe



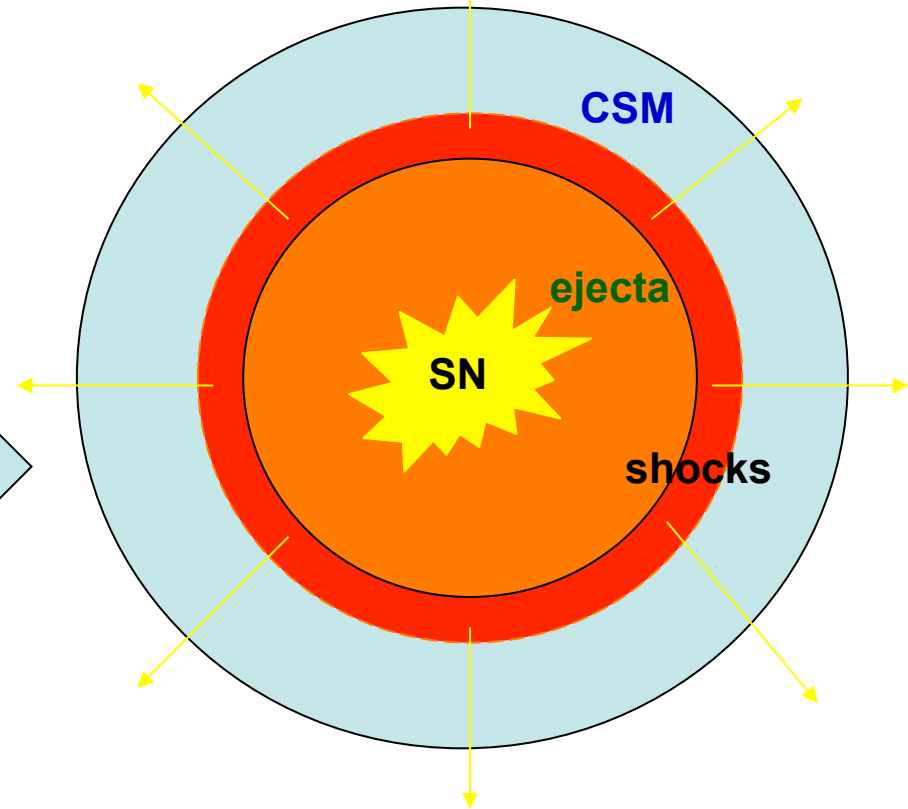
Circumstellar-Material-Collision Scenario

CSM eruptions before explosion



Similar to SNe IIn mechanism
for luminous SNe,
Smith & McCray 07 ApJL, Woosley+ 07 Nature

True SN explosion



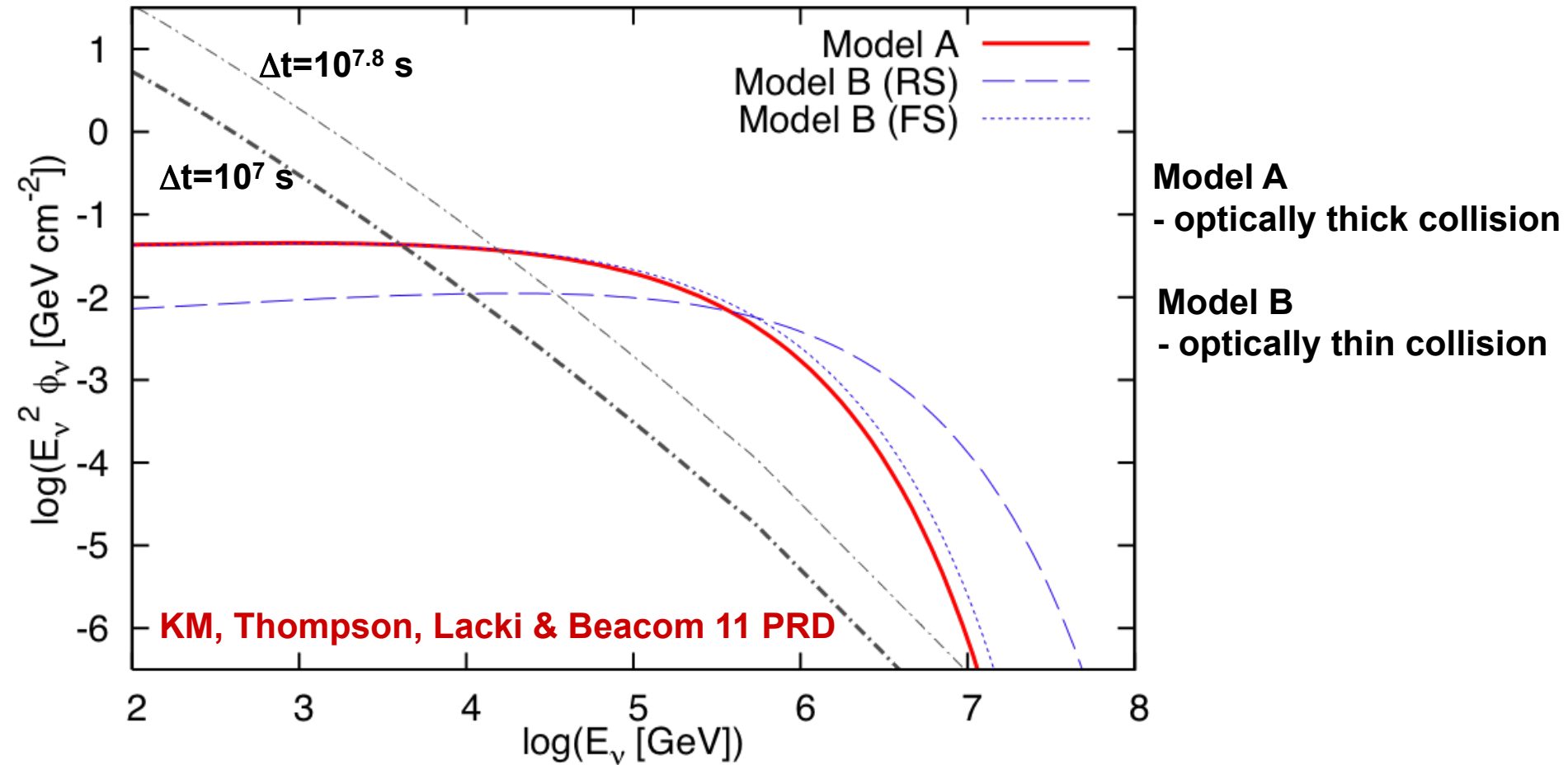
kinetic energy \rightarrow thermal + radiation

$$E_{\text{rad}} \simeq \frac{\alpha M_{\text{CSM}}}{M_{\text{ej}} + M_{\text{CSM}}} E_{\text{ej}}.$$

From SNe IIn to Luminous SNe

- $\tau_T \gg 1$ collision \rightarrow luminous SNe
strong thermalization (optical, infrared)
ex. SN 2006gy
 $R \sim 3 \times 10^{15}$ cm, $V \sim 5000$ km/s
 $n_{\text{CSM}} \sim 3 \times 10^{10}$ cm $^{-3}$ ($M_{\text{CSM}} \sim 10 M_{\text{sun}}$)
characteristic timescale: $t_{\text{bo}} \sim t_{\text{diff}} \sim t_{\text{dyn}} \sim 60$ day
- $\tau_T < 1$ collision \rightarrow SNe IIn
weaker thermalization (optical + x rays, radio)
ex. SN 2006jd
 $R \sim 3 \times 10^{16}$ cm, $V \sim 5000$ km/s
 $n_{\text{CSM}} \sim 3 \times 10^6$ cm $^{-3}$ ($M_{\text{CSM}} \sim 1 M_{\text{sun}}$)
characteristic timescale: $t_{\text{dyn}} \sim 2$ yr

Neutrinos from SNe Colliding with Massive CSM



$V_s \sim 10^{3.5} - 10^4 \text{ km/s}$

If $\varepsilon_B \sim 10^{-3} - 10^{-2} \rightarrow E_{\text{max}} \sim \text{PeV}$

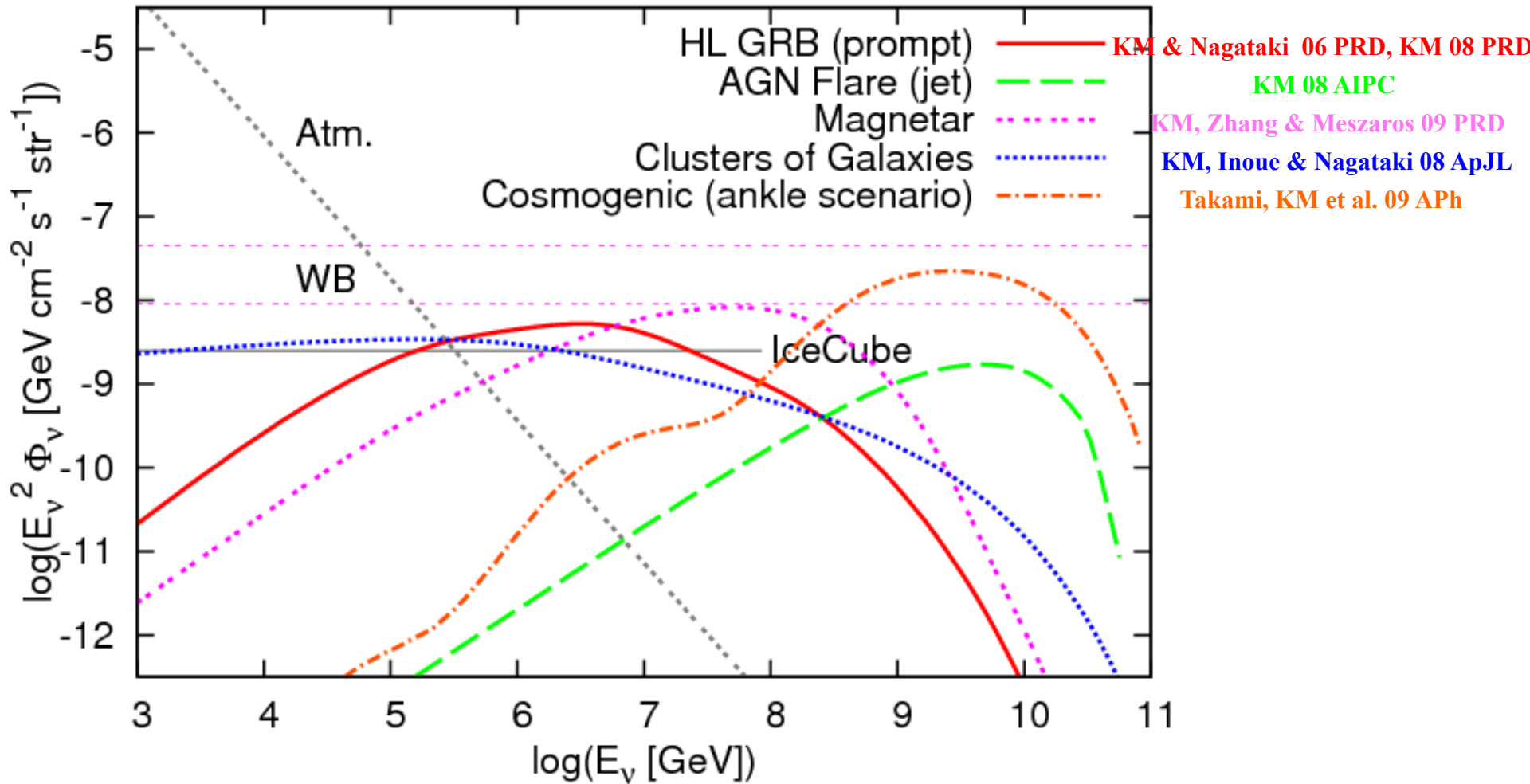
If CRs carry $\sim 10\%$ ($E_{\text{CR}} \sim 10^{50} \text{ erg}$ c.f. SNR)

\rightarrow # of $\mu\text{s} \sim$ a few for SN@10Mpc

Summary: SN Neutrinos

- Shock acc. may occur after $\tau_T \sim c/V \rightarrow$ **TeV-PeV** ν s
(ν s from slow jets are unlikely since $\tau_T \gg 1$)
- Possible non-shock acc. in magnetars \rightarrow **PeV-EeV** ν s
- Detectable typically up to **$< \sim 10$ Mpc**
- Timescales longer than GRBs \rightarrow fight w. atm. ν s...
 - trans-relativistic SNe (\sim **hr**)
 - magnetar-driven SNe (\sim **day**)
 - SN colliding with CSM (\sim **month-to-year**)
 - \rightarrow counterparts in opt./IR (+ x rays, radio and γ rays)
 - \rightarrow probes of emission mechanisms, mass loss, progenitors and GRB-SN connection

Various Predictions for Neutrino Background





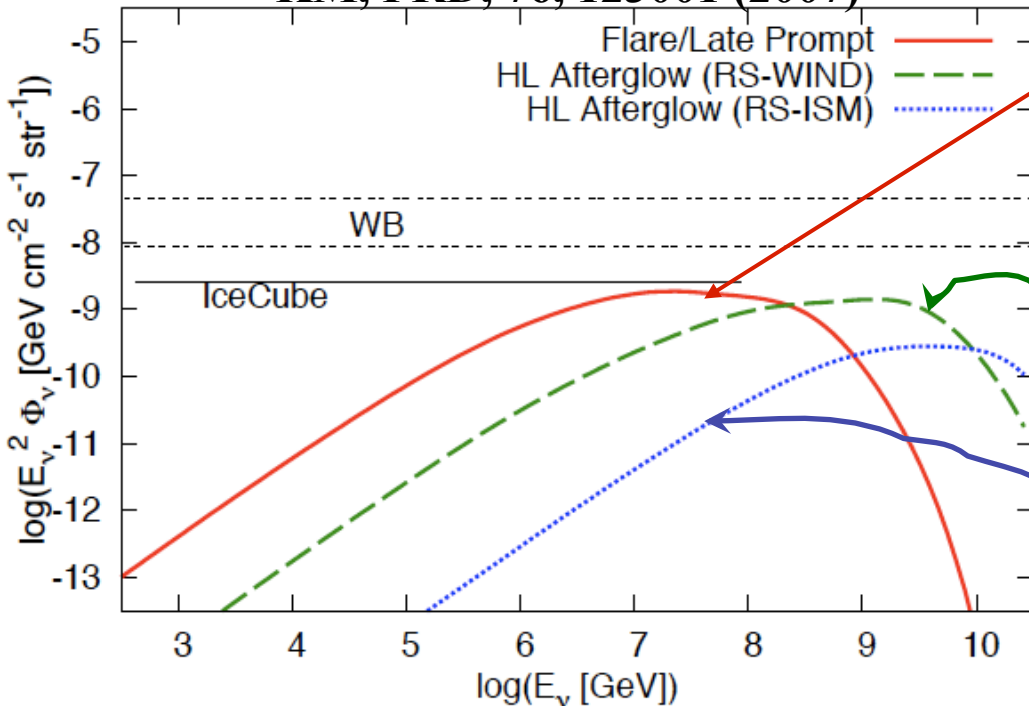
Backup Slides



GRB Early Afterglow Emission

- Most ν s are radiated in **$\sim 0.1-1$ hr** (physically $\max[T, T_{\text{dec}}]$)
- Afterglows are typically explained by **external shock scenario**
- But flares and early afterglows may come from **internal dissipation**

KM, PRD, 76, 123001 (2007)



Inner jet protons + flare x rays
(normalized by 10% of UHECR budget)

KM & Nagataki, PRL, 97, 051101 (2006)

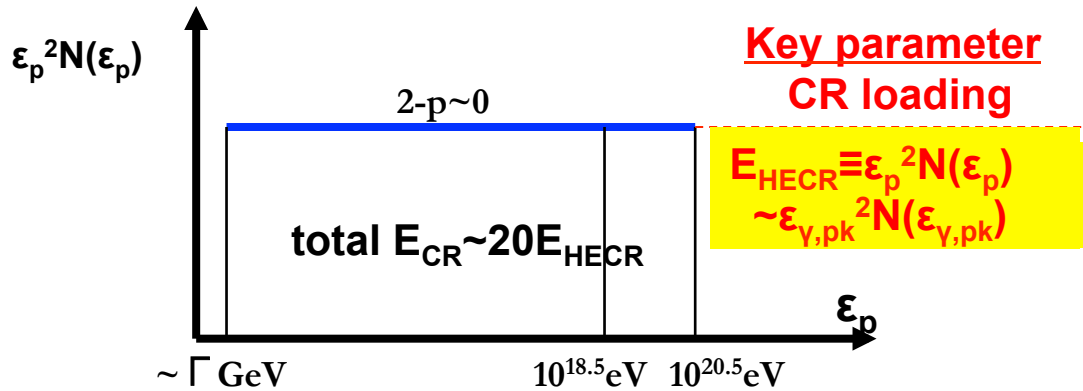
ES protons + ES opt-x rays
stellar wind medium
(normalized by UHECR budget)

ES protons + ES opt-x rays
interstellar medium
(normalized by UHECR budget)

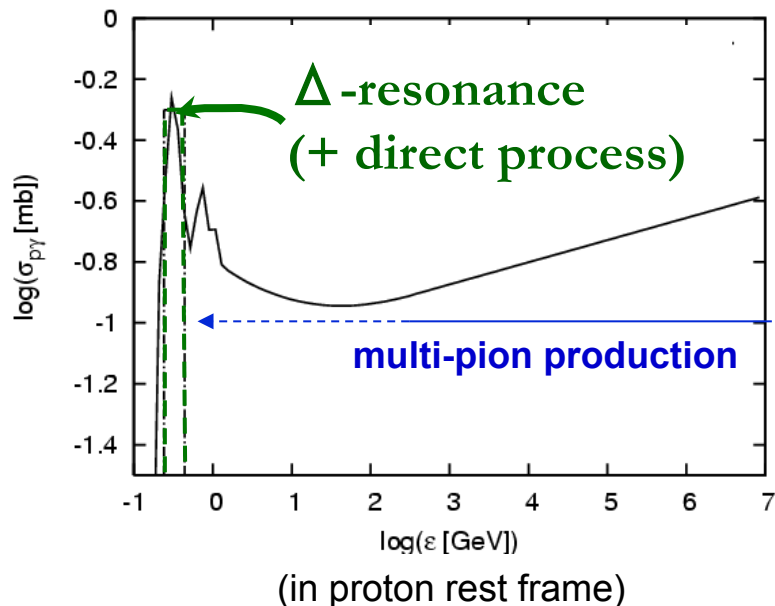
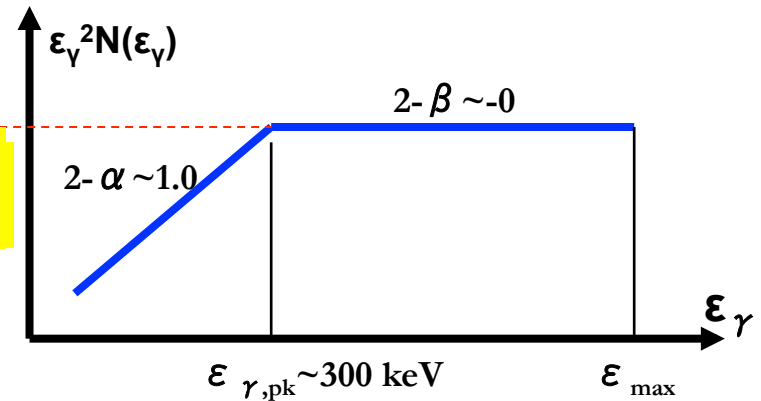
- **Flares** – efficient meson production ($f_{\text{py}} \sim 1-10$), maybe detectable
- **External shock** – not easy to detect both ν s and hadronic γ rays

Basics of Neutrino Emission

CR Spectrum (Fermi mechanism)



Photon Spectrum (observed)



Photomeson Production

at Δ -resonance

$$\epsilon_p \epsilon_\gamma \sim 0.3 \Gamma^2 \text{ GeV}^2$$

$$\epsilon_p^b \sim 0.15 \text{ GeV } m_p c^2 \Gamma^2 / \epsilon_{\gamma, pk} \sim 50 \text{ PeV}$$

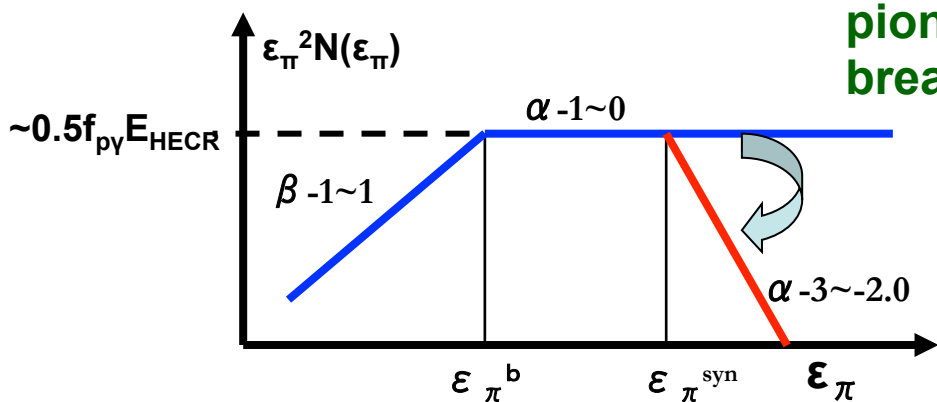
Photomeson production efficiency

~ effective optical depth for $p\gamma$ process

$$f_{p\gamma} \sim 0.2 n_\gamma \sigma_{p\gamma} (r/\Gamma)$$

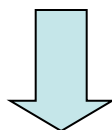
✳ multi-pion effect $\sim 3-6$ (KM & Nagataki 06 PRD)

Meson Spectrum

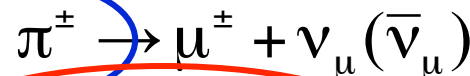
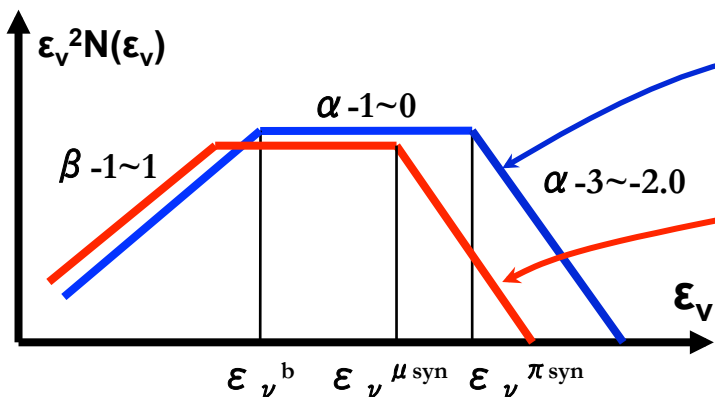


pion energy $\epsilon_\pi \sim 0.2 \epsilon_p$
 break energy $\epsilon_{\pi^b} \sim 0.07 \text{ GeV}^2 \Gamma^2 / \epsilon_{\gamma, pk} \sim 10 \text{ PeV}$

meson cooling before decay
 (meson cooling time) \sim (meson life time)
 \rightarrow break energy in neutrino spectra



Neutrino Spectrum “Waxman-Bahcall” type spectrum (Waxman & Bahcall 97 PRL)



neutrino energy $\epsilon_\nu \sim 0.25 \epsilon_\pi \sim 0.05 \epsilon_p$
 •v lower break energy $\epsilon_{\nu^b} \sim 2.5 \text{ PeV}$
 •v higher break energy $\epsilon_{\nu^{\pi^{syn}}} \sim 25 \text{ PeV}$

Neutrino oscillation

$$\nu_e : \nu_\mu : \nu_\tau = 1 : 2 : 0 \begin{cases} \rightarrow \nu_e : \nu_\mu : \nu_\tau = 1 : 1 : 1 \\ \rightarrow \nu_e : \nu_\mu : \nu_\tau = 1 : 1.8 : 1.8 \end{cases}$$

No loss

High ϵ_ν
 Loss limit

(e.g., Kashti & Waxman 05 PRL)