

# Quarkonium Production at the Tevatron

Robert Harr

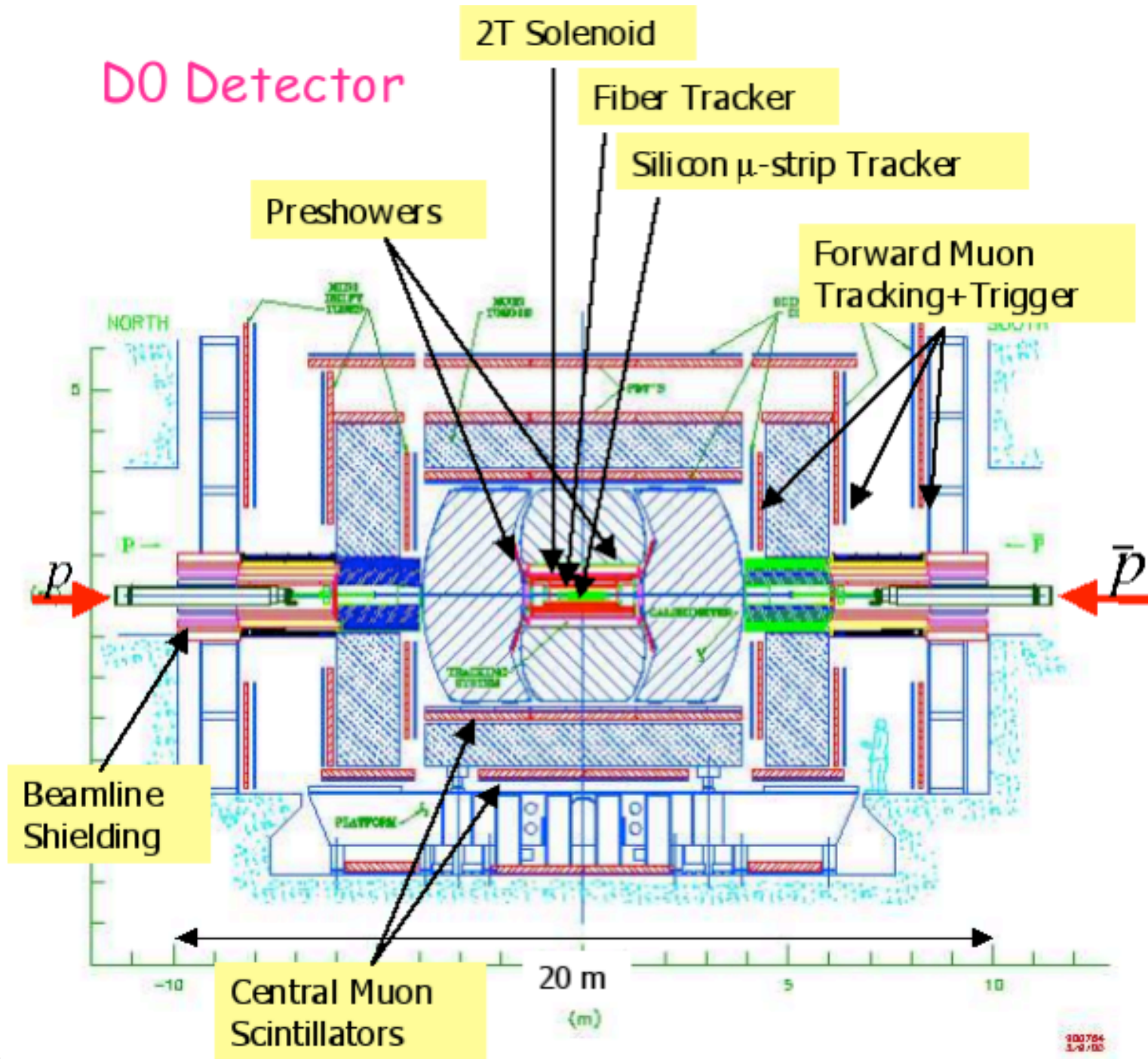
Wayne State University

Workshop on Heavy Flavor Production at Hadron  
Colliders

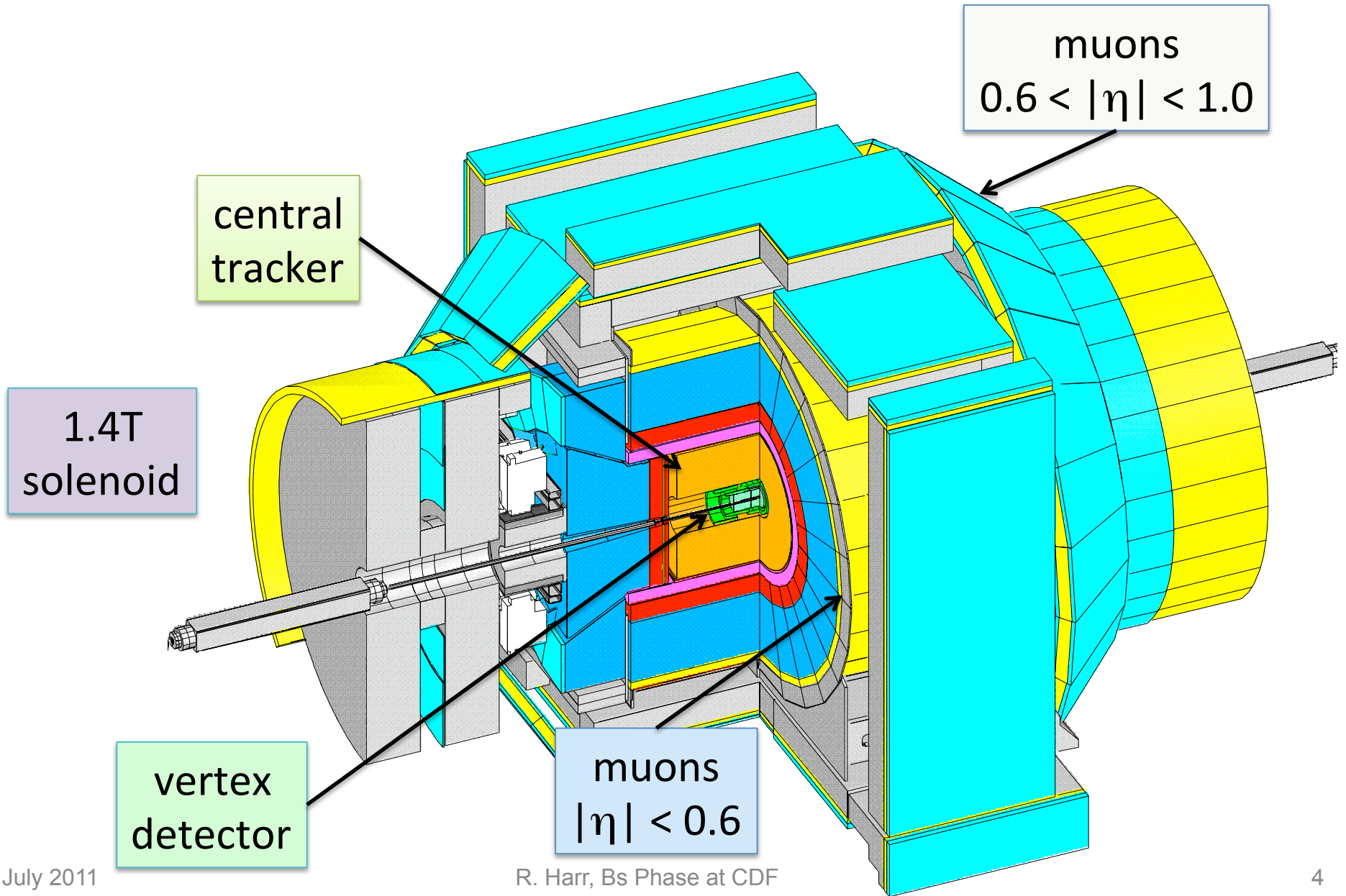


# Outline

- Brief review of issues with quarkonium production
- Upsilon polarization
- $\chi_b$  states
- Some topics omitted for lack of time:
  - $J/\psi$  and  $\psi(2S)$  cross section and polarization
  - $\chi_c$  states
  - $X(3872)$



# The CDF Detector

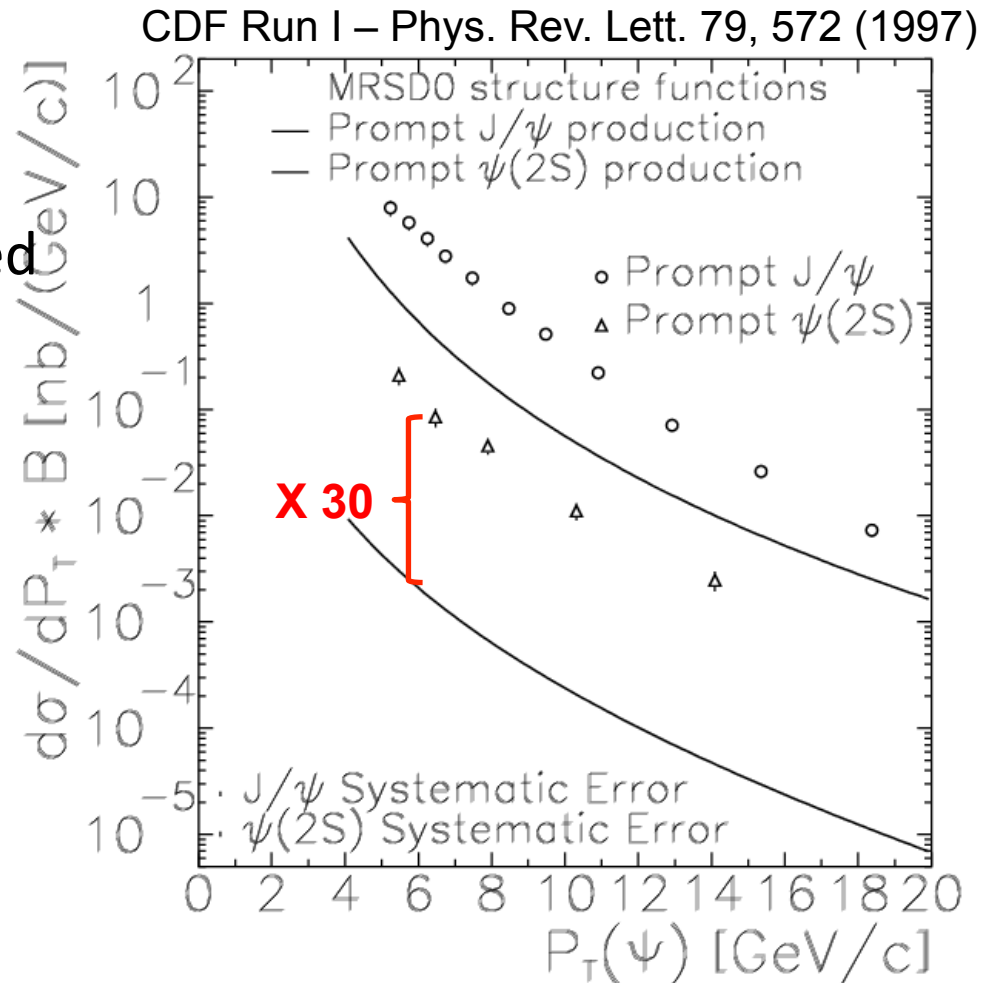


# Basic Issues

- Understanding the production cross section has been a long road.
- The quarkonium acceptance depends on the polarization at production.
- The determination of the polarization has been tricky, with unknown feed down, and inconsistent results.
- This has limited the precision of the production cross section measurements.
- Cross section alone has limited power to discriminate among theoretical models.

# CDF $J/\psi$ Cross Section

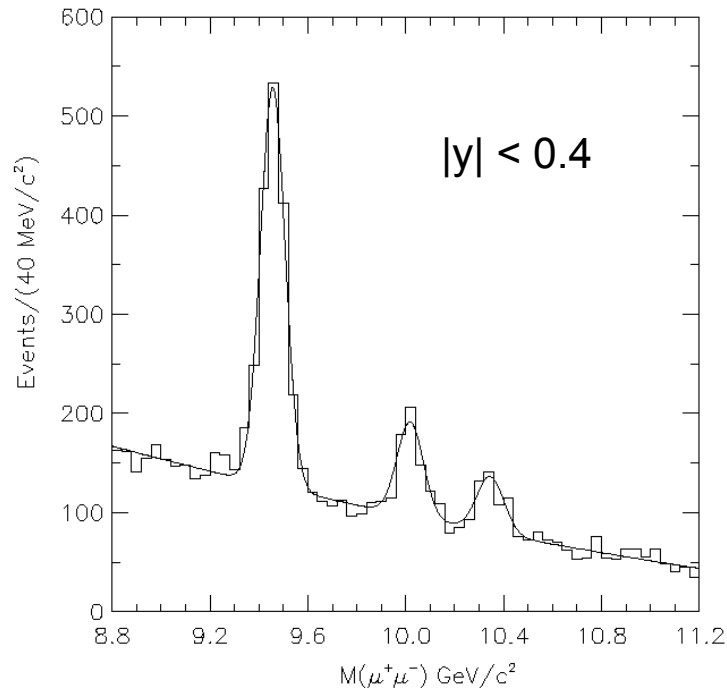
- Run I measurement
  - about  $18 \text{ pb}^{-1}$
  - non-prompt prod. estimated from displaced vertices
- $\times 30$  not explained by
  - structure functions
  - secondary production
  - feed down from  $\chi_c$
- $Y$  system is “easier”
  - no secondary component
  - calculations more reliable for heavy quarks



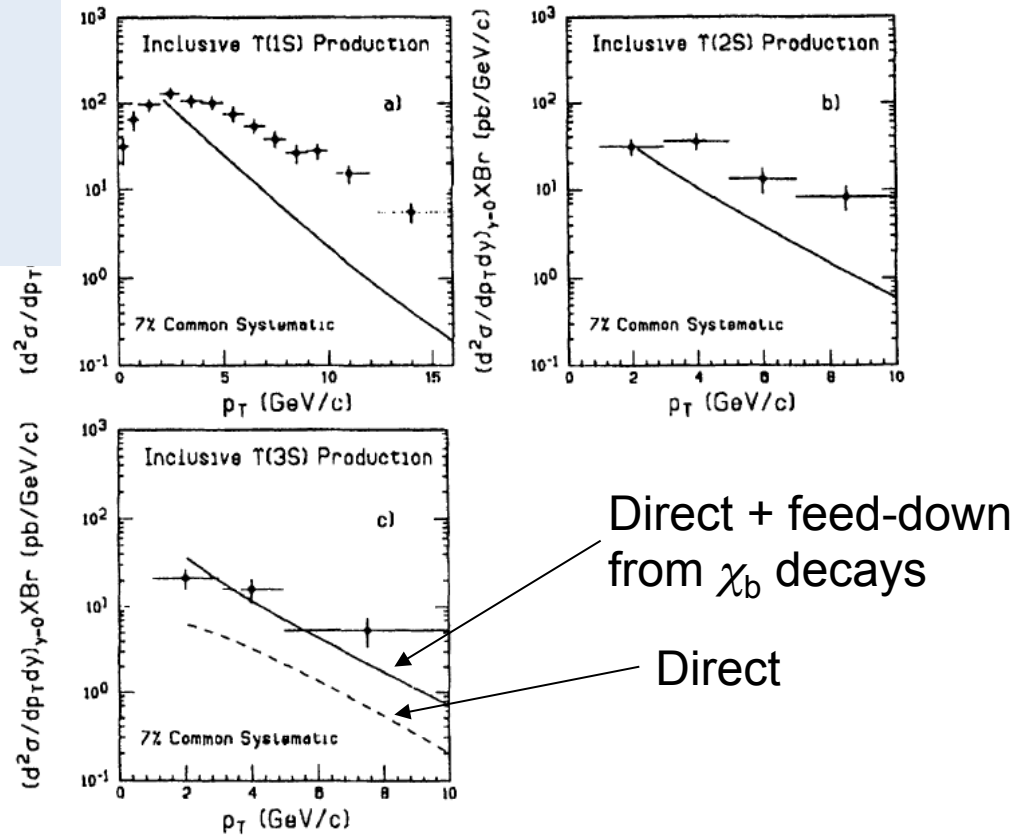
Prompted theory beyond CSM

# CDF $\Upsilon(nS)$ Cross Section

- Run I measurement
  - 17 pb<sup>-1</sup> of int. lumi.
  - no secondary production

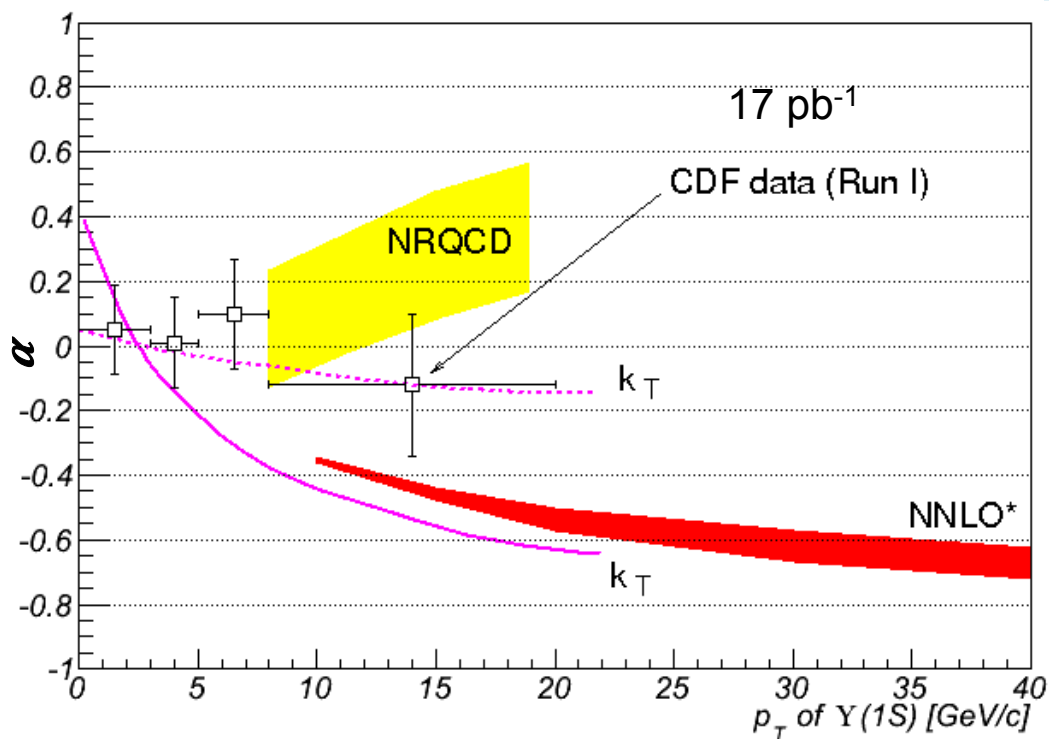


CDF Run I – Phys. Rev. Lett. **75**, 4358 (1995)



- Excess is a factor of a few.
- Still significant.

# $\Upsilon(1S)$ Polarization in Run I



CDF Run I: [Phys. Rev. Lett. 88, 161802 \(2002\)](#).  
 NRQCD: [Phys. Rev. D63, 071501\(R\) \(2001\)](#).  
 -factorization: [JETP Lett. 86, 435 \(2007\)](#).  
 NNLO\*: [Phys. Rev. Lett. 101, 152001 \(2008\)](#).

Different feed down assumptions in  $k_T$  calculations:

- $\chi_b$  decays preserve polarization
- $\chi_b$  decays destroy all polarization

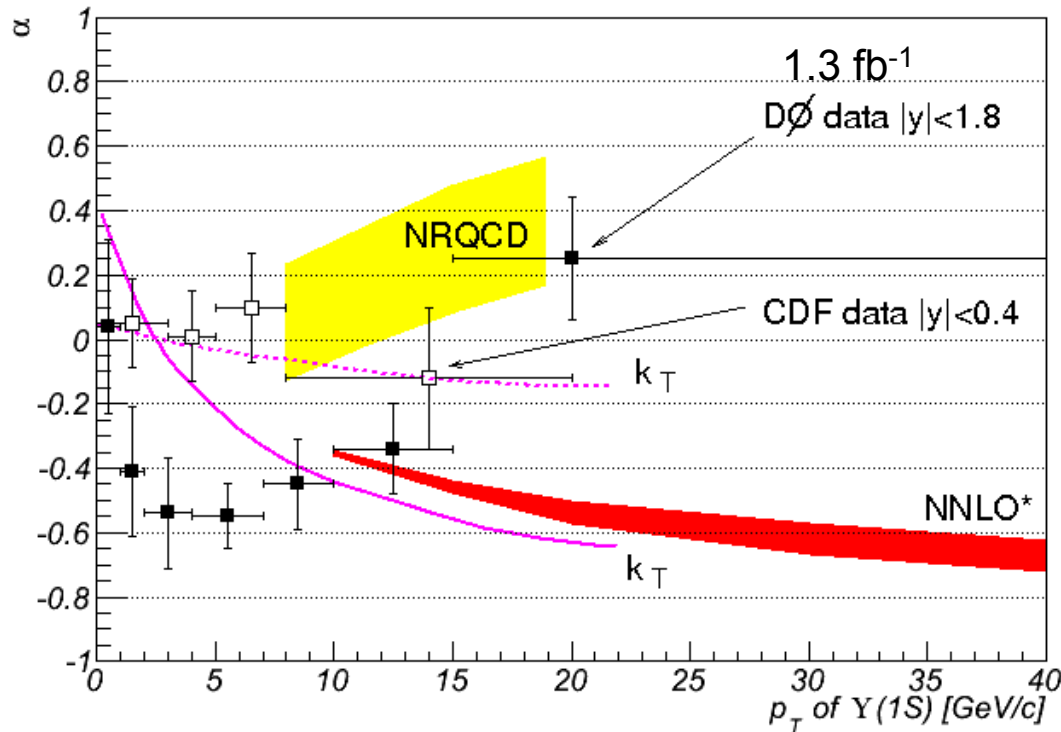
No strong polarization observed in  $\Upsilon(1S)$  decays

...

- What happens at high  $p_T$ ?
- Feed down from  $\chi_b$  states?
- Presumably less feed down for  $\Upsilon(2S)$  and  $\Upsilon(3S)$  states



# $\Upsilon$ Polarization from $D\bar{D}$ in Run II



$D\bar{D}$  Run II: [Phys. Rev. Lett. 101, 182004 \(2008\)](#).  
 CDF Run I: [Phys. Rev. Lett. 88, 161802 \(2002\)](#).  
 NRQCD: [Phys. Rev. D63, 071501\(R\) \(2001\)](#).  
 $k_T$ -factorization: [JETP Lett. 86, 435 \(2007\)](#).  
 NNLO\*: [Phys. Rev. Lett. 101, 152001 \(2008\)](#).

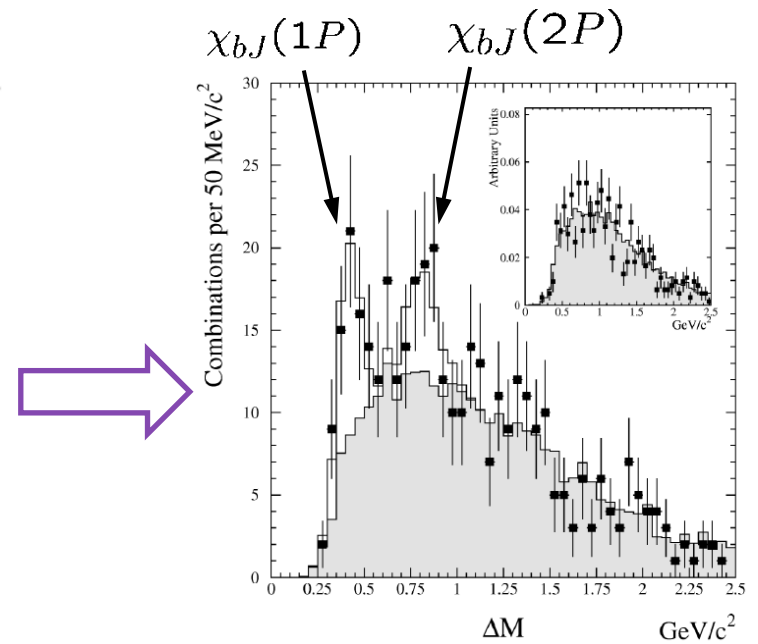
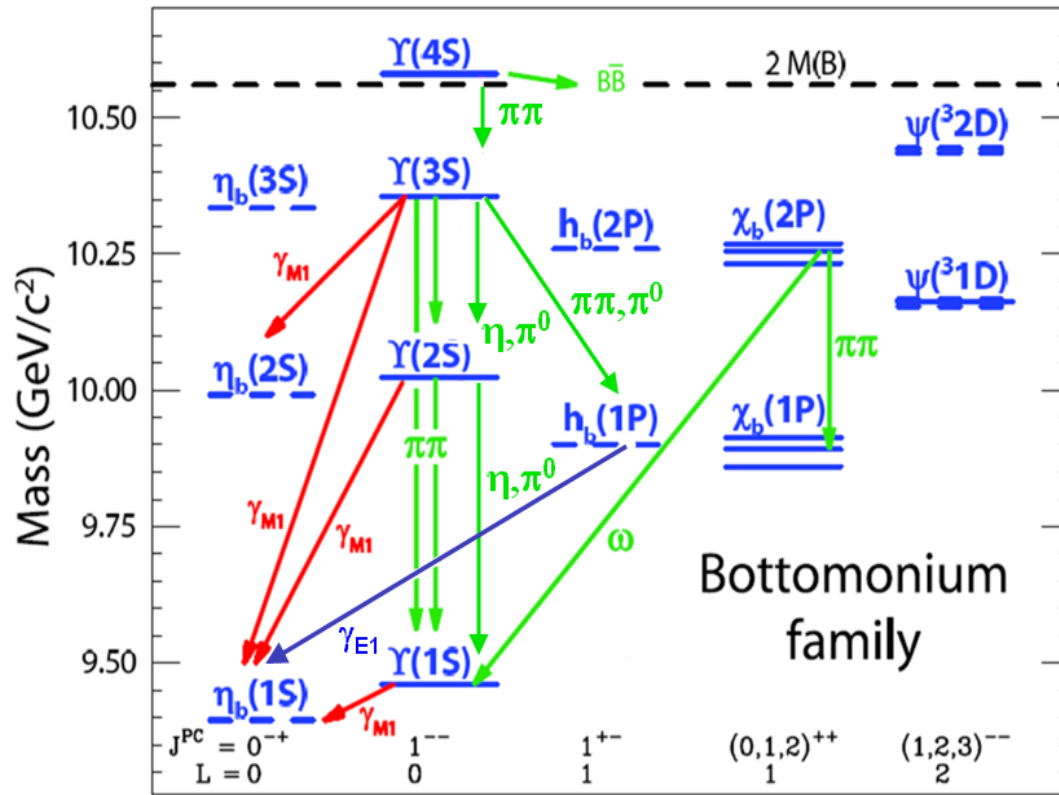
Similar analysis technique:

- Fit  $\mu^+\mu^-$  mass distribution to get  $Y$  yield in bins of  $\cos \theta$
- Correct for detector acceptance
- Fit to  $1 + \alpha \cos^2 \theta$

Results are inconsistent ...

*...why?!?*

# Feed Down from Higher States



CDF Run I - [Phys. Rev. Lett. 84, 2094 \(2000\)](#)

- P-wave states feed down to S-wave states and can influence polarization measurement.
- Need feed down fractions to account for effect.

# Suggested New Paradigm

- Angular distribution for decays to fermions:

$$\frac{d\Gamma}{d\Omega} \sim 1 + \lambda_\theta \cos^2 \theta + \lambda_\phi \sin^2 \theta \cos 2\phi + \lambda_{\theta\phi} \sin 2\theta \cos \phi + \dots$$

Faccioli *et al.* PRL 102, 151802 (2009)

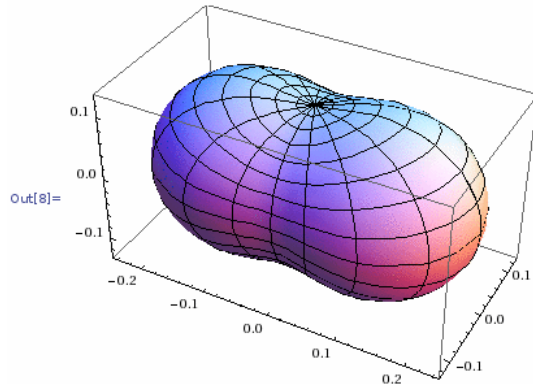
- A pure state cannot have all  $\lambda = 0$  simultaneously.
- Measured values can depend on det. acceptance
- Use of different ref. frames could help check acceptance corrections
- **We need to measure more than just  $\lambda_\theta$ !**

# Need for full polarization analysis

Transverse:  $a_0 = 0$

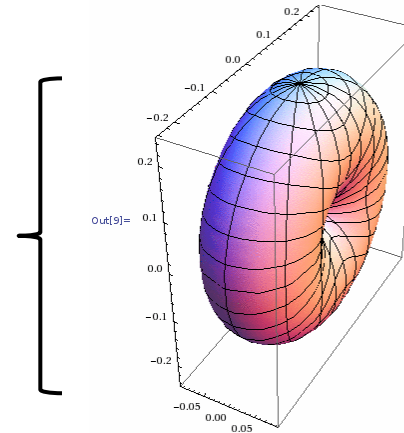
Longitudinal:  $a_{\pm 1} = 0$

```
In[8]:= SphericalPlot3D[
  g[ $\theta$ ,  $\phi$ ] /. { $\lambda_{\frac{1}{2}}^2 \rightarrow 1$ ,  $a_1 \rightarrow (1 + \text{Cos}[\pi/2]) / 2$ ,
   $a_0 \rightarrow -\text{Sin}[\pi/2] / \text{Sqrt}[2]$ ,  $a_{-1} \rightarrow (1 - \text{Cos}[\pi/2]) / 2$ },  $\theta$ ,  $\phi$ ]
```



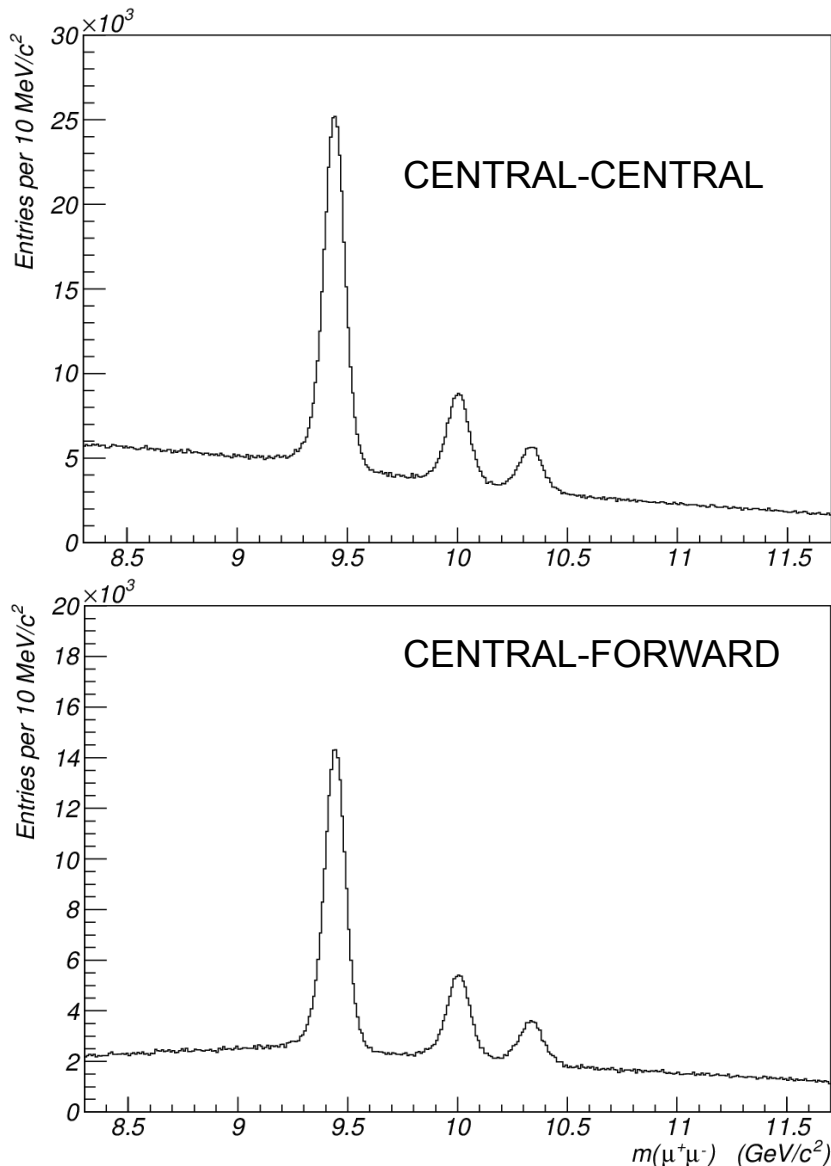
$\hat{z}$

```
In[9]:= SphericalPlot3D[
  g[ $\theta$ ,  $\phi$ ] /. { $\lambda_{\frac{1}{2}}^2 \rightarrow 1$ ,  $a_1 \rightarrow -\text{Sin}[\pi/2] / \text{Sqrt}[2]$ ,  $a_0 \rightarrow \text{Cos}[\pi/2]$ ,
   $a_{-1} \rightarrow \text{Sin}[\pi/2] / \text{Sqrt}[2]$ },  $\theta$ ,  $\phi$ ]
```



- The templates for  $dN/d\Omega$  are more complicated than simply  $1 \pm \cos^2\theta$
- Need to measure  $\lambda_\theta$ ,  $\lambda_\phi$ , and  $\lambda_{\theta\phi}$  simultaneously
- Check invariant  $\tilde{\lambda} = (\lambda_\theta + 3\lambda_\phi) / (1 - \lambda_\phi)$  w/2 frames

# The CDF Upsilon Sample



- Two triggers used:
  - central  $\mu^+\mu^-$  pair (CC)
  - central-forward  $\mu^+\mu^-$  (CF)
- Rapidity coverage:
  - C:  $|\eta(\mu)| \leq 0.6$
  - F:  $0.6 \leq |\eta(\mu)| \leq 1$
- Good signal separation
  - $\sigma_m \approx 5 \text{ MeV}/c^2$
- Yields in 6.7 fb<sup>-1</sup>:
  - $Y(1S)$  550,000
  - $Y(2S)$  150,000
  - $Y(3S)$  76,000

# Analysis Method

- Measure distribution of  $(\cos \theta, \phi)$  for all  $\mu^+\mu^-$  pairs in Upsilon signal regions
- Measure same distributions in background regions.
- Observed distribution depends on physics dist. modified by the detector acceptance:

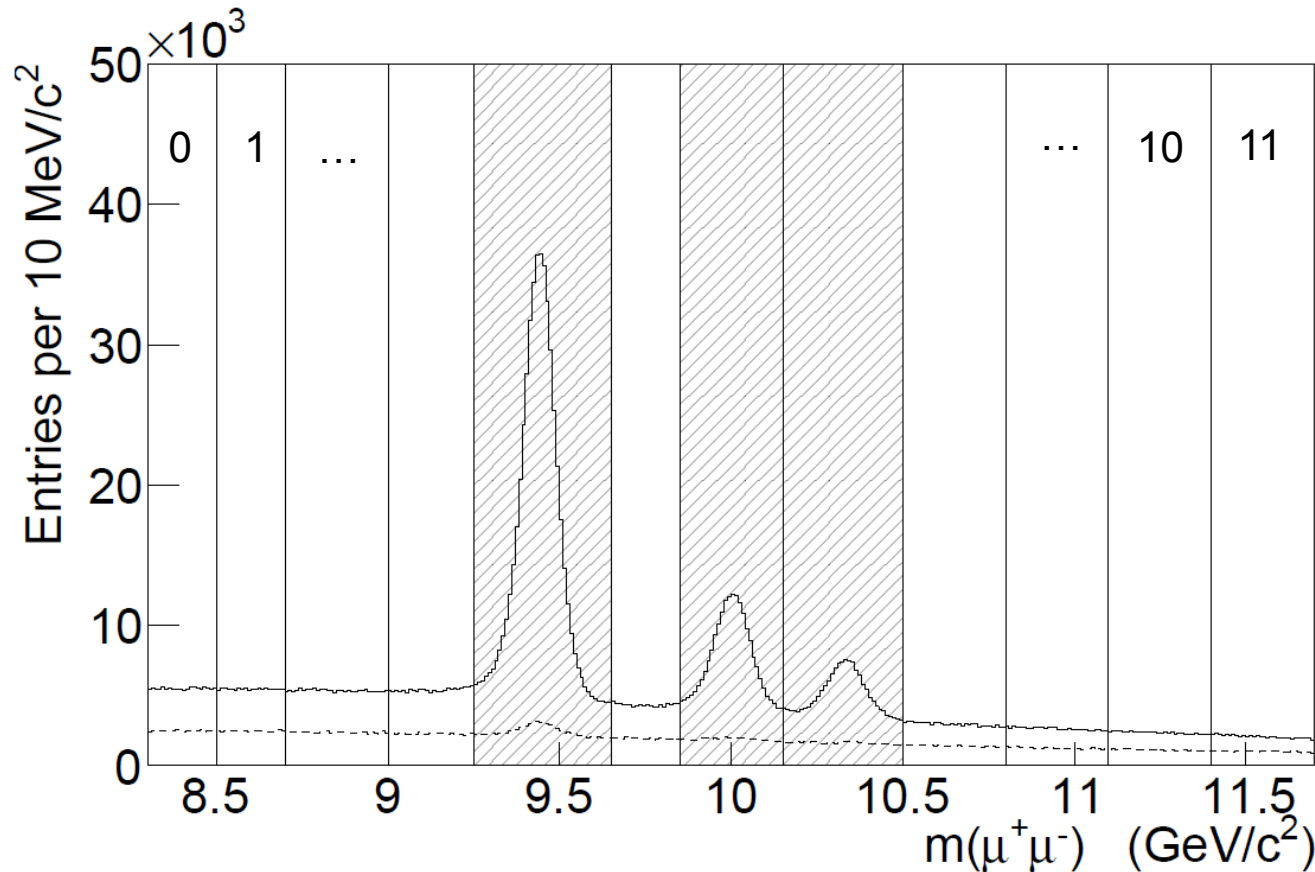
$$\frac{dN}{d\Omega} \sim f_s A_s(\cos \theta, \phi) \times w_s(\cos \theta, \phi; \lambda_s) + (1 - f_s) A_b(\cos \theta, \phi) \times w_b(\cos \theta, \phi; \lambda_b)$$

- $A(\cos \theta, \phi)$  determined from MC for sig. & bkg.

$$w(\cos \theta, \phi) \sim 1 + \lambda_\theta \cos^2 \theta + \lambda_\phi \sin^2 \theta \cos 2\phi + \lambda_{\theta\phi} \sin 2\theta \cos \phi$$

- Fit for  $\lambda_\theta$ ,  $\lambda_\phi$ , and  $\lambda_{\theta\phi}$  in both components.

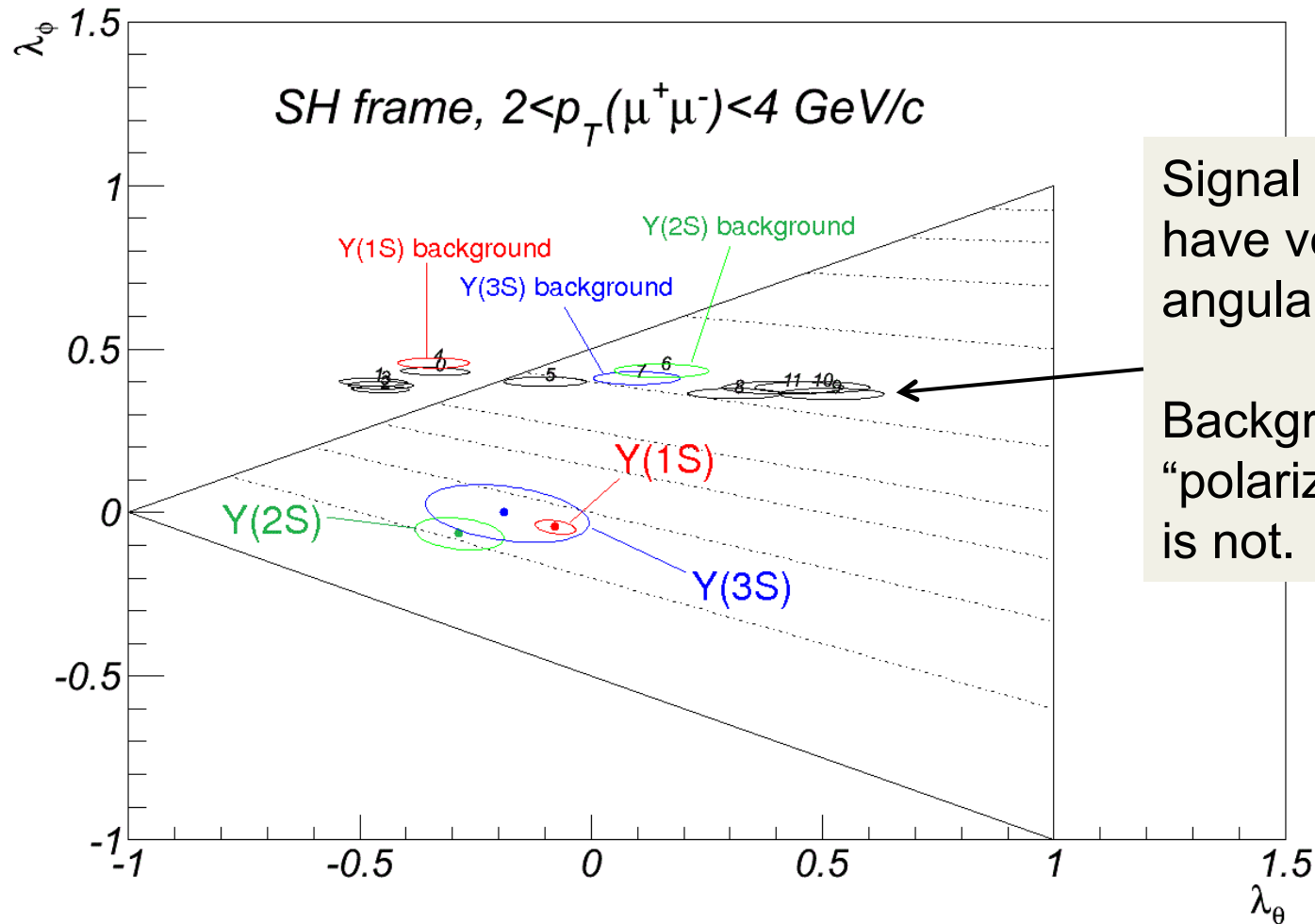
# Analysis Method



- Two components in each mass range: signal + background

$$\lambda_{\text{observed}} = f_{\text{sig}} \vec{\lambda}_{\text{sig}} + (1 - f_{\text{sig}}) \vec{\lambda}_{\text{bkg}}$$

# Fitted Parameters



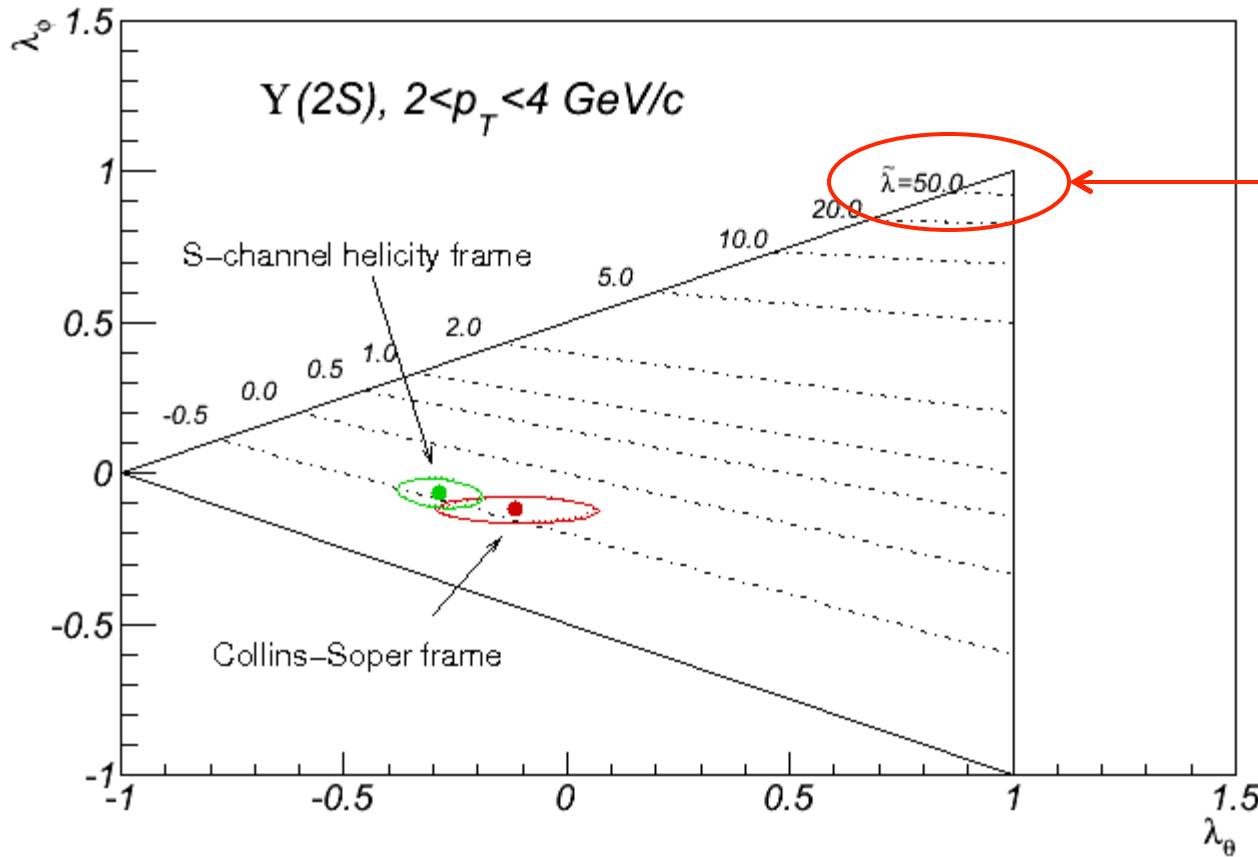
Signal and background have very different angular distributions.

Background is highly “polarized” but the signal is not.



# Consistency Tests

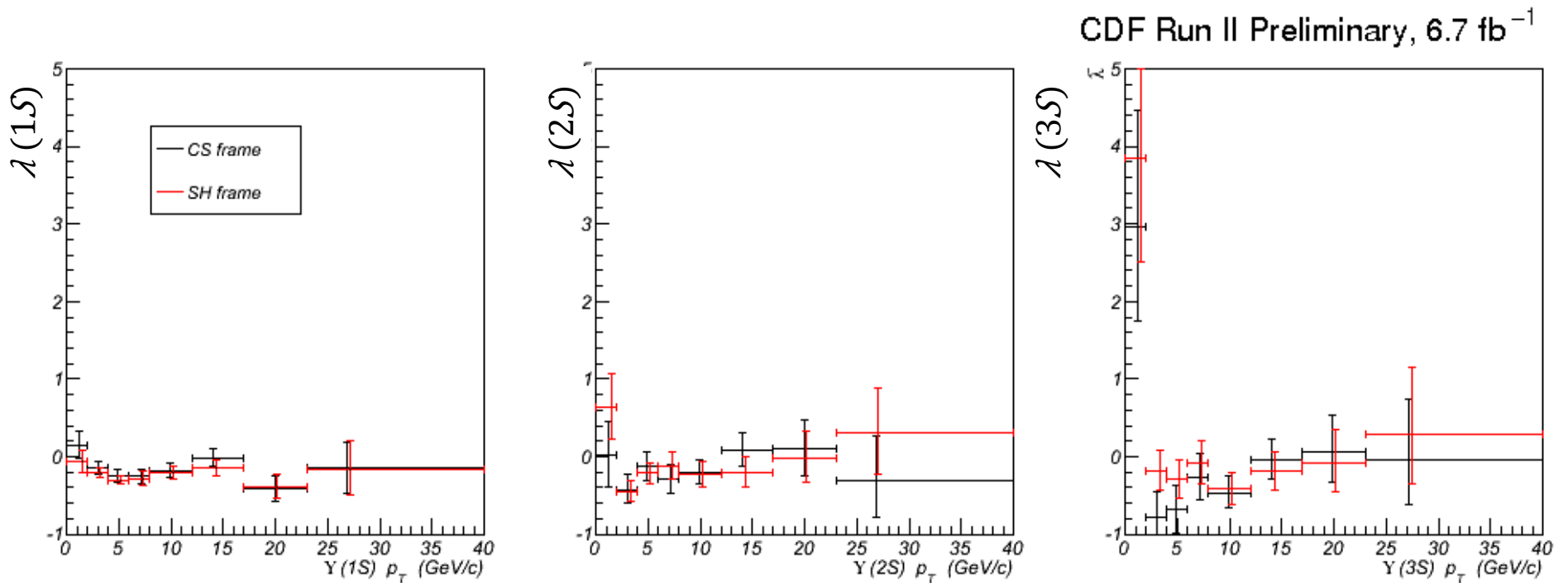
CDF Run II Preliminary,  $6.7 \text{ fb}^{-1}$



It can be shown that the expression  $\bar{\lambda} = (\lambda_\theta + 3\lambda_\phi) / (1 - \lambda_\phi)$  is the same in all reference frames.

*We observe that indeed it is.*

# Frame Invariance Tests

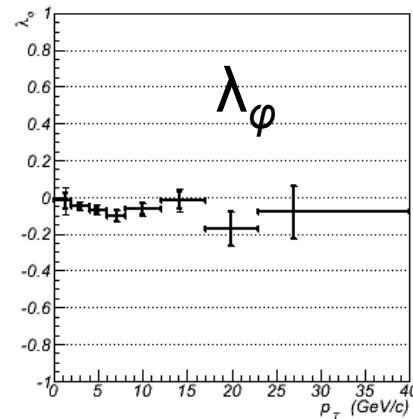
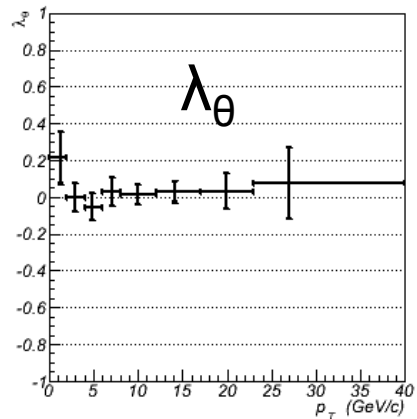


- Differences generally consistent with fluctuation predicted from simulation
- Difference used to quantify systematic uncertainties on  $\lambda_\theta$ ,  $\lambda_\phi$ , and  $\lambda_{\theta\phi}$

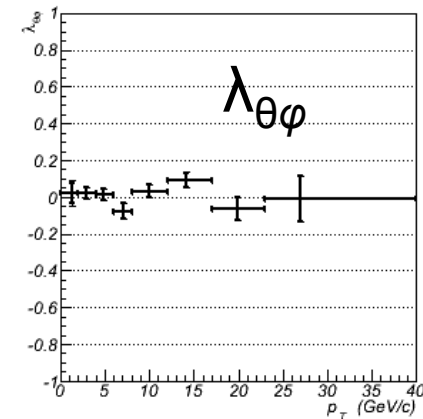
# Results for $\Upsilon(1S)$ state

CS

$\Upsilon(1S)$  - Collins-Soper frame

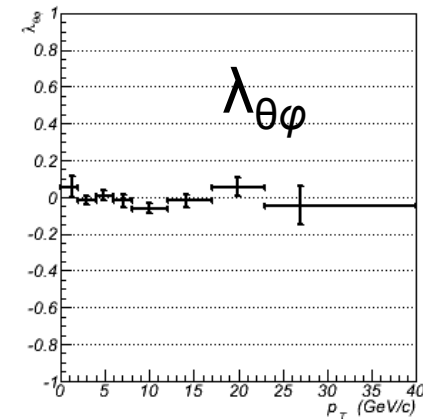
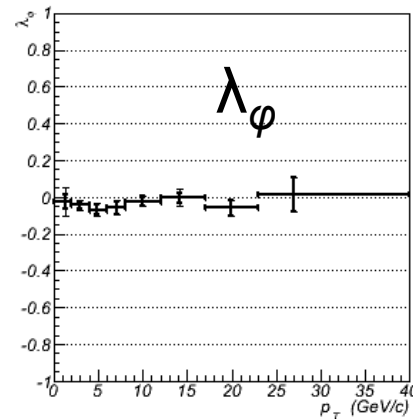
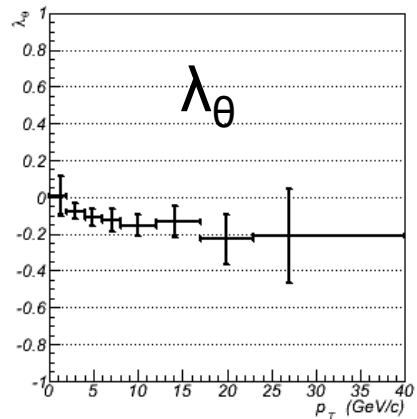


CDF Run II Preliminary,  $6.7 \text{ fb}^{-1}$



S-H

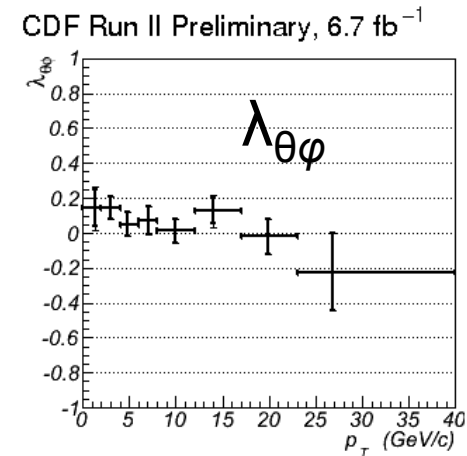
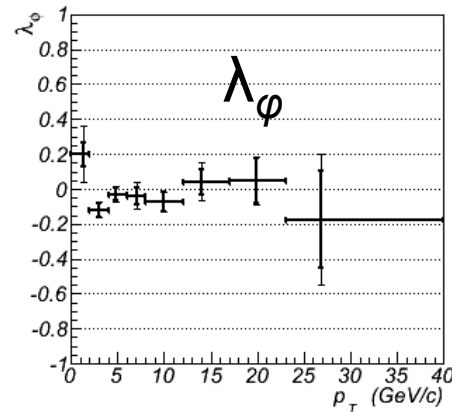
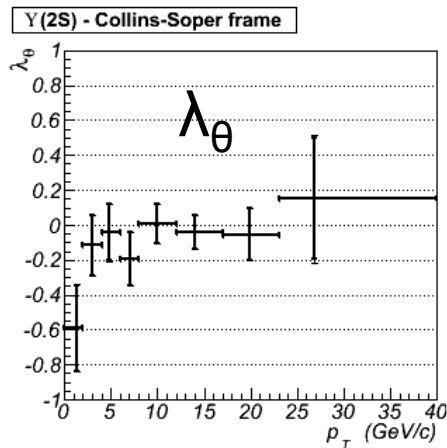
$\Upsilon(1S)$  - S-channel helicity frame



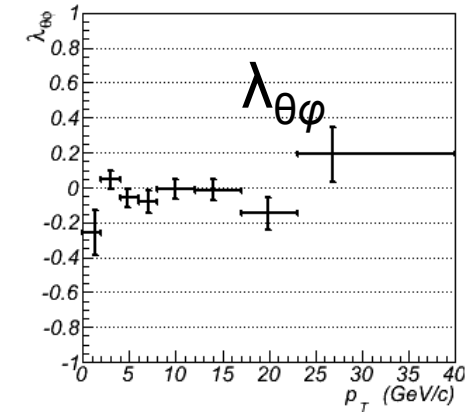
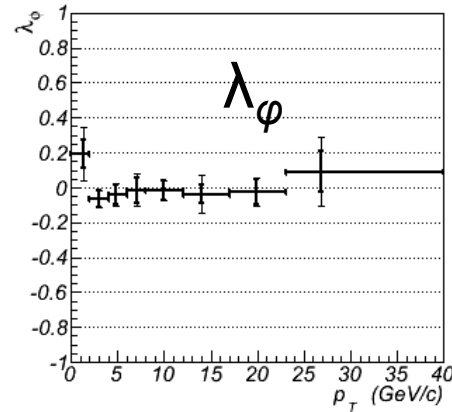
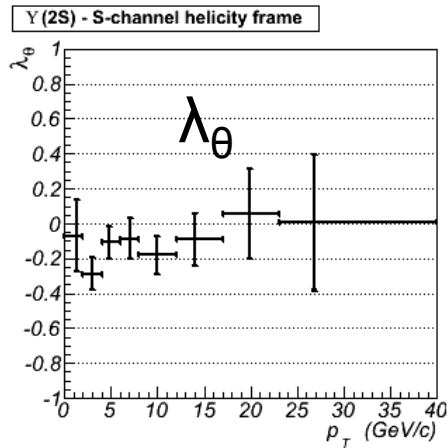
- Nearly isotropic... what about the  $\Upsilon(2S)$  and  $\Upsilon(3S)$  states?

# Results for $\Upsilon(2S)$ state

CS



S-H

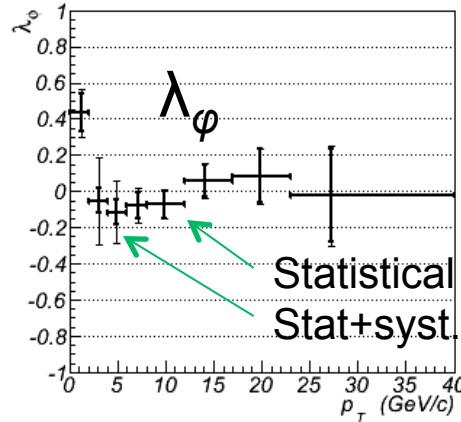
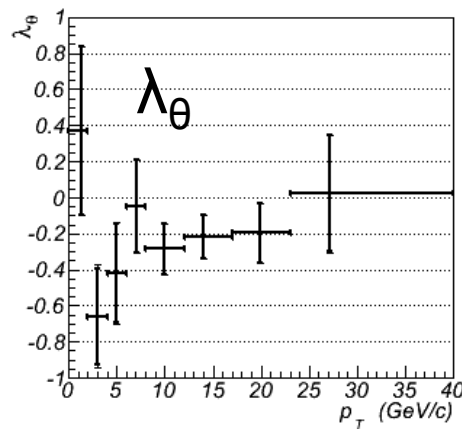


- Looks isotropic, even at large values of  $p_T$  ...

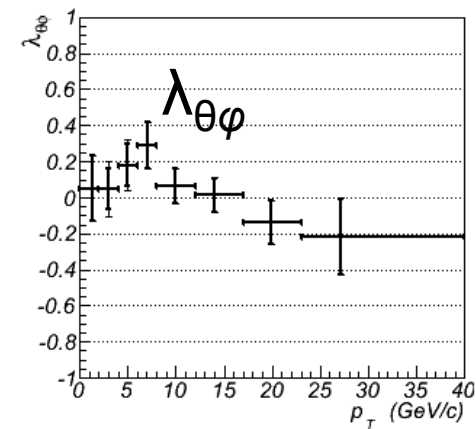
# First measurement of $\Upsilon(3S)$ spin alignment

CS

$\Upsilon(3S)$  - Collins-Soper frame

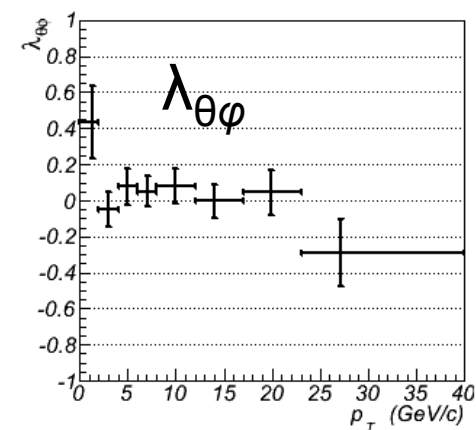
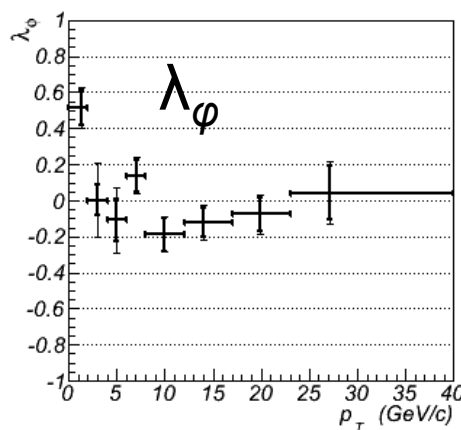
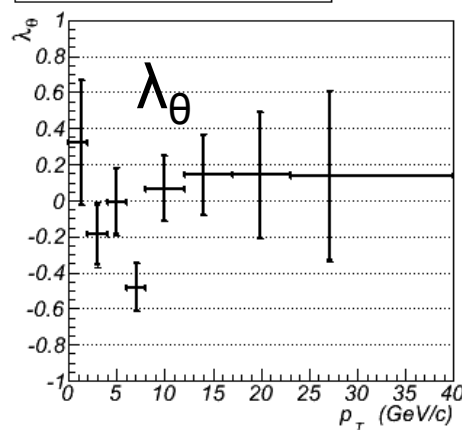


CDF Run II Preliminary,  $6.7 \text{ fb}^{-1}$



S-H

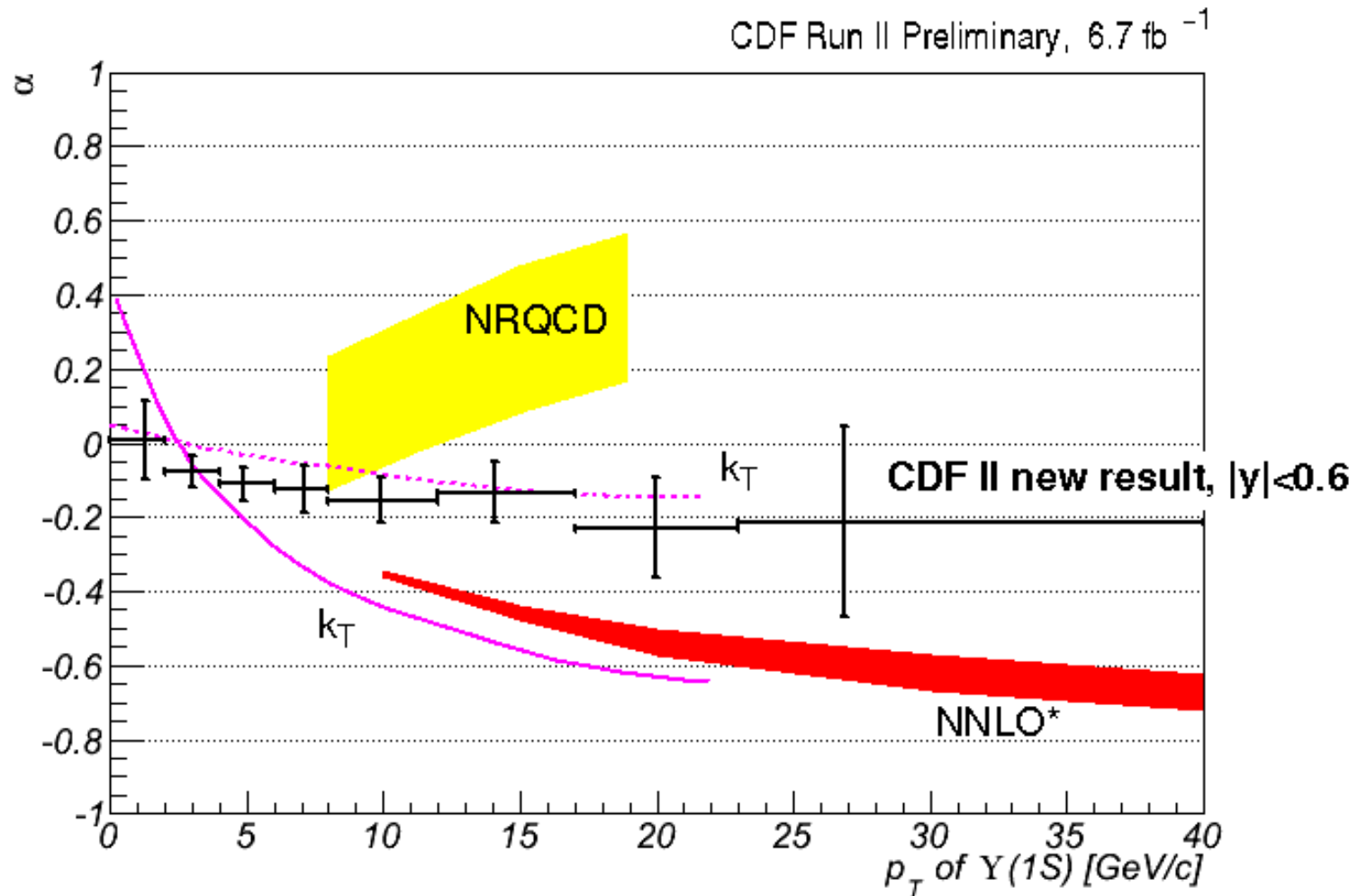
$\Upsilon(3S)$  - S-channel helicity frame



- No evidence for significant polarization.

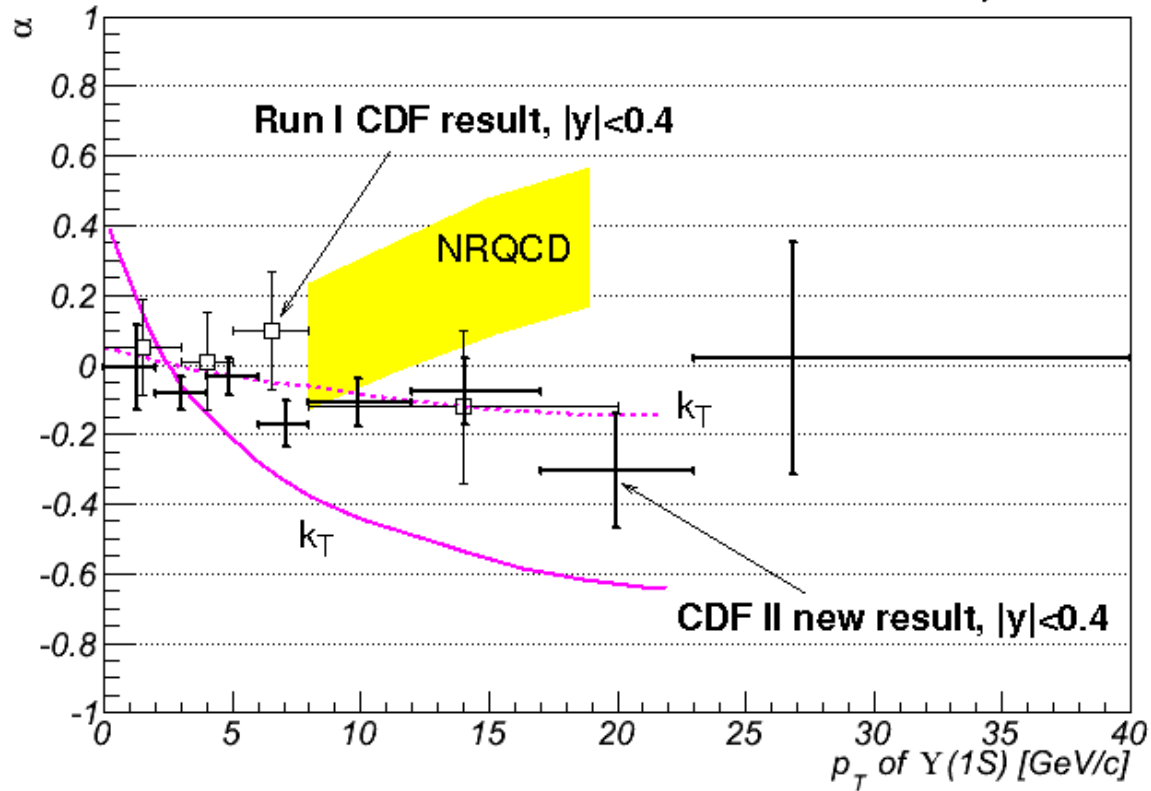
# Comparison with Models

- Old predictions for  $\lambda_\theta$  in the S-channel helicity frame



# Comparison with previous results

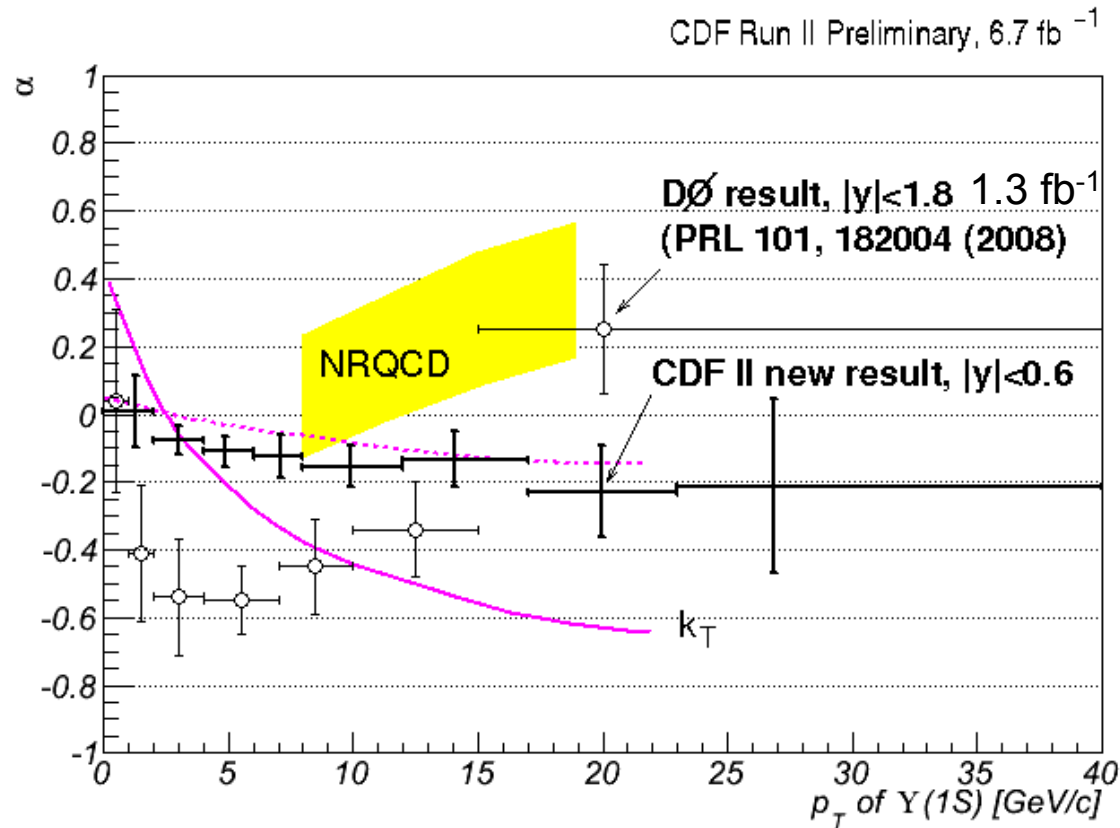
CDF Run II Preliminary,  $6.7 \text{ fb}^{-1}$



NRQCD – Braaten & Lee, Phys. Rev. D63, 071501(R) (2001)  
 $k_T$  – Baranov & Zotov, JETP Lett. 86, 435 (2007)

Agrees with previous CDF publication from Run I

# Comparison with previous results



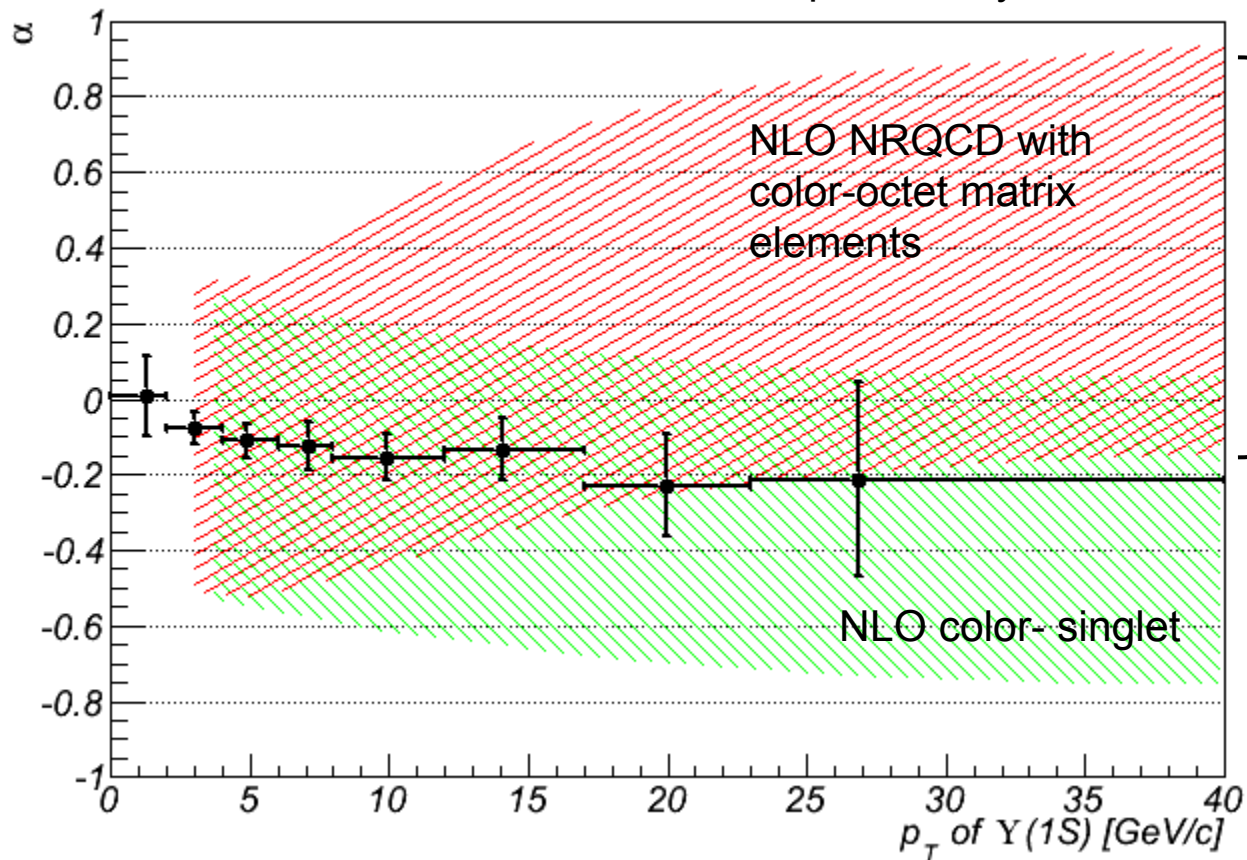
NRQCD – Braaten & Lee, Phys. Rev. D63, 071501(R) (2001)  
 $k_T$  – Baranov & Zotov, JETP Lett. 86, 435 (2007)

- Does not agree with result from DØ at about the  $4.5\sigma$  level



# Comparisons with newer calculations

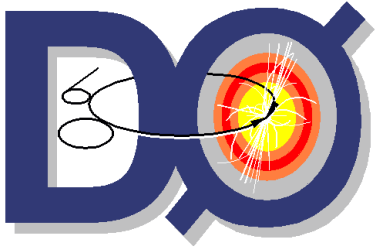
CDF Run II preliminary – 6.7 fb<sup>-1</sup>



Significant uncertainty due to feed-down from states (conservative assumptions)

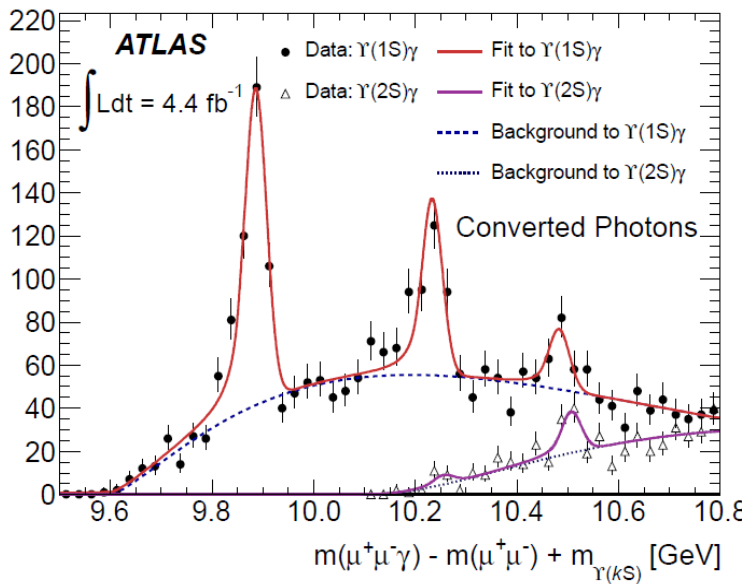
Nucl. Phys. B 214, 3 (2011) summary:

- NLO NRQCD – Gong, Wang & Zhang, Phys. Rev. D83, 114021 (2011)
- Color-singlet NLO and NNLO\* - Artoisenet, *et al.* Phys. Rev. Lett. 101, 152001 (2008)

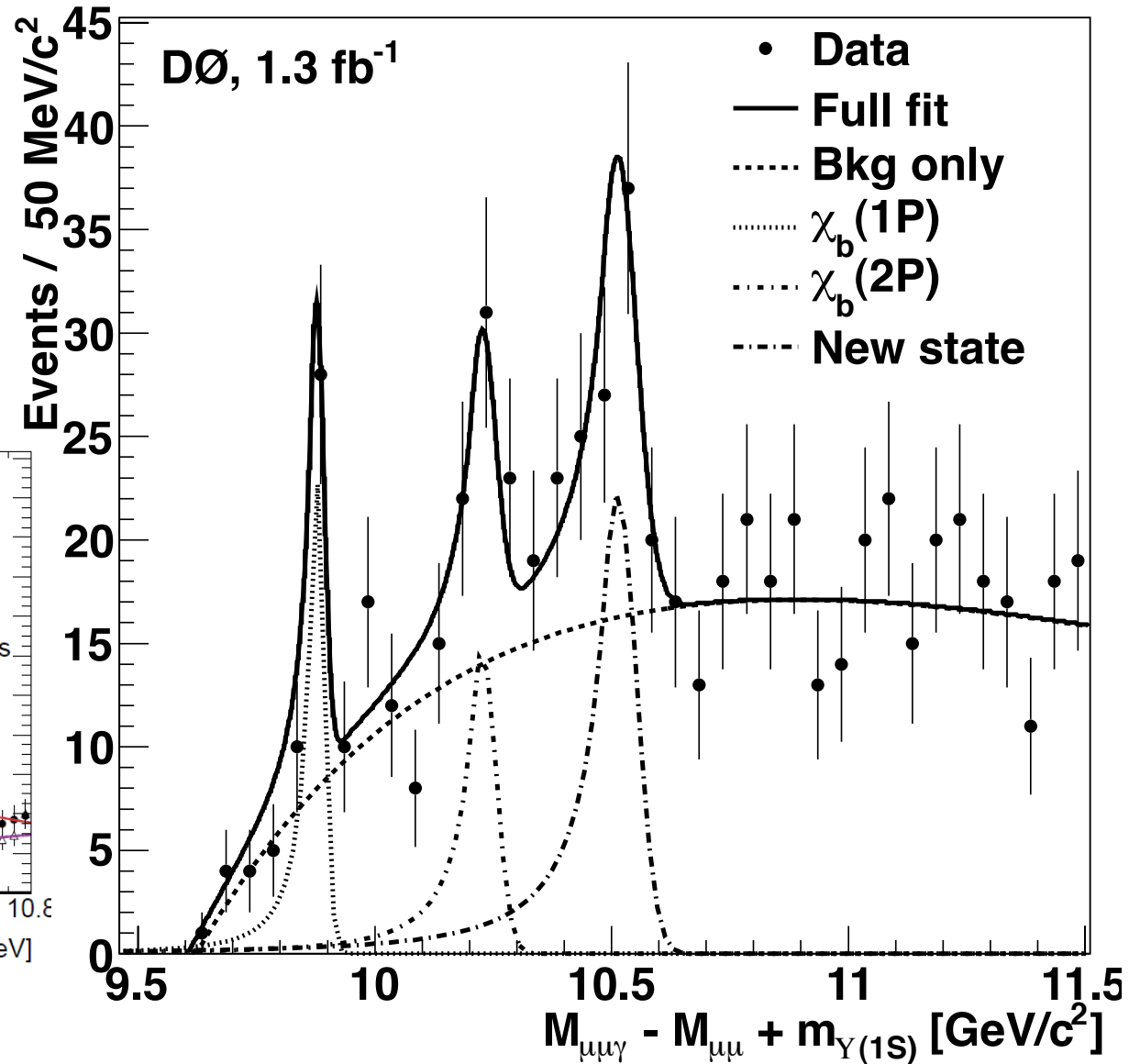


# $\chi_b$ States in $Y(1S) + \gamma$

Consistent with ATLAS observation



ATLAS – [arXiv:1112.5154](https://arxiv.org/abs/1112.5154)



# Summary (1)

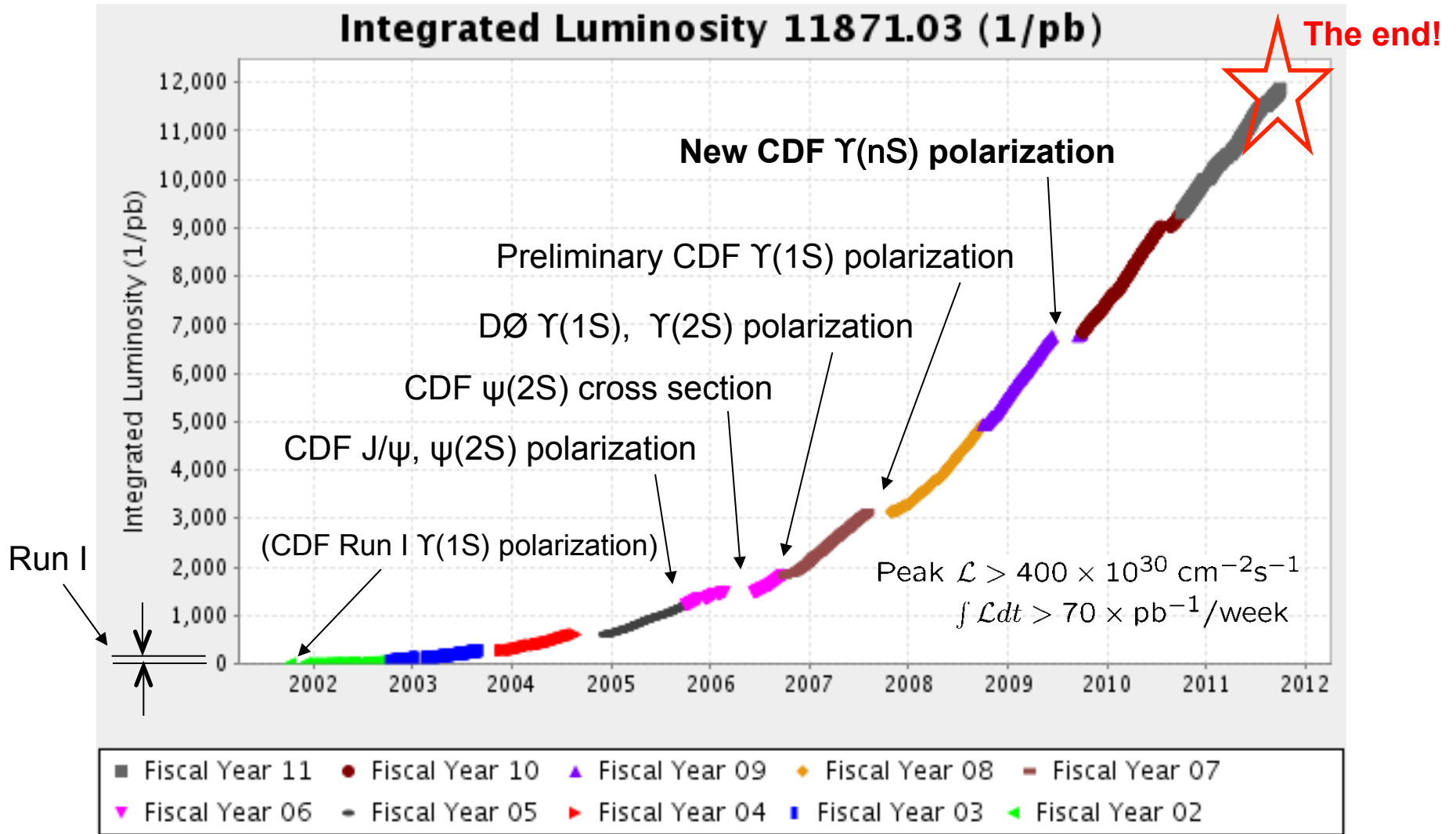
- Polarizations and feed-down provide important tests of production models.
- CDF measurement of Upsilon polarization
  - for 1S, 2S, and 3S states,
  - in Collins-Soper and helicity frames, and
  - using full 3-D measurement,
- Result indicates that production is isotropic, for all  $p_T$  and even for  $Y(3S)$ .
- D0 measurements of  $\chi_b(3P)$  state help establish it as a possible source of feed-down.

# Summary (2)

- More effort is needed to
  - complete measurements on bottomonium system
    - better understand feed-down fractions
    - measure the Upsilon cross section
  - and extend the techniques to the charmonium system
    - handle additional non-prompt production

# Additional Material

# Tevatron Run II

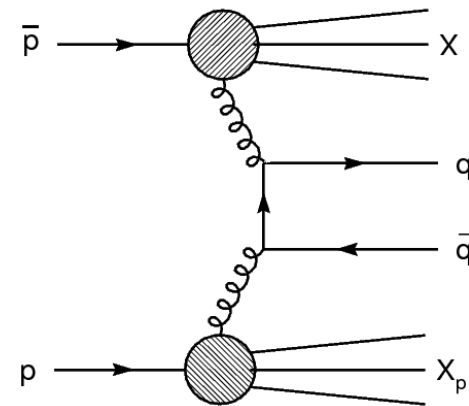


# Another Model: “ $k_T$ factorization”

$$\sigma_{pp} = \int G(x_1, \mu^2) G(x_2, \mu^2) \sigma_{gg}(x_1, x_2) dx_1 dx_2$$

$$G(x, \mu^2) \rightarrow \mathcal{F}_g(x, k_T^2, \mu^2)$$

“un-integrated gluon densities”



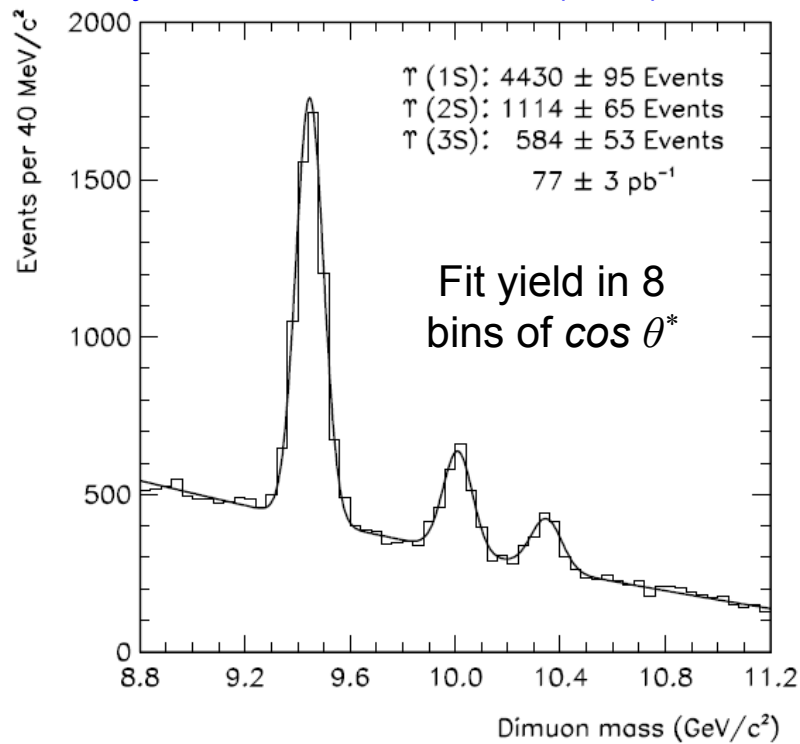
$$\overline{\epsilon_g^\mu \epsilon_g^{*\nu}} = k_T^\mu k_T^\nu / |k_T|^2$$

$\Rightarrow$  Initial state gluon polarization related to  $k_T$

- No need for color-octet terms...
- Predicted **longitudinal**  $\Upsilon$  polarization for  $p_{\perp T} \gg m_{\perp Q}$

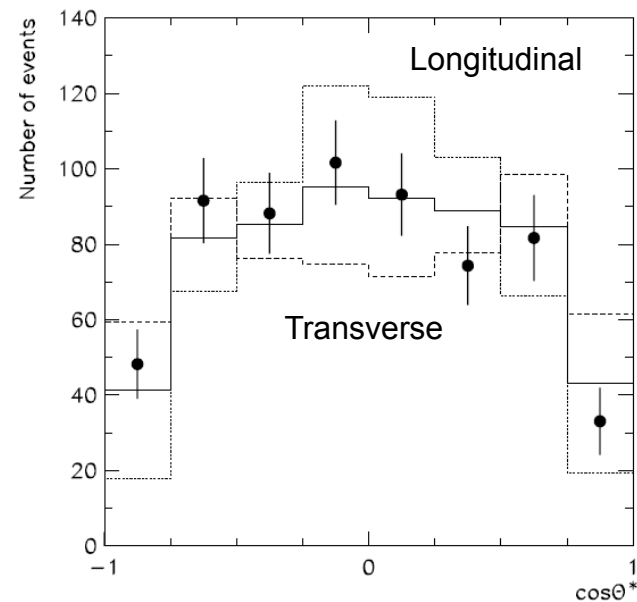
# CDF Measurement

[Phys. Rev. Lett. 88, 161802 \(2002\).](#)



Transverse:  $1 + \cos^2 \theta^*$   
 Longitudinal:  $1 - \cos^2 \theta^*$

Template distributions for transverse/longitudinal polarization strongly influenced by detector acceptance.



- Observed distribution is **isotropic** - neither longitudinal nor transverse.



# $\Upsilon$ Polarization from $D\bar{D}$ in Run II

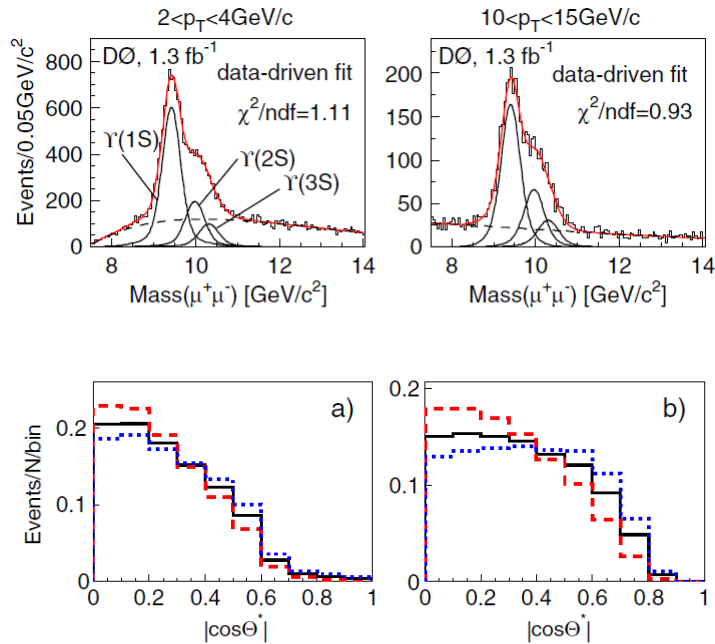
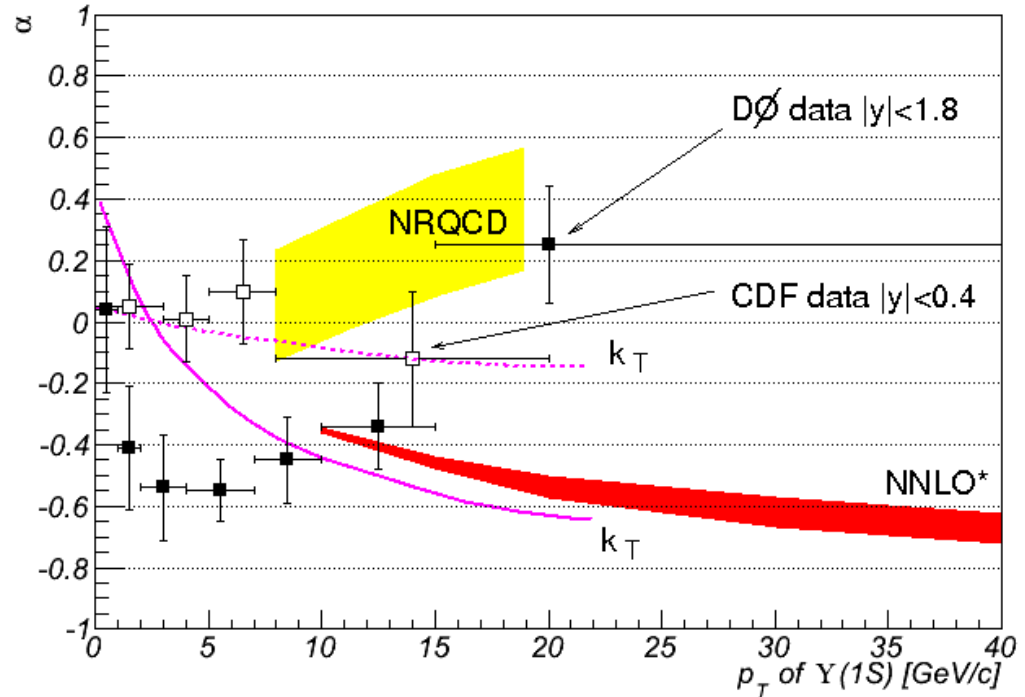


FIG. 2 (color online). Monte Carlo  $|\cos\theta^*|$  distributions after all selection requirements for different  $\alpha$  values:  $-1$  (dashed histogram),  $0$  (solid histogram), and  $+1$  (dotted histogram). (a)  $0 < p_T^\Upsilon < 1$  GeV/c, (b)  $p_T^\Upsilon > 15$  GeV/c.

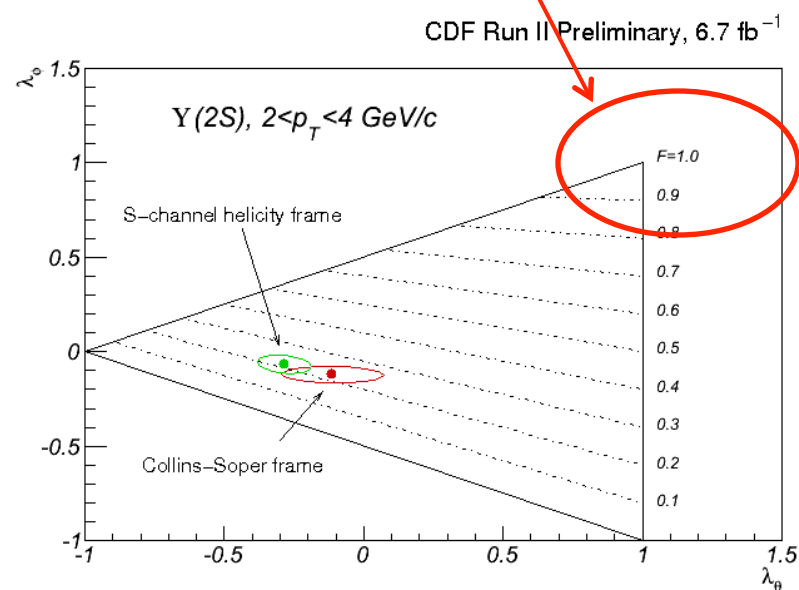
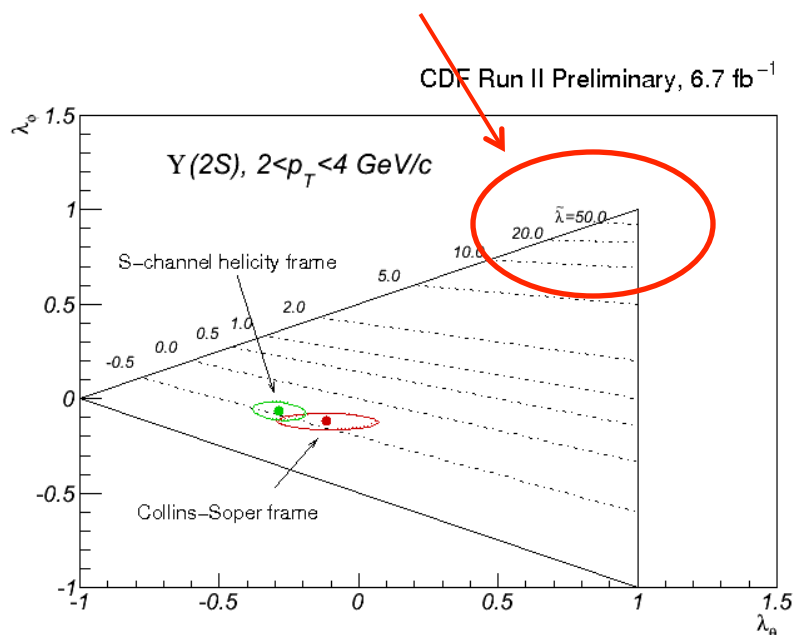


- DØ Run II: [Phys. Rev. Lett. 101, 182004 \(2008\)](#).
- CDF Run I: [Phys. Rev. Lett. 88, 161802 \(2002\)](#).
- NRQCD: [Phys. Rev. D63, 071501\(R\) \(2001\)](#).
- $k_T$ -factorization: [JETP Lett. 86, 435 \(2007\)](#).
- NNLO\*: [Phys. Rev. Lett. 101, 152001 \(2008\)](#).

# Other Rotational Invariants

$$\lambda = \lambda_{\theta} + 3\lambda_{\phi} / 1 - \lambda_{\phi}$$

$$F = \frac{1 + \lambda_{\theta} + 2\lambda_{\phi}}{3 + \lambda_{\theta}}$$



$$\lambda = \lambda_{\theta} + 3\lambda_{\phi} / 1 - \lambda_{\phi} = 4 / 1 + |\cos\theta| / 2 - \cos\theta \cdot \cos\theta - \cos\theta \cdot \cos\theta - 3$$

This is the part that is invariant under rotations.

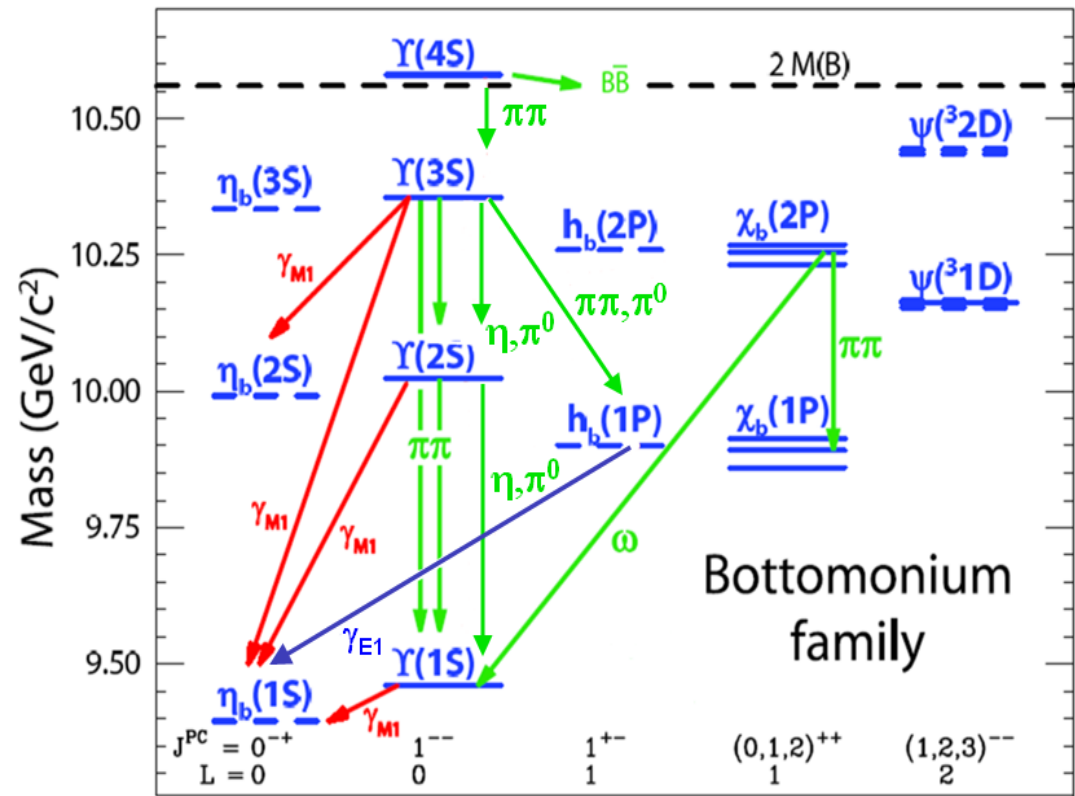
# Bottomonium Spectroscopy

$$\eta_b(nS) = n^1S_0$$

$$\Upsilon(nS) = n^3S_1$$

$$h_b(nP) = n^1P_1$$

$$\chi_{bJ}(nP) = n^3P_J$$



# Theoretical Description

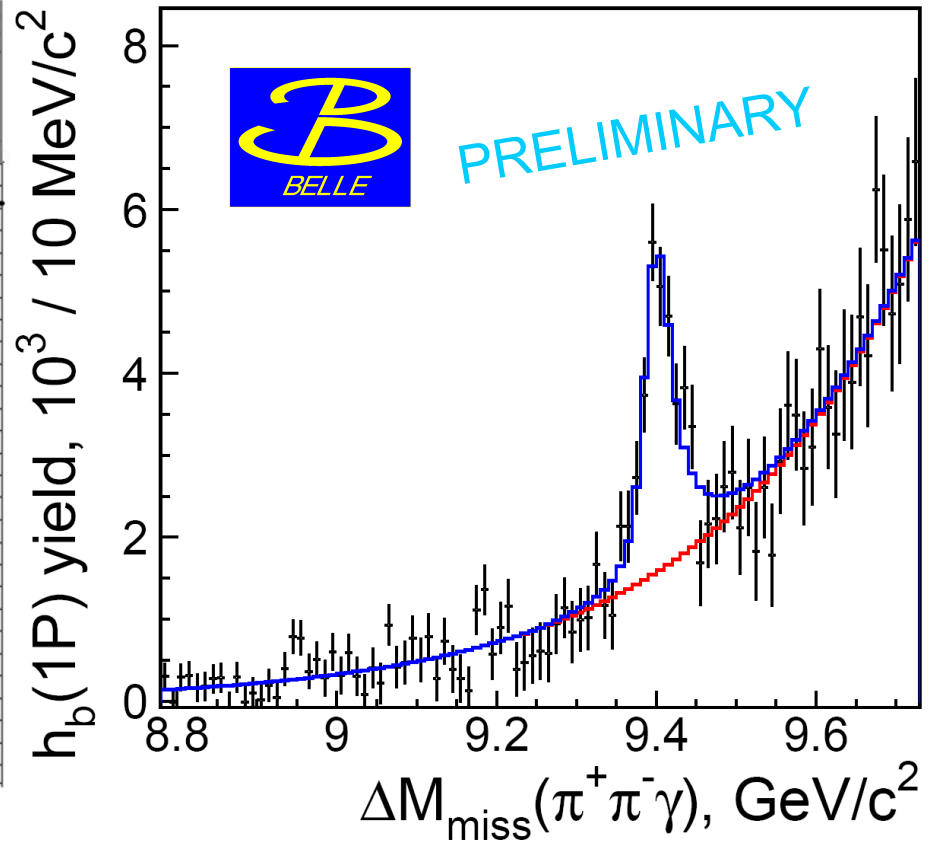
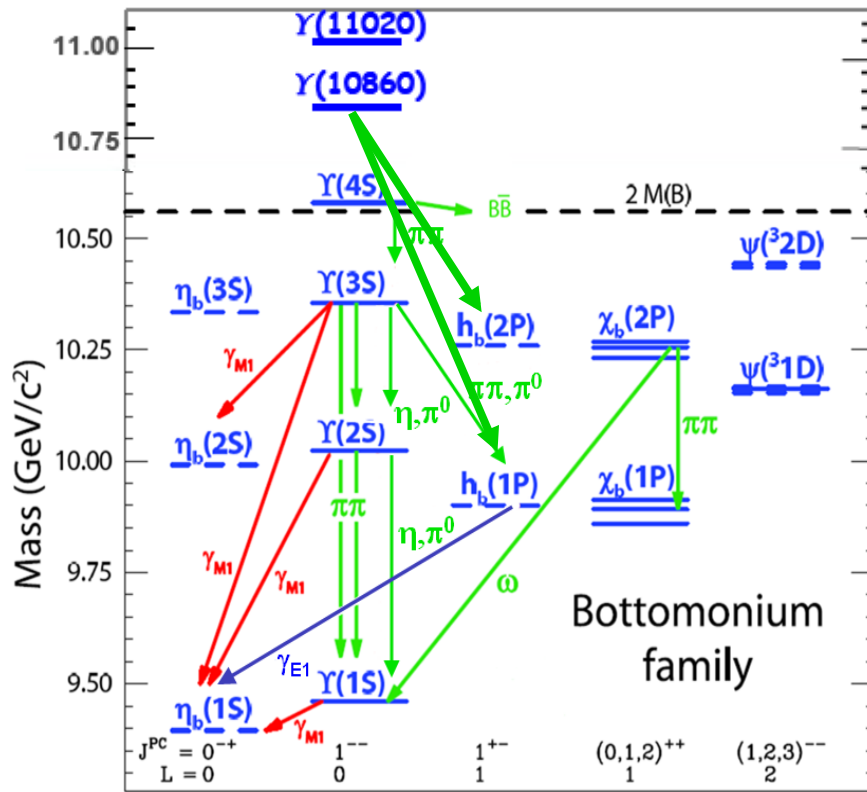
- Heavy quarks  $\rightarrow$  non-relativistic mechanics
- Potential models:

$$V_0(r) = -\frac{4}{3}\frac{\alpha_s}{r} + br + \frac{32\pi\alpha_s}{9m_Q^2}\delta(r)\vec{S}_Q \cdot \vec{S}_{\bar{Q}}$$

$$V_{spin-dep} = \frac{1}{m_Q^2} \left[ \left( \frac{2\alpha_s}{r^3} - \frac{b}{2r} \right) \vec{L} \cdot \vec{S} + \frac{4\alpha_s}{r^3} T \right]$$

- Reasonably good empirical description of spectrum and transitions.
- Small  $1/m_Q \rightarrow$  Effective field theories
  - HQET:  $1/m_Q$
  - NRQCD:  $\alpha_s, v : (M_Q v^2)^2 \ll (M_Q v)^2 \ll M_Q^2$

# Bottomonium Spectroscopy



$$Y(5S) \rightarrow Z_b^+ \pi^-$$

$$\hookrightarrow h_b(nP) \pi^+$$

$$\hookrightarrow \eta_b(mS) \gamma$$

QWG, October 2011

# Color Evaporation Model

- $c\bar{c}$  pairs produced with  $2m_c < m < 2m_D$  must eventually form a bound state.

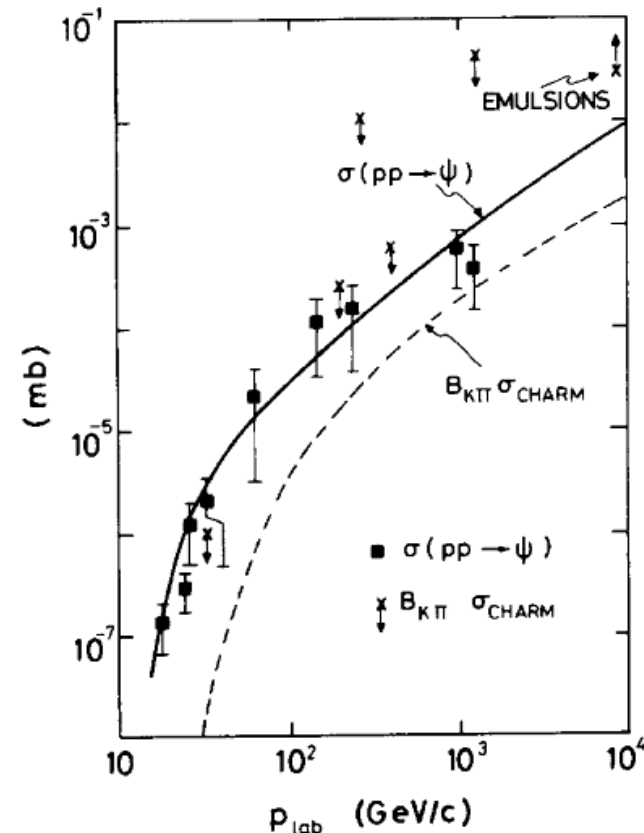
According to those the cross section for producing any  $c\bar{c}$  state below charm threshold is approximately equal to the cross section for producing a free  $c\bar{c}$  pair in the energy interval 3 ... 3.8 GeV:

$$\sum_{c\bar{c}} \sigma(p_1 + p_2 \rightarrow (c\bar{c}) + X) \approx \int_3^{3.8} \frac{d\sigma}{dM}(p_1 + p_2 \rightarrow \mu^+\mu^- + X) \frac{2\kappa^2}{3\alpha^2 e^2} dM. \quad (6)$$

Fritzsch - [Phys. Lett. B 67, 217 \(1977\)](#)

- Unable to predict polarization...

Halzen - [Phys. Lett. B 69, 105 \(1977\)](#)



# Color Evaporation Model

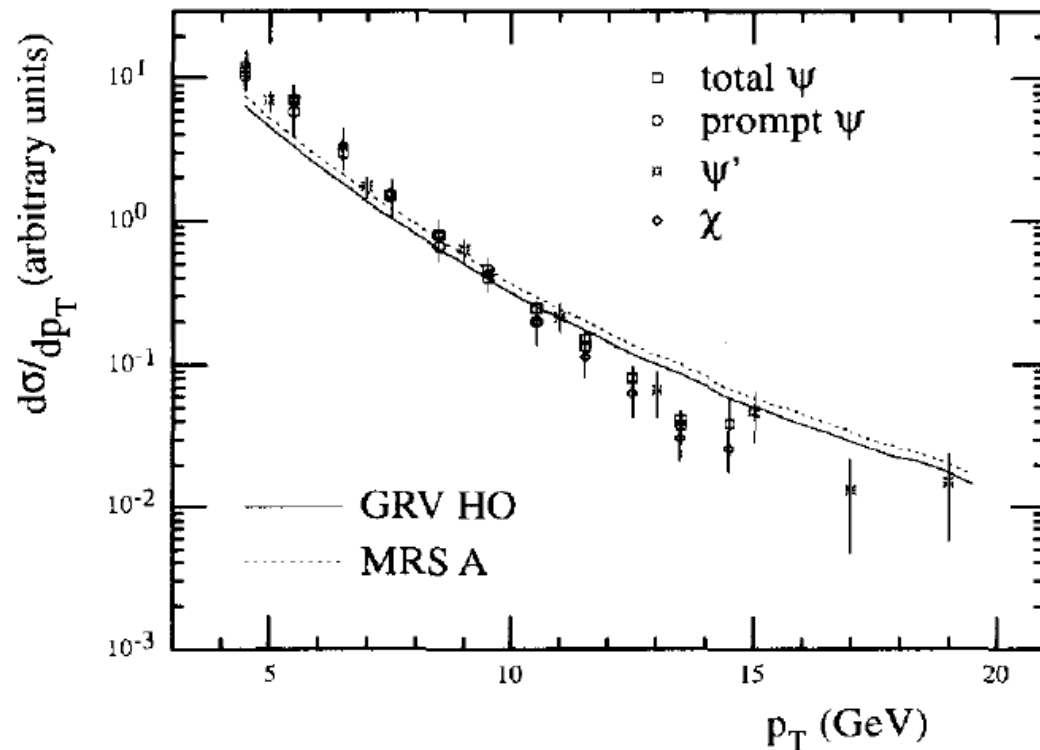


Fig. 6. Data from the CDF Collaboration [23], shown with arbitrary normalization. The curves are the predictions of the color evaporation model at tree level, also shown with arbitrary normalization. The normalization is correctly predicted within a  $K$  factor of 2.2.

Compare the overall shape of the  $p_T$  spectrum...

Maybe okay?

...but everything has been scaled...

# Heavy Quarkonium: $\psi(cc)$ and $\Upsilon(bb)$

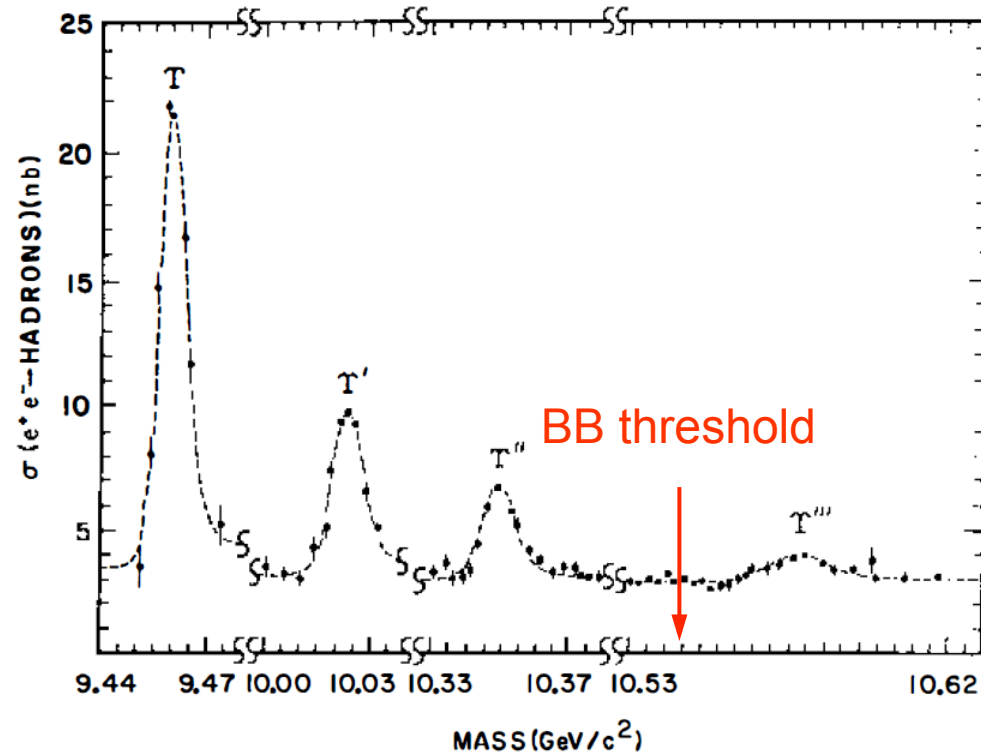


Figure 3 Cross section for  $e^+e^-$  annihilations into hadrons at CESR (CUSB data).

- Very simple system – non-relativistic QM works:

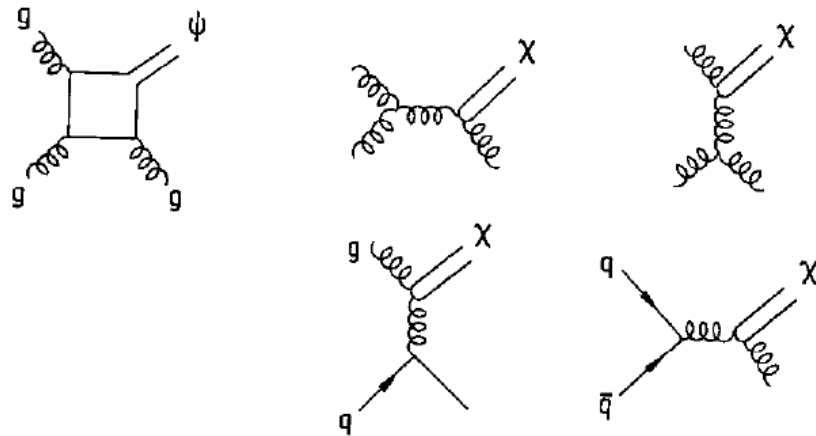
$$E\psi_n(x) = \left(-\frac{\hbar^2 \nabla^2}{2m} + V(x)\right)\psi_n(x)$$



# Can QCD Describe Heavy Quark Production?

Einhorn & Ellis: [Phys. Rev. D12, 2007 \(1975\)](#).

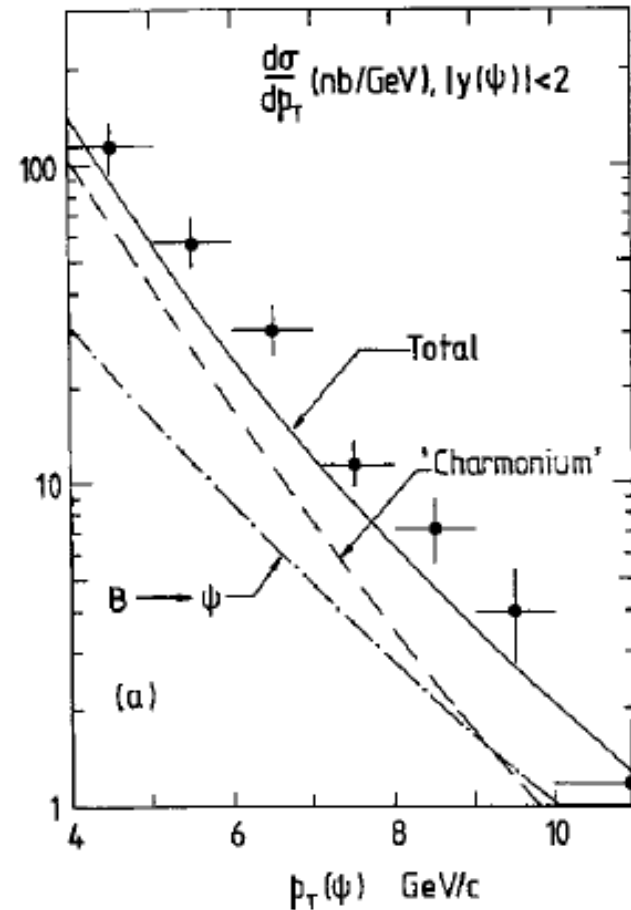
Glover, Martin & Stirling: [Z. Phys. C38, 473 \(1988\)](#).



The observed shape agrees well with the QCD expectations and the normalisation is within the error associated with the QCD calculation.

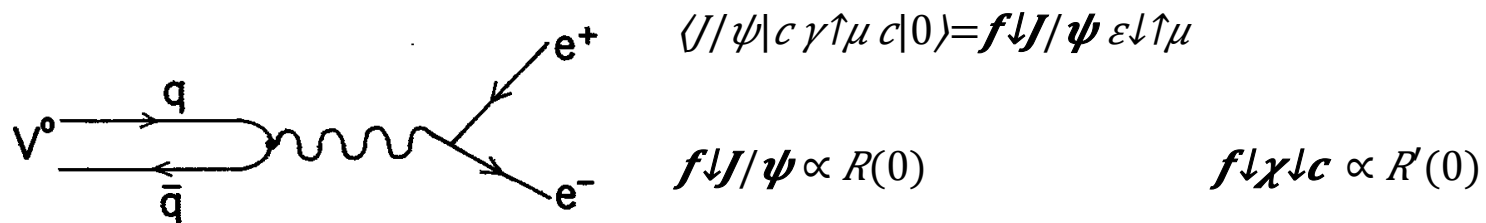
Or not...

Comparison with UA1 data at  $\sqrt{s} = 630 \text{ GeV}$

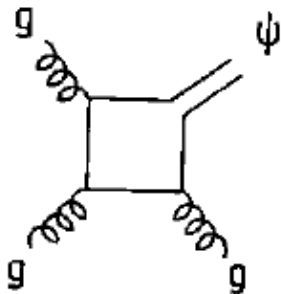


# Color-Singlet Production Model

- Production/decay via  $e^+e^-$  :



- Production at hadron colliders:



$$\frac{d\hat{\sigma}}{d\hat{t}} = \alpha_s^3(Q^2, \Lambda^2) |R(0)|^2 |\mathcal{M}|^2$$


- Matrix elements also predict **polarization**.

# Non-Relativistic QCD

Caswell & Lepage – [Phys. Lett. 167B, 437 \(1986\)](#)

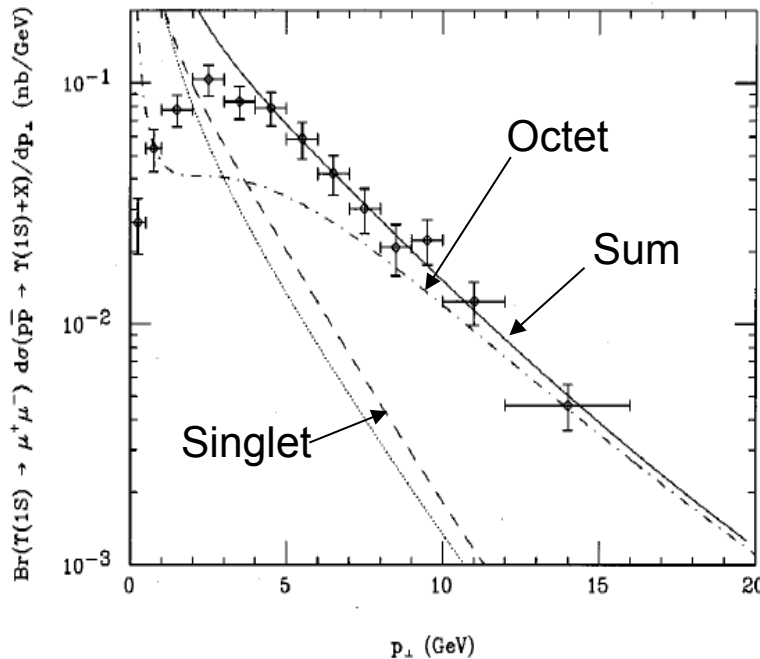
Bodwin, Braaten & Lepage – [Phys. Rev. D 51, 1125 \(1995\)](#)

- Expansion in powers of  $\alpha_s$  and  $1/m$
- Factorization of different energy scales:

- Bound states are “color singlets” – no net color charge.  

- 
- *Color-octet terms might be really important!*

# NRQCD + Color-Octet Models

- Matrix elements tuned to accommodate Tevatron results



Cho & Leibovich, PRD 53, 6203 (1996).

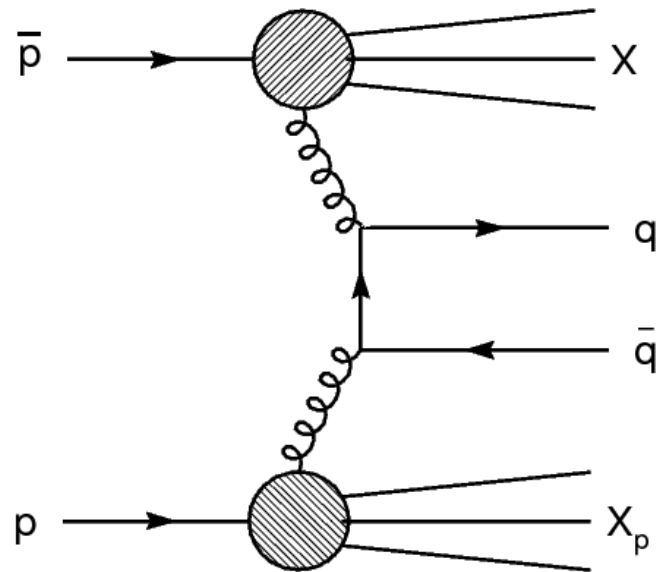
Unknown NRQCD Matrix Elements adjusted to match data.

Agreement with cross section is not too surprising now.

***We need an independent observable to really test the model.***

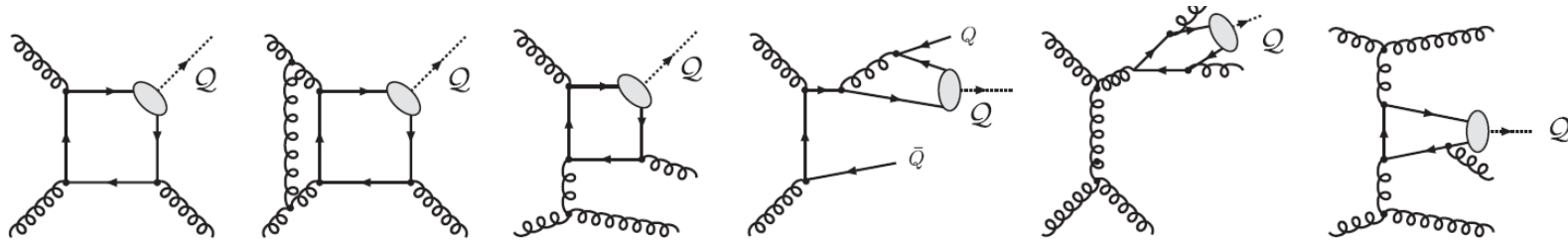
- Nearly on-shell gluons can fragment to form  $\Upsilon$
- Predicted **transverse**  $\Upsilon$  polarization for  $p_{\perp T} \gg m_{\perp Q}$

# Another Model: “ $k_T$ factorization”

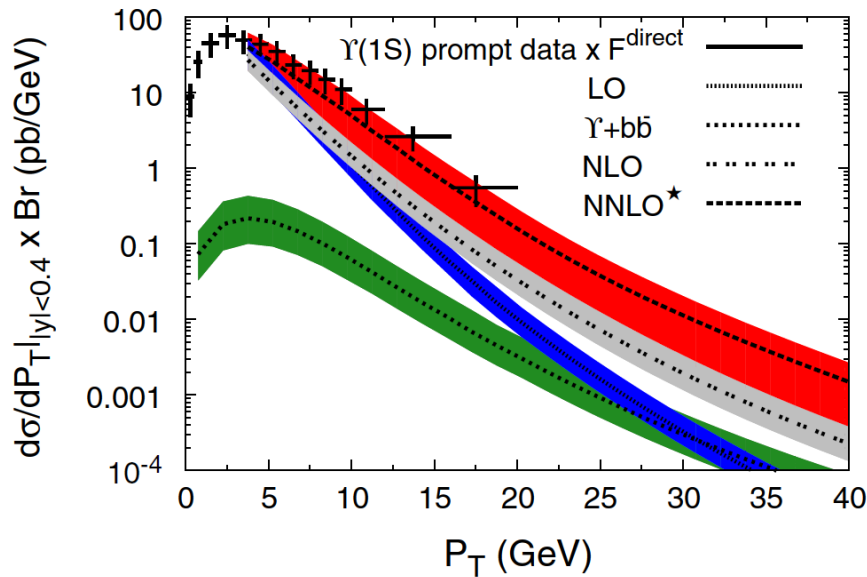


- Initial state gluon polarization related to their transverse momentum,  $k_T$ .
- No need for color-octet terms...
- Predicted **longitudinal**  $\Upsilon$  polarization for  $p_{\perp T} \gg m_{\perp Q}$

# Higher-order QCD calculations



Artoisenet, et al – [Phys. Rev. Lett. 101, 152001 \(2008\)](#).

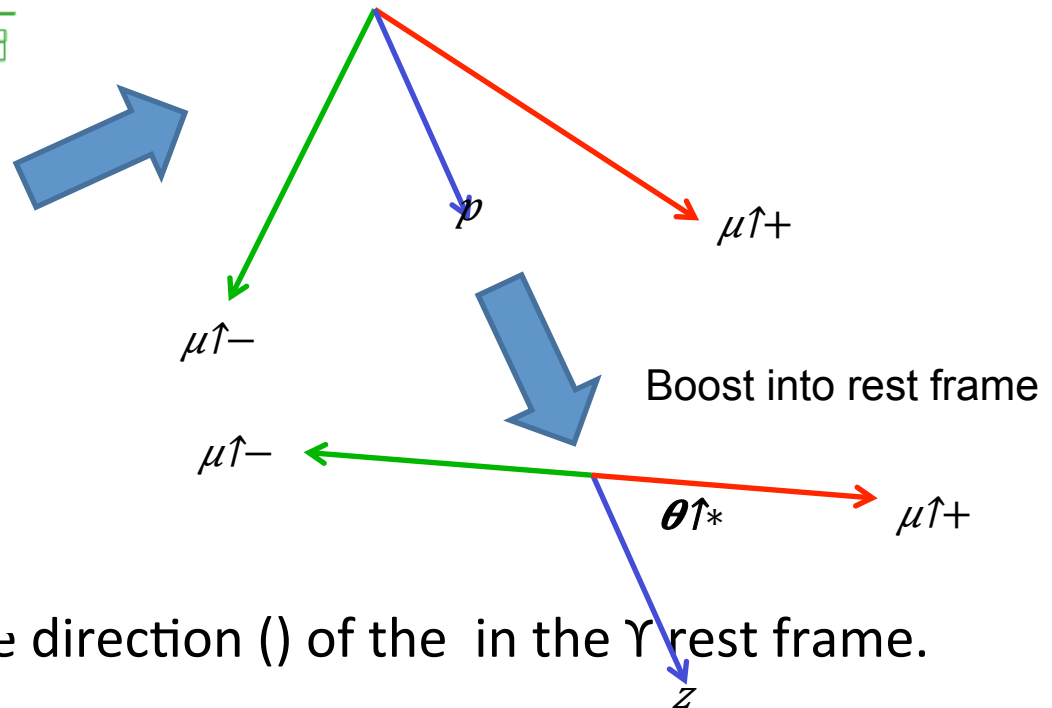
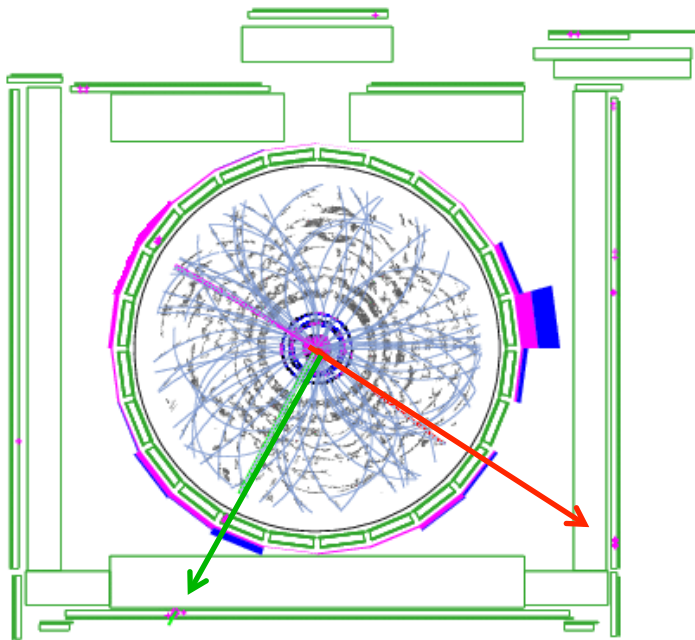


- *Partial* calculation including terms up to  $\alpha_s^5$  ...
- Large increase in cross section compared with LO calculation
- No need for color-octet contributions

- Predicts **longitudinal**  $\Upsilon$  polarization for  $p_T \gg m_Q$

# Measuring “Polarization”

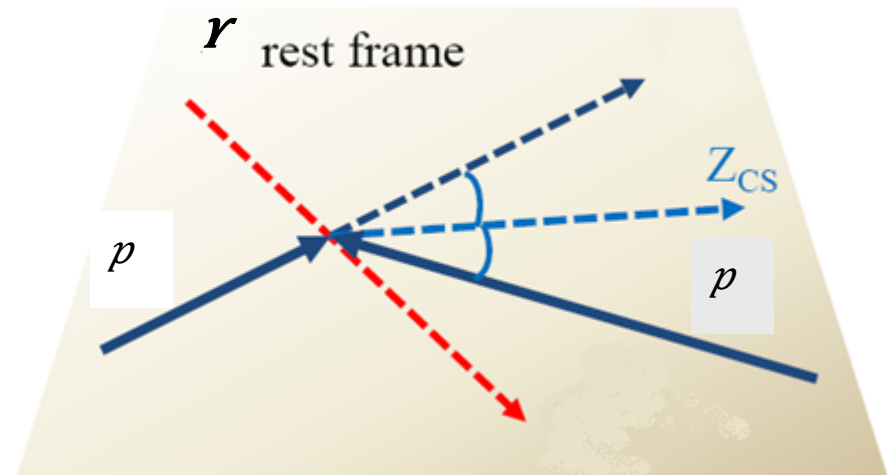
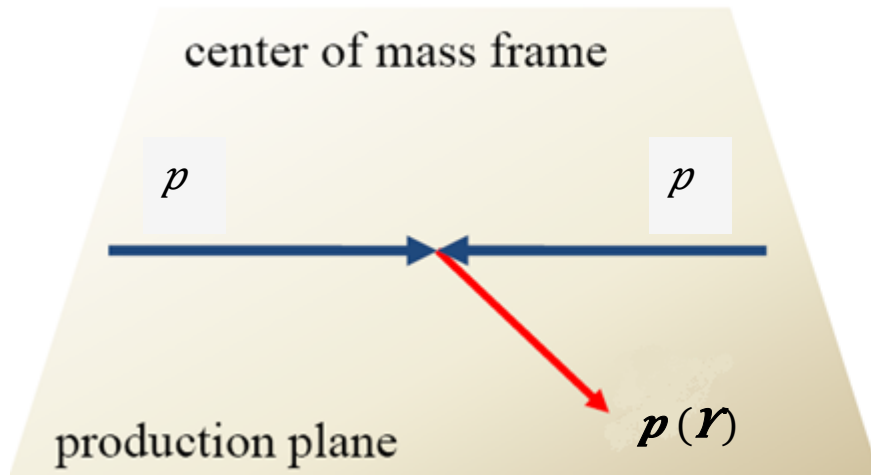
- We don't really measure polarization...



- we actually measure the direction ( $\theta\uparrow^*$ ) of the  $\mu\uparrow+$  in the  $\Upsilon$  rest frame.

# Which coordinate system?

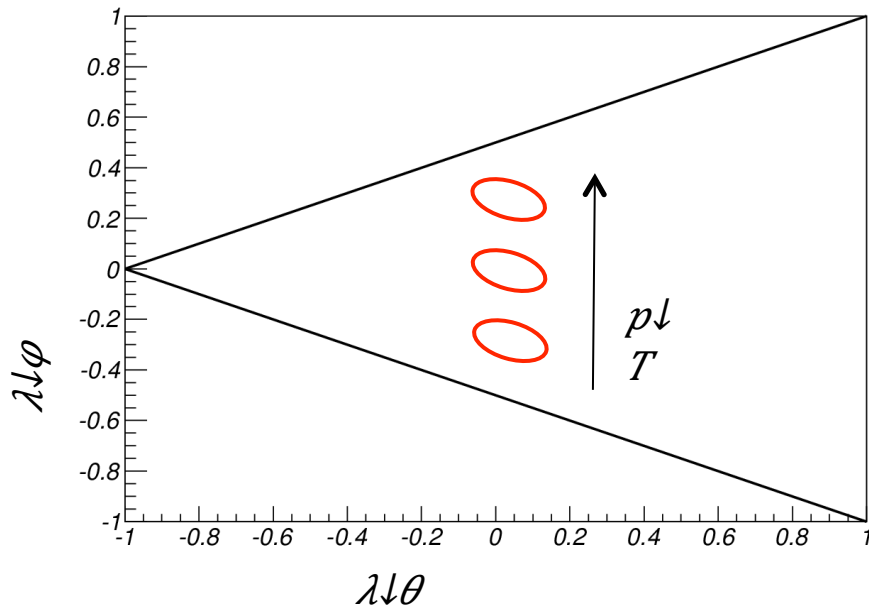
- S-channel Helicity (SH) –  $\mathcal{Y}$  momentum vector defines the z-axis, the x-axis is in the production plane
- Collins-Soper (CS) – z-axis bisects beam momentum vectors in  $\mathcal{Y}$  rest frame, x-axis in the production plane:



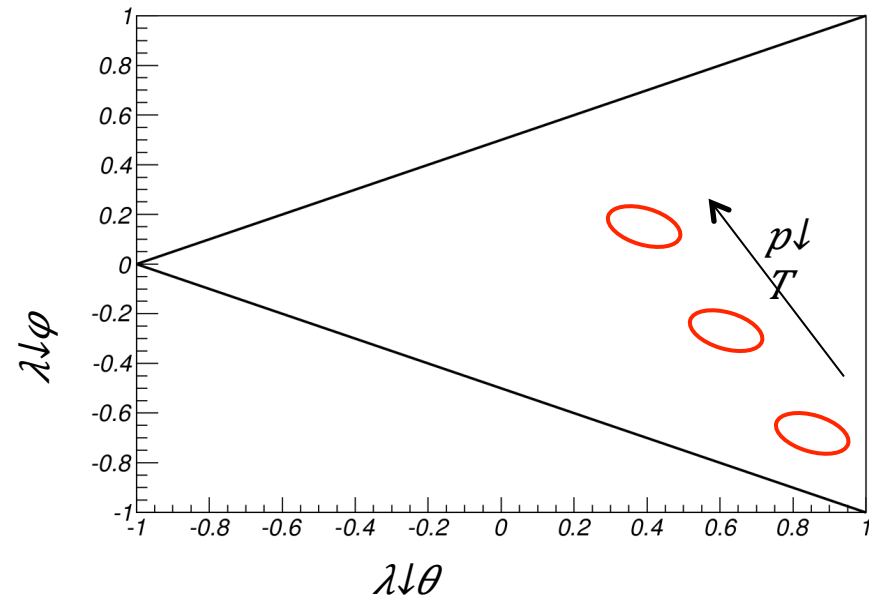


# Could it be possible?

**S-channel helicity  
frame**



**Collins-Soper  
frame**

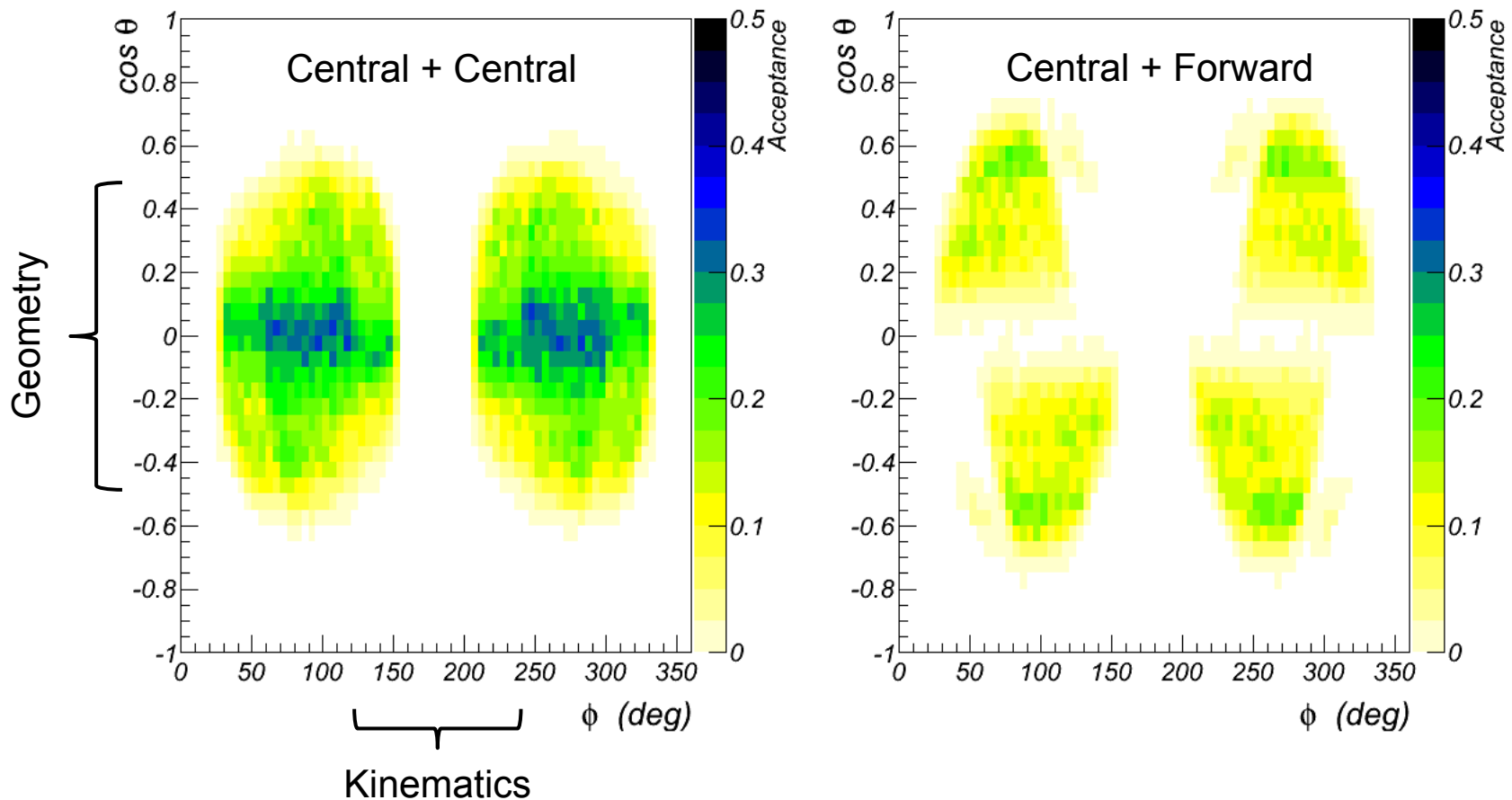


If  $\lambda\downarrow\theta$  is zero in one coordinate frame, then  
it **must** be non-zero in another frame!

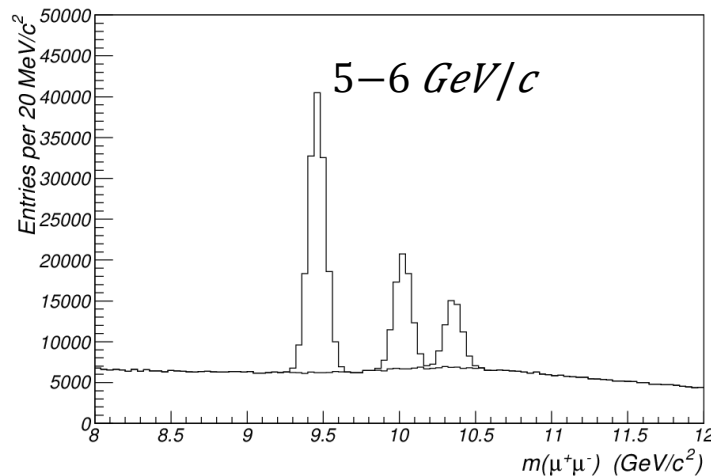
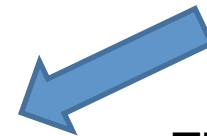
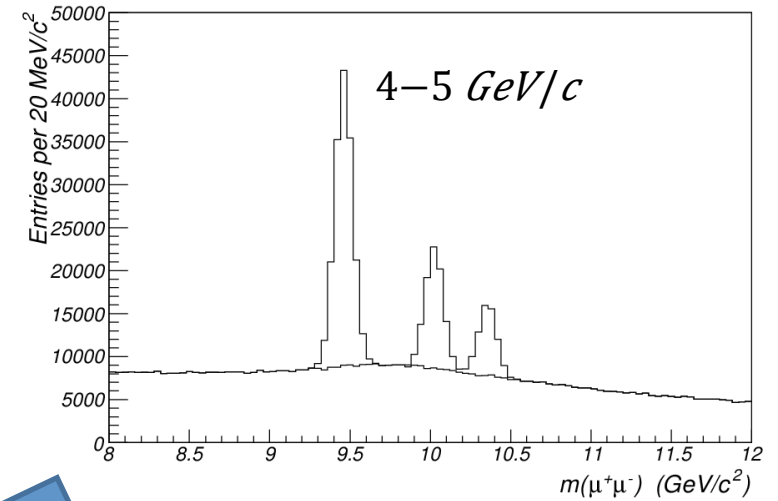
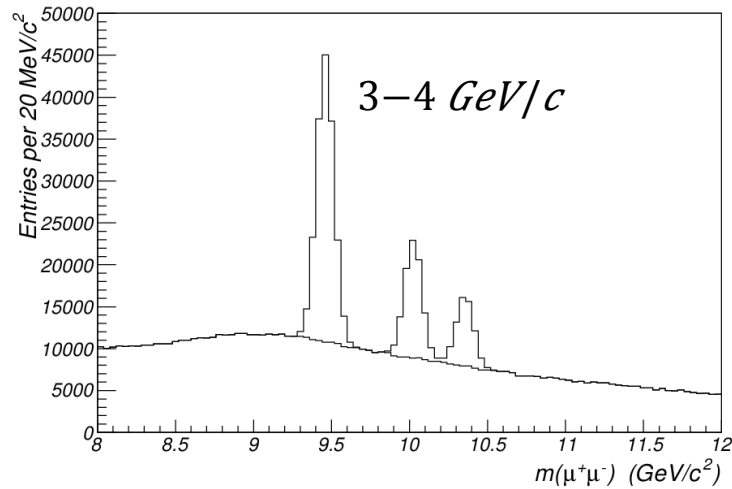
(provided  $\lambda\downarrow\varphi$  is not also zero)

# Geometric Acceptance

- Geometric acceptance calculated with full detector simulation for each  $p_T$  range analyzed
- Muon detectors simulated with 100% efficiency



# Background Structure



This is just all toy Monte Carlo but it makes us worried...

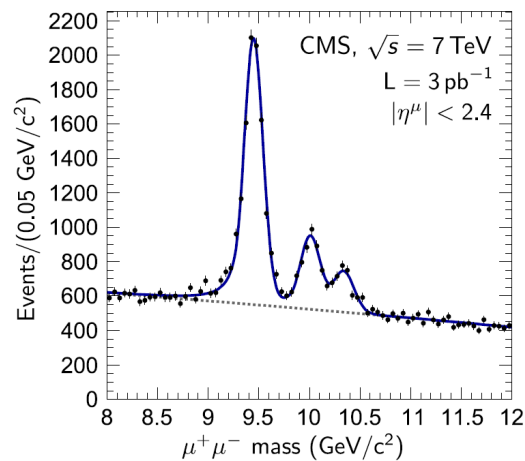
A polynomial may not describe the mass distribution under the signal when fitted using just the sidebands.

# Systematic Uncertainties

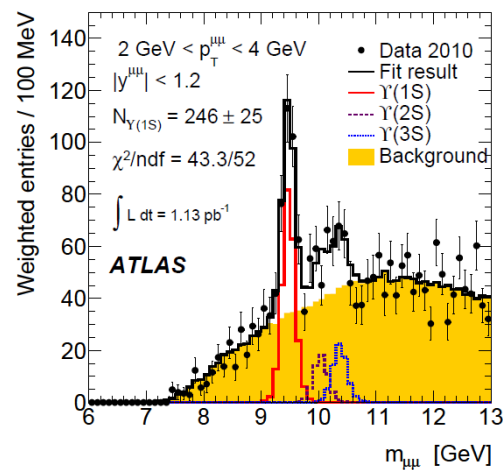
- Efficiency measurement:
  - Vary measured trigger efficiencies by
- Monte Carlo statistics:
  - Impact of finite sample sizes in acceptance calculated using toy Monte Carlo experiments
- Background scale factor:
  - Compare linear and quadratic interpolation from sidebands into signal region
- Frame invariance tests:
  - Treat as a systematic uncertainty
  - Consistent with statistical fluctuations in almost all cases
- All are generally much smaller than statistical uncertainty

# New Cross Section Measurements

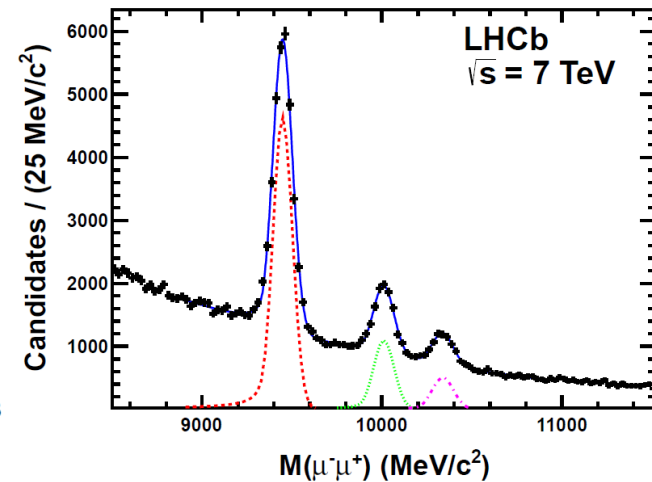
[Phys. Rev. D83, 112004 \(2011\)](#)



[Phys. Lett. B 705 \(2011\), 9](#)



[arXiv:1202.6579](#)



- 10-20% systematic uncertainty due to unknown polarization.