Connecting Collider Data with Air-Shower Physics

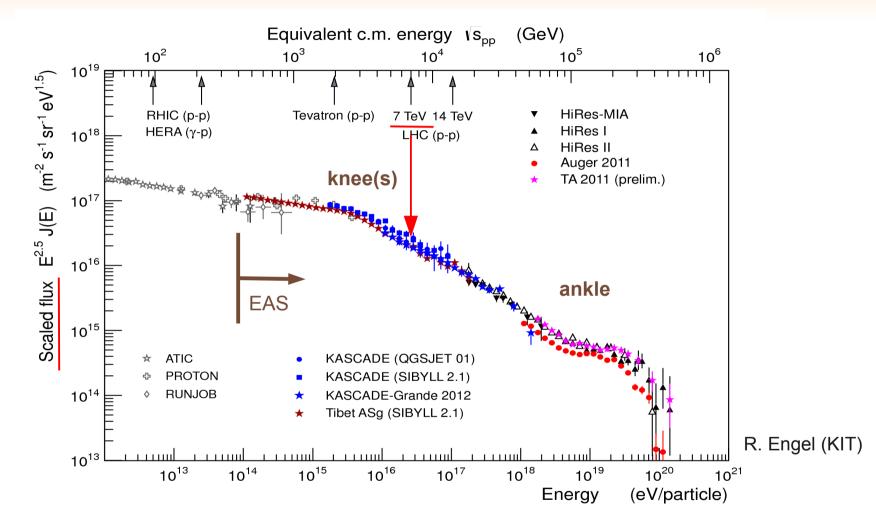
Tanguy Pierog

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LHC-CR Workshop, CERN February the 12th 2013

Cosmic Ray Spectrum



- Origins of spectrum properties
 - mostly unknown
 - depend on primary CR mass

- Most of analysis based on EAS simulations
 - CORSIKA AIRES
 - ➡ COSMOS
 ➡ CONEX, ...

Muons in EAS



Hadronic Interaction Models for CR

New models

Connection with Cosmic Rays (CR)

Heitler model

- LHC and Xmax
 - Longitudinal development
- LHC and muons in Extensive Air Showers (EAS)
 - Particles at ground
- Summary

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Hadronic Interaction Models

- Theoretical basis :
 - → pQCD (large p_t)

Pb : CR physic dominated by soft interactions

- Gribov-Regge (cross section with multiple scattering)
- energy conservation
- Phenomenology (models) :
 - string fragmentation
 - beam remnants
 - diffraction (Good-Walker, ...)
 - higher order effects
- Comparison with data to fix parameters :
 - minimum theory requirement with few parameters and limited data set (QGSJET approach) : better predictive power (if nothing forgotten...)

... or ...

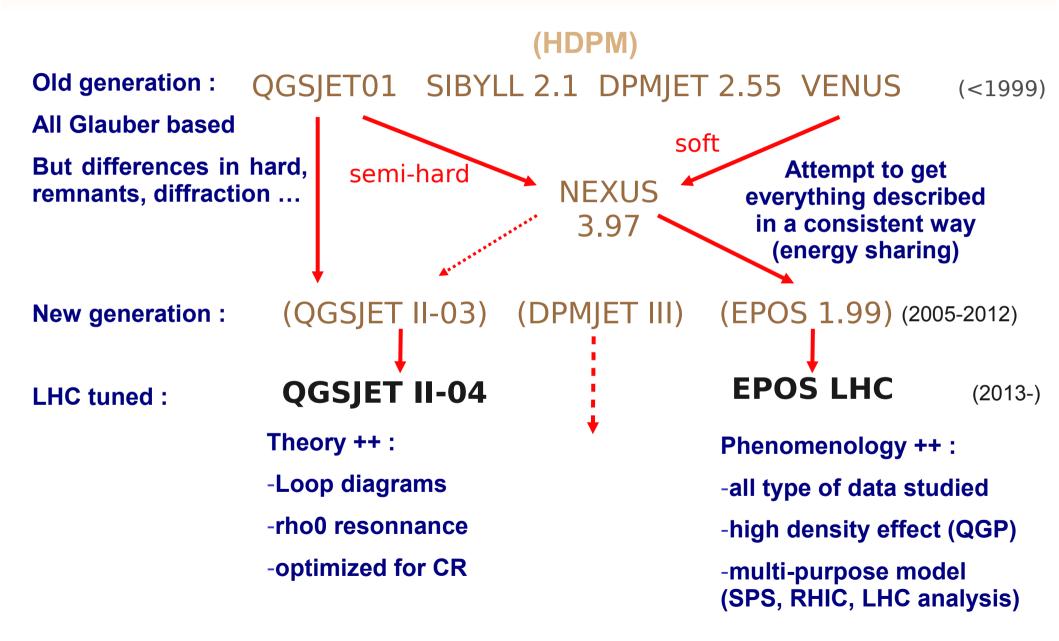
more detailed data with more parameters (EPOS approach) : nothing neglected (but more uncertainties in extrapolation)

- Pb : Gribov-Regge do not take into account energy conservation ...
- **Need Parameters !**



Muons in EAS

Hadronic Interaction Models in CORSIKA



Muons in EAS

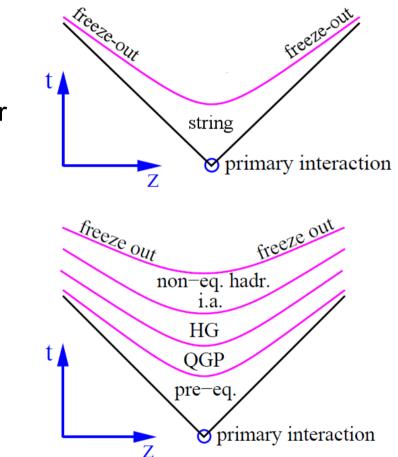
New Models

QGSJETII-04 :

- loop diagrams
- rho0 forward production in pion interaction
- re-tuning some parameters for LHC and lower energies

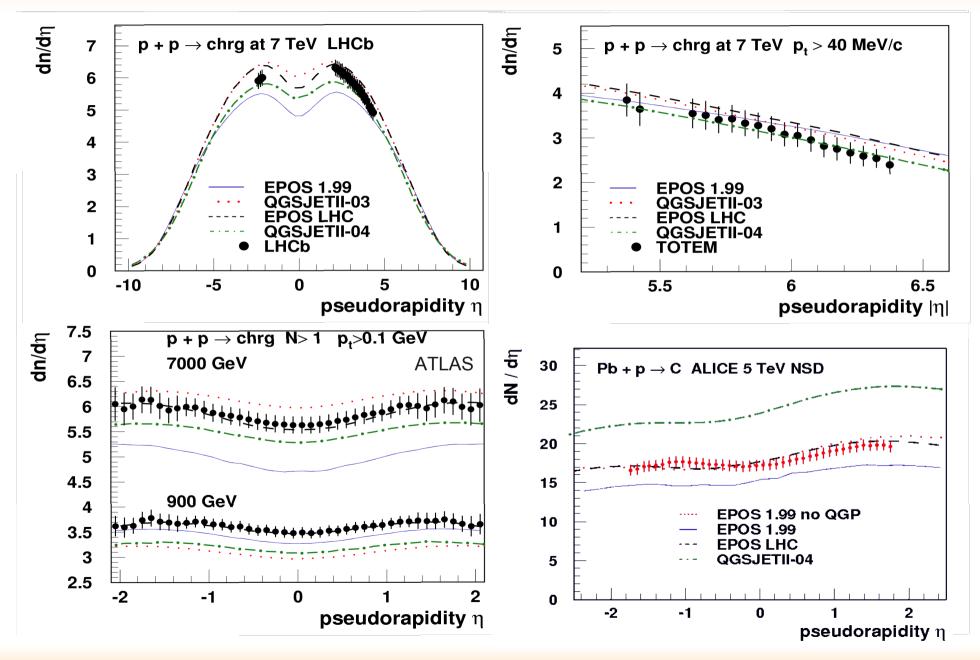
EPOS LHC

- tune cross section to TOTEM value
- change old flow calculation to a more realistic one
- introduce central diffraction
- keep compatibility with lower energies



Direct influence on EAS simulations to be shown but important to compare to LHC and set parameters properly (<pt>, ...).

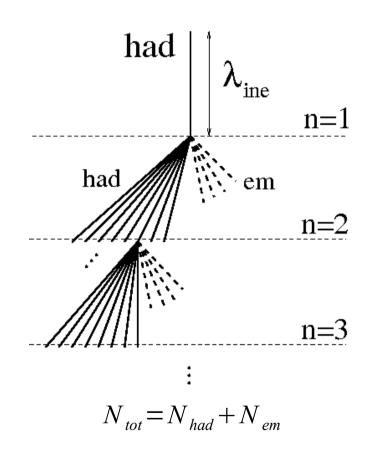
Pseudorapidity Distributions





Simplified Shower Development

Using generalized Heitler model and superposition model :



J. Matthews, Astropart.Phys. 22 (2005) 387-397

$$X_{max} \sim \lambda_e \ln \left((1-k) \cdot E_0 / (2 \cdot N_{tot} \cdot A) \right) + \lambda_{ine}$$

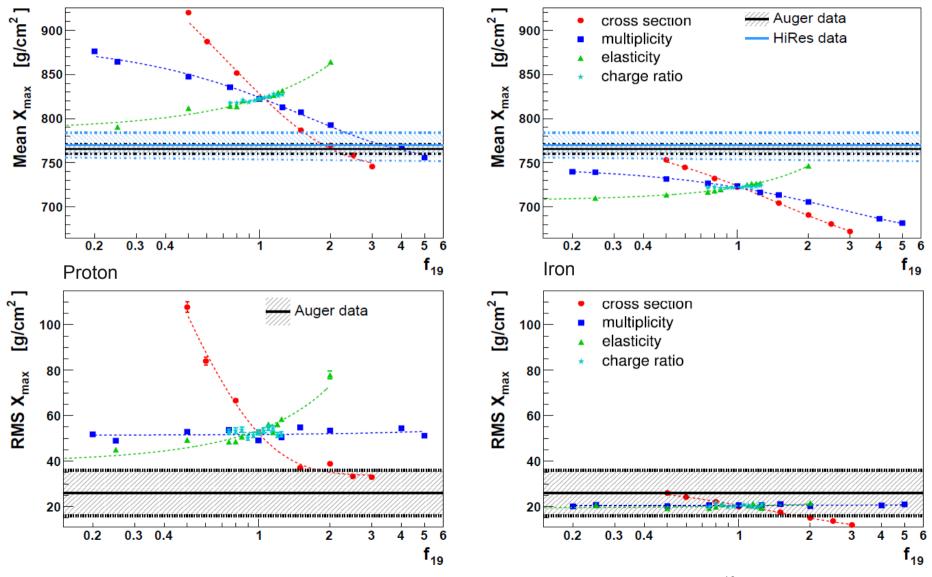
- Model independent parameters :
 - \bullet E₀ = primary energy
 - A = primary mass
 - $\lambda_{\rm a}$ = electromagnetic mean free path
- Model dependent parameters :
 - k = elasticity
 - N_{tot} = total multiplicity
 - λ_{ine} = hadronic mean free path (cross section)



Muons in EAS

Effects of Parameters

Sensibility depends on observable and parameter :



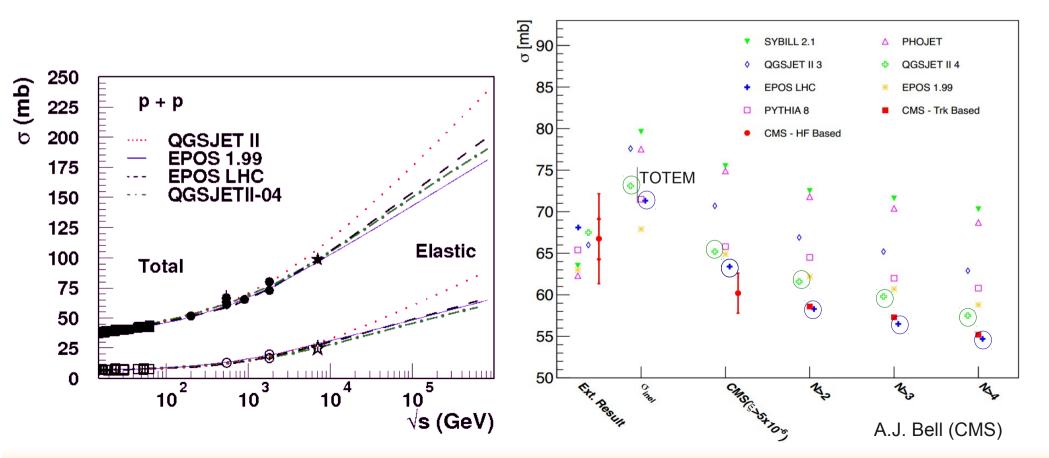
Plots by R. Ulrich (KIT) with Sibyll model and PAO data @ 10¹⁹ eV

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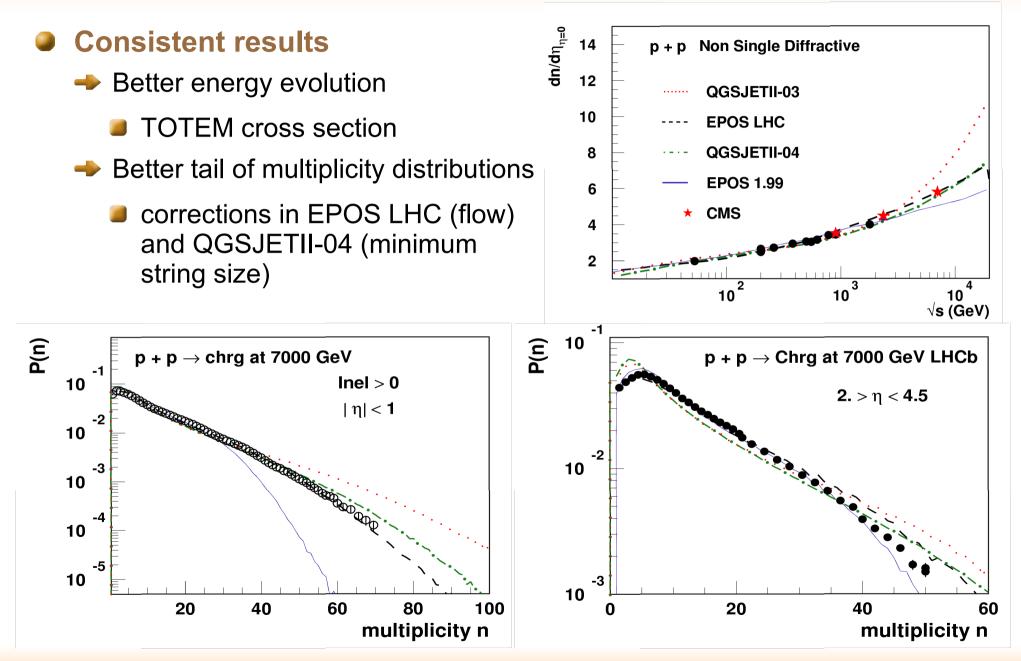
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Cross Sections

- Same cross sections at pp level up to LHC
- extrapolation to pA or to high energy
 - different amplitude and scheme : different extrapolations
- LHC measurements test the difference between models (diffraction)



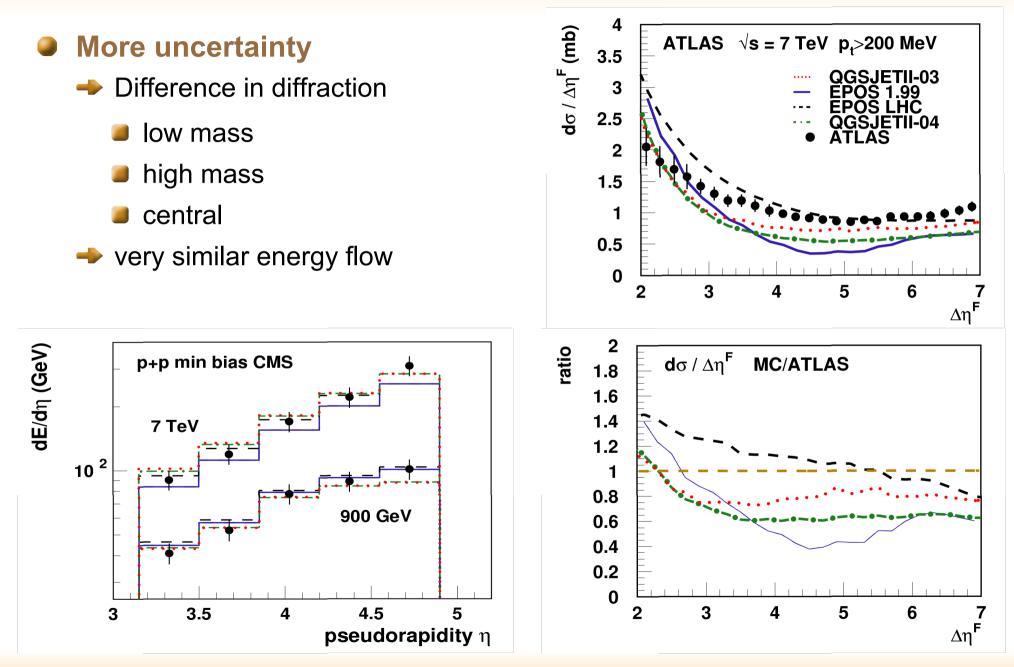
Multiplicity



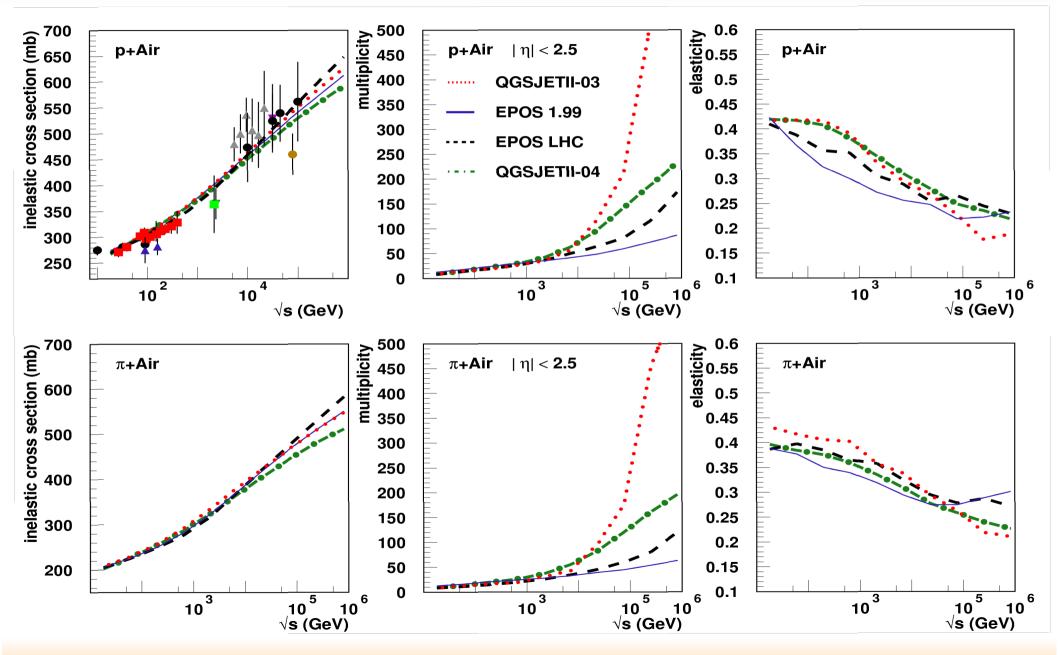
T. Pierog, KIT - 11/27

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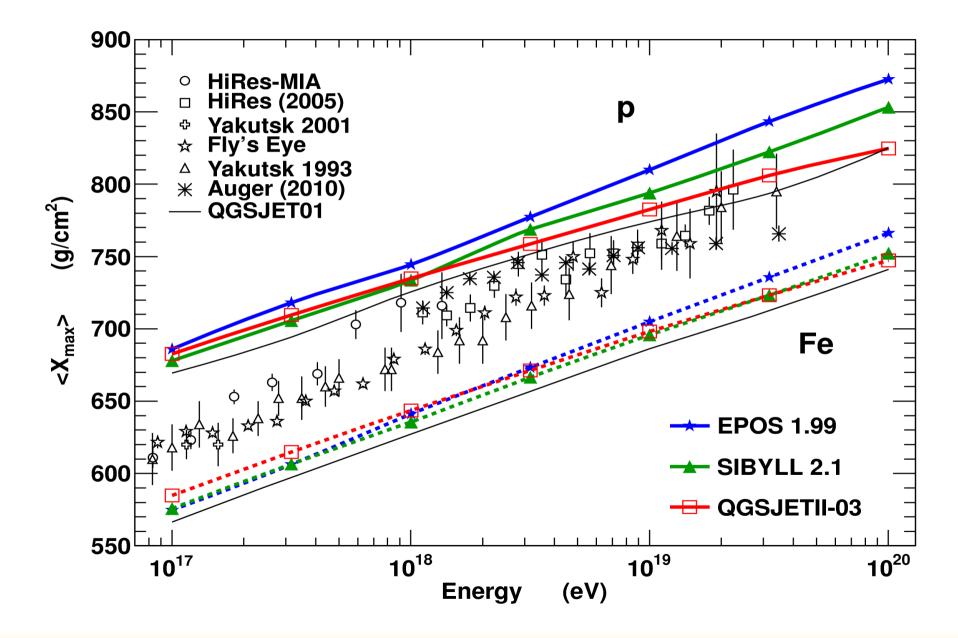
Inelasticity



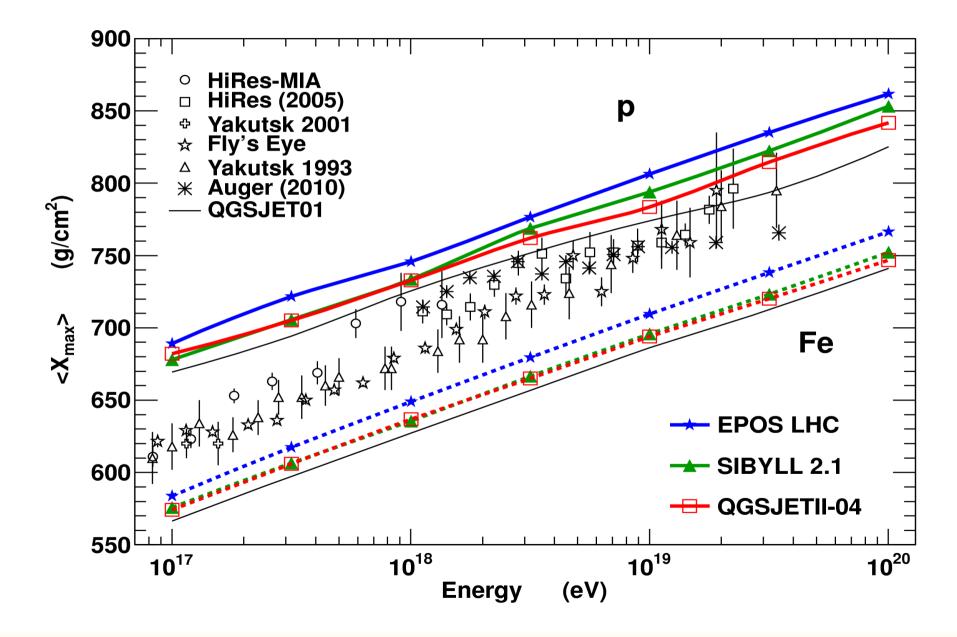
Predictions with retuned models



EAS with old CR Models : X_{max}



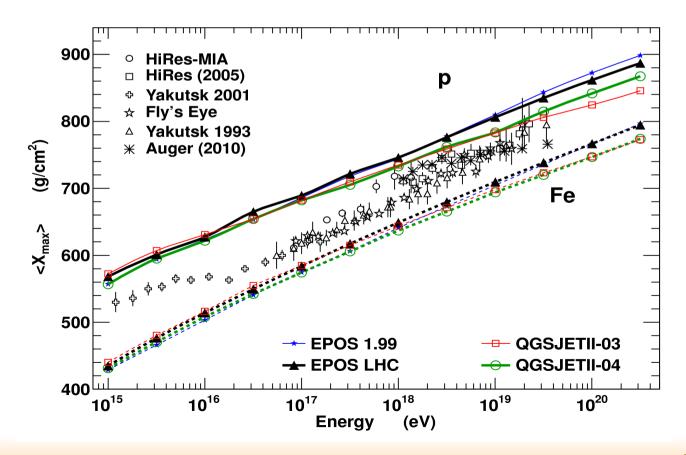
EAS with Re-tuned CR Models : X_{max}



EAS with Re-tuned CR models : X_{max}

Cross section and multiplicity fixed at 7 TeV

- smaller slope for EPOS and larger for QGSJETII
- re-tuned model converge to old Sibyll 2.1 predictions
 - reduced uncertainty from ~50 g/cm² to ~20 g/cm²
 (difference proton/iron is about 100 g/cm²)



LHC and Xmax

Muons in EAS

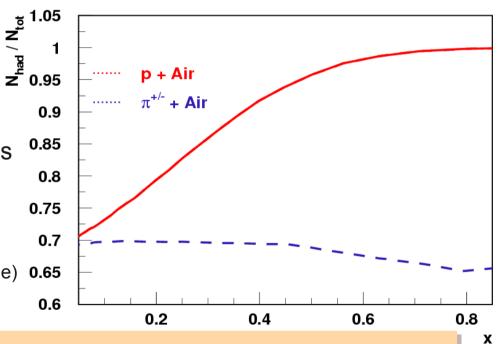
Muon Number

From Heitler

$$N_{\mu} = \left(\frac{E_0}{E_{dec}}\right)^{\alpha}, \quad \alpha = \frac{\ln N_{\pi^{ch}}}{\ln \left(N_{\pi^{ch}} + N_{\pi^0}\right)}$$

In real shower, not only pions : Kaons and (anti)Baryons (but 10 times less ...)

- \rightarrow Baryons do not produce leading π^0
- With leading baryon, energy kept in hadronic channel = muon production
- Cumulative effect for low energy muons
- High energy muons
 - important effect of first interactions
 and baryon spectrum (LHC energy range)



Muon number depends on the number of (anti)B in p- or π -Air interactions at all energies

More fast (anti)baryons = more muons

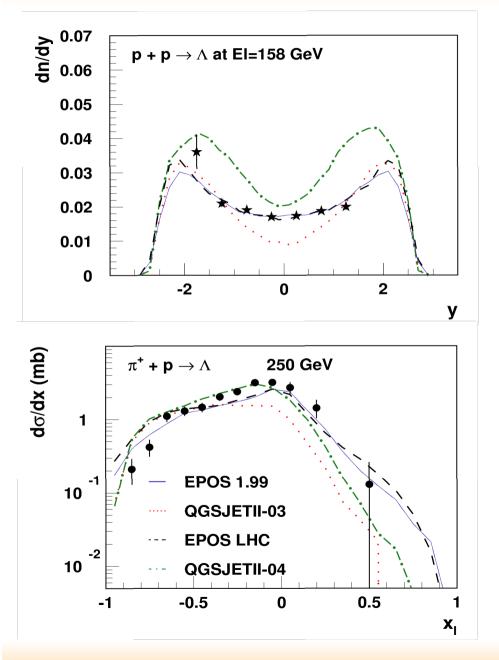
T. Pierog et al., Phys. Rev. Lett. 101 (2008) 171101

LHC-CR – February 2013

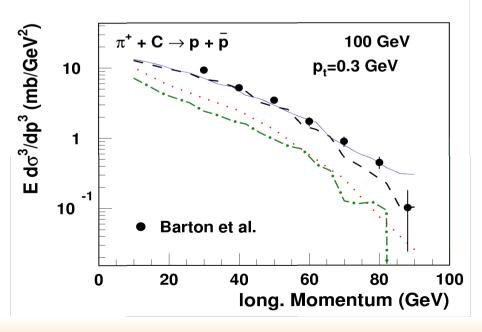
T. Pierog, KIT - 17/27

Muons in EAS

Baryon Forward Spectra



- Large differences between models
- Need a new remnant approach for a complete description (EPOS)
- Problems even at low energy
 - check data with NA61
- No measurement at high energy !
 - neutron spectra from LHCf ?



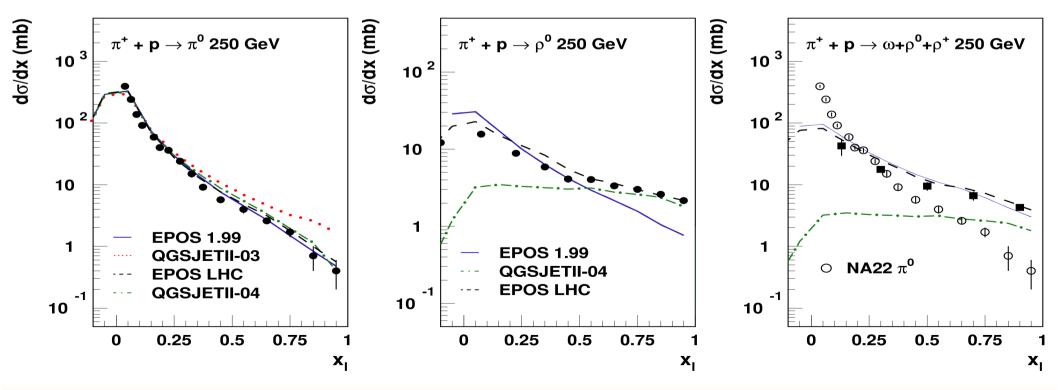
Muons in EAS

Pion Leading Particle Effect

- Rho meson production added in QGSJETII to take into account leading particle effect in pion-Air interaction
 - same effect as baryon production : forward π⁰ replaced by charged pions (reduced leading π⁰)
 - increase muon production

already in EPOS





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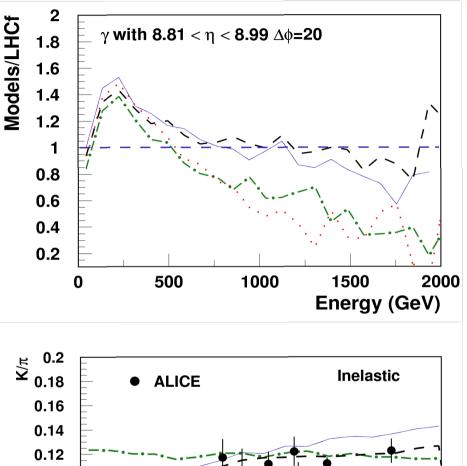
Identified particles

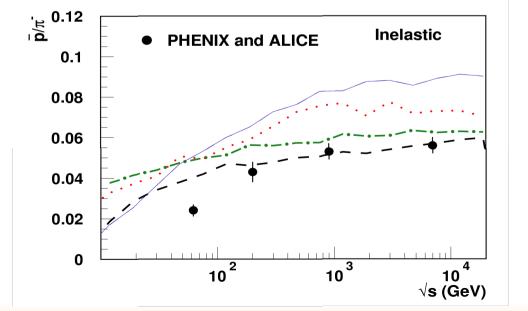
Large improvement at mid-rapidity

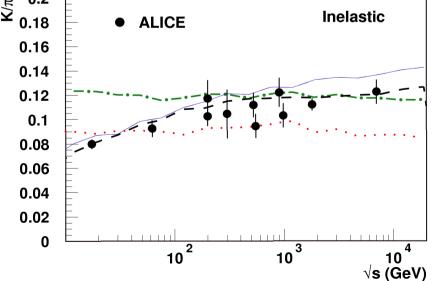
- very similar results for particle ratios
- overestimation of baryon production before due to wrong interpretation of Tevatron data

Only small changes very forward

no try to tune LHCf data (difficult)

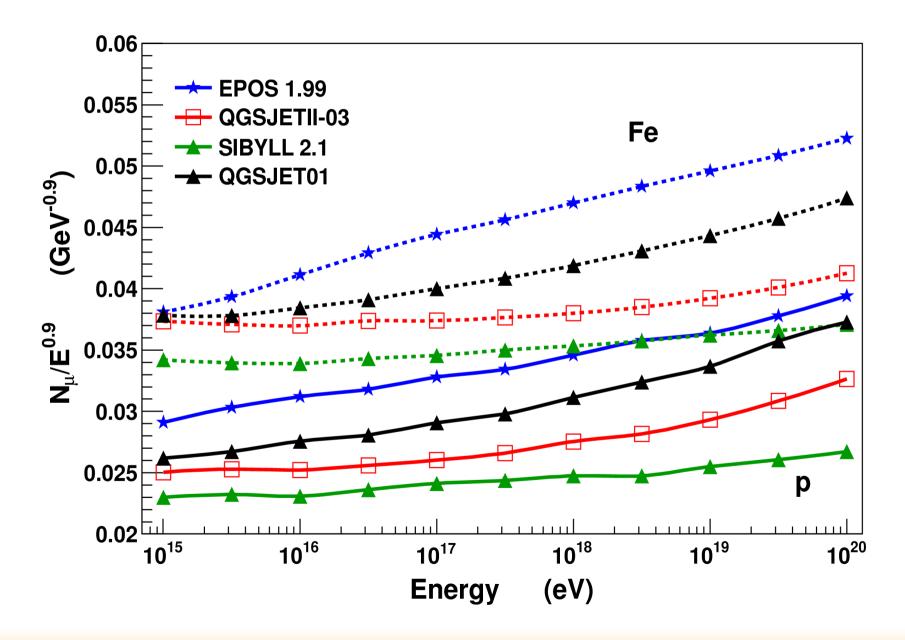






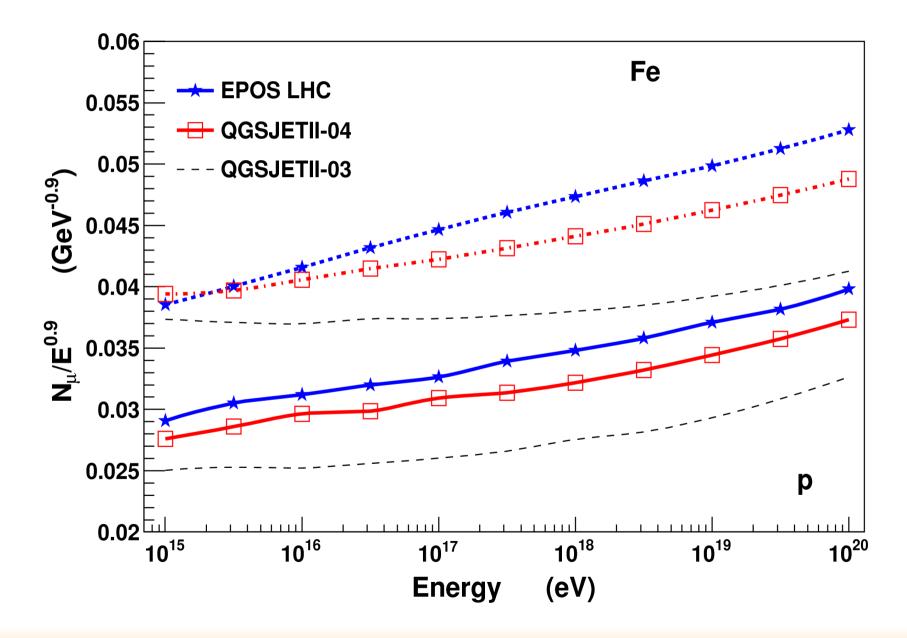
Muons in EAS

EAS with old CR models : Muons



Muons in EAS

EAS with Re-tuned CR models : Muons

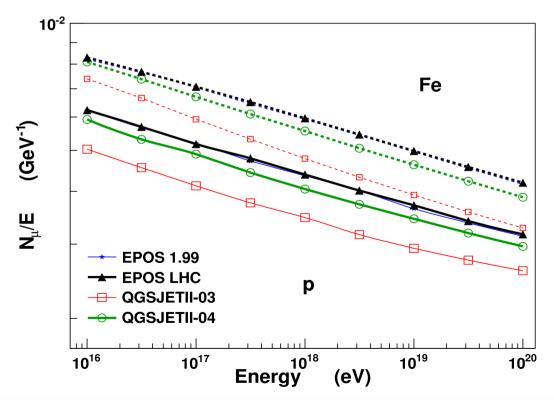


EAS with Re-tuned CR models : Muons

Weak effect of LHC

- Corrections at mid-rapidity thanks to LHC
- Changes in QGSJETII motivated by pion induced data
- Changes for forward production in EPOS LHC can not be checked by LHC (yet ?) (motivated by model consistency)

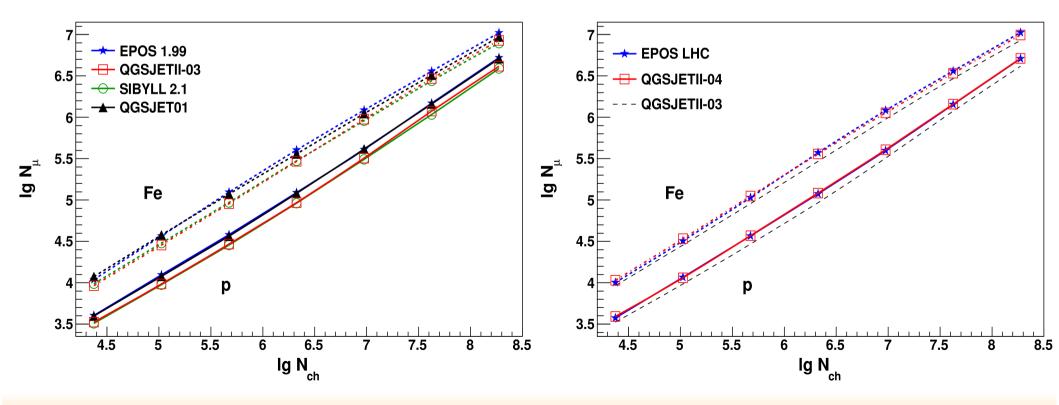
NA61 data wanted to check old data set



EAS with Re-tuned CR Models : Correlations

QGSJETII-04 and EPOS LHC similar to EPOS 1.99

- More muons AND more electrons in EPOS LHC than in QGSJETII-04
- More muons and less electrons in QGSJETII-04 than in QGSJETII-03
- Same correlations in EPOS LHC and QGSJETII-04
- Lighter composition than with QGSJETII-03

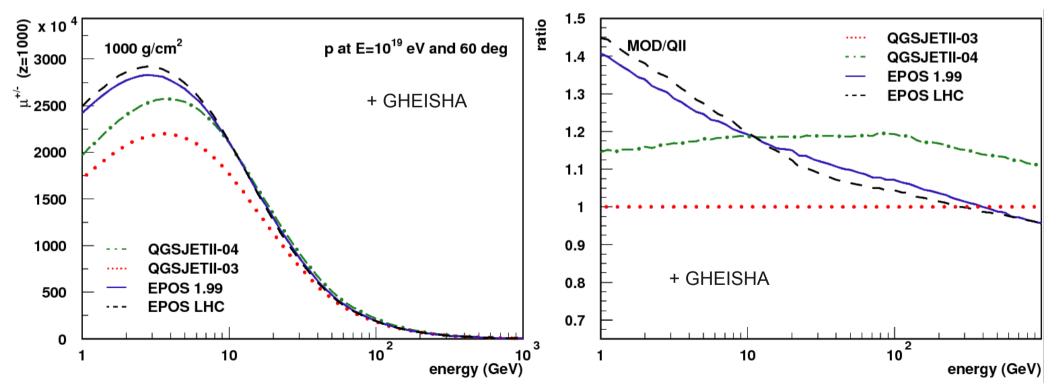


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Muon Energy Spectra

Total number of muons in QGSJETII-04 (@60°) closer to EPOS BUT

- muons with different energy (hadronic energy stored in mesons or baryons ?)
- different zenith angle dependence (attenuation length depends on muon energy spectrum)
- effect of low energy hadronic interaction models (Gheisha, Fluka, UrQMD) ?
 - muon production dominated by last hadronic interaction(s) !

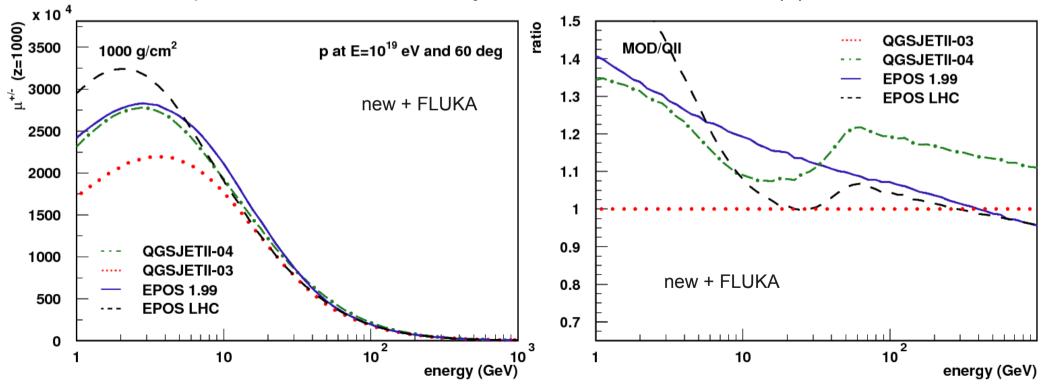


Muon Energy Spectra

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muon production dominated by last hadronic interaction(s) !



Summary

Hadronic interaction models for CR reproduce LHC data in a reasonable way

➡ No change of hadronic physics around the knee (10¹⁵ eV)

Large uncertainties in <X_{max}> simulations due to hadronic models reduced by precise fit of LHC data to the value of the exp. resolution

Low mass composition UNLIKELY at highest energy

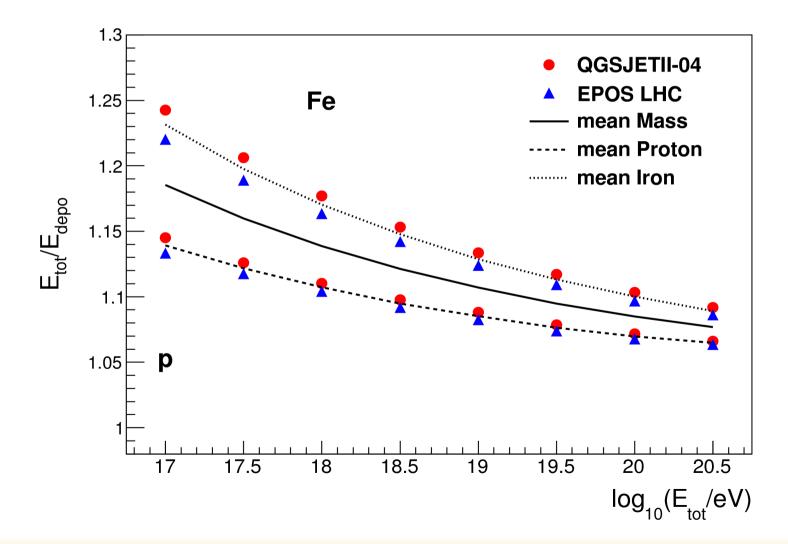
- Muon number converging to high value
 - NA61 will help for precise muon energy spectrum below 100 GeV.
 - LHC energies important for high energy muons
 - Low mass composition LIKELY at lowest energy
- Hadronic interaction models for CR <u>can</u> be re-tuned to LHC data without too many changes
 - Better predictive power than HEP MC models
 - All CR models available with hepMC interface to be compared with LHC !
 - CR models for LHC

Muons in EAS

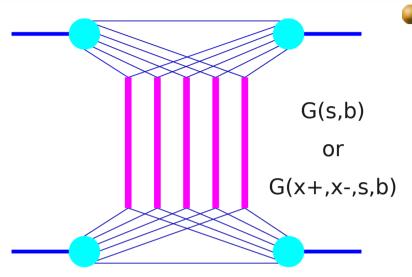
EAS Energy Deposit

Increase of muons in QII04

larger correction factor from missing energy



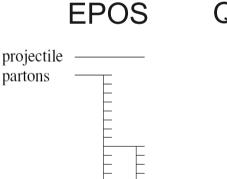
Differences between Models



Gribov-Regge and optical theorem

- Basis of all models (multiple scattering) but
 - Classical approach for QGSJET and SIBYLL (no energy conservation for cross section calculation)
 - Parton based Gribov-Regge theory for EPOS (energy conservation at amplitude level)

pQCD



QGSJET II

>=1

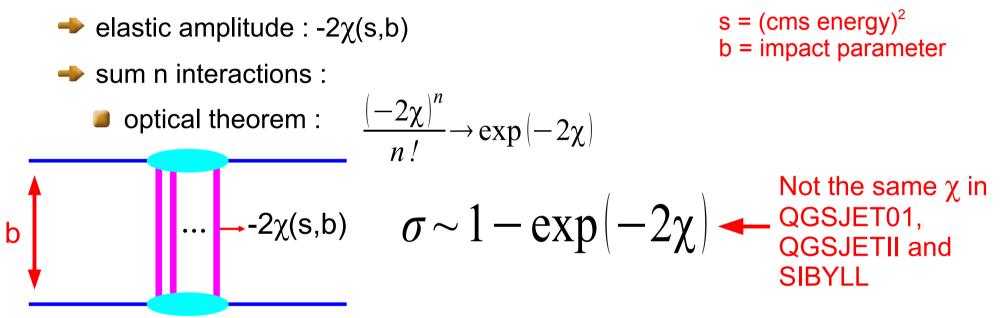
- Minijets with cutoff in SIBYLL
- Same semi-hard Pomeron (DGLAP convoluted with soft part : not cutoff) in QGS and EPOS but
 - No enhanced diagram in Q01 (old PDF)
 - Generalized enhanced diagram in QII
 - Simplified non linear effect in EPOS
 - Phenomenological approach

target

partons

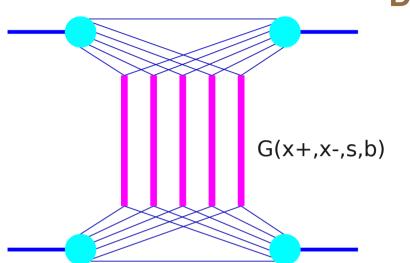
Cross Section Calculation : SIBYLL / QGSJET

Interaction amplitude given by parameterization (soft) or pQCD (hard) and Gribov-Regge for multiple scattering :



- $\rightarrow \chi(s,b)$ parameters for a given model fixed by pp cross-section
- pp to pA or AA cross section from Glauber
- energy conservation not taken into account at this level

Cross Section Calculation : EPOS



Different approach in EPOS :

- Gribov-Regge but with energy sharing at parton level : MPI with energy conservation !
- amplitude parameters fixed from QCD and pp cross section
- cross section calculation take into account interference term

$$\Phi_{\rm pp}\left(x^+, x^-, s, b\right) = \sum_{l=0}^{\infty} \int dx_1^+ dx_1^- \dots dx_l^+ dx_l^- \left\{ \frac{1}{l!} \prod_{\lambda=1}^l -G(x_\lambda^+, x_\lambda^-, s, b) \right\}$$

$$\times F_{\rm proj}\left(x^+ - \sum x_\lambda^+\right) F_{\rm targ}\left(x^- - \sum x_\lambda^-\right).$$

- $\sigma_{\rm ine}(s) = \int d^2b \left(1 \Phi_{\rm pp}(1, 1, s, b)\right) \rightarrow {\rm can not use complex diagram like QII}$ with energy sharing
 - non linear effects taken into account as correction of single amplitude G

Particle Production in SIBYLL and QGSJET

Number n of exchanged elementary interaction per event fixed from elastic amplitude (cross section) :

🔶 n from :

$$P(n) = \frac{(2\chi)^n}{n!} \cdot \exp(-2\chi)$$

- no energy sharing accounted for (interference term)
- → 2n strings formed from the n elementary interactions
 - in QGSJET II, n is increased by the sub-diagrams
 - energy conservation : energy shared between the 2n strings
 - particles from string fragmentation
- inconsistency : energy sharing should be taken into account when fixing n
 - EPOS approach

Particle Production in EPOS

m number of exchanged elementary interaction per event fixed from elastic amplitude taking into account energy sharing :

 \rightarrow m from :

$$\Omega_{AB}^{(s,b)}(m,X^+,X^-) = \prod_{k=1}^{AB} \left\{ \frac{1}{m_k!} \prod_{\mu=1}^{m_k} G(x_{k,\mu}^+, x_{k,\mu}^-, s, b_k) \right\} \Phi_{AB} \left(x^{\text{proj}}, x^{\text{targ}}, s, b \right)$$

m and X fixed together by a complex Metropolis (Markov Chain)

→ 2m strings formed from the m elementary interactions

- energy conservation : energy fraction of the 2m strings given by X
- consistent scheme : energy sharing reduce the probability to have large m
- modified hadronization due to high density effect
 - statistical hadronization instead of string fragmentation

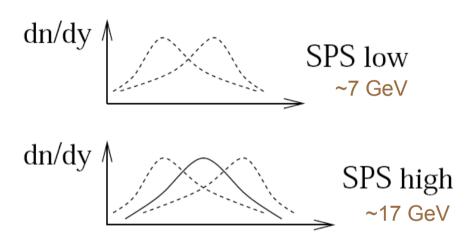
larger Pt (flow)

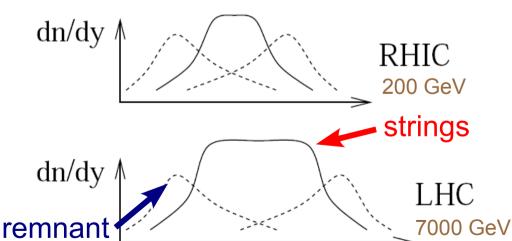
LHC and Xmax

Muons in EAS

Forward Spectra

Forward particles mainly from projectile remnant





The (in)elasticity is closely related to diffraction and forward spectra

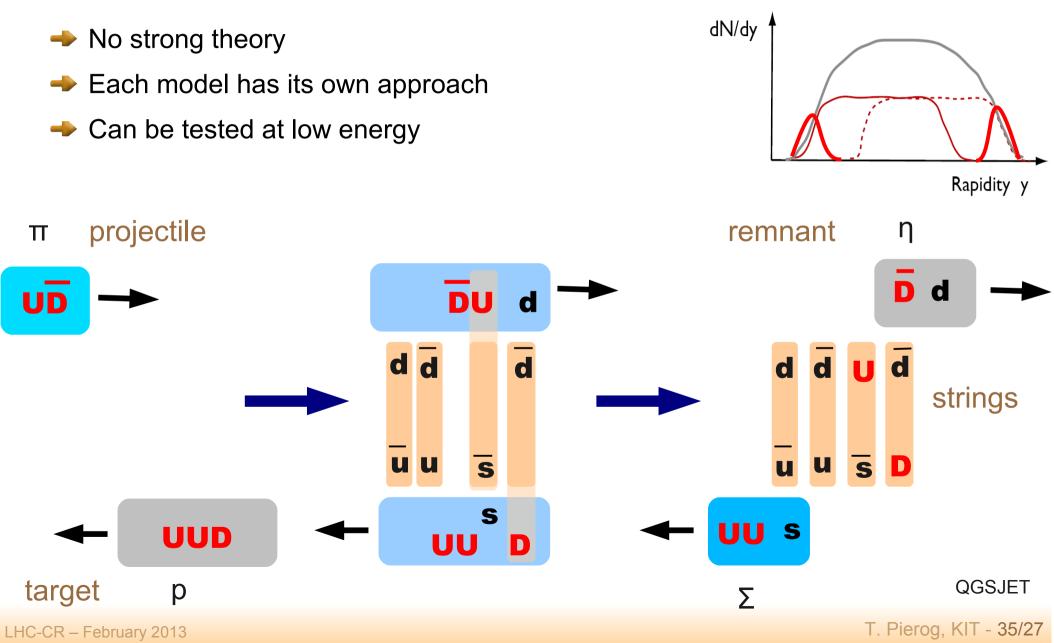
- At very low energy only particles from remnants
- At low energy (fixed target experiments) (SPS) strong mixing
- At intermediate energy (RHIC) mainly string contribution at mid-rapidity with tail of remnants.
- At high energy (LHC) only strings at mid-rapidity (baryon free)

Different contributions of particle production at different energies or rapidities

Muons in EAS

Beam Remnants

Forward particle production dominated by beam remnants

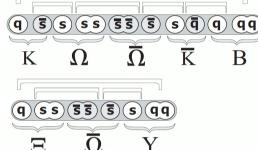


Baryons and Remnants

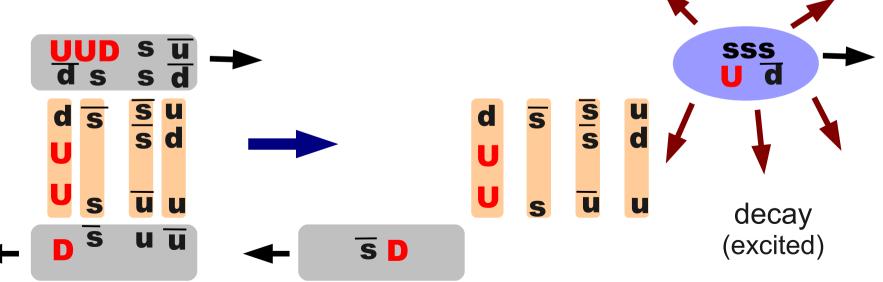
Parton ladder string ends :

Problem of multi-strange baryons at low energy (Bleicher et al., Phys.Rev.Lett.88:202501,2002)

- 2 strings approach :
- $\Rightarrow \Omega / \Omega$ always > 1
- But data < 1 (Na49)</p>
- EPOS

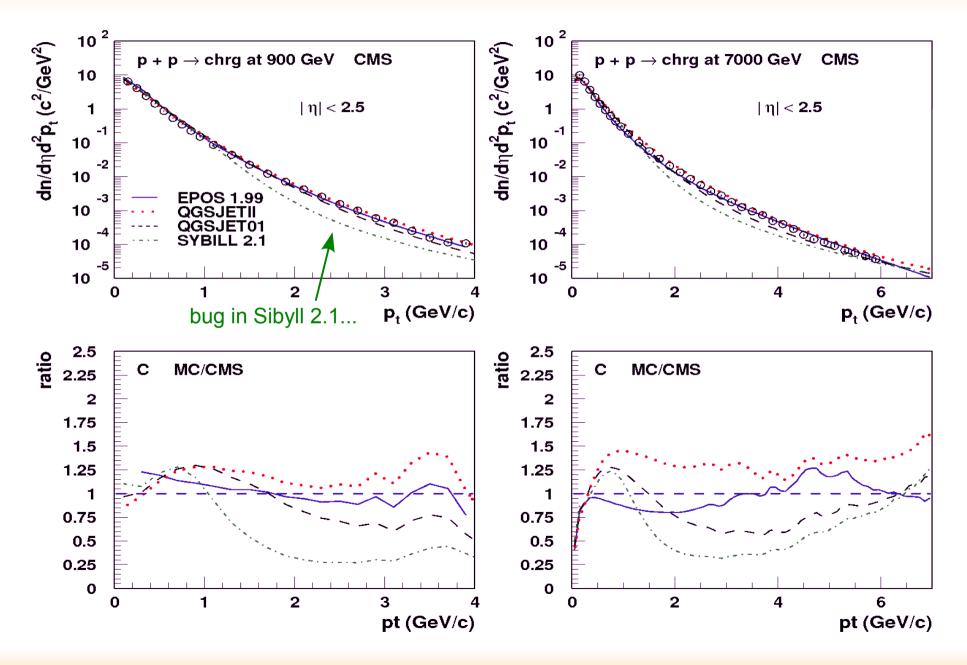


- No "first string" with valence quarks : $\overline{\overline{\Omega}}$ strings equivalent
- Wide range of excited remnants (from light resonances to heavy quark-bag)
- $\Rightarrow \Omega / \Omega$ always < 1 _



Muons in EAS

Pt @ LHC



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Muons in EAS

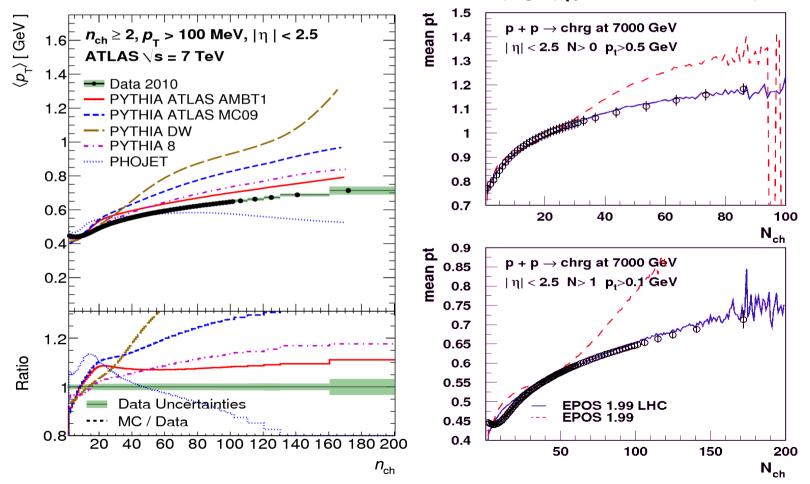
EPOS LHC

Detailed description can be achieved

better than HEP MC used by LHC collaborations

can be used as min bias generator at LHC

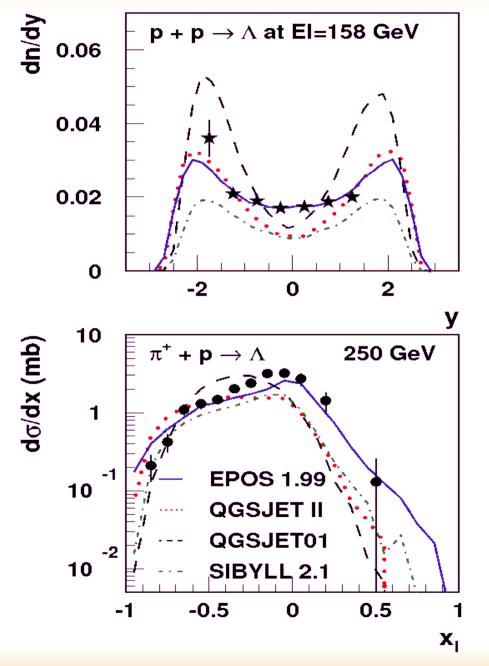
not suitable for rare events (high pt jets or electroweak)



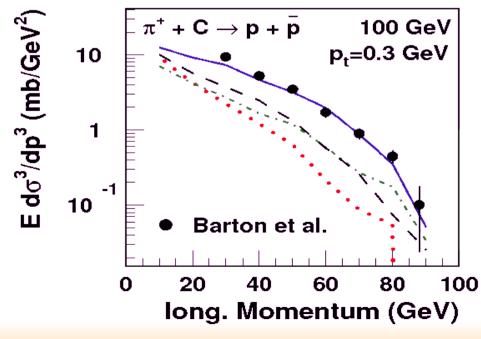
EAS

Muons in EAS

Baryon Forward Spectra



- Large differences between models
- Need a new remnant approach for a complete description (EPOS)
- Problems even at low energy
- No measurement at high energy !



EAS

Basic Observables

Pseudorapidity

 emission angle of a particle from interaction point ("mid-rapidity" : η=0) :

$$\eta = -\ln\left[\tan\left(\frac{\theta}{2}\right)\right]$$
 $\eta = \frac{1}{2}\ln\left(\frac{|\mathbf{p}| + p_{\mathrm{L}}}{|\mathbf{p}| - p_{\mathrm{L}}}\right)$

when the mass of the particle is known the rapidity is used :

$$y = \frac{1}{2} \ln \left(\frac{E + p_{\rm L}}{E - p_{\rm L}} \right)$$

 for EAS development, "forward" particles (with large η) are most important

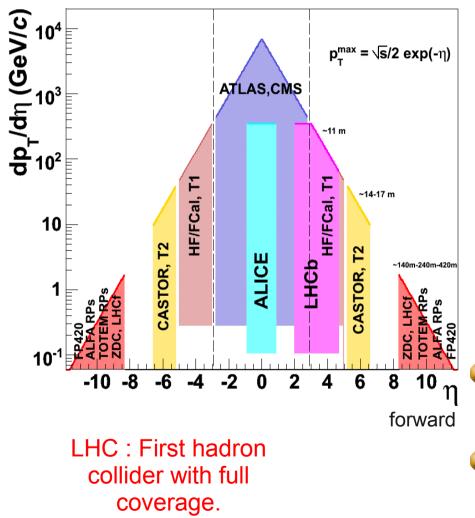
Transverse momentum

$$p_t = \sqrt{p_x^2 + p_y^2}$$

Multiplicity

 number of particles in a given η and p_t range
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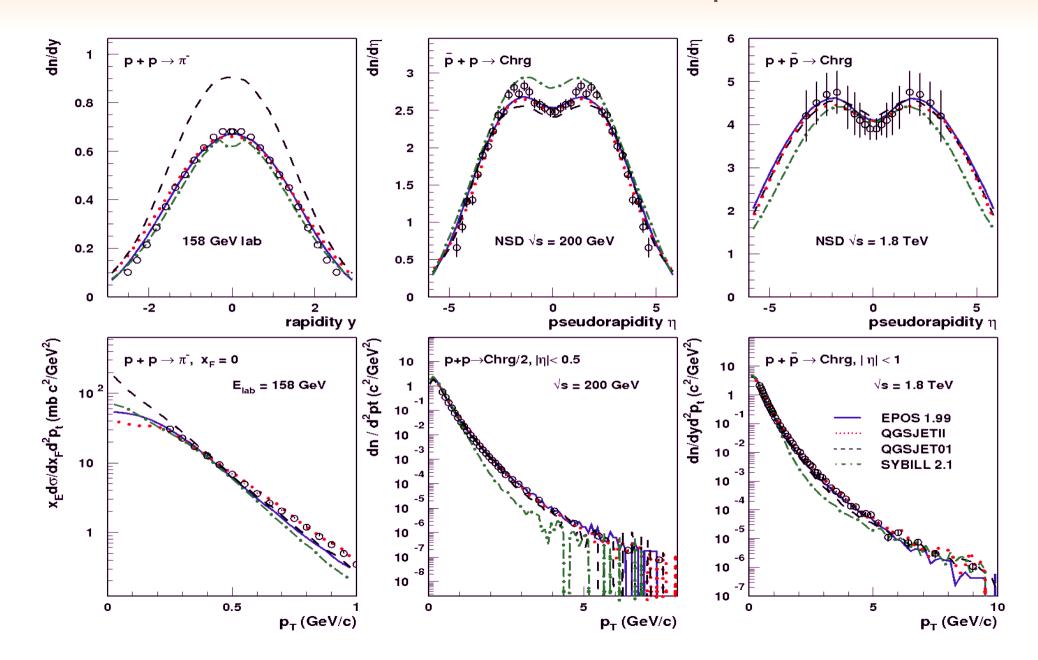
mid-rapidity



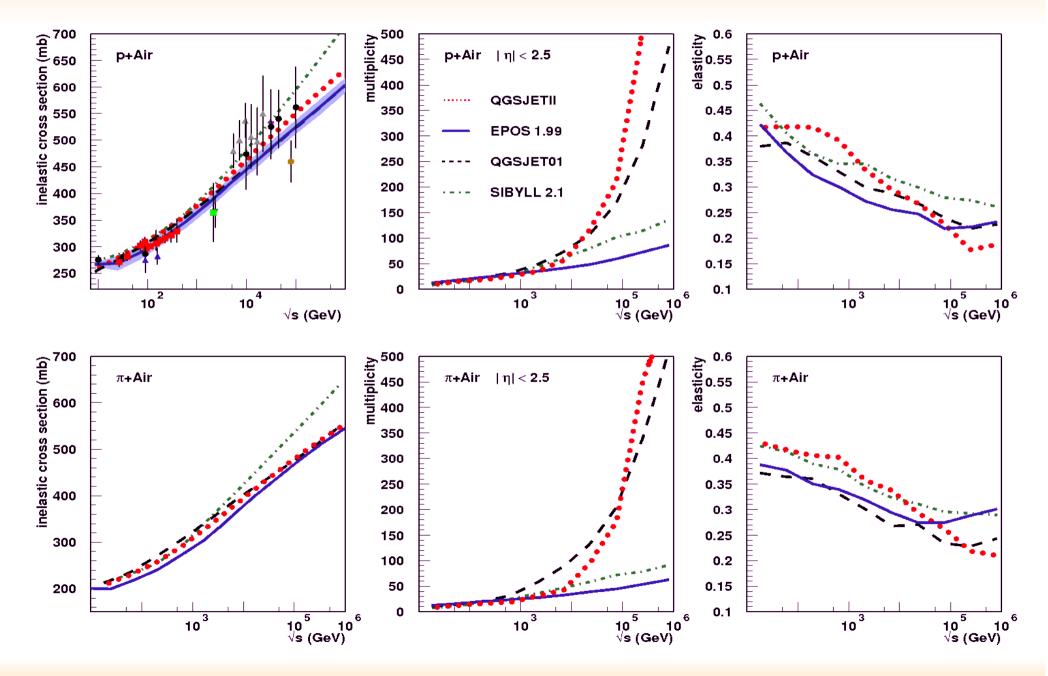
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Muons in EAS

Pseudorapidity and p_T

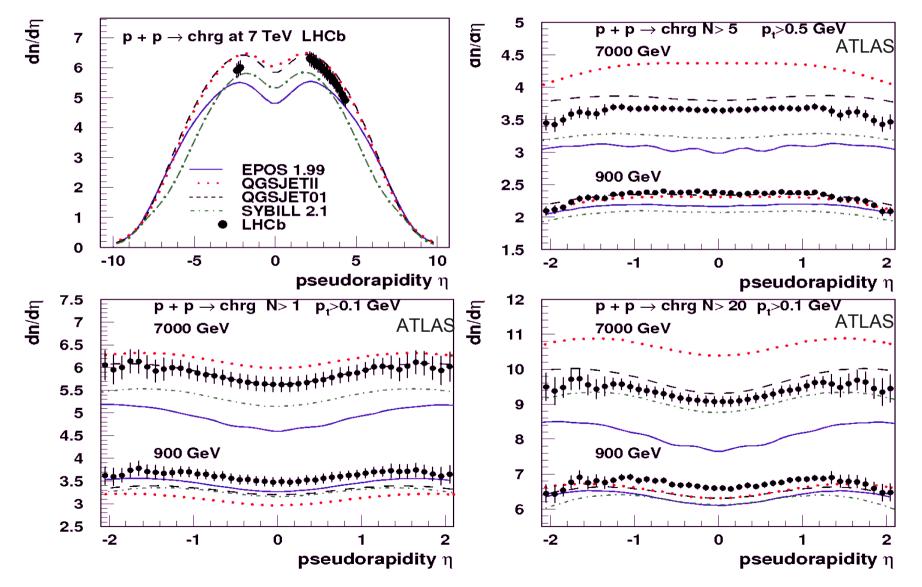


Ultra-High Energy Hadronic Model Predictions



Pseudorapidity Distributions

No model with perfect prediction : but data well bracketed



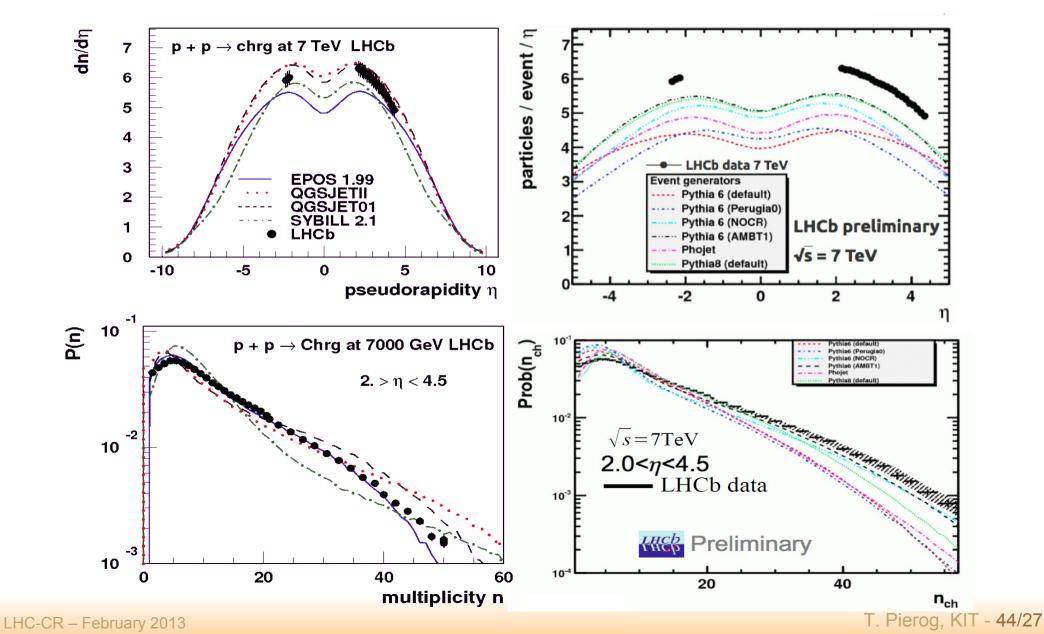
Predictions ! ... newest model released in march 2009

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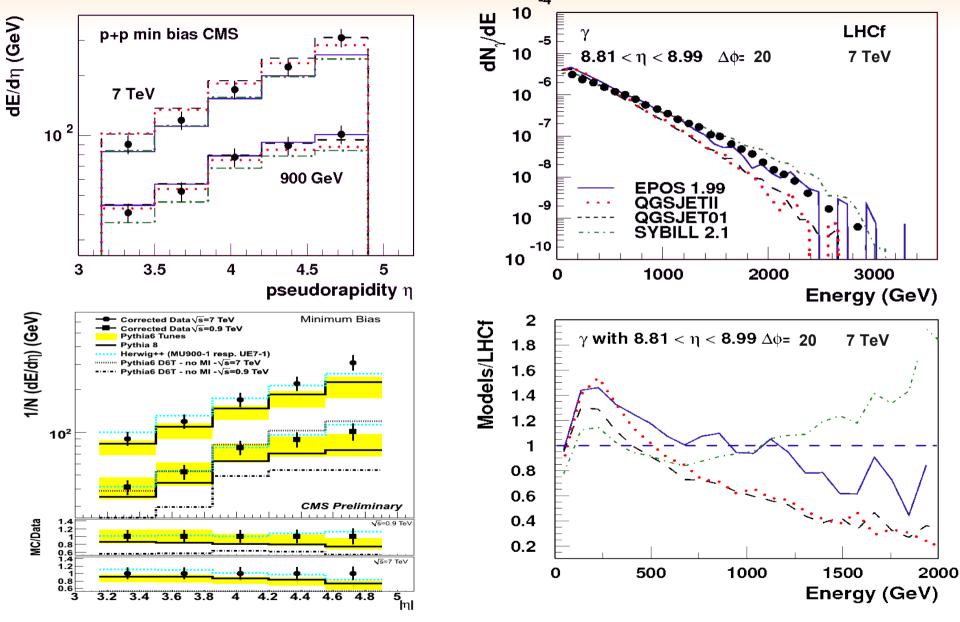
T. Pierog, KIT - 43/27

Pseudorapidity Distributions

No model with perfect prediction : but better than HEP MC

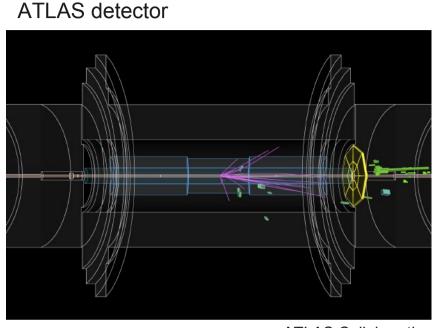


CMS and LHCf Forward Spectra



■ Forward calorimeter → better than HEP models

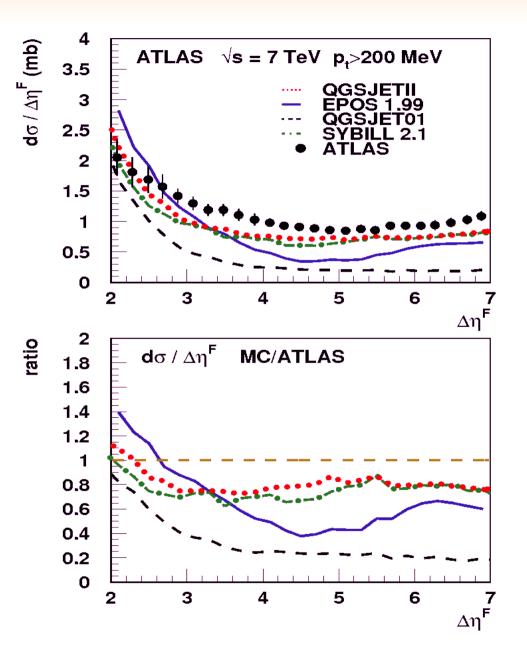
Rapidity Gap



ATLAS Collaboration

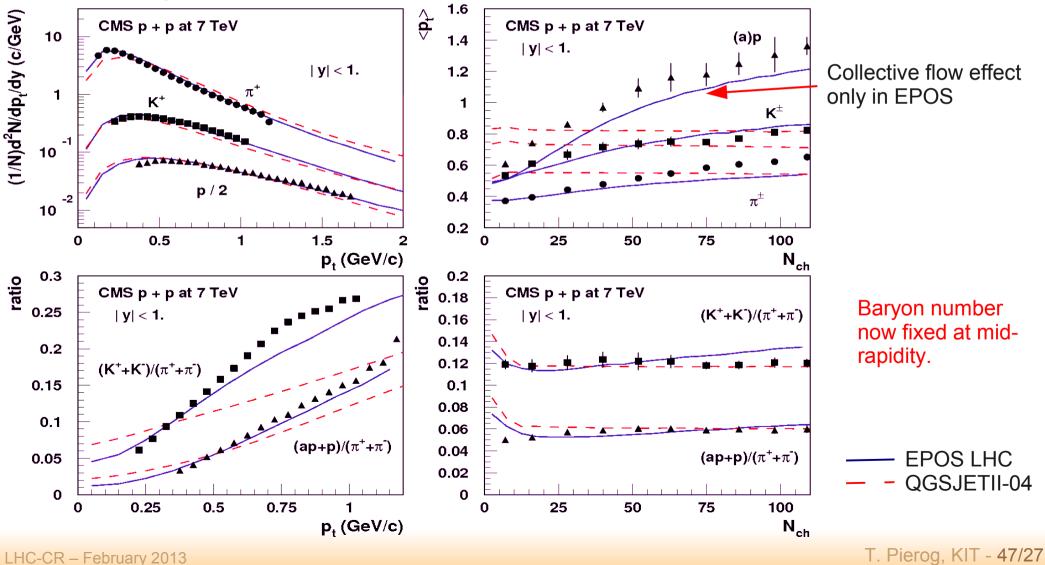
Rapidity gap closely related to diffraction

- diffractive cross-section
- AND diffractive mass distribution
- Hard constraint for CR
 - change elasticity



Identified Particle Spectra

- Detailed description can be achieved (tested by ATLAS for publications)
 - identified spectra
 - p_t behavior driven by collective effects (statistical hadronization + flow)



Particle Spectra

