



Inclusive muon measurement: atmospheric muon charge ratio with the OPERA detector

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Results and prospects of forward physics at the LHC

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Atmospheric muon charge ratio

- The atmospheric muon charge ratio $R_\mu \equiv N_{\mu^+}/N_{\mu^-}$ is being studied and measured since many decades
 - Depends on the **chemical composition** and energy spectrum of the primary cosmic rays
 - Depends on the **hadronic interaction features**
 - At high energy, depends on the **prompt component**
- It provides the possibility to check HE hadronic interaction models ($E > \text{TeV}$) in the **fragmentation region**, where no data exist
- Since atmospheric muons are kinematically related to atmospheric neutrinos (same sources), R_μ provides a benchmark for **atmospheric ν flux computations** (e.g. background for neutrino telescopes)

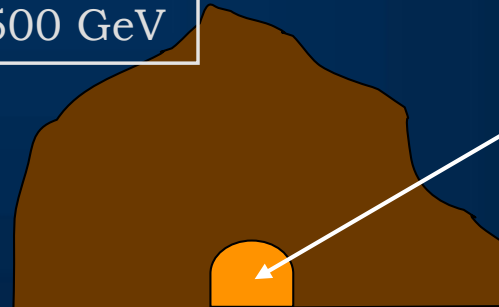


Atmospheric muon charge ratio

The atmospheric muon charge ratio R_μ mainly depends on:

- 1) the composition of primary cosmic rays (n/p ratio in the all-nucleon spectrum)
- 2) the hadronic interactions in atmosphere (π^+/π^- and K^+/K^- ratios)

$\langle E_\mu \rangle$ underground: ~ 270 GeV
 $\langle E_\mu \rangle$ surface cut-off: ~ 1500 GeV



μ

Underground experiments measure the charge ratio in a high energy region: muons are energy-selected with the overburden

About 1): OPERA can provide R_μ separately for single muons and multiple muon bundles, testing the hypothesis of “dilution” of charge excess in neutron-enriched primaries

About 2): The underground charge ratio measurement gives useful information on K production in the TeV fragmentation region, where there are no data from accelerators
→ this measure can provide constraints on theoretical hadronic models, in particular for fluxes of atmospheric TeV ν 's (dominated by K production)

OPERA can measure the atmospheric muon charge ratio in the TeV energy range, where the kaon contribution becomes significant



Analytic predictions

Naïve prediction

- Assume only primary protons with a spectrum $dN/dE = N_0 E^{-(1+\gamma)}$
- Assume only pions and neglect muon decays (HE limit)
- Consider the inclusive cross-section for pions

$$f_{p\pi}^{\pm}(E_{\pi}, E_p) \equiv \frac{E_{\pi}}{\sigma_{pp}^{inel}} \frac{d\sigma_{p \rightarrow \pi}^{\pm}}{dE_{\pi}}$$

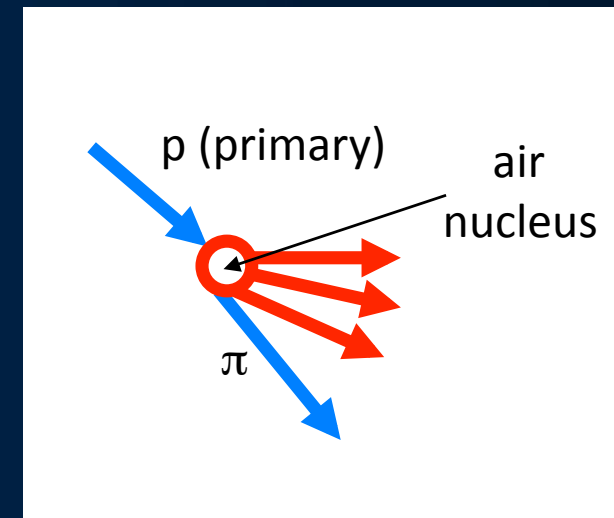
The pion spectrum is then

$$\pi(E_{\pi}) = \frac{(const)}{E_{\pi}} \int_{E_{\pi}}^{\infty} dE E^{-(1+\gamma)} f_{p\pi}^{\pm}(E_{\pi}, E_p)$$

Assuming Feynman scaling

$$f_{p\pi}^{\pm}(E_{\pi}, E_p) \equiv \frac{E_{\pi}}{\sigma_{pp}^{inel}} \frac{d\sigma_{p \rightarrow \pi}^{\pm}}{dE_{\pi}} \xrightarrow{E \rightarrow \infty} \tilde{f}_{p\pi}^{\pm}(x)$$

Feynman
scaling





Analytic predictions

The pion spectrum becomes:

$$\pi^\pm(E_\pi) = (\text{const}) E_\pi^{-(1+\gamma)} Z_{p\pi^\pm}$$

$$Z_{p\pi^\pm} \equiv \int_0^1 \tilde{f}_{p\pi^\pm}^\pm(x) x^{\gamma-1} dx$$

Spectrum weighted moments

Finally, the muon charge ratio prediction:

$$R_\mu = \frac{\mu^+(E_\mu)}{\mu^-(E_\mu)} = \frac{\pi^+(E_\pi)}{\pi^-(E_\pi)} = \frac{Z_{p\pi^+}}{Z_{p\pi^-}}$$

Interpretation of the prominent features:

- The result is valid only in the fragmentation region, since in the central region Feynman scaling is strongly violated
- But the steeply falling primary spectrum ($\gamma \sim 1.7$) in the SWM suppresses the contribution of the central region in the secondary production \rightarrow scaling holds (at least for $E < 1$ TeV);

In other words: each pion is likely to have an energy close to the one of the projectile (primary CR proton) and comes from its fragmentation (valence quarks) \rightarrow positive charge

- R_μ does not depend on E_π (or E_μ) nor on the target nature
- R_μ depends on the primary spectrum γ



Kaon contribution

- At higher energy (>100 GeV) the contribution of K becomes important
- In general, the contribution of each component to the muon flux $N_{\text{par}} = (\pi, K, \text{ charmed, etc.})$ depends on the relative contribution of decays and interaction probabilities:

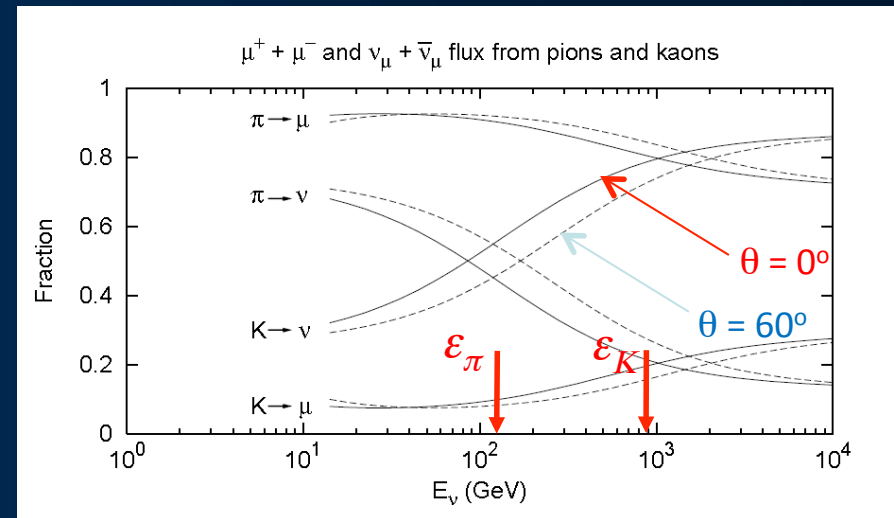
$$\Phi_{\mu} = \frac{\Phi_N(E_{\mu})}{1 - Z_{NN}} \sum_{i=1}^{N_{\text{par}}} \frac{a_i Z_{Ni}}{1 + b_i E_{\mu} / \varepsilon_i(\theta)}$$

where

$$a_i = a_i(\gamma) = \frac{(1 - r_i^{\gamma+1}) Br(i \rightarrow \mu)}{(1 - r_i)(\gamma + 1)}$$

$$b_i = b_i(\gamma) = \frac{(\gamma + 2)(1 - r_i^{\gamma+1})(\lambda_i - \lambda_N)}{(\gamma + 1)(1 - r_i^{\gamma+2})(\lambda_i \log \lambda_i / \lambda_N)}$$

$$r_i = (m_{\mu} / m_i)^2$$



$\varepsilon_i = \varepsilon_i(\theta)$ is the “critical energy”, i.e. the energy above which interactions dominate over decays. Along the vertical ($\theta = 0^\circ$) $\varepsilon_i(0) = m_i c h / \tau_i$ ($h = 6.5 \text{ km}$)

$$\begin{aligned} \varepsilon_{\pi} &= 115 \text{ GeV} \\ \varepsilon_K &= 850 \text{ GeV} \\ \varepsilon_X &> 10^7 \text{ GeV} \end{aligned}$$



Kaon contribution to R_μ

For kaons:

$$Z_{pK^+} \gg Z_{nK^-} \approx Z_{pK^-}$$

because the reaction

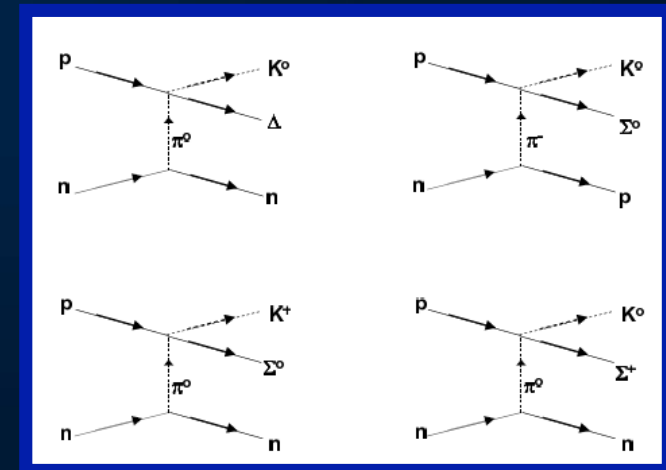
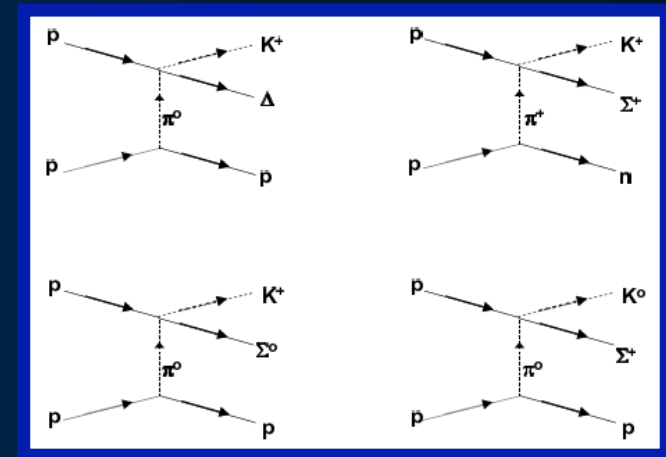


is favoured.

Conservation of strangeness requires production of K^+ and K^0 in association with a Λ or Σ .

On the other hand, the production of K^-, \bar{K}^0 requires the creation of an associated baryon and an additional strange meson.

This leads to a larger R_μ ratio at high energy





General form for R_μ

- Let us consider again the general form for the muon flux

$$\Phi_{\mu^\pm} = \frac{\Phi_N(\mathbf{E}_\mu)}{1 - Z_{NN}} \sum_{i=1}^{N_{par}} \frac{a_i Z_{Ni}^\pm}{1 + b_i \mathbf{E}_\mu \cos \theta^* / \varepsilon_i(0)}$$

where we have explicited the $\varepsilon_i(\theta)$ dependence on θ

$$\varepsilon_i(\theta) = \varepsilon_i(0) / \cos \theta^*$$

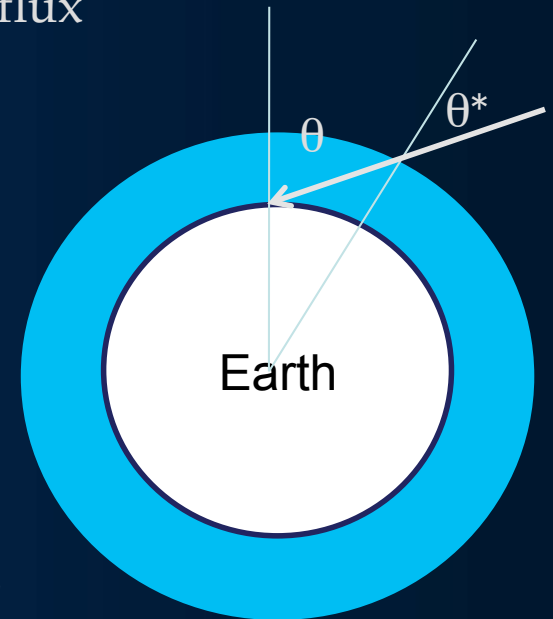
where θ^* is the zenith angle at the production point

- The correct variable to describe the evolution of R_μ is therefore $\mathbf{E}_\mu \cos \theta^*$
- The R_μ evolution as a function of $\mathbf{E}_\mu \cos \theta^*$ spans over the different sources

$$R_\mu = w_\pi R_\mu^\pi + w_K R_\mu^K + w_{charm} R_\mu^{charm} + \dots \longrightarrow \text{POWERFUL HANDLE TO DISCRIMINATE MODELS}$$

Measuring R_μ as a function of $\mathbf{E}_\mu \cos \theta^*$ we can infer the parameters $R_\pi = Z_{N\pi^+} / Z_{N\pi^-}$ and $R_K = Z_{NK^+} / Z_{NK^-}$.

OPERA: $\langle \mathbf{E}_\mu \cos \theta^* \rangle \approx 2000 \text{ GeV}$ \longrightarrow The (magnetized) experiment with the largest $\mathbf{E}_\mu \cos \theta^*$





R_μ measurements with $E_\mu \cos \theta^* > 1 \text{ TeV}$

- Utah:
G. K. Ashley et al., Phys. Rev. D12 (1975) 20
 - Underground at Utah University, flat surface above ~ 1400 m.w.e., magnetic spectrometer (1.63 T) + spark chambers, six bins with $46^\circ < \theta < 78^\circ$
- Kamiokande-II:
M. Yamada et al., Phys. Rev. D44 (1991) 617
 - Underground Cherenkov detector at Kamioka ~ 2700 m.w.e., delayed events on stopping muons, one bin with $0^\circ < \theta < 90^\circ$
- MINOS:
P. Adamson et al., Phys. Rev. D76 (2007) 052003 + arXiv:0906.3726
 - Underground at Soudan, magnetized steel, flat surface above ~ 2000 m.w.e., $0^\circ < \theta < 90^\circ$
- LVD:
N. Agafonova et al., Proc. 31th ICRC, ŁÓDZ 2009
 - Underground at LNGS, average overburden ~ 3800 m.w.e., scintillators, delayed events on stopping muons, one bin with $\theta < 15^\circ$
- OPERA:
N. Agafonova et al., Eur. Phys. J. C67 (2010) 25
 - Underground magnetic spectrometer (1.53 T) at LNGS, average overburden ~ 3800 m.w.e., drift tubes + RPC + scintillators, $0^\circ < \theta < 90^\circ$



The OPERA detector

Brick: ECC target basic unit
(57 nuclear emulsion films
+ 56 lead plates)

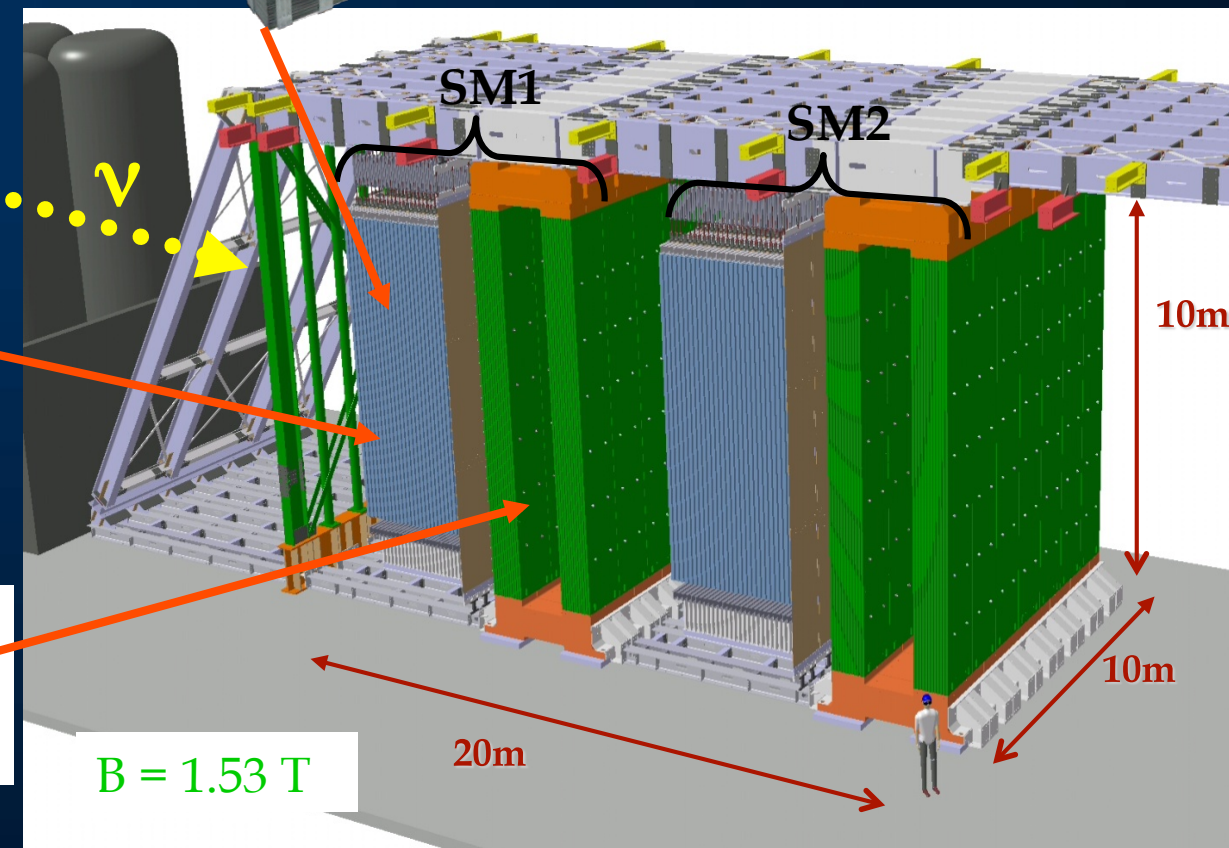


Target section:
27 brick walls (75000 bricks)
31 Target Tracker walls (TT)

Neutrino interaction trigger
Brick selection
Calorimetry

Magnetic spectrometer:
22 RPC planes
6 drift tube layers (PT stations)

μ ID, charge, momentum



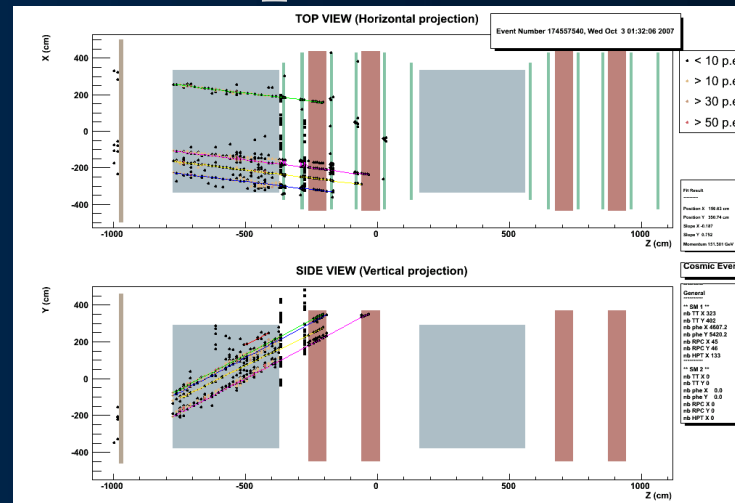
$B = 1.53 \text{ T}$

Total target mass = 1.25 ktons



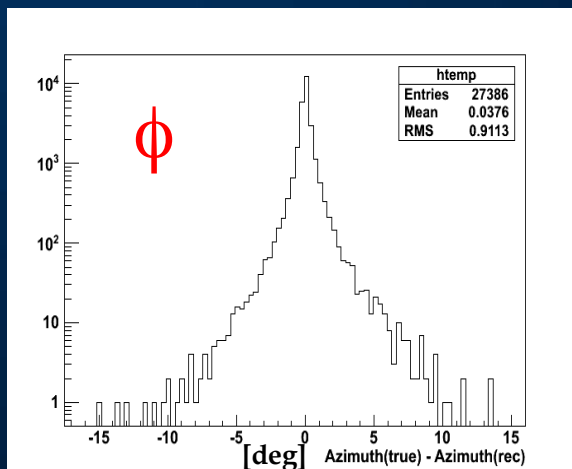
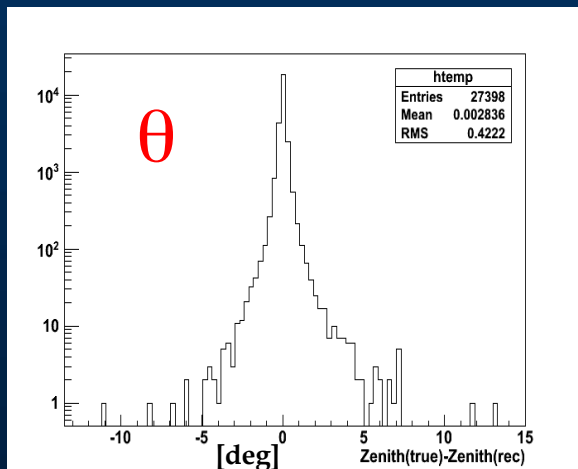
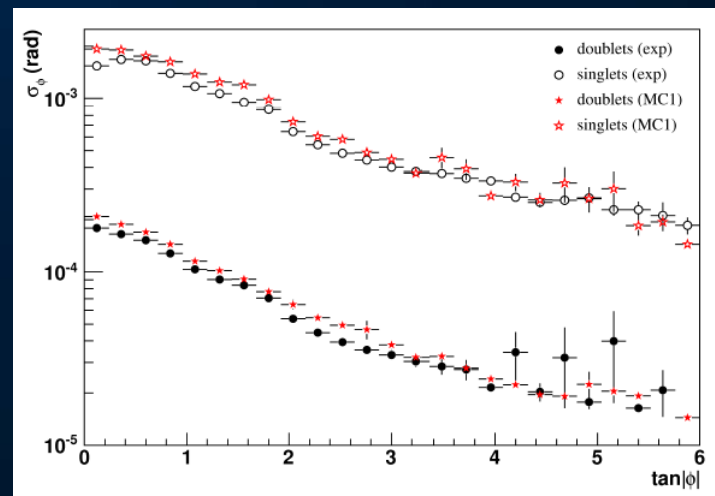
Cosmic event reconstruction performance

➤ Multiple muon events well reconstructed



➤ Good overall angular resolution

“resolutions” < 1 deg both for zenith and azimuth direction reconstruction



➤ High angular resolution in the PT system



New OPERA data

Major improvements w.r.t. 2010 published analysis (Eur. Phys. J. C67 (2010) 25)

- **Statistics $\times 3$:**
1454057 μ collected during
2008-2009-2010 CNGS runs corresponding to 407.1 days of livetime
- Underground $R_\mu = 1.403 \pm 0.008$ (stat) + 0.017 - 0.015 (syst)
- Separation of R_μ for single and multiple muons, different at 7.2σ level

N_μ	$\langle A \rangle$	$\langle E/A \rangle_{\text{primary}}$ [TeV]	H fraction	N_p/N_n	R_μ^{unf}
= 1	3.35 ± 0.09	19.4 ± 0.1	0.667 ± 0.007	4.99 ± 0.05	1.403 ± 0.008
> 1	8.5 ± 0.3	77 ± 1	0.352 ± 0.012	2.09 ± 0.07	1.18 ± 0.03

- Convolution of two effects:
larger n/p ratio in the all-nucleon spectrum \otimes different X_F region
- Cosmic events selected outside the CNGS spill window
- Atmospheric neutrino events rejected on the basis of ToF
- Analysis cuts and details in:
http://www.bo.infn.it/opera/docs/phd_thesis-BO-2011_05_20-mauri.pdf

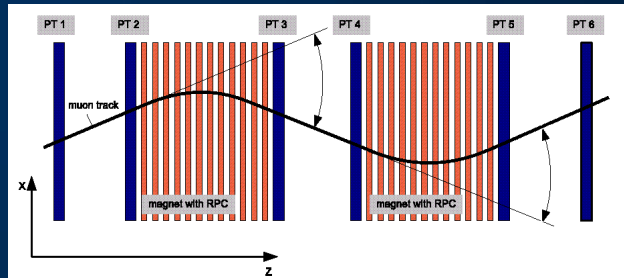


Systematic error estimate

Two main sources for the systematic error

- PT station misalignment: Two-arm test

$$\delta(\Delta\phi) = \Delta\phi_1 - \Delta\phi_2 = 0$$



- computation of the difference between $\Delta\phi$ angles for tracks crossing both magnet arms in each spectrometer

Propagating the value $\delta(\Delta\phi) = 0.08$ mrad in the R_μ computation

$$\delta R_\mu = 0.015$$

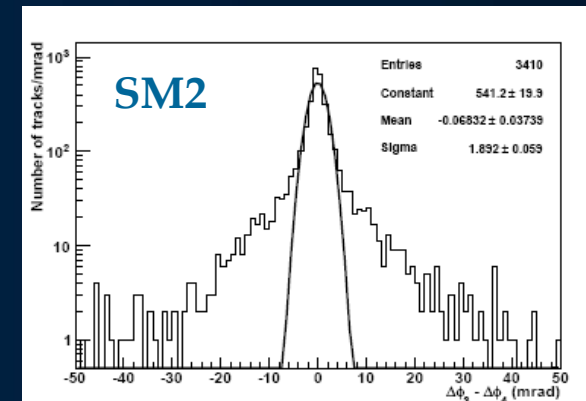
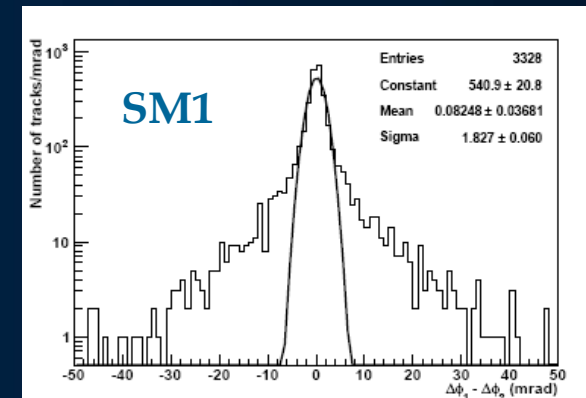
- Misidentification probability: only **from experimental data**

- for tracks crossing both magnet arms (subsample of events), computation of the probability of reconstructing opposite charges \rightarrow we find η_{data}

The difference $\delta\eta = \eta_{data} - \eta_{MC} = 0.007$ is quoted as systematic error on $\eta \rightarrow$ propagating in R_μ

$$\delta R_\mu = 0.007$$

only positive contribution! ($\eta_{real} \geq \eta_{MC}$)





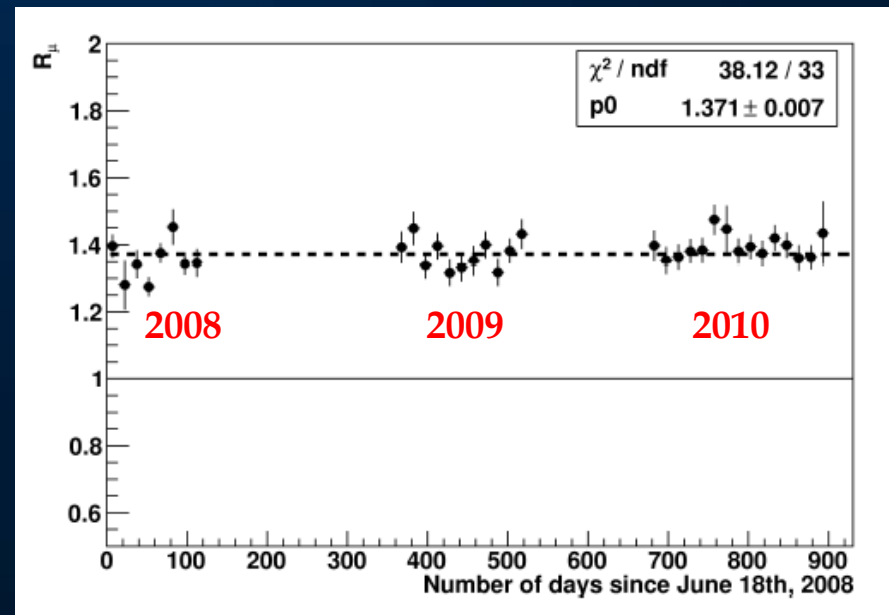
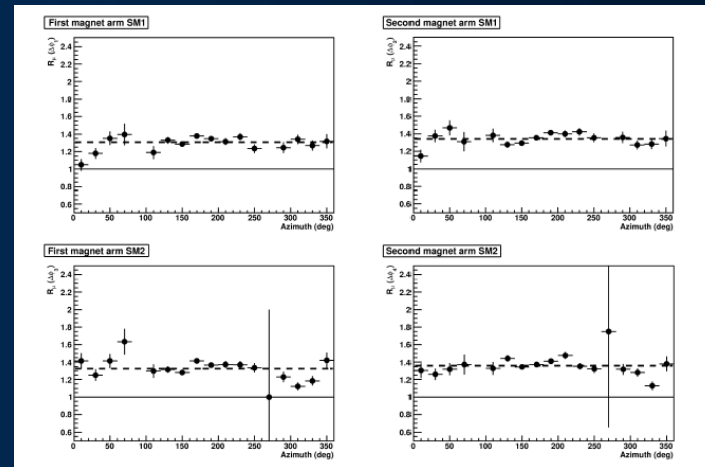
Consistency checks



The total systematic uncertainty is $\delta R_\mu^{\text{unf}}(\text{syst}) = \begin{matrix} +0.017 \\ -0.015 \end{matrix}$

Consistency checks:

- R_μ as a function of the azimuth angle φ (mis-alignment check) \longrightarrow
- The 4 R_μ values, computed separately for each magnet arm, fluctuate around the average of 0.016, which is within their statistical accuracy (0.018)
- Run with **inverted polarity** (9+7 days): $R_\mu^{\text{inverted}} = 1.39 \pm 0.04$
- Run with **magnet off** (13 days): $R_\mu^{\text{off}} = 1.02 \pm 0.04$
- Measurement stability as a function of the **data taking**: R_μ remains constant \longrightarrow





R_μ as a function of $E_\mu \cos\theta^*$

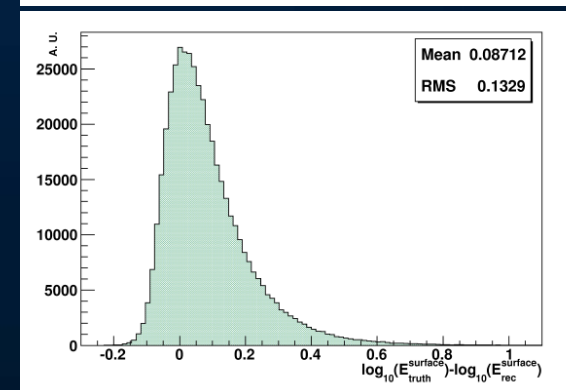
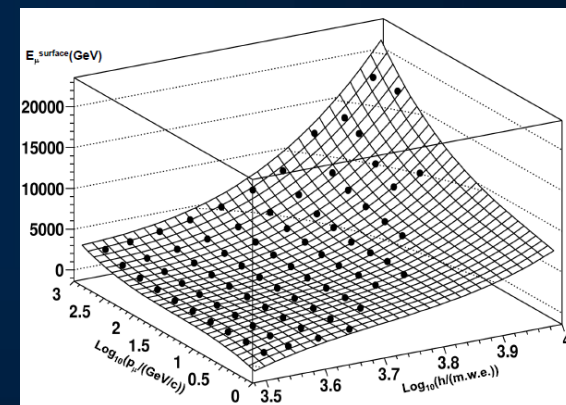
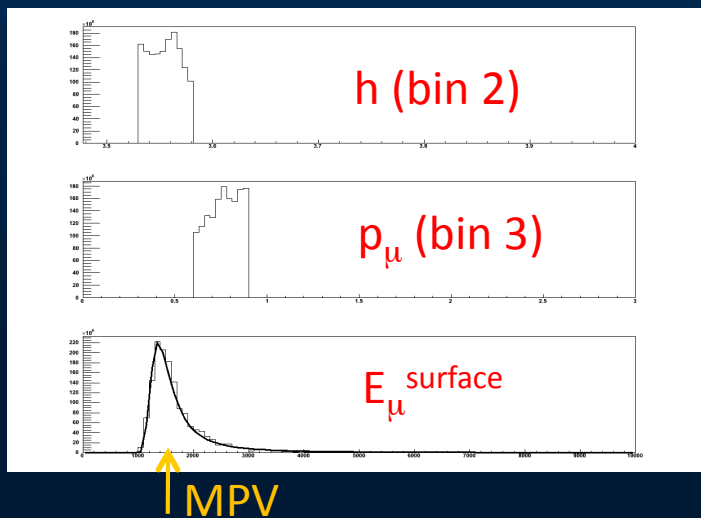
The resolution on E_μ is dominated by fluctuations in the stochastic term β of the energy loss

$$E_\mu = (E_\mu + \alpha / \beta)e^{\beta h} - \alpha / \beta$$

$$\alpha = \alpha(E)$$

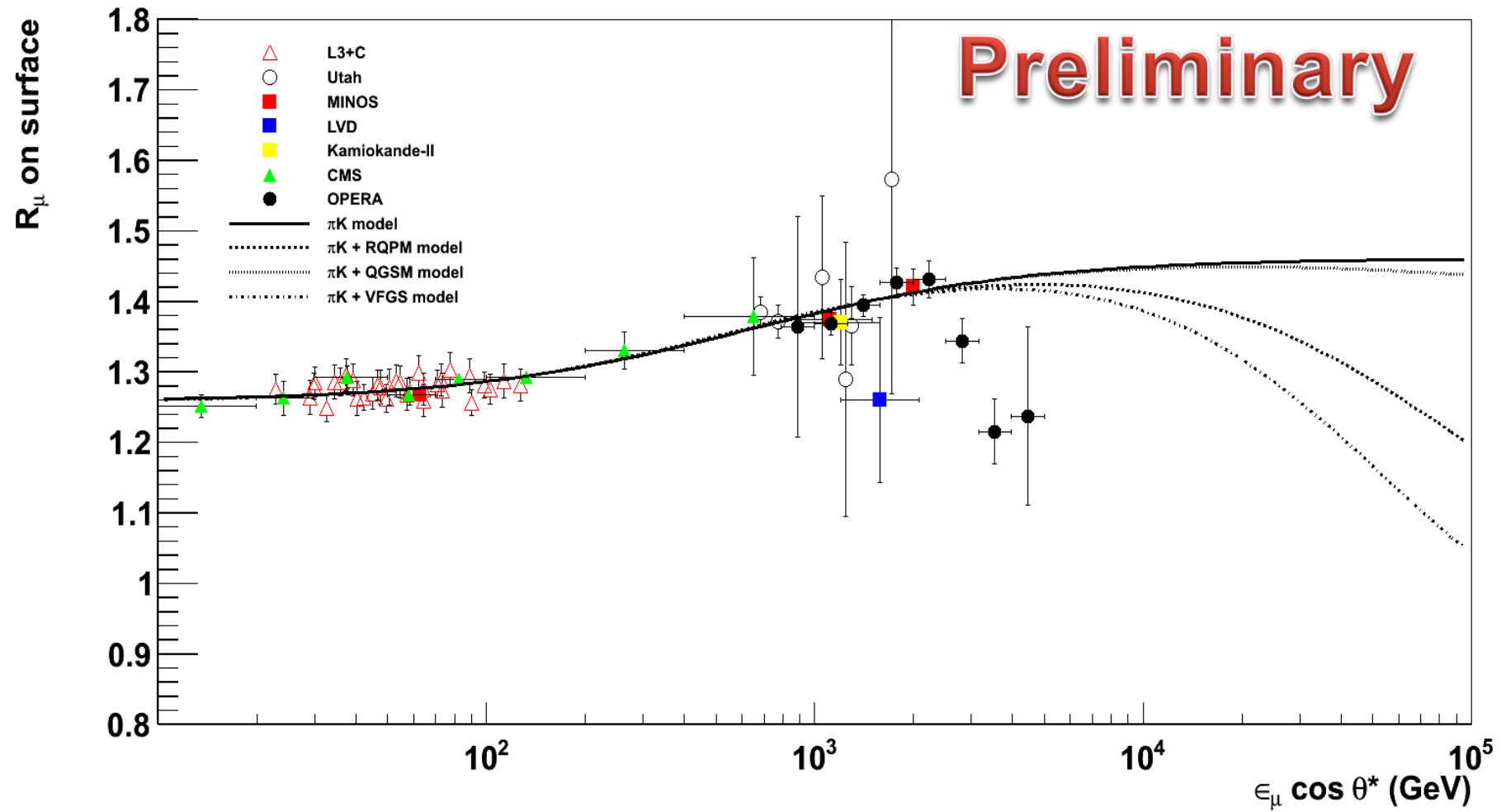
$$\beta = \beta(E)$$

→ Build a table $\langle E_\mu \rangle_{\text{MPV}} = f(h, p_\mu)$ using a full MC simulation based on FLUKA: take the MPV of the Landau distribution → better resolution and residuals well centered





R_μ measurements vs $E_\mu \cos\theta^*$

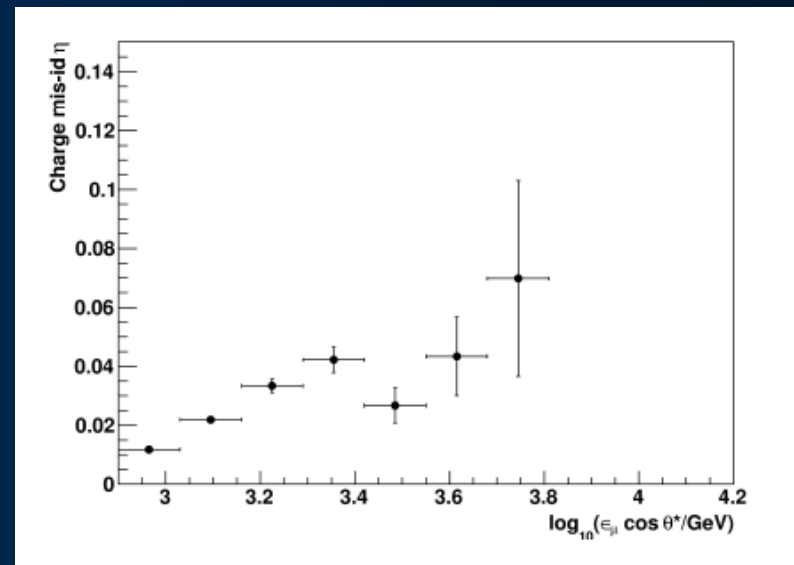
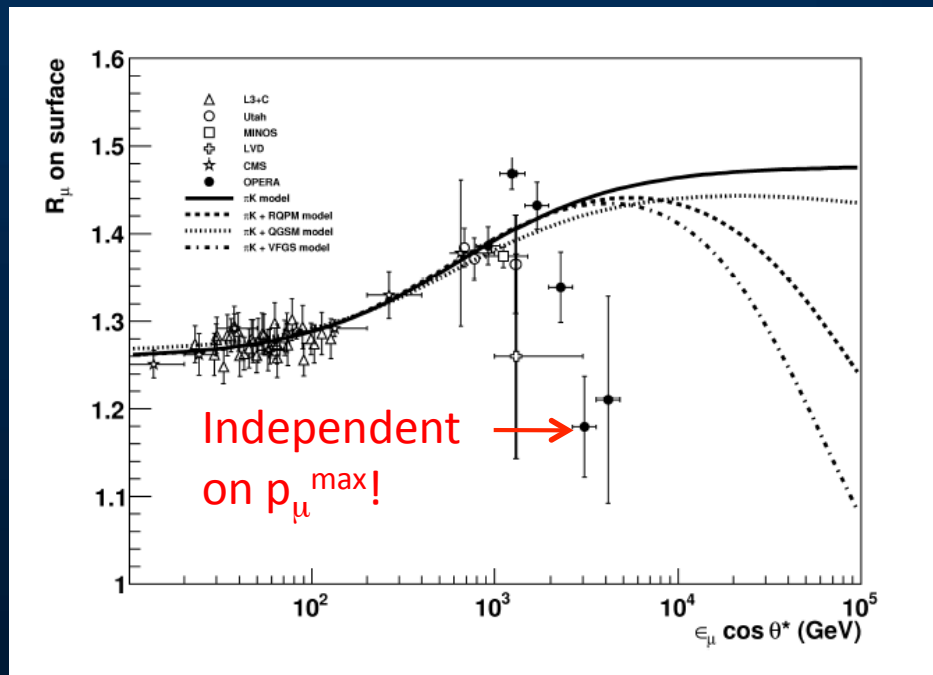




R_μ vs $E_\mu \cos\theta^*$: systematics check

Due to the increased statistics, it is now possible to use only *doublet* data, for which the charge mis-identification can be extracted directly from data (\rightarrow 2-arm test)

only single muons



Check two sources of systematics:

- **unfolding** based on charge misidentification extracted from MC \rightarrow here extracted from data
- **misalignment** of the PT system \rightarrow cut on max μ momentum underground



OPERA measurement result

- There is a clear indication of a smooth transition between the low and high energy regimes as predicted by the π -K model
 - ✓ $R_\mu \approx 1.25$ around 100 GeV \rightarrow
smooth transition to $R_\mu \approx 1.4$ in TeV region
- Sharp drop for $E_\mu \cos\theta^* > 2-3$ TeV in the OPERA data
- In 2012 we collected one year statistics with **inverted magnet polarity**, in order to check the alignment systematics
 - ✓ We are analyzing the new data
 - ✓ We plan to publish in the near future the OPERA measurement with the whole statistics (2008-2012)
- If this drop is not explained by any systematic effect we can speculate on its possible physical nature:
 - Sudden change in the n/p ratio in the all-nucleon spectrum?
 - Strong Feynman scaling violation in the forward fragmentation region for π /K mesons?
 - Upraise of a new muon component?



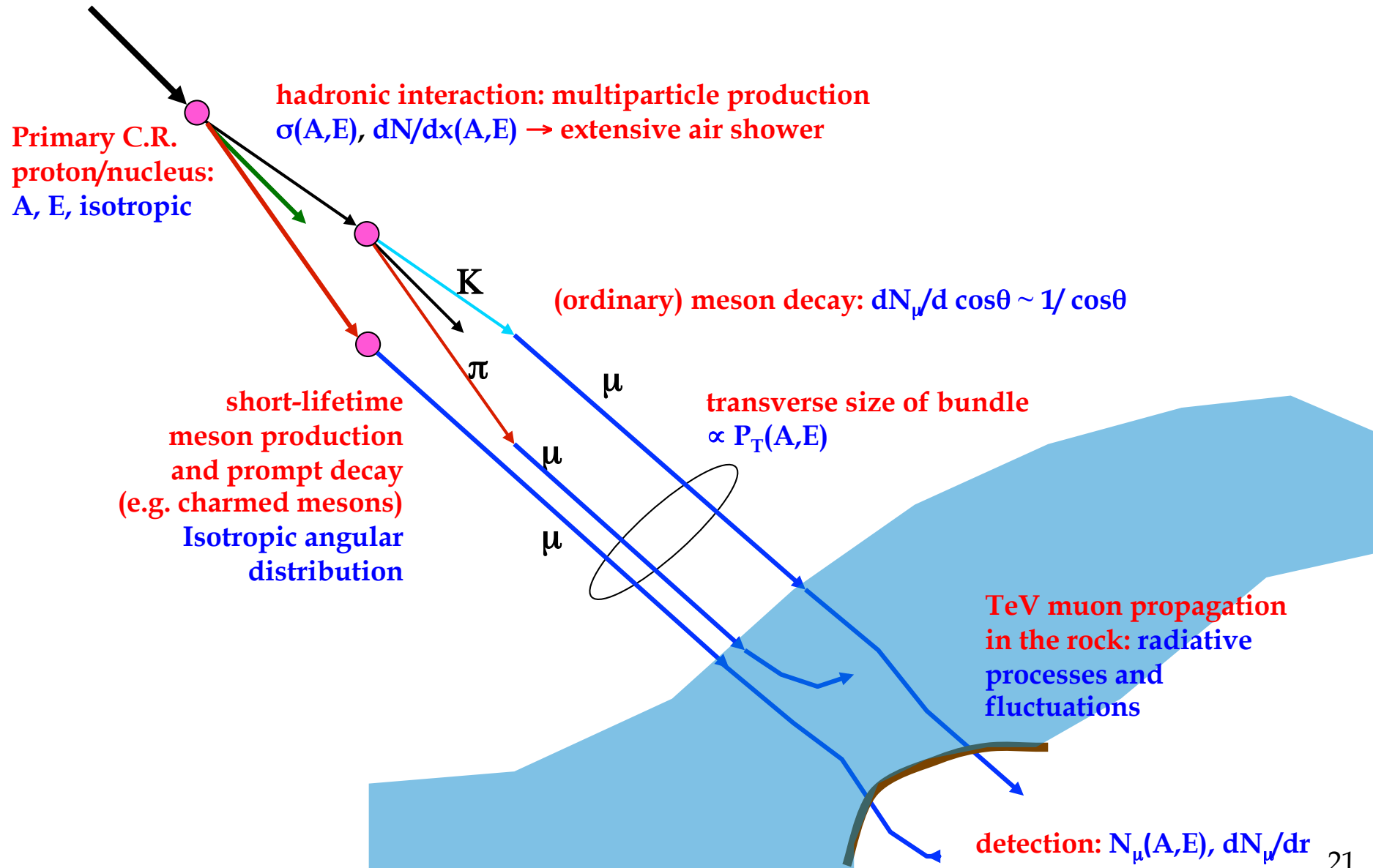
Conclusions

- We reviewed the present knowledge of the atmospheric muon charge ratio for $E_\mu \cos\theta^* > 1$ TeV
- The new OPERA data with a largely increased statistics were presented: OPERA measured the charge ratio in the highest energy region
- Found a strong reduction of the charge ratio for multiple muon events
- The integral value is compatible with expectation from a simple π -K model, as the transition between the pion and the kaon region
- Unexpected behavior in the last HE points of the energy distribution:
 - Further checks mandatory, analysis on going on the 2012 data with inverted magnet polarity
- If the drop will not be explained by still unknown systematic effects, it would be difficult to accommodate the effect within the present knowledge of the primary radiation and of the hadronic interaction features for light mesons in the kinematical region of interest [see J. Phys.: Conf. Ser. 375 (2012) 052018]



Spares

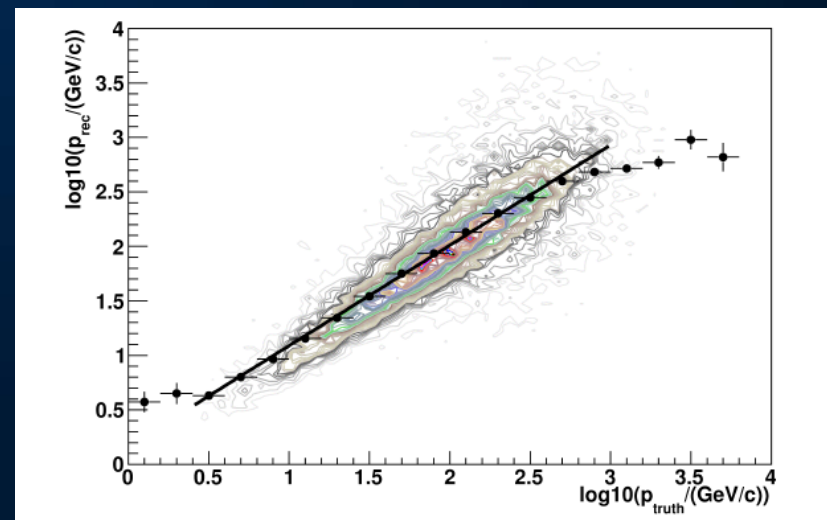
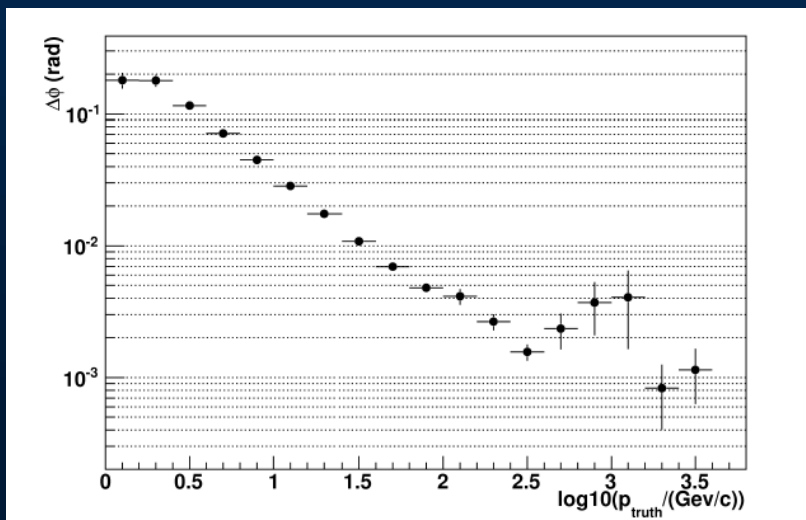
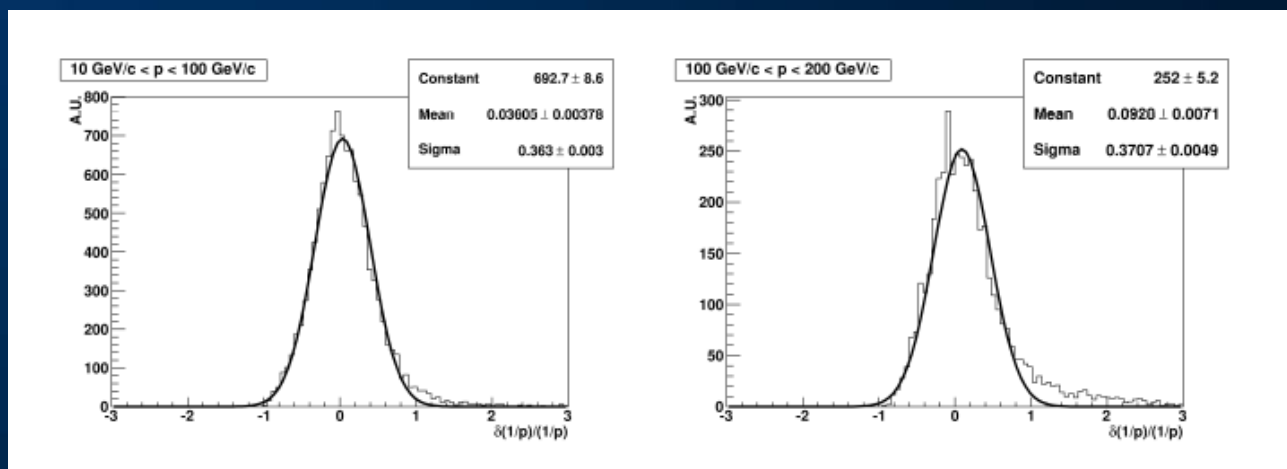
The physics of cosmic ray TeV muons





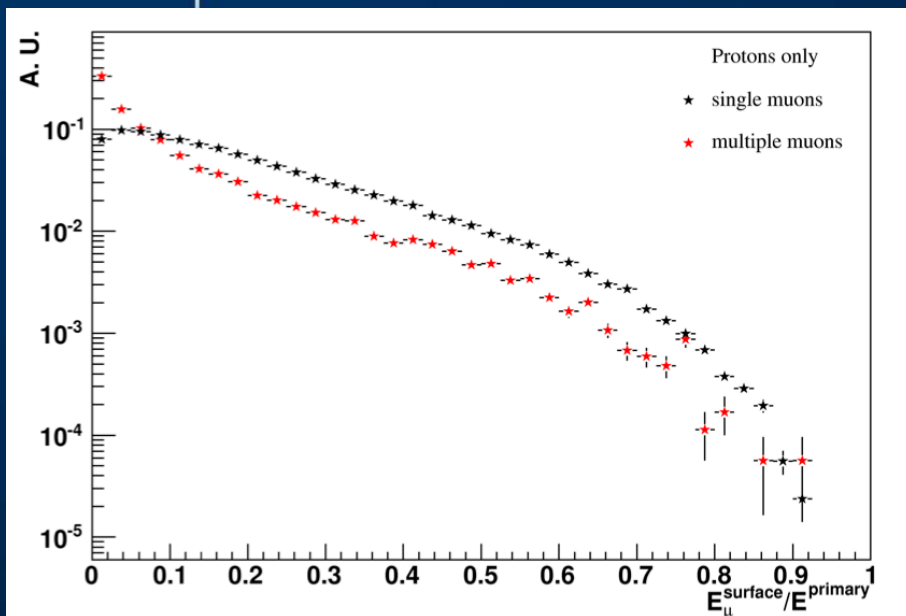
Momentum reconstruction performance

Momentum resolution



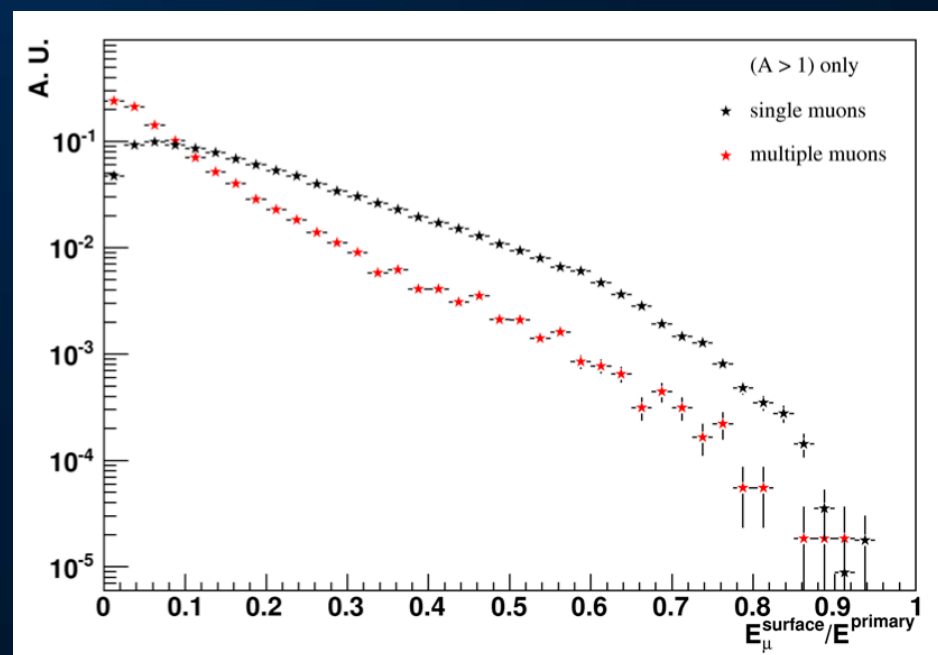


Fraction of primary energy taken by μ



$$E_{\mu}^{\text{surface}} / E^{\text{primary}}$$

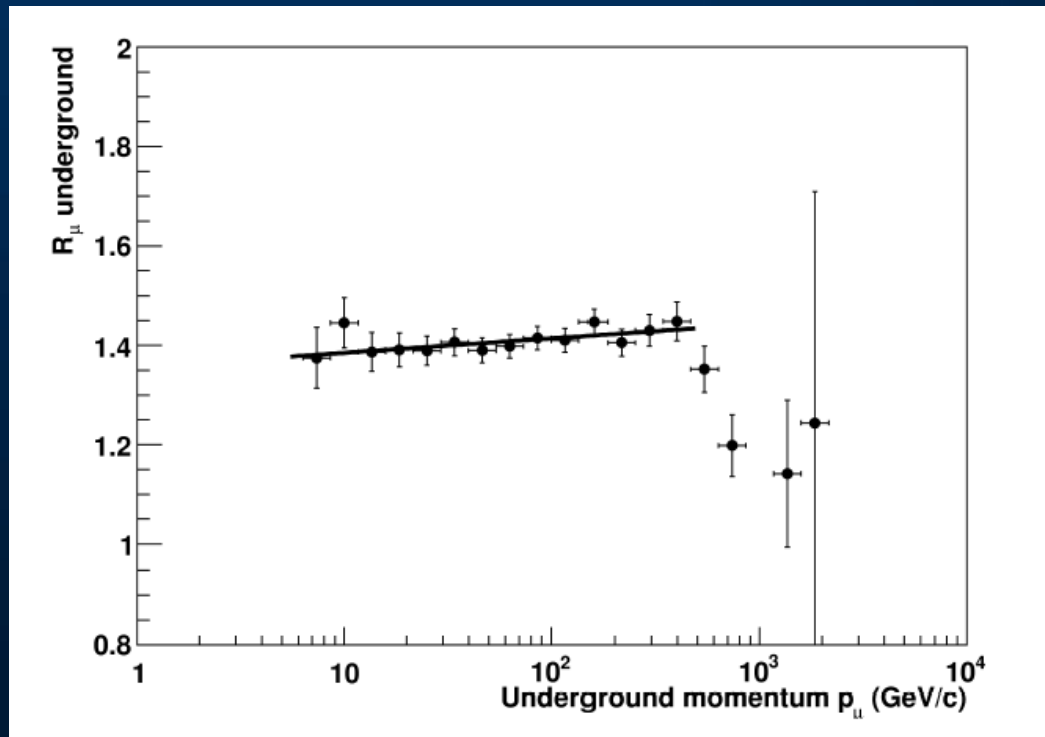
Different distributions for **single** and for **multiple** muons





R_μ as a function of p_μ

R_μ for single muons was computed in bins of p_μ (underground momentum) and unfolded (bin migrations taken into account with the Bayes method)
Evolution with p_μ is compatible with a constant



Linear fit in the range (5,500) GeV/c

$$R_\mu = a_0 + a_1 \log_{10} p_\mu$$

$$\rightarrow a_0 = 1.36 \pm 0.03$$

$$\rightarrow a_1 = 0.029 \pm 0.018$$

$$R_\mu = c_0$$

$$\rightarrow c_0 = 1.409 \pm 0.008$$

$\Delta\chi^2/\text{dof} = 2.48/1$ ($\sim 1.6 \sigma$)
(compatible with a constant)



Maximum detectable momentum

Consider the magnetic field bending and the total deflection spoiling (detector + MCS)

$$\Delta\phi_B = \frac{0.3Bd}{p}$$

$$\sigma_{\Delta\phi} = \sqrt{\sigma_{\phi_1}^2 + \sigma_{\phi_2}^2 + \left(\frac{0.0136}{p}\right)^2 \frac{d}{X_0}}$$

Requiring $\Delta\phi_B/\sigma_{\Delta\phi} > 1$ we obtain (for $\phi=0$):

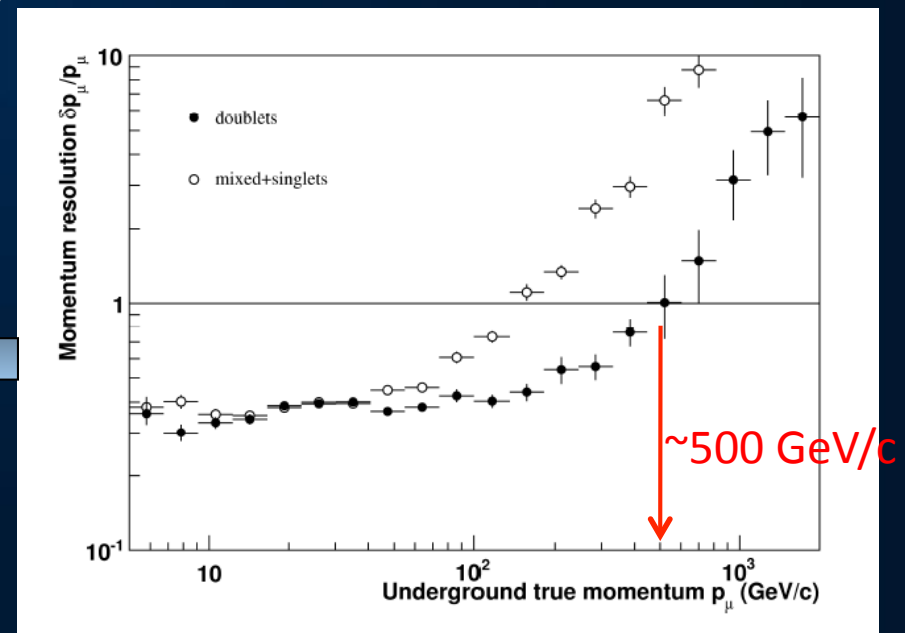
$$p_{\max}(\text{doublets}) = 1.25 \text{ TeV}/c$$

$$p_{\max}(\text{singlets}) = 150 \text{ GeV}/c$$

$$p_{\max}(\text{mixed}) = 260 \text{ GeV}/c$$

Exact computation at all angles (MC)

$$\delta p_{\mu}/p_{\mu} = 1$$





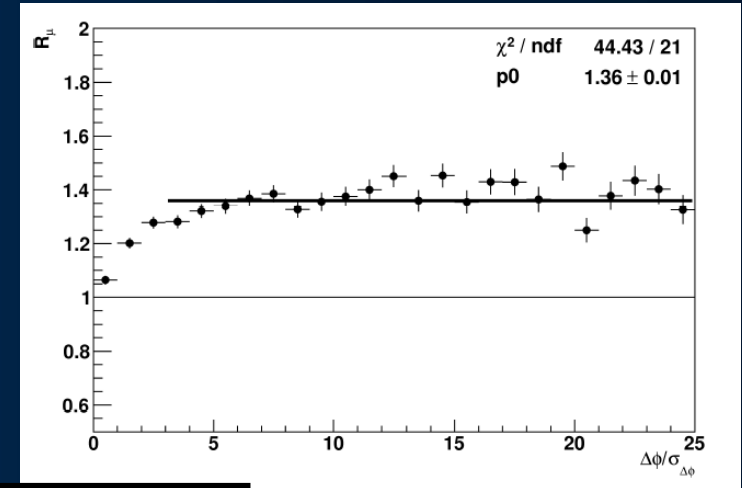
Quality cuts

Deflection Cut: Cut on bending angle information below the PT resolution

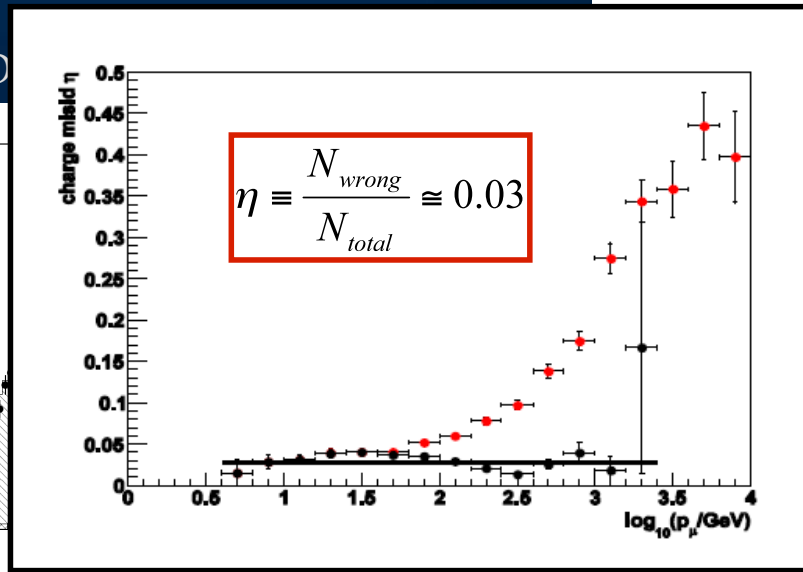
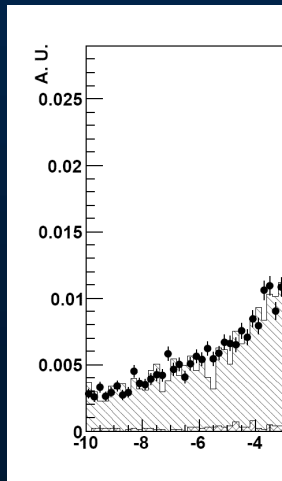
We choose the cut: $\Delta\phi/\sigma_{\Delta\phi} > 3$

R_{μ} saturates above 3σ ; as expected, for lower number of σ , $R_{\mu} \rightarrow 1$ (charge misidentification $\sim 50\%$, total randomization of charge reconstruction)

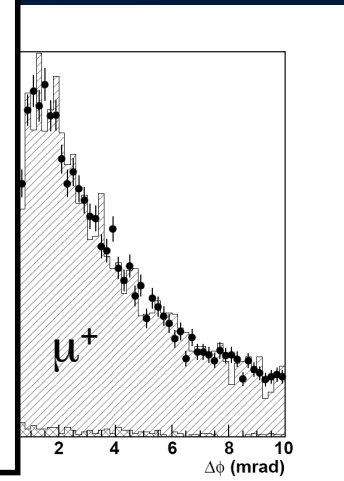
We also require $\Delta\phi < 100$ mrad, to reject isolated secondary particles produced at large angles by high energy muons



before D



after cut



Exp. data
MC data

MC distributions are split in the 2 regions μ^+_{true} and μ^-_{true}

misidentification probability



Phenomenological background

Taking into account the contribution of kaons and pions to the muon flux, a simplified model of the atmospheric muon charge ratio (Gaisser) is given

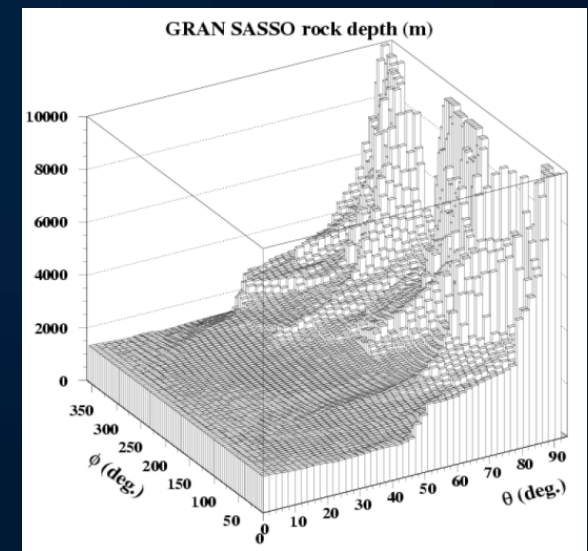
$$\phi_{\mu^\pm} = \frac{\phi_0 \mathcal{E}_\mu^{-\gamma}}{1 - Z_{NN}} \left\{ \frac{a_\pi Z_{N\pi^\pm}}{1 + b_\pi \mathcal{E}_\mu \cos \theta / \epsilon_\pi} + \frac{a_K Z_{NK^\pm}}{1 + b_K \mathcal{E}_\mu \cos \theta / \epsilon_K} \right\}$$

where \mathbf{E}_μ is the muon energy in atmosphere, θ is the zenith angle, $\gamma \sim 2.7$ is the primary cosmic ray spectral index, $a_{\pi/K}$ and $b_{\pi/K}$ are known kinematical factors, $\epsilon_{\pi/K}$ are π/K critical energies

Measuring R_μ as a function of $E_\mu \cos \theta$ we can infer the parameters $R_\pi = Z_{N\pi^+}/Z_{N\pi^-}$ and $R_K = Z_{NK^+}/Z_{NK^-}$.

OPERA is the magnetized experiment with the largest $\langle E_\mu \cos \theta \rangle$

- UTAH (1975): $\langle E_\mu \cos \theta \rangle \sim 500$ GeV
- MINOS (2007): $\langle E_\mu \cos \theta \rangle \sim 1000$ GeV
- OPERA (2009): $\langle E_\mu \cos \theta \rangle \sim 2000$ GeV





The $E_\mu \cos\theta$ regions

SOUDAN

LNGS

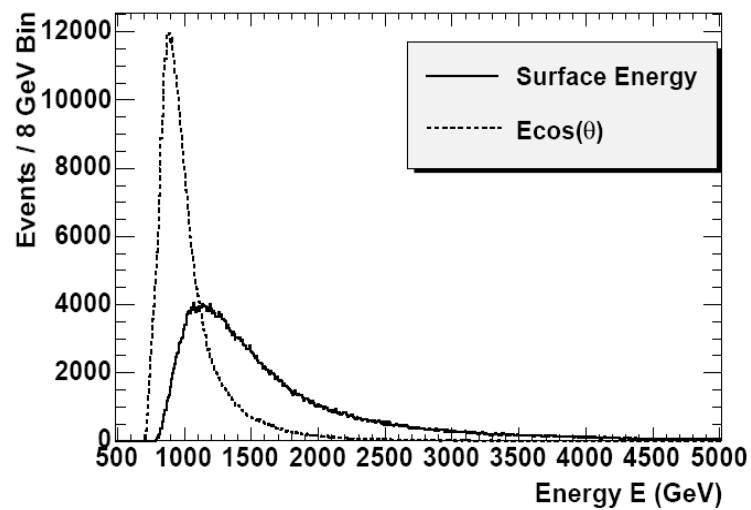
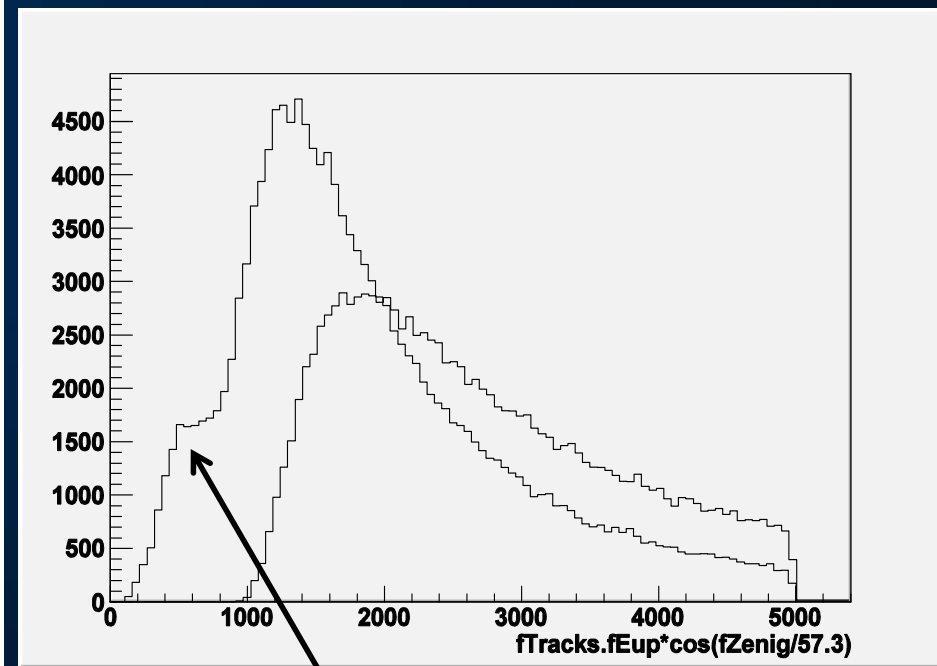


Figure 1: Distribution of E and $E_\mu^{surface} \cos\theta$ for MINOS data muons in the far detector, after cuts.



“Teramo Valley”

From: [MINOS Collaboration], in ICRC '07 proceedings (arXiv: 0706.0538v1); also in arXiv:0906.3726v1



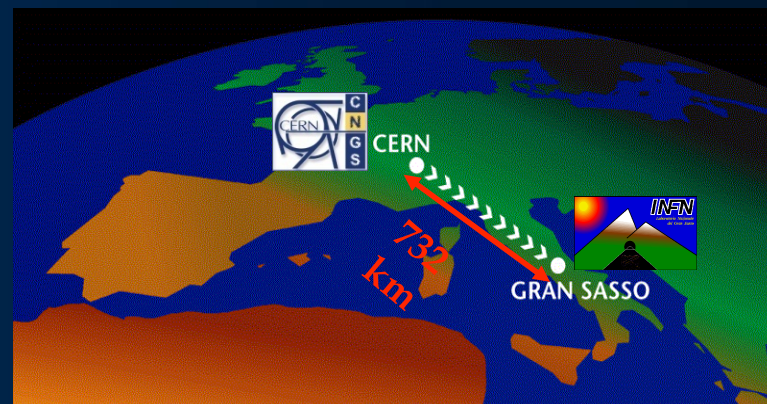
The OPERA experiment

Oscillation Project with Emulsion tRacking Apparatus

OPERA is a long baseline neutrino oscillation experiment aiming at the detection of ν_τ appearance in an almost pure ν_μ beam (CNGS)

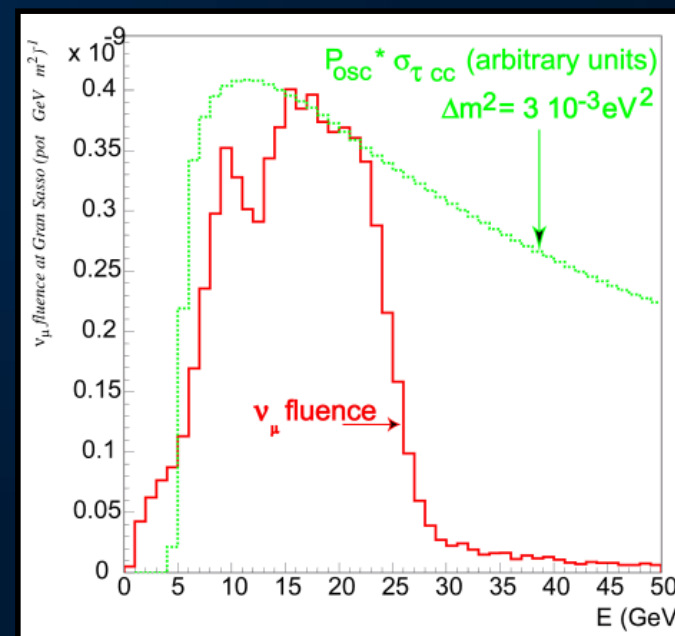
CNGS beam optimized to maximize the number of ν_τ CC interactions: taking into account the ν_τ CC cross section and the τ production threshold

→ ν_μ flux “off peak” w.r.t the maximum oscillation probability (~ 1.5 GeV)



L	732 km
$\langle E_\nu \rangle$	17 GeV
$(\nu_e + \bar{\nu}_e) / \nu_\mu$	0.87%
$\bar{\nu}_\mu / \nu_\mu$	2.1%
ν_τ prompt	negligible

(ν_μ CC)





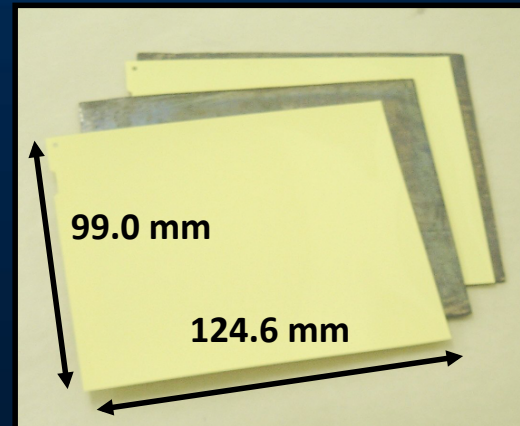
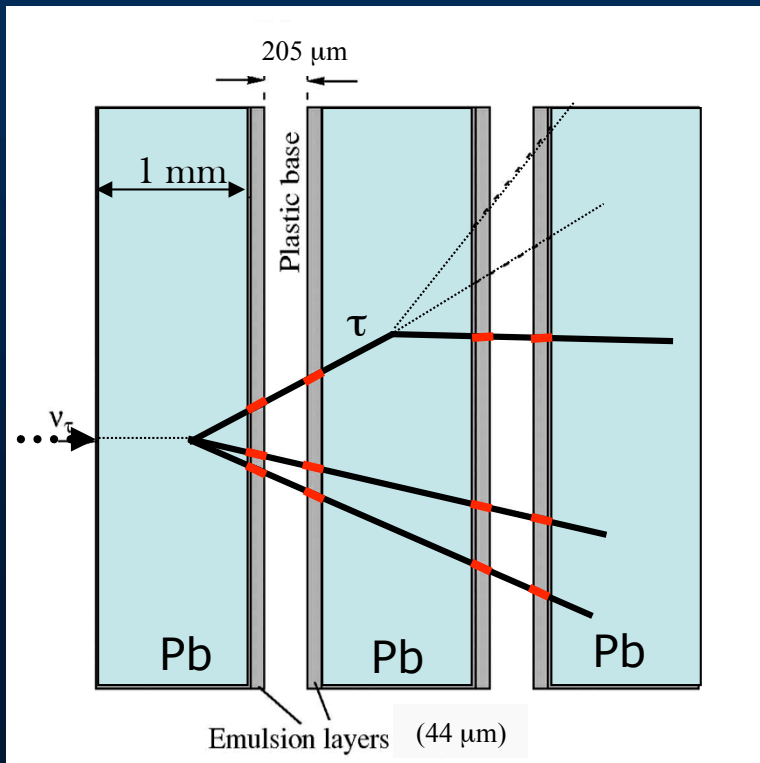
Emulsion Cloud Chamber

ECC = sequence of emulsion-lead layers

→ High resolution and large mass in a modular way

The brick is the target basic component

- 57 nuclear emulsion films interleaved with 1 mm thick lead plates



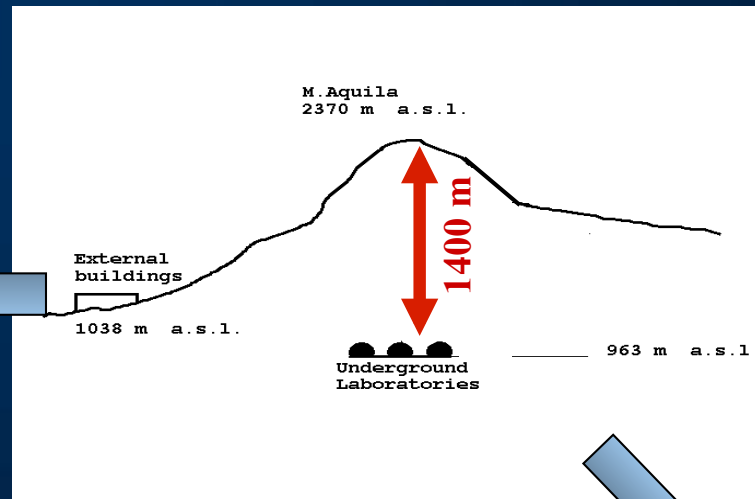
Emulsion
resolution:
 $\delta x = 1 \mu\text{m}$
 $\delta\theta = 2 \text{ mrad}$

~150000 bricks (1.25 ktons)



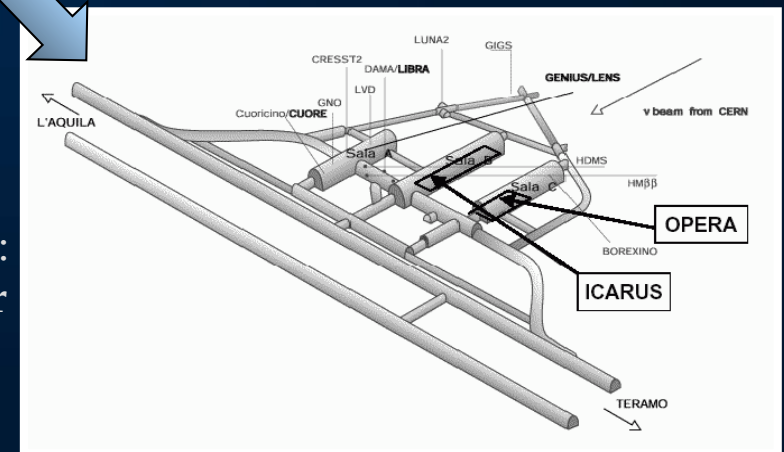
Cosmic ray studies with OPERA

CNGS beam events are identified through a timing coincidence with CERN
→ possibility to collect cosmic events during the physics run



Gran Sasso underground lab:
on average 1400 m of rock
shielding (3800 m w.e.),
cosmic ray flux reduced by a
factor 10^6 w.r.t. surface, very
reduced environmental
radiactivity

OPERA vs. previous and current underground experiments:
charge and momentum measurement with the spectrometer
at large depth and excellent timing capabilities (~ 10 ns)





Monte Carlo simulation

Two Monte Carlo simulations for different purposes:

1) a fast tool for the cross-check with experimental data, for the validation of the analysis software and for unfolding

→ a **parameterized generator (MC1)** with a primary chemical composition based on the MACRO fit model [Phys. Rev. D56 (1997) 1418]

Approximate but fast

$$\frac{Rate_{DATA}}{Rate_{MC1}} = (96.3 \pm 0.3)\%$$

2) a reliable tool for surface muon energy estimation and a link between underground variables and primary cosmic ray parameters

→ a FLUKA-based **full Monte Carlo simulation (MC2)** with primary composition model from Astropart. Phys. 19 (2003) 193

Reliable but slow

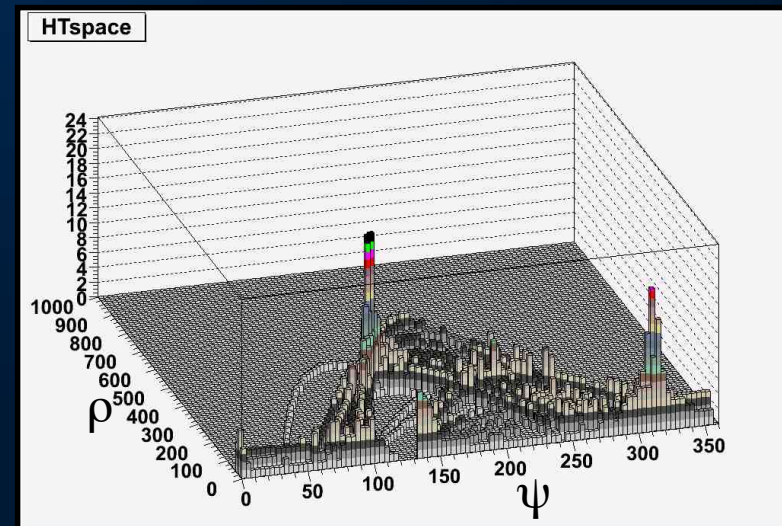
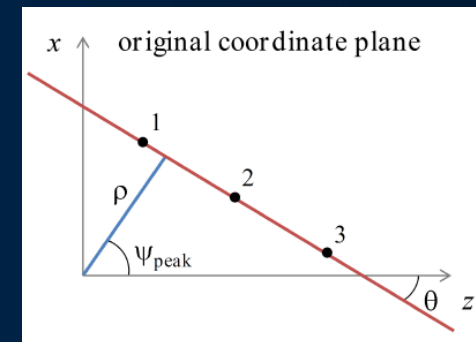
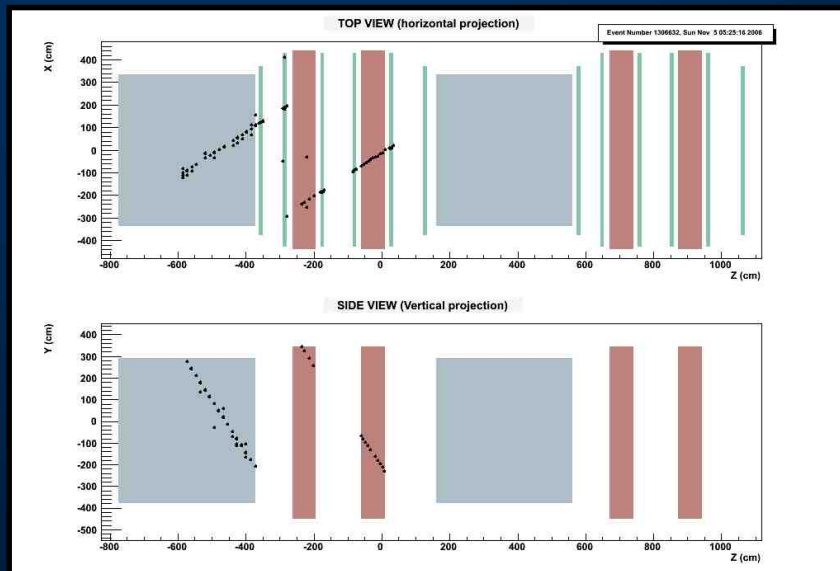


Event reconstruction: Pattern recognition

Dedicated pattern recognition for cosmic events (\equiv outside CNGS spill window):
takes into account event topology and multiple events (μ bundles):

➤ hybrid strategy:

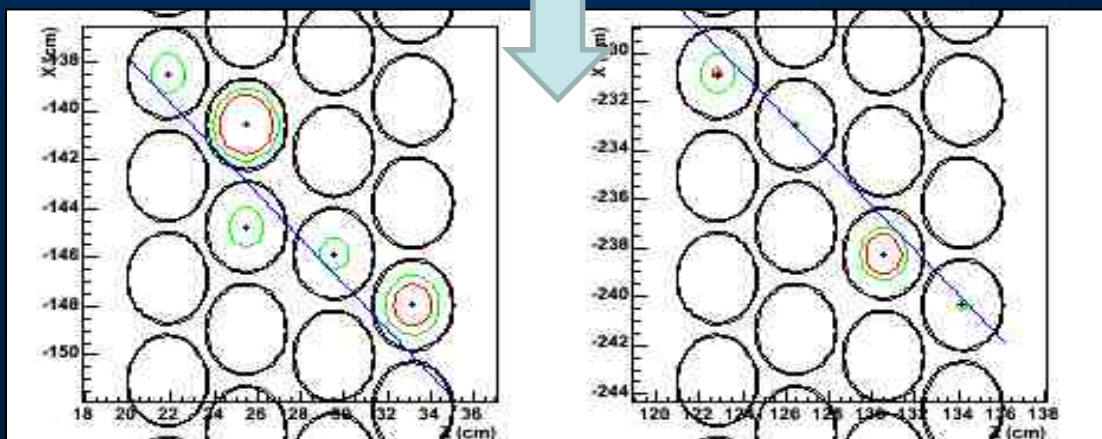
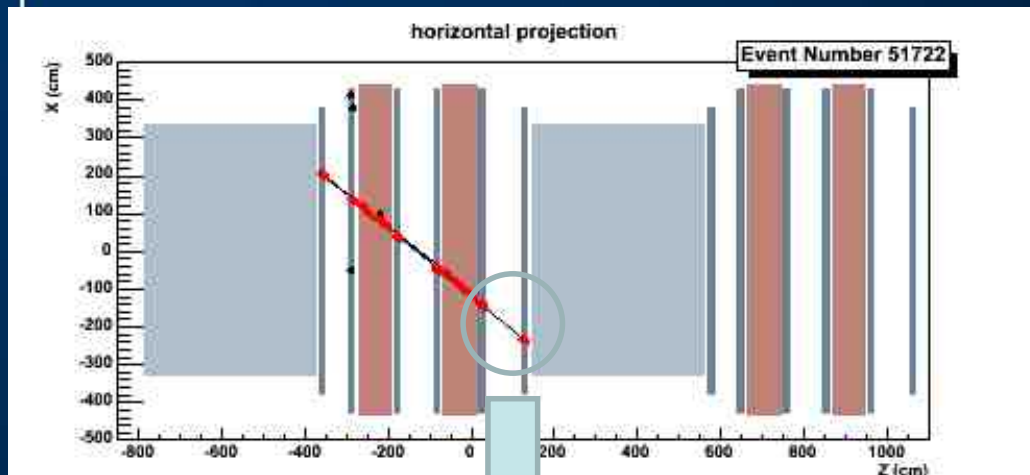
global method (Hough Transform) + **local method** (pivot points)



Real double muon event



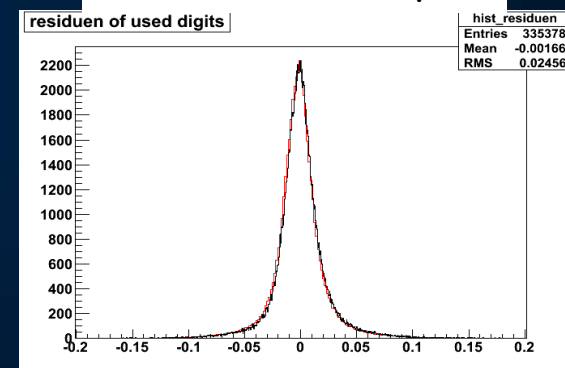
PT track reconstruction



Global tracks are used to “guide” track finding and fitting in the PT system:

- Find the line tangent to the drift circles with the best χ^2
- 250 μm position resolution
- 0.15 (1) mrad angular resolution for doublets (singlets) for $\phi = 0$ (improve for $\phi > 0$)

Residuals $\sim 250 \mu\text{m}$



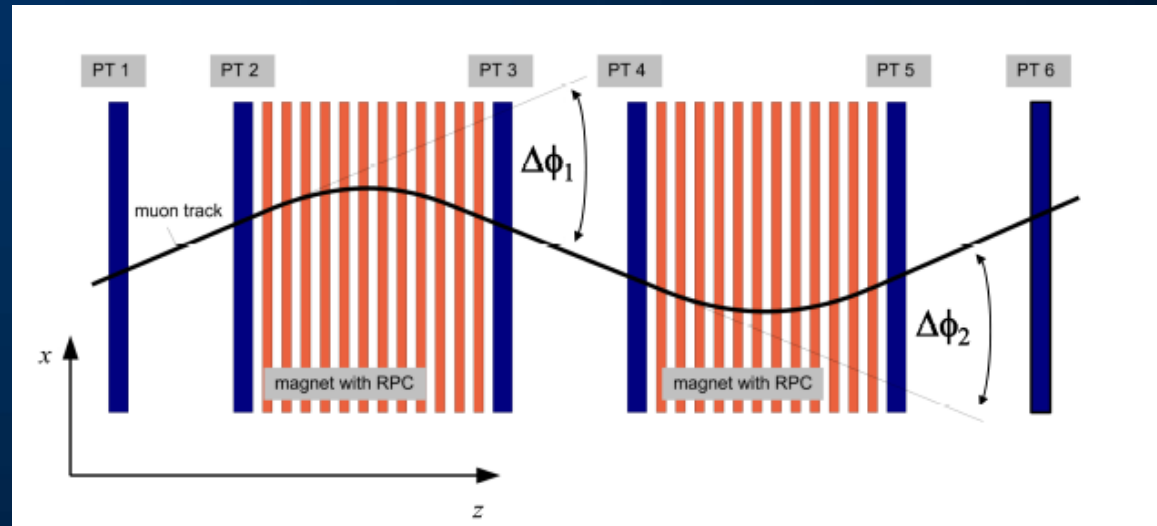
Track reconstructed in 1 PT station (in a pair) \equiv **singlet**
Track reconstructed in 2 PT station (in a pair) \equiv **doublet**



Charge and momentum reconstruction

The charge and momentum reconstruction is performed for tracks crossing at least one magnet arm using the bending angle information $\Delta\phi$

Top view of the OPERA detector



- In each side of the magnet arm we can reconstruct an independent angle ϕ_j , $j=1,\dots,6$.
- Each ϕ_j can be reconstructed with one station (*singlets*) or two stations (*doublers*)
- We compute $\Delta\phi_k = \phi_i - \phi_j$, $k=1,\dots,4$, angle differences between adjacent station-pairs
 - 55% of $\Delta\phi$'s comes from doublers
 - 9% from singlets
 - 36% are mixed



$$p_k = \frac{l(dE/dx)}{1 - \exp[\Delta\phi(dE/dx)/eB]} \sqrt{1 + \frac{s_{yz}^2}{1 + s_{xz}^2}}$$

$B \equiv Bd/l =$
effective magnetic field



Data reduction

A set of progressive cuts applied in order to isolate a clean data sample:

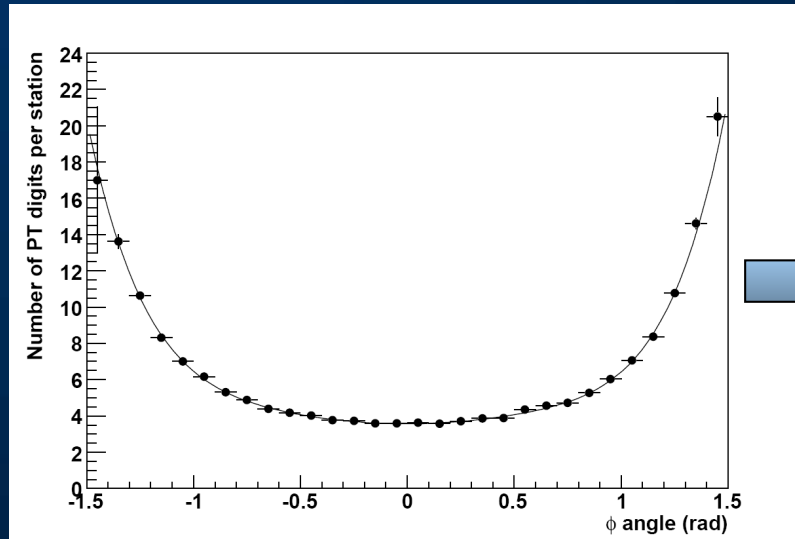
- a) at least one reconstructed $\Delta\phi$ angle
(*acceptance cut*)
- b) Remove events with large number of PT hits
(*clean PT cut*)
- c) Remove events with bending smaller than the experimental resolution (*deflection cut*)
- c') Remove events with very large bending
(effective for $p_{\mu} < 5 \text{ GeV}/c$)



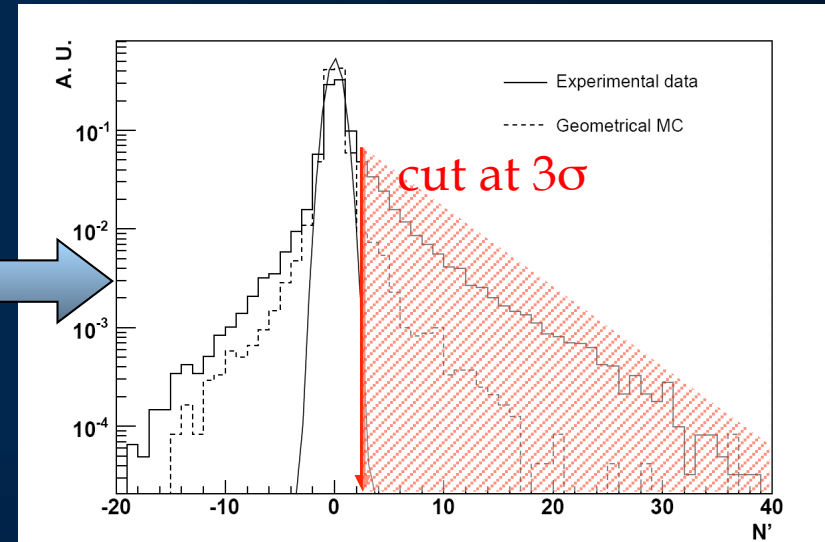
Quality cuts

Clean PT Cut: number of digits allowed from geometrical considerations:

$$N_{\text{data}}(\text{re-scaled}) = N_{\text{data}} - N_{\text{MC}}(\text{geometrical})(\phi)$$



Geometrical MC dependence on the ϕ angle



Re-scaled distribution of the experimental number of PT digits

The cut is chosen at 3σ of the gaussian fit on MC data.

We verified by visual inspection that events removed by this cut have many PT digits fired due to interactions, showers, δ -rays

