

Recent developments in POWHEG and MINLO

Emanuele Re*

Rudolf Peierls Centre for Theoretical Physics, University of Oxford



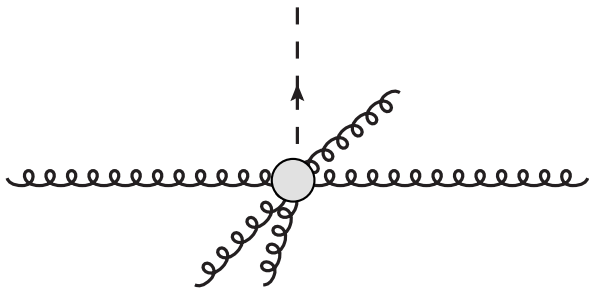
LoopFest XII

Florida State University, 13 May 2013

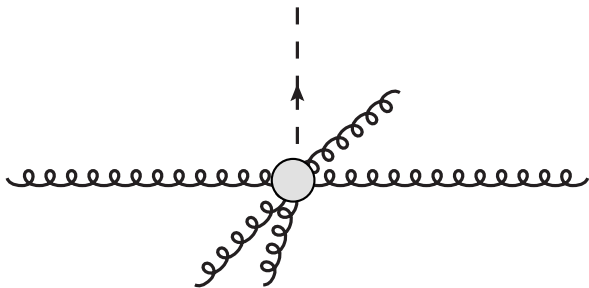
*based mainly on Hamilton et al. [1206.3572,1212.4504]
+ many others for POWHEG improvements [POWHEG BOX homepage]

- Multiscale Improved NLO
- Merging of B at NLO and $B + 1j$ at NLO, and possible further developments
- (selected) new processes/improvements in POWHEG BOX

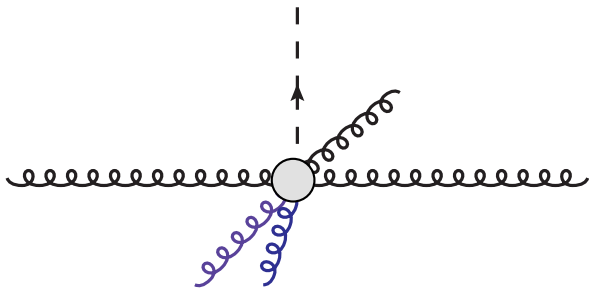
- scale often chosen a posteriori, requiring typically
 - NLO corrections to be small
 - sensitivity upon scale choice to be minimal (\rightarrow plateau in $\sigma(\mu)$ vs. μ)
 - Motivation: bad scale \rightarrow large logs \rightarrow potentially large K factors and large scale dependence
 - However large NLO effects **can be genuine**: new channels (and gluon fluxes), Sudakov logarithms
-
- Typically these issues show up when approaching PS regions with widely different scales (e.g. $B+$ jets).
 - At LO, the **CKKW** procedure allows to **take these effects into account**, modifying the LO weight $B(\Phi_n)$ obtained from ME in order to include LL effects.



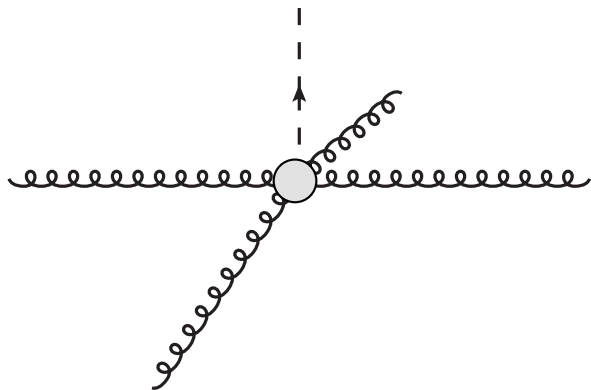
- start from ME weight $B(\Phi_n)$



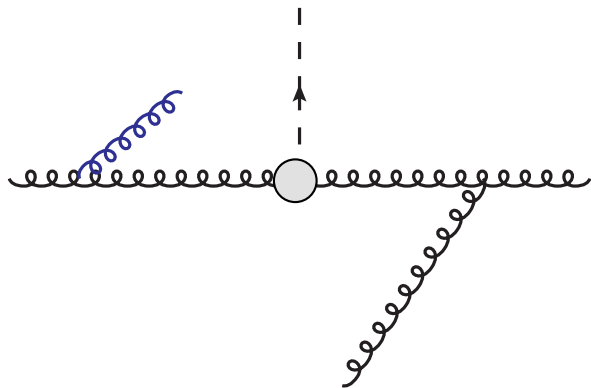
- find “most-likely” shower history (via k_T -algo)



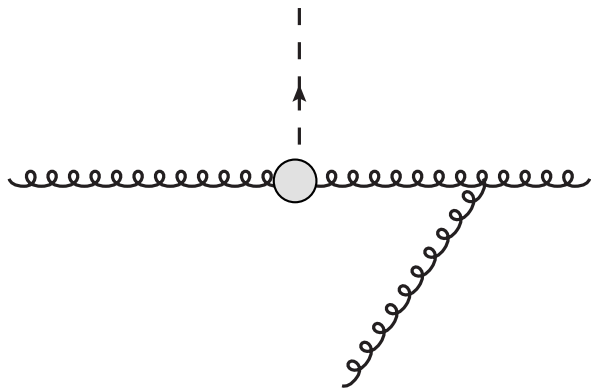
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- clustering scale $q_1 = k_T$



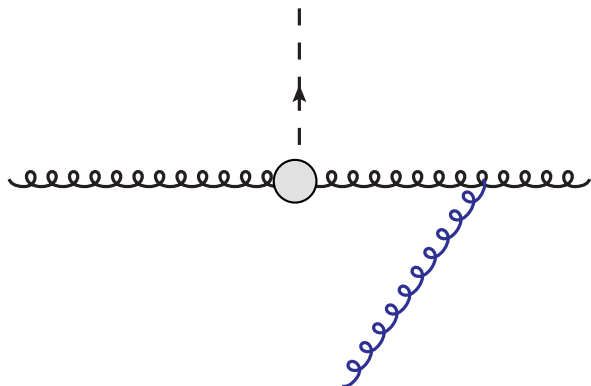
- find “most-likely” shower history (via k_T -algo)



- find “most-likely” shower history (via k_T -algo)
- clustering scale $q_2 = k_T$

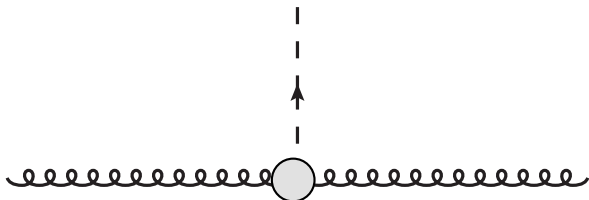


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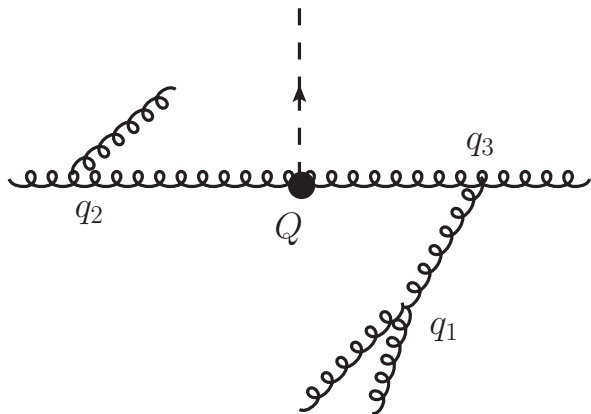


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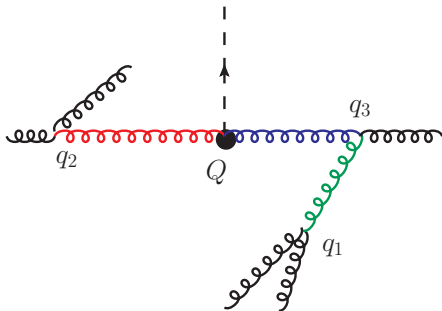
- clustering scale
 $q_3 = k_T$



- Hard process scale Q



- most-likely shower history



- New weight:

$$\alpha_S^5(Q)B(\Phi_3) \rightarrow \alpha_S^2(Q)B(\Phi_3) \frac{\Delta_g(Q_0, Q)}{\Delta_g(Q_0, q_2)} \frac{\Delta_g(Q_0, Q)}{\Delta_g(Q_0, q_3)} \frac{\Delta_g(Q_0, q_3)}{\Delta_g(Q_0, q_1)}$$

$$\Delta_g(Q_0, q_2)\Delta_g(Q_0, q_2)\Delta_g(Q_0, q_3)\Delta_g(Q_0, q_1)\Delta_g(Q_0, q_1)$$

$$\alpha_S(q_1)\alpha_S(q_2)\alpha_S(q_3)$$

where $Q > q_3 > q_2 > q_1 \equiv Q_0$ and, at LL,

$$\log \Delta_f(q_T, Q) = - \int_{q_T^2}^{Q^2} \frac{dq^2}{q^2} \frac{\alpha_S(q^2)}{2\pi} \left[A_{1,f} \log \frac{Q^2}{q^2} + B_{1,f} \right]$$

- Fill phase space below Q_0 with **vetoed** shower

MiNLO: Multiscale Improved NLO

[Hamilton,Nason,Zanderighi,1206.3572]

- ⇒ Use CKKW on top of NLO computation that potentially involves many scales.
- ✓ Scales assigned a priori

$$\alpha_S^n(Q) \Rightarrow \alpha_S(q_1)\alpha_S(q_2)\dots\alpha_S(q_n)$$

- ✓ include Sudakov FFs in internal and external lines ⇒ improved description of Sudakov regions

$$B(\Phi_n) \Rightarrow B(\Phi_n) \times \{\Delta(Q_0, Q)\Delta(Q_0, q_i)\dots\}$$

- ✓ $B+$ jets cross-section finite without generation cuts

Next-to-Leading Order accuracy needs to be preserved

- 1 Scale dependence shows up at NNLO [“scale compensation”]:

$$O(\mu') - O(\mu) = \mathcal{O}(\alpha_S^{n+2}) \quad \text{if} \quad O \sim \alpha_S^n \quad \text{at LO}$$

- 2 Away from soft-collinear regions, exact NLO recovered:

$$O_{\text{MiNLO}} = O_{\text{NLO}} + \mathcal{O}(\alpha_S^{n+2})$$

- 1 MiNLO: All α_S in Born term are chosen with CKKW (local) scales q_1, \dots, q_n

$$\alpha_S^n(\mu_R)B \Rightarrow \alpha_S(q_1)\alpha_S(q_2)\dots\alpha_S(q_n)B$$

- Normal NLO structure ($\mu = \mu_R$):

$$\sigma(\mu) = \underbrace{\alpha_S^n(\mu)B}_{\text{Born}} + \underbrace{\alpha_S^{n+1}(\mu)\left(C + nb_0 \log(\mu^2/Q^2)\right)B}_{\text{Virtual}} + \underbrace{\alpha_S^{n+1}(\mu)R}_{\text{Real}}$$

- Explicit μ dependence of virtual term as required by RG invariance:

$$\alpha_S^n(\mu')B = \left[\alpha_S(\mu) \left(-nb_0\alpha_S^{n+1}(\mu) \log(\mu'^2/\mu^2) \right) \right] B + \mathcal{O}(\alpha_S^{n+2})$$

$$\text{Virtual}(\mu') = \text{Virtual}(\mu) + \alpha_S^{n+1}(\mu)nb_0 \log(\mu'^2/\mu^2) B + \mathcal{O}(\alpha_S^{n+2})$$

$$\Rightarrow \sigma(\mu') - \sigma(\mu) = \mathcal{O}(\alpha_S^{n+2})$$

- In MiNLO “scale compensation” kept if

$$\left(C + nb_0 \log(\mu_R^2/Q^2) \right) B \Rightarrow \left(C + nb_0 \log(\bar{\mu}_R^2/Q^2) \right) B$$

$$\text{with } \bar{\mu}_R^2 = (q_1 q_2 \dots q_n)^{2/n}$$

- 2 MiNLO: internal and external lines of Born “skeleton” are assigned Sudakovs

$$B(\Phi_n) \Rightarrow B(\Phi_n) \times \{\Delta(Q_0, Q)\Delta(Q_0, q_i)\dots\}$$

- Expanded Sudakovs generate IR logs, that hit $\alpha_S^n B(\Phi_n)$

$$\begin{aligned}\Delta_f(q', Q) &= 1 + \Delta_f^{(1)}(q', Q) + \mathcal{O}(\alpha_S^2) \\ \Delta_f^{(1)}(q', Q) &= -\frac{\alpha_S}{2\pi} \left[\frac{1}{2} A_{1,f} \log^2 \frac{Q^2}{q'^2} + B_{1,f} \log \frac{Q^2}{q'^2} \right]\end{aligned}$$

- These terms are $\mathcal{O}(\alpha_S^{n+1})$, so already present in original NLO!
- Subtract them to keep NLO accuracy (but keep overall CKKW Sudakovs!):

$$B(\Phi_n) \Rightarrow B(\Phi_n) \left(1 - \Delta^{(1)}(Q_0, Q) - \Delta^{(1)}(Q_0, q_i) + \dots \right)$$

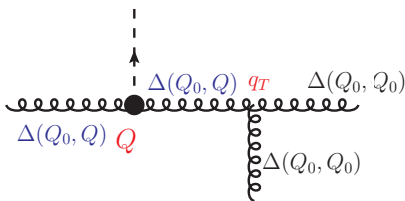
Example, in 1 line: $H + 1$ jet

- Pure NLO:

$$\bar{B} = \alpha_s^3(\mu_R) \left[B + \alpha_s V(\mu_R) + \alpha_s \int d\Phi_{\text{rad}} R \right]$$

- MiNLO:

$$\bar{B} = \alpha_s^2(M_H) \alpha_s(q_T) \Delta_g^2(q_T, M_H) \left[B \left(1 - 2\Delta_g^{(1)}(q_T, M_H) \right) + \alpha_s V(\bar{\mu}_R) + \alpha_s \int d\Phi_{\text{rad}} R \right]$$



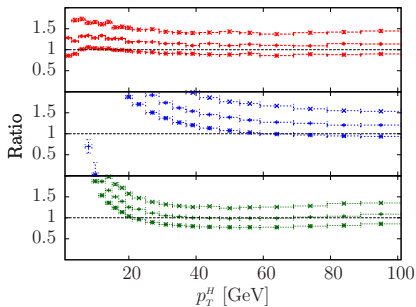
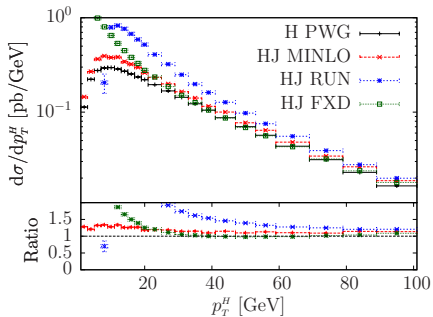
It can be shown that this corresponds to use $\alpha_s^3(q_T) B \left[1 - C_A \frac{\alpha_s}{2\pi} \log^2 \frac{Q^2}{q_T^2} \right]$

MiNLO+POWHEG

- \bar{B} with MiNLO prescription can be used as starting point of POWHEG
- For $B + n$ jets, this modified \bar{B} will be finite even if we are inclusive on radiation (also totally inclusive)
- MiNLO for $B + n$ jets improves description of observables made out of n jets, because Sudakov logs are taken into account
- (One can build observables s.t. MiNLO fails)
- POWHEG resums extra radiation in Φ_{n+1} , without spoiling $B + n$ jets NLO

⇒ Many benefits when combined!

Moreover MiNLO+POWHEG is the starting point to merge NLO B with NLO $B + j$ (without a merging scale).



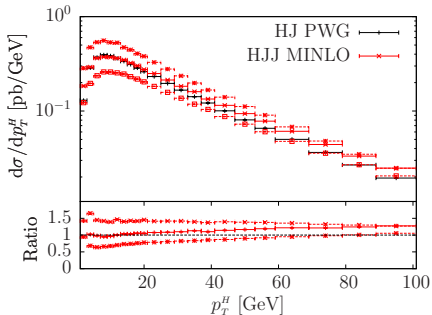
- “H PWG”: POWHEG ($gg \rightarrow H$) + PYTHIA (PS level)
- “HJ RUN(FXD)”: NLO Hj with $\mu = p_{T,H}(M_H)$ (7pt envelopes)

✗ NLO's band don't overlap if $p_{T,H} < 30$ GeV

✓ At small p_T , MinLO is finite and \sim H PWG (both include Sudakov)

✓ At large p_T , MinLO \sim Hj NLO

- HJ FXD higher scale “compensate” for absence of Sudakov (more than HJ RUN, that misses it)

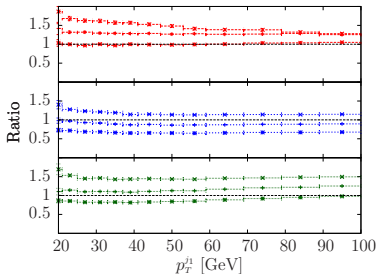
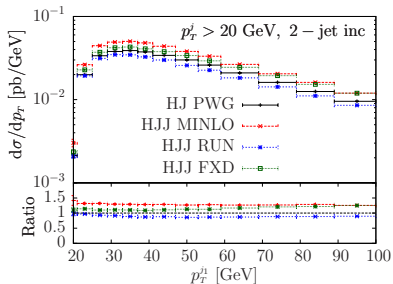


LEFT:

- “HJ PWG”: HJ POWHEG (w MiNLO) + PYTHIA (PS level)
- “HJJ MINLO”: finite even if fully inclusive + compatible with POWHEG down to very low p_T

BOTTOM:

- All results predictive (2 jets required)
- “HJJ MINLO” \sim “HJJ RUN” seems to support running scale H_T ; for $H_T/2$, MINLO and RUN would overlap even better



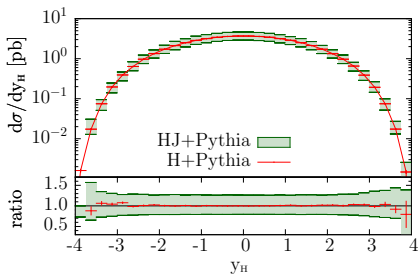
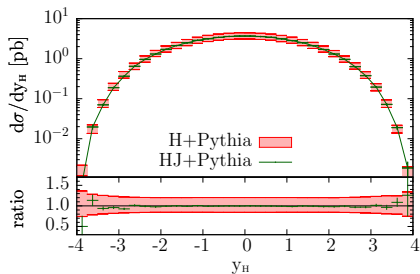
- Accuracy of $B_J + \text{MiNLO}$ for inclusive observables carefully investigated in arXiv:1212.4504 (Hamilton, Nason, Oleari, Zanderighi)
- $B_J + \text{MiNLO}$ describes inclusive boson observables at relative order α_s wrt $B + 0j$ at LO
- However, to reach genuine NLO, higher terms must be order α_s^2 , *i.e.*

$$O_{VJ+\text{MiNLO}} = O_{V@NLO} + \mathcal{O}(\alpha_s^2)$$

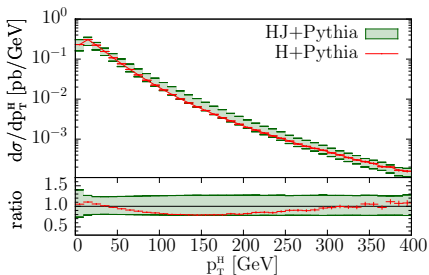
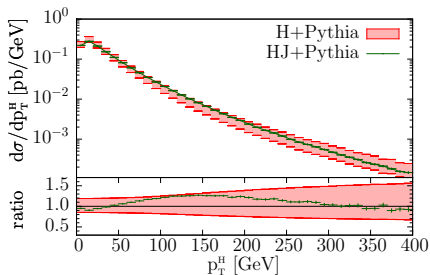
if O is inclusive. “Original MiNLO” contains **ambiguous $\mathcal{O}(\alpha_s^{3/2})$ terms**

- Possible to modify $B_J + \text{MiNLO}$ such that NLO $B + 0j$ is recovered, without spoiling NLO accuracy for $B + 1j$.
 - proof based on careful comparisons of general resummation formula with MiNLO
 - need to include B_2 in Sudakovs
 - need to evaluate $\alpha_s^{(\text{NLO})}$ in $B_J + \text{MiNLO}$ at scale q_T

In the end is like if we merged $\text{NLO}^{(0)}$ and $\text{NLO}^{(1)}$ samples, **without merging** different samples (no merging scale used).



- “H+Pythia”: standalone POWHEG ($gg \rightarrow H$) + PYTHIA (PS level) [7pts band]
- “HJ+Pythia”: HJ-MiNLO* + PYTHIA (PS level) [7pts band]
- ✓ very good agreement (both value and band)



- Good agreement
- At high p_T , bands as expected (LO vs NLO)
- Low p_T shape difference: different NNLL terms in MiNLO Sudakovs
- Bands at low p_T : “H+Pythia” band spurious (S-events, i.e. inherit property of full \bar{B})
 “HJ+Pythia” widens as expected (strong coupling regime)

- NNLO+PS

- HJ-MiNLO* differential cross section $(d\sigma/dy)_{\text{HJ-MiNLO}}$ is NLO accurate

$$\frac{\left(\frac{d\sigma}{dy}\right)_{\text{NNLO}}}{\left(\frac{d\sigma}{dy}\right)_{\text{HJ-MiNLO}}} = \frac{c_2\alpha_S^2 + c_3\alpha_S^3 + c_4\alpha_S^4}{c_2\alpha_S^2 + c_3\alpha_S^3 + d_4\alpha_S^4} \simeq 1 + \frac{c_4 - d_4}{c_2}\alpha_S^2 + \mathcal{O}(\alpha_S^3)$$

- thus, reweighting with this factor, we get NNLO+PS
 - notice that α_S^4 accuracy of HJ-MiNLO* in 1-jet region not spoiled
 - variants possible
-
- merging for higher multiplicity needs more study

Automation:

- [Interface to MadGraph 4](#) (Frederix): automatically builds subprocesses list, B , B_{ij} , $B^{\mu\nu}$, R and large- N Born color structures.
 - Used to build the code for Hj and Hjj (Campbell, Ellis, Frederix, Nason, Oleari), with virtuals from MCFM.
 - Work is in progress (Luisoni, Nason, Oleari, Tramontano) to build an [interface to GoSam](#), to automatically write the code for 1-loop amplitudes
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Complicated processes & parallelization:

- Possible to run POWHEG BOX in parallel. Not only event generation, but also grid computation.
- details in repository [[Parallel-grids](#)]

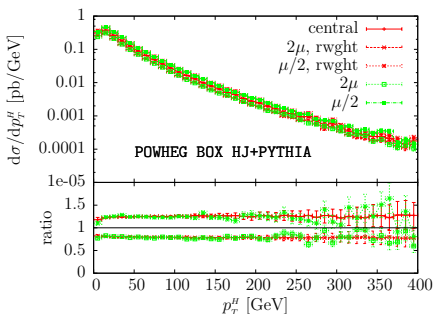
PDF and scale uncertainties:

- Generate MC samples for different scale choices, and, even more, for different PDFs, is very time consuming
- Primitive reweighting facility now superseded by new mechanism

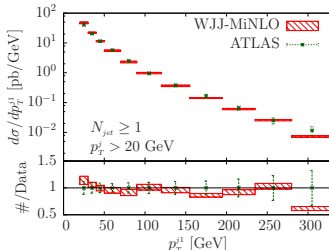
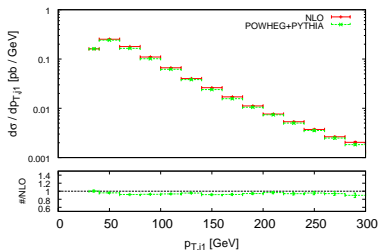
(Hamilton,Nason,ER):

- instead of tracking coefficients of μ logs, for each event save **random number seeds** associated to PS point and **value of \tilde{B}** (+ some extra info)
- for each event in already existing LH file, use stored random number seeds to generate the same PS points, and compute new value of \tilde{B} with new scales and/or pdf
- append new weight, rescaling with

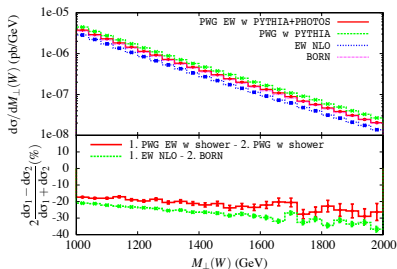
$$w' = w \times (\text{new integrand}/\text{old integrand})$$



- Zjj (ER, '12), Vjj -MiNLO (Campbell, Ellis, Nason, Zanderighi '13)



- NLO QCD and NLO EW, for DY (Bernaciak, Wackerath '12 - Barze et al '12-'13)



1 MiNLO:

- assign scales and Sudakov FF in $B + n$ jets NLO computations
- well-behaved in Sudakov regions
- NLO away from Sudakov regions
- ideal as starting point for POWHEG

2 Improved MiNLO:

- $B + 1$ jet improved MiNLO allows to merge $\text{NLO}^{(0)}$ and $\text{NLO}^{(1)}$ samples, **without merging** (no merging scale used)
- seems possible to reach NNLO+PS
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3 Progress in POWHEG:

- list of processes steadily increasing
- automation
- keep up with theoretical and experimental needs
- ...
- ongoing work on angular correlations in heavy-quark decays
- ...other talks here !

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Thanks for your attention!

Few technicalities for MiNLO:

- $\mu_F = Q_0$ (as in CKKW)
- Cluster with CKKW also V and R kinematics
 - Actual implementation uses FKS mapping for first cluster of Φ_{n+1}
 - Ignore CKKW Sudakov for 1^{st} clustering of Φ_{n+1} (inclusive on extra radiation)
- Some freedom in choice of $\alpha_S^{(n+1)}$ (entering V , R and $\Delta^{(1)}$) (not free for MiNLO merging)
- Used full NLL Sudakovs

Improved MiNLO:

- Resummation formula

$$\frac{d\sigma}{dq_T^2 dy} = \sigma_0 \frac{d}{dq_T^2} \left\{ [C_{ga} \otimes f_a](x_A, q_T) \times [C_{gb} \otimes f_b](x_B, q_T) \times \exp S(q_T, Q) \right\} + R_f$$

- NLO⁽⁰⁾ if $C_{ij}^{(1)}$ included and R_f is LO⁽¹⁾
- Take derivative:

$$\sim \sigma_0 \frac{1}{q_T^2} [\alpha_S, \alpha_S^2, \alpha_S^3, \alpha_S^4, \alpha_S L, \alpha_S^2 L, \alpha_S^3 L, \alpha_S^4 L] \exp S(q_T, Q)$$

- can be shown that

$$\int^{Q^2} \frac{dq_T^2}{q_T^2} L^m \alpha_S^n(q_T) \exp S \sim (\alpha_S(Q^2))^{n-(m+1)/2}$$

- if I drop B_2 in MiNLO, I miss a term $(1/q_T^2) \alpha_S^2 B_2 \exp S$
- upon integration, violate NLO⁽⁰⁾ by a term $\mathcal{O}(\alpha_S^{3/2})$
- wrong scale in NLO α_S in MiNLO produces again same error

More “analytic” proof with all details also available in the paper.