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# LHC LUMINOSITY UPGRADE PLAN

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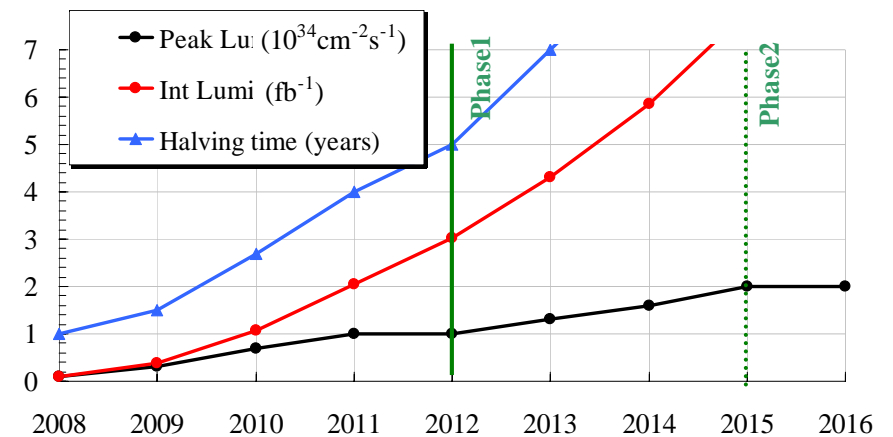
# CONTENTS

- Motivations for the luminosity upgrade
- How to remove bottlenecks
- Actions
- Technologies



# MOTIVATIONS

- Aim of the phase I (2012)
  - If nominal is reached, go up to the maximum peak luminosity that can be tolerated by experiments **without any detector upgrade** (so-called ultimate)  $L \sim 2 \times 10^{34} [\text{cm}^{-2} \text{s}^{-1}]$
  - If nominal is not reached, remove the bottlenecks so that the nominal can be **recovered**
    - For instance: if beam current is not reachable, increase focusing from  $\beta^* = 55 \text{ cm}$  to 35-25 cm to reach nominal
- Aim of phase II (2016-17)
  - After some years ( $\sim 5$ ) at nominal/ultimate luminosity, a big boost is needed  $\rightarrow$  **factor 5-10 is needed**
    - Otherwise the time to halve the statistical error becomes huge)
  - Go up to  $L \sim 10 \times 10^{34} [\text{cm}^{-2} \text{s}^{-1}]$
  - This involves **detector and injector upgrade**
  - All solutions that can be envisaged should be adopted
  - Challenge: energy deposition



A guess of luminosity versus time in absence of phase II upgrade inspired by J. Strait work in 2002



# PEAK LUMINOSITY

- Peak luminosity equation

$$L = \frac{f_{rev} \gamma}{4\pi \epsilon_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*}$$

Constants      Beam intensity      Beam focusing

- Constants

- $\epsilon_n$ : transverse normalized emittance (LHC aperture, injectors) [3.75×10<sup>-4</sup> cm rad]
- $\gamma$ : relativistic factor (energy of the machine, type of particles) [7461]
- $f_{rev}$ : revolution frequency (size of the machine) [1.12 ×10<sup>4</sup> s<sup>-1</sup>]

- Beam intensity

- $N_b$ : number of particle per bunch [1.15 ×10<sup>11</sup>]
- $n_b$ : number of bunches [2808]

- Beam focusing

- $\beta^*$ : beta function in the IP (transverse size of the beam) [55 cm]
- $F$ : geometrical loss reduction factor [0.86]



# BOTTLENECKS

- Present baseline  $L=10^{34}$  [ $\text{cm}^{-2} \text{s}^{-1}$ ]
- Will we reach the baseline ? Bottlenecks
  - The collimation scheme shows that the **impedance is not tolerable**
    - If this estimate is confirmed we have either to
      - Keep  $\beta^*=55$  cm and reduce beam intensity of 60%  $\rightarrow$  go to  $0.16 \times 10^{34}$  (catastrophe!)
      - Keep the intensity and **reduce  $\beta^*=80$  cm** to be able to open the collimator gap  $\rightarrow$  go to  $0.76 \times 10^{34}$  (much better) - today we expect to **lose 25% w.r.t. nominal peak luminosity**
  - Beam focusing: the **triplet aperture** is today, by design, the **bottleneck**
    - LHC was designed to go **up to  $\beta^*=25$  cm** except the IR triplet and D1
    - Enlarge the IR triplet aperture would also ease collimation [R. Assman et al, LIUWG 11 November 2007]
  - **Beam current**: will we reach the nominal values ?
    - Luminosity  $\propto$  square of the no. of particle per bunch  $\rightarrow$  if we just lose 20% we have 2/3 of the nominal !
    - This parameter is determined both by **injectors** and by the **LHC** performances
    - Example: an alternative filling scheme having  $n_b=2592$  is being considered  $\rightarrow$  8% luminosity loss [W. Herr et al, LTC 14 February 2007]
- Experiments can bear up to  $L \sim 2 \times 10^{34}$  [ $\text{cm}^{-2} \text{s}^{-1}$ ]



# PHYSICAL LIMITATIONS: NUMBER OF PARTICLES PER BUNCH

$$L = \frac{f_{rev} \gamma}{4\pi \epsilon_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*}$$

- Number of particle per bunch  $N_b$  is limited by the beam-beam effect, i.e., the **Coulomb interaction** between colliding bunches

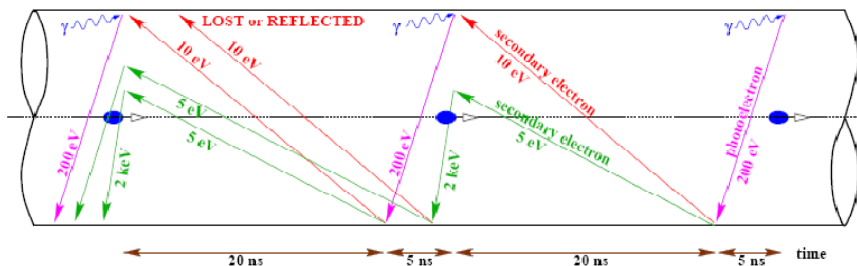
- The beam-beam parameter is defined as  $\xi = \frac{r_p}{4\pi} \frac{N_b}{\epsilon_n} F(\beta^*)$
- Empirical experience on machines prove that one can run as long as  $\xi < 0.015$  – having three experiments this means  $\xi < 0.005$
- For the LHC baseline a **safety margin with  $F=1$**  (instead of 0.86 as we have) has been taken
- This gives **an upper bound on  $N_b = 1.15 \times 10^{11}$**
- Switching off one experiment, we can go up to  **$N_b = 1.7 \times 10^{11}$  (ultimate)**, but we need an **upgrade of the injectors**



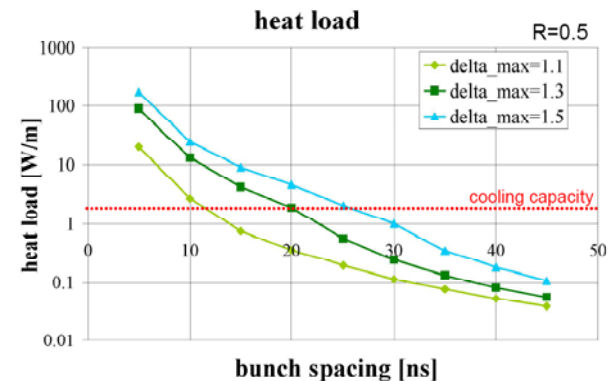
# PHYSICAL LIMITATIONS: NUMBER OF BUNCHES

$$L = \frac{f_{rev} \gamma}{4\pi\epsilon_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*}$$

- The number of bunches  $n_b$  (i.e. bunch spacing) is limited by different factors, among them the **electron cloud effect**
  - Nominal  $n_b=2808$  bunches → bunch spacing of 25 ns
  - Limits given by simulations and some experiences on machines: shorter spacing has to be excluded - larger can be possible
    - Doubling  $n_b$ , the bunch spacing becomes 12.5 ns and it looks not feasible (heat load for e-cloud + image current)



Mechanism of electron cloud build up, courtesy of F. Ruggiero



Heat load induced by electron cloud versus bunch spacing, F. Zimmermann LTC, 6<sup>th</sup> April 2005



# PHYSICAL LIMITATIONS: FOCUSING

$$L = \frac{f_{rev} \gamma}{4\pi\epsilon_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*}$$

- The focusing is presently limited by the **aperture of the quadrupoles** Q1-Q3 around the IP (the so-called triplet)

- The beta function of the beam in the quadrupoles is  $\propto 1/\beta^*$
- The present aperture of 70 mm limits  $\beta^*=0.55$  cm
- Changing the triplet, one hits the **hard limit of the chromaticity correction** at

- Nb-Ti triplet  $\beta^*=0.17$  cm

Nb<sub>3</sub>Sn triplet  $\beta^*=0.14$  cm

[E. Todesco et al, CARE LUMI-06 J. P. Koutchouk et al., PAC 07]

- If the distance of the triplet from the IP is reduced from 23 m to 13 m (**extreme case**), one can further improve by  $\sim 25\%$

- Nb-Ti triplet  $\beta^*=0.14$  cm

Nb<sub>3</sub>Sn triplet  $\beta^*=0.11$  cm





# PHYSICAL LIMITATIONS: FOCUSING VS GEOMETRICAL FACTOR

$$L = \frac{f_{rev} \mathcal{Y}}{4\pi \mathcal{E}_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*} \quad F(\beta^*) = \frac{1}{\sqrt{1 + \phi^2(\beta^*)}}$$

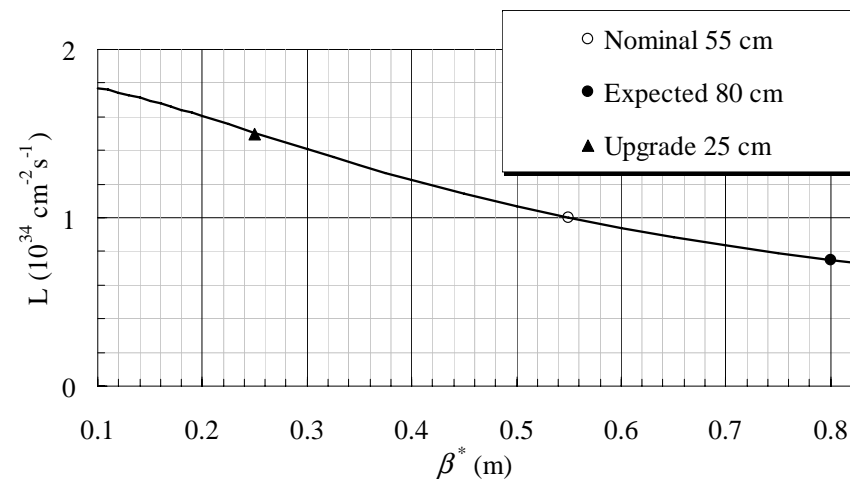
- Luminosity is not  $\propto 1/\beta^*$  - for  $\beta^* < 25$  cm the gain is marginal if the beam current is kept constant

- Empirical scaling law: larger focusing induces large crossing angle and therefore more diluted collisions [Y. Papaphilippou, F. Zimmermann, Phys. Rev. STAB 10 (1999) 104001]

- Going to 25 cm one gains ~50% w.r.t. 55 cm

- $\phi$  is called **Piwinski angle**  
 $\phi \sim 0.64$  for nominal  $\rightarrow F \sim 0.84$

$$\phi = \frac{\vartheta(\beta^*) \sigma_z}{2\sigma(\beta^*)} \propto \frac{\sigma_z}{\beta^*}$$



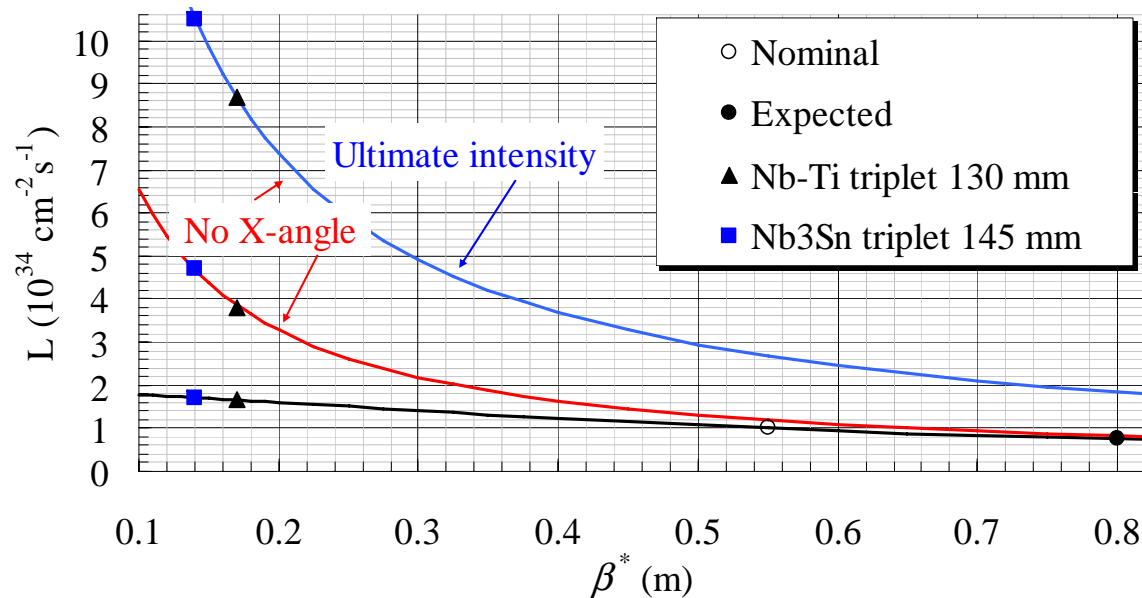
Luminosity versus focusing in case of constant beam current



# UPGRADE BASED ON STRONGER FOCUSING (EARLY SEPARATION/CRAB CAVITY)

- Reduce  $\beta^*$  as much as possible
- Hypothesis: cannot increase bunch charge  $N_b$  beyond nominal
- Kill the crossing angle effect  $\rightarrow$  set the factor  $F=1$  through two means
  - D0: dipole to provide an **early separation** of the beams but head-on-collision with small crossing angle
  - **Crab cavities**: RF rotating the bunch to maximize the interaction area

$$L = \frac{f_{rev} \gamma}{4\pi \epsilon_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*}$$



Luminosity versus focusing in case of constant beam current

- If these ideas work, one could gain a factor 4-5 with Nb-Ti or Nb<sub>3</sub>Sn, and get to a **factor 10 with ultimate intensity**
- Challenge: D0 **integration**, proof of crab cavity

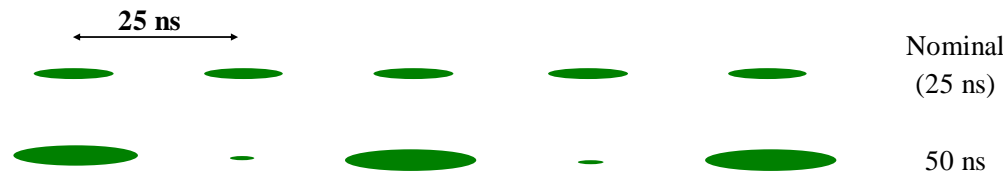


# UPGRADE BASED ON LARGER CURRENT (LARGE PIWINSKI ANGLE SCHEME)

$$L = \frac{f_{rev} \gamma}{4\pi \epsilon_n} (N_b)^2 n_b \frac{F(\beta^*)}{\beta^*}$$

$$\xi = \frac{r_p}{4\pi} \frac{N_b}{\epsilon_n} F(\beta^*) < 0.015$$

- Reduce  $F$  and increase  $N_b$  keeping the beam-beam limit - we will gain in luminosity since it is proportional to  $F (N_b)^2$
- Hypothesis: **focusing not possible** below  $\beta^*=25$  cm  $\rightarrow F$  is made small by **increasing** bunch length  $\sigma_z$  by 60% and doubling bunch space



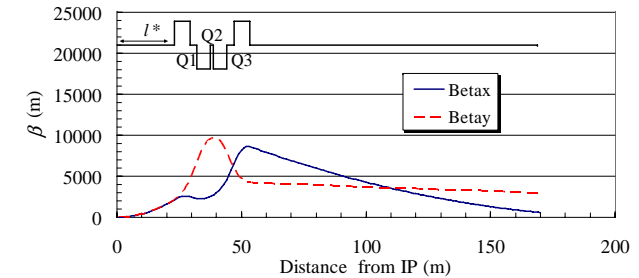
$$F(\beta^*) = \frac{1}{\sqrt{1 + \frac{1}{\beta^*} \frac{\gamma}{\epsilon_n} \left( \frac{\vartheta(\beta^*) \sigma_z}{2} \right)^2}}$$

- $N_b$  increased by 4.2  $\rightarrow N_b=4.9 \times 10^{11} \rightarrow$  gain a factor 16
- $n_b$  decreased by 2  $\rightarrow n_b=1404 \rightarrow$  lose a factor 2
- Challenges:
  - Upgrade of the injectors needed to provide this large bunch intensity ( $\times 4$ )
  - 50% larger beam current: need **long-range beam-beam wire compensation**
  - Machine protection, collimation

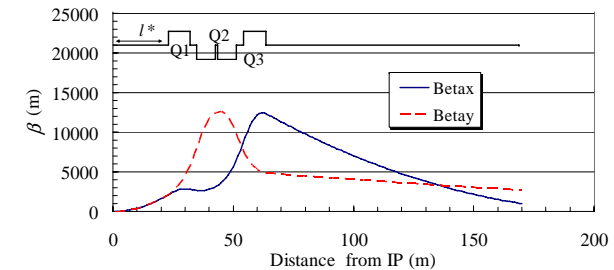


# TECHNOLOGIES: IR TRIPLET

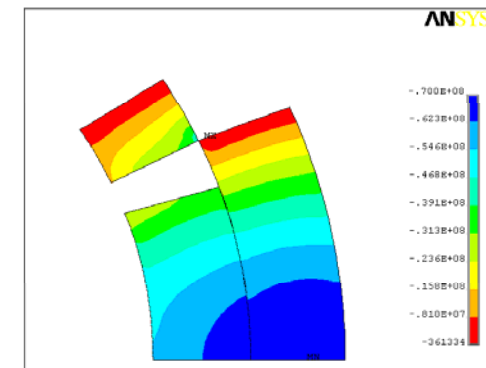
- A larger (longer) triplet
  - Aim: have a **larger aperture** to be able to go at  $\beta^*=25$  cm or down to the limit imposed by chromaticity ( $\beta^*=14-17$  cm)
- Two solutions
  - **Nb-Ti magnets** around 130 mm aperture, with a total triplet length of 40 m (10 m more than today) – to be used for phase I
  - **Nb<sub>3</sub>Sn magnets**, a bit wider, more compact, to be used for phase II
    - Smaller  $\beta^*$
    - **Better tolerance to energy deposition**
- General challenges
  - Large aperture, **large stress**
  - **Energy deposition**
  - Good field quality



Today baseline of IP



Upgrade with 40m Nb-Ti triplet



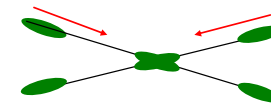
Estimated forces in the coil



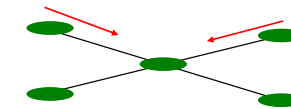
# TECHNOLOGIES: CRAB CAVITY AND EARLY SEPARATION DIPOLE

## ● Crab cavity

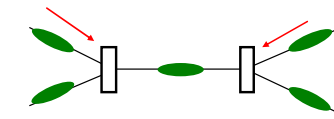
- Aim: kill the geometrical reduction factor that reduces luminosity for  $\beta^* < 25$  cm
- Idea: the bunch is rotated longitudinally to maximize the collision area
- Status: tested at KEK on **electron machine**: cavity works but no improvement on beam. More investigation and test needed. Scaled up for proton possible but not trivial.



Collision with finite crossing angle



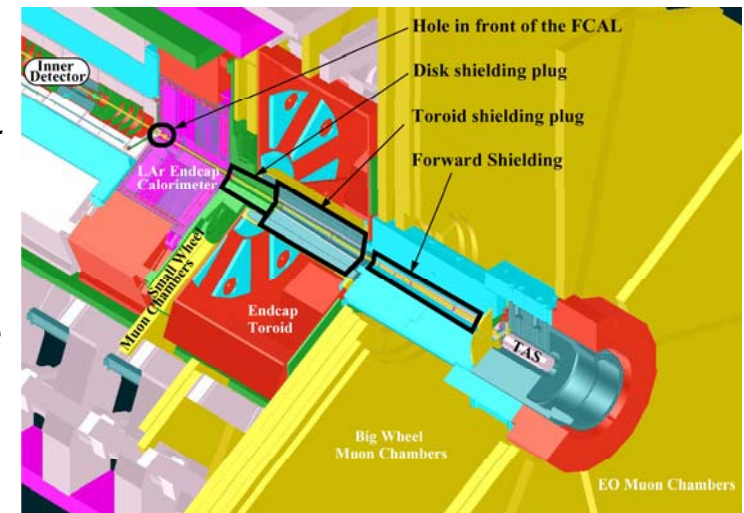
Collision with finite crossing angle and crab cavity



Collision with finite crossing angle and separation dipole

## ● Early separation dipole

- Aim: as crab cavity
  - Idea: Have zero crossing angle but separate the beams as soon as possible to avoid parasitic beam-beam interaction with a dipole ( $\sim 5$  Tm)
  - Challenges: has to be in the detector, **in a high radiation environment**
  - **It's a small magnet: you can take as a consumable**
  - Status: integration studies ongoing
- Each technologies could not completely set  $F=1$   
→ **both could solve it completely**



Positions where D0 could be integrated



# CONCLUSIONS

- Phase I upgrade (~2012)
  - To recover nominal or reach ultimate by being able to go up to  $\beta^* < 25$  cm
  - New IR triplet with ~130 mm aperture needed
- Phase II upgrade (~2016)
  - Common features to both schemes
    - Go at least to  $\beta^* = 25$  cm - this already done by phase I
    - Go at least to ultimate bunch charge [ $1.7 \times 10^{11}$ ] - this need upgrade of injectors
    - For this we need High FIELD (Gradient) Magnets with heat depo capability
  - Scenario I: early separation, i.e., strong focusing
    - Going to ultimate bunch charge
    - Push to  $\beta^*$  down 14 cm with Nb<sub>3</sub>Sn triplet
    - Remove the adverse effect of crossing angle by
      - Crab cavity
      - Early separation dipole - High FIELD Magnets with huge heat depo capability
  - Scenario II: Large Piwinski angle
    - Go ~3 times the ultimate bunch charge [ $4.7 \times 10^{11}$ ]
    - Limit with beam-beam wire compensation the effect of larger charge



## EXPENDITURE 2007

- 334,240CHF (commissioned research)
- 1,920 CHF Trip Tatsushi Nakamoto to RAL (CARE-NED meeting)
- 5,495 CHF Trip Tatsushi Nakamoto to America labs for high field magnet technology
  
- Balance: 1,847,851 CHF