

# A Baryonic Impact Factor

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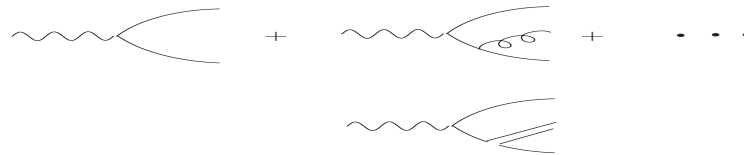
Krakov, October 19, 2007

- Introduction
- The baryon wave function
- Decomposition of the Impact Factor
- Evolution
- Conclusions

Collaboration with **Leszek Motyka**

## Introduction

Very popular: color dipole picture (large- $N_c$ ), much used in DIS:



Question: how much of the dipole picture can be used in hadron-hadron (hadron-nucleus) scattering? Obviously:

- proton is 3 quark state - not a dipole
- In large- $N_c$ : proton would be  $N_c$ -quark state

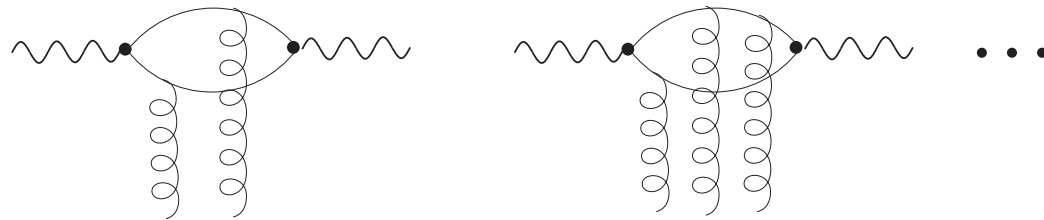
About ten years ago ([Praszalowicz, Rostworowski](#)):

Problem with color dipole picture for the proton,

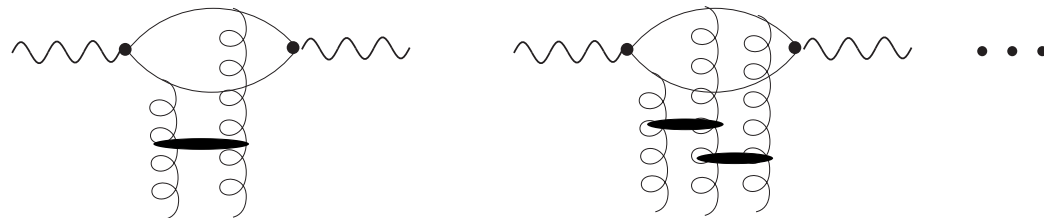
'Each step in rapidity evolution creates new color configuration'

Two steps: baryon impact factor and evolution.  
How to address: t-channel approach, BFKL approximation

Impact factor:

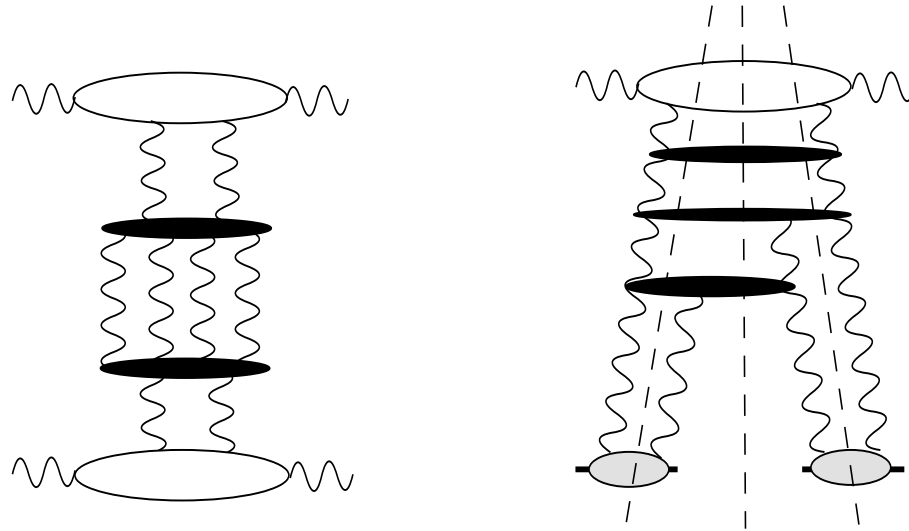


Gluon emission  $\rightarrow$  evolution:



Repeat the analysis of DIS: use virtual photon  $\rightarrow$  heavy baryon.

How to perform systematic study: consider scattering on nucleus, multiple energy discontinuity:

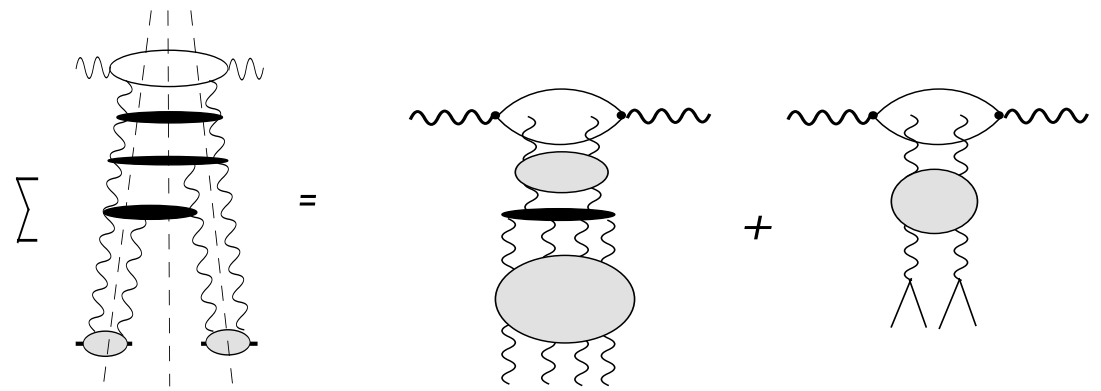


Advantage: independent energy variables, stay in leading log approximation.

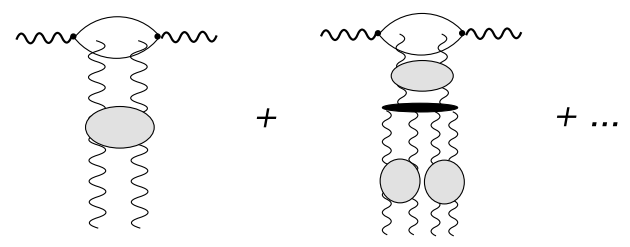
In order to go from multiple discontinuities to amplitudes: use signature factors, e.g.

$$\text{disc } T(s, t) = s \int \frac{d\omega}{2\pi i} s^\omega F(\omega, t) \leftrightarrow T(s, t) = s \int \frac{d\omega}{2\pi i} s^\omega \frac{e^{-i\pi\omega} - 1}{\sin \pi\omega} F(\omega, t)$$

In the photon case one finds a remarkable simplification: after reordering (bootstrap), a fan-like structure emerges:



Reggeization, bootstrap.  
 Contains elastic unitarity. Large  $N_c$ : BK equation.

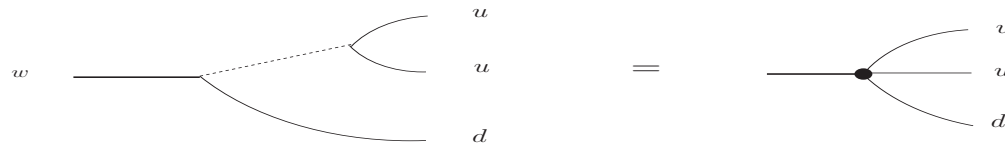


In the following: apply the same analysis to a baryonic impact factor.

## The baryon wave function

A few technicalities:

start from 4-Fermi operator (Ioffe):

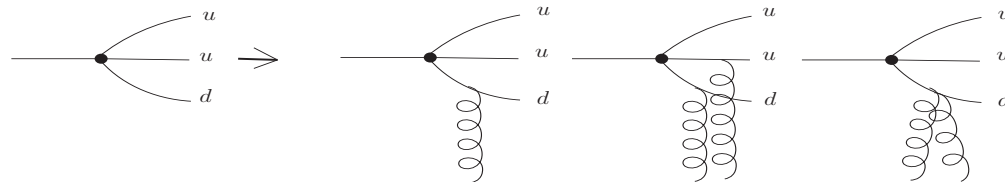


$$\epsilon_{abc} u(x)^{aT} C \gamma^\mu u(x)^b w(x) \gamma_5 \gamma_\mu d(x)^c$$

More realistic: include nonperturbative wave function;

Borel transform, from QCD sum rules.

Use helicity basis, infinite momentum frame, eikonal approximation:

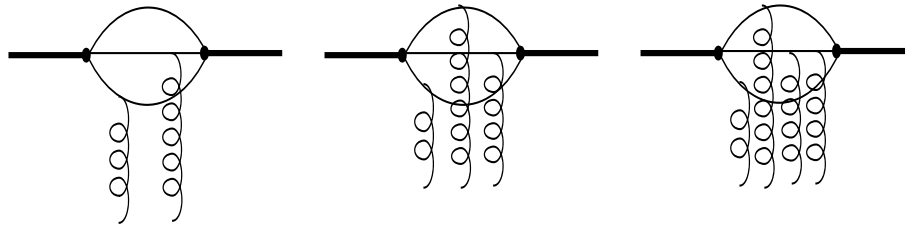


$$\Theta_{\lambda}^{(\lambda_1, \lambda_2)}(\{\alpha_i\}, \{\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3\}, \mathbf{P}) = \lambda \mathcal{N} \frac{2 \sqrt{\alpha_1 \alpha_2 \alpha_3}}{M^2 + \mathbf{P}^2 - \frac{\mathbf{p}_1^2}{\alpha_1} - \frac{\mathbf{p}_2^2}{\alpha_2} - \frac{\mathbf{p}_3^2}{\alpha_3}} \delta_{-\lambda_1, \lambda_2} \delta^{(2)}(\mathbf{p}_1 + \mathbf{p}_2 + \mathbf{p}_3 - \mathbf{P})$$

$$\cdot \left\{ \delta_{\lambda_1, \lambda} \left[ \left( \frac{\mathbf{p}_1}{\alpha_1} - \mathbf{P} \right) \cdot \left( \frac{\mathbf{p}_2}{\alpha_2} - \frac{\mathbf{p}_3}{\alpha_3} \right) - i\lambda \left( \frac{\mathbf{p}_1}{\alpha_1} - \mathbf{P} \right) \times \left( \frac{\mathbf{p}_2}{\alpha_2} - \frac{\mathbf{p}_3}{\alpha_3} \right) \right] \right.$$

$$\left. + \delta_{\lambda_2, \lambda} \left[ \left( \frac{\mathbf{p}_2}{\alpha_2} - \mathbf{P} \right) \cdot \left( \frac{\mathbf{p}_1}{\alpha_1} - \frac{\mathbf{p}_3}{\alpha_3} \right) - i\lambda \left( \frac{\mathbf{p}_2}{\alpha_2} - \mathbf{P} \right) \times \left( \frac{\mathbf{p}_1}{\alpha_1} - \frac{\mathbf{p}_3}{\alpha_3} \right) \right] \right\}$$

Square, sum over helicities, compute color traces, sum over all possibilities of attaching gluons:



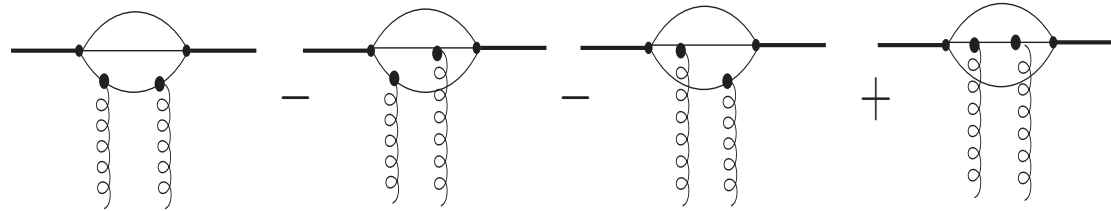
Results: decompose the impact factor into irreducible (under evolution) and gauge invariant pieces:

$$T = T_{dipole} + T_{odderon} + T_{new}$$

Main task: color structure.

## Decomposition of the impact factor

A. Two Gluons (C even): 'normal' dipole structure. The pair (23):



Dipole structure, but lines 2 and 3 are in antitriplet.

'Antitriplet dipole'

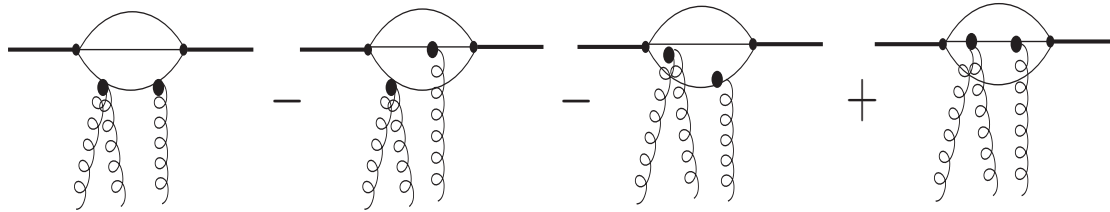
Contains diquark configuration, depends upon dynamics of wave function.

$$\left( D_{20}^{23}(k_1, k_2) + D_{20}^{13}(k_1, k_2) + D_{20}^{12}(k_1, k_2) \right) \delta_{ab}$$

Satisfies Ward identities.



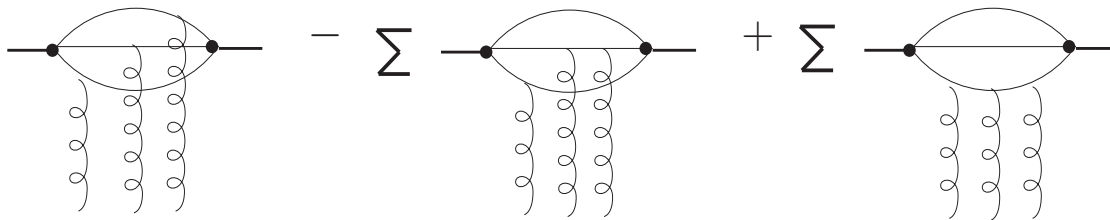
B. Three gluons (C even): reggeization of the 2 gluon system



$$\left( D_{20}^{23}(k_1 + k_2, k_3) + D_{20}^{23}(k_1 + k_3, k_2) + D_{20}^{23}(k_1, k_2 + k_3) \right) f_{abc} + D_{20}^{12} + D_{20}^{13}$$

Reggeization + bootstrap, similar to photon impact factor.

Three gluons (C odd): Odderon (C.Ewerz)

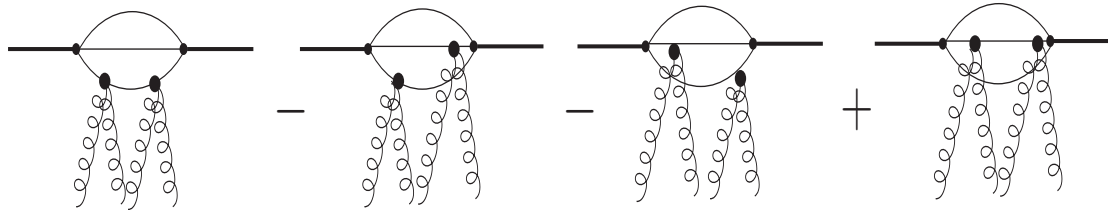


New function:  $E_{30}(k_1, k_2, k_3)$ . Satisfies Ward identities. Nonabelian charge configuration.

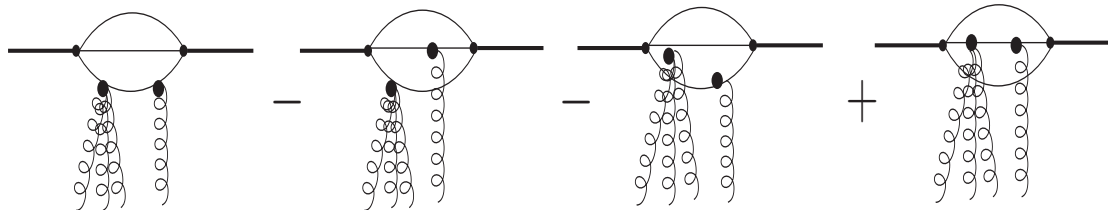
C. Four gluons (C even): different pieces

- reggeization of the 2 gluon system
- new configuration

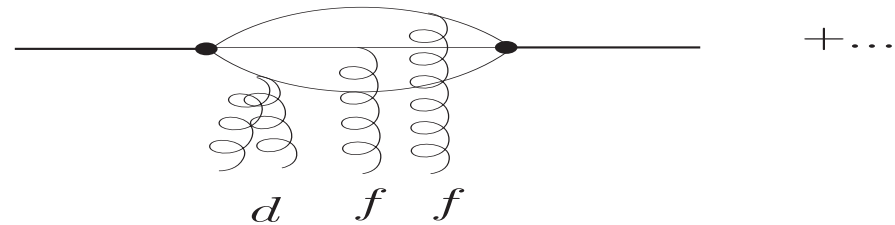
$$D_{20}^{23}(k_1 + k_2, k_3 + k_4) :$$



$$D_{20}^{23}(k_1 + k_2 + k_3, k_4) :$$



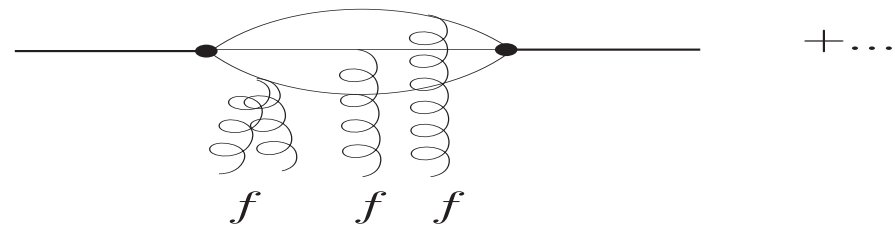
The new configuration: C even, even signature



$$Q_{4,0}(1, 2, , 3, 4) = \left( d_{a_1 a_2 c} d_{a_1 a_2 c} - \frac{1}{3} \delta_{a_1 a_2} \delta_{a_3 a_4} \right) \left[ E_{30}(1 + 2, 3, 4, ) + E_{30}(1, 2, 3 + 4) \right]$$

Sum over all permutations satisfies Ward identities. Needs all 3 quarks.

Four gluons (C odd): reggeization of the odderon (C.Ewerz)



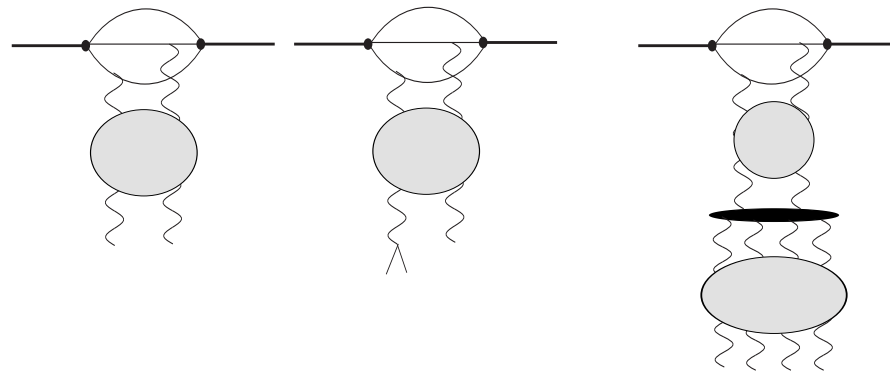
$$E_{30}(k_1 + k_2, k_3, k_4) d_{a_1 a_2 c} f_{a_1 a_2 c}$$

S-channel picture: Fourier transform to transverse coordinates.

## Evolution, gluon radiation

The decomposition of the impact factor is preserved under evolution.

Dipole-like term (color anti-triplet):

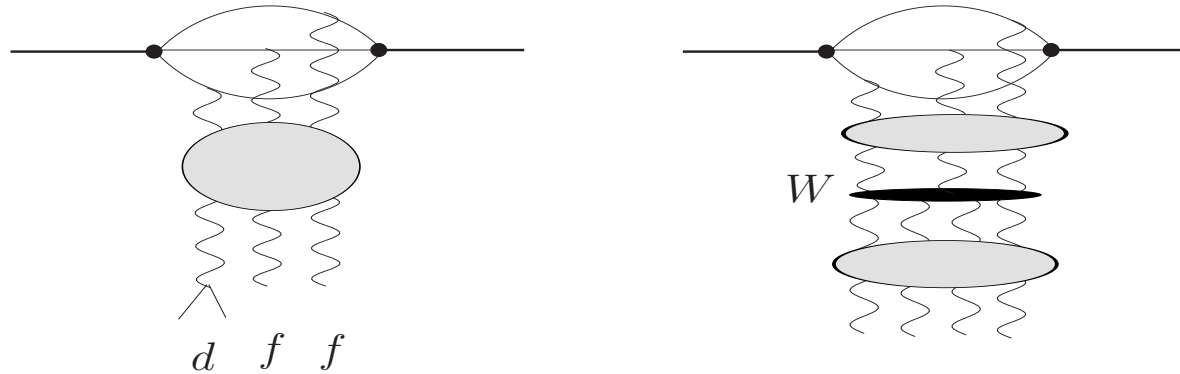


Looks like photon impact factor, fan structure... (BK equation?)

Contains the diquark configuration, but it also allows for spacial separation.

The new piece:

a new vertex appears (good properties, e.g. Möbius invariant):



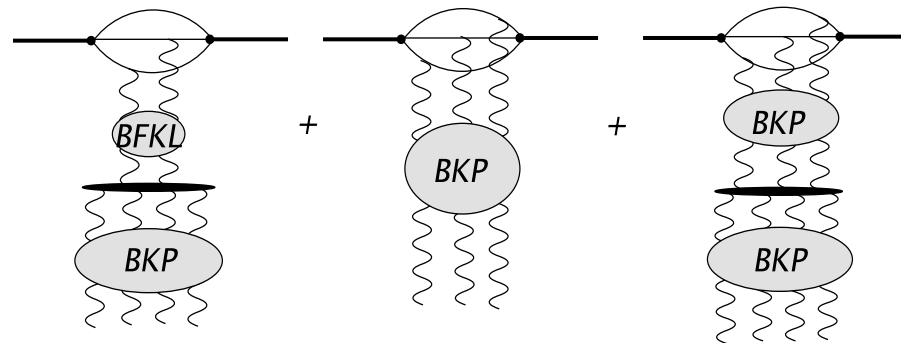
s-channel picture: radiation from all three quarks, new evolution kernel.

In addition:

C-odd, Odderon: both the BLV and the WJ solutions couple.

## Summary and Discussion

In comparison with photon or meson (=color dipole), baryon has more complex pattern of gluon radiation:



+ reggeizing pieces

- Baryon contains dipole-like configuration (diquark): same pattern as in the photon case, but with relative coupling  $1/2$ .
- Baryon couples to the Odderon ( $C$ -odd), both WJ and BLV solutions

- New vertex: transition from  $C$ -even three gluon state  $\rightarrow$  two-Pomerons
- In LO no direct coupling of two Pomerons: 'disappearance of elastic intermediate state' (reggeization and bootstrap).

Can we neglect the 'new' contributions? Most likely not.

Unitarity requires:

$$|T(s, b)| = \left| \frac{1}{2} \sum T_{dipole} \pm T_{odderon} + T_{new} \right| \leq 1$$

Assume: each dipole-like configuration saturates at high energies

$$T_{dipole} \rightarrow 1,$$

sum of three dipole configurations violates bound!

New contribution is needed for unitarity!

Deeper connection between all pieces.



Consequences:

QCD reggeon field theory contains more than BFKL Pomerons + interactions, importance of three gluon state (not necessarily odderon!) and  $d$ -reggeons.

Interesting possibility: connection with Casimir Operators.

In  $SU(2)$ : two Casimirs, baryon couples to states with two gluons (BFKL) and three gluons (Odderon).

Two 'fundamental excitations'.

In  $SU(N)$ :  $N - 1$  Casimir operators.

Baryon consists of  $N$  quarks, couples to  $t$ -channel states with 2, 3, ... $N$  gluons (in fundamental representation):

Baryon excites all the 'fundamental excitations of high energy QCD'.

Applications:

- Model for nucleon (spin and color structure).
- elastic  $pp$  scattering at intermediate  $t$  (ISR: 3 gluon model)
- Spin physics?

Most important:

Think about finding solutions to QCD reggeon field theory!