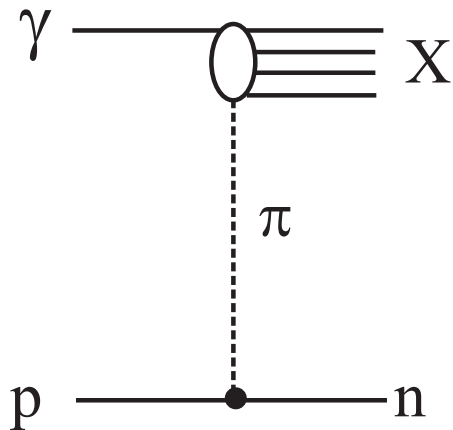


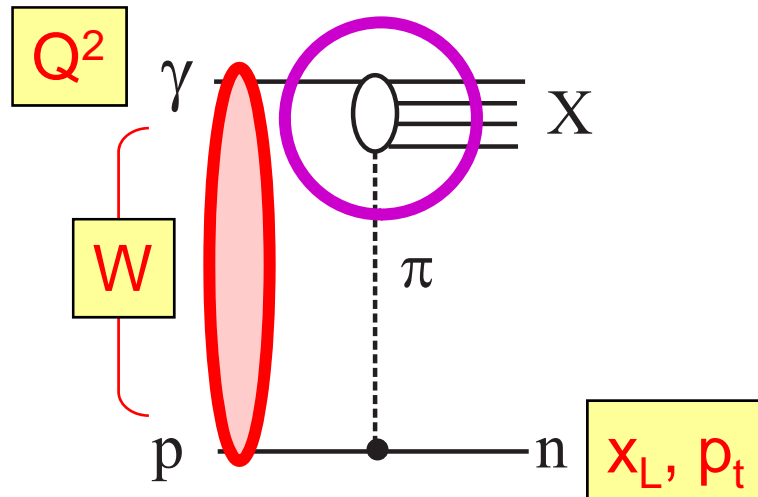
# Leading neutrons and protons

V.A. Khoze, A.D. Martin, M.G. Ryskin

For Forward Physics at LHC it is useful to start with **leading neutrons** observed at HERA --- prelimin. ZEUS data --- good example of problems to be faced.



Alan Martin (Durham)  
Forward Physics Workshop  
Manchester, Dec. 2006



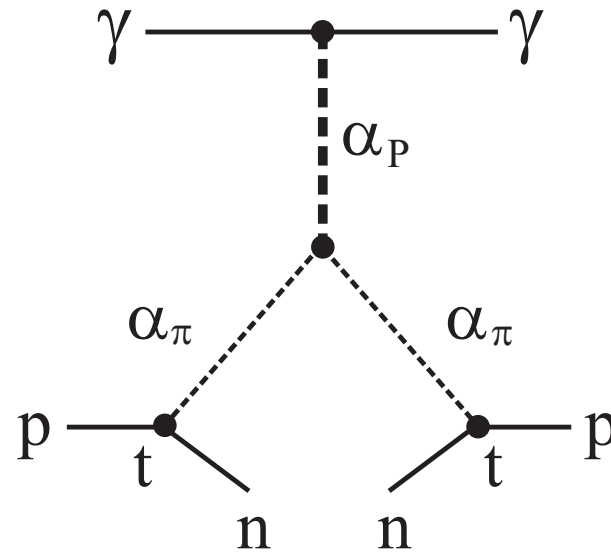
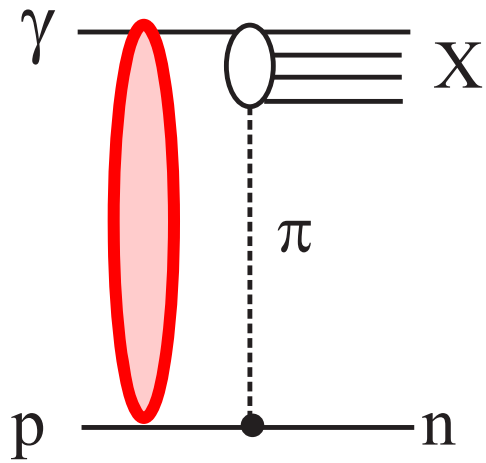
$\pi$  exchange dominates  
for  $x_L > 0.6$ ,  
but absorptive effects

Leading neutron data  $(Q^2, x_L, p_t) \rightarrow$

$\pi$  structure fn,  $F_2^\pi(x, Q^2)$  at small  $x \rightarrow f_{q,g}^\pi$

Here, we study re-scatt. effects in photoprod:

$\sigma_{\text{abs}}(q\bar{q}-N) \rightarrow \text{check of } S^2$



$$\frac{d\sigma(\gamma p \rightarrow XN)}{dx_L dt} = S^2 \frac{G_{\pi+pn}^2}{16\pi^2} \frac{(-t)}{(t - m_\pi^2)^2} F^2(t) (1 - x_L)^{1-2\alpha_\pi(t)} \sigma_{\gamma\pi}^{\text{tot}}(M^2)$$

gap survival factor

$S^2 \sim 0.48$  for photoproduction

$$\sigma_{\gamma\pi}^{\text{tot}} = \frac{2}{3} \sigma_{\gamma p}^{\text{tot}}$$

Add<sup>ve</sup>QM:

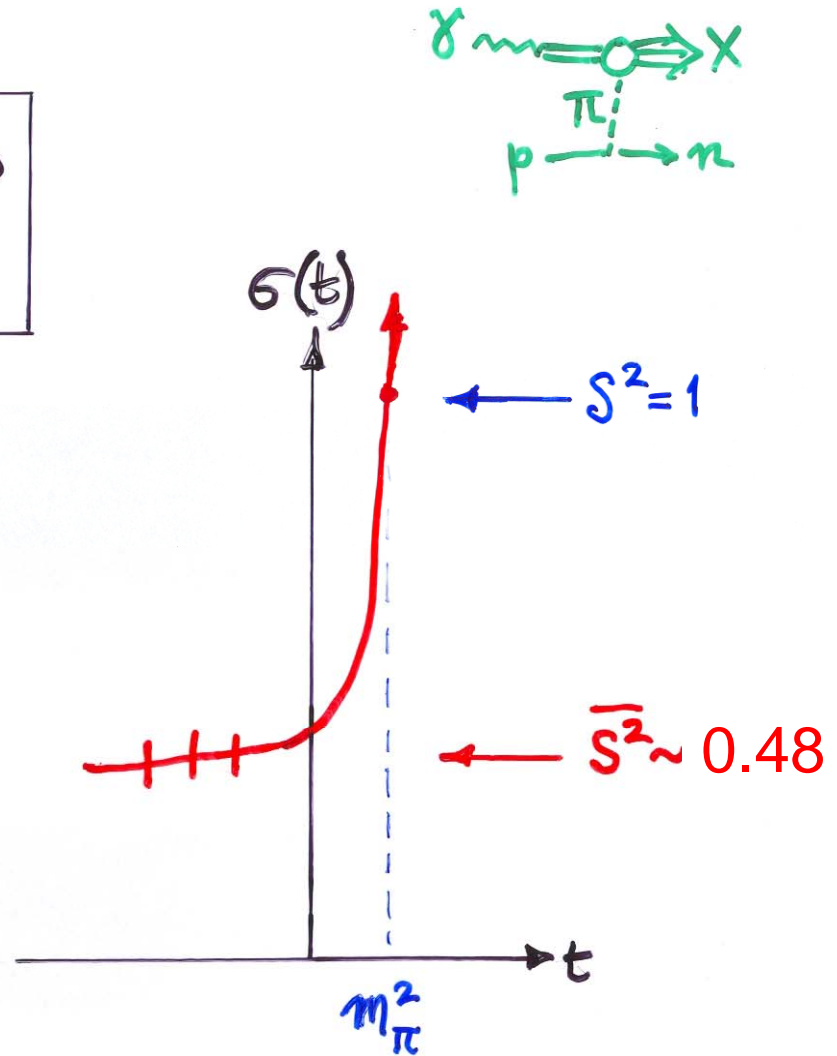
$$\overline{S^2} \sim \frac{\int e^{-\Omega(s,b)} d\sigma_0 d^2b}{\int d\sigma_0 d^2b}$$

$$\sigma(q\bar{q}-N) \sim \sigma(pp) \sim \sigma(\pi p)$$

$$\Omega(s,b) = \frac{\sigma(s)}{4\pi B} e^{-b^2/4B}$$

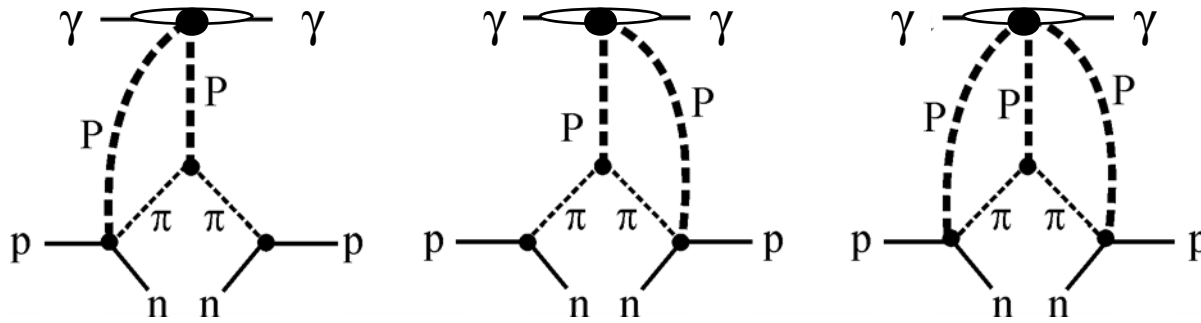


$$T_{el} \sim e^{Bt}$$



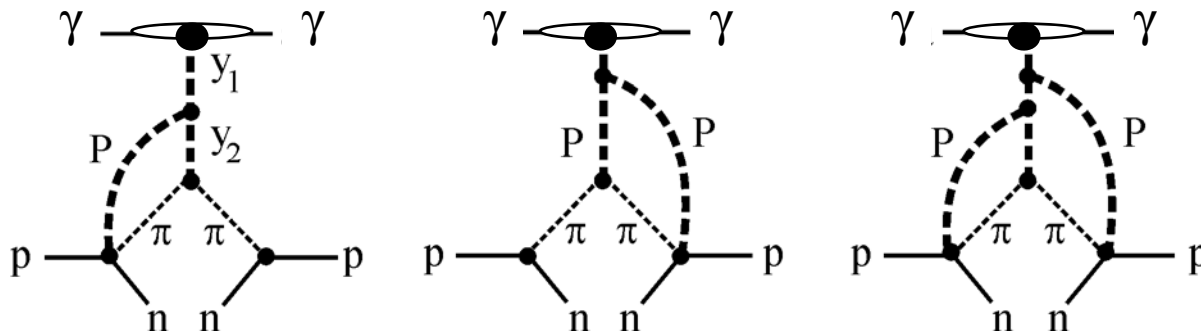
# Calculation of survival factor, $S^2(x_L, p_t^2, Q^2)$

(a)



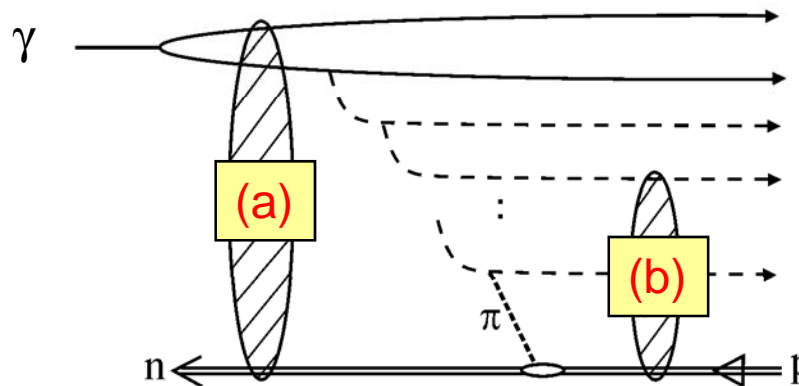
normal eikonal diagrams

(b)



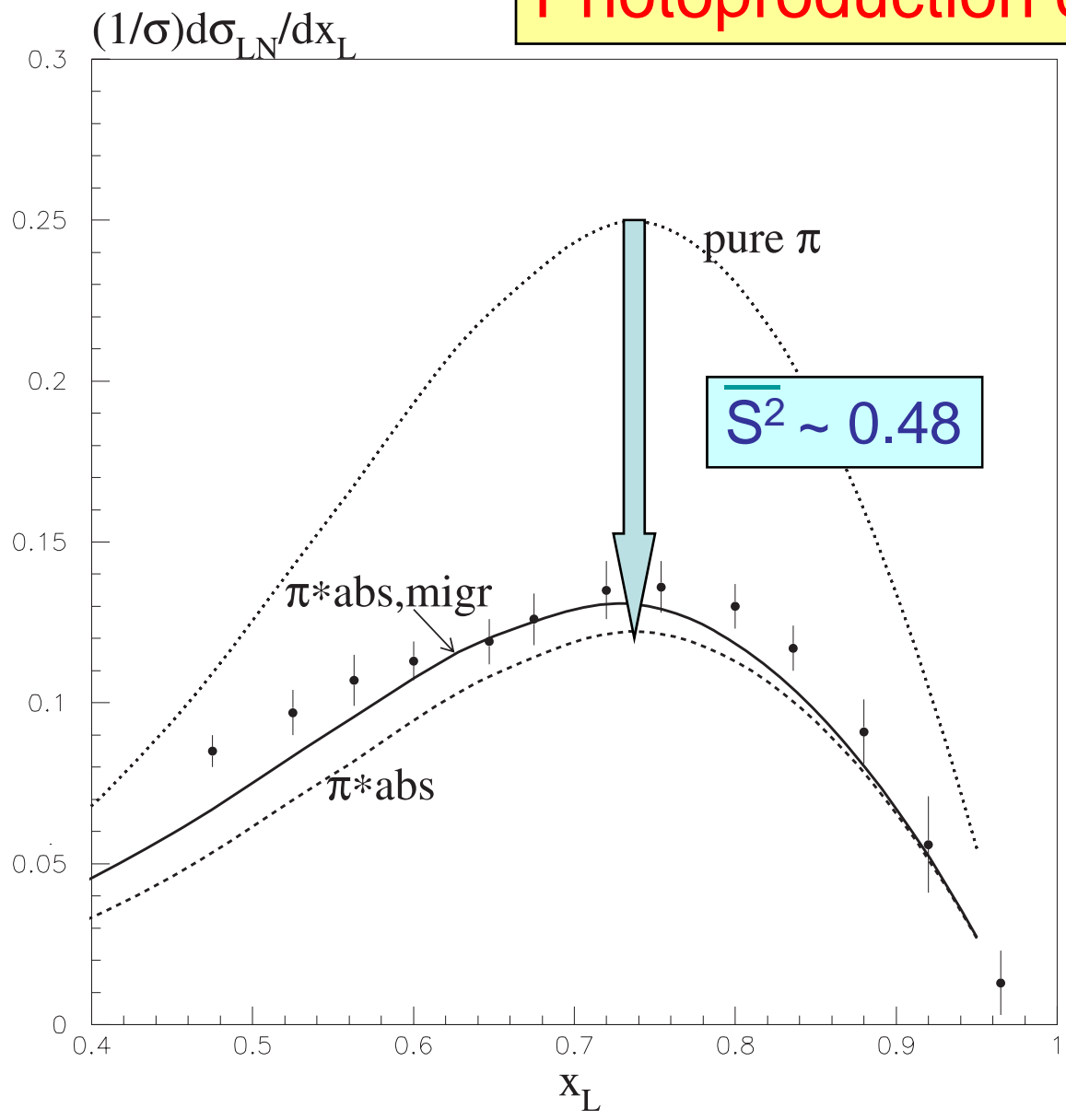
enhanced diag.  
need  $y_i > 2-3$ :  
result  $\sim 15\%$

space-time picture

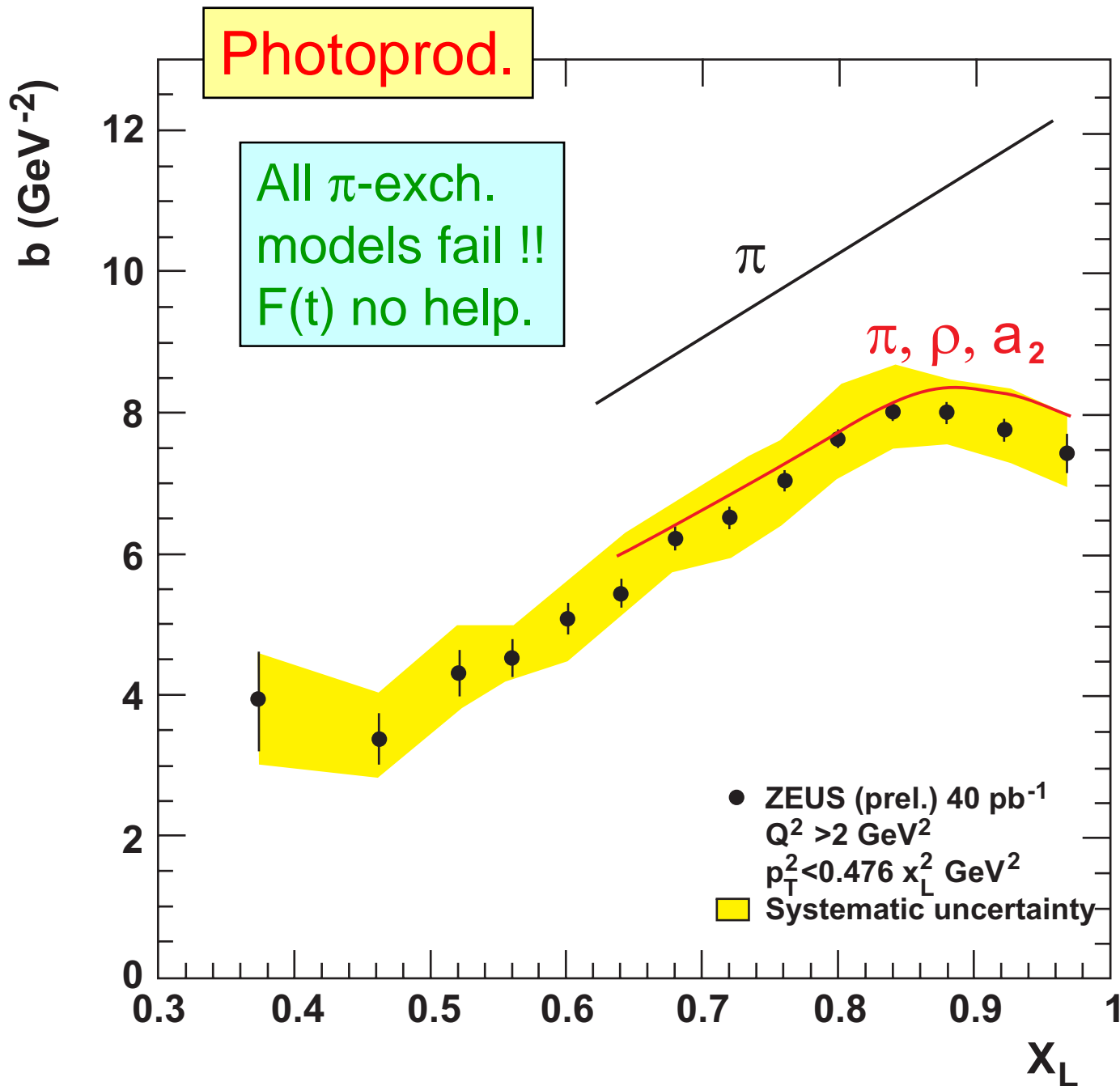


If enhanced diag were important, then n yield would be energy dep.  
Not seen in data.

# Photoproduction of leading n



Prelim. ZEUS data  
(DIS2006)



$p_t^2$  dep.  
 $\sim \exp(-b p_t^2)$

led to include  $\rho, a_2$ -exch. →

Including  $\rho$  and  $a_2$  exchange

Irving,  
Worden

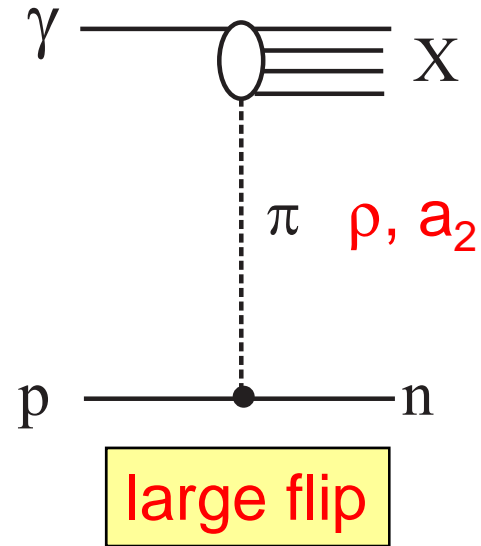
$$\frac{V_{\rho}^{\text{flip}}}{V_{\rho}^{\text{non-flip}}} \approx 8 \frac{p_t}{2m_N}$$

enlarges  $\sigma$ , more at larger  $p_t$   
leading to **smaller** slope, b

Use  $\rho$ ,  $a_2$  exch. degeneracy, additive QM

Slope b now OK ---  $\sigma$  too large --- adjust parameters to  
attempt to **simultaneously describe  $\sigma$  and b**

diff<sup>ve.</sup>  $\rightarrow$  1.3  $\left[ \sigma(q\bar{q}-p) \sim \sigma(\rho p) \sim \sigma(\pi p) \sim 31 \text{ mb} \right]$   
excit. 1.6 34mb

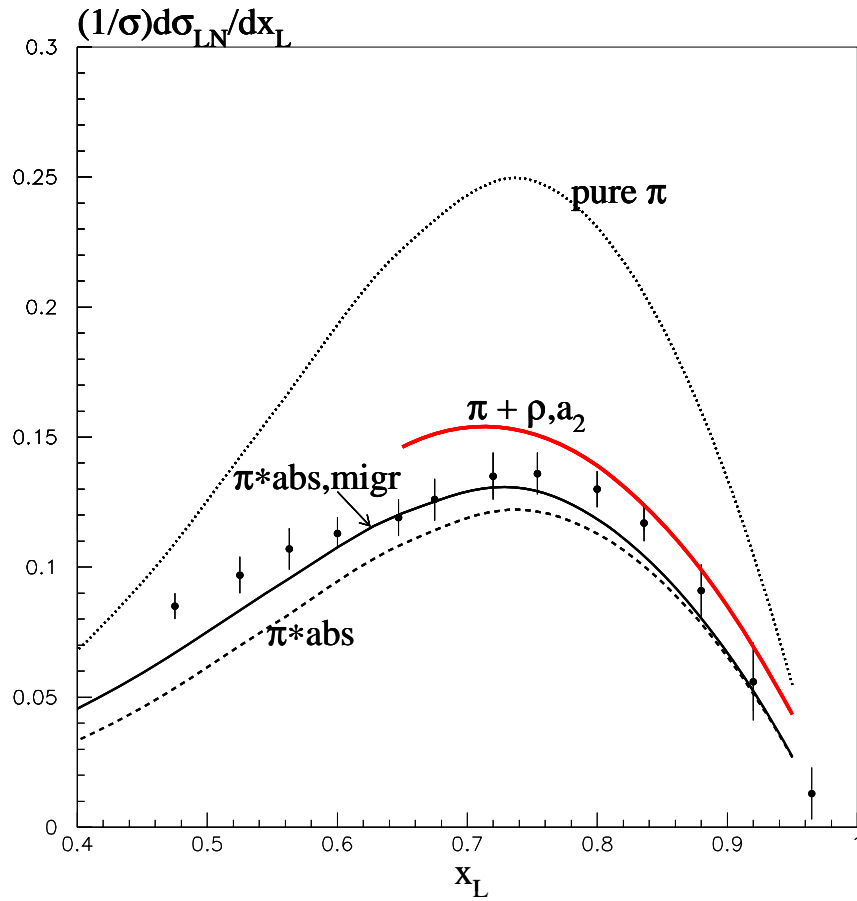


$$\bar{S}^2 \sim 0.48$$

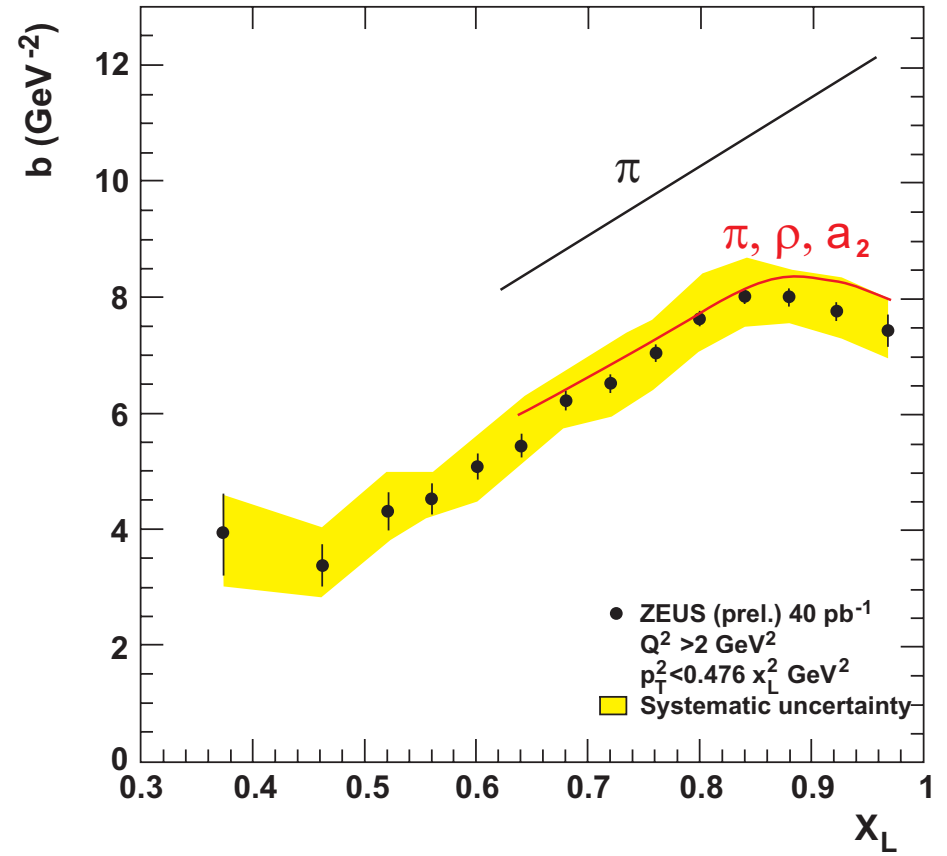
$$\text{now } \bar{S}^2 \sim 0.4$$



# cross section



# $p_t^2$ slope



now  $\overline{S^2} \sim 0.4$ , instead of 0.48

## Conclusions on leading neutrons at HERA

- Exploratory study of prelim. ZEUS data ( $Q^2$ ,  $x_L$ ,  $p_t$ ,  $W$ ) informative for forward physics at LHC
- $\pi$  exch (with abs.) describes  $\sigma$ , but not  $p_t^2$  slope  $b$   
→ need also  $\rho$ ,  $a_2$  exchange

turnover of slope as  $x_L \rightarrow 1$  ( $t_{\min} \rightarrow 0$ ) may be used to determine  $\rho, a_2$  versus  $\pi$  exchange contributions

- Absorptive corrections important  $\bar{S}^2 \sim 0.4$   
Small contrib. from enhanced diagrams
- Simultaneous description all data ( $Q^2$ ,  $x_L$ ,  $p_t$  dep.) difficult
- This is good. Precise LN data should determine  $F_2^\pi(x, Q^2)$  at small  $x$  and  $S^2(x_L, p_t, Q^2)$

important  
for LHC

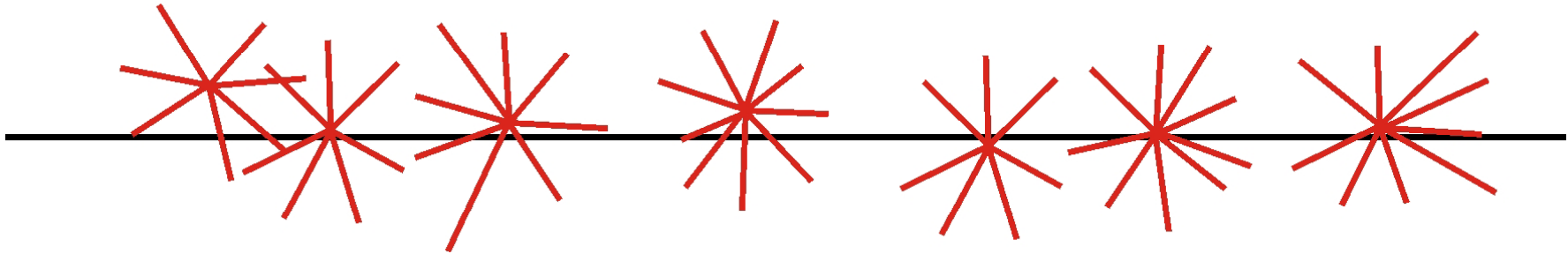
LHC

Pile-up

$$\mathcal{L} = 10^{33} \text{ cm}^{-2} \text{ sec}^{-1}$$
$$10^{34}$$

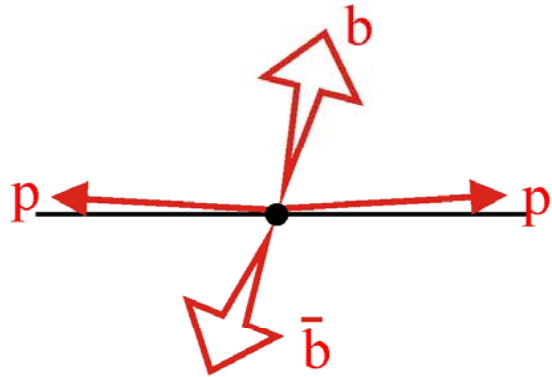
$$N = 3.5$$

$$N = 35$$



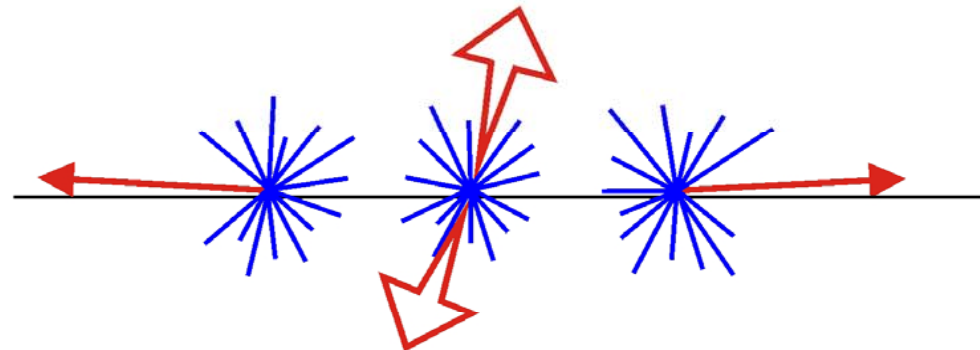
every bunch crossing !

exclusive  $H \rightarrow b\bar{b}$  signal



**1**

Background due to pile-up

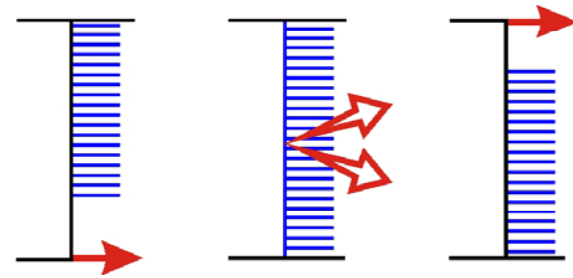


$\sim 0.02N$

$10^8$

$\sim 0.02N$

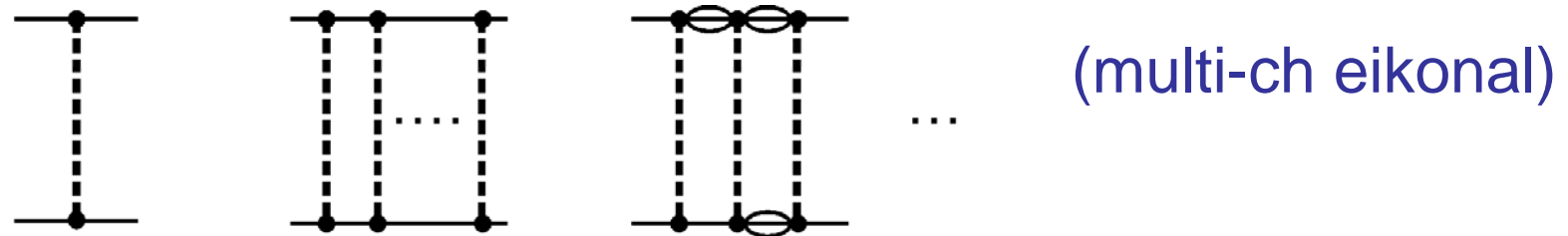
$\sim 0.01N$  for 420m  
 $\sim 0.03N$  for 220m



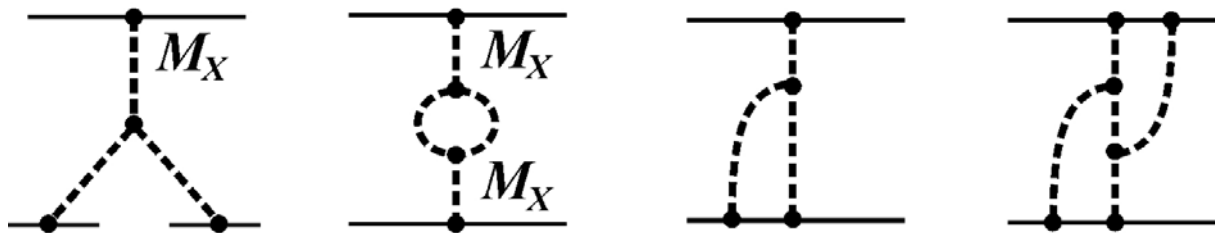
need to predict SD....see Misha Ryskin's talk

Forward physics at LHC needs predictions for leading protons

s-channel unitarity generates a whole sequence of multi-Pomeron diagrams:  $\rightarrow$  (low-mass) SD, DD

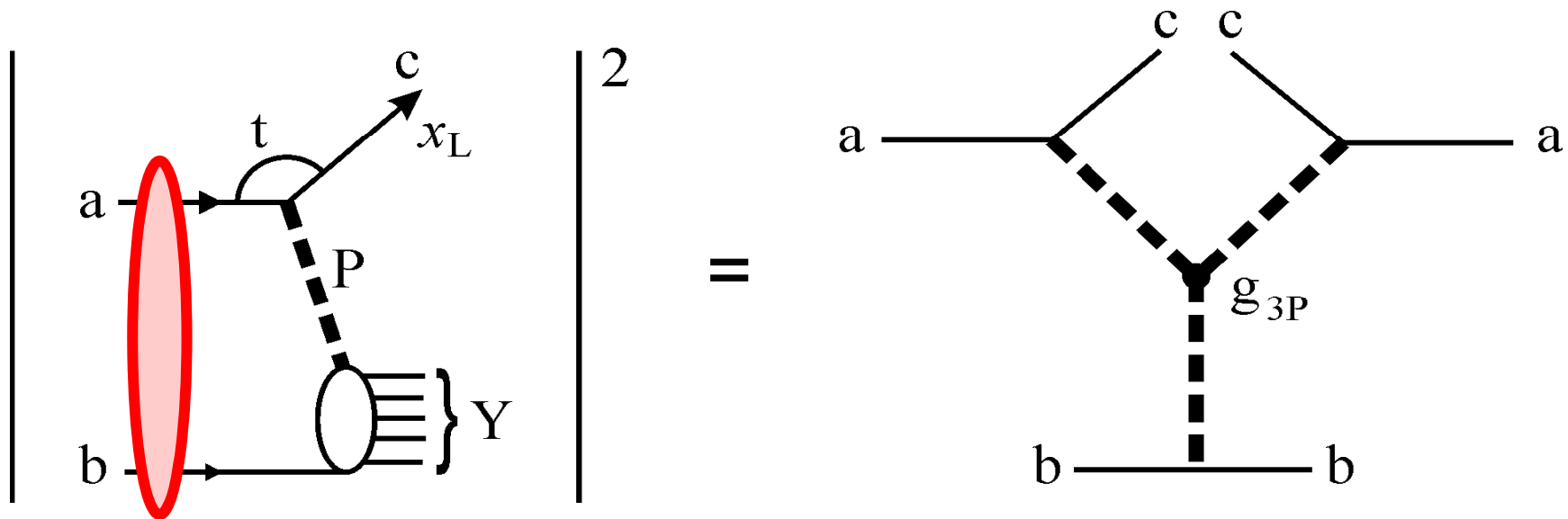


Also have non-eikonal high-mass SD, DD :



need triple-Pomeron coupling  $g_{3P}$

# Determination of triple-Pomeron coupling $g_{3P}$

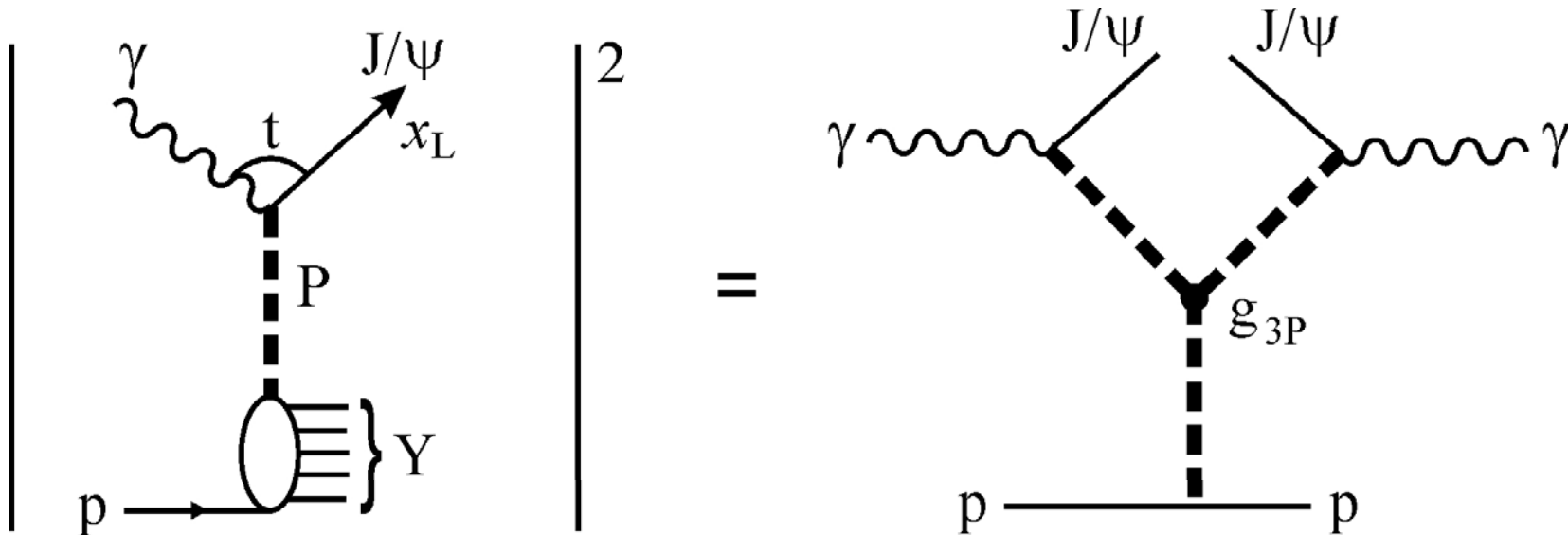


e.g.  $pp \rightarrow pY$ ,  $\pi p \rightarrow \pi Y \dots$  at large  $M_Y$  and  $s/M_Y^2$

Soft rescatt: leading hadron  $\rightarrow$  secondaries  $\rightarrow$  smaller  $x_L$   
 $\rightarrow$  populate/destroy rapidity gap

old hadron data  $\rightarrow$  **effective**  $g_{3P}$  (which embodies  $S^2$ )

# Estimate of bare triple-Pomeron coupling $g_{3P}$



$c\bar{c}$  system v. compact  $\rightarrow$  rescatt. suppressed

so need  $\gamma p \rightarrow J/\psi Y$  data as a func. of large  $M_Y$

some info. exists as bkgd to  $\gamma p \rightarrow J/\psi p$  ZEUS, H1

# ZEUS, H1

elastic  $\delta p \rightarrow J/\psi p$   
 incl.  $\Upsilon$   $\delta p \rightarrow J/\psi \Upsilon$

$$\sigma_i \sim W^{\delta_i}$$

$$\delta_{el} \sim 0.7 \pm 0.05$$

$$\delta_{\Upsilon} \sim 0.7 \pm 0.2$$

$$\sigma \sim e^{bt}$$

$$b_{el} \sim 4.5$$

$$b_{\Upsilon} \sim 0.7$$

$$\sigma_{\Upsilon} \sim (M_{\Upsilon}^2)^{0.08}$$

← expected for PPP

PPP more compact than proton

$$\delta p \rightarrow J/\psi : \left. \begin{array}{l} \sigma_{\Upsilon} \\ \sigma_{el} \end{array} \right\} \approx 1 \pm 0.3$$

$$\delta p \rightarrow \rho : \left. \begin{array}{l} \sigma_{\Upsilon} \\ \sigma_{el} \end{array} \right\} \approx 0.6 \pm 0.2$$

stronger abs. for  $\rho$

$$\frac{\int \frac{d\sigma_{\Upsilon}}{dt dM_{\Upsilon}^2} dM_{\Upsilon}^2}{\frac{d\sigma_{el}}{dt}}$$

$$\Rightarrow \frac{g_{3P}^{bare}}{g_N} \sim \frac{1}{3}$$

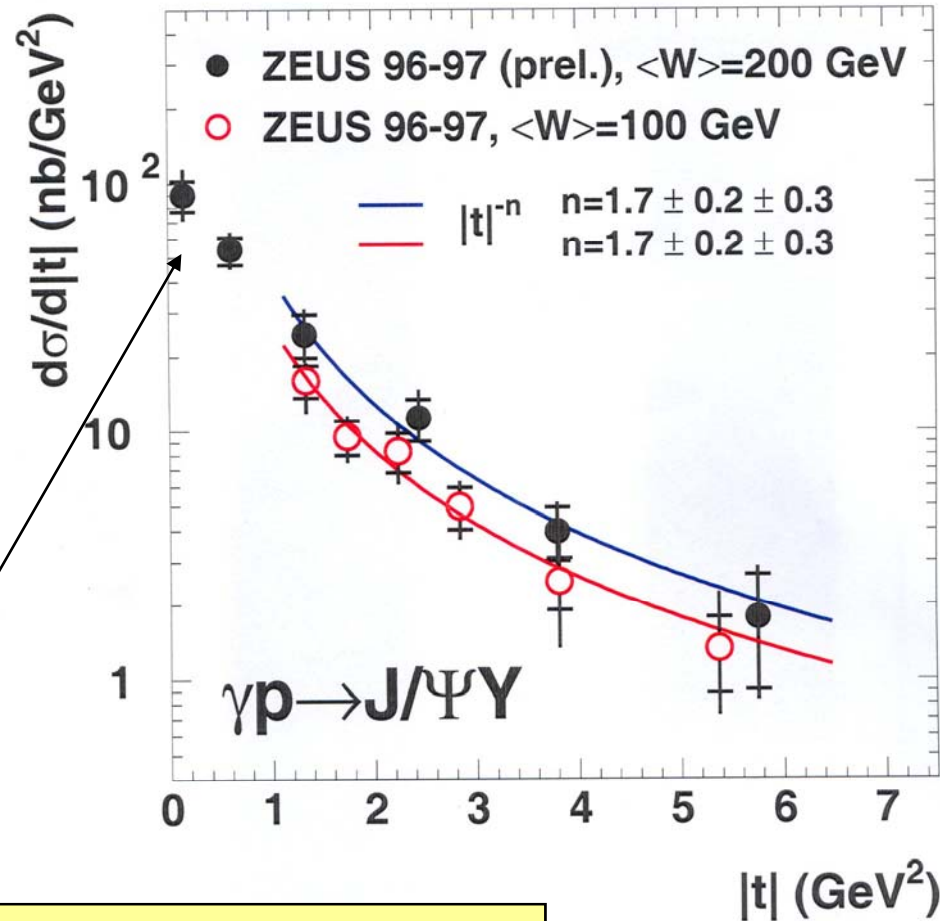
(new)

$$\left( \frac{g_{3P}^{eff}}{g_N} \sim \frac{1}{10} \right)$$

(old)



# ZEUS



first evidence that  $g_{3P}(t)$  does not vanish as  $t \rightarrow 0$ .

first exptal. evidence of “strong coupling” of Pomeron

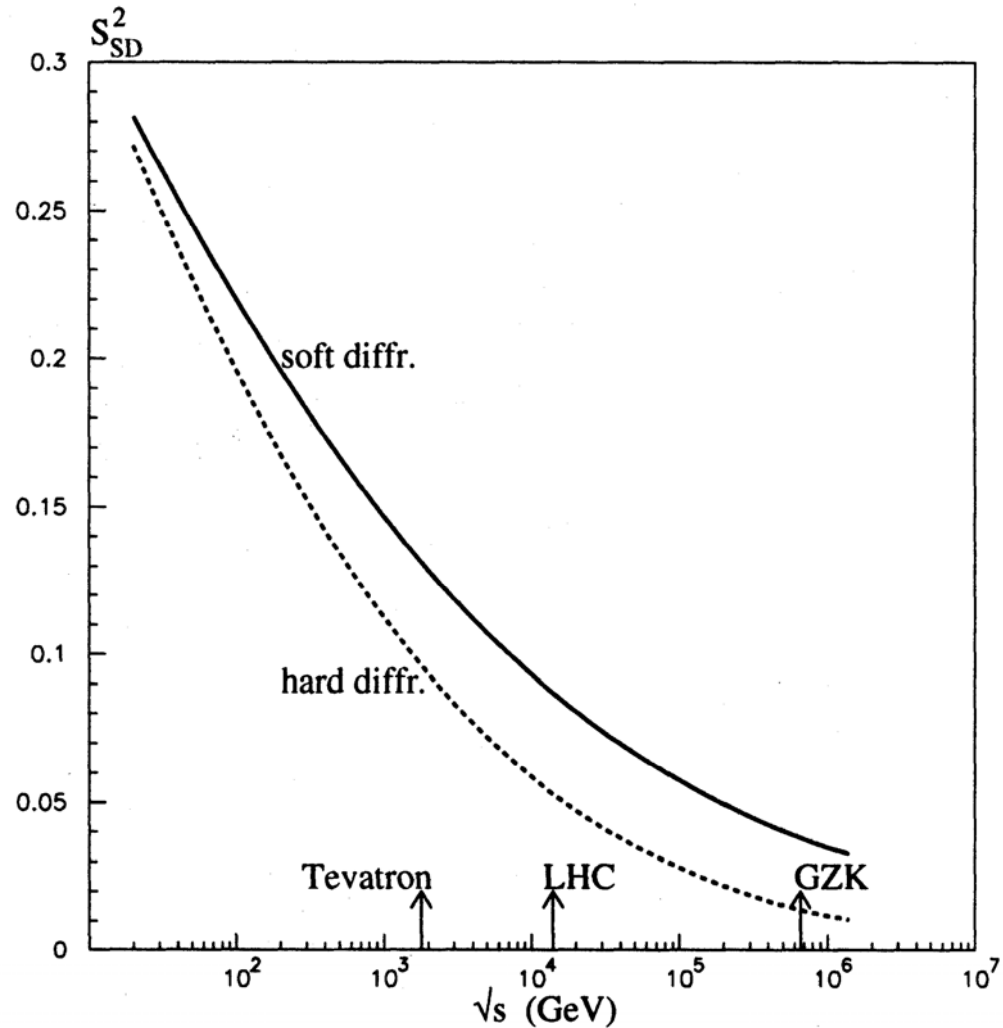
$$\frac{1}{\sigma_{\text{tot}}} \frac{d\sigma^{SD}}{dt dx_L}$$

$$\sim \frac{g_N^2(t) g_{3P}(t)}{g_N(0)} (1-x_L)^{\alpha_P(0) - 2\alpha_R(t)} S_{SD}^2(s, t)$$

factors  $\sim 0.5$   $\sim 0.5$   
 from data in going  
 $W = 20 \rightarrow 1800$  GeV

↑  
 consistent with  
 decrease of  $S^2(s)$ .

Unlike limiting  
 frag<sup>n</sup> hypothesis,  
 normalised  $\sigma$  depends  
 on energy due to  $S^2$



## Conclusions

Leading neutrons/forward physics:  $S^2$  is important  
but no effect from enhanced diagrams

Forward physics at LHC needs prediction of SD  
to quantify “pile-up” backgrounds  
see following talks

From inclusive  $\gamma p \rightarrow J/\psi \ Y$  photoprod. we estimate  
the **bare**  $g_{3P}$  coupling to be about 3 times larger  
than old hadron estimate of the **effective**  $g_{3P}$ .  
(already anticipated---but is a direct estimate)

## Conclusions continued...

Is “soft” physics a speciality at the LHC, since first interest is in high  $p_t$  leptons, photons and jets ??

No --- perhaps it is generally important !

At LHC need good knowledge of underlying event.

e.g. consider  $H \rightarrow \gamma\gamma$  with small signal on huge bkgd.  
Accurate  $\gamma\gamma$  mass is crucial – but what about  $\pi^0$ 's  
from underlying/pile-up events !! Correlations.

(Moreover some believe there is deep theoretical link between “soft” and “hard” physics.)