

# Novel detection techniques in FP420



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UCL

Intro: Detectors at 420m (for high lumi)  
HECTOR: LHC beamline simulation  
Detection system: Hamburg pipes  
Detectors: Tracking and timing  
Outlook

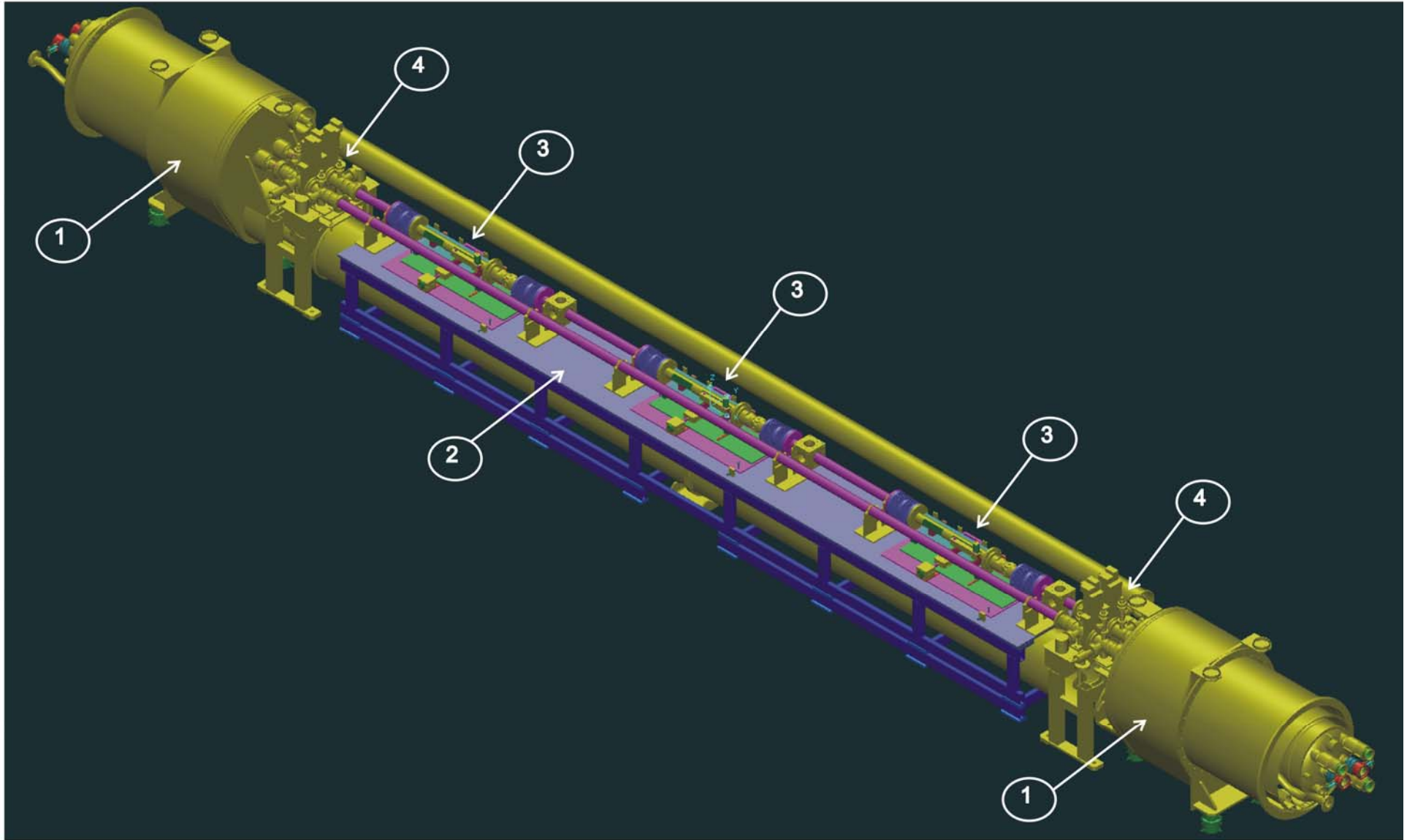
## FP420 R&D issues (NB. Report in preparation):

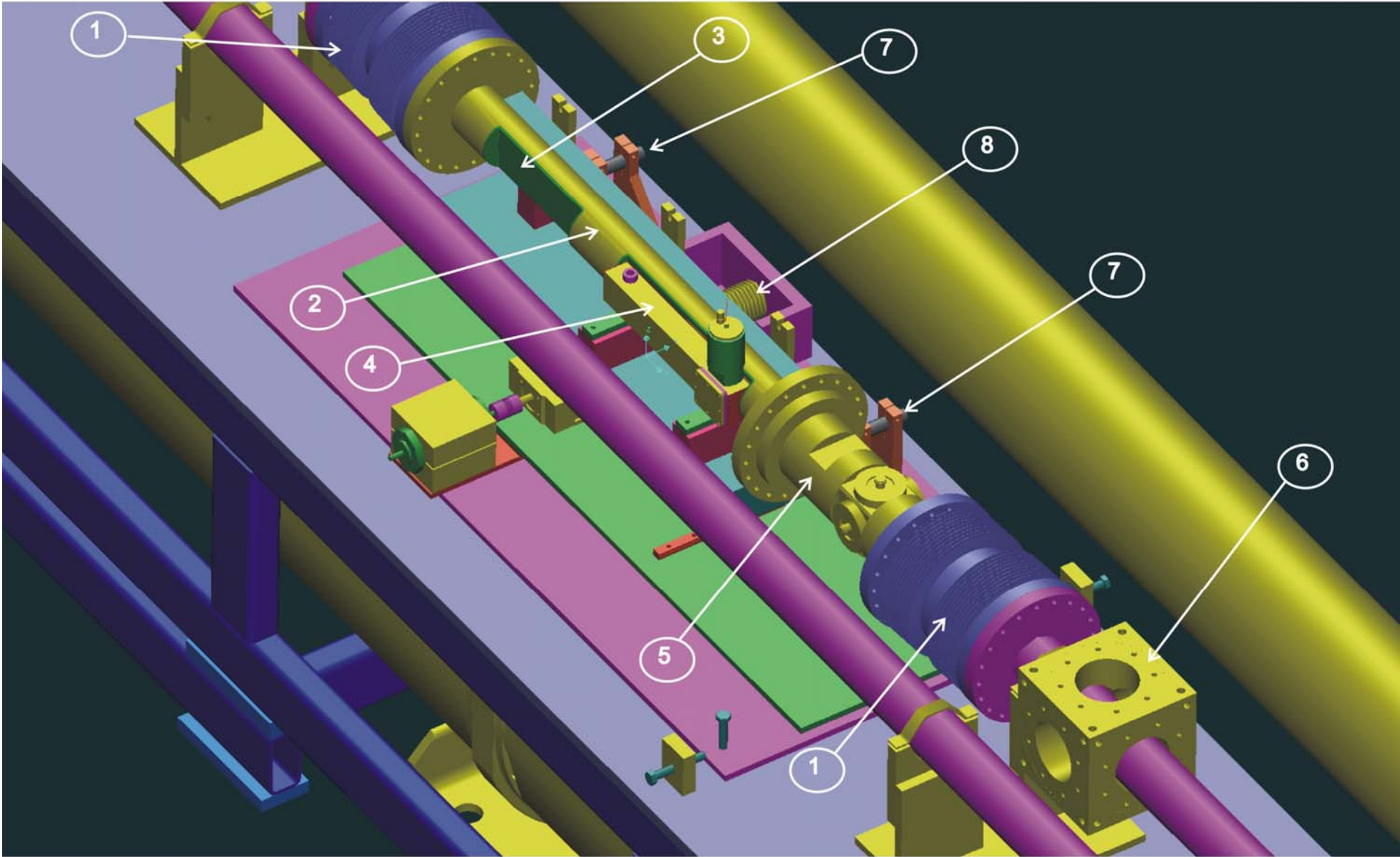
- Modification of LHC connecting cryostat to allow for near-beam detectors at ~420 m
- Optimization of detector setup for resolutions and acceptances
- Design detector system to allow for reliable and precise close approach to beam
- Study and tests detectors for precise proton tracking and timing
- Study and design required infrastructure in LHC tunnel
- Study backgrounds and radiations hardness

## Moving Hamburg pipe detector system

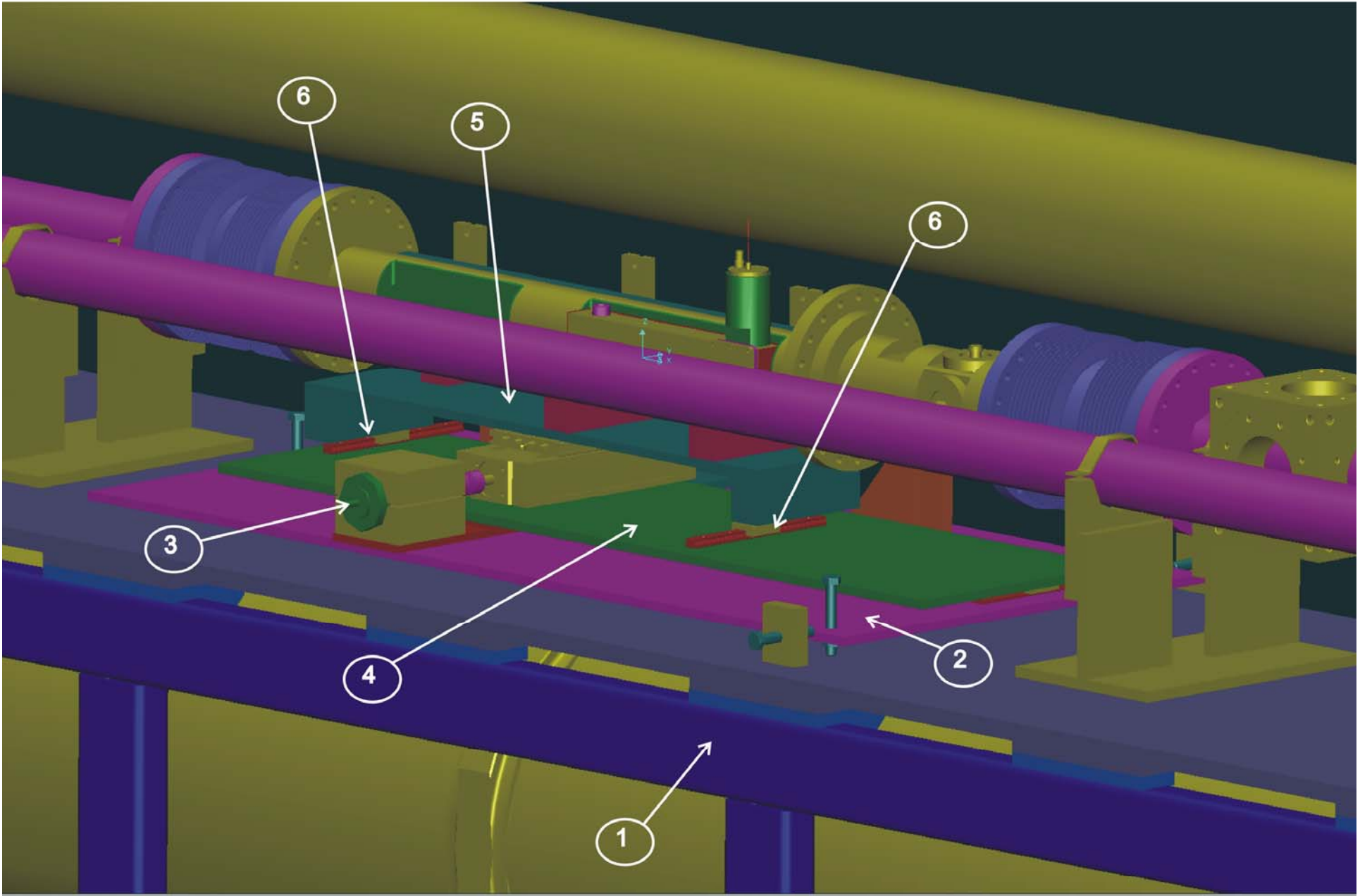
### 1. Introduction

Detection of diffractive protons at 420 m from the IP is particularly challenging since it has to be done between two LHC beam pipes, which are separated mechanically only by about 140 mm (the distance between the pipe axes is of nominal 194 mm). In addition, the nearby cryo-lines limit strongly the available free space. This is why the traditional Roman Pot technique is not possible to use, and another concept for near beam detectors, pioneered at DESY is proposed. This technique of moving Hamburg pipes was developed within the ZEUS collaboration in 1994/5 in the context of photoproduction tagging at HERA using very forward-scattered electrons [1], and was inspired by the moving pipes used in the PETRA wiggler line to allow for the beam-line aperture changes. Electron detectors (as small sandwich calorimeters) attached to the moving pipe some 44 m from the IP could approach the coasting beam and register off-beam energy electrons, which left the pipe by special exit windows. The detectors could be easily maintained and were successfully and routinely run for six years, providing data essential for several publications [2]. Each time after setting HERA in the collision mode, HERA shift crew remotely inserted detectors at the working position; about 15 mm from the coasting electron beam, using the HERA slow control system. Before installation, the Hamburg pipe system was tested by making several thousands displacement operations. Finally, no significant RF effects on the electron beam were observed due to the changed beam-pipe geometry. It should be noted that no special RF screening was used, only the inclined at  $45^\circ$  exit windows and the so-called RF fingers providing good electrical contact across the connecting bellows.









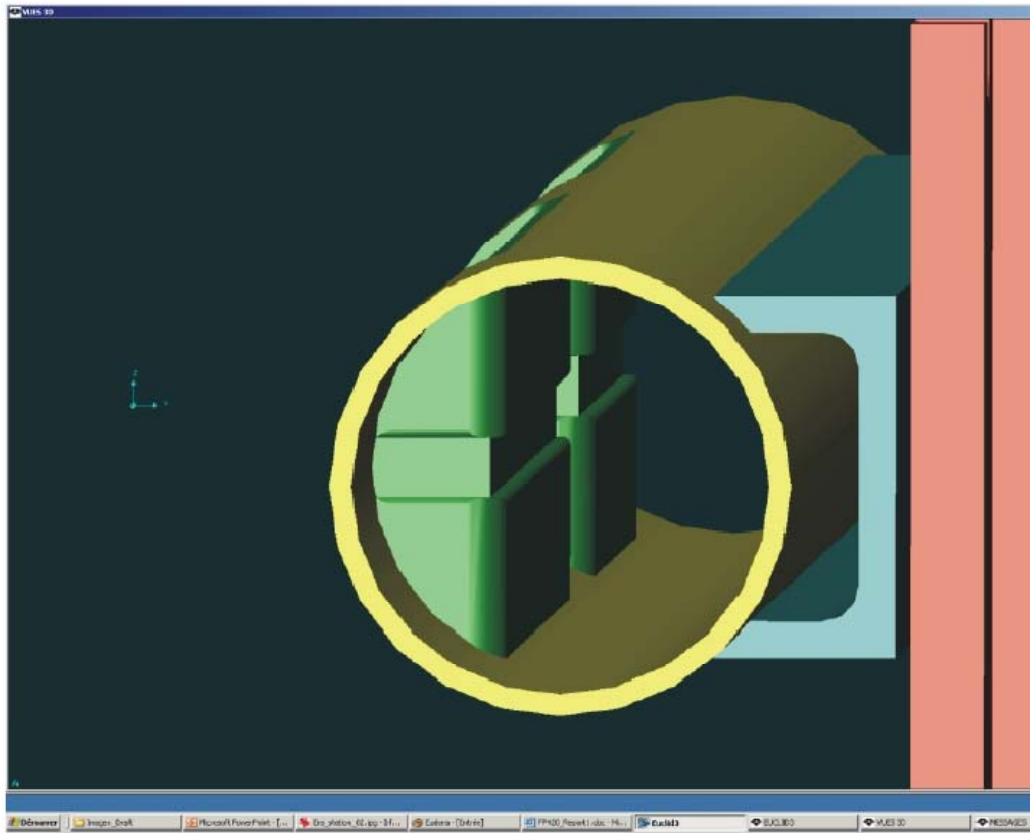
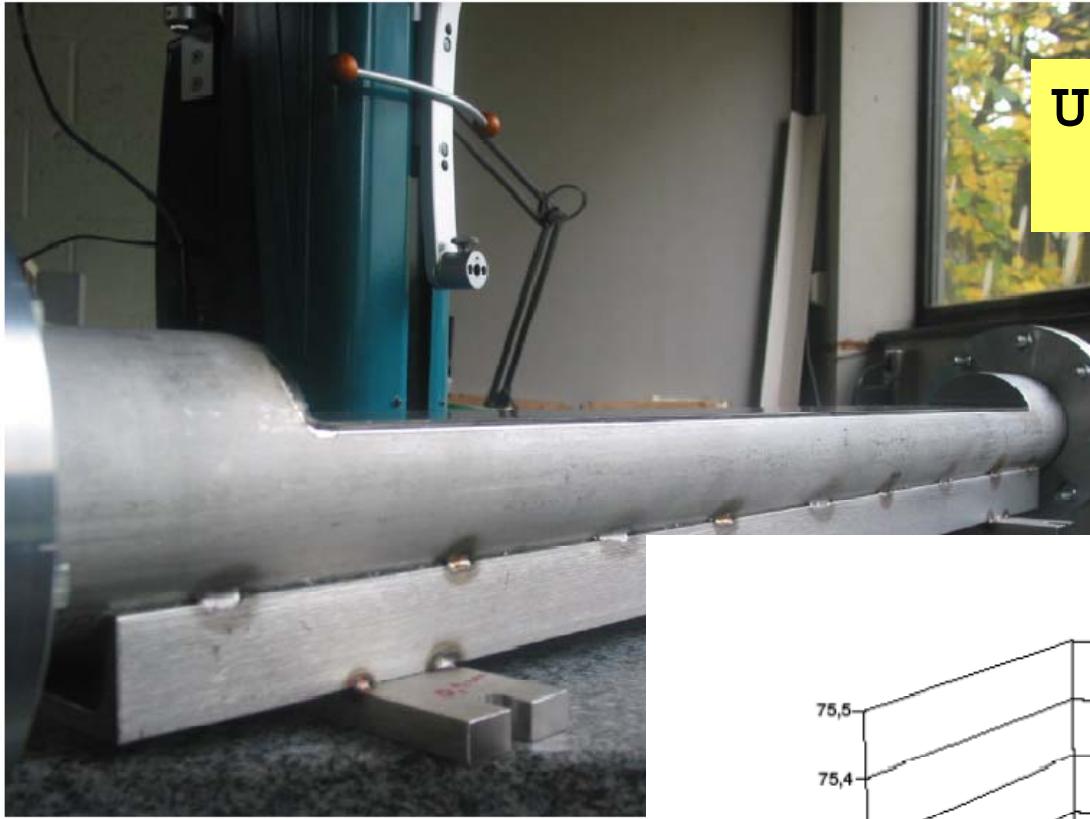


Figure 4: interior of the moving beam pipe as seen by the particles.



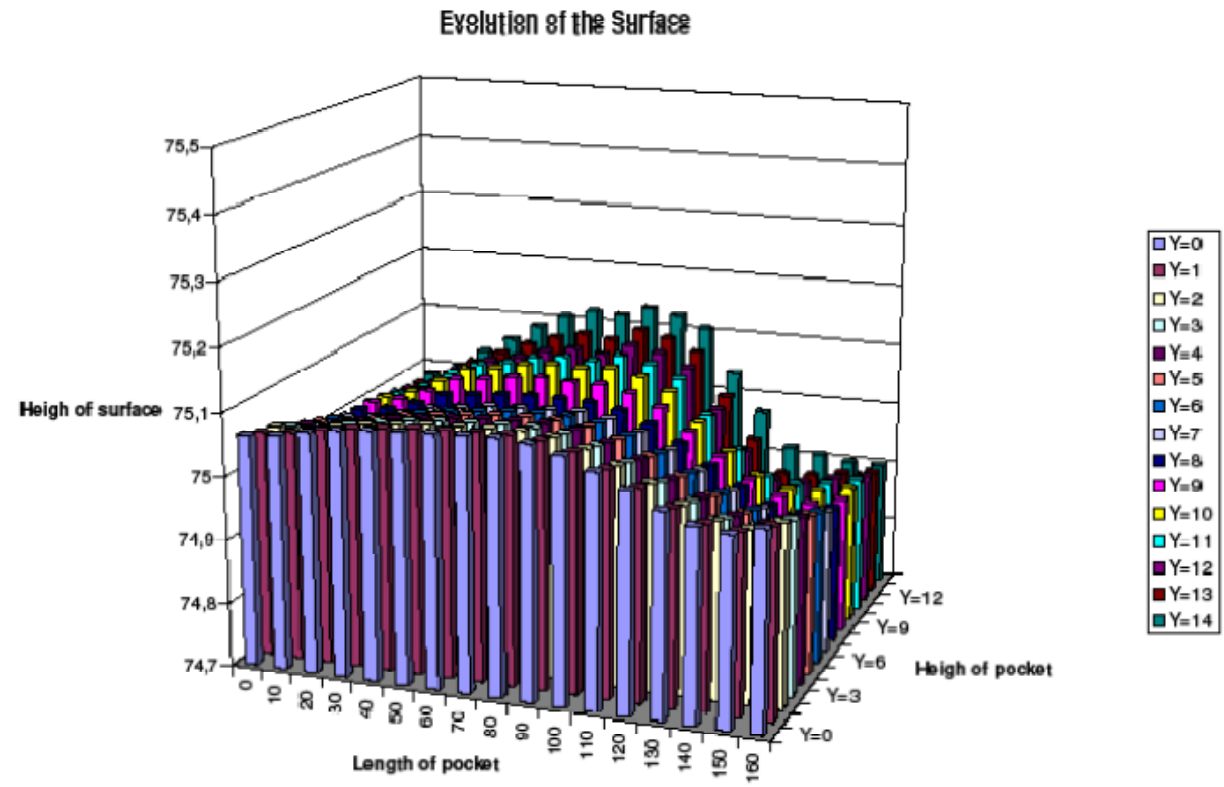
Figure 5: end view of a machined window before welding on the beam pipe slo

Figure 8: fabrication drawing for 600 mm long pocket welded in a reinforced tube.



Used for RF tests at Cockroft

Aim at  $< 50 \mu\text{m}$  pocket surface flatness





# HECTOR, a fast simulator for the transport of particles in beamlines

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## Abstract

Computing the trajectories of particles in generic beamlines is an important ingredient of experimental particle physics, in particular regarding near-beam detectors. A new tool, HECTOR, has been built for such calculations, using the transfer matrix approach and energy corrections. The limiting aperture effects are also taken into account. As an illustration, the tool was used to simulate the LHC beamlines, in particular around the high luminosity interaction points (IPs), and validated with results of the MAD-X simulator. The LHC beam profiles, trajectories and beta functions are presented. Assuming certain forward proton detector scenarios around the IP5, acceptance plots, irradiation doses and chromaticity grids are produced. Furthermore, the reconstruction of proton kinematical variables at the IP (energy and angle) is studied as well as the impact of the misalignment of beamline elements.

*PACS:* 29.27.-a ; 41.85.Ja

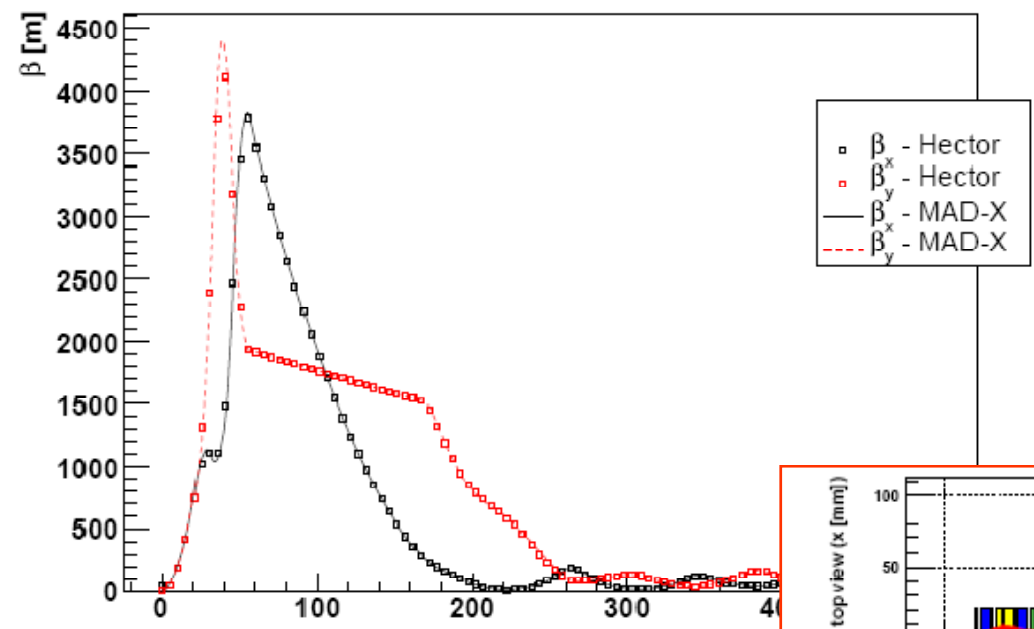
*Keywords:* Particle transport, transfer matrix, LHC beamline, forward detectors, misalignment

## 1 Introduction

In the context of the LHC beamlines and in particular of very forward, near beam detectors, we introduce a new fast simulator, called HECTOR, for the transport of single particles through generic beamlines, although primarily dedicated to the LHC. The simulator is based on a linear approach of the beamline optics, implementing transport matrices from the optical element magnetic effective length, and with correction factors on magnetic strength for particles with non nominal energy. HECTOR deals with the computation of the position and angle of beam particles, and the limiting aperture of the optical elements. It has been designed to be fast, light and user friendly and its object oriented structure, using the ROOT framework [1], helps its usage, its maintenance and the future improvements.

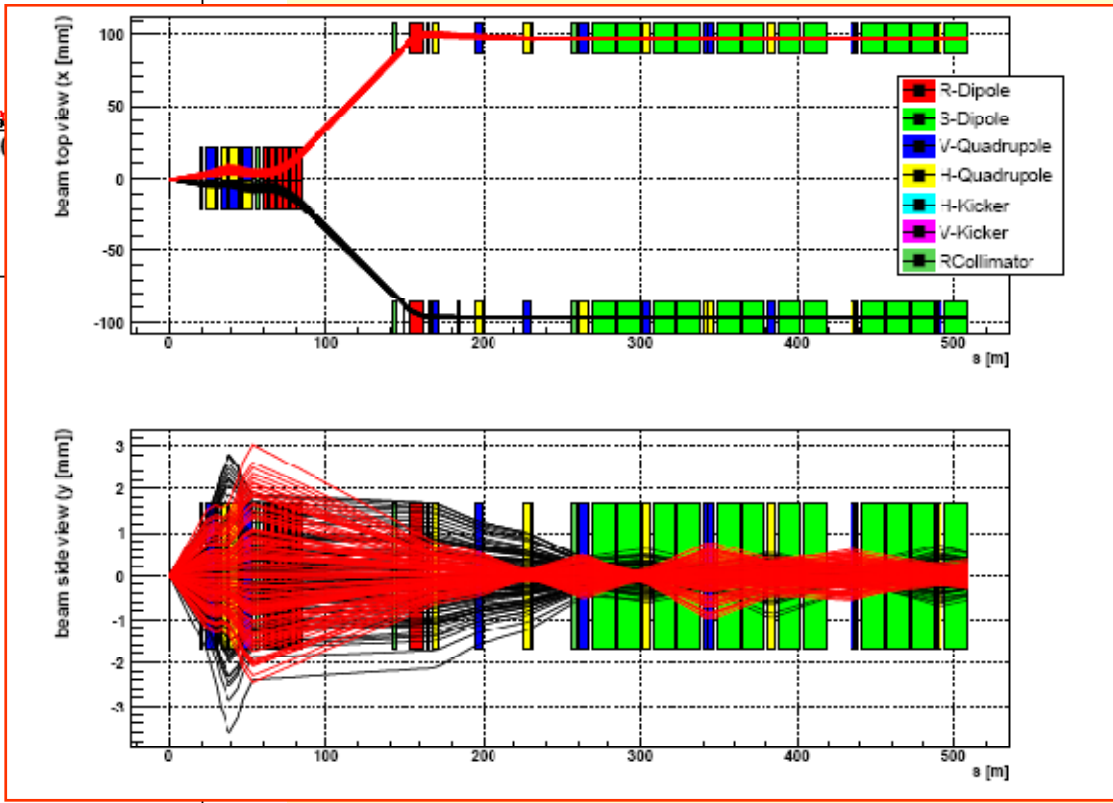
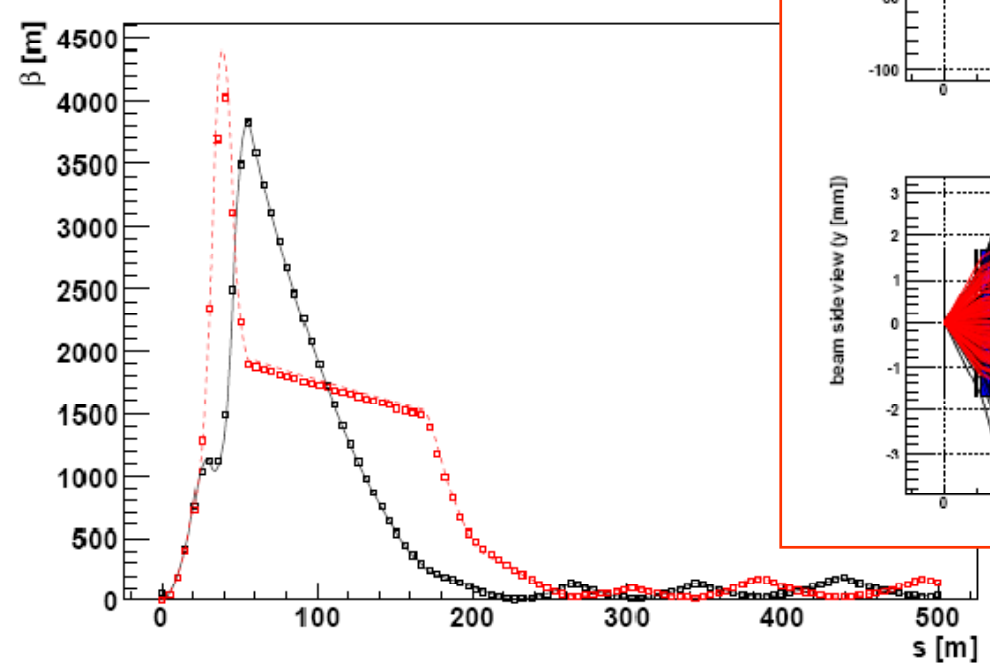
arXiv:0707.1198v1 [physics.acc-ph] 9 Jul 2007

$\beta$  functions - beam 1, forward



Hector verified with Mad-X: it is fast, versatile, precise tool for FP420/220/LHC studies

$\beta$  functions - beam 2, forward



# Acceptance driven mostly by p energy loss, and by detector/beamline geometry/aperture

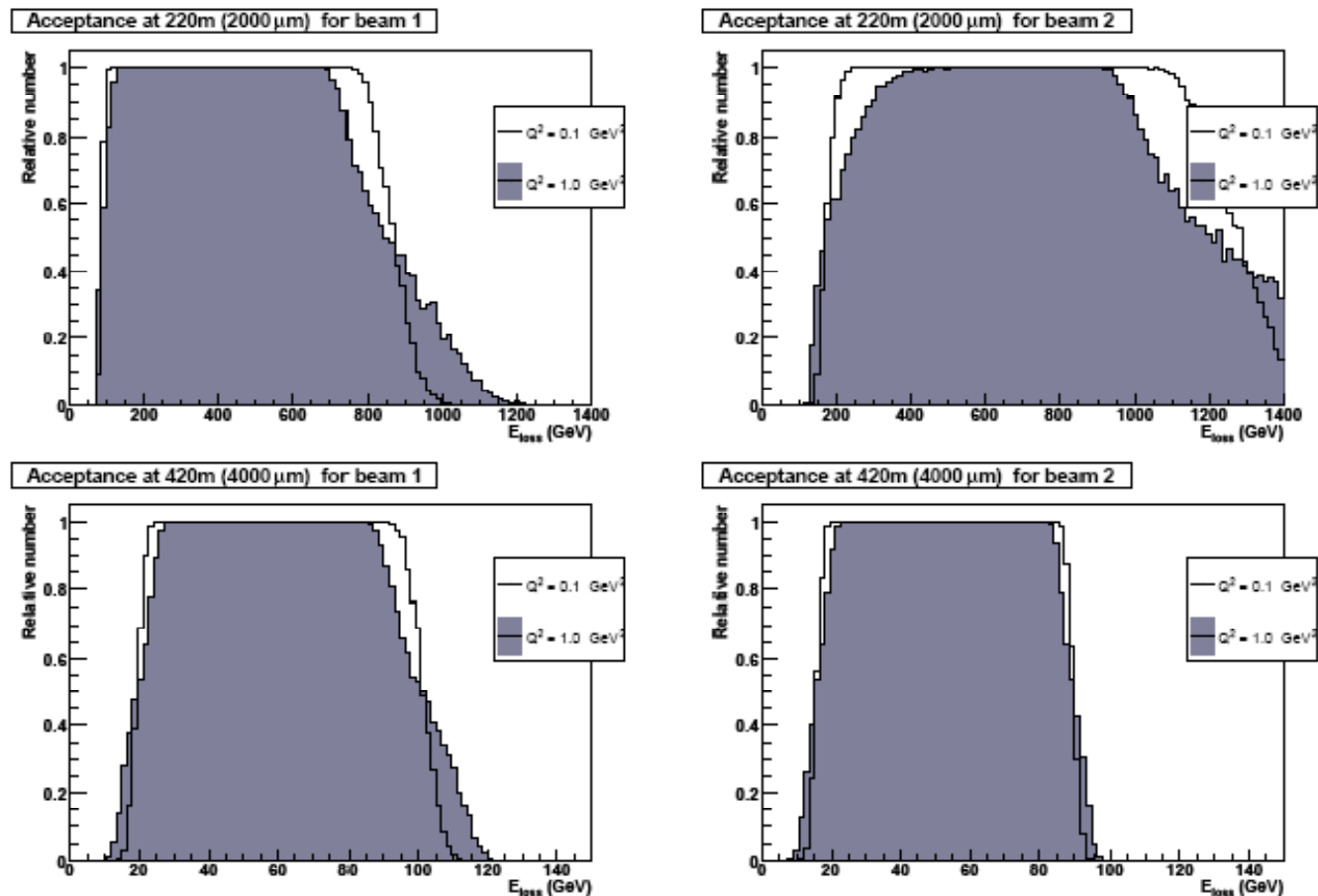
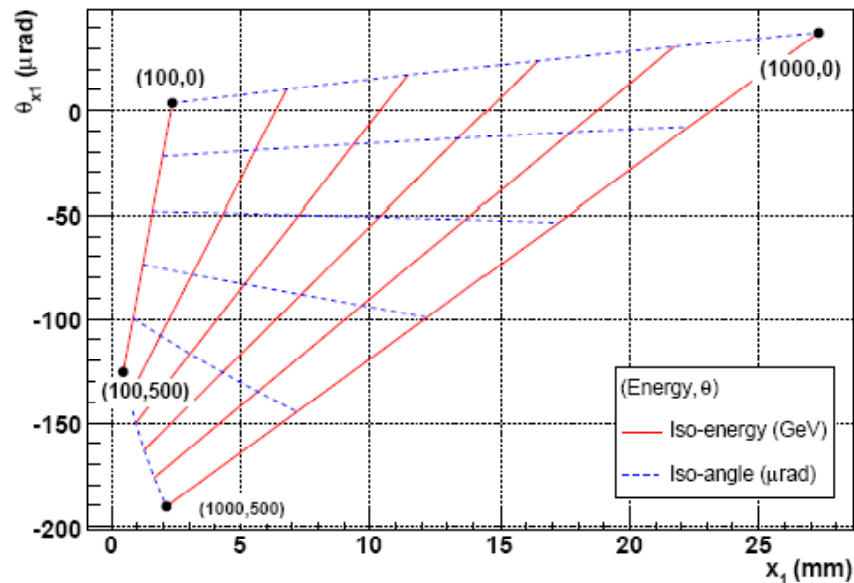
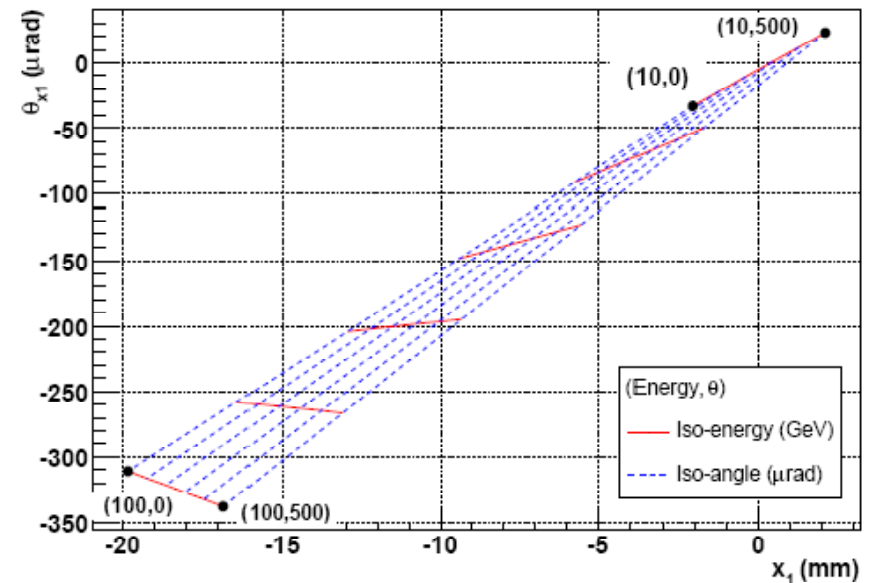


Figure 9: VFD acceptance as a function of energy loss, for two fixed virtualities (Reminder:  $Q^2 = -t$ ). See previous figure for more details.

Chromaticity grid at 220 m

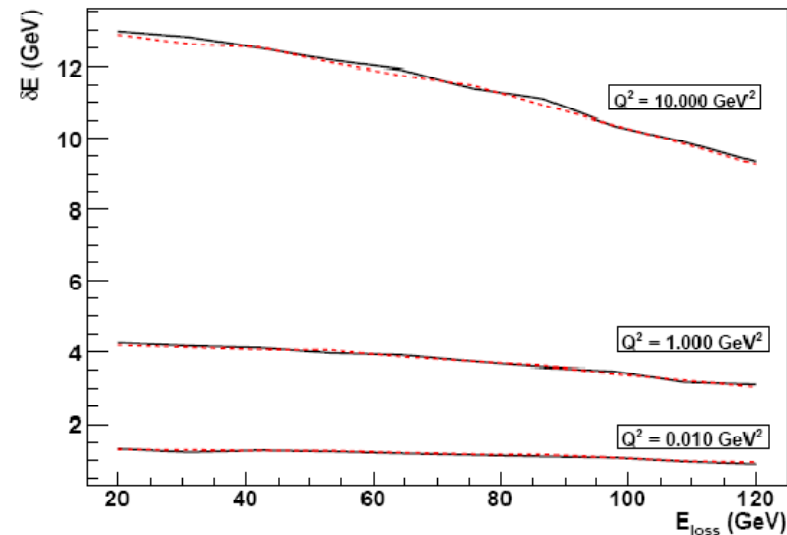


Chromaticity grid at 420 m



- Optics non-linear
- Effects of energy and angles mixed
- Transverse momentum measurement  
difficult

Energy resolution at 420 m - trivial method



Energy resolution at 420 m - trivial method

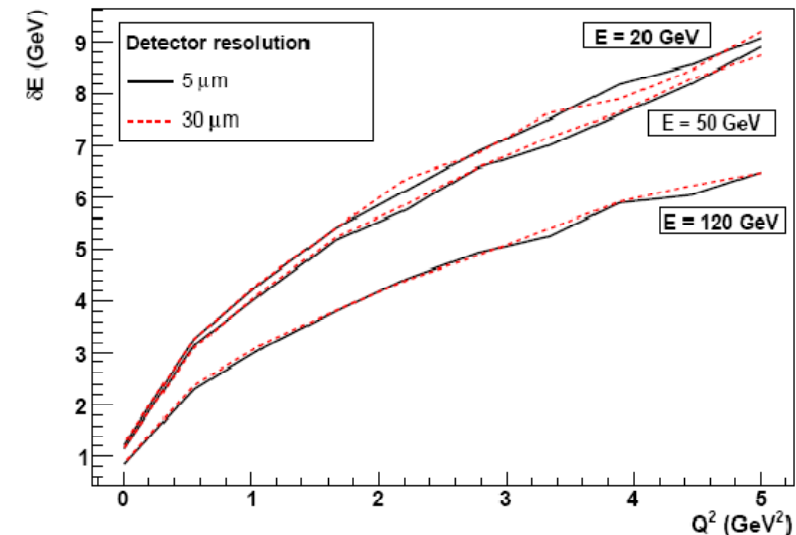


Figure 13: Energy resolution for the trivial reconstruction, for protons measured at 220 m and 420 m from the IP5. The protons have exchanged a particle taking away some energy, at a given virtuality. The resolution is shown as a function of the energy loss and the virtuality. The effects of some error on position measurement, due for instance to the spatial resolution of the detector stations (5  $\mu\text{m}$  and 30  $\mu\text{m}$ ), are taken into account. The beam energy dispersion is not included.

**However, energy measurement is good at 420 m even for poor spatial resolutions  $\sim 100 \mu\text{m}$  (for scattering angles not too large)**



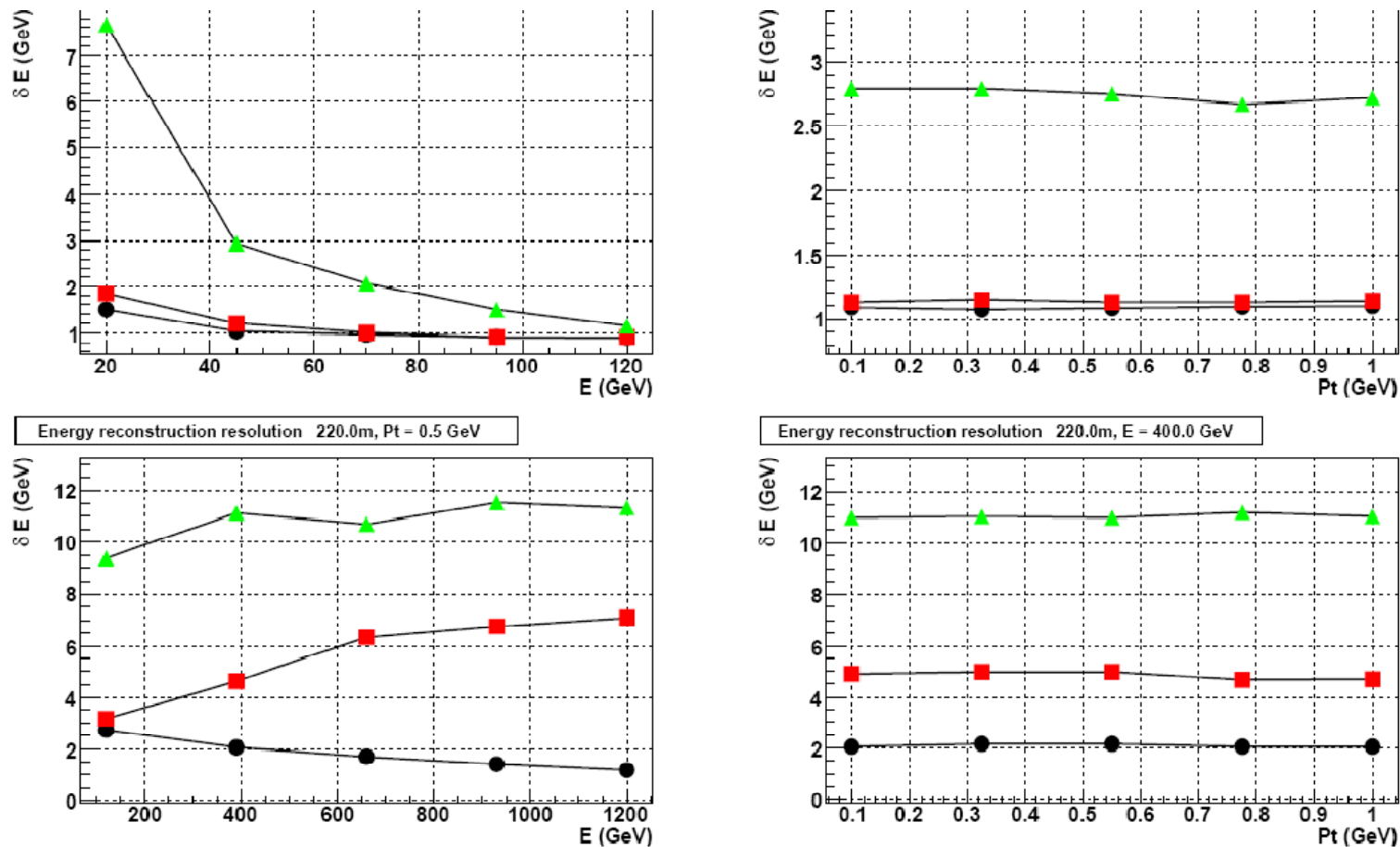


Figure 18: Resolution of the reconstruction of the particle energy loss  $E$ , as a function of the energy loss (*left*) and of the transverse momentum (*right*), for VFDs at 420 (*upper plot*) and 220 m (*lower plot*) from IP5. The advanced method, using full detector information, is used. Dots correspond to different scenarios of detector resolutions, namely perfect detectors (*circles*), 5  $\mu\text{m}$  (*squares*) and 30  $\mu\text{m}$  (*triangles*) spatial resolution.

**Ultimate full reconstruction requires extremely good resolutions and alignment (esp. at 220 m), below 10  $\mu\text{m}$**

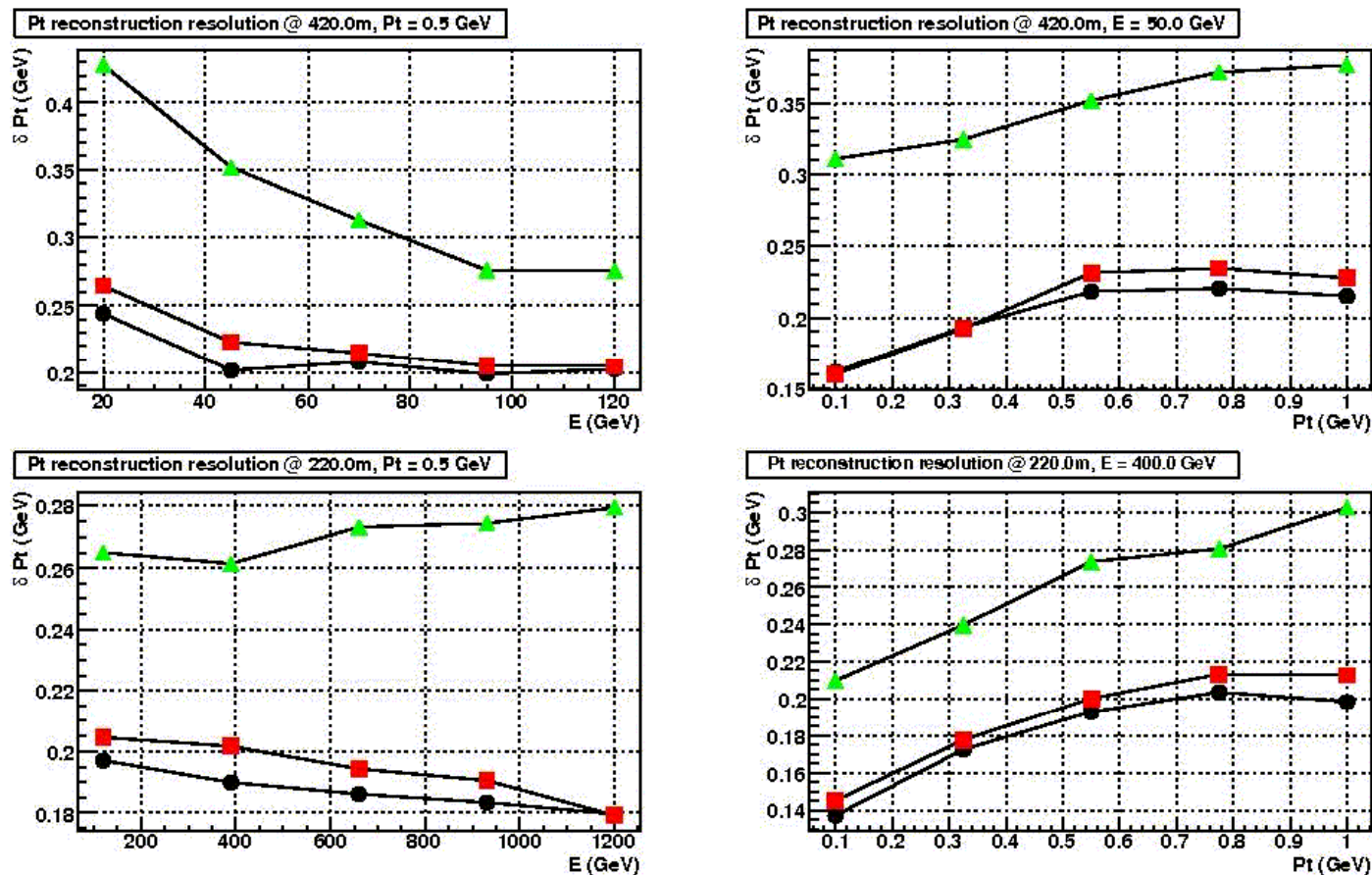


Figure 17: Resolution of the reconstruction of the particle transverse momentum  $p_T$ , as a function of the energy loss (*left*) and of the transverse momentum (*right*), for VFDs at 220 (*upper plot*) and 420 m (*lower plot*) from IP5. The advanced method, using full detector information, is used. Dots correspond to different scenarios of detectors resolutions, namely perfect detectors (*circles*), 5  $\mu\text{m}$  (*squares*) and 30  $\mu\text{m}$  (*triangles*) spatial resolution.

Ultimate full reconstruction requires extremely good resolutions and alignment, below 10  $\mu\text{m}$ ;  $p_T$  better measured at 220 m

## Misalignment impact on Higgs mass reconstruction

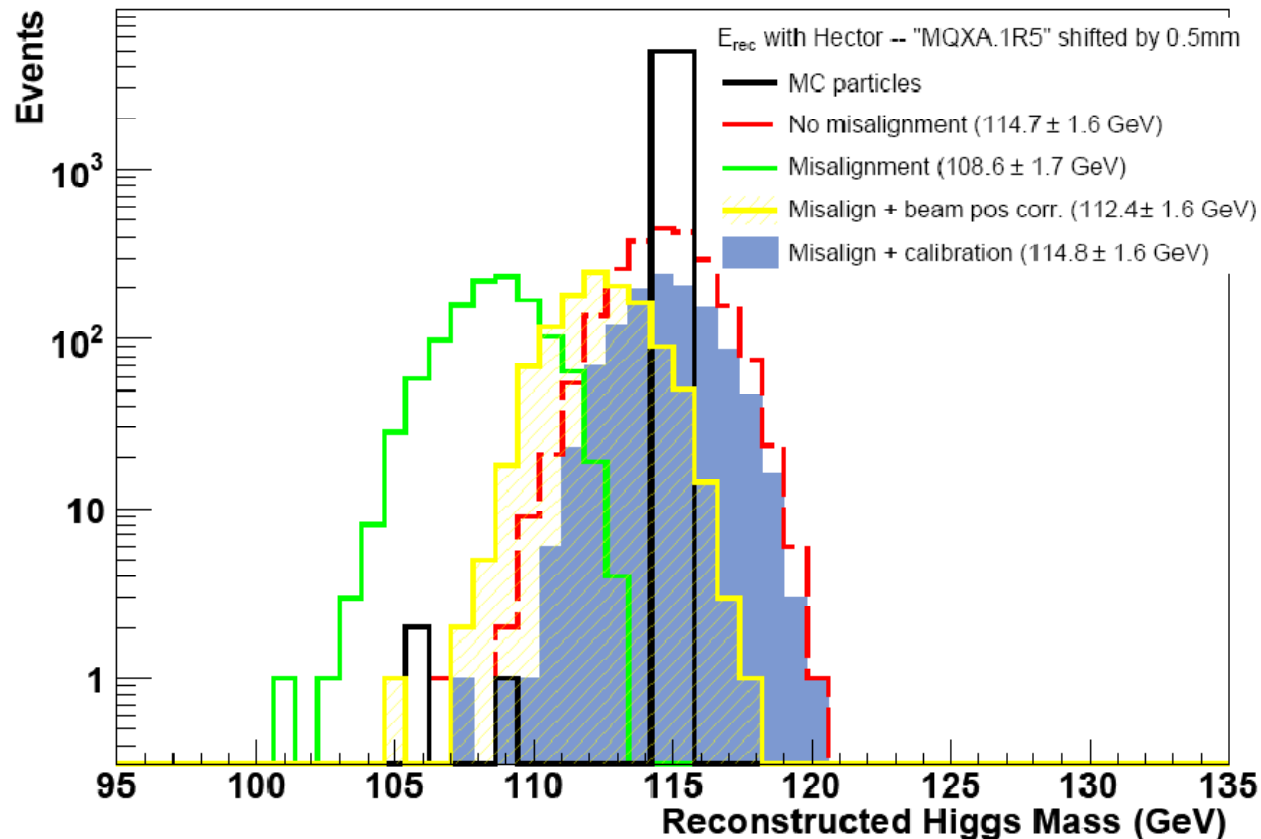


Figure 20: Illustration of the effects in the energy reconstruction due to the misalignment of LHC quadrupoles. The graphs show the reconstructed Higgs boson mass in the two-photon exclusive production, using energy of two forward scattered protons. In the upper plot, a quadrupole (MQM9R5,  $s = 347$  m) close to the detector has been shifted by  $100 \mu\text{m}$ . Misaligning an optical element (MQXA1R5,  $s = 29$  m) close to the IP leads to a loss of acceptance (lower plot). The reconstructed values including the correction due to the dimuon calibration is also plotted. In brackets, the average reconstructed mass and its resolution are given, without including the beam energy dispersion.

# Standard Candle: Exclusive di-muons

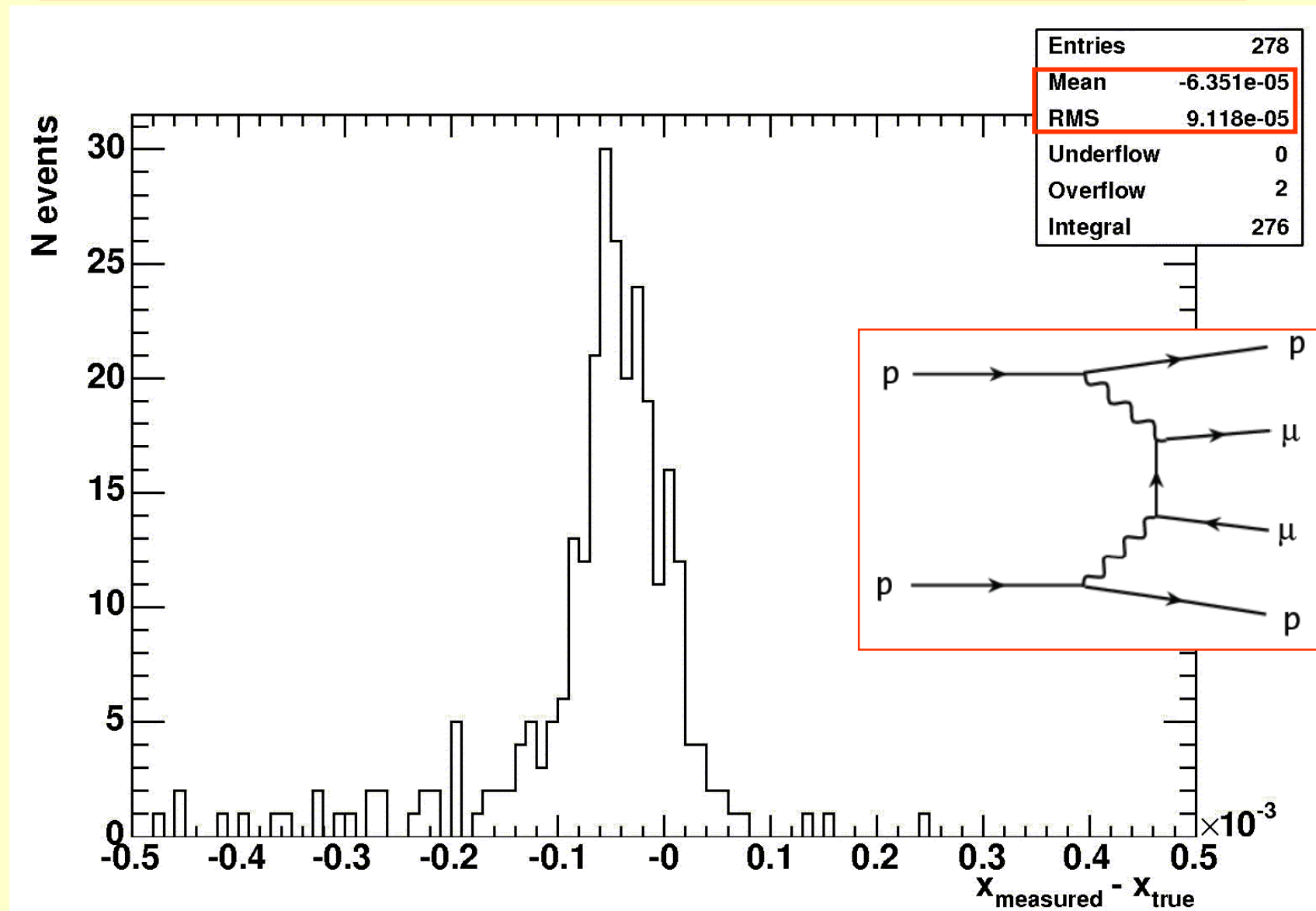
Y. Liu (UCLouvain)

## Reconstruction & trigger efficiencies:

- Sample of 5000  $\mu\mu$  pairs at  $p_T > 3$  GeV generated with LPAIR ( $\sigma \approx 50$  pb)
- Reconstructed with ORCA - efficiency for muons reaches plateau of 90% at  $p_T = 7$  GeV (at  $p_T = 5$  GeV it is 60%)
- Cross-section for reconstructed pairs is 6 pb, L1 efficiency is high but high  $p_T$  cuts at HLT results in 40% global trigger efficiency

**Note: Roman Pots are NOT used for trigger!**

Resolution of the proton energy loss for the reconstructed dimuon pairs:





## Hits in VFD at 420m ( $L=20 \text{ fb}^{-1}$ )

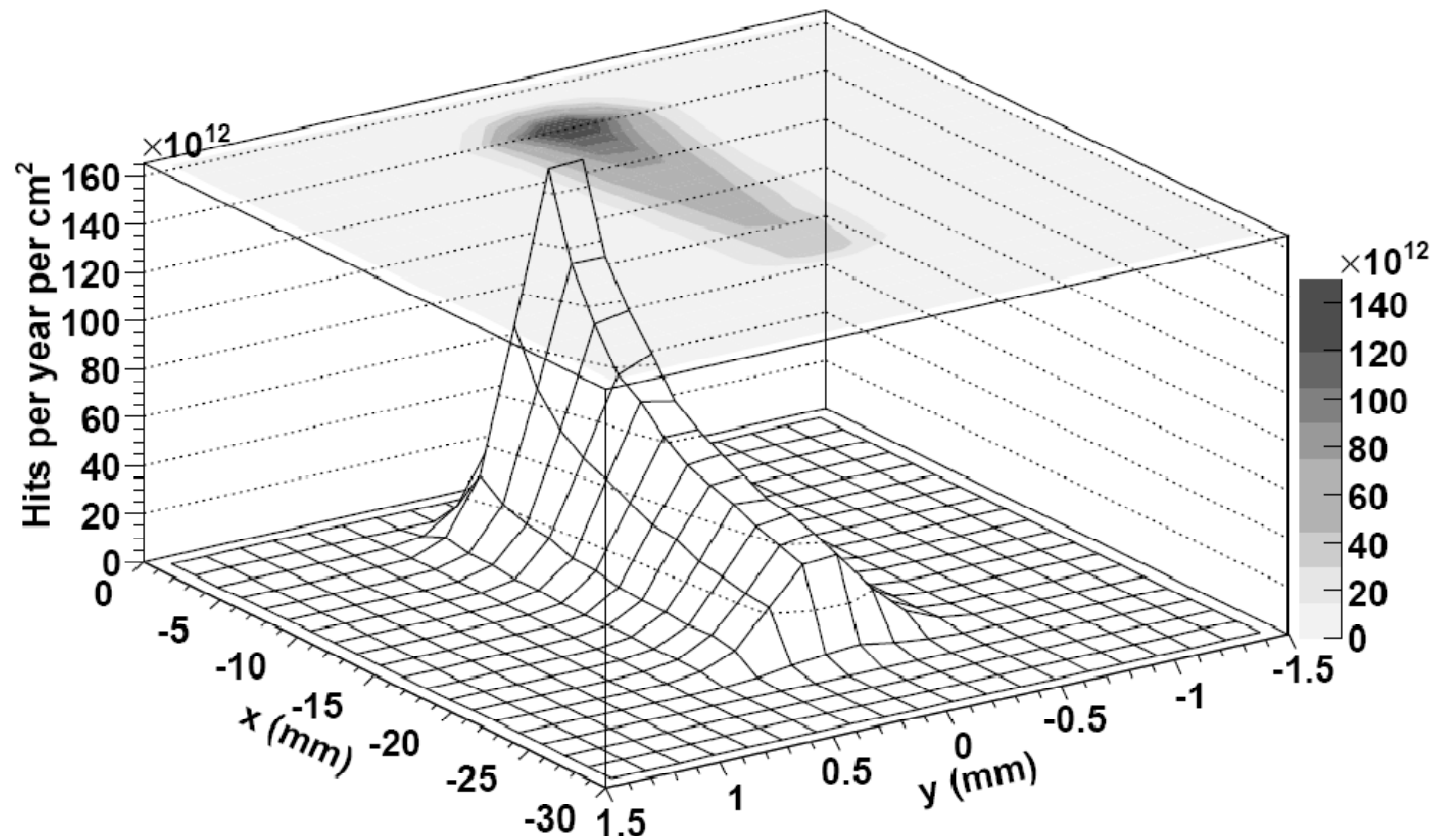


Figure 12: Irradiation levels due to  $pp \rightarrow pX$  processes, of VFDs located at  $s = 220 \text{ m}$  (*above*) and  $s = 420 \text{ m}$  (*below*) from the IP5. The horizontal position of the detector edge is respectively  $x = 2 \text{ mm}$  and  $x = 4 \text{ mm}$  from the nominal position of the center of the corresponding beam. The fluence is given per year and square centimeter. The integrated luminosity of  $L = 20 \text{ fb}^{-1}$  was assumed. For this analysis, HECTOR has propagated protons with 4-momentum generated by PYTHIA 6.2.10 (process number 93). The irradiation levels can locally exceed  $1 \times 10^{14}$  protons per square centimeter.

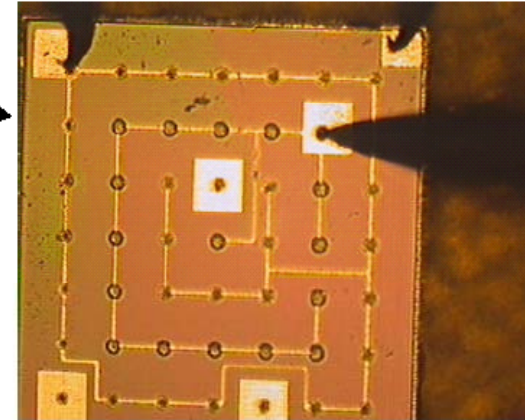
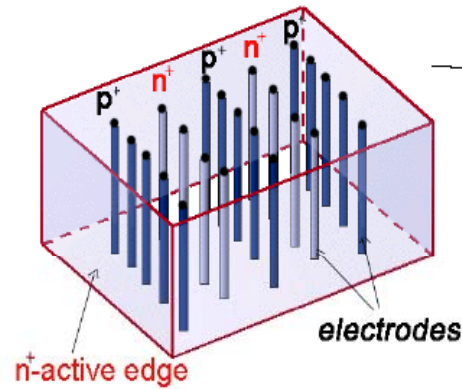
## Conclusions from Hector:

- Radiation hard detectors are needed
- Spatial resolution is essential for  $p_T$  reconstruction
- Di-muon essential for understanding energy scales, acceptances and luminosity of FP420 events
- Absolute energy resolution worse but  $p_T$  resolution better at 220m
- Calibration of detectors at 220m more tricky: need to study electron pairs with forward CALs, and bremsstrahlung in ZDCs

# 3D DETECTORS AND ACTIVE EDGES

Brunel, Hawaii, Stanford

## 3D Si Tracking + Cerenkov Fast Timing

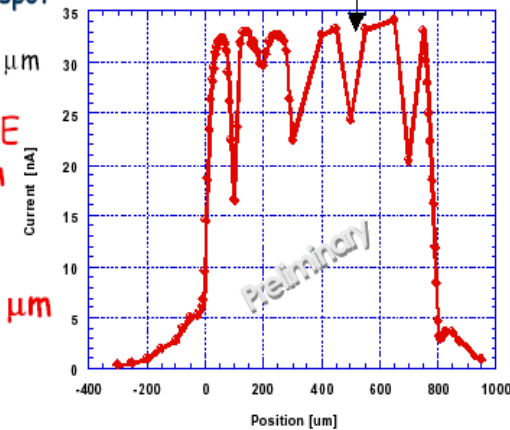


❖ EDGE SENSITIVITY	< 10 $\mu\text{m}$
❖ COLLECTION PATHS	~50 $\mu\text{m}$
❖ SPATIAL RESOLUTION	10-15 $\mu\text{m}$
❖ DEPLETION VOLTAGES	< 10 V
❖ DEPLETION VOLTAGES at $10^{15}\text{n/cm}^2$	~105 V
❖ SPEED AT RT	3.5 ns
❖ AREA COVERAGE	3X3 $\text{cm}^2$
❖ SIGNAL AMPLITUDE before Irradiation	24 000 e
❖ SIGNAL AMPLITUDE at $10^{15}\text{n/cm}^2$	15 000 e

- ❖ 15  $\mu\text{m}$  InfraRed beam spot
- ❖ FWHM = 772  $\mu\text{m}$
- ❖ Edge Al strip width = 16  $\mu\text{m}$

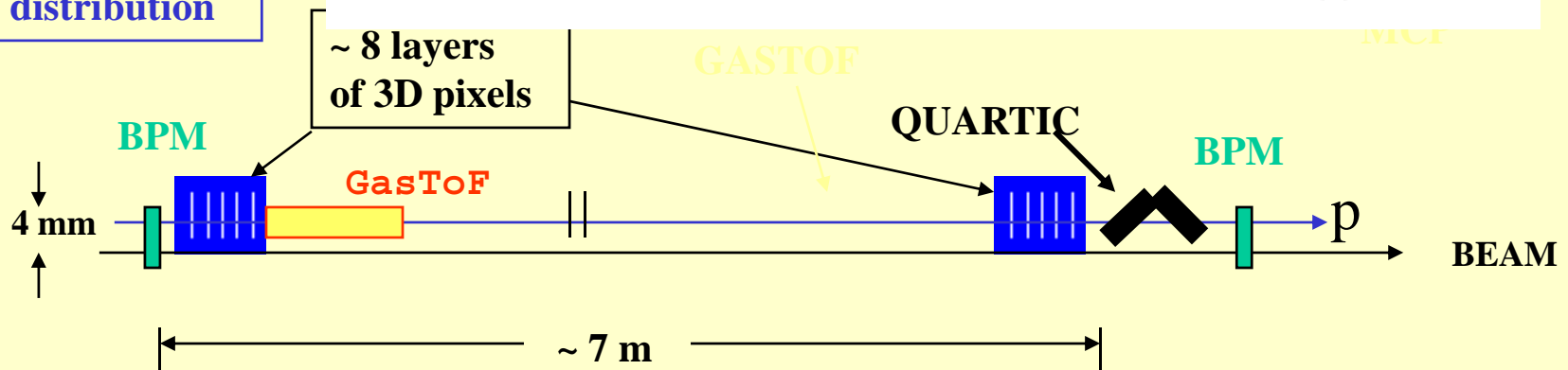
INSENSITIVE EDGE (INCLUDING 16  $\mu\text{m}$  Al STRIP):

$$(813 - 772) / 2 = 21 \mu\text{m}$$



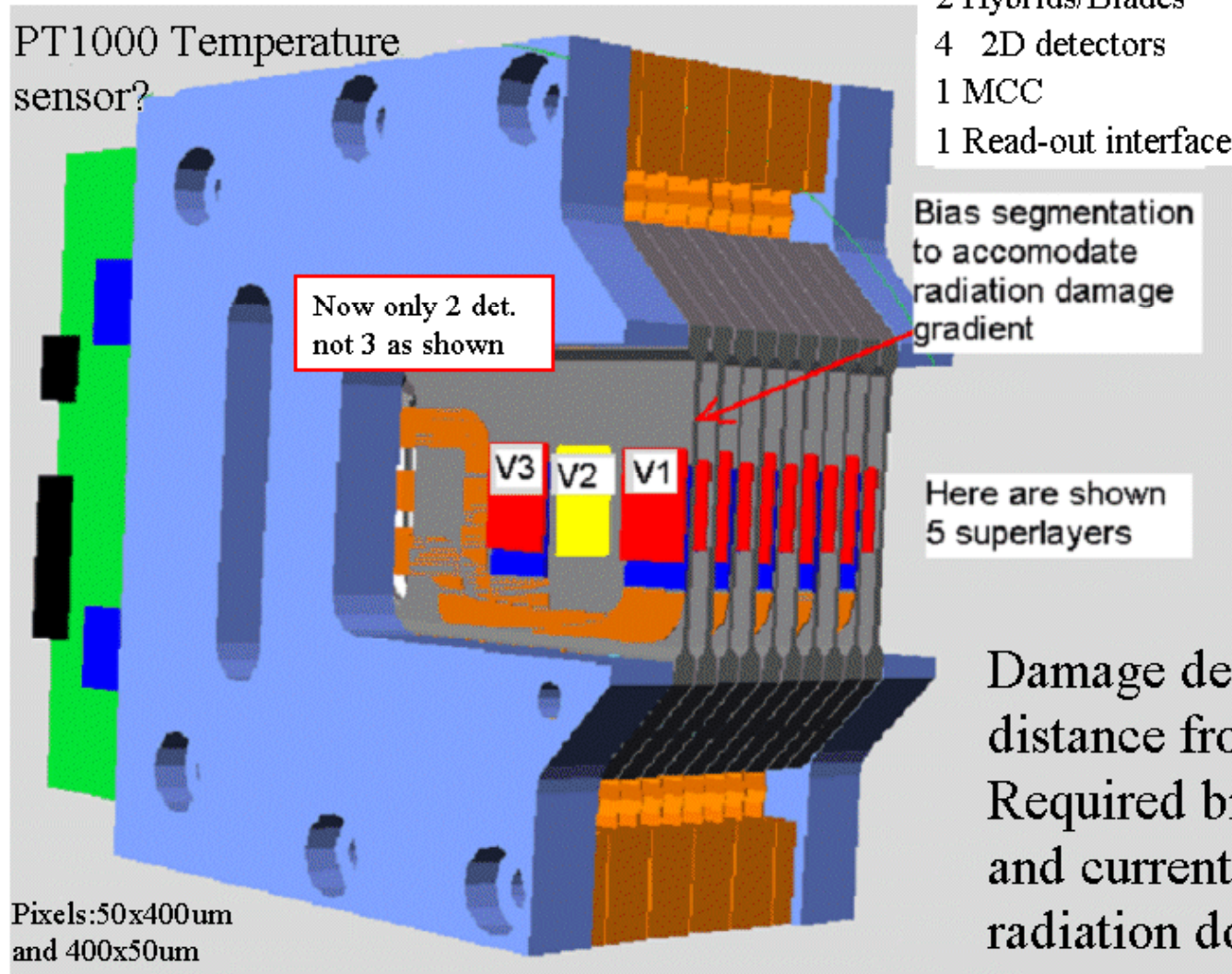
6mm(y) x 24mm (x) covers distribution

❖ CERN Courier, Vol 43, Number 1, Jan 2003



# HV-LV supply segmentation

PT1000 Temperature sensor?



Pixels: 50x400um and 400x50um

Drawing: From Ray Thompson

MCC: Module Controller Chip

# Why need s(uperfast)TOF?

Z-by-timing is crucial for running at high LHC luminosity, by suppressing accidental backgrounds:

If  $\delta t = 10$  ps can be achieved for a single ToF, then z-vertex resolution is 2 mm (from time difference for two arms) to be compared with >50 mm RMS of IR!

Note: Background suppression power is ~ inversely proportional to resolution of ToF!



## gastof: Basic idea

Consider gas Cerenkov as complementary solution to Quartic:

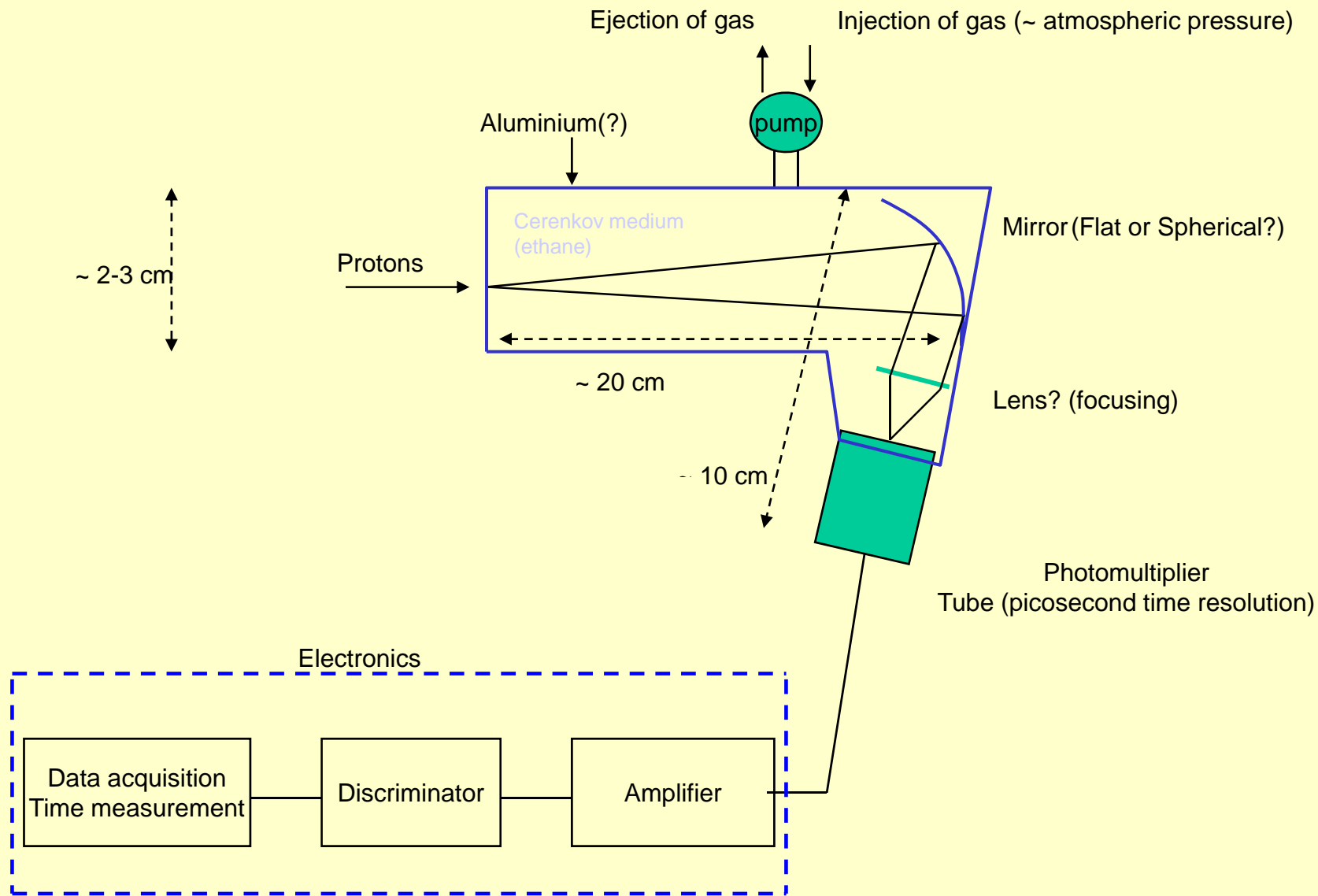
- Very simple and robust design
- Very thin and light detector - can be used before the tracking part
- (Very) radiation hard
- High energy threshold

Basic formula:  $N_{pe} \approx 100 \sin^2\theta_c L[\text{cm}]$

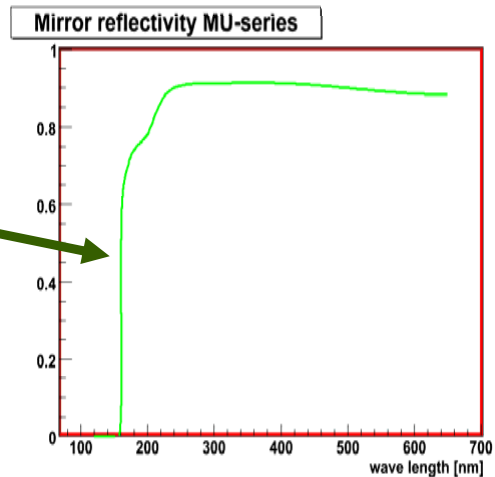
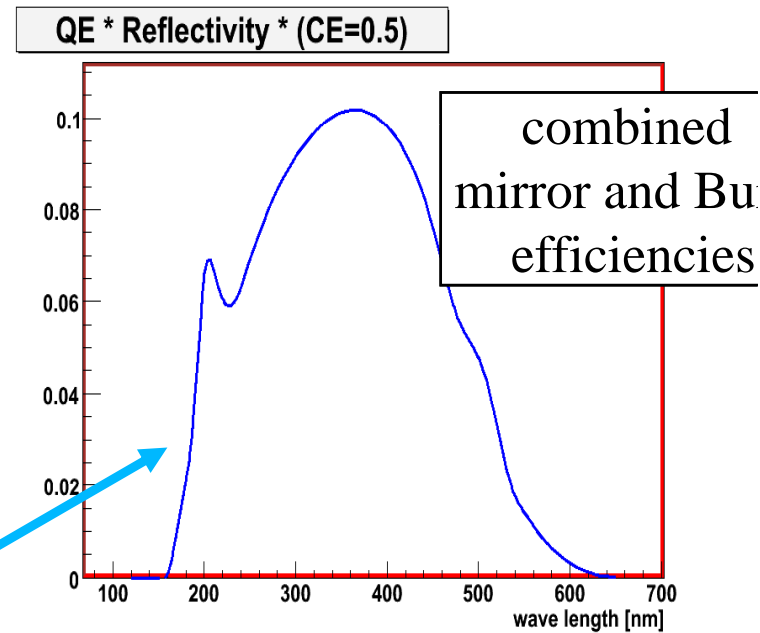
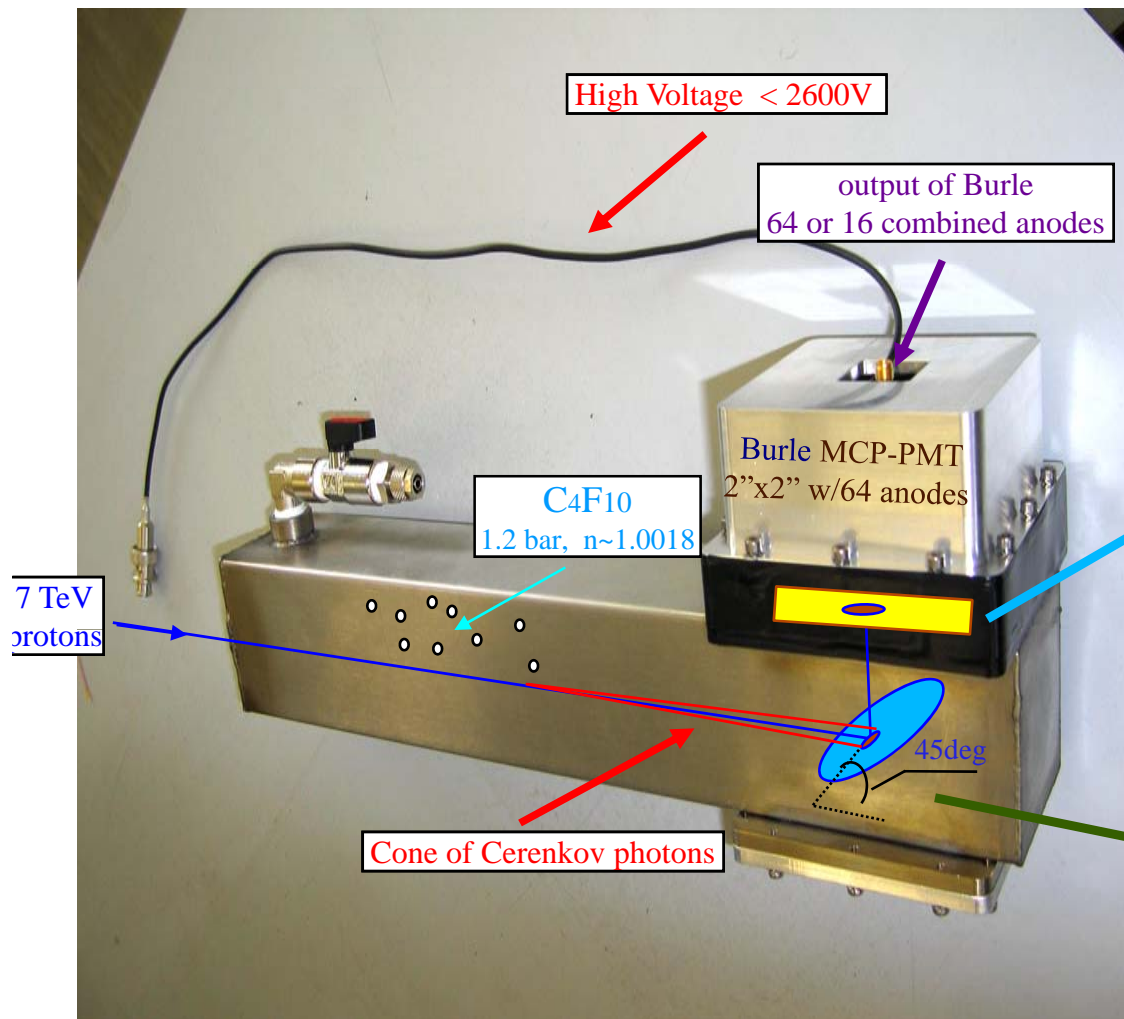
To estimate position sensitivity estimate average light spot radius  $\langle r \rangle$ , at radiator exit:

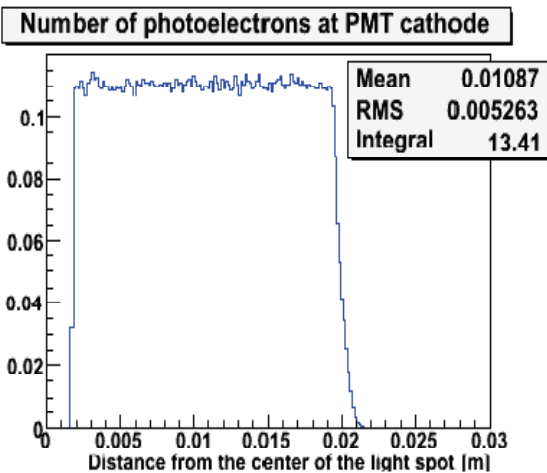
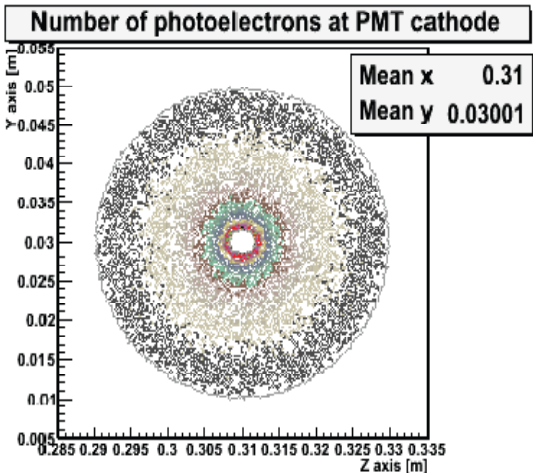
$$\langle r \rangle \approx 0.5 L \tan\theta_c \approx \sin\theta_c L/2$$

$$N_{pe} \approx 200 \langle r \rangle[\text{cm}] \sin\theta_c$$

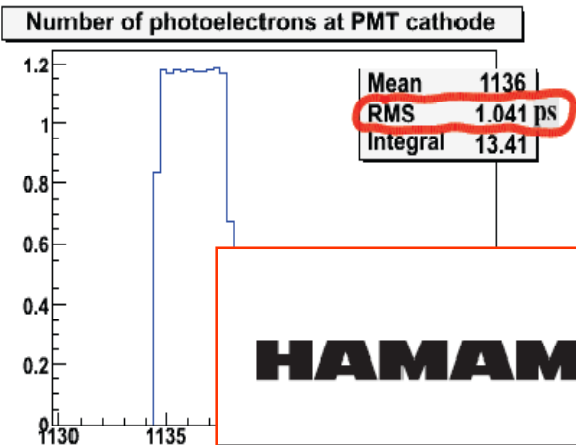
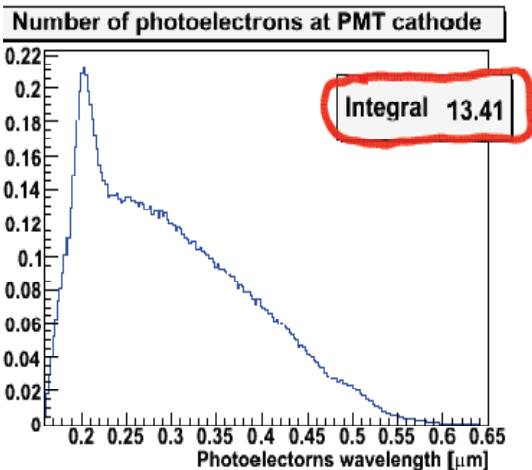


# gastof prototyping





- simulation input :
- ◆ length 31cm
  - ◆ C4F10 at 1.2 bar
  - ◆ protons at 6.9 TeV



- results :
- ◆ Arrival Time RMS ~ 1ps ! (without PMT)
  - ◆ Spot Size < 2 cm
  - ◆ Avg. number of photoelectrons = 13.4

**HAMAMATSU**

**MICROCHANNEL PLATE  
PHOTOMULTIPLIER TUB  
(MCP-PMTs)  
R3809U-50 SERIES**

Compact MCP-PMT Series Featuring  
Variety of Spectral Response with Fast Time Response

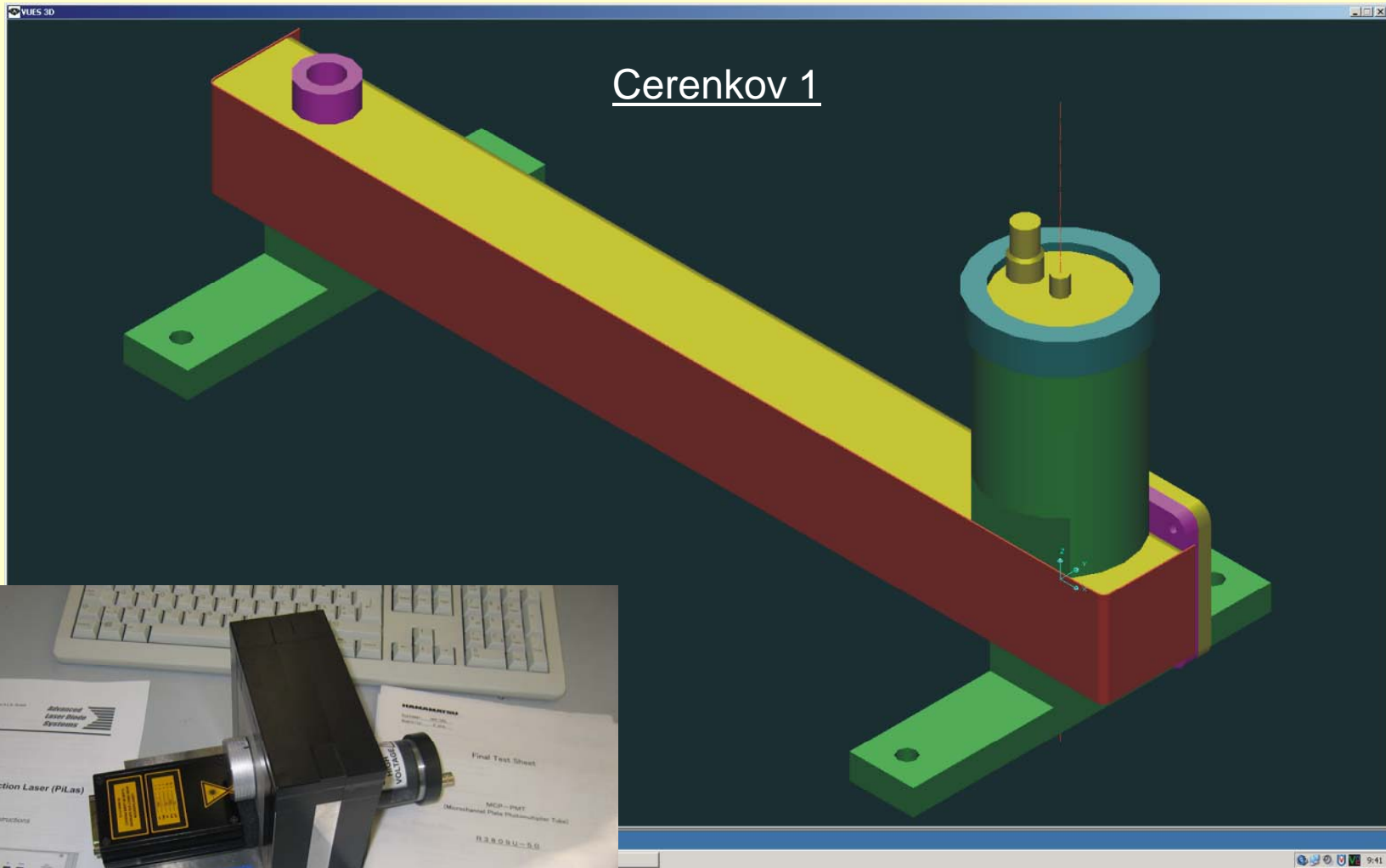
**FEATURES**

- High Speed  
Rise Time: 150ps  
T.T.S. (Transit Time Spread)<sup>1)</sup>: ≤ 25ps(FWHM)
- Low Noise
- Compact Profile  
Useful Photocathode: 11mm diameter  
(Overall length: 70.2mm Outer diameter: 45.0mm)



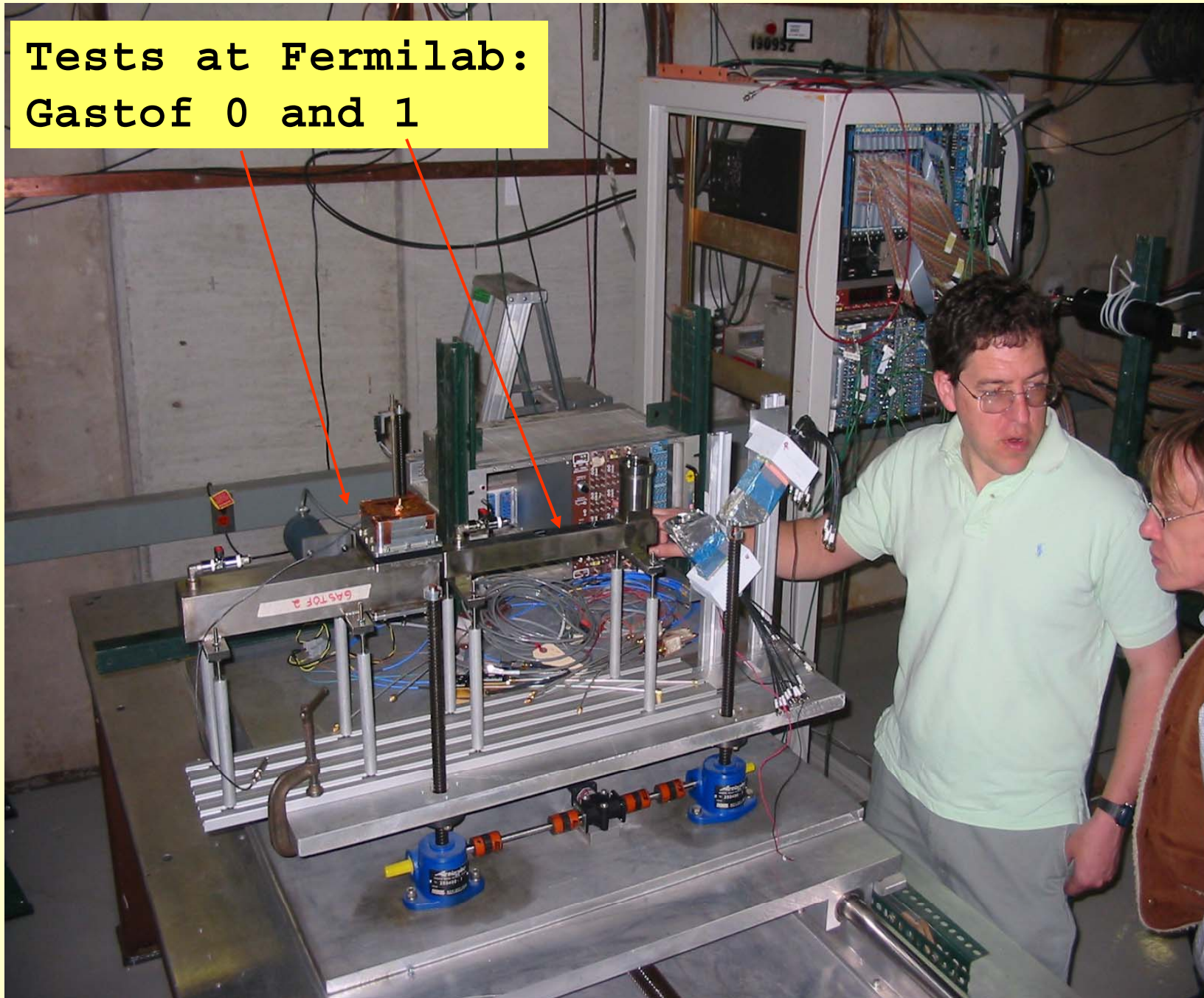
Prototype 1 with  
Hamamatsu MCP-PMT:

# GasToF

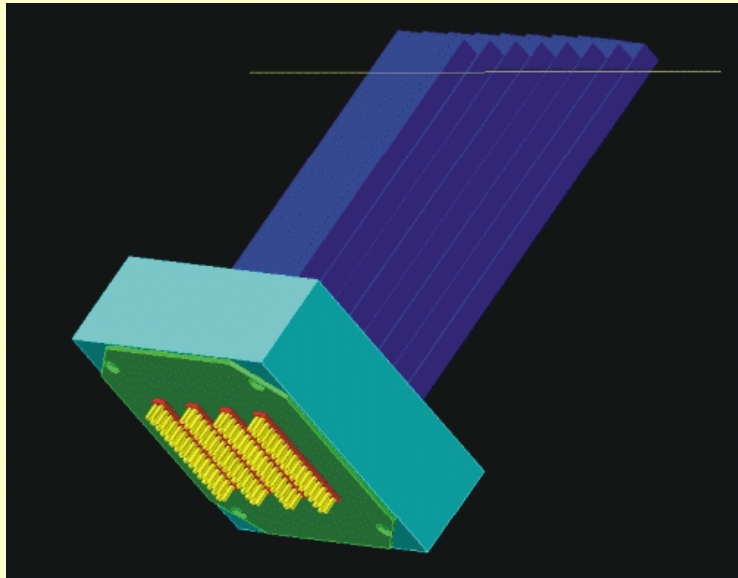




Tests at Fermilab:  
Gastof 0 and 1



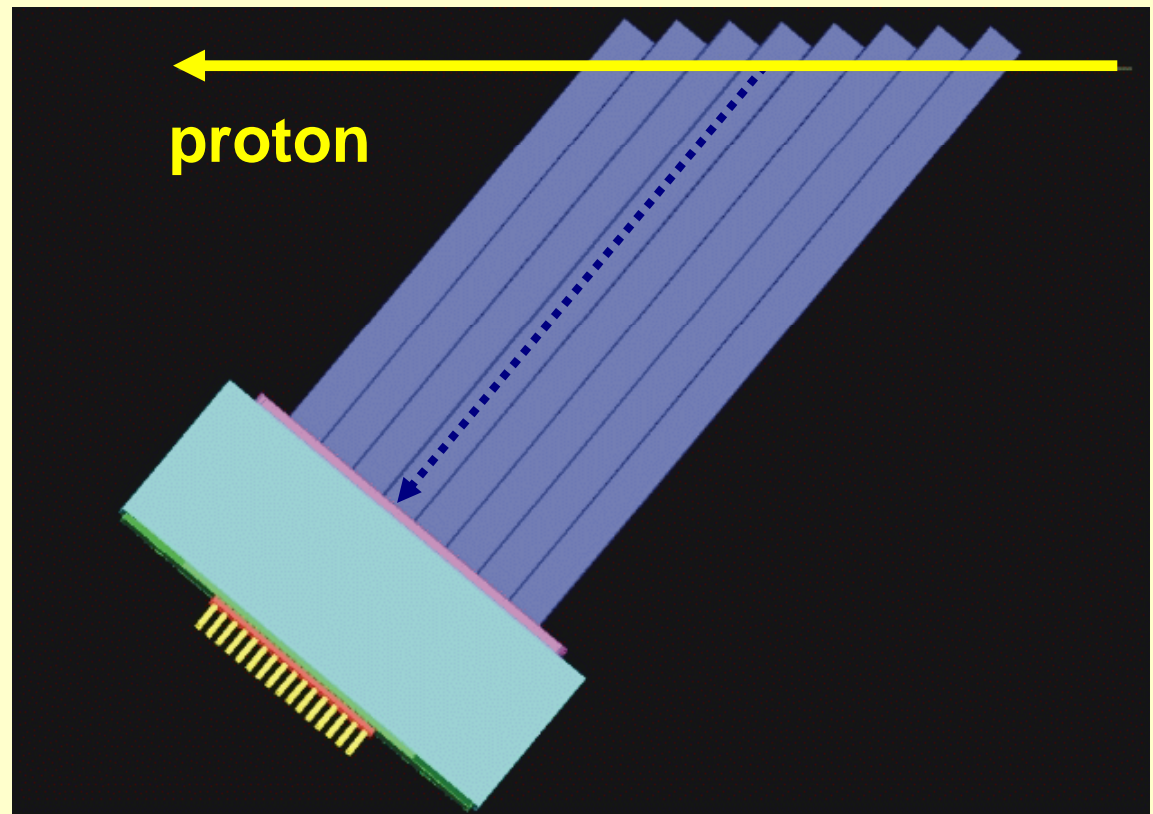
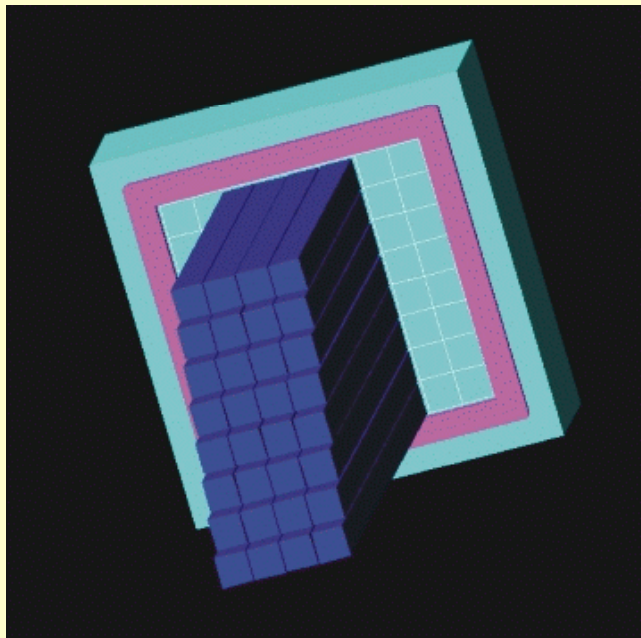
# The Detectors : 2) QUARTIC



Mike's idea

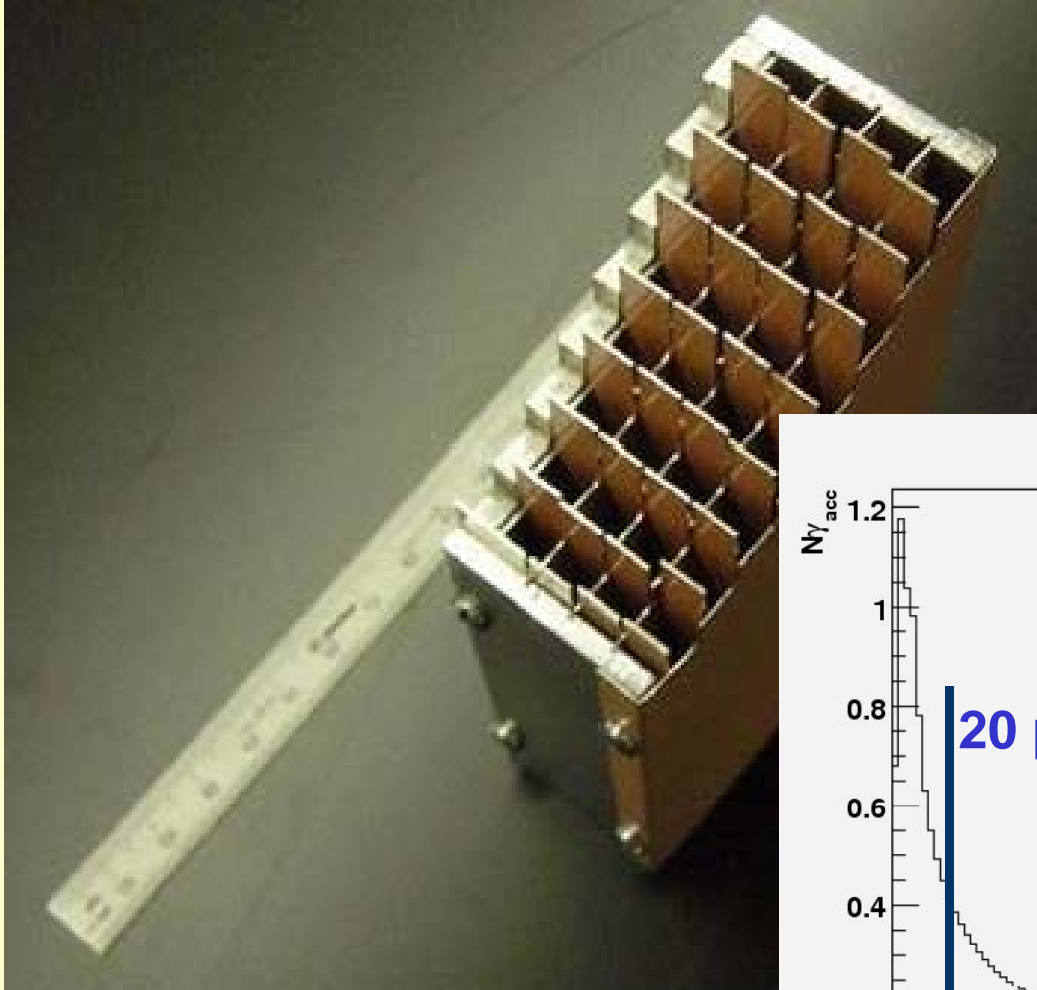
Jaak's drawing

8 fused silica rods  
in z (my fused  
silica!)





# QUARTIC (V2)

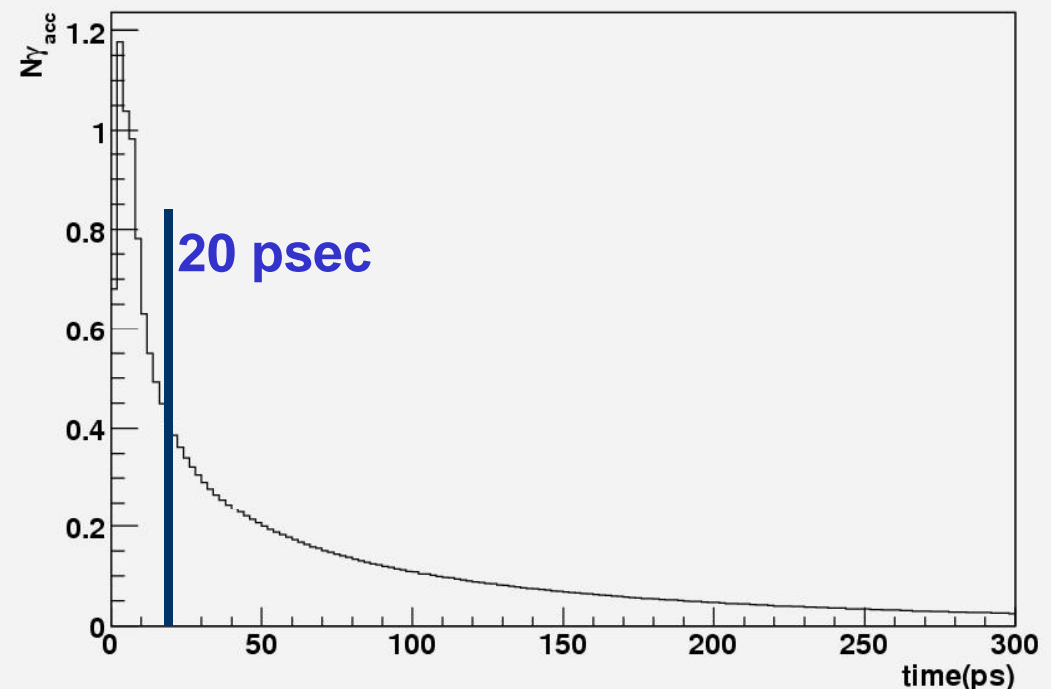


## Advantages:

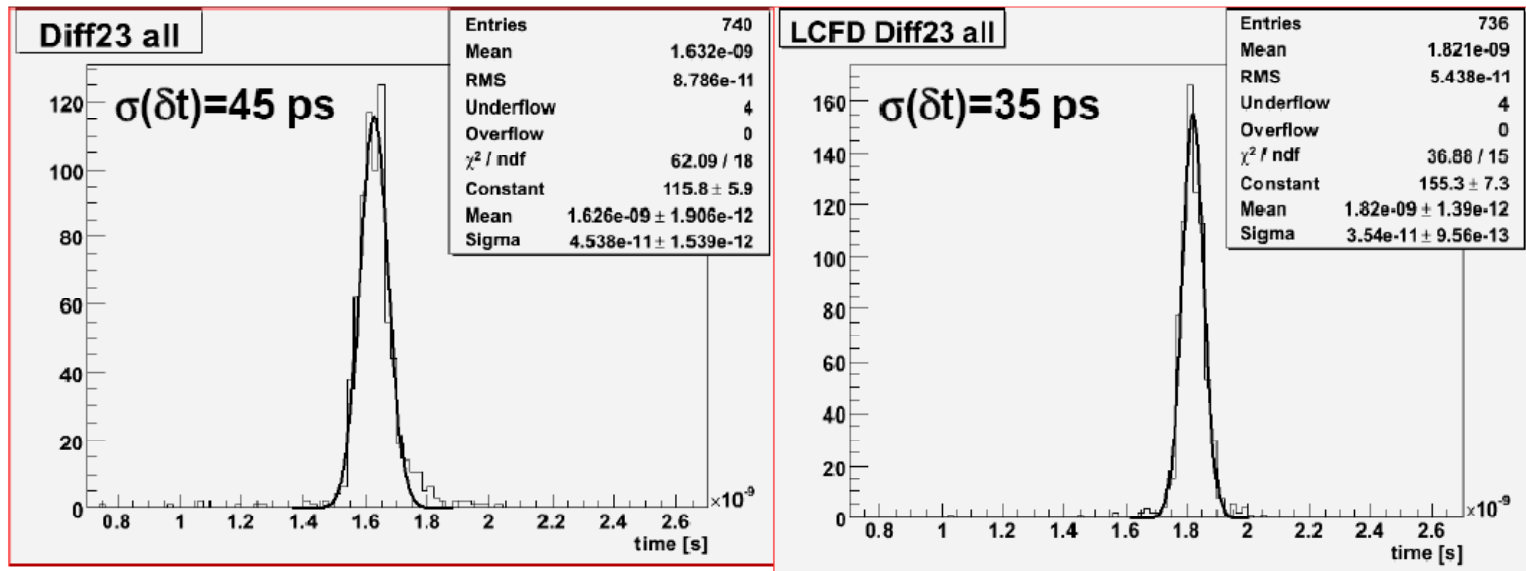
- Segmentation (8 x 4)
- Compact

## Disadvantages:

- More material
- Not as fast



# Scope Analysis (G1-G2)



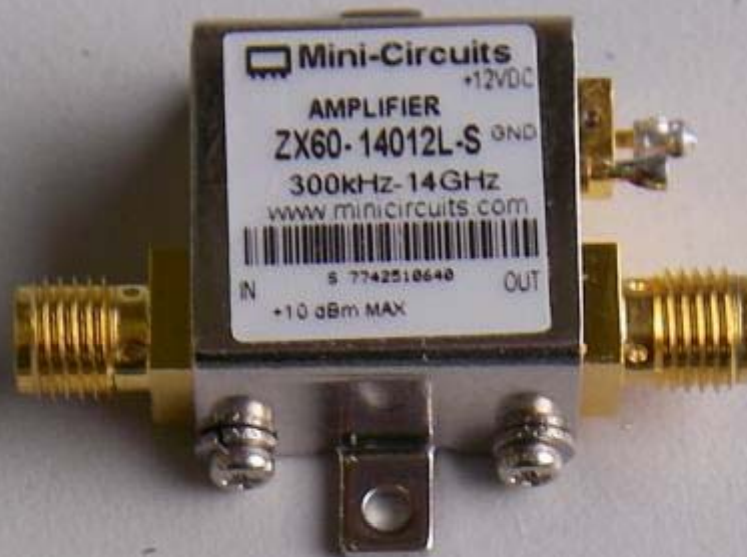
**Threshold discriminaton**

**CFD algo simulated**

Individual detectors	
$\sigma(\text{G01})$	32 ps
$\sigma(\text{G02})$	13 ps
$\sigma(\text{QBE})$	68 ps
$\sigma(\text{QBD})$	52 ps

**Target detector resolutions achieved. Still have to study efficiencies, and tune front-end electronics...**

# Electronics for Cerenkovs:





## Example of NIM module with mini-modules

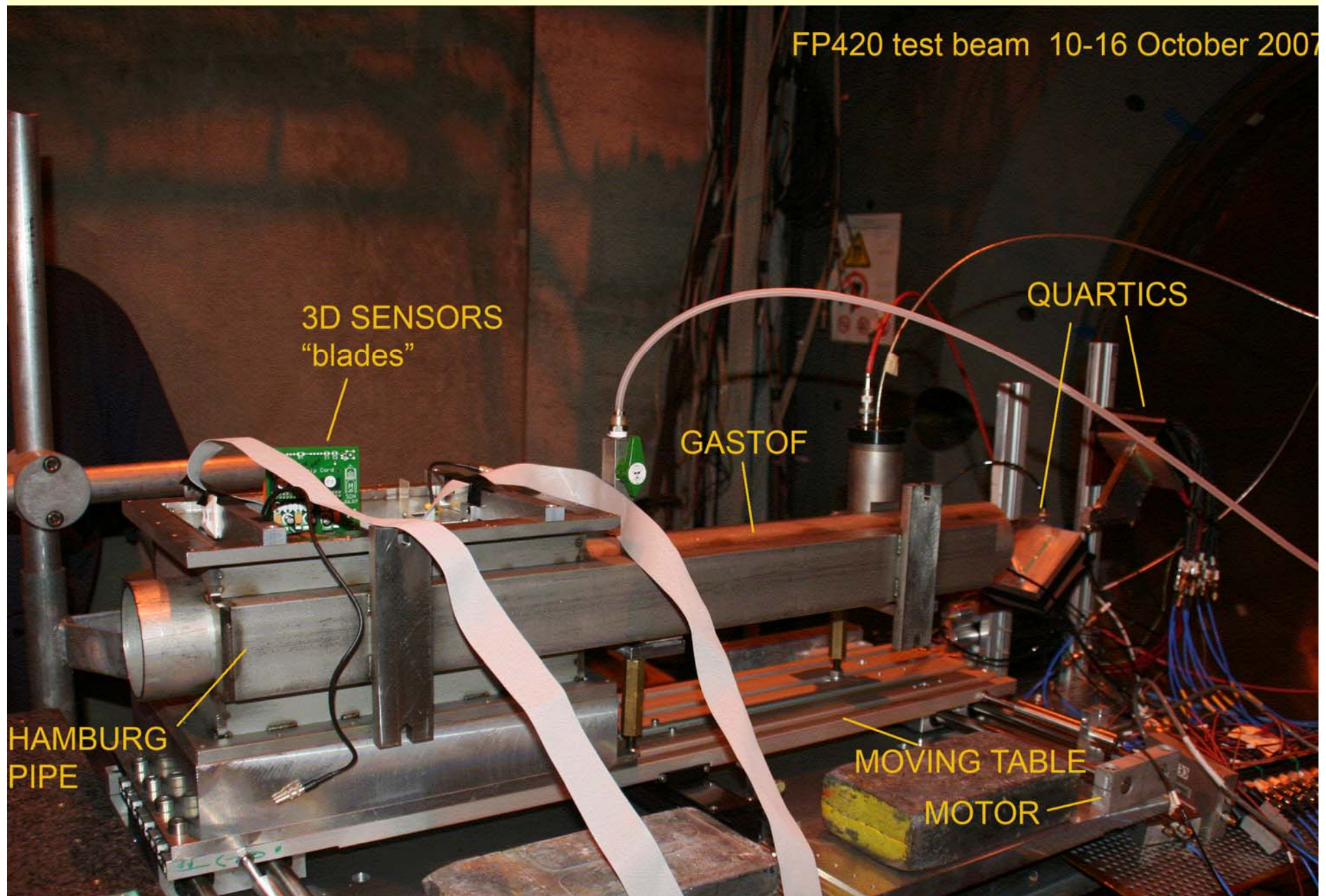
Two such modules with ~8 CFD mini-modules were built and are under tests

Note:

- It has remote control
- It is tuned for Burle&Hamamatsu MCP-PMTs
- It has double output







3D SENSORS  
"blades"

GASTOF

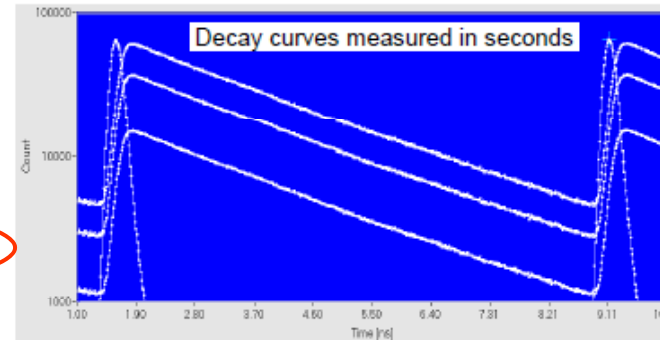
QUARTICS

HAMBURG  
PIPE

MOVING TABLE  
MOTOR

## Four Channel Time-Correlated Single Photon Counting Module

- ◆ Four Completely Parallel TCSPC Channels
- ◆ Ultra-High Data Throughput
- ◆ Overall Count Rate 32 MHz
- ◆ Channel Count Rate 8 MHz (Dead Time 125ns)
- ◆ Dual Memory Architecture: Readout during Measurement
- ◆ Reversed Start/Stop. Repetition Rates up to 200 MHz
- ◆ Electrical Time Resolution down to 8 ps FWHM / 5 ps rms
- ◆ Channel Resolution down to 813 fs
- ◆ Up to 4096 Time Channels / Curve
- ◆ Measurement Times down to 0.1 ms
- ◆ Software Versions for Windows 95 / 98 / NT
- ◆ Direct Interfacing to most Detector Types
- ◆ Single Decay Curve Mode
- ◆ Oscilloscope Mode
- ◆ Sequential Recording Mode
- ◆ Spectrum Scan Mode with 8 Independent Time Windows
- ◆ Continuous Flow Mode for Single Molecule Detection
- ◆ FIFO / Time Tag Mode for Single Molecule Detection



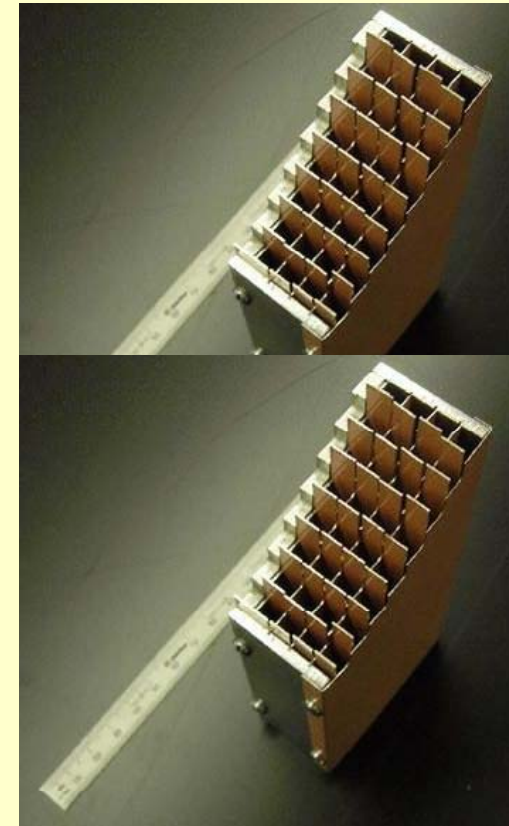
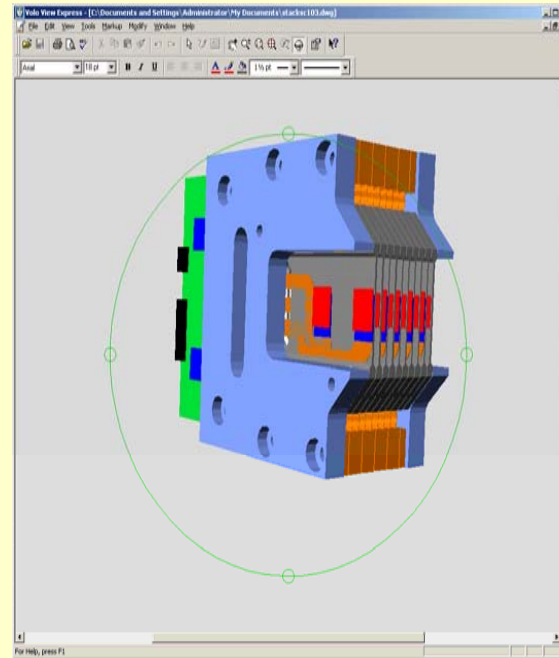
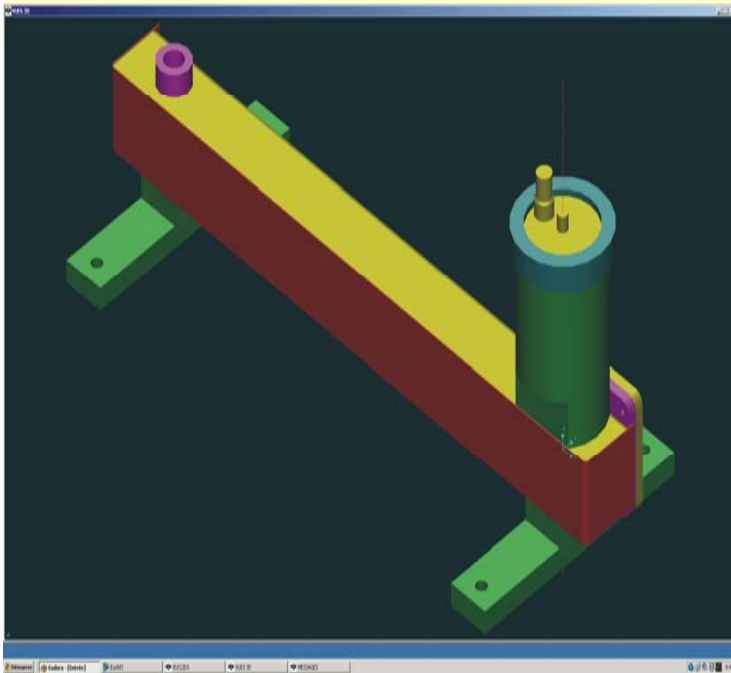


# Baseline Plan

1 GASTOF

Lots of silicon

2 QUARTICs



# Summary/Outlook

- Final cryostat design should be ready this Fall, and production could start early 2008... If accepted and allowed by LHC schedules installation in 2008/9 of one cryostat very profitable! Aim for full system installation in 2009/10.
- Moving HH pipe system well advanced; first irradiation tests underway; this October tests at CERN of a motorized and fully equipped (3D silicon + Cerenkov) section; complete RF studies show full compatibility with LHC
- Both Gastof and Quartic detector designs well advanced and with target resolutions (~10 and 50 ps per channel, respectively) achieved

# Summary/Outlook

- Need to finalize FE electronics and DAQ studies; complete background estimates/radiation hardness studies should be available by this Fall
- Tunnel infrastructure preparations has started as it requires long lead time for cabling and shieldings!
- All these studies in very close contact with machine vacuum, RF, collimator, instrumentation, ... LHC groups
- Many (all?) of these developments are very interesting and applicable at 220m!

Thanks