

Kinematics of Particle Beams

SPECIAL RELATIVITY

(in less than 60 minutes ...)

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(http://cern.ch/Werner.Herr/CAS/CAS2013_Chavannes/lectures/rel.pdf)



Why Special Relativity ?

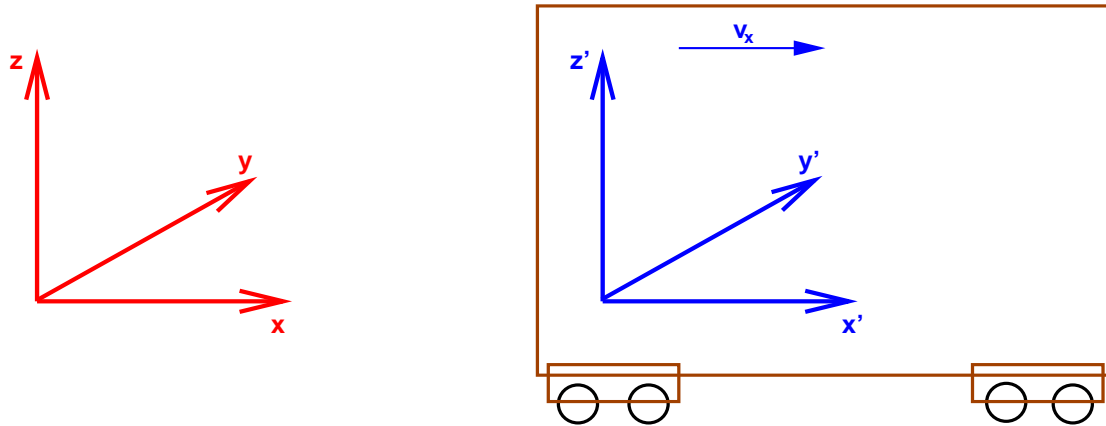
- Most beams at CERN are relativistic
 - Strong implications for beam dynamics:
 - Transverse dynamics (e.g. momentum compaction, radiation, ...)
 - Longitudinal dynamics (e.g. transition, ...)
 - Collective effects (e.g. space charge, beam-beam, ...)
 - Luminosity in colliders
 - Particle lifetime and decay (e.g. μ , π , Z_0 , Higgs, ...)
-

Small history

- 1678 (Römer, Huygens): Speed of light c is finite ($c \approx 3 \cdot 10^8$ m/s)
 - 1687 (Newton): **Principles of Relativity**
 - 1863 (Maxwell): Electromagnetic theory, light are waves moving through static ether
 - 1887 (Michelson, Morley): Speed c independent of direction, → ether theory R.I.P.
 - 1904 (Lorentz, Poincaré): **Lorentz transformations**
 - 1905 (Einstein): **Principles of Special Relativity**
 - 1907 (Minkowski): Concepts of Spacetime
-

Principles of Relativity (Newton)

- Assume a frame at rest (F) and another frame moving in x -direction (F') with constant velocity $\vec{v} = (v_x, 0, 0)$



Principles of Relativity (Newton)

- Assume a frame at rest (F) and another frame moving in x -direction (F') with constant velocity $\vec{v} = (v_x, 0, 0)$
 - Classical laws (mechanics) are the same in all frames (Newton, Poincaré)
 - No absolute space possible, but absolute time
Time is the same in all frames
 - Physical laws are invariant ...
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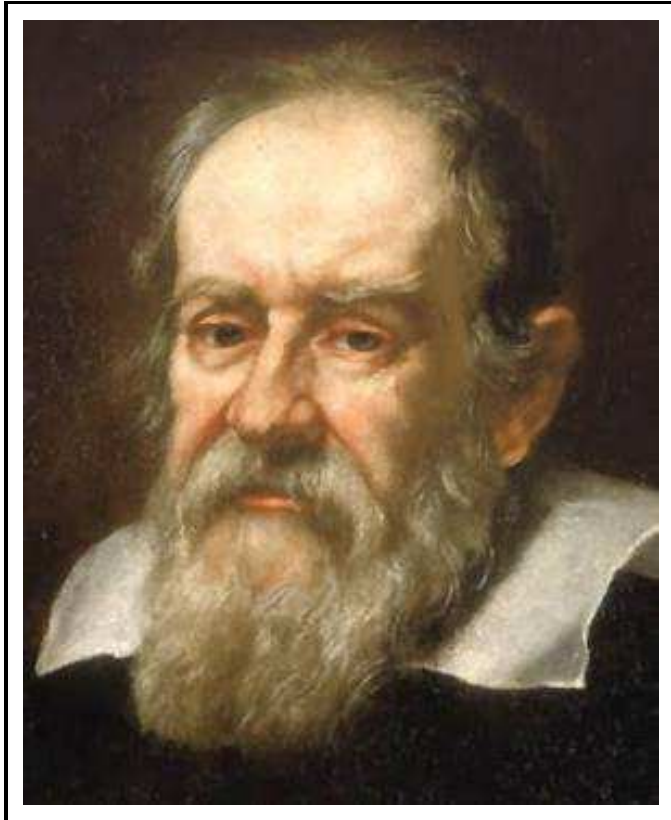
What does it mean ?

■ Invariance of physical laws:

- How is a physical process observed in F described (observed) in the moving frame F' ?
- Need transformation of coordinates (x, y, z) to describe (translate) results of measurements and observations to the moving system (x', y', z') .
- For Poincaré's, Newton's principle of relativity need Galilei transformation for
 $(x, y, z) \rightarrow (x', y', z')$



Galilei transformation



$$x' = x - v_x t$$

$$y' = y$$

$$z' = z$$

$$t' = t$$

Consequences of Galilei transformation

▣ Velocities can be added

➤ From Galilei transformation, take derivative:

$$x' = x - v_x t$$

$$\dot{x}' = \dot{x} - v_x \quad \rightarrow \quad v' = v - v_x$$

➤ A car moving with speed v' in a frame moving with speed v_x we have in rest frame $v = v' + v_x$

➤ But: if $v' = 0.75c$ and $v_x = 0.75c$
do we get $v = 1.5c$?

Problems with Galilei transformation

- Maxwell's equations are wrong when Galilei transformations are applied (because they predict the speed of light, see later)
 - First solution: introduction of "ether"
 - But: speed of light the same in all frames and all directions (no "ether")
 - Need other transformations for Maxwell's equations
 - Introduced principles of special relativity
-

Principles of Special Relativity (Einstein)

- A frame moving with constant velocity is called an "inertial frame"
- All (not only classical) physical laws in related frames have equivalent forms, in particular:
 - speed of light c the same in all frames
- Cannot distinguish between inertial frames, in particular:
 - Cannot determine absolute speed of an inertial frame
 - No absolute space, no absolute time

All you need to know !

Principles of Invariance (Poincaré)

Concept of **Invariance**:

The laws of Physics are **invariant** under a transformation between two coordinate frames moving at constant velocities with respect to each other

(The **world** is not invariant, but the laws of physics are !)

Poincaré + Einstein:

Need **Transformations** (not Galilean) which make the physics laws the same everywhere !

Coordinates must be transformed differently

- Transformation must keep speed of light constant
- Time must be changed by transformation as well as space coordinates
- Transform $(x, y, z), t \rightarrow (x', y', z'), t'$

Constant speed of light requires:

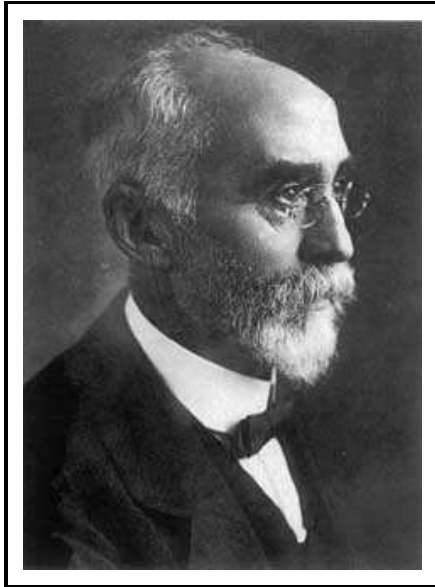
$$x^2 + y^2 + z^2 - c^2t^2 = 0 \rightarrow x'^2 + y'^2 + z'^2 - c^2t'^2 = 0$$

(front of a light wave)

- Defines the Lorentz transformation
(but established by Poincaré !)



Lorentz transformation



$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}} = \gamma \cdot (x - vt)$$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v \cdot x}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}} = \gamma \cdot \left(t - \frac{v \cdot x}{c^2}\right)$$

- Transformation for constant velocity along x-axis



Definitions: relativistic factors

$$\beta_r = \frac{v}{c}$$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{1}{\sqrt{1 - \beta_r^2}}$$

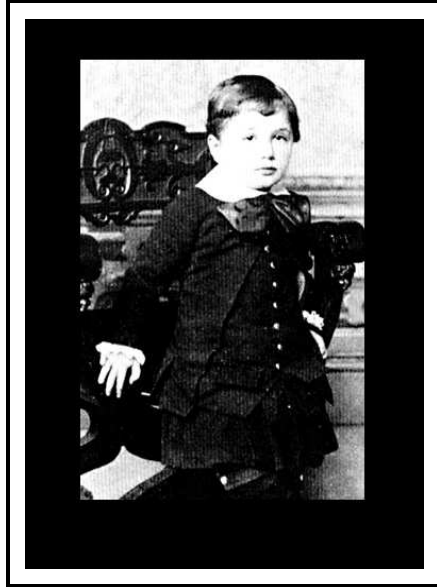
➤ β_r relativistic speed: $\beta_r = [0, 1]$

➤ γ relativistic factor: $\gamma = [1, \infty]$

(unfortunately, you will also see other β and γ ... !)



Einstein's contributions



$$x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$y' = y$$

$$z' = z$$

$$t' = \frac{t - \frac{v \cdot x}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

$$(x, y, z) \rightarrow (x, y, z, ct)$$

- Time has no absolute meaning
- Simultaneity has no absolute meaning
- Combine time with the 3 dimensions of space
- Energy and mass equivalence



Consequences of Einstein's interpretation

■ Relativistic phenomena:

- (Non-) Simultaneity of events in independent frames
- Lorentz contraction
- Time dilation

■ Formalism with four-vectors introduced

- Invariant quantities
- Mass - energy relation



Simultaneity between moving frames

- Assume two events in frame F at positions x_1 and x_2 happen simultaneously at times $t_1 = t_2$:

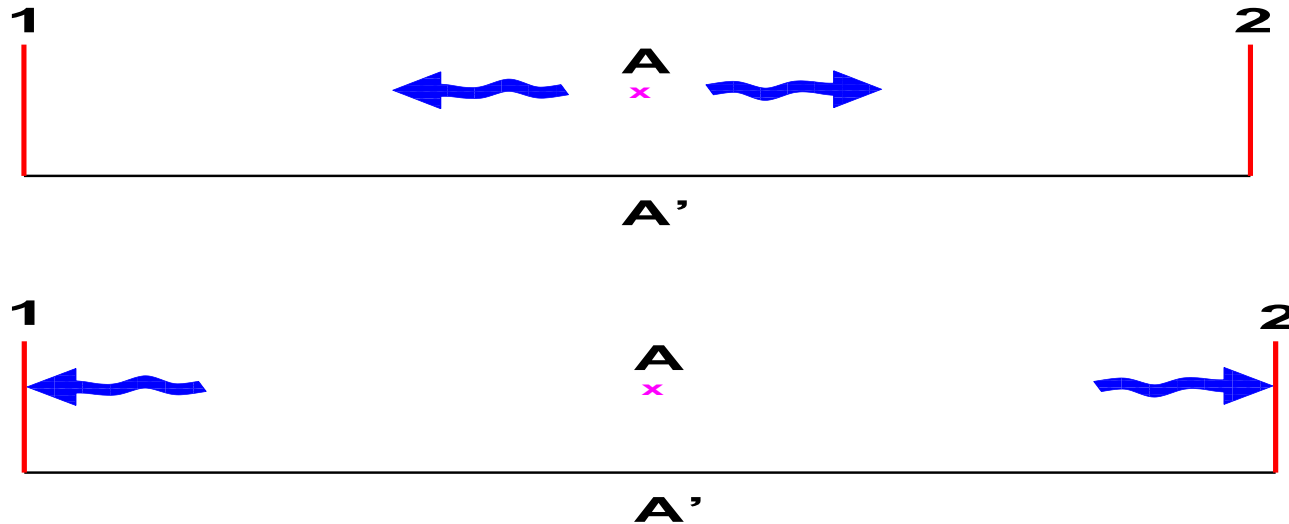
$$t'_1 = \frac{t_1 - \frac{v \cdot x_1}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}} \quad \text{and} \quad t'_2 = \frac{t_2 - \frac{v \cdot x_2}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}}$$

implies that $t'_1 \neq t'_2$ in frame F' !!

- Two events simultaneous at positions x_1 and x_2 in F are not simultaneous in F'

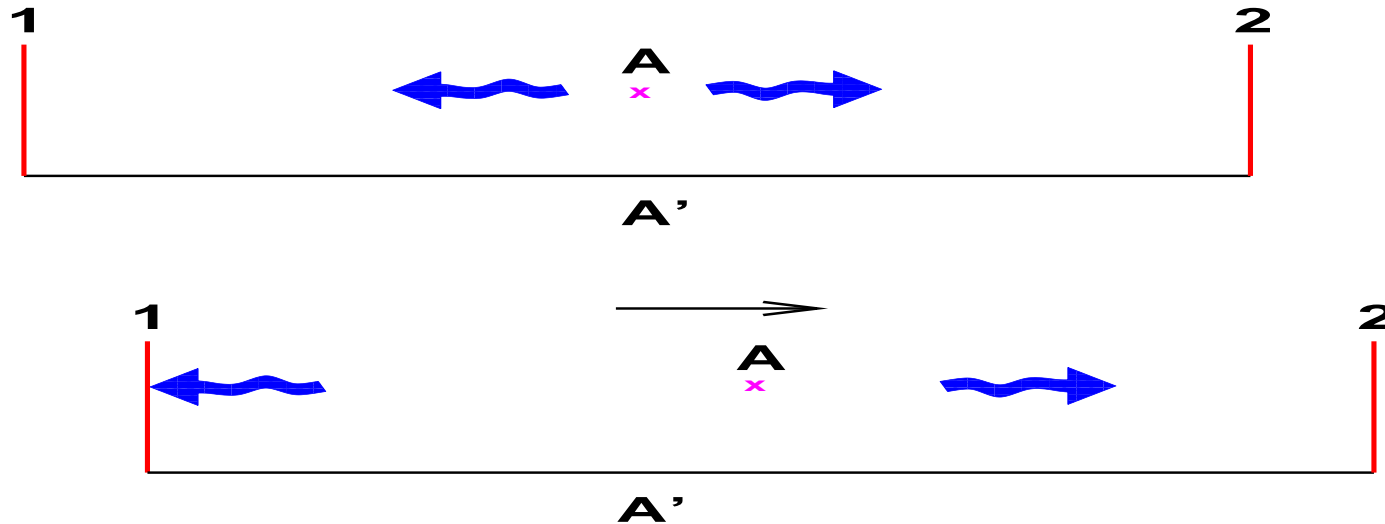


Simultaneity between moving frames



- System with a light source (x) and detectors ($1, 2$) and one observer (A) in this frame, another (A') outside
- System at rest \rightarrow observation the same in A and A'
- What if system with A is moving ?

Simultaneity between moving frames



- For A: both flashes arrive simultaneously in 1,2
- For A': flash arrives first in 1, later in 2
- A simultaneous event in F is not simultaneous in F'
- Why do we care ??

Why care about simultaneity ?

- Simultaneity is **not** frame independent
- This is a key in special relativity
- Most paradoxes are explained by that (although not the twin paradox) !
- Different observers see a different reality

→ relativity is not a spectator sport ...

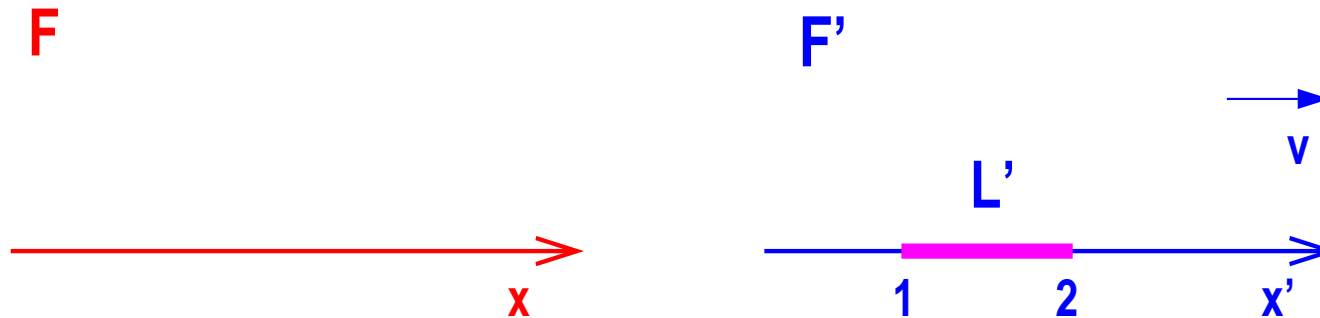


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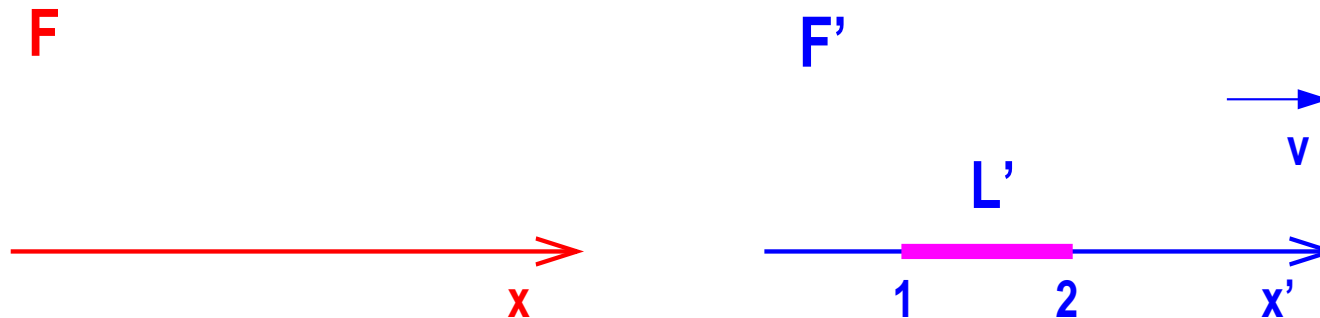
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- More important: sequence of events can change !
 - For $t_1 < t_2$ we may find (not always !) a frame where $t_1 > t_2$ (concept of **before** and **after** depends on the observer)
 - Requires introduction of "antiparticles" in relativistic quantum mechanics
-

Consequences: length measurement



Length of a rod in F' is $L' = x'_2 - x'_1$, measured simultaneously (!) at a fixed time t' in frame F', what is the length L seen in F ??

Consequences: length measurement



We have to measure simultaneously (!) the ends of the rod at a fixed time t in frame **F**, i.e.: $L = x_2 - x_1$ →

$$x'_1 = \gamma \cdot (x_1 - vt) \quad \text{and} \quad x'_2 = \gamma \cdot (x_2 - vt)$$

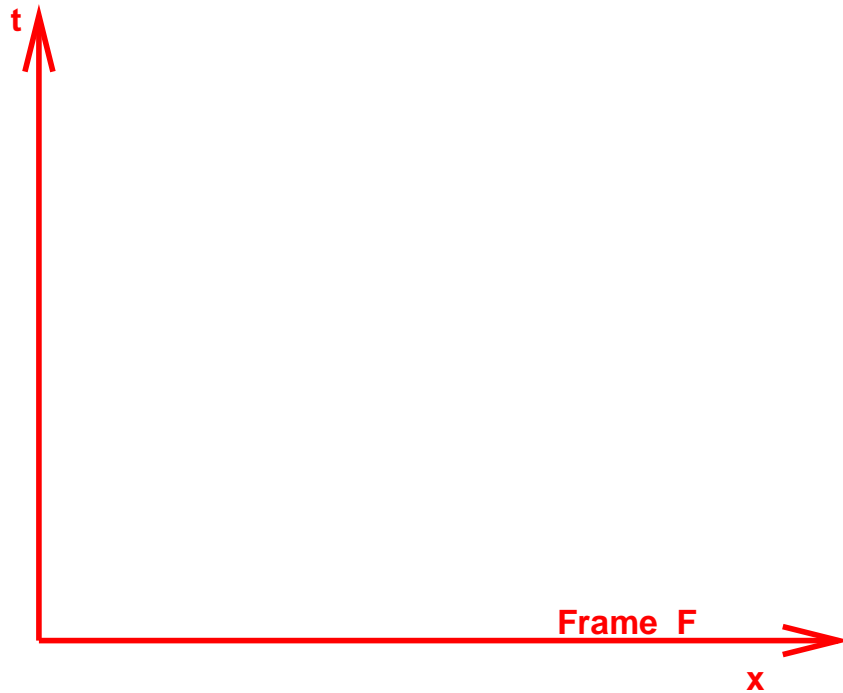
$$L' = x'_2 - x'_1 = \gamma \cdot (x_2 - x_1) = \gamma \cdot L$$

$$\rightarrow L = L' / \gamma$$

Lorentz contraction

- In moving frame an object has always the same length (it is invariant, our principle !)
- From stationary frame moving objects appear contracted by a factor γ (Lorentz contraction)
- Why do we care ?
- Turn the argument around: assume length of a proton bunch appears always at 0.1 m in laboratory frame (e.g. in the RF bucket), what is the length in its own (moving) frame ?
 - At 5 GeV ($\gamma \approx 5.3$) $\rightarrow L' = 0.53$ m
 - At 450 GeV ($\gamma \approx 480$) $\rightarrow L' = 48.0$ m

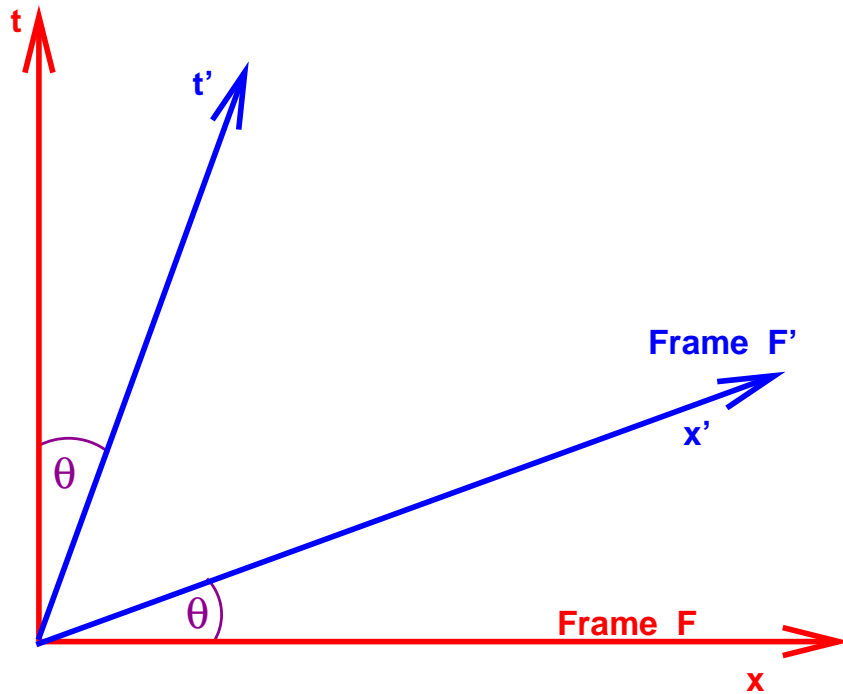
Lorentz transformation - schematic



- Rest frame (x only, difficult to draw many dimensions)
y and z coordinates are not changed (transformed)



Lorentz transformation - schematic

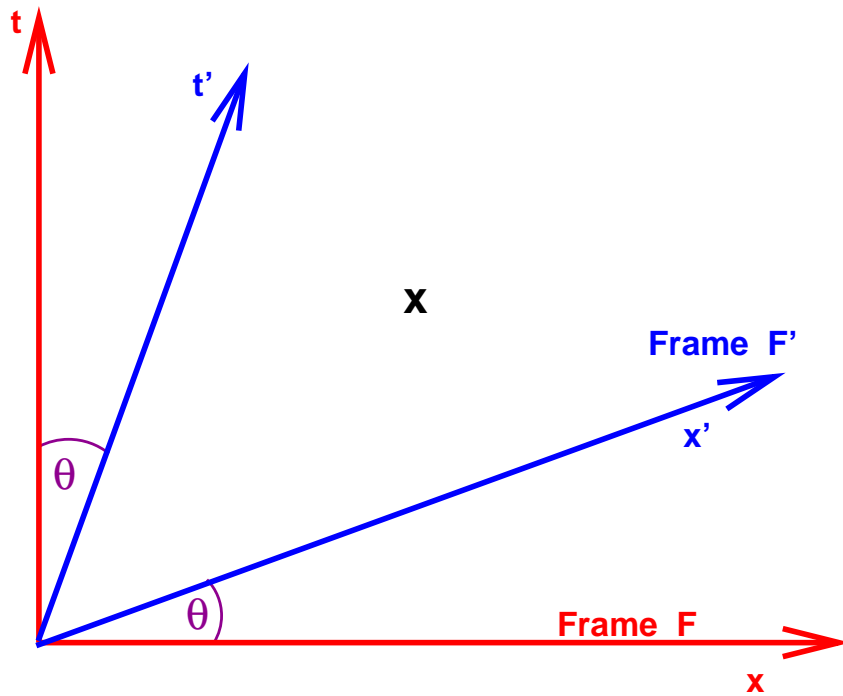


➤ Rest frame and moving frame

➤ $\tan(\theta) = \frac{v}{c}$



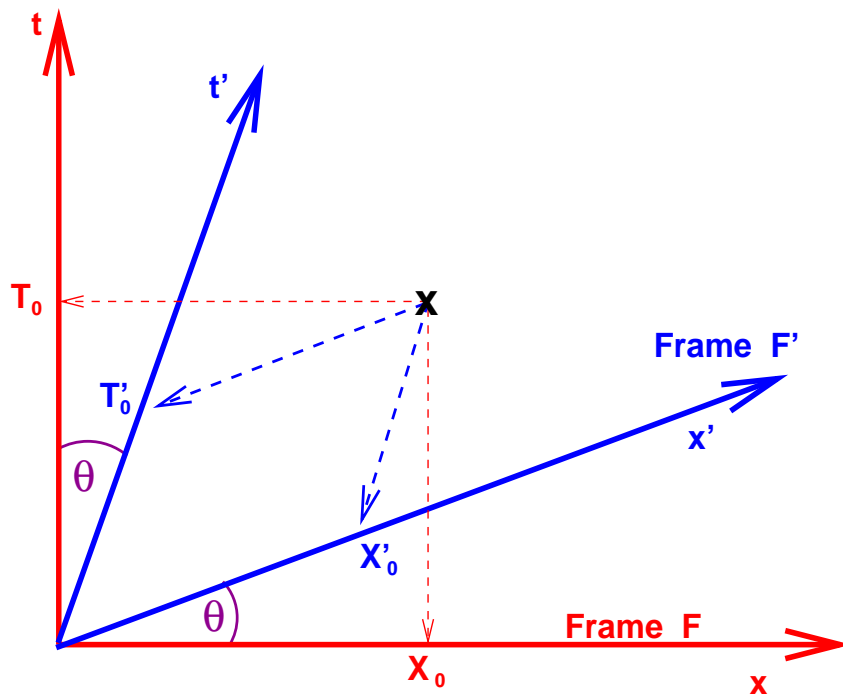
Lorentz transformation - schematic



➤ An event X



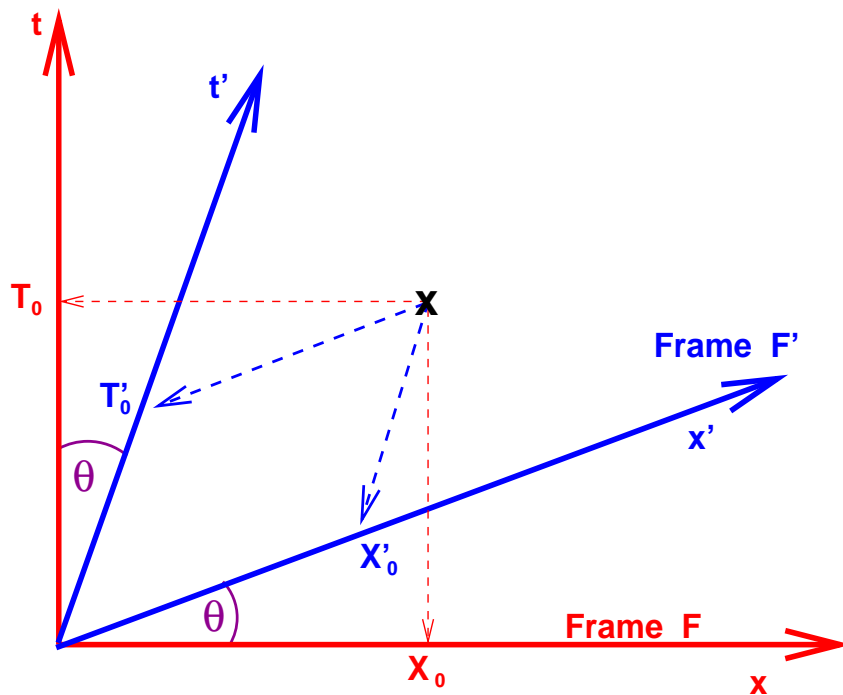
Lorentz transformation - schematic



- Event X seen at different time and location in the two frames, projected on axes of F and F'



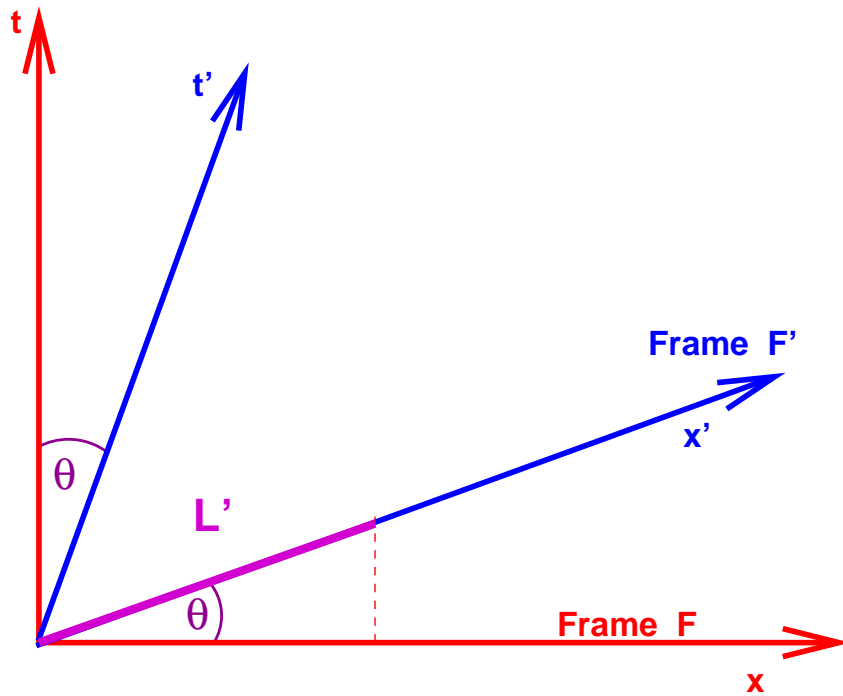
Lorentz transformation - schematic



➤ Q: How would a Galilei-transformation look like ??



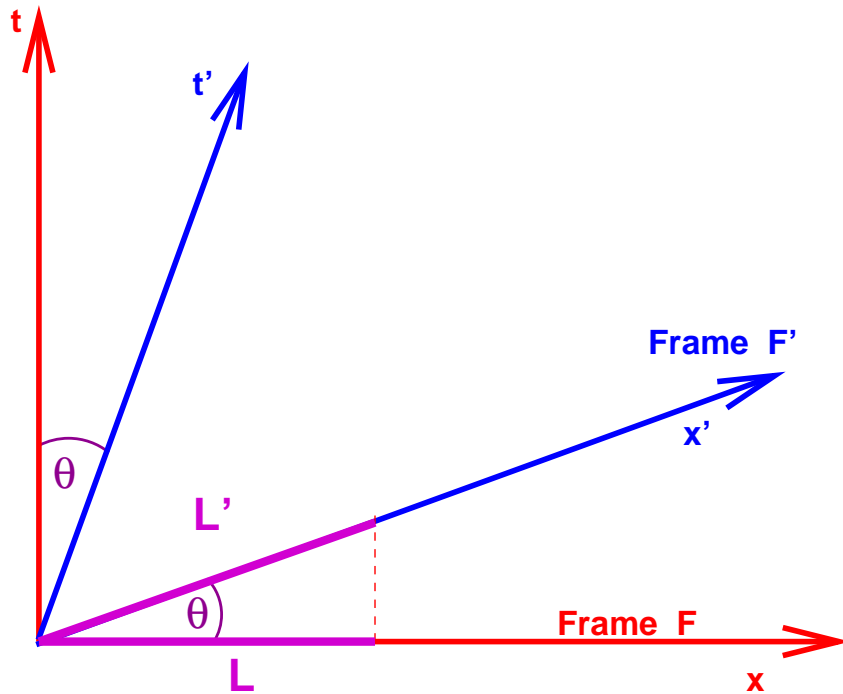
Lorentz contraction - schematic



➤ Length L' as measured in moving frame



Lorentz contraction - schematic



- From moving frame: L appears shorter in rest frame
- Length is maximum in frame (F') where object is at rest

Lorentz contraction

For the coffee break and lunch:



Could you "see" (visually) a Lorentz contraction ??
(if you run fast enough ...)



Time dilation

A clock measures time difference $\Delta t = t_2 - t_1$ in frame F, measured **at fixed position x**, what is the time difference $\Delta t' = t'_2 - t'_1$ as measured from the moving frame F' ??

For Lorentz transformation of time in moving frame we have:

$$t'_1 = \gamma\left(t_1 - \frac{v \cdot x}{c^2}\right) \quad \text{and} \quad t'_2 = \gamma\left(t_2 - \frac{v \cdot x}{c^2}\right)$$

$$\Delta t' = t'_2 - t'_1 = \gamma \cdot (t_2 - t_1) = \gamma \cdot \Delta t$$

$$\rightarrow \Delta t' = \gamma \Delta t$$

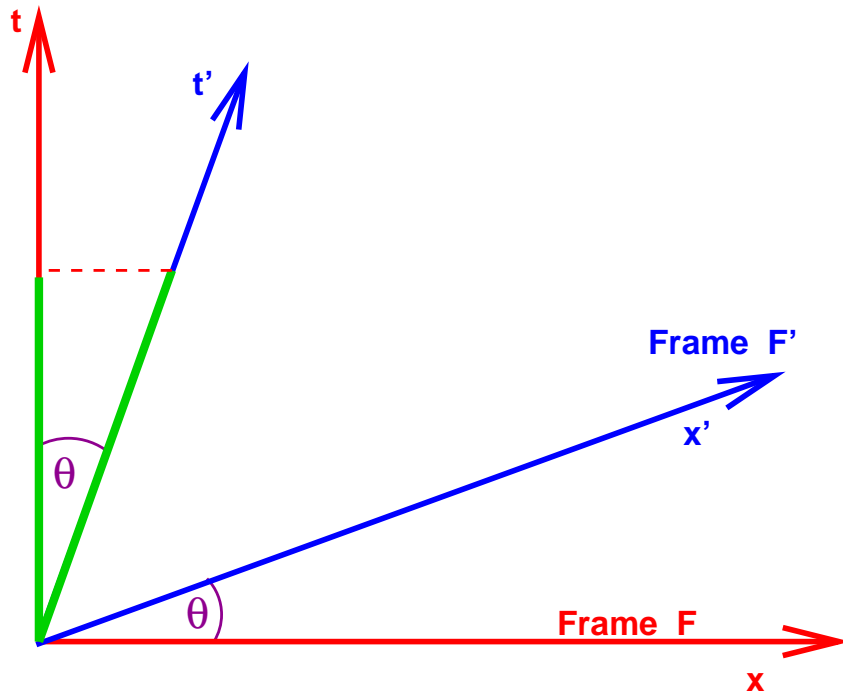


Time dilation

- In moving frame time appears to run slower
- Why do we care ?
 - μ have lifetime of $2 \mu\text{s}$ ($\equiv 600 \text{ m}$)
 - For $\gamma \geq 150$, they survive 100 km to reach earth from upper atmosphere
 - They can survive more than $2 \mu\text{s}$ in a μ -collider
 - Generation of neutrinos from the SPS beams



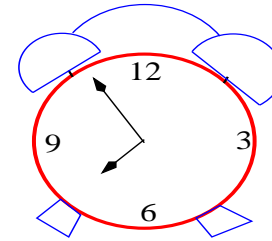
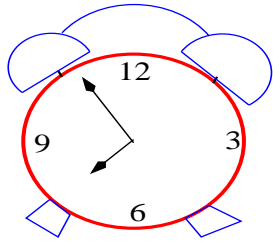
Time dilation - schematic



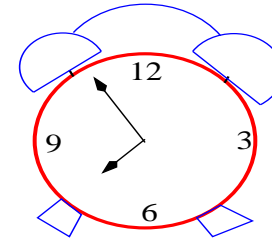
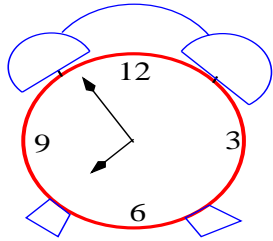
- From moving frame: time goes slower in rest frame
- Time shortest in frame (F') where object is at rest



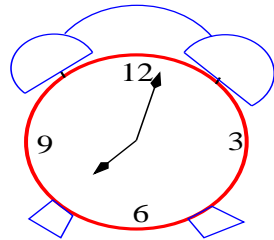
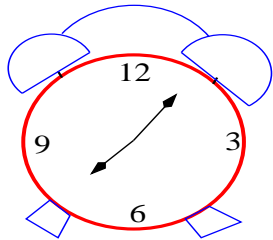
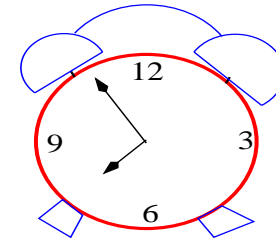
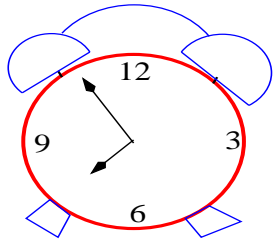
Moving clocks go slower



Moving clocks go slower



Ten minutes later ...



Travel by airplane:

On a flight from Montreal to Geneva, the time is slower by
25 - 30 ns !*)

*) (unfortunately there is a catch 22 ...)



Addition of velocities

➤ Galilei: $v = v_1 + v_2$

➤ With Lorentz transform we have:

$$v = \frac{v_1 + v_2}{1 + \frac{v_1 v_2}{c^2}} \quad \text{or equivalently :} \quad \beta = \frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2}$$

for $\beta = 0.5$ we get:

$$0.5c + 0.5c = 0.8c$$

$$0.5c + 0.5c + 0.5c = 0.93c$$

$$0.5c + 0.5c + 0.5c + 0.5c = 0.976c$$

$$0.5c + 0.5c + 0.5c + 0.5c + 0.5c = 0.992c$$

➔ Nothing can go faster than the speed of light ...



First summary

- Physics laws the same in different moving frames ...
 - Speed of light is maximum possible speed
 - Constant speed of light requires Lorentz transformation
 - Moving objects appear shorter
 - Moving clocks seem to go slower
 - No absolute space or time !
- Now: applications and how to calculate something ...



Introducing four-vectors

Four-vector: $F = (f_1, f_2, f_3, f_4)$

a vector with four components

Example: position four-vector $X = (ct, x, y, z) = (ct, \vec{x})$

This mathematical setting is called **Minkowski space** and Lorentz transformation can be written in matrix form:

$$\begin{pmatrix} ct' \\ x' \\ y' \\ z' \end{pmatrix} = \begin{pmatrix} \gamma & \frac{-\gamma v}{c} & 0 & 0 \\ \frac{-\gamma v}{c} & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \circ \begin{pmatrix} ct \\ x \\ y \\ z \end{pmatrix}$$

$$X' = M_L \circ X$$


Introducing four-vectors

Define a scalar product^{*)} like: $X \diamond Y$

$$X = (x_0, \vec{x}), \quad Y = (y_0, \vec{y}) \quad \rightarrow \quad X \diamond Y = x_0 \cdot y_0 - \vec{x} \cdot \vec{y}$$

For example try $X \diamond X$ $(ct, \vec{x}) \diamond (ct, \vec{x})$:

$$X \diamond X = c^2 t^2 - x^2 - y^2 - z^2$$

This product is an **invariant**, i.e.:

$$X \diamond X = c^2 t^2 - x^2 - y^2 - z^2 = X' \diamond X' = c^2 t'^2 - x'^2 - y'^2 - z'^2$$

Invariant Quantities have the **same value** in all inertial frames

***) definition of product not unique ! (I use PDG 2008)**



Why bother about four-vectors ?

- We have seen the importance of **invariants**:
 - Ensure equivalence of physics laws in different frames
 - ➔ The solution: write the laws of physics in terms of **four vectors**
 - ➔ Any four-vector (scalar) product $F \diamond F$ has the same value in all coordinate frames moving at constant velocities with respect to each other ...
(remember that phrase ?)
-

Using four-vectors

We can describe a **distance** in the spacetime between two points X_1 and X_2 :

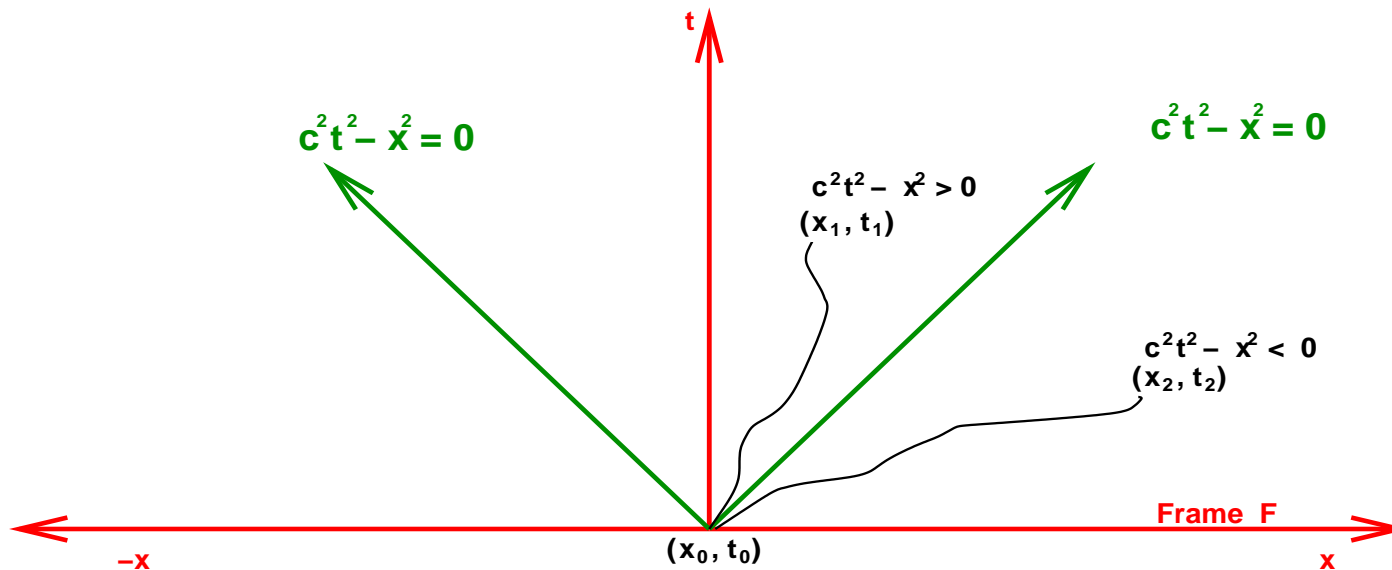
$$\Delta X = X_2 - X_1 = (ct_2 - ct_1, x_2 - x_1, y_2 - y_1, z_2 - z_1)$$

$$\Delta s^2 = \Delta X \diamond \Delta X = c^2 \Delta t^2 - \Delta x^2 - \Delta y^2 - \Delta z^2$$

Δs^2 can be positive (timelike) or negative (spacelike)



Moving in Minkowski space



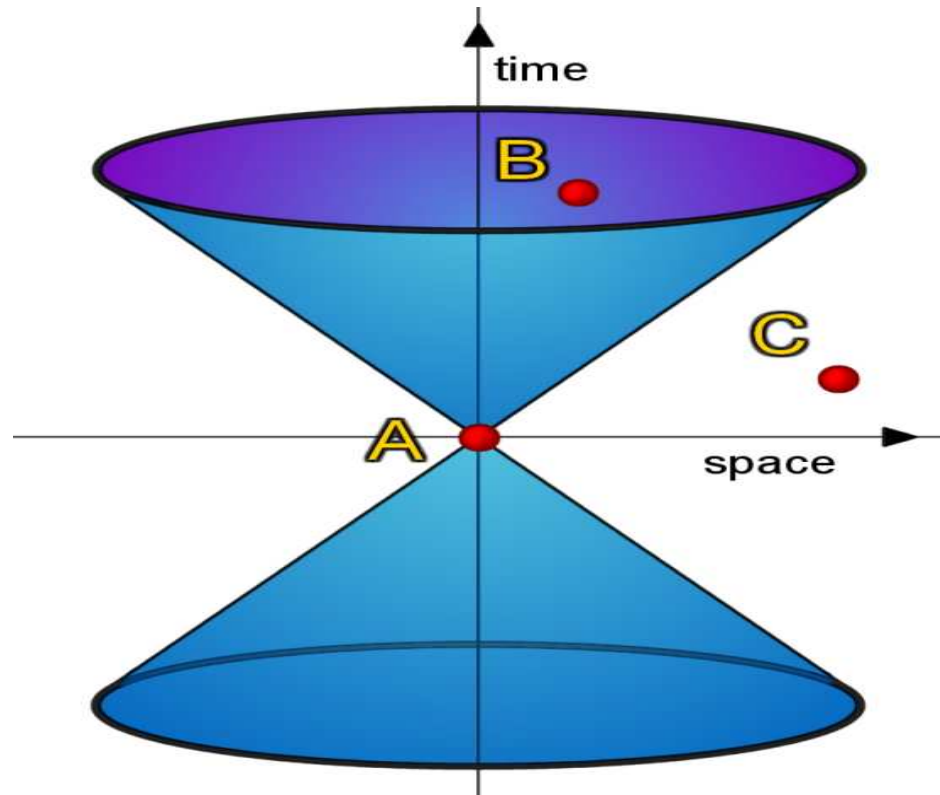
➤ Light travels with $c^2 \cdot t^2 - x^2 = 0$

➤ Particle travels with $c^2 \cdot t^2 - x^2 > 0$ (allowed)

➤ Particle travels with $c^2 \cdot t^2 - x^2 < 0$ (not allowed)

➡ Allowed region defines **light cone**

Light cone ...



➤ Distances: timelike (AB), spacelike (AC)

Using four-vectors

Special case (time interval $\vec{x}_2 = \vec{x}_1 + \vec{v}\Delta t$):

$$c^2\Delta t^2 - \Delta x^2 - \Delta y^2 - \Delta z^2 = c^2\Delta t^2\left(1 - \frac{v^2}{c^2}\right) = c^2\left(\frac{\Delta t}{\gamma}\right)^2 = c^2\Delta\tau^2$$

- $\Delta\tau$ is the time interval measured in the moving frame
- τ is a fundamental time: **proper time** τ



The meaning of "proper time"

$\Delta\tau$ is the time interval measured inside the moving frame

Back to μ -decay

- μ lifetime is $\approx 2 \mu\text{s}$
 - μ decay in $\approx 2 \mu\text{s}$ in their frame, i.e. using the "proper time"
 - μ decay in $\approx \gamma \cdot 2 \mu\text{s}$ in the laboratory frame, i.e. earth
 - μ appear to live longer than $2 \mu\text{s}$ in the laboratory frame, i.e. earth
-

The meaning of "proper time"

■ How to make neutrinos ?? Let pions decay: $\pi \rightarrow \mu + \nu_\mu$

➤ π -mesons have lifetime of $2.6 \cdot 10^{-8}$ s (i.e. 7.8 m)

➤ For 40 GeV π -mesons: $\gamma = 288$

➤ In laboratory frame: decay length is 2.25 km
(required length of decay tunnel)

■ VERY intuitive (quote A. Einstein, modified):

➤ It's 7:00 a.m. in bed, you close your eyes for $\Delta\tau = 10$ minutes, it's 9:00 a.m.

➤ It's 7:00 a.m. in a meeting, you close your eyes for $\Delta\tau = 10$ minutes, it's 7:01 a.m.

More four-vectors

Position four-vector X :

$$X = (ct, x, y, z) = (ct, \vec{x})$$

Velocity four-vector V :

$$V = \frac{dX}{d\tau} = \gamma \frac{dX}{dt} = \gamma \dot{X} = \gamma \left(\frac{d(ct)}{dt}, \dot{x}, \dot{y}, \dot{z} \right) = \gamma(c, \vec{\dot{x}}) = \gamma(c, \vec{v})$$

Please note that:

$$V \diamond V = \gamma^2(c^2 - \vec{v}^2) = c^2!!$$

- c is an invariant (of course), has the same value in all inertial frames



More four-vectors

Momentum four-vector P :

$$P = m_0 V = m_0 \gamma (c, \vec{v}) = (\mathbf{m}c, \vec{p})$$

using:

m_0 (mass of a particle)

$\mathbf{m} \equiv m_0 \cdot \gamma$ (relativistic mass)

$\vec{p} = \mathbf{m} \cdot \vec{v} = m_0 \gamma \vec{v}$ (relativistic 3-momentum)

We can get another invariant: $P \diamond P = m_0^2 (V \diamond V) = m_0^2 c^2$

Invariant of the four-momentum vector is the mass m_0

→ The rest mass is the same in all frames (thanks a lot ..)

Still more four-vectors

Force four-vector F :

$$F = \frac{dP}{d\tau} = \gamma \frac{dP}{dt} = \gamma \frac{d}{dt}(mc, \vec{p}) = \gamma \left(c \frac{dm}{dt}, \frac{d\vec{p}}{dt} \right) = \gamma \left(\frac{1}{c} \frac{dE}{dt}, \frac{d\vec{p}}{dt} \right)$$

and we had already more four-vectors:

Coordinates : $X = (ct, x, y, z) = (ct, \vec{x})$

Velocities : $V = \gamma(c, \vec{\dot{x}}) = \gamma(c, \vec{v})$

Momenta : $P = m_0 V = m_0 \gamma(c, \vec{v}) = (\mathbf{m}c, \vec{p})$

$X \diamond X, V \diamond V, P \diamond P, P \diamond X, V \diamond F, \dots$ are **ALL** invariants

Dynamics with four-vectors

We compute: $V \diamond F = 0$

All right, 0 is the same in all frames, sounds useless, but:

$$V \diamond F = 0 \quad \longrightarrow \quad \frac{d}{dt}(mc^2) - \vec{f}\vec{v} = 0$$

Now $\vec{f}\vec{v}$ is rate of change of kinetic energy dT/dt
after integration:

$$T = \int \frac{dT}{dt} dt = \int \vec{f}\vec{v} dt = \int \frac{d(mc^2)}{dt} dt = mc^2 + const.$$

$$T = mc^2 + const. = mc^2 - m_0c^2$$



Relativistic energy

Interpretation:

$$E = mc^2 = T + m_0c^2$$

- Total energy E is $E = mc^2$
- Sum of kinetic energy plus rest energy
- Energy of particle at rest is $E_0 = m_0c^2$

$$E = m \cdot c^2 = \gamma m_0 \cdot c^2$$

using the definition of relativistic mass again: $m = \gamma m_0$

Still more four-vectors

Equivalent four-momentum vector (using E instead of m):

$$P = (mc, \vec{p}) \rightarrow (E/c, \vec{p})$$

then:

$$P \diamond P = m_0^2 c^2 = \frac{E^2}{c^2} - \vec{p}^2$$

follows:

$$E^2 = \vec{p}^2 c^2 + m_0^2 c^4$$

another familiar expression ...

Relativistic energy

These units are not very convenient:

$$m_p = 1.672 \cdot 10^{-27} \text{ Kg}$$

$$\rightarrow m_p c^2 = 1.505 \cdot 10^{-10} \text{ J}$$

$$\rightarrow m_p c^2 = 938 \text{ MeV} \rightarrow m_p = 938 \text{ MeV}/c^2$$

$$\rightarrow m_p c^2 \cdot \gamma(7 \text{ TeV}) = 1.123 \cdot 10^{-6} \text{ J}$$

$$\rightarrow m_p c^2 \cdot \gamma(7 \text{ TeV}) \cdot 1.15 \cdot 10^{11} \cdot 2808 = 360 \cdot 10^6 \text{ J}$$



Practical units

In particle physics: omit c and dump it into the units:

$$[E] = \text{eV} \quad [p] = \text{eV}/c \quad [m] = \text{eV}/c^2$$

Four-vectors get an easier form:

$$P = (m, \vec{p}) = (E, \vec{p})$$

and from $P \diamond P = E^2 - p^2 = m_0^2$ follows directly:

$$E^2 = \vec{p}^2 + m_0^2 \quad (= m^2 = \gamma^2 m_0^2)$$



Relativistic energy

Note:

$$E = mc^2 = \gamma \cdot m_0 c^2 \quad \rightarrow \quad E = \gamma m_0 c^2$$

$$p = m_0 \gamma v = \gamma m_0 \cdot \beta c \quad \rightarrow \quad p = \gamma m_0 \cdot \beta c$$

$$T = m_0(\gamma - 1) \cdot c^2 \quad \rightarrow \quad T = \gamma m_0 c^2 - m_0 c^2$$



Interpretation of relativistic energy

- For any object, $m \cdot c^2$ is the total energy
 - Object can be composite, like proton ..
 - m is the mass (energy) of the object "in motion"
 - m_0 is the mass (energy) of the object "at rest"
- For discussion: what is the mass of a photon ?



Relativistic mass

The mass of a fast moving particle is increasing like:

$$m = \gamma m_0 = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

assume a 75 kg heavy man:

- Rocket at 100 km/s, $\gamma = 1.00000001$, $m = 75.000001$ kg
 - PS at 26 GeV, $\gamma = 27.7$, $m = 2.08$ tons
 - LHC at 7 TeV, $\gamma = 7642$, $m = 573.15$ tons
 - LEP at 100 GeV, $\gamma = 196000$, $m = 14700$ tons
-

Relativistic mass

➤ Why do we care ?

$$m = \gamma m_0 = \frac{m_0}{\sqrt{1 - \frac{v^2}{c^2}}}$$

➤ Particles cannot go faster than c !

➤ What happens when we accelerate ?



Relativistic mass

When we accelerate:

■ For $v \ll c$:

➤ E, m, p, v increase ...

■ For $v \approx c$:

➤ E, m, p increase, but v does not !

➤ Remember that for later



Relativistic energy

Since we remember that:

$$T = m_0(\gamma - 1)c^2$$

therefore:

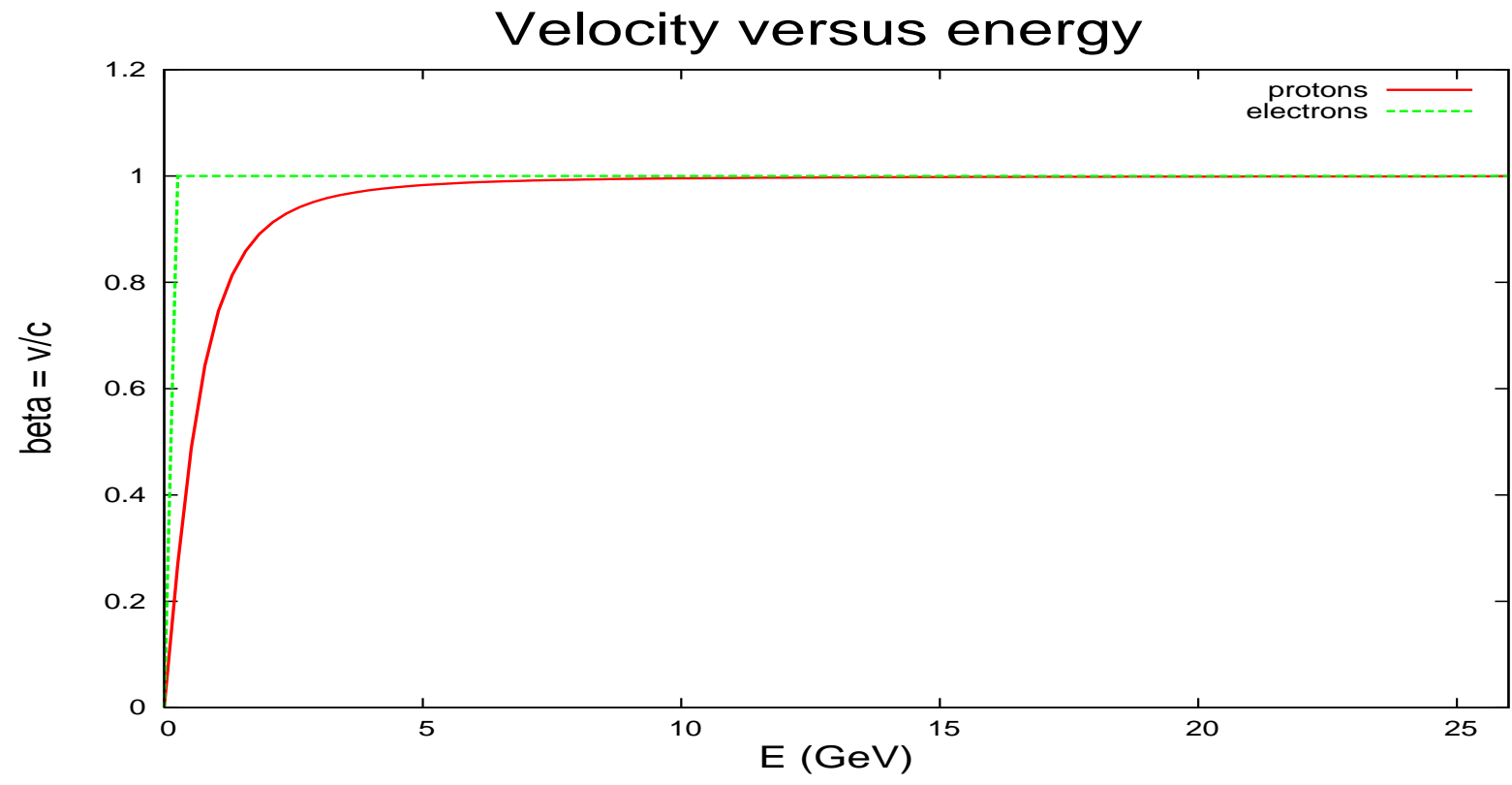
$$\gamma = 1 + \frac{T}{m_0c^2}$$

we get for the speed v , i.e. β :

$$\beta = \sqrt{1 - \frac{1}{\gamma^2}}$$



Velocity versus energy (protons)



Why do we care ??

E (GeV)	v (km/s)	γ	β	T (LHC)
450	299791.82	479.74	0.99999787	88.92465 μs
7000	299792.455	7462.7	0.999999999	88.92446 μs

- For identical circumference very small change in revolution time
- If path for faster particle slightly longer, the faster particle arrives later !



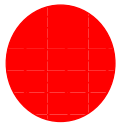
Four vectors

- Use of four-vectors simplify calculations significantly
- Follow the rules and look for invariants
- In particular kinematic relationships, e.g.
 - Particle decay (find mass of parent particle)
 - Particle collisions →

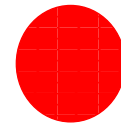


Particle collisions

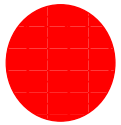
P1



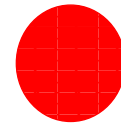
P2



P1



P2



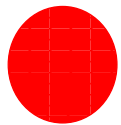
➤ What is the available collision energy ?



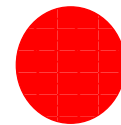
Particle collisions - collider

Assume identical particles and beam energies, colliding head-on

P1



P2



The four momentum vectors are:

$$P1 = (E, \vec{p}) \quad P2 = (E, -\vec{p})$$

The four momentum vector in centre of mass system is:

$$P^* = P1 + P2 = (E + E, \vec{p} - \vec{p}) = (2E, \vec{0})$$



Particle collisions - collider

The four momentum vector in centre of mass system is:

$$P^* = P_1 + P_2 = (E + E, \vec{p} - \vec{p}) = (2E, \vec{0})$$

The square of the total available energy s in the centre of mass system is the momentum invariant:

$$s = P^* \diamond P^* = 4E^2$$

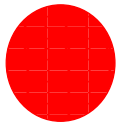
$$E_{cm} = \sqrt{P^* \diamond P^*} = 2E$$

i.e. in a (symmetric) collider the total energy is twice the beam energy

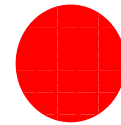


Particle collisions - fixed target

P1



P2



The four momentum vectors are:

$$P1 = (E, \vec{p}) \quad P2 = (m_0, \vec{0})$$

The four momentum vector in centre of mass system is:

$$P^* = P1 + P2 = (E + m_0, \vec{p})$$



Particle collisions - fixed target

With the above it follows:

$$P^* \diamond P^* = E^2 + 2m_0E + m_0^2 - \vec{p}^2$$

since $E^2 - \vec{p}^2 = m_0^2$ we get:

$$s = 2m_0E + m_0^2 + m_0^2$$

if E much larger than m_0 we find:

$$E_{cm} = \sqrt{s} = \sqrt{2m_0E}$$



Particle collisions - fixed target

Homework: try for $E1 \neq E2$ and $m1 \neq m2$

Examples:

collision	beam energy	\sqrt{s} (collider)	\sqrt{s} (fixed target)
pp	315 (GeV)	630 (GeV)	24.3 (GeV)
pp	7000 (GeV)	14000 (GeV)	114.6 (GeV)
e+e-	100 (GeV)	200 (GeV)	0.320 (GeV)
TLEP	175 (GeV)	350 (GeV)	0.423 (GeV) !



Kinematic invariant

We need to make cross sections (and therefore luminosity) invariant !

This is done by a calibration factor which is (without derivation):


$$K = \sqrt{(\vec{v}_1 - \vec{v}_2)^2 - (\vec{v}_1 \times \vec{v}_2)^2 / c^2}$$

Here \vec{v}_1 and \vec{v}_2 are the velocities of the two (relativistic) beams.

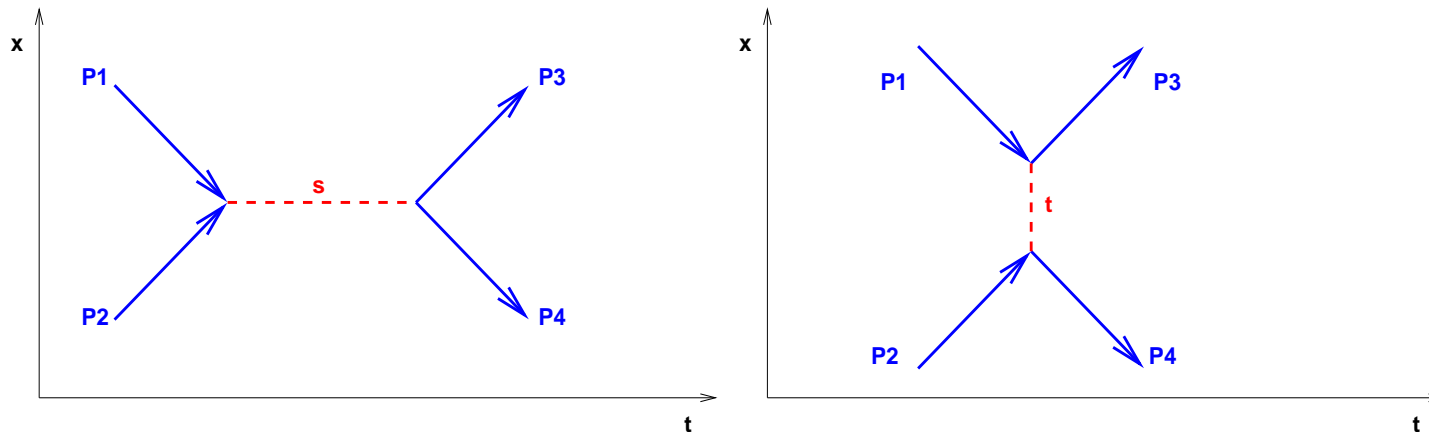
For a (symmetric) collider, e.g. LHC, we have:

$$\vec{v}_1 = -\vec{v}_2, \quad \vec{v}_1 \times \vec{v}_2 = 0 \quad \text{head-on!}$$

→ $K = 2 \cdot c !$



For completeness ...



Squared centre of mass energy:

$$s = (P1 + P2)^2 = (P3 + P4)^2$$

Squared momentum transfer in particle scattering
(small t - small angle, see again lecture on Luminosity):

$$t = (P1 - P3)^2 = (P2 - P4)^2$$

Kinematic relations

We have already seen a few, e.g.:

➤ $T = E - E_0 = (\gamma - 1)E_0$

➤ $E = \gamma \cdot E_0$

➤ $E_0 = \sqrt{E^2 - c^2 p^2}$

➤ etc. ...

Very useful for everyday calculations →



Kinematic relations

	cp	T	E	γ
$\beta =$	$\frac{1}{\sqrt{(\frac{E_0}{cp})^2 + 1}}$	$\sqrt{1 - \frac{1}{(1 + \frac{T}{E_0})^2}}$	$\sqrt{1 - (\frac{E_0}{E})^2}$	$\sqrt{1 - \gamma^{-2}}$
cp =	cp	$\sqrt{T(2E_0 + T)}$	$\sqrt{E^2 - E_0^2}$	$E_0 \sqrt{\gamma^2 - 1}$
$E_0 =$	$\frac{cp}{\sqrt{\gamma^2 - 1}}$	$T/(\gamma - 1)$	$\sqrt{E^2 - c^2 p^2}$	E/γ
T =	$cp \sqrt{\frac{\gamma - 1}{\gamma + 1}}$	T	$E - E_0$	$E_0(\gamma - 1)$
$\gamma =$	$cp/E_0\beta$	$1 + T/E_0$	E/E_0	γ



Kinematic relations

➤ Example: CERN Booster

At injection: $T = 50 \text{ MeV}$


➔ $E = 0.988 \text{ GeV}$, $p = 0.311 \text{ GeV}/c$

➔ $\gamma = 1.0533$, $\beta = 0.314$

At extraction: $T = 1.4 \text{ GeV}$

➔ $E = 2.338 \text{ GeV}$, $p = 2.141 \text{ GeV}/c$

➔ $\gamma = 2.4925$, $\beta = 0.916$



Kinematic relations - logarithmic derivatives

	$\frac{d\beta}{\beta}$	$\frac{dp}{p}$	$\frac{dT}{T}$	$\frac{dE}{E} = \frac{d\gamma}{\gamma}$
$\frac{d\beta}{\beta} =$	$\frac{d\beta}{\beta}$	$\frac{1}{\gamma^2} \frac{dp}{p}$	$\frac{1}{\gamma(\gamma+1)} \frac{dT}{T}$	$\frac{1}{(\beta\gamma)^2} \frac{d\gamma}{\gamma}$
$\frac{dp}{p} =$	$\gamma^2 \frac{d\beta}{\beta}$	$\frac{dp}{p}$	$[\gamma/(\gamma + 1)] \frac{dT}{T}$	$\frac{1}{\beta^2} \frac{d\gamma}{\gamma}$
$\frac{dT}{T} =$	$\gamma(\gamma + 1) \frac{d\beta}{\beta}$	$(1 + \frac{1}{\gamma}) \frac{dp}{p}$	$\frac{dT}{T}$	$\frac{\gamma}{(\gamma-1)} \frac{d\gamma}{\gamma}$
$\frac{dE}{E} =$	$(\beta\gamma)^2 \frac{d\beta}{\beta}$	$\beta^2 \frac{dp}{p}$	$(1 - \frac{1}{\gamma}) \frac{dT}{T}$	$\frac{d\gamma}{\gamma}$
$\frac{d\gamma}{\gamma} =$	$(\gamma^2 - 1) \frac{d\beta}{\beta}$	$\frac{dp}{p} - \frac{d\beta}{\beta}$	$(1 - \frac{1}{\gamma}) \frac{dT}{T}$	$\frac{d\gamma}{\gamma}$

Example LHC (7 TeV): $\frac{\Delta p}{p} \approx 10^{-4} \rightarrow \frac{\Delta\beta}{\beta} = \frac{\Delta v}{v} \approx 2 \cdot 10^{-12}$



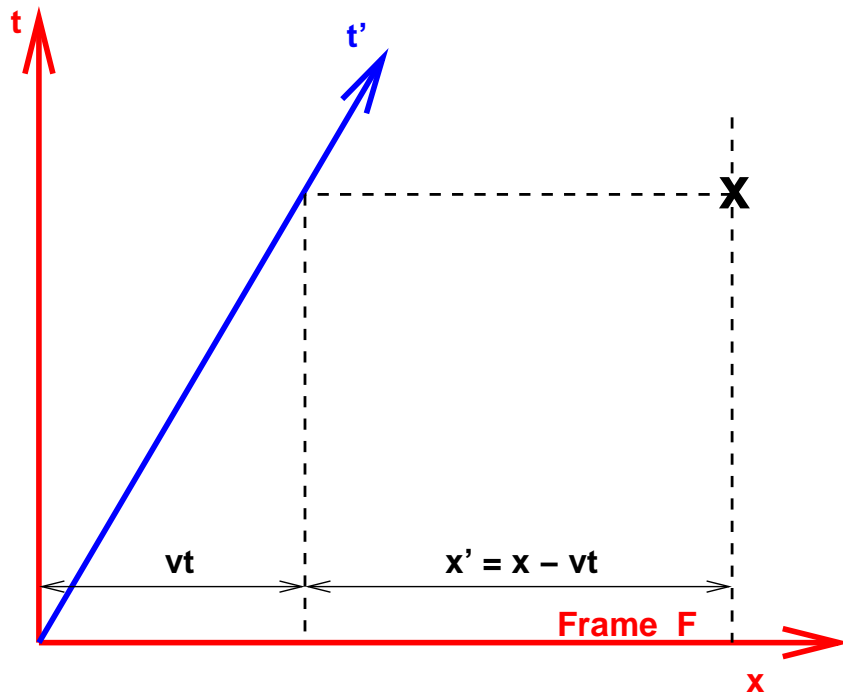
Summary

- Special Relativity is very simple, derived from basic principles
- Relativistic effects vital in accelerators:
 - Lorentz contraction and Time dilation
 - Invariants !
 - Relativistic mass effects
 - Modification of electromagnetic field
- Find back in later lectures ...



- BACKUP SLIDES -

Galilei transformation - schematic



➤ Rest frame and Galilei transformation ...



Forces and fields

Motion of charged particles in electromagnetic fields \vec{E} , \vec{B} determined by Lorentz force

$$\vec{f} = \frac{d}{dt}(m_0\gamma\vec{v}) = q \cdot (\vec{E} + \vec{v} \times \vec{B})$$

or as four-vector:

$$F = \frac{dP}{d\tau} = \gamma \left(\frac{\vec{v} \cdot \vec{f}}{c}, \vec{f} \right) = \gamma \left(\frac{1}{c} \frac{dE}{dt}, \frac{d\vec{p}}{dt} \right)$$



Field tensor

Electromagnetic field described by field-tensor $F^{\mu\nu}$:

$$F^{\mu\nu} = \begin{pmatrix} 0 & \frac{-E_x}{c} & \frac{-E_y}{c} & \frac{-E_z}{c} \\ \frac{E_x}{c} & 0 & -B_z & B_y \\ \frac{E_y}{c} & B_z & 0 & -B_x \\ \frac{E_z}{c} & -B_y & B_x & 0 \end{pmatrix}$$

derived from four-vector $A_\mu = (\Phi, \vec{A})$ like:

$$F^{\mu\nu} = \delta^\mu A^\nu - \delta^\nu A^\mu$$



Lorentz transformation of fields

$$\vec{E}'_{\perp} = \gamma(\vec{E}_{\perp} + \vec{v} \times \vec{B})$$

$$\vec{B}'_{\perp} = \gamma\left(\vec{B}_{\perp} - \frac{\vec{v} \times \vec{E}}{c^2}\right)$$

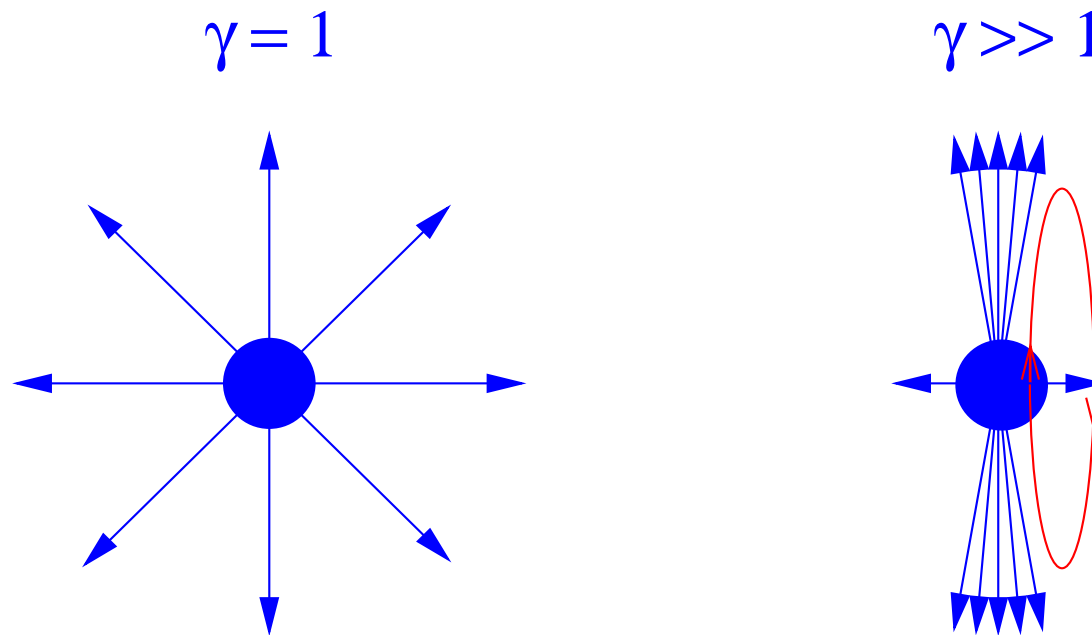
$$\vec{E}'_{\parallel} = \vec{E}_{\parallel}$$

$$\vec{B}'_{\parallel} = \vec{B}_{\parallel}$$

➤ Field perpendicular to movement transform



Lorentz transformation of fields



- In rest frame purely electrostatic forces
- In moving frame \vec{E} transformed and \vec{B} appears