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Gravitational Quantum States of Antihydrogen

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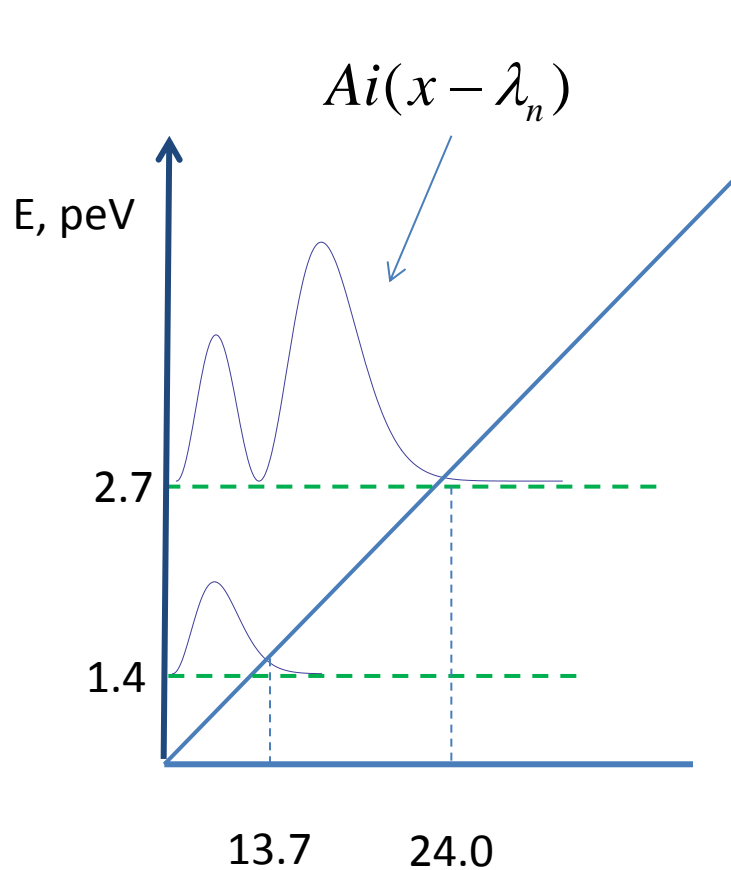


Plan of the talk

- Gravitational states of antihydrogen: Is it possible?
- How can we get gravitational mass out of gravitational states?
- Properties of gravitational states
- Spectroscopy, interference and time-spatial resolution of gravitational states

Gravitational quantum states?

State of motion of a quantum particle, which is localized near reflecting surface in a gravitational field of the Earth.



$$\left[-\frac{d^2}{dx^2} + x - \lambda \right] F(x) = 0, \quad F(0) = 0$$

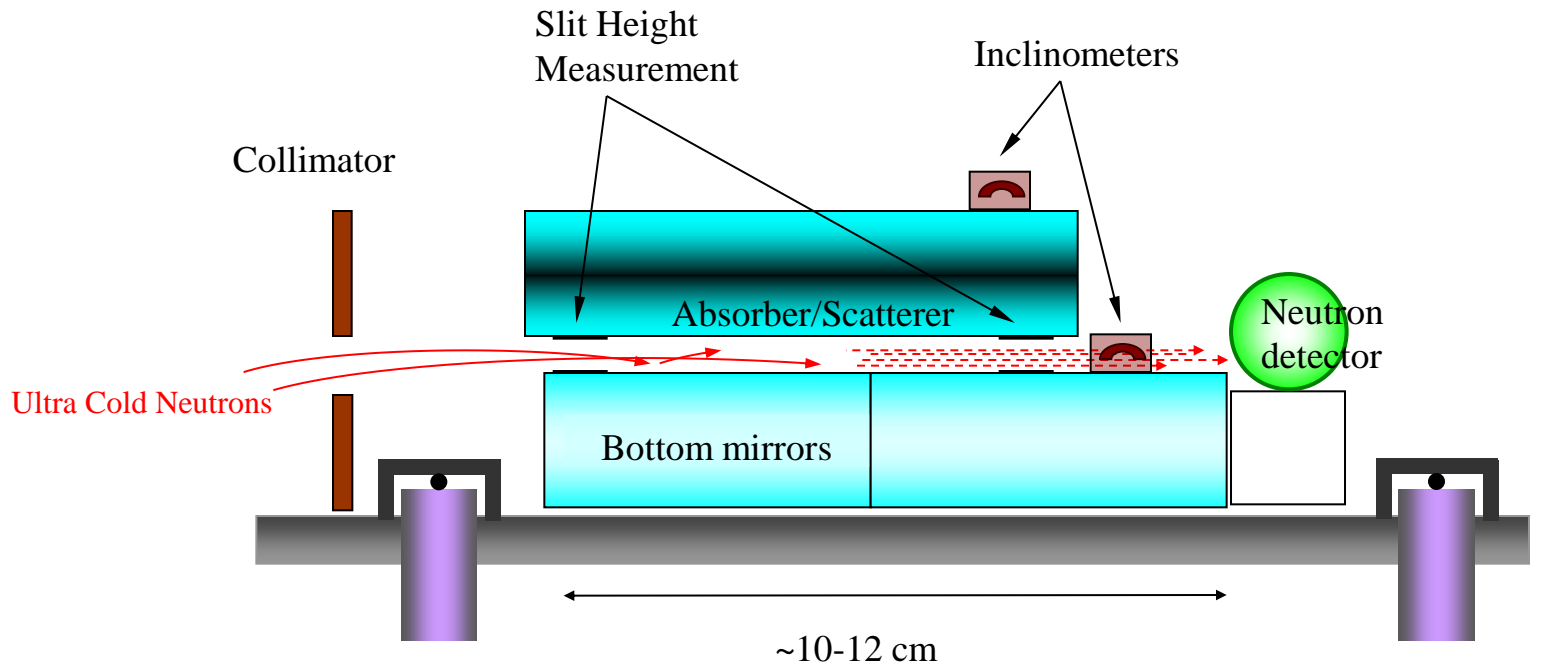
TABLE I. The eigenvalues, gravitational energies, and classical turning points of a quantum bouncer with the mass of (anti)hydrogen in the Earth's gravitational field.

n	λ_n^0	E_n^0 (peV)	z_n^0 (μm)
1	2.338	1.407	13.726
2	4.088	2.461	24.001
3	5.521	3.324	32.414
4	6.787	4.086	39.846
5	7.944	4.782	46.639
6	9.023	5.431	52.974
7	10.040	6.044	58.945

$Z, \mu\text{m}$

First Observation: Gravitational States of Neutrons

Nesvizhevsky et al. Nature 415, 297 (2002)

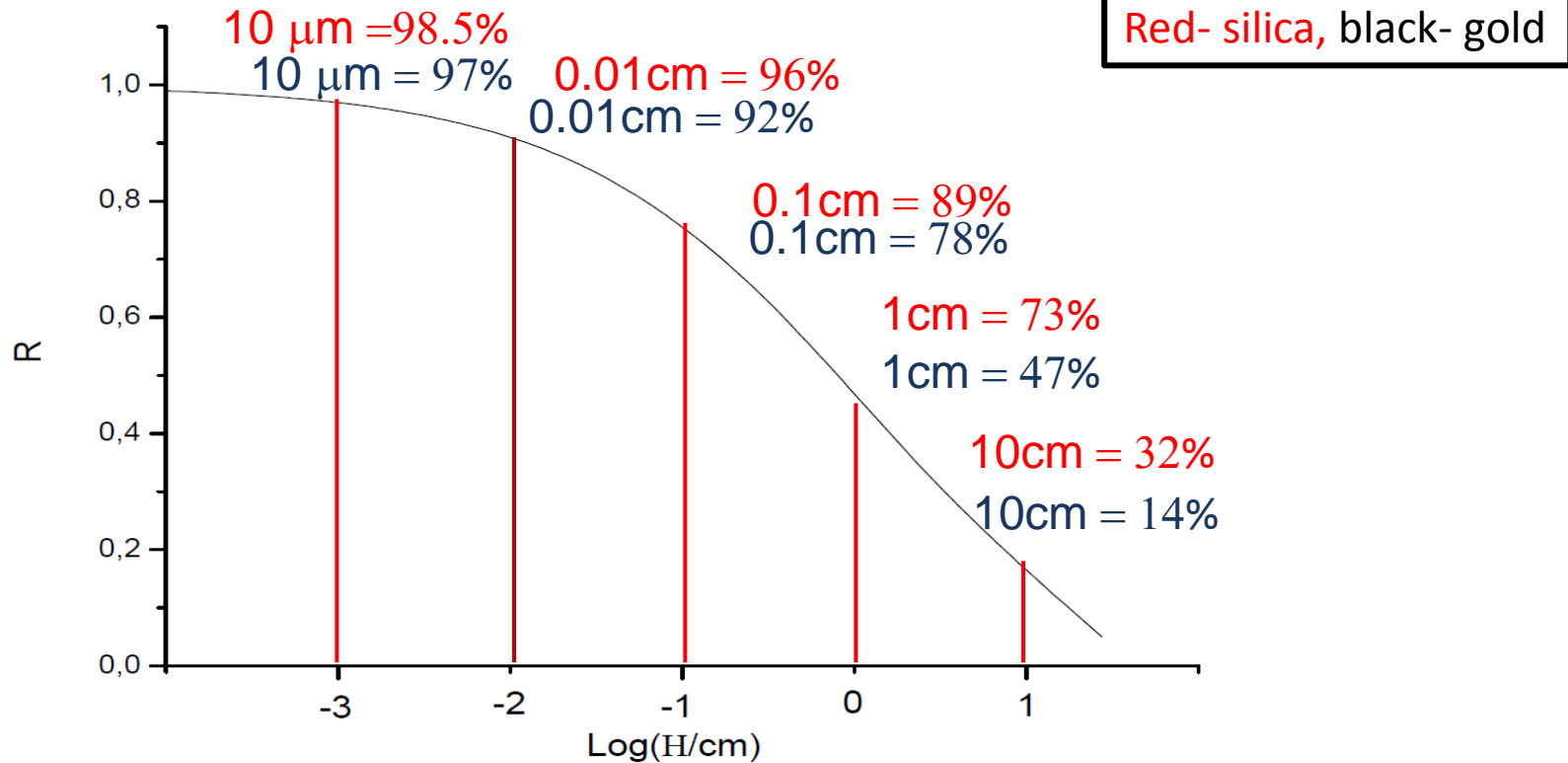


- Count rates at ILL turbine: $\sim 1/s$ to $1/h$
- Effective (vertical) temperature of neutrons is ~ 20 nK
- Background suppression is a factor of $\sim 10^8$ - 10^9
- Parallelism of the bottom mirror and the absorber/scatterer is $\sim 10^{-6}$



GRANIT SPECTROMETER

Gravitational states of Antihydrogen: Seems Impossible? Quantum Reflection!

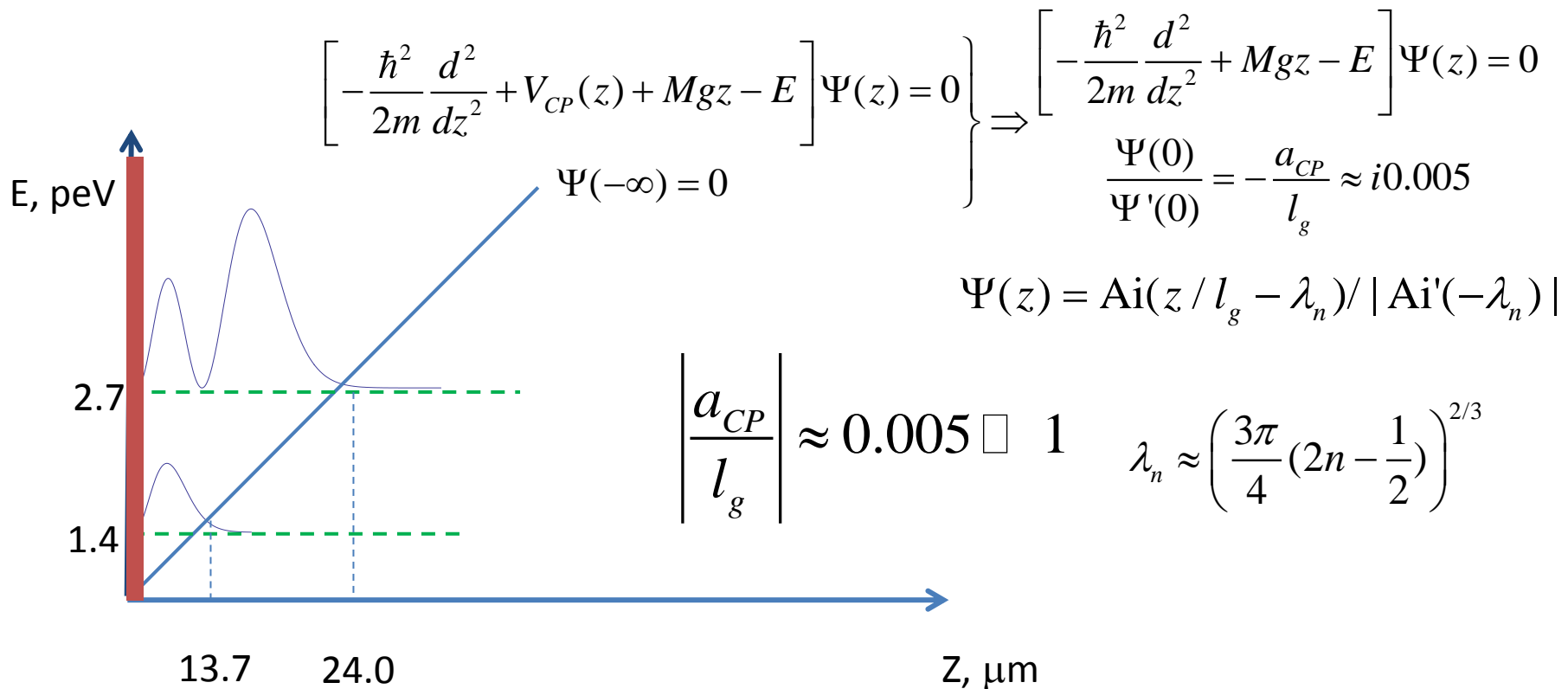


A. Yu. Voronin, P. Froelich, and B. Zygelman, *Phys. Rev. A* **72**, 062903 (2005).

[G. Dufour](#), [A. Gérardin](#), [R. Guérout](#), [A. Lambrecht](#), [V. V. Nesvizhevsky](#), [S. Reynaud](#), [A. Yu. Voronin](#) *Phys. Rev. A* **87**, 012901 (2013)

Gravitational states of antihydrogen

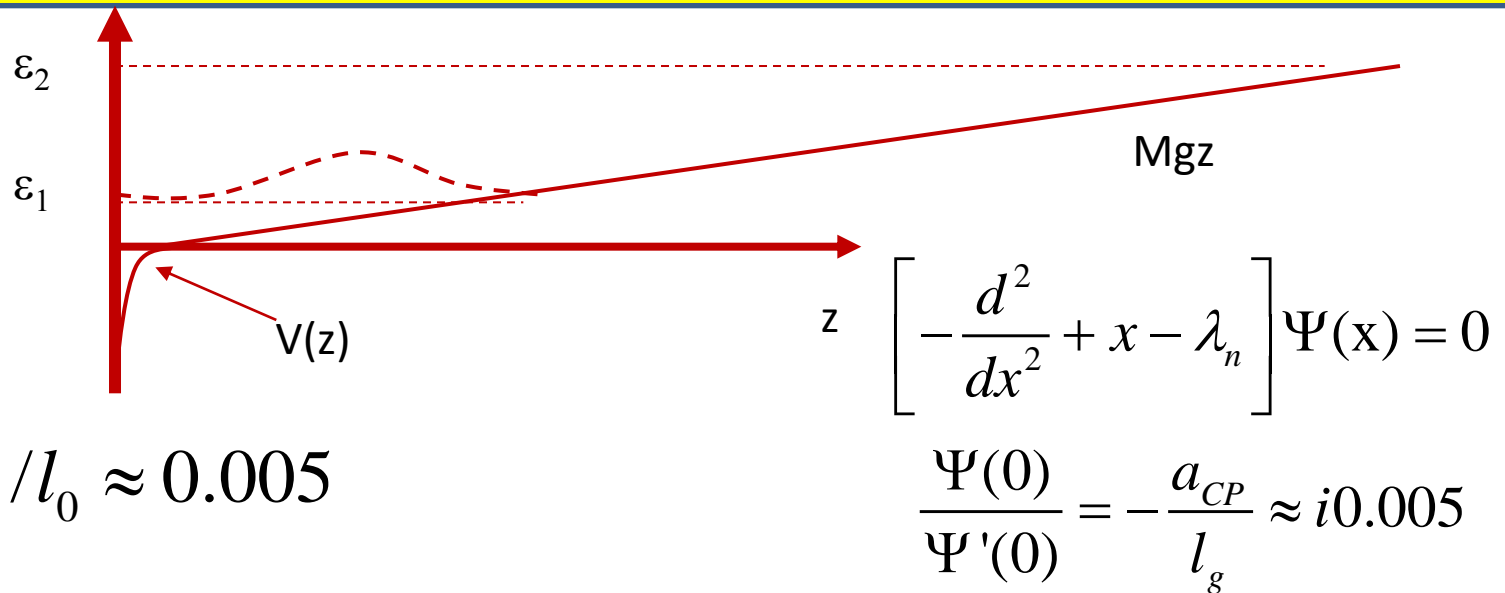
Quantum reflection is about 97% - it works like a reflecting wall



Effects of surface

- Hierarchy of scales $l_g \ll |a_{CP}|$: gravity and surface-atom interaction are factorized
- Annihilation in the bulk of the wall: short – range atom-wall interactions are washed out
- Small annihilation width of gravitational states: compromise between long life-time and observation

Correction by Casimir-Polder potential + annihilation



$$|a|/l_0 \approx 0.005$$

Correction by Casimir-Polder and annihilation:

$$\tilde{\lambda}_n = \lambda_n + a/l_g$$

$$\varepsilon_n = \varepsilon_0(\lambda_n + \text{Re } a/l_g) \quad \Gamma = 2\varepsilon_0 |\text{Im } a|/l_g$$

$$\tau = \frac{l_g}{\varepsilon_g} \frac{\hbar}{2|\text{Im } a|} = \frac{\hbar}{2Mg|\text{Im } a|} \approx \underline{0.1s}$$

All states have equal shift and lifetime $\sim 0.1s \Rightarrow$

No surface effects in transition frequencies

Gravitational states and Gravitational mass

$$\text{Classical: } m\ddot{z} = M/g \rightarrow \ddot{z} = g \rightarrow T = \sqrt{2H/g}$$

$$\text{Quantum: } \left[-\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + Mgz - E \right] \Psi(z) = 0 \Rightarrow \left[-\frac{d^2}{dx^2} + x - \lambda_n \right] F(x) = 0$$

$$\varepsilon_g = \sqrt[3]{\frac{\hbar^2 M^2 g^2}{2m}} = 0.61 \cdot 10^{-12} \text{ eV}; \quad l_g = \sqrt[3]{\frac{\hbar^2}{2Mmg}} = 5.87 \cdot 10^{-6} \text{ m}$$

$$E_n = \varepsilon_g \lambda_n \quad z = l_g x \quad \Psi_n(z) = \frac{1}{\sqrt{l_g}} F(z/l_g - E_n/\varepsilon_g)$$

$$m = \frac{\hbar^2}{2\varepsilon_g l_g^2}; \quad M = \frac{\varepsilon_g}{g l_g}$$

$$M = m \Rightarrow \frac{\hbar}{\varepsilon_g} = \sqrt{\frac{2l_g}{g}} \quad \text{or} \quad T = \sqrt{\frac{2H}{g}}$$

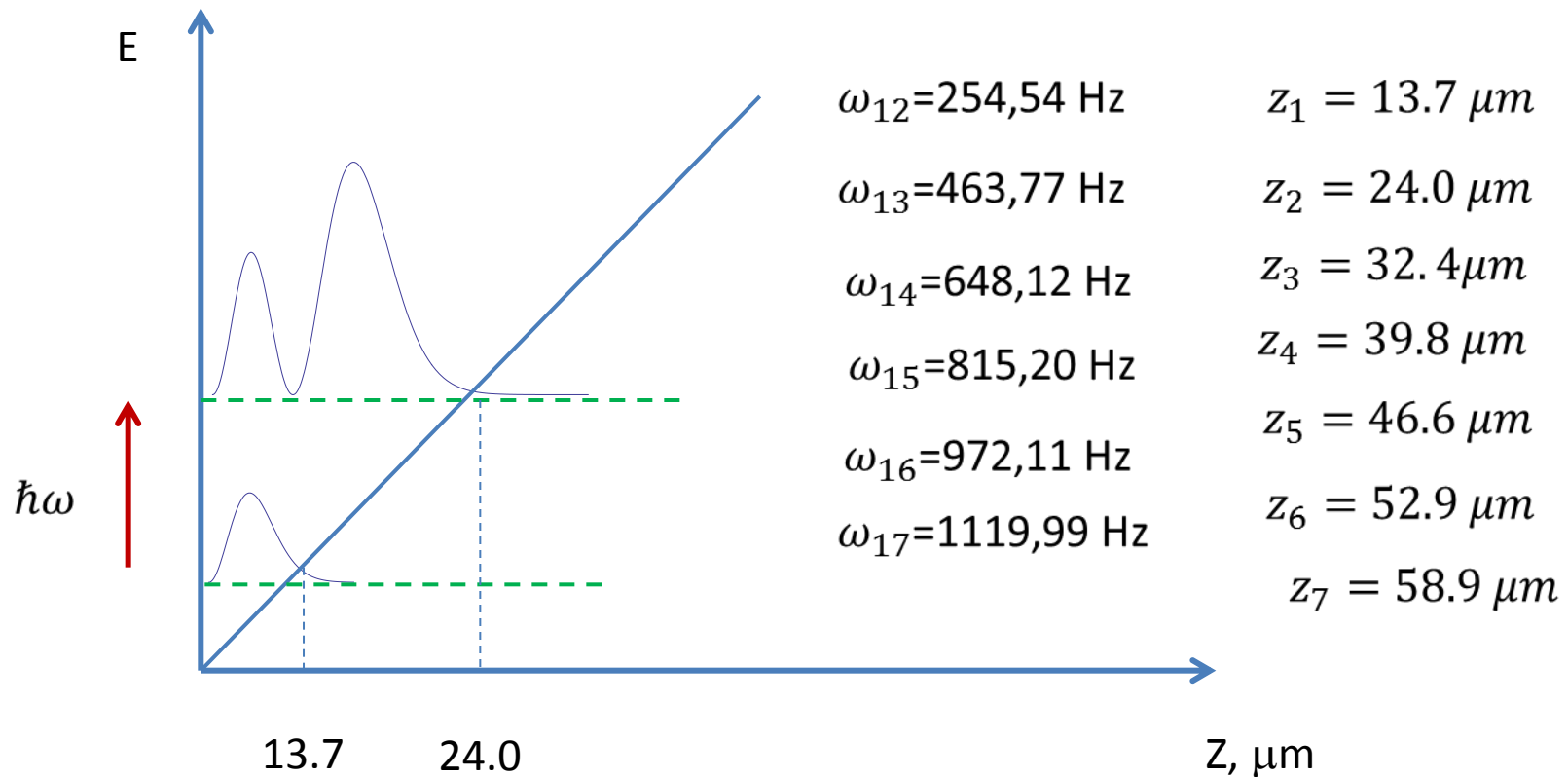
Gravitational states are all about energy and spatial scales

Methods of observation

- Spectroscopy: induced transition between gravitational states
- Interference: temporal and spatial oscillations of annihilation signal of superposition of gravitational states
- Time and spatial resolution of free-fall events: mapping of momentum distribution of gravitational state into time-of-fall or spatial distribution

Spectroscopy- to induce transitions between gravitational states with alternating magnetic field

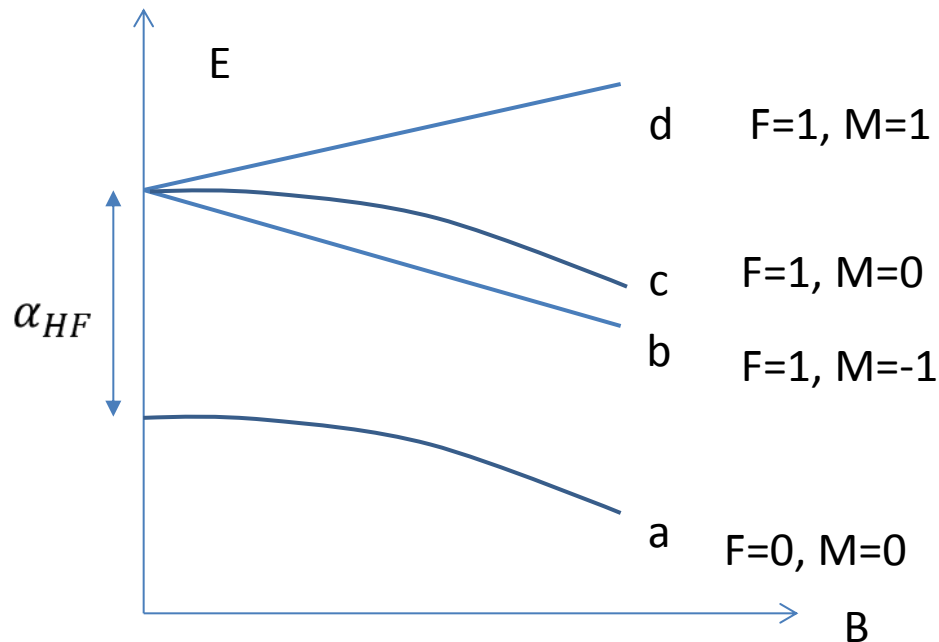
Developed for neutrons by V. Nesvizhevsky, S. Baessler, G. Pignol, K. Protassov, A.Voronin



Antihydrogen in Magnetic Field

$$\hat{T}_z + Mgz + \hat{H}_c + \frac{\alpha_{HF}}{2} (\mathbf{F}^2 - \frac{3}{2}) + \mathbf{B}(z, t) (\mu_B \hat{\mathbf{s}}_e + \mu_p \hat{\mathbf{s}}_p) + \hat{H}_{st}$$

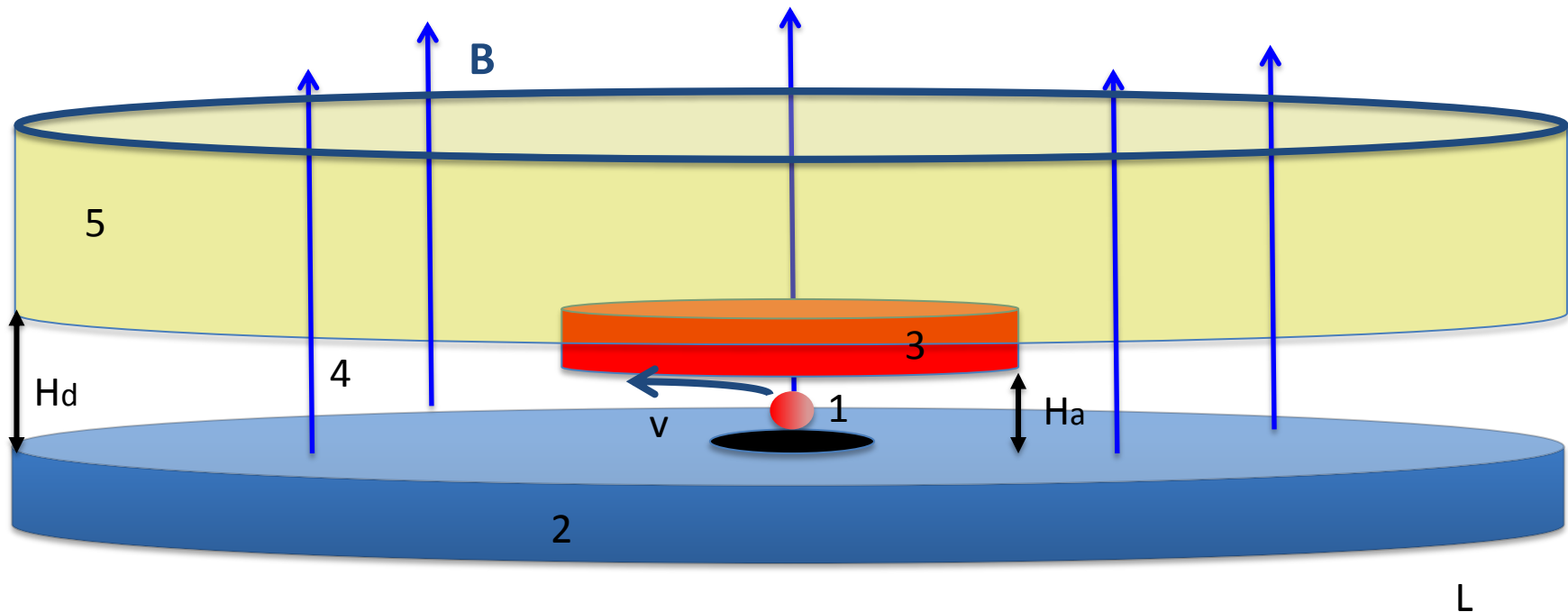
\hat{T}_z → Center-of-mass fall
 Mgz → Coulomb
 \hat{H}_c → HF splitting
 $\mathbf{B}(z, t) (\mu_B \hat{\mathbf{s}}_e + \mu_p \hat{\mathbf{s}}_p)$ → Field-magnetic moment interaction
 \hat{H}_{st} → Effective Stark-effect



$$\frac{\alpha_{HF}}{\hbar} \cong 1420 \text{ Mhz}$$

$$\epsilon_{grav} \ll \mu_B \ll \alpha_{HF}$$

Possible scheme of flow-throw experiment

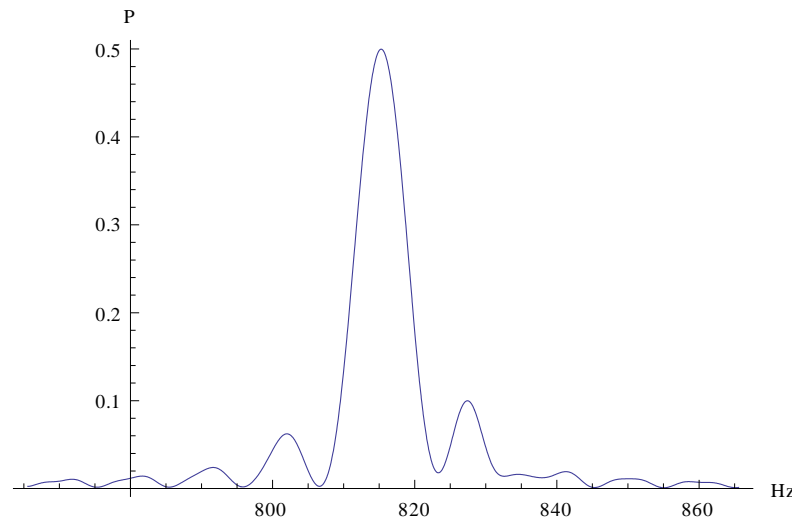


1-source of ultracold antihydrogen, 2-mirror, 3- absorber, 4- magnetic field, 5- detector

$$v \approx 1m / s, H_a = 15 \mu m, H_d = 25 \mu m, B_0 \approx 10Gs, \beta \approx 10Gs / m, L = 30cm$$

Transition probability

$$P_{ik} = \frac{1}{2} \frac{\Omega_{ik}^2}{\Omega_{ik}^2 + \hbar^2 \Delta^2} \sin^2\left(\frac{t}{\hbar} \sqrt{\Omega_{ik}^2 + \hbar^2 \Delta^2}\right) \exp(-\Gamma t)$$
$$\Omega_{ik} = \frac{(\mu_B + \mu_{\bar{p}}) \beta l_g^3}{\hbar (z_k - z_i)^2}$$



Transition probability as a function of frequency. Transition 1- \rightarrow 5

Time of observation $t=0.1$ s

EP and gravitational mass

$$\varepsilon_g = \sqrt[3]{\frac{\hbar^2 M^2 g^2}{2m}} = 0.61 \cdot 10^{-12} \text{ eV}$$

$$\hbar\omega_{ik} = \sqrt[3]{\frac{\hbar^2 M^2 g^2}{2m}} (\lambda_k - \lambda_i) \Rightarrow M = \sqrt{\frac{2m\hbar\omega_{ik}^3}{g^2 (\lambda_k - \lambda_i)^3}}$$

$$M = m \Rightarrow M = \frac{2\omega_{ik}^3}{(\lambda_k - \lambda_i)^3} \frac{\hbar}{g^2}$$

PRECISION

$$\varepsilon \approx \frac{3\gamma}{2\sqrt{N_{\bar{H}}}\omega}; N_{\bar{H}} = 10^2$$

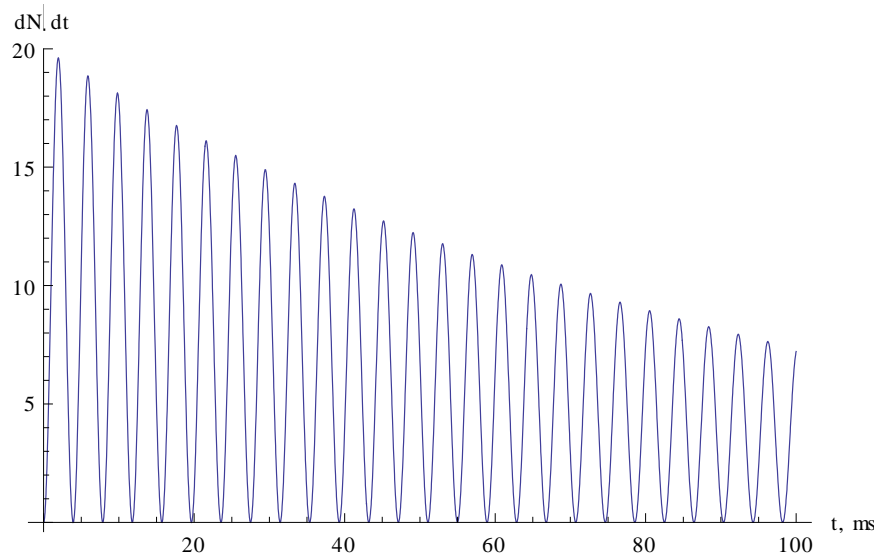
$$\varepsilon_{1-5} \approx 10^{-3}$$

Interference of gravitational states

$$\Psi = \sum_{i=1}^N C_i \Psi_i$$

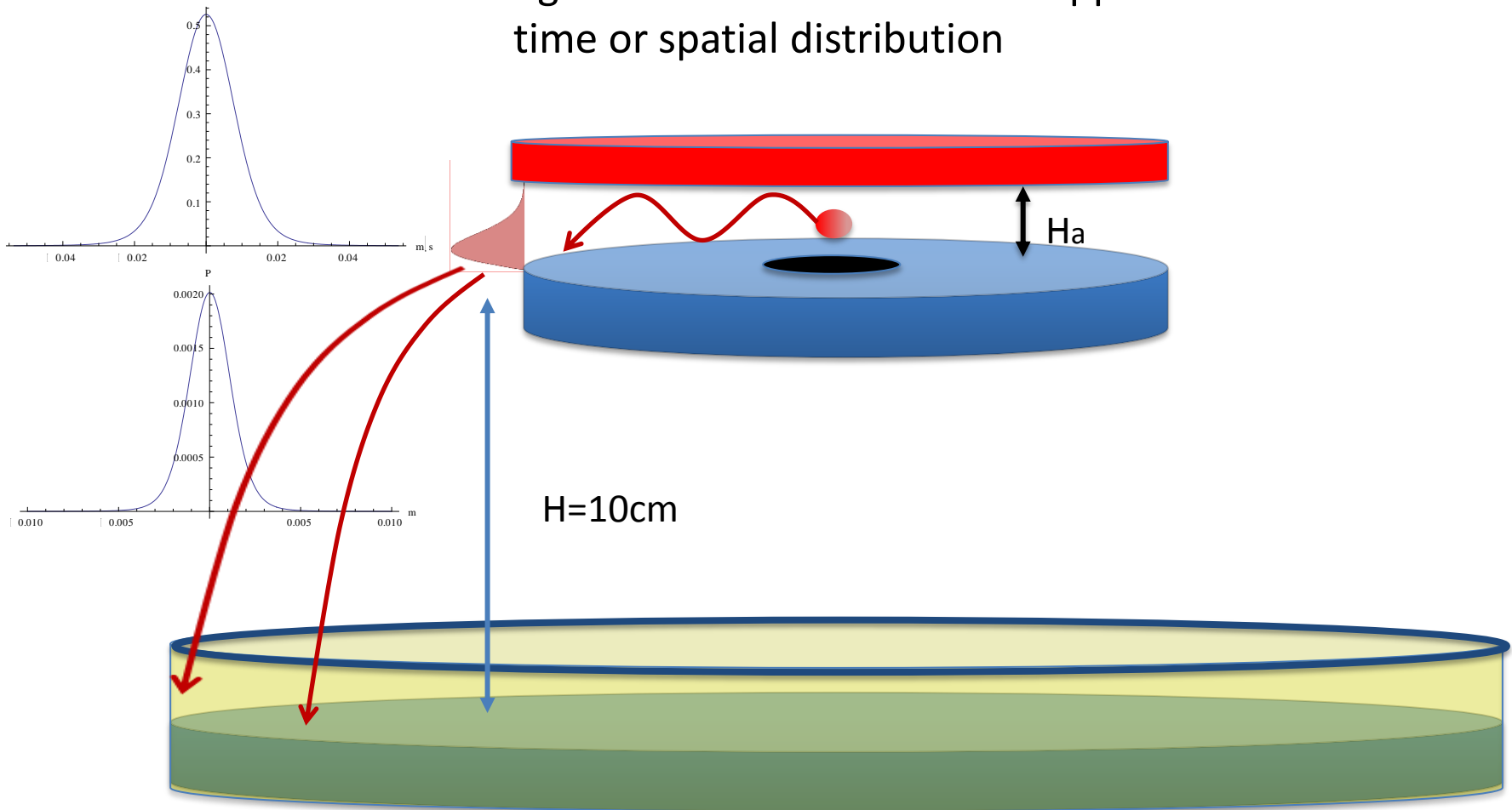
$$j(z,t) = \frac{i\hbar}{2m} \left[\Psi(z,t) \frac{d\Psi^*(z,t)}{dz} - \Psi^*(z,t) \frac{d\Psi(z,t)}{dz} \right]$$

$$\frac{dF_{12}}{dz} = -\frac{\Gamma}{\hbar} \exp\left(-\frac{\Gamma}{\hbar} t\right) (1 - \cos(\omega_{12}t)) \quad \omega_{12} = \frac{E_2 - E_1}{\hbar} \quad \hbar\Gamma^{-1} \approx 0.1s$$



Time and Spatial Resolving of Gravitational States

Momentum distribution of gravitational state can be mapped into measurable time or spatial distribution



Mapping of momentum distribution

$$\Psi(z, t) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{ipz/\hbar} G(p, t, p') F_0(p') dp dp'$$

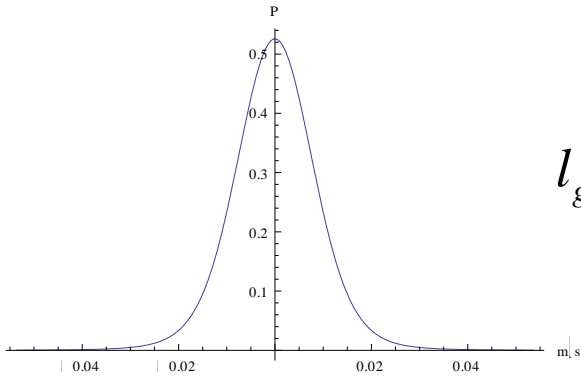
$$G(p, t, p') = \exp\left[-\frac{it}{2m\hbar} (p^2 - Mgpt + M^2 g^2 t^2 / 3)\right] \delta(p - Mgt - p')$$

$$\Psi(z, t) \approx \sqrt{\frac{m}{t}} e^{\frac{imz^2}{2\hbar} + \frac{it^3 M^2 g^2}{2m\hbar}} F_0(p_0 - Mgt); \quad p_0 = \left(z + \frac{gt^2}{2}\right) \frac{m}{t}$$

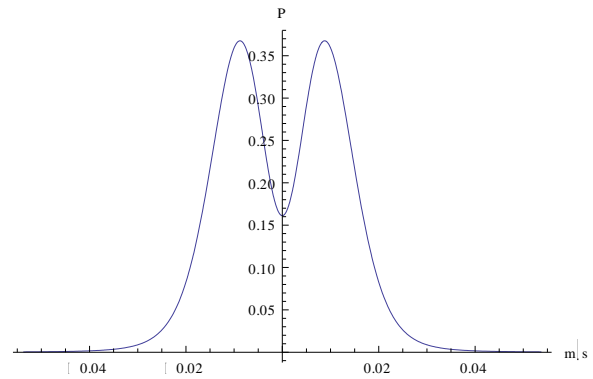
$$|\Psi(z, t)|^2 \approx \frac{m}{t} |F_0(k)|^2$$

$$1) z = z_0 : k = mg(t - t_0), \quad t_0 = \sqrt{2g / z_0}$$

$$2) t = t_0 = L / v : k = \frac{m(z - z_0)}{t_0}$$

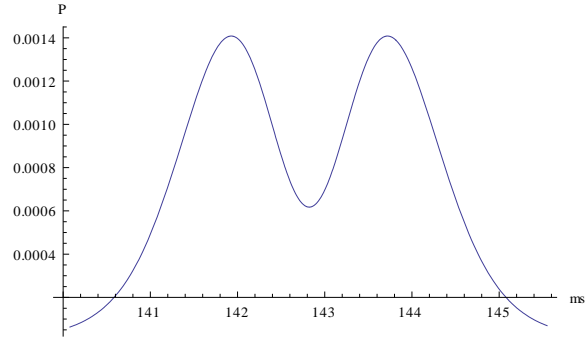
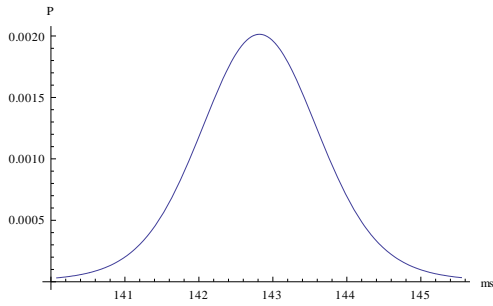


$$l_g = \frac{\delta(n)\hbar}{\Delta k_n}$$

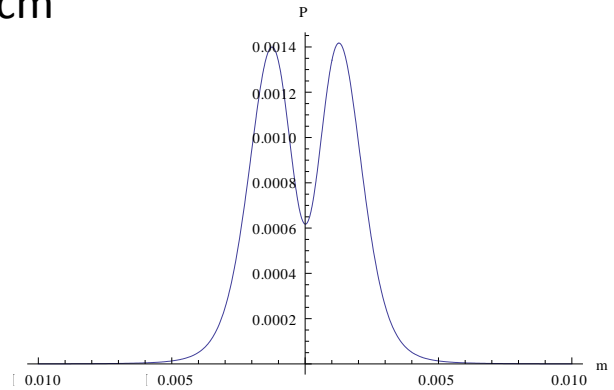
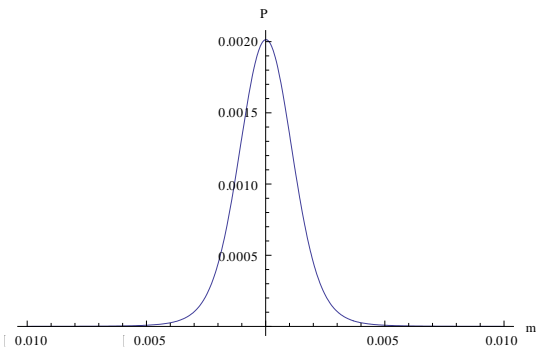


1 state velocity distribution

2 state



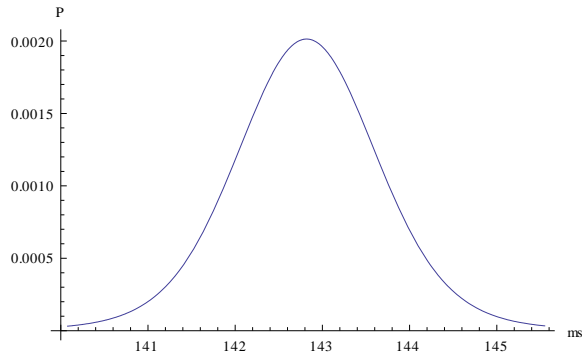
Time-of-fall distribution H= 10 cm



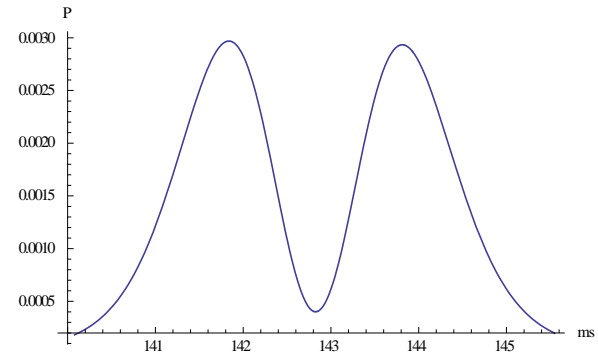
Spatial distribution time-of-flight T=0.1s

Phase monitoring

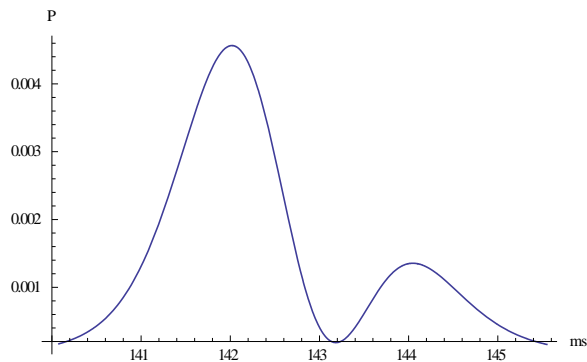
$$\Psi = \Psi_1 + \text{Exp}(-i\omega_{21}t)\Psi_2$$



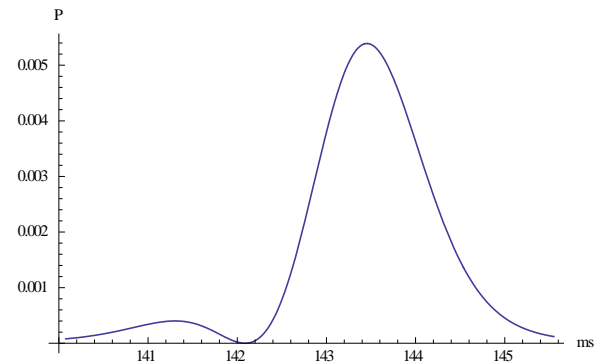
$t = 0$



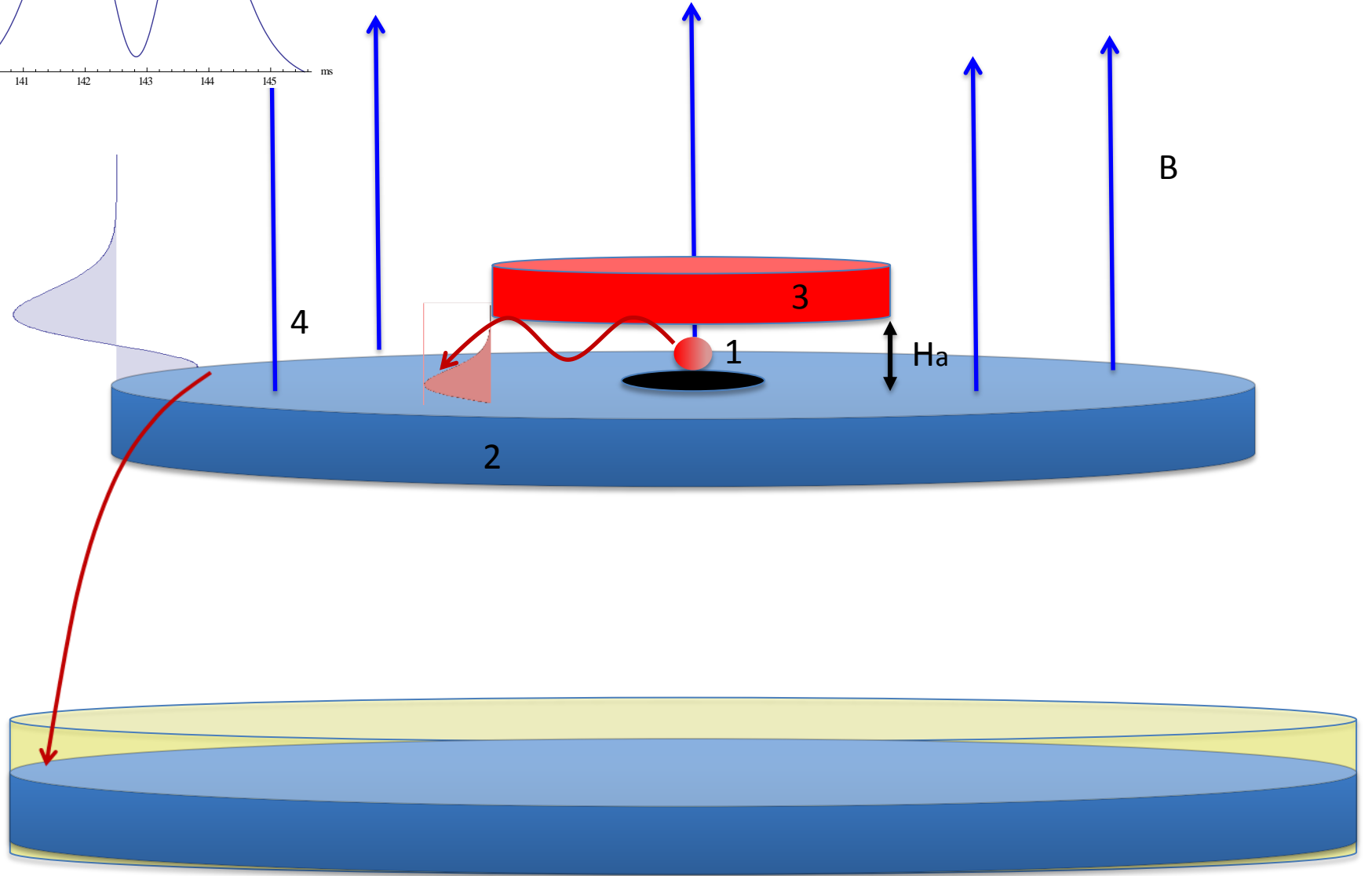
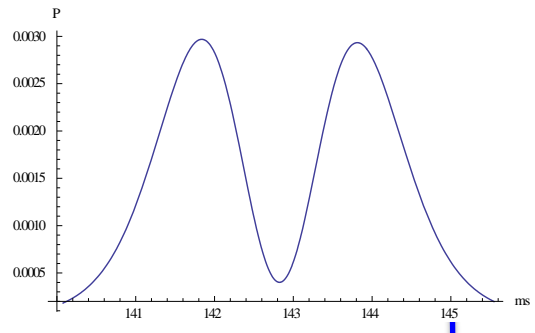
$t = 0.0019s$



$t = 0.0014s$



$t = 0.0029s$



Conclusions

- Gravitational states of Antihydrogen: simplest bound antimatter quantum system, determined by gravity. Effects of surface are canceled out.
- Gravitational states of Antihydrogen- metastable and long-living, easy to study due to annihilation signal
- Gravitational states- a way to precision measurement of the gravitational mass M