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## LARGE volume scenario in 5D SUGRA

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based on arXiv:1307.5585
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## Introduction

## LARGE volume scenario (LVS)

 in string theoryModuli stabilization: exponentially large extra dimension

> SUSY breaking: The scale is much smaller than the Planck scale

Without fine-tuned small parameters !
But the sring theoretical effects and the non-trivial geometry are required...
V. Balasubramanian, P. Berglund, J. P. Conlon, and F. Quevedo (2005)
J. P. Conlon, F. Quevedo, and K. Suruliz (2005)

## Introduction

Q. Can we realize the LVS with the simpler theory?
A. Yes,

I'Il show that the LVS can be realized
in the 5D supergravity on $S^{1} / Z_{2}$
without the stringy effects.

## Multiple moduli in 5D SUGRA

## General 5D SUGRA on $S^{1} / Z_{2} \longrightarrow \quad$ General 4D effective theory <br> T. Kugo and K. Ohashi (2001) <br> reduction

5 D vector multiplet $\longrightarrow 4 \mathrm{D}$ vector multiplet

Well known set-up : one modulus
= radion

## Multiple moduli in 5D SUGRA

## General 5D SUGRA on $S^{1} / Z_{2} \longrightarrow \quad$ General 4D effective theory <br> T. Kugo and K. Ohashi (2001)

General set-up : multiple moduli
= radion + non-geometric moduli
H. Abe, H. Otsuka, Y. Sakamura and Y.Y (2011)

## The norm function and Kahler potential in 5D SUGRA

The size of the extra dim. $\Rightarrow L_{\text {phys }}=\left\langle\mathcal{N}^{\frac{1}{3}}\right\rangle$

$$
\mathcal{N}=C_{I J K} \operatorname{Re} T^{I} \operatorname{Re} T^{J} \operatorname{Re} T^{K}
$$

## The norm function and Kahler potential in 5D SUGRA

The size of the extra dim. $\rightarrow L_{\text {phys }}=\left\langle\mathcal{N}^{\frac{1}{3}}\right\rangle$

$$
\mathcal{N}=C_{I J K} \operatorname{Re} T^{I} \operatorname{Re} T^{J} \operatorname{Re} T^{K}
$$

Kähler potential: $K=-\log \mathcal{N}$

$$
K_{I} K^{I \bar{J}} K_{\bar{J}}=3
$$

## The norm function and Kahler potential in 5D SUGRA

The size of the extra dim. $\longrightarrow L_{\mathrm{phys}}=\left\langle\mathcal{N}^{\frac{1}{3}}\right\rangle$

$$
\mathcal{N}=C_{I J K} \operatorname{Re} T^{I} \operatorname{Re} T^{J} \operatorname{Re} T^{K}
$$

1-loop corrected Kähler potential: $K=-\log (\mathcal{N}+\xi)$


## LVS in 5D SUGRA

$$
K=-\log (\mathcal{N}+\xi) \quad W=W_{0}+A e^{-a T_{s}}
$$

where

$$
\mathcal{N}=\left(\operatorname{Re} T_{b}\right)^{3}-C_{s}\left(\operatorname{Re} T_{s}\right)^{3}
$$

$$
\begin{aligned}
a & =\mathcal{O}\left(4 \pi^{2}\right) \\
W_{0} & =\mathcal{O}\left(M_{p l}^{3}\right)
\end{aligned}
$$

$$
\begin{aligned}
& T_{b}=\tau_{b}+i \rho \quad \sim \text { radion } \\
& T_{s}=\tau_{s}+i \sigma \quad \sim \text { non-geometric modulus }
\end{aligned}
$$

## Moduli stabilization

$$
\begin{aligned}
V & \sim \frac{1}{\mathcal{N}}\left(\frac{2 \mathcal{N}}{3 C_{s} \tau_{s}}(a A)^{2} e^{-2 a \tau_{s}}+4 a \tau_{s} W_{0} A e^{-a \tau_{s}} \cos (a \sigma)\right)+\frac{6 \xi W_{0}^{2}}{\mathcal{N}^{2}} \\
& =\frac{2(a A)^{2}}{3 C_{s} \tau_{s}} e^{-2 a \tau_{s}}+4 a \tau_{s} W_{0} A \cos (a \sigma) \frac{e^{-a \tau_{s}}}{\mathcal{N}}+6 \xi W_{0}^{2} \frac{1}{\mathcal{N}^{2}}
\end{aligned}
$$

Non-perturbative term vs Volume suppressed term

$$
\begin{gathered}
\langle\mathcal{N}\rangle \sim \frac{3 \xi W_{0} e^{a\left\langle\tau_{s}\right\rangle}}{a\left\langle\tau_{s}\right\rangle A} \quad\left\langle\tau_{s}\right\rangle \sim\left(\frac{\xi}{C_{s}}\right)^{\frac{1}{3}} \\
L_{\text {phys }}=\left\langle\mathcal{N}^{\frac{1}{3}}\right\rangle \gg 1 \quad \text { (In Planck unit) }
\end{gathered}
$$

Exponentially large extra dimension!

## SUSY breaking in 5D LVS

Gravitino mass: $m_{3 / 2}=\frac{W_{0}}{\sqrt{\mathcal{N}}} \sim \mathcal{O}\left(M_{p l} / \sqrt{\mathcal{N}}\right) \ll M_{p l}$

$$
\text { F-terms: } \begin{aligned}
& \frac{F^{T_{b}}}{T_{b}+\bar{T}_{b}} \\
\sim & \sim \frac{W_{0}}{\sqrt{\mathcal{N}}}=m_{3 / 2} \\
& \frac{F^{T_{s}}}{T_{s}+\bar{T}_{s}}
\end{aligned} \sim \frac{W_{0}}{\left(a \tau_{s}\right) \sqrt{\mathcal{N}}} \sim \frac{m_{3 / 2}}{\log \mathcal{N}} .
$$

Small SUSY breaking scale can be realized naturally!

## Comparison of the 5D LVS with the string LVS


"Swiss-cheese" Calabi-Yau manifold

$$
\mathcal{V}=\left(T_{b}+\bar{T}_{b}\right)^{3 / 2}-\left(T_{s}+\bar{T}_{s}\right)^{3 / 2}
$$

$$
K=-2 \log (\mathcal{V}+\xi) \longleftarrow \alpha^{\prime} \text {-correction }
$$

Non-geometric modulus $\qquad$ Small cycle modulus

## Summary

We construct the LVS in 5D SUGRA without stringy effects.

- General set-up of 5D SUGRA $\rightarrow$ multi-moduli
- Casimir term $\rightarrow$ role of the $\alpha^{\prime}$ correction in string theory


## Future work

Construction of a realistic model in 5D LVS ( Dark matter, inflation, Higgs mass...etc)

## Thank you.

## Appendix

## Anomaly mediation in 5D LVS

$$
\frac{F^{\phi_{C}}}{\phi_{C}}=\frac{m_{3 / 2}}{\mathcal{N}} \ll m_{3 / 2}
$$

Anomaly mediation is much suppressed by the leading no-scale structure.

M.A. Luty and N. Okada (2002)<br>N. Arkani-Hamed and S. Dimopoulos (2005)

## The value of $\xi$

$$
\xi \equiv \frac{\left(\bar{n}_{H}-n_{V}-1\right) \zeta(3)}{32 \pi^{2}}
$$

$n_{V}$ : The number of the vector multiplets
$\bar{n}_{H}$ : The effective number of the hypermultiplets
where
$\bar{n}_{H}=\sum_{a} n_{a} \frac{\mathcal{Z}\left(d_{a} \cdot \operatorname{Re} T / 2\right)}{\mathcal{Z}(0)} \quad \mathcal{Z}(x)=-\int_{0}^{\infty} d \lambda \lambda \ln \left(2 e^{-\sqrt{\lambda^{2}+x^{2}}} \sinh \sqrt{\lambda^{2}+x^{2}}\right)$

## Multiplets in 4D effective theory

|  | 5D vector <br> multiplet (even) |  | 5D vector <br> multiplet(odd) |  | Hypermultiplet |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 4D multiplet | $V^{I}$ <br> vector | $\tilde{T}^{I}$ <br> chiral | $\tilde{V}^{I^{\prime}}$ <br> vector | $T^{I^{\prime}}$ <br> chiral | $Q_{a}$ <br> chiral | $Q_{a}^{\prime}$ <br> chiral |
| parity | + | - | - | + | + | - |
| Zero mode | $V^{I}$ |  |  | $T^{I^{\prime}}$ | $Q_{a}$ |  |
| Role in 4D | vector <br> (gauge) |  |  | moduli | matter |  |

## KK mass \& moduli masses

Gravitino mass: $m_{3 / 2}=\frac{W_{0}}{\sqrt{\mathcal{N}}} \sim \mathcal{O}\left(M_{p l} / \sqrt{\mathcal{N}}\right)$
KK mass \& Moduli mass: $m_{\pi_{0}} \sim \frac{m_{3 / 2}}{\sqrt{\mathcal{N}}} \quad m_{\rho} \sim 0$

$$
\begin{aligned}
& m_{\tau_{s}} \sim m_{\sigma} \sim(\log \mathcal{N}) m_{3 / 2} \quad m_{K K} \sim \frac{M_{p l}}{\mathcal{N}^{1 / 3}} \\
& M_{p l} \gg m_{K K} \gg m_{3 / 2}
\end{aligned}
$$

Analysis by the effective theory is valid!

