

Dark Matter and Entropy Production in the ν MSM

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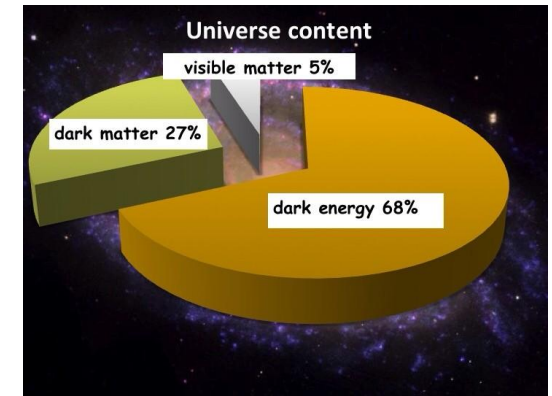
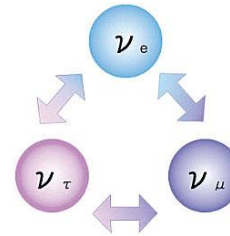
In collaboration with Takehiko Asaka (Niigata University)

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Introduction

■ Problem of the Standard Model

{
 Neutrino mass
 Dark Matter
 BAU



■ νMSM (Neutrino minimal standard model)

[T.Asaka, S.Blanchet, M.Shaposhnikov ('05),
T.Asaka, M.Shaposhnikov ('05)]

νMSM: SM + 3 right-handed neutrinos (ν_R)

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{SM}} + \bar{\nu}_R i \partial_\mu \gamma^\mu \nu_R - F_{\alpha I} \bar{L}_\alpha \tilde{\Phi} \nu_{RI} - \frac{1}{2} (M_M)_{IJ} \bar{\nu}_{RI}^c \nu_{RJ} + h.c.$$

■ Seesaw mechanism $|M_D| = |F\langle\Phi\rangle| \ll M_M < \Lambda_{\text{EW}} \sim \mathcal{O}(100\text{GeV})$

3 heavy neutrino

N_1 is DM candidate

N_2, N_3 explain neutrino mass and BAU

Purpose of study : Impact of entropy production by $N_{2,3}$ decay on DM physics in the νMSM

Motivation 1 – Constraint on DM candidate

The DM candidate N_1

Mass : $O(1\text{keV})$

N_1 behaves as **WDM**.

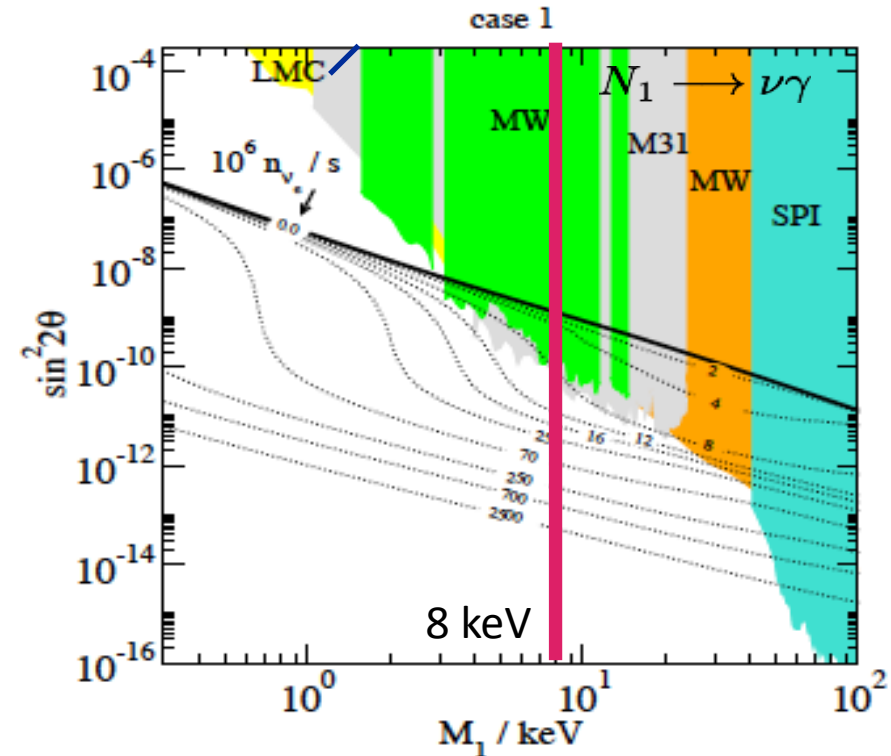
■ Free-streaming length

$$\lambda_{\text{FS}} = 0.2 \text{Mpc} \left(\frac{1\text{keV}}{M_{\text{DM}}} \right) \left(\frac{T_{\text{DM}}}{T_\nu} \right) \left[\ln \left(\frac{t_{\text{EQ}}}{t_{\text{NR}}} \right) + 2 \right]$$

($t_{\text{EQ}} \sim 10^{11}\text{sec}$, $t_{\text{NR}} \sim 10^7\text{sec}$)

■ The lower limit of the DM mass (Lyman- α)

$$M_1 > 8\text{keV} \quad [\text{A. Boyarsky, et al ,arXiv.0812.0010}]$$



[M. Laine, M. Shaposhnikov, ('08)]

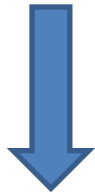
Considering the X-ray and Lyman- α constraints, the allowed region of DM is very limited.

Motivation 1 – Impact of Entropy Production

There is a possibility that the entropy production is induced by the $N_{2,3}$ decay before BBN.

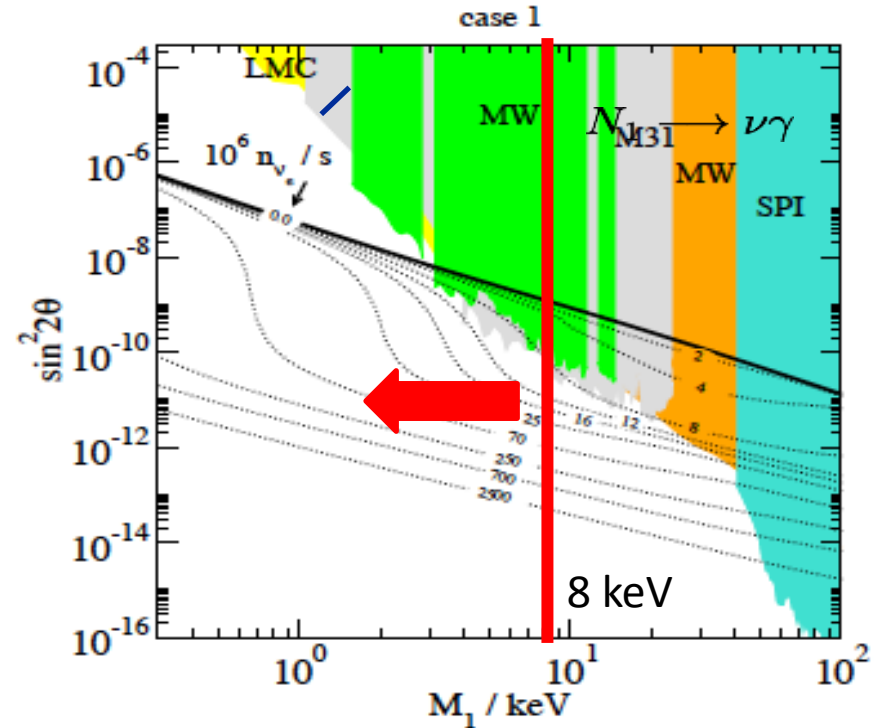
If $N_{2,3}$ decay increase the entropy by ΔS
the λ_{FS} becomes suppressed as,

$$\lambda_{\text{FS}}^{\text{wEP}} = \frac{\lambda_{\text{FS}}^{\text{woEP}}}{\Delta S^{\frac{1}{3}}}$$



The lower limit of DM mass from structure formation is relaxed.

$$M_N > 8\text{keV} \frac{1}{\Delta S^{1/3}}$$



[M. Laine, M. Shaposhnikov, ('08)]

We would like to know how large ΔS is obtained in the νMSM .

Motivation 2 – Evaluation of the previous work

Evaluation of the entropy production rate ΔS in the vMSM have been carried out.

[T.Asaka, M. Shaposhnikov, A. Kusenko ('06) }

■ The maximal value of ΔS in the previous work

- $\Delta S = 29$ [$T_R > 0.7\text{MeV}$]
- $\Delta S = 10$ [$T_R > 4\text{MeV}$]

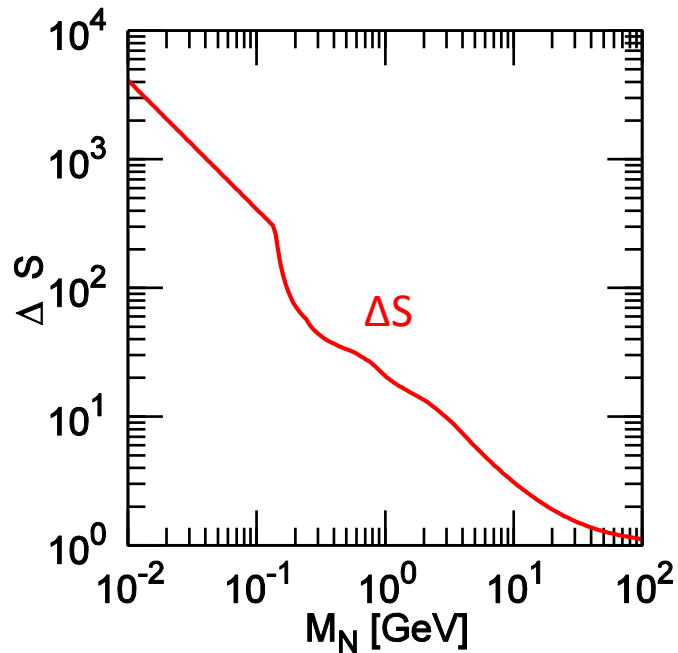
However, there were the unsatisfactory points .

- 1. Evaluation of the lifetime of the $N_{2,3}$ was incomplete.
(In particular the decay modes into meson were not included.)
- 2. Only specific Yukawa coupling constant is considered.

We would like to solve these points.

- 1. All decay modes Including the decay into meson.
(example : $N \rightarrow \pi\nu$, Kl , Bl)
- 2. All parameter region in the Yukawa coupling constant .

Results



To evaluate the maximal of ΔS

mass : $M_2 = M_3 = M_N$

All parameter region in Yukawa coupling constant

■ The constraint on the reheating temperature by the $N_{2,3}$ decay

$$T_R > 2\text{MeV} \quad (\text{WMAP5})$$

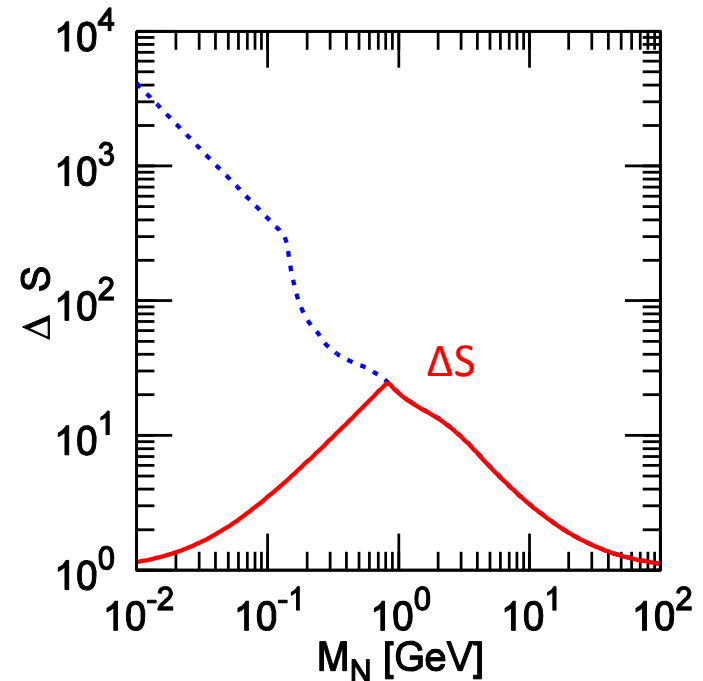
[F. De. Bernardis et al. ('08)]

When M is lighter than 800MeV, the maximum of ΔS becomes suppressed.

■ The maximal value of ΔS in this work

$$\Delta S = 24 \quad \text{at} \quad M_N = 800\text{MeV}$$

[T. Asaka, T.K. in preparation]



Results

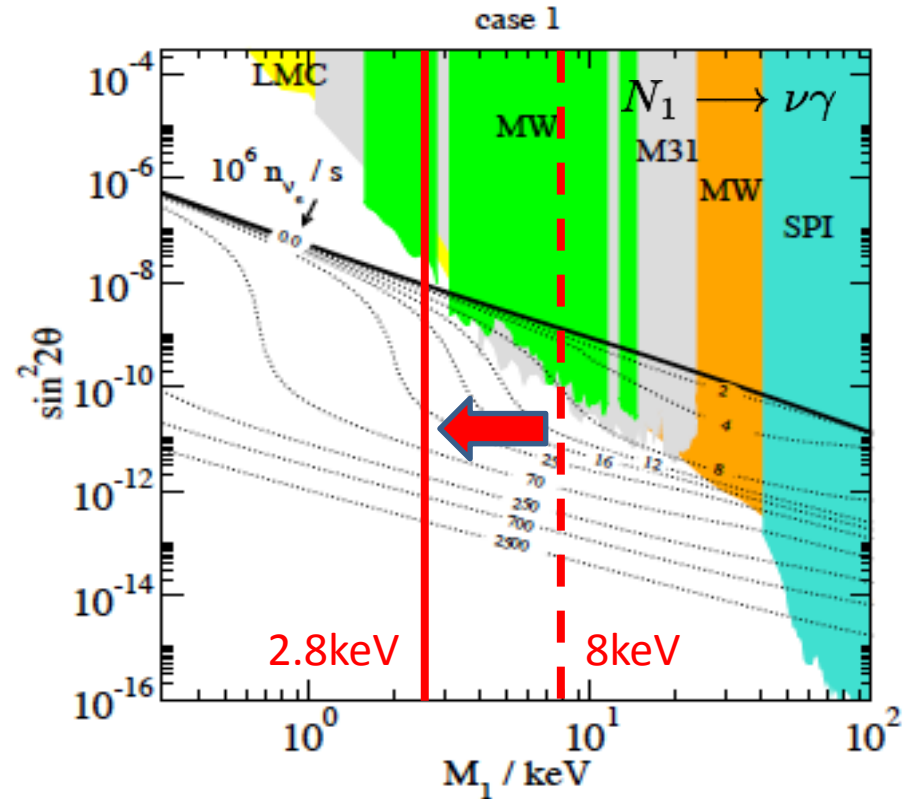
■ Impact to dark matter physics

$$M_1 > 8\text{keV}$$

[A. Boyarsky, et al ,arXiv.0812.0010]



$$M_1 > 2.8\text{keV} \left(\frac{24}{\Delta S} \right)^{\frac{1}{3}}$$



The lower limit of DM mass is relaxed and the allowed region of DM becomes wider.

Summary

In this talk, We discussed **the Entropy Production** by $N_{2,3}$ decay in the vMSM.

■ In the Evaluation of the Entropy Production

- We calculated the decay width for the possible decay modes of $N_{2,3}$.
- We evaluated ΔS with all free parameter in the Yukawa coupling constant.

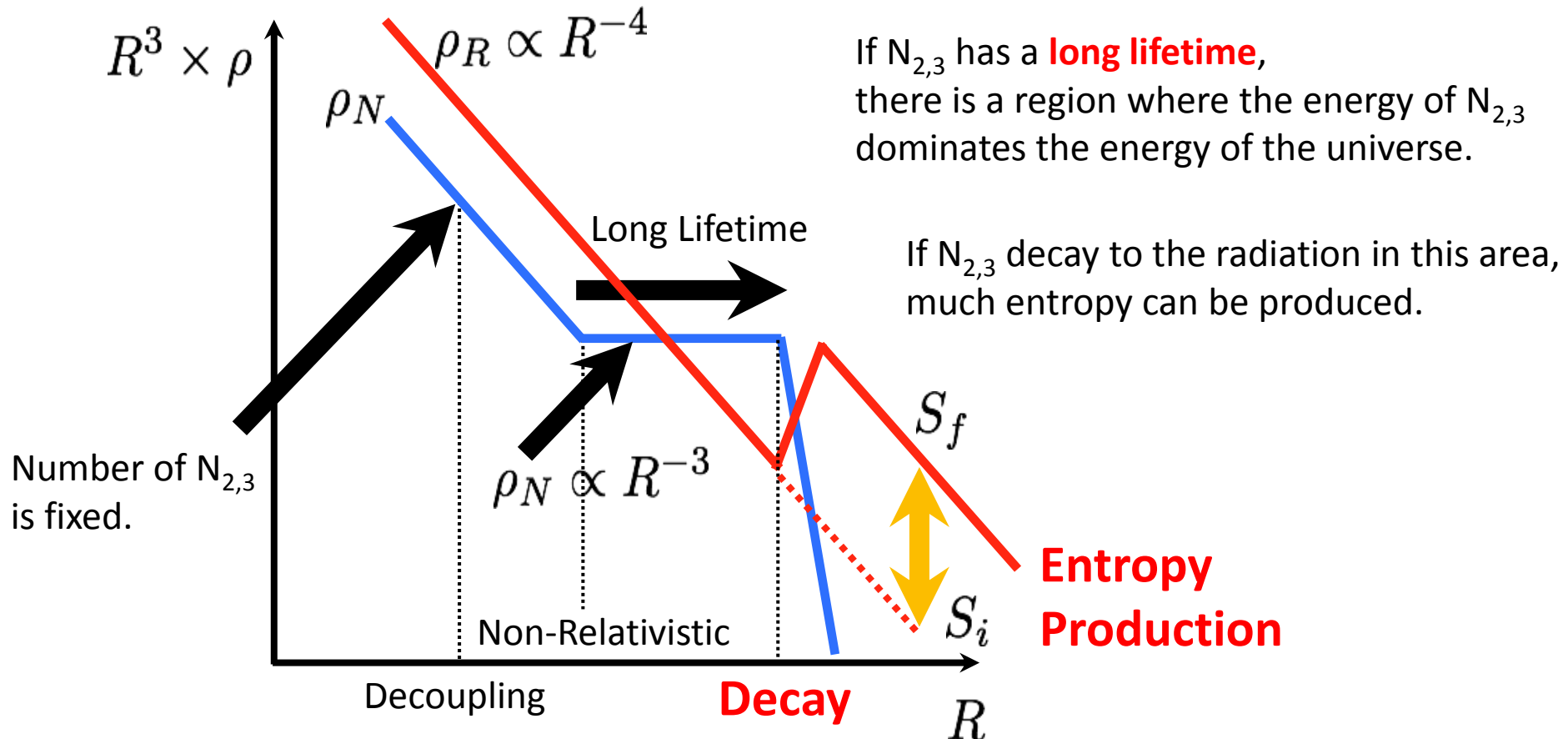
■ Results

- By constraint from the cosmic background radiation, we got $\Delta S < 24$.
- The DM mass limit is relaxed by the influence to free-Streaming length.
- By the entropy production, the baryon number asymmetry and the dark matter abundance are diluted. These values have to be produced more larger by ΔS .

Back Up

The Entropy Production by $N_{2,3}$ Decay

Consider the case that $N_{2,3}$ decoupled from heat bath and behave Non-relativistic.



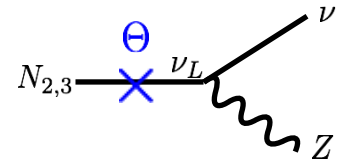
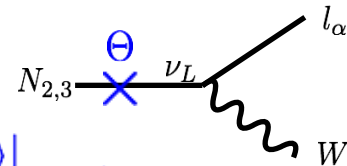
We had better check if $N_{2,3}$ lifetime is sufficiently long to evaluate the produced entropy of the universe.

Decay of $N_{2,3}$

■ neutrino mixing: $\nu_{L\alpha} = U_{\alpha i} \nu_i + \Theta_{\alpha I} N_I^c$

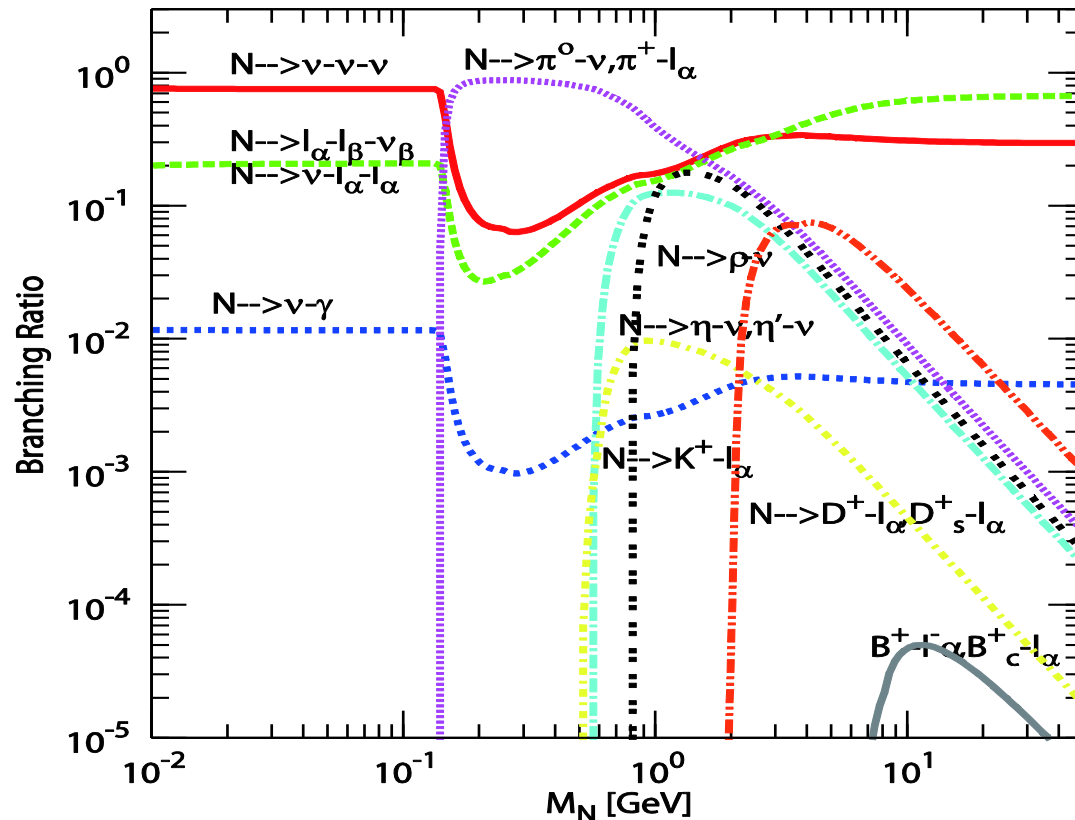
$$|F\langle\Phi\rangle| \ll M_M < \mathcal{O}(100\text{GeV})$$

$$\Theta = \frac{|F\langle\Phi\rangle|}{M_M} \ll 1$$



In the vMSM, $N_{2,3}$ has a weak interaction that suppressed by mixing matrix Θ .

\Rightarrow We can expect that $N_{2,3}$ lifetime is sufficiently long.



■ The decay modes

$$\begin{aligned}
 N_{2,3} &\rightarrow \nu\nu\nu \\
 &\rightarrow \nu l^+ l^- \\
 &\rightarrow \nu\gamma \\
 &\rightarrow \pi^0 \nu, \pi^+ l^- \\
 &\rightarrow K^+ l^- \\
 &\rightarrow \rho\nu \\
 &\rightarrow \eta\nu, \eta'\nu \\
 &\rightarrow D^+ l^-, D_s^+ l^- \\
 &\rightarrow B^+ l^-, B_c^+ l^-
 \end{aligned}$$

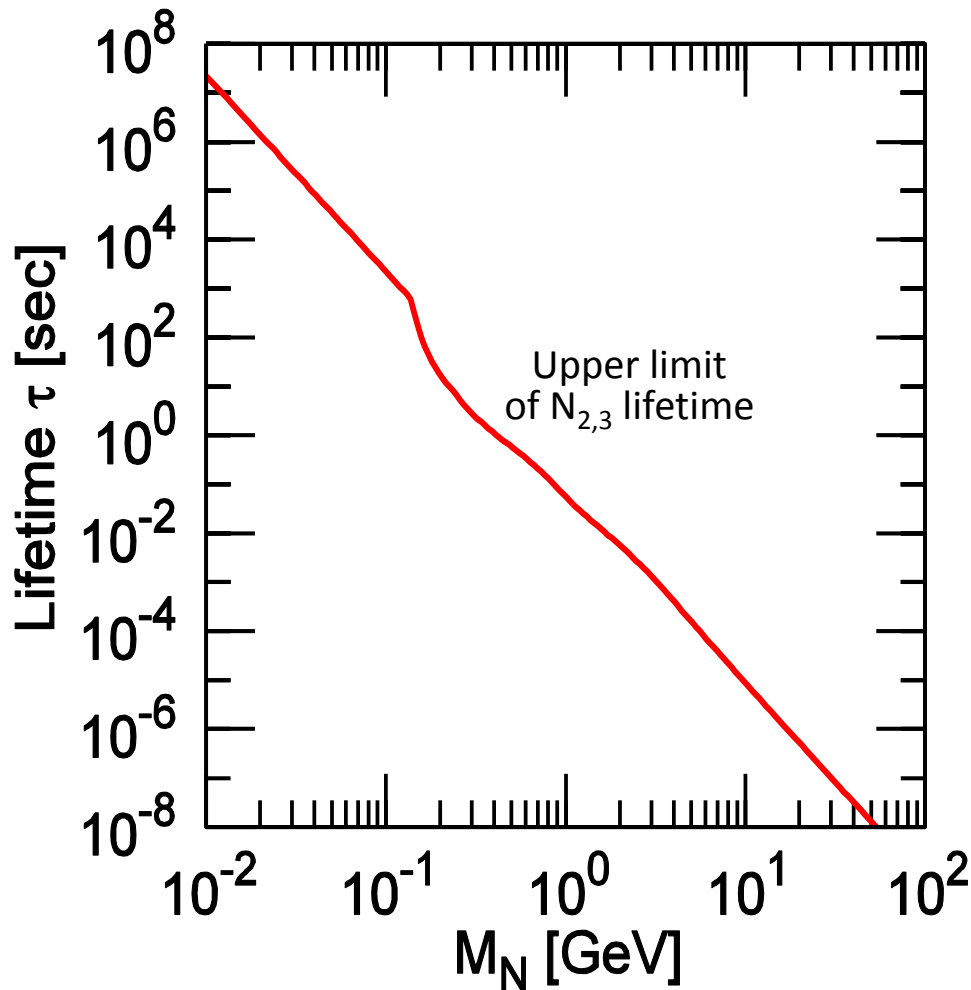
\Rightarrow The main decay modes change according to $N_{2,3}$ mass.

Upper limit of $N_{2,3}$ lifetime

We have to fix Yukawa coupling constant and $N_{2,3}$ mass to calculate their lifetime.

$$N_{2,3} \text{ mass: } M_{N2} = M_{N3} = M_N$$

Yukawa coupling : Consider all free parameter region



■ The rough behavior of lifetime τ

$$\tau \propto \begin{cases} |\Theta|^{-2} M_N^{-5} \text{ (3-body decay)} \\ |\Theta|^{-2} M_N^{-3} \text{ (2-body decay)} \end{cases}$$

If $N_{2,3}$ mass is light,
their lifetime is sufficiently long.

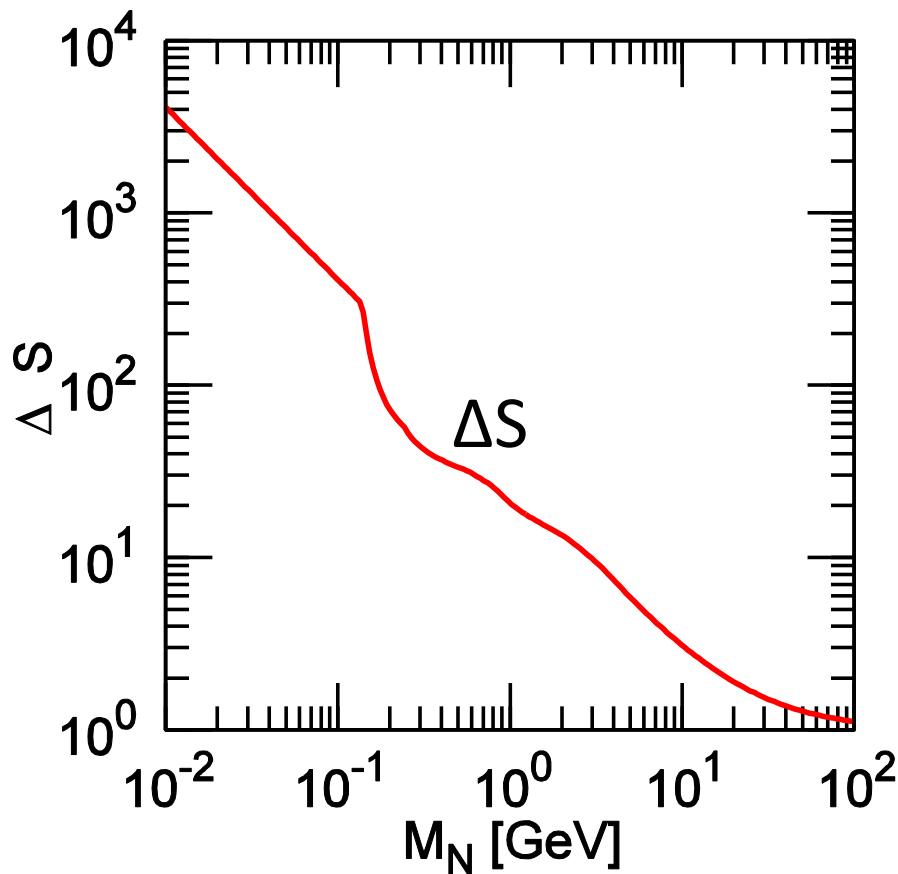
Evaluation of The entropy Production rate

■ The Evaluation formula of entropy production rate in the vMSM

$$\Delta S = \frac{S_f}{S_i} = \left[1 + \left(\frac{1.37 \mathcal{T}}{M_{pl} \left(\frac{g_*(T_D)}{M_N} \right)^2} \right)^{\frac{2}{3}} \right]^{\frac{3}{4}} \propto (M_N^2 \mathcal{T})^{\frac{1}{2}}$$

T_D : temperature of $N_{2,3}$ decoupling

[R.J. Scherrer, M.S. Turner, Phys. Rev. D31 (1985) 681]



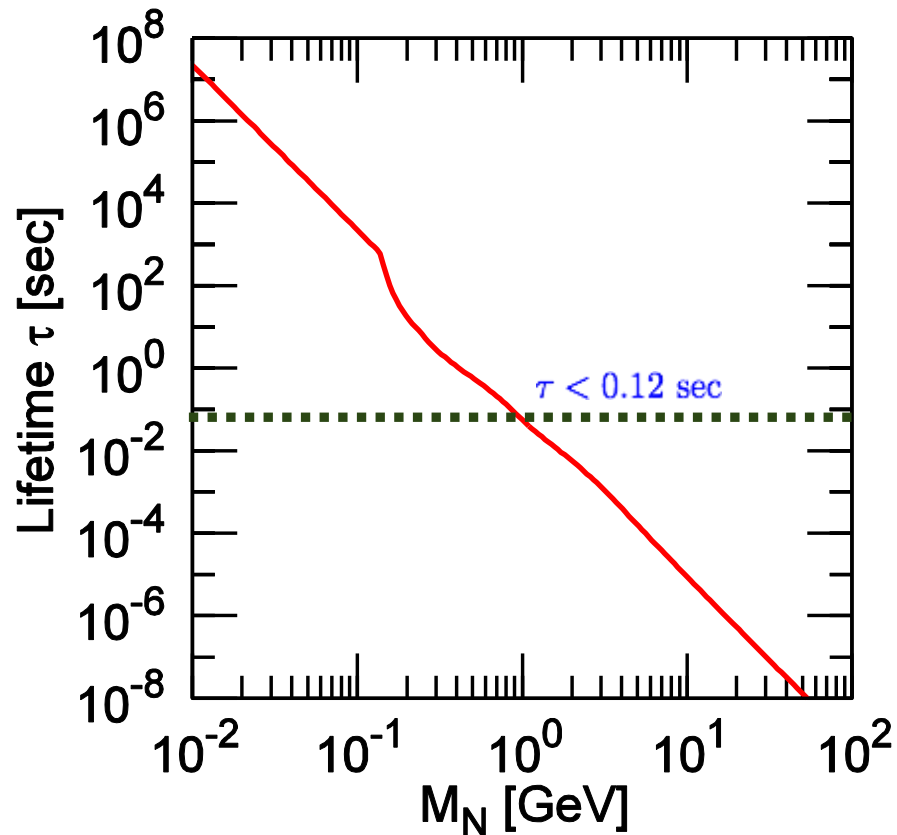
Assumption : $g_*(T_D) \approx 10.75$

In the region where $N_{2,3}$ mass is light, a lot of Entropy is produced.

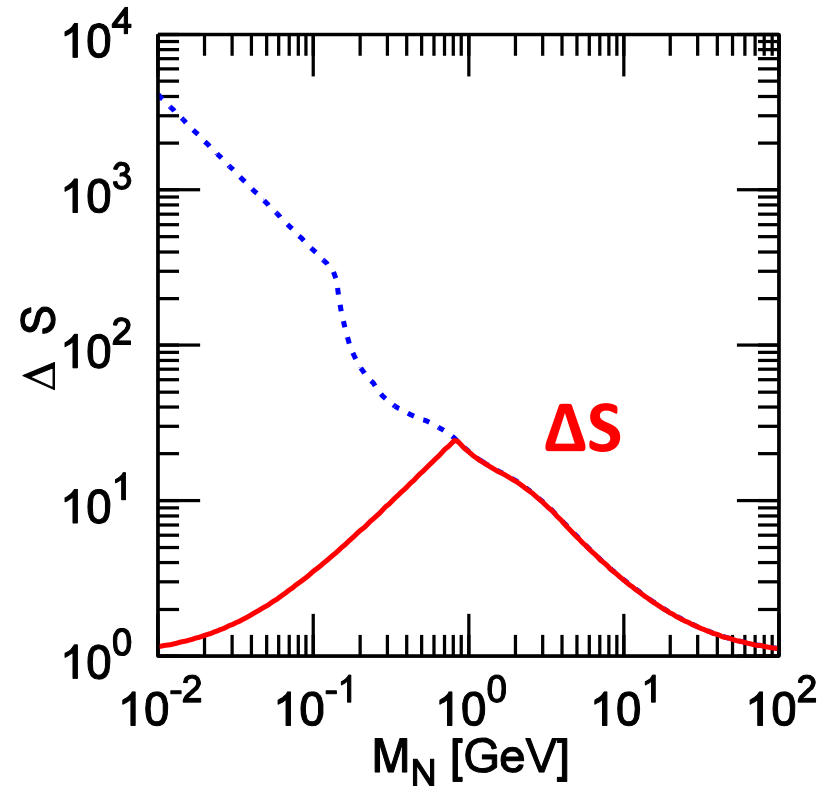
Constraint from the cosmic background radiation

■ The upper limit of $N_{2,3}$ lifetime

$$\tau < 0.12 \text{ sec}$$



In the vMSM, the upper limit of $N_{2,3}$ lifetime is lower than 0.12s in the region $M_N > 800 \text{ MeV}$.



In the region $M_N < 800 \text{ MeV}$, we fixed $\tau = 0.12 \text{ sec}$.

The upper limit of ΔS

$$\Delta S < 24$$

[T. Asaka, T.K. ('13)]

The Yukawa coupling constant

$$F = \frac{i}{v} U D_\nu^{1/2} \Omega D_N^{1/2}$$

U : PMNS matrix (3 mixing angles and 2 CP phases)

$$D_\nu = \text{diag}(m_1, m_2, m_3) \quad (\text{m is the mass of light neutrinos.})$$

$$D_N = \text{diag}(M_2, M_3) \quad (\text{M is the mass of heavy neutrinos.})$$

$$\Omega_{\text{NH}} = \begin{pmatrix} 0 & 0 \\ \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \end{pmatrix} \quad \Omega_{\text{IH}} = \begin{pmatrix} \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \\ 0 & 0 \end{pmatrix}$$

Free parameter : $M_2, M_3, \text{Re}\omega, \text{Im}\omega, \xi, \text{Dirac phase, Majorana phase}$