

# Markov Chain Monte Carlo solution of BK equation

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# Introduction

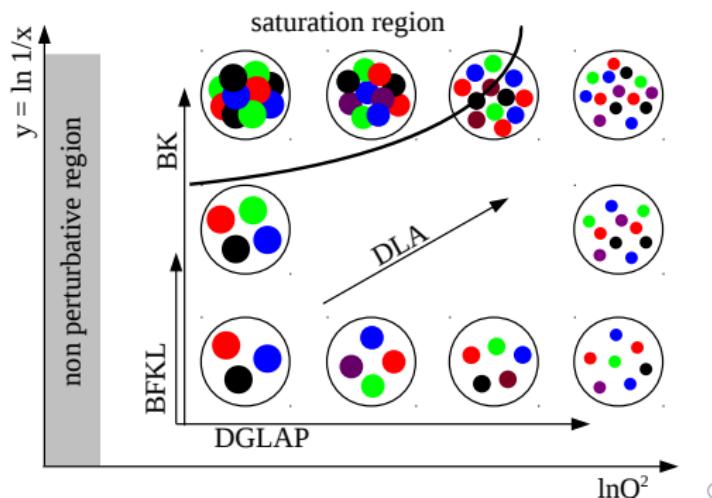
- To do any physics @LHC we need good understanding of PDFs.
- Several parton density evolution equations on the market.
- Current MC generators based on linear evolution equations (PYTHIA, Herwig++, SHERPA, CASCADE) with exception of DIPSY.

## Linear evolution

- DGLAP  
Evolution in  $Q^2$
- BFKL  
Evolution in  $1/x$

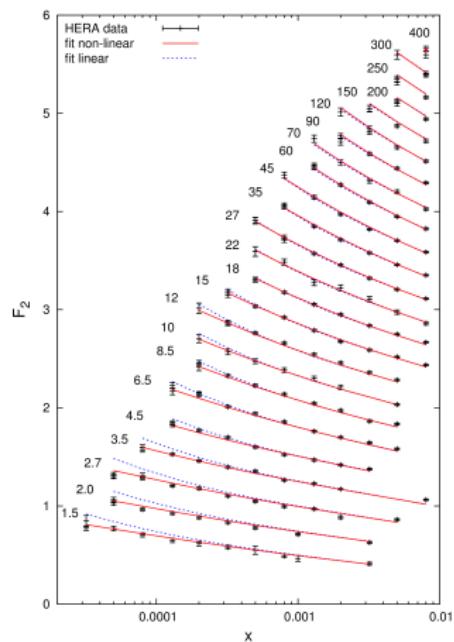
## Non linear evolution

- Balitsky-Kovchegov (BK)  
BFKL + saturation



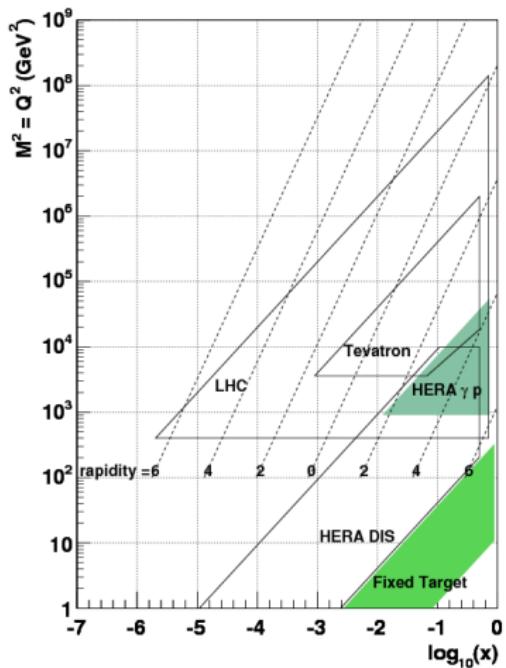
# Introduction

- Hints of saturation



[Kutak, Sapeta (2012)]

- LHC opens new kinematic region



[Butterworth et al. (2004)]

# BK vs BFKL

General structure of QCD equations for parton densities

$$\Phi(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \Phi(., .)), \quad (1)$$

where  $\mathbb{K}$  and  $\mathbb{G}$  are  $N$ -dim (for some  $N$ ).

## BFKL equation

- Kernel  $K$  linear in  $\Phi$ .
- Range of numerical methods.
- Among them Monte Carlo (MC) methods as well, e.g. **HEJ** [Andersen, Smillie], [Schmidt (1996)].

## BK equation

- Kernel  $K$  non-linear in  $\Phi$ ,
- Range of numerical methods **BKsolver** [Enberg, et al. (2005)], [Golec-Biernat, et al. (2001)]
- Convergence of Monte Carlo methods is not guaranteed,
- No known MC solutions.

# Motivation

- MC method for the BK equation.
- New approach to the integral equations.

$$\Phi(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \Phi(., .))$$

This form allows to extend the BK equation, e.g. KGBJS (nonlinear extension of CCFM) equation [Kutak, et al. (2012)], NLO BK equation [Balitsky, Chirilli (2010)].

$$\begin{aligned}\mathcal{E}(x, k^2, p) &= \mathcal{E}_0(x, k^2, p) \\ &+ \bar{\alpha}_s \int_x^1 dz \int \frac{d^2 \bar{\mathbf{q}}}{\pi \bar{q}^2} \theta(p - z \bar{q}) \Delta_s(p, z \bar{q}) \left( \frac{\Delta_{ns}(z, k, q)}{z} + \frac{1}{1-z} \right) \times \\ &\left[ \mathcal{E}\left(\frac{x}{z}, k'^2, \bar{q}\right) - \bar{q}^2 \delta(\bar{q}^2 - k^2) \mathcal{E}^2\left(\frac{x}{z}, \bar{q}^2, \bar{q}\right) \right].\end{aligned}$$

# Newton–Kantorovich method

How to make the BK equation solvable by a MC method?



$$\Phi(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \Phi(., .)),$$

Linearize non-linear kernel!

Taylor expansion of the kernel

$$K(\Phi(\mathbb{X})) = K(\bar{\Phi}(\mathbb{X})) + K'_{\Phi}(\bar{\Phi}(\mathbb{X})) [\Phi(\mathbb{X}) - \bar{\Phi}(\mathbb{X})] + \mathcal{O}\left([\Phi(\mathbb{X}) - \bar{\Phi}(\mathbb{X})]^2\right). \quad (2)$$

where  $K'_{\Phi}(\bar{\Phi}(\mathbb{X})) = \frac{\delta K(\Phi(\mathbb{X}))}{\delta \Phi(\mathbb{X})}$  is a functional derivative.

If we neglect higher-order terms we are done.

## Newton–Kantorovich method

Assuming  $|\Psi(\mathbb{X})| := |\Phi(\mathbb{X}) - \bar{\Phi}(\mathbb{X})| \ll 1$

$$K(\Phi(\mathbb{X})) \approx K(\bar{\Phi}(\mathbb{X})) + K_{\Phi}'(\bar{\Phi}(\mathbb{X})) \cdot \Psi(\mathbb{X}).$$

the initial equation

$$\Phi(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \Phi(., .)),$$

transforms to the *set of equations*:

$$\Phi(y, \mathbb{K}) = \bar{\Phi}(y, \mathbb{K}) + \Psi(y, \mathbb{K}), \quad (3)$$

$$\Psi(y, \mathbb{K}) = \Lambda(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K_{\Phi}'(y, t, \mathbb{K}, \mathbb{G}, \bar{\Phi}(., .)) \Psi(t, \mathbb{G}), \quad (4)$$

$$\Lambda(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) - \bar{\Phi}(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \bar{\Phi}(t, \mathbb{G})) \quad (5)$$

# Newton–Kantorovich method

Assuming  $|\Psi(\mathbb{X})| := |\Phi(\mathbb{X}) - \bar{\Phi}(\mathbb{X})| \ll 1$

$$K(\Phi(\mathbb{X})) \approx K(\bar{\Phi}(\mathbb{X})) + K'_{\Phi}(\bar{\Phi}(\mathbb{X})) \cdot \Phi(\Psi(\mathbb{X})).$$

the initial equation

$$\Phi(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \Phi(., .)),$$

transforms to the *set of equations*:

"given"  $\bar{\Phi}(y, \mathbb{K})$

$$\Phi(y, \mathbb{K}) = \bar{\Phi}(y, \mathbb{K}) + \Psi(y, \mathbb{K}),$$

Integral equation on  $\Psi$ ,  
but *linear* (3)

$$\Psi(y, \mathbb{K}) = \Lambda(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K'_{\Phi}(y, t, \mathbb{K}, \mathbb{G}, \bar{\Phi}(., .)) \Psi(t, \mathbb{G}), \quad (4)$$

$$\Lambda(y, \mathbb{G}) = \Phi^0(y, \mathbb{K}) - \bar{\Phi}(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(y, t, \mathbb{K}, \mathbb{G}, \bar{\Phi}(t, \mathbb{G})), \quad (5)$$

Simple Integration

# Newton–Kantorovich method

- Unfortunately,  $\bar{\Phi}(X)$  is not given...      Solution – Iteration

$$\Phi_n(y, \mathbb{K}) = \Phi_{n-1}(y, \mathbb{K}) + \Psi_{n-1}(y, \mathbb{K}) \quad (6)$$

$$\Psi_{n-1}(y, \mathbb{K}) = \Lambda_{n-1}(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K_{\Phi}'(., \Phi_{n-1}(.)) \Psi_{n-1}(t, \mathbb{G}) \quad (7)$$

$$\Lambda_{n-1}(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) - \Phi_{n-1}(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(., \Phi_{n-1}(.)) \quad (8)$$

- To compare with, a straight forward iterative method on the initial equation reads:

Iteration method on initial equation

$$\Phi_n(y, \mathbb{K}) = \Phi_{n-1}(y, \mathbb{K}) + \Psi_{n-1}(y, \mathbb{K}) \quad (9)$$

$$\Psi_{n-1}(y, \mathbb{K}) = \Lambda_{n-1}(y, \mathbb{K}) \quad (10)$$

$$\Lambda_{n-1}(y, \mathbb{K}) = \Phi^0(y, \mathbb{K}) - \Phi_{n-1}(y, \mathbb{K}) + \int_{y_0}^y dt \int d\mathbb{G} K(., \Phi_{n-1}(.)) \quad (11)$$

- The Newton–Kantorovich method more complicated, but accessible by MC methods.

# Markov Chain Monte Carlo (MCMC)

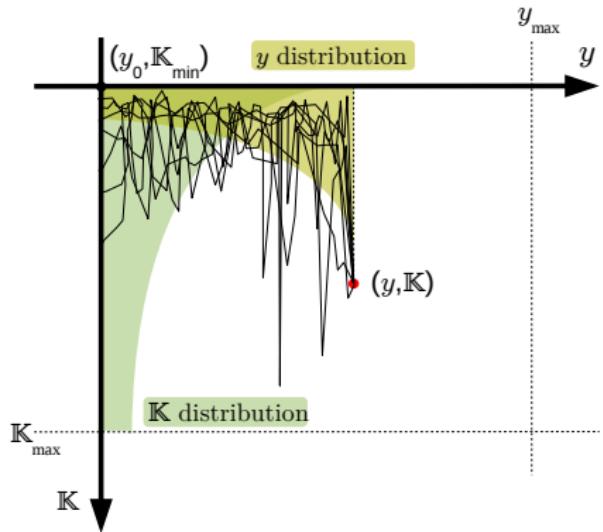
We can solve the linear equation that emerges in the Newton–Kantorovich method using a Monte Carlo method.

$$\Psi_{n-1}(\mathbf{y}, \mathbb{K}) = \Lambda_{n-1}(y, \mathbb{K}) + \int_{y_0}^{\mathbf{y}} dt \int d\mathbb{G} K_{\Phi}'(., \Phi_{n-1}(.)) \Psi_{n-1}(t, \mathbb{G})$$

For example by a random walk  
(Markov chain) ordered in  $y$ .

For each point  $(y, \mathbb{K})$ :

- generate a trajectory  $(y, \mathbb{K}) \rightarrow (y_0, \mathbb{X})$  with points generated according to some distributions.
- average over a number  $M$  of trajectories.



## Example

To test numerical feasibility of our method we solve the leading-order BK equation for the Weizsäcker–Williams gluon density  $\Phi(y, k^2)$ :

$$\Phi(y, k^2) = \Phi^0(y, k^2) + \int_{y_0}^y dt \int_0^\infty dl^2 K(y, t, k^2, l^2, \Phi(t, l^2)), \quad (12)$$

where the non-linear kernel  $K$  reads:

$$K = \frac{\bar{\alpha}_s}{l^2} \left[ \frac{l^2 \Phi(t, l^2) - k^2 \Phi(t, k^2)}{|k^2 - l^2|} + \frac{k^2 \Phi(t, k^2)}{\sqrt{4l^4 + k^4}} \right] - \bar{\alpha}_s \delta(l^2 - k^2) \Phi^2(t, l^2). \quad (13)$$

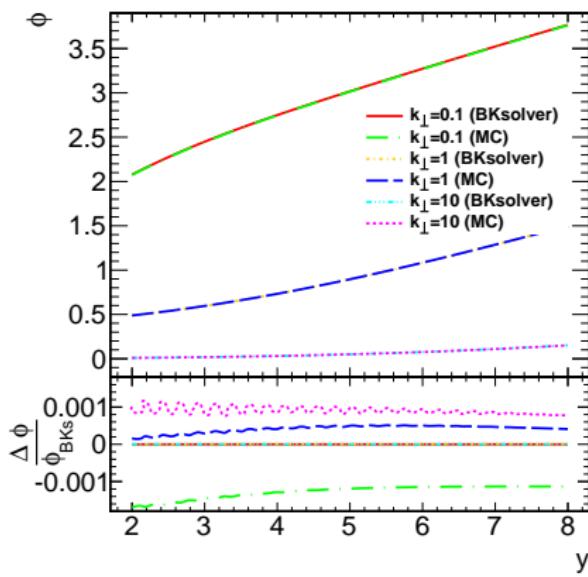
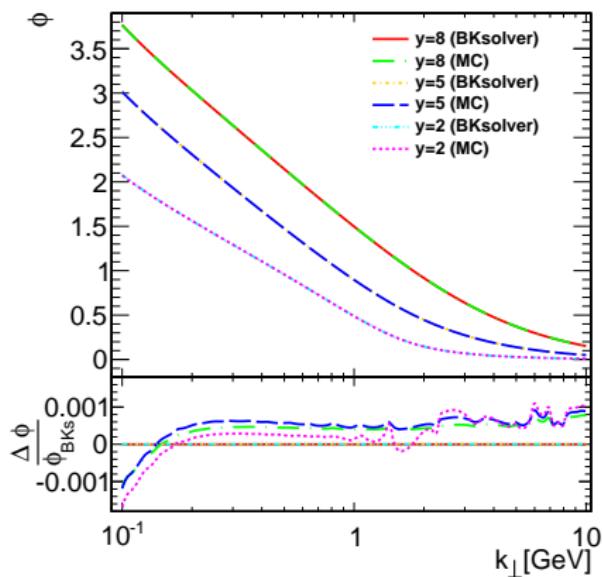
We took  $\Phi^0(y, k^2) = \exp(-\frac{k^2}{\text{GeV}})$ , and generated the MC trajectories according to the following distributions:

$$\eta(l_i^2) = \frac{\mu^2}{(l_i^2)^2} \implies \text{for } l^2 \quad (14)$$

$$\rho(\Delta t_i) = e^{\Delta t_i} \implies \text{for } t, \Delta t_i = t_i - t_{i-1} \leq 0 \quad (15)$$

# Example

- Solutions of the BK equation obtained by our MCMC method are compared with results from the external package BKSolver.



- The agreement is at the level of 0.1%!

# Final words

## Conclusion

- We proposed a MC method to solve the non-linear BK equation in integral form which is the most general form of the PDF evolution equations.
- The result from the new method agree at the level of 0.1% with the solution from the independent program **BKSolver**.
- The proposed method is feasible to handle efficiently multi-dimentional and complex problems.

## Outlook

- A code to solve an arbitrary high-dimensional integral equation
  - ongoing.
- Solving other BK-like equations, e.g. the exclusive version of the BK equation, KGBJS, NLO BK equation.
- Possibility to construct a Monte Carlo event generator based on non-linear integral equations?