

SUSY Dark Matter Annihilation in the Galactic halo

V. Berezinsky², V. Dokuchaev¹ and Yu. Eroshenko¹

¹*Institute for Nuclear Research, RAS, Russia*

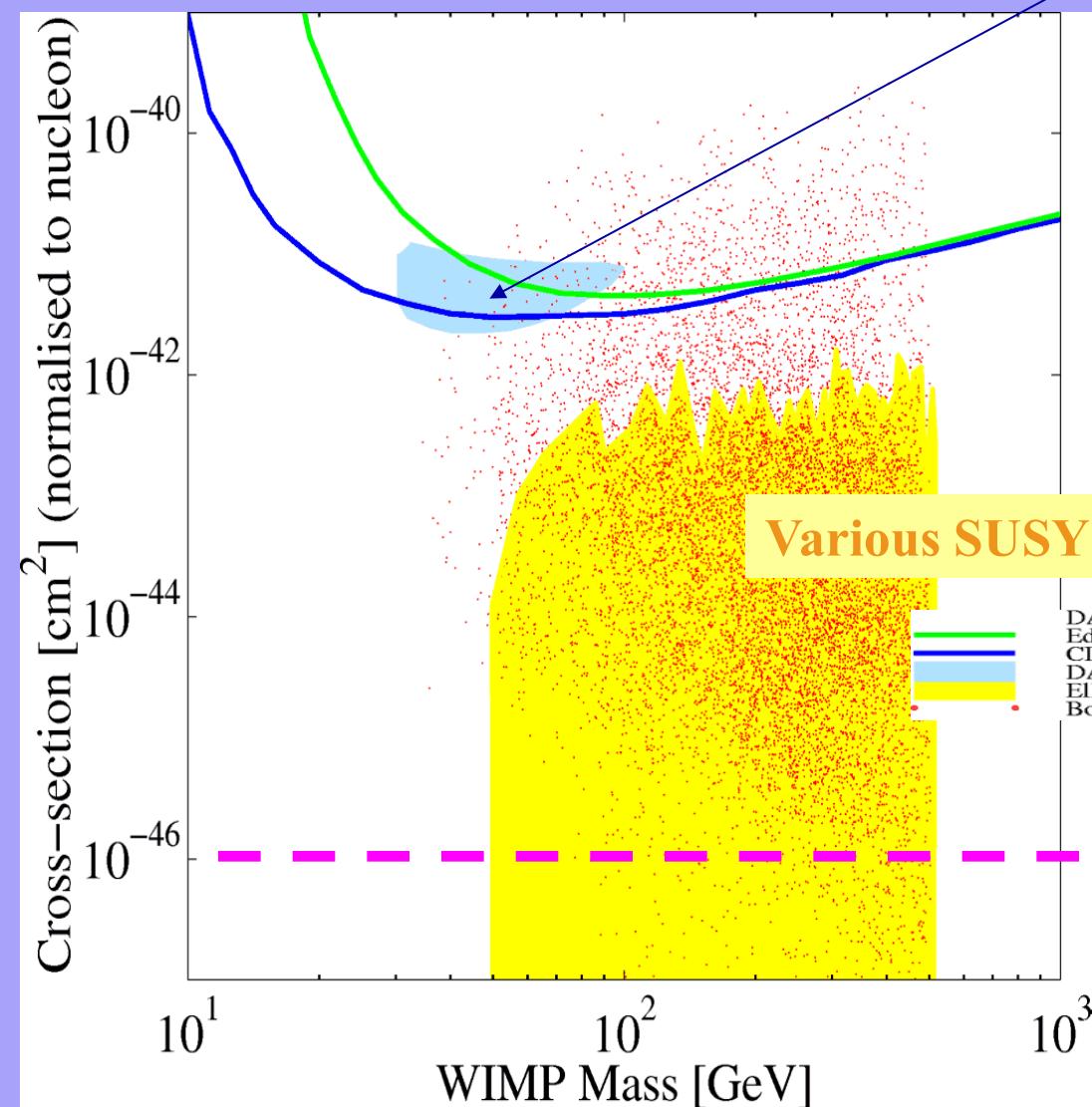
²*Laboratori Nazionali del Gran Sasso, INFN, Italy*

SUSY neutralino DM detection

- Direct neutralino χ detection
 - * Recoil detectors: DAMA $m_\chi \approx 60$ GeV?, Edelweiss, CDMS, XENON
 - LHC $m_\chi = \mathcal{O}(100)$ GeV?
- Indirect χ detection
 - * $\chi \chi$ annihilation in the galactic haloes
 - * WIMP annihilation in the Sun, Earth, Moon...
 - ...
- ◆ Gamma-rays, positrons, antiprotons, neutrino
- * Halo structure: DM profile $\rho_{\text{halo}}(r)$ – NFW, Moore, Einasto...
- * Halo clumpiness: mass fraction of clumps ξ
- * Minimal mass of clumps M_{\min}
- * Clump distribution function $\xi(M_{\text{cl}}, \rho_{\text{cl}}, r)$

Direct DM detection

Experiments: CDMS, CRESST, CUORE, DAMA, EDELWEISS, EURECA, PICASSO, XENON, XMASS-DM, ZEPLIN...



DATA listed top to bottom on plot
Edelweiss, 4.5 kg-days Ge(320g) June 2001 limit
CDMS Feb. 2000 ver. sub. to PRL
DAMA 2000 58k kg-days NaI Ann.Mod. 3sigma,w/o DAMA 1996 limit
Ellis et al., Spin indep. sigma in MSSM
Bottino et al., hep-ph/0001309 SUSY

10^{-10} pb

<http://dmtools.berkeley.edu>
(Gaitskell/Mandic)

Neutralino as WIMP

χ is LSP in mSUGRA: MSSM with SUGRA inspired breaking

spin 1/2 Majorana particle: $|\chi\rangle = N_1|B_0\rangle + N_2|W_3\rangle + N_3|H_1\rangle + N_4|H_2\rangle$

B_0, W_3 – gauginos, H_1, H_2 - higgsinos

χ is almost pure bino: $(N_1, N_2, N_3, N_4) = (0.95, -0.10, 0.27, -0.09)$

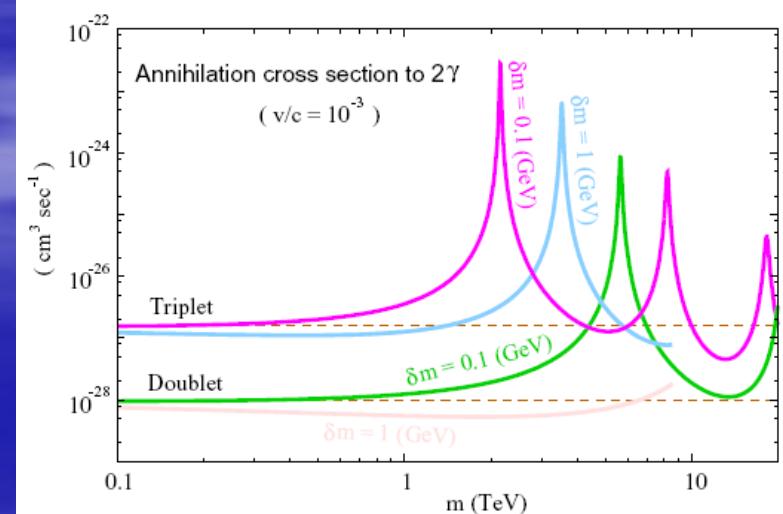
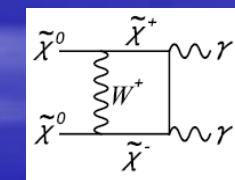
Only 5 parameters:

Parameter	Value
m_0	1500 GeV
$m_{1/2}$	170 GeV
A_0	$0 \cdot m_0$
$\tan \beta$	52.2
sign μ	+
$\alpha_s(M_Z)$	0.122
$\alpha_{em}(M_Z)$	0.0078153697
$1/\alpha_{em}$	127.953
$\sin^2(\theta_W)_{MS}$	0.2314
m_t	175 GeV
m_b	4.214 GeV

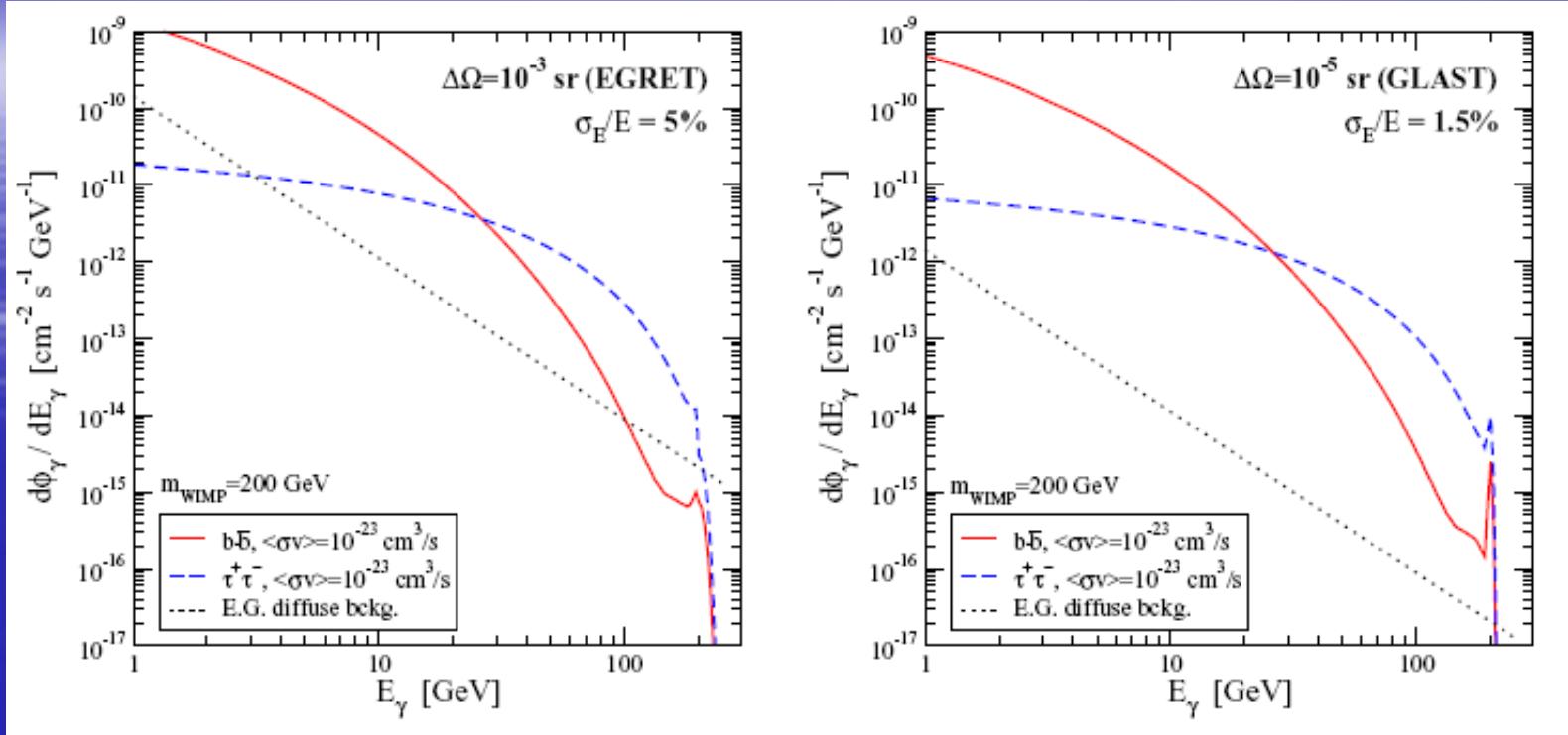
Particle	Mass [GeV]
$\tilde{\chi}_{1,2,3,4}^0$	64, 113, 194, 229
$\tilde{\chi}_{1,2}^\pm, \tilde{g}$	110, 230, 516
$\tilde{e}_{1,2} = \tilde{c}_{1,2}$	1519, 1523
$\tilde{d}_{1,2} = \tilde{s}_{1,2}$	1522, 1524
$\tilde{t}_{1,2}$	906, 1046
$\tilde{b}_{1,2}$	1039, 1152
$\tilde{e}_{1,2} = \tilde{\mu}_{1,2}$	1497, 1499
$\tilde{\tau}_{1,2}$	1035, 1288
$\tilde{\nu}_e, \tilde{\nu}_\mu, \tilde{\nu}_\tau$	1495, 1495, 1286
h, H, A, H^\pm	115, 372, 372, 383
Observable	Value
$Br(b \rightarrow X_s \gamma)$	$3.02 \cdot 10^{-4}$
Δa_μ	$1.07 \cdot 10^{-9}$
Ωh^2	0.117

Dominant annihilation into quark pairs

Annihilation to VV is suppressed by a loop factor



Photon spectra of SUSY $\chi\chi$ annihilation



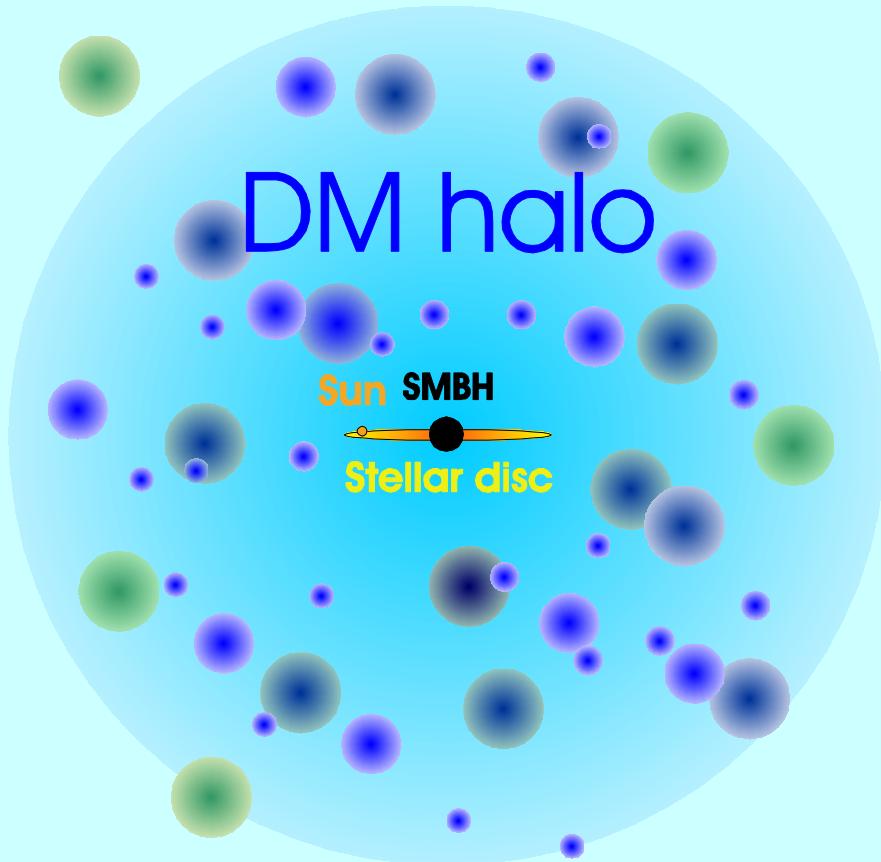
μ	m_1	m_2	m_3	m_A	$m_{\tilde{S}}$	$A_{\tilde{S}_3}$	$\tan\beta$
$30 \div 1200$	$2 \div 1200$	$50 \div 1200$	$m_{\text{LSP}} \div 20000$	$100 \div 10m_{\text{LSP}}$	$(1 \div 10)m_{\text{LSP}}$	$(-3 \div 3)m_{\tilde{S}}$	$1 \div 60$

Table 2: Ranges of the MSSM parameters used to generate the models shown in Figs. 6 and 8. All masses are in GeV, and $m_{\text{LSP}} \equiv \min(\mu, m_1, m_2)$. The quantity $m_{\tilde{S}}$ indicates the following scalar masses (which were independently sampled): $m_{\tilde{Q}_{1,3}}, m_{\tilde{u}_{1,3}}, m_{\tilde{d}_{1,3}}, m_{\tilde{L}_{1,2,3}}, m_{\tilde{e}_{1,2,3}}$. To avoid FCNC constraints, we assumed the squark soft supersymmetry breaking terms of the first two generations to be equal. $A_{\tilde{S}_3}$ stands for the third generation sfermion trilinear terms: those of the first two generations were taken to vanish.

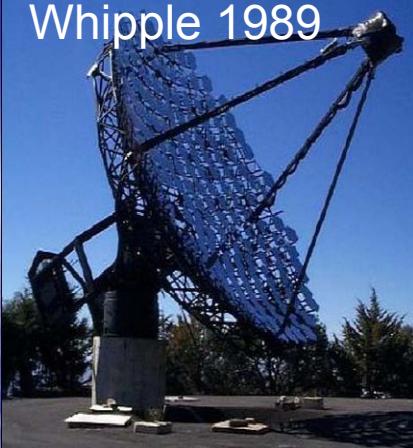
Boosting factor $\eta \sim 100$ is required!

Clumpiness of DM.

Small-scale DM clumps in the Galactic halo



Whipple 1989



Atmospheric Cherenkov Telescopes

CANGAROO III



MAGIC



H.E.S.S. 2002



HEGRA 1997



CELESTE



VERITAS
photomontage



CACTUS Gamma-Ray Excess from Draco

CACTUS: Converted Atmospheric Cherenkov Telescope Using Solar-2

Solar 2 Heliostat Array CACTUS. Barstow, California, effective area $5 \times 10^4 \text{ m}^2$, 144 heliostats, each 42 m^2

Dwarf spheroidal (dSph) galaxies within 100 kpc from the Milky Way center:

Carina, Draco, Ursa Minor and Sextans

Boosting factor ~ 40 insufficient for detection

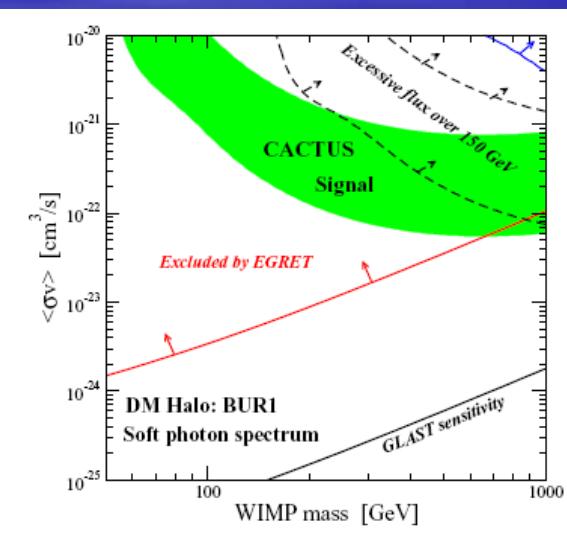
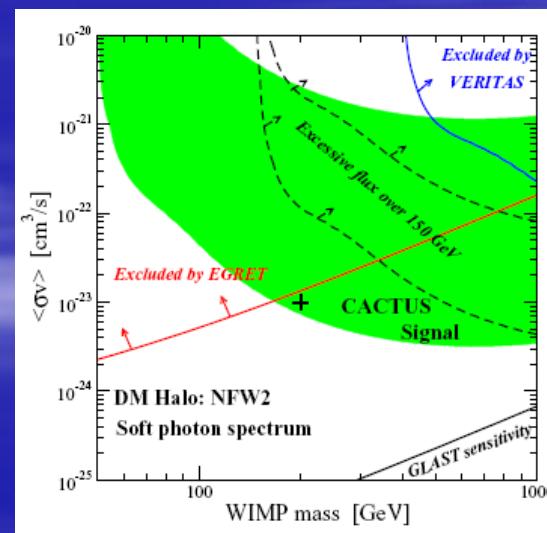
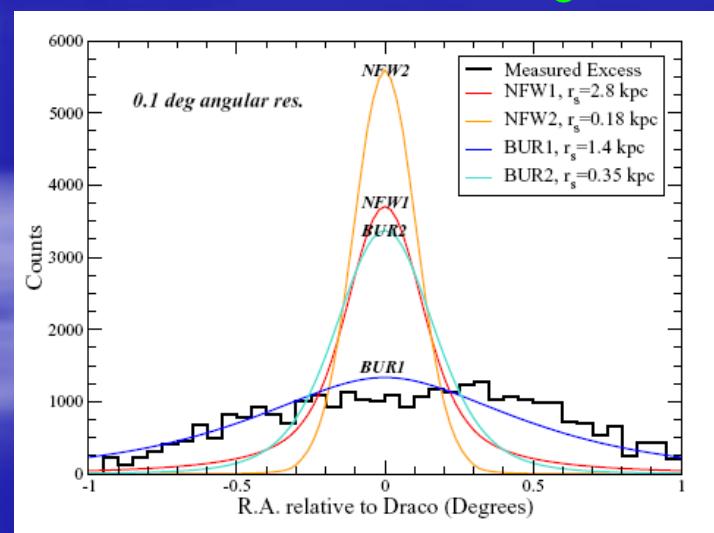
Draco heliocentric distance $75.8 \pm 0.7 \pm 5.4 \text{ kpc}$, $R=0.5 \text{ kpc}$



Threshold: 50 GeV

Background: γ -like hadronic EASs, CR electrons, diffuse extragalactic background >50 GeV

Excess $\sim 20\%$ of the background >50 GeV

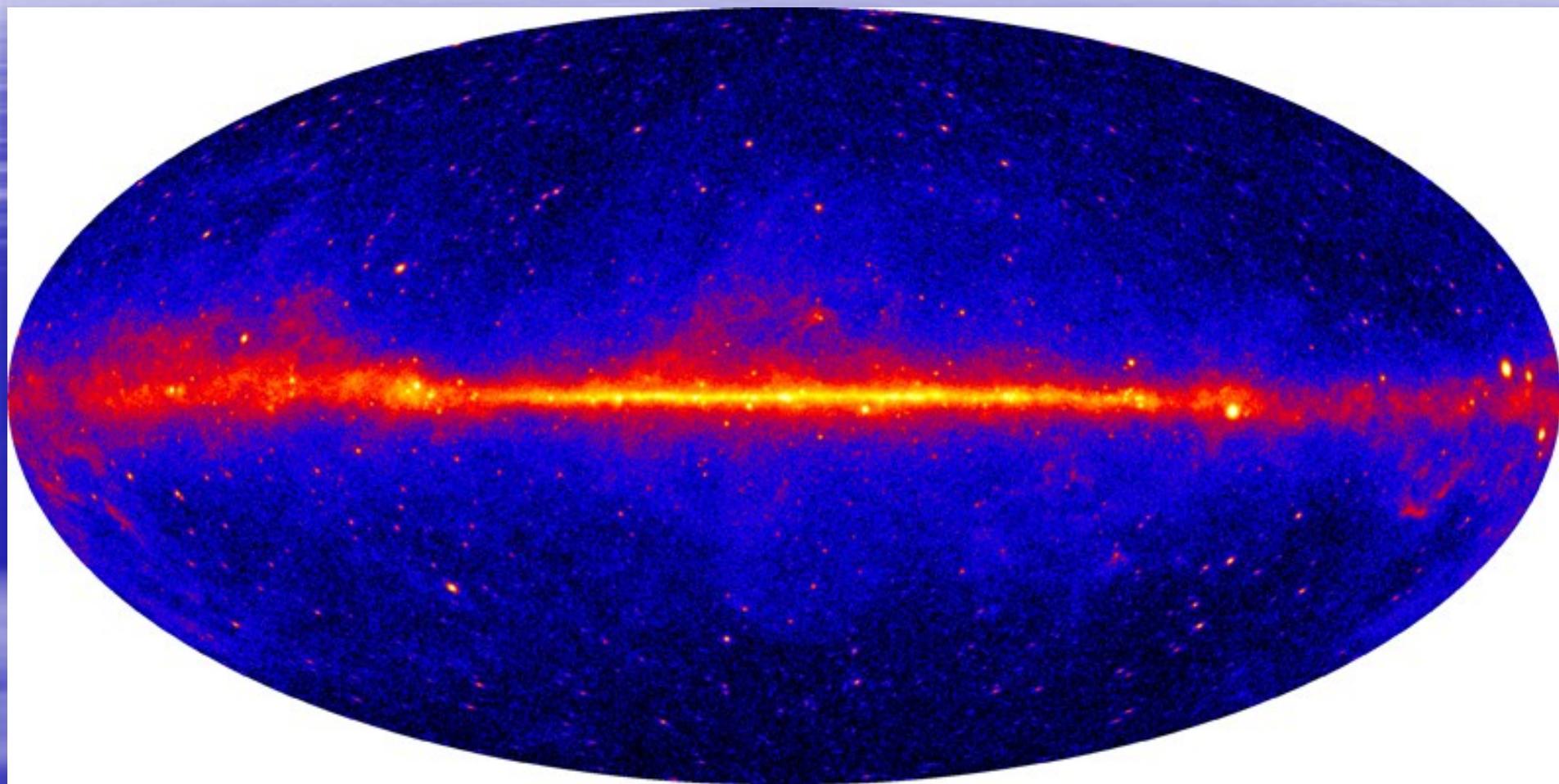


angular resolution $\Delta\theta=1^\circ$?

Profumo & Kamionkowski, 2006

Opportunity for Fermi (formerly GLAST)?

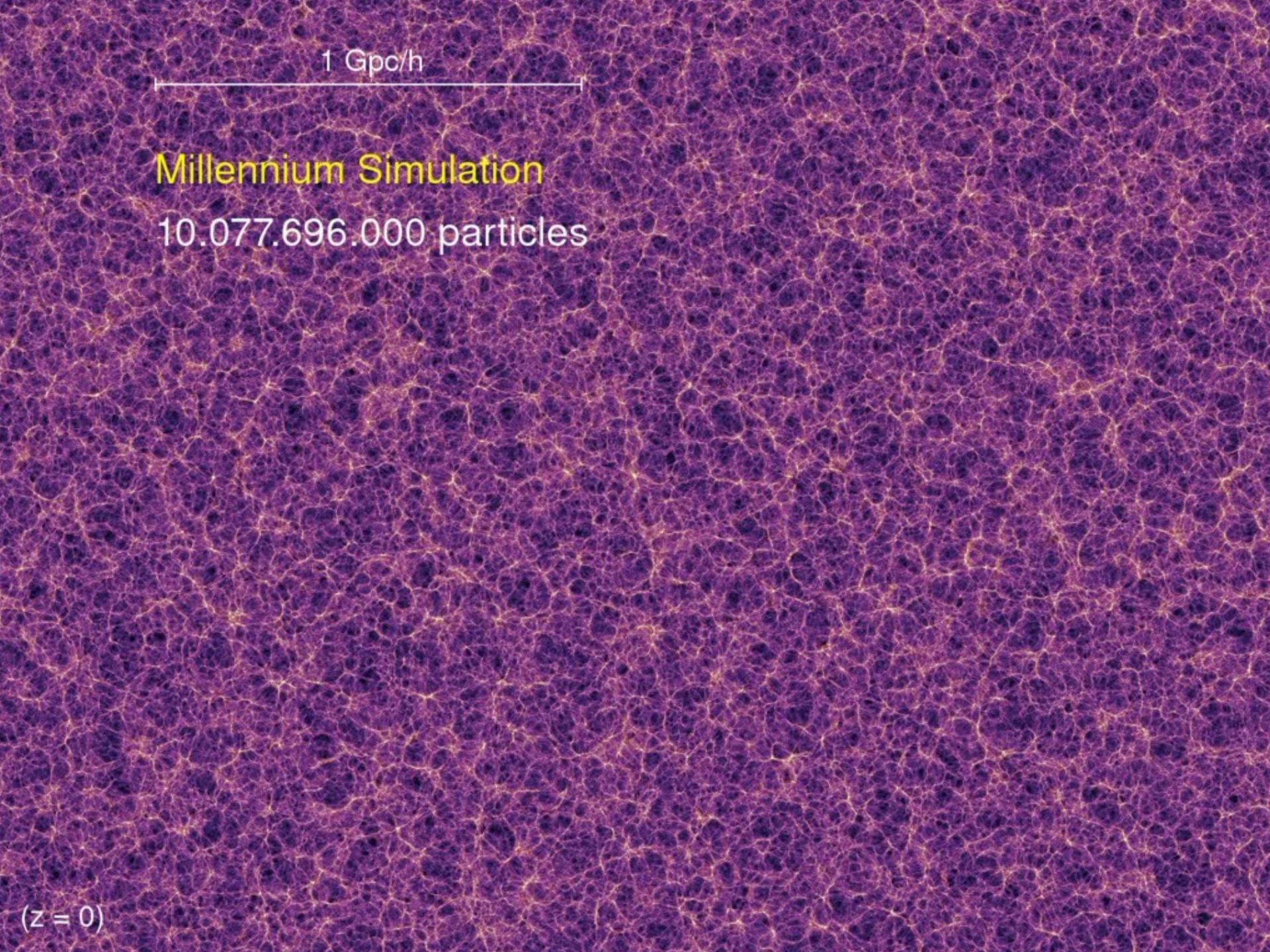
All-sky view from Fermi



<http://fermi.gsfc.nasa.gov/ssc/>



Project Columbia supercomputer (NASA)



1 Gpc/h

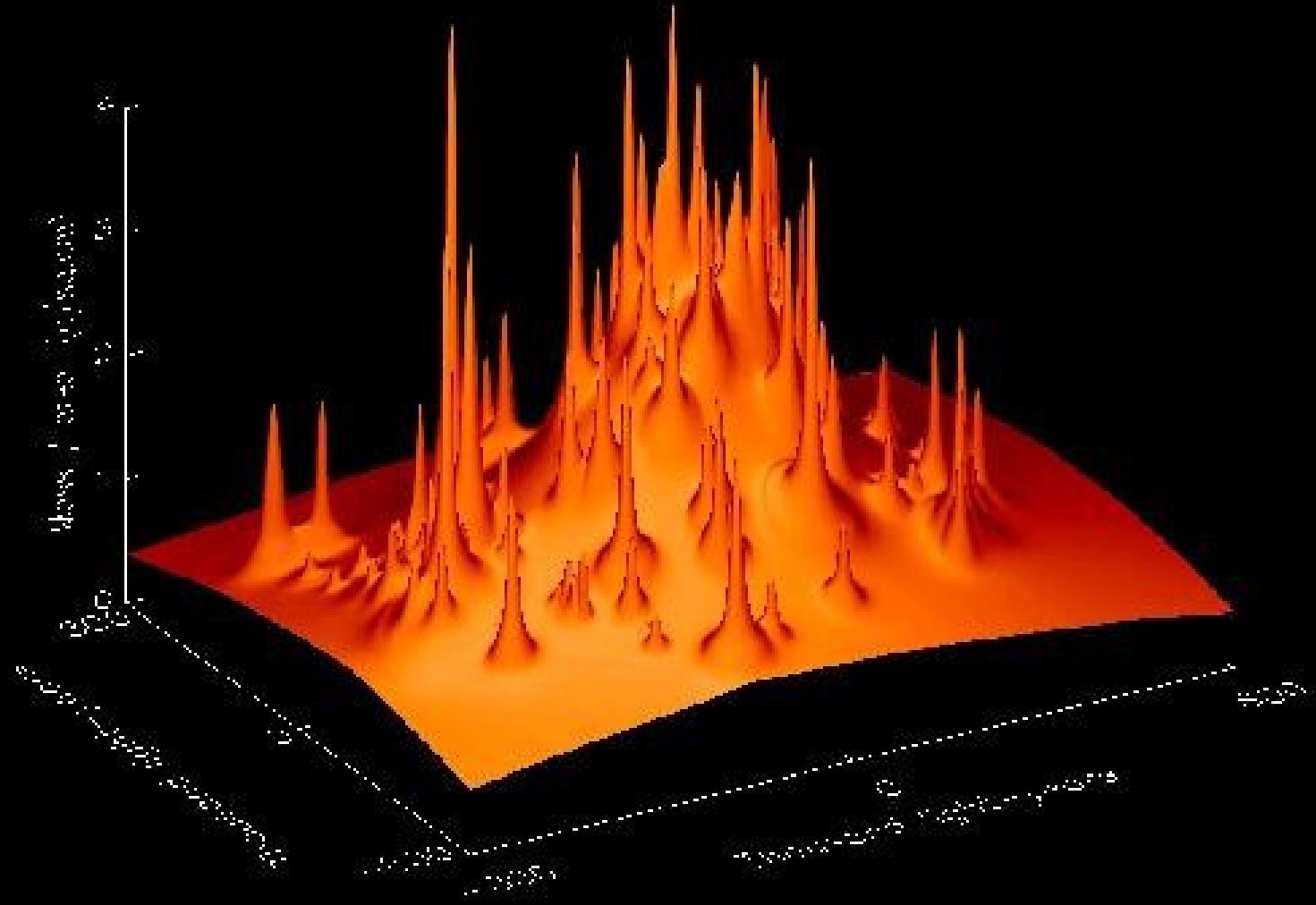
Millennium Simulation

10.077.696.000 particles

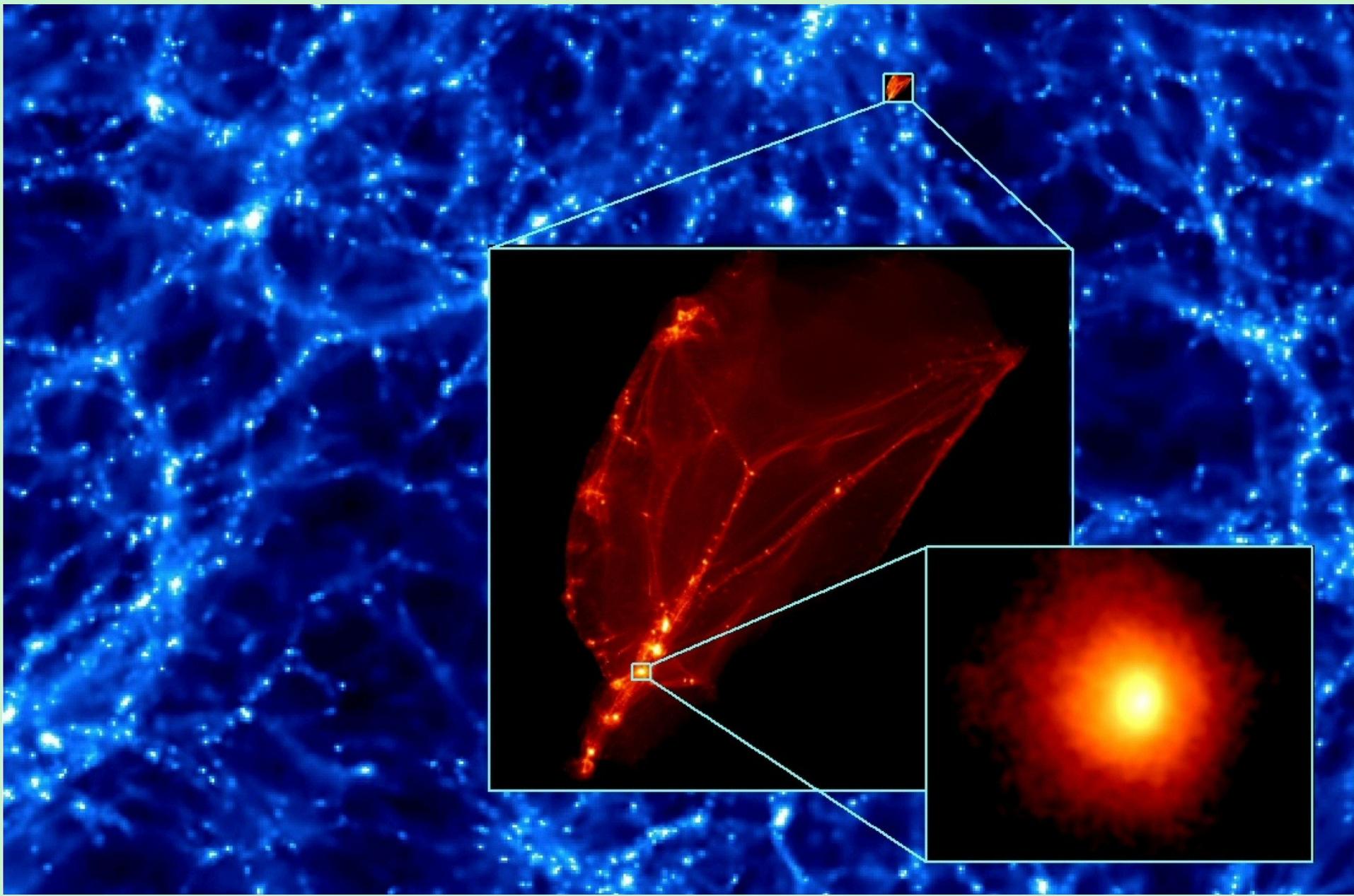
(z = 0)

$z=0.0$

80 kpc



Mass reconstruction of the cluster. Note the large, smooth distribution of (apparently invisible) matter.



$N=62 \cdot 10^6$, $m=1.2 \cdot 10^{-10} M_{\odot}$ $z=350 \rightarrow z=26$, grav. softening of 10^{-2} pc

★ Clumps of Minimal Masses

$$M_{\min} \sim 10^{-12} M_{\odot}$$

Zybin Vysotsky Gurevich 99

$$M_{\min} \sim (10^{-7} - 10^{-6}) M_{\odot}$$

Schwarz Hofmann Stocker 01

Kinetic decoupling

$$\frac{1}{\tau_{rel}} \simeq H(t), \quad n_0(\omega) = \frac{1}{2\pi^2} \frac{\omega^2}{e^{\omega/T} + 1}$$

Energy relaxation time

$$\frac{1}{\tau_{rel}} = \frac{1}{E_k} \frac{dE_k}{dt} = \frac{40}{2E_k m_\chi} \int d\Omega \int d\omega n_0(\omega) \left(\frac{d\sigma_{el}}{d\Omega} \right)_{f_L \chi} (\delta p)^2$$

$$t_d = 3 \cdot 10^{-5} \left(\frac{m_\chi}{100 \text{ GeV}} \right)^{-1/2} \left(\frac{\tilde{M}}{1 \text{ TeV}} \right)^{-2} \left(\frac{g_*}{10} \right)^{-3/4} \text{ s}$$

$$T_d = 150 \left(\frac{m_\chi}{100 \text{ GeV}} \right)^{1/4} \left(\frac{\tilde{M}}{1 \text{ TeV}} \right) \left(\frac{g_*}{10} \right)^{1/8} \text{ MeV}$$

★ Diffusion cutoff of the mass spectrum

Diffusion equation

$$\frac{\partial \delta(\vec{x}, t)}{\partial t} = \frac{D(t)}{a^2(t)} \Delta_{\vec{x}} \delta(\vec{x}, t)$$

$$\delta_{\vec{k}}(t) = \delta_{\vec{k}}(t_f) \exp \left\{ -k^2 C \tilde{M}^4 \left(t^{5/2} - t_f^{5/2} \right) \right\}$$

$$M_D = \frac{4\pi}{3} \rho_\chi(t_d) \lambda_D^3(t_d) = 10^{-13} \left(\frac{m_\chi}{100 \text{ GeV}} \right)^{-15/8} \left(\frac{\tilde{M}}{1 \text{ TeV}} \right)^{-3/2} \left(\frac{g_*}{10} \right)^{-15/16} M_\odot$$

★ Free streaming cutoff of the mass spectrum

$$\vec{x} = \vec{f}(\vec{q}, \vec{v}_d, t) = \vec{q} + \int_{t_d}^t \frac{\vec{v}(t') dt'}{a(t')} = \vec{q} + g(t) \vec{v}_d$$

$$n(\vec{x}, t) = \int d^3 v_d \phi(\vec{v}_d) \sum_{\vec{q}_*} n(\vec{q}_*, t_d) \left| \frac{D\vec{f}}{D\vec{q}} \right|_{\vec{q}=\vec{q}_*}$$

$$= \int d^3 v_d \phi(\vec{v}_d) \int d^3 q n(\vec{q}, t_d) \delta^{(3)}(\vec{x} - \vec{f}(\vec{q}, \vec{v}_d, t))$$

$$n_{\vec{k}}(t) = n_{\vec{k}}(t_d) e^{-\frac{1}{2} k^2 g^2(t) \frac{T_d}{m \chi}}, \quad g(t) = a(t_d) \int_{t_d}^t \frac{dt'}{a^2(t')}$$

★ Minimal mass of the clump

$$M_{\min} = \frac{\pi^{1/4}}{2^{19/4} 3^{1/4}} \frac{\rho_{\text{eq}}^{1/4} t_d^{3/2}}{G^{3/4}} \left(\frac{T_d}{m_\chi} \right)^{3/2} \ln^3 \left\{ \frac{24}{\pi G \rho_{\text{eq}} t_d^2} \right\}$$
$$= 10^{-8} \left(\frac{m_\chi}{100 \text{ GeV}} \right)^{-15/8} \left(\frac{\tilde{M}}{1 \text{ TeV}} \right)^{-3/2} \left(\frac{g_*}{10} \right)^{-15/16} M_\odot$$

★ Density perturbation spectrum

$$\delta(\vec{r}) = (\rho(\vec{r}) - \bar{\rho}) / \bar{\rho}$$

Power spectrum $P(k)$:

$$\langle \delta_{\vec{k}}^* \delta_{\vec{k}'} \rangle = (2\pi)^3 P(k) \delta_D^{(3)}(\vec{k} - \vec{k}'), \quad \delta_{\vec{k}} = \int \delta(\vec{r}) e^{i\vec{k}\vec{r}} d^3 r$$

Transfer function $T(k)$

$$P_{\text{eq}}(k) = P_p(k) T^2(k)$$

Moments of spectrum $P(k)$

$$\sigma_{(j)}^2 = \frac{1}{2\pi^2} \int_0^\infty k^2 dk P(k) k^{2j}, \quad \langle T_{ij} T_{ji} \rangle = \frac{2}{3} s_{(2)}^2 = \frac{2}{3} (4\pi)^2 G^2 \bar{\rho}^2 \sigma_{(0)}^2$$

Power-law spectrum $P_{\text{eq}}(k) \propto k^n$, $\sigma_{\text{eq}}(M) \propto M^{-(n+3)/6}$

Effective power-law index $n = -3 - 6 \frac{\partial \ln \sigma_{\text{eq}}(M)}{\partial \ln M}$, $n_p \approx 1$, $n_p \approx 1 \pm 0.1$

$$\sigma_{\text{eq}}(M) \simeq \frac{2 \cdot 10^{-4}}{\sqrt{f_s(\Omega_\Lambda)}} \left[\ln \left(\frac{k}{k_{\text{eq}}} \right) \right]^{3/2} \left(\frac{k}{k_{h0}} \right)^{(n_p-1)/2}$$

★ Core of a Dark Matter Clump

$$\phi(\vec{r}, t) = \phi_0 + \frac{\partial \phi}{\partial r^i} \Big|_0 r^i + \frac{1}{6} \Phi_{ll} \Big|_0 \delta_{ij} r^i r^j + \frac{1}{2} T_{ij} \Big|_0 r^i r^j + \dots$$

$$\Phi_{ij} = \frac{\partial^2 \phi(\vec{r})}{\partial r^i \partial r^j}, \quad T_{ij} = \Phi_{ij} - \frac{1}{3} \Phi_{ll} \delta_{ij}$$

Peak height

$$\nu = \delta_{\text{eq}} / \sigma_{\text{eq}}(M)$$

Tidal velocity

$$\frac{dv_{tid,i}}{dt} = -T_{ij}(t) r^j$$

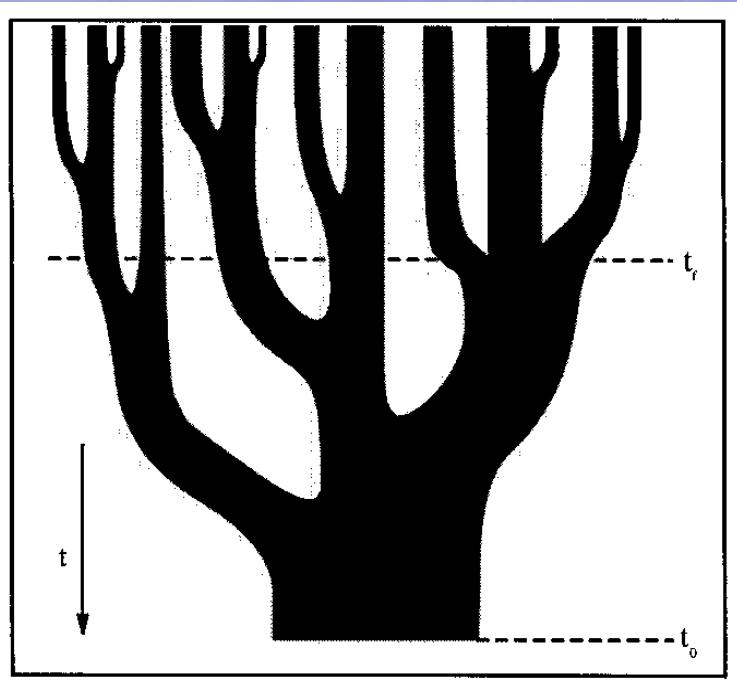
$$\boxed{\Delta E \simeq \Delta V}$$

$$\Delta E \simeq \int d^3r \rho_{\text{int}}(t_s) v_{tid}^2(t_s)/2, \quad \Delta V \simeq \frac{GM M_o}{R}$$

Clump core radius

$$\frac{R_c}{R} \simeq \frac{\pi 2^{5/3} 3^{13/3}}{5^3} G \rho_{\text{eq}} t_{\text{eq}}^2 \frac{f^2}{\nu^2}(\delta_{\text{eq}}) \simeq 0.3 \nu^{-2} f^2(\delta_{\text{eq}})$$

Press-Schechter formalism:



‘Merger tree’

Mass function of unconfined clumps

$$\xi_{\text{PS}}(t) \frac{dM}{M} = \frac{2\delta_c}{\sqrt{2\pi}\sigma_{\text{eq}}^2 D(t)} \frac{d\sigma_{\text{eq}}}{dM} \exp\left[-\frac{\delta_c^2}{2\sigma_{\text{eq}}^2 D^2(t)}\right] dM,$$

$$\delta_c = 3(12\pi)^{2/3}/20 \simeq 1.686$$

HIERARHICAL STRUCTURING

Press & Schechter 1974
Lacey & Cole 1993

★ Tidal destruction of clumps in hierarchical model

Clump destruction at

$$\Delta E \geq |E| \sim GM^2/2R, \quad \delta(M, t_f) = \delta_c$$

Number density of unconfined (free) clumps

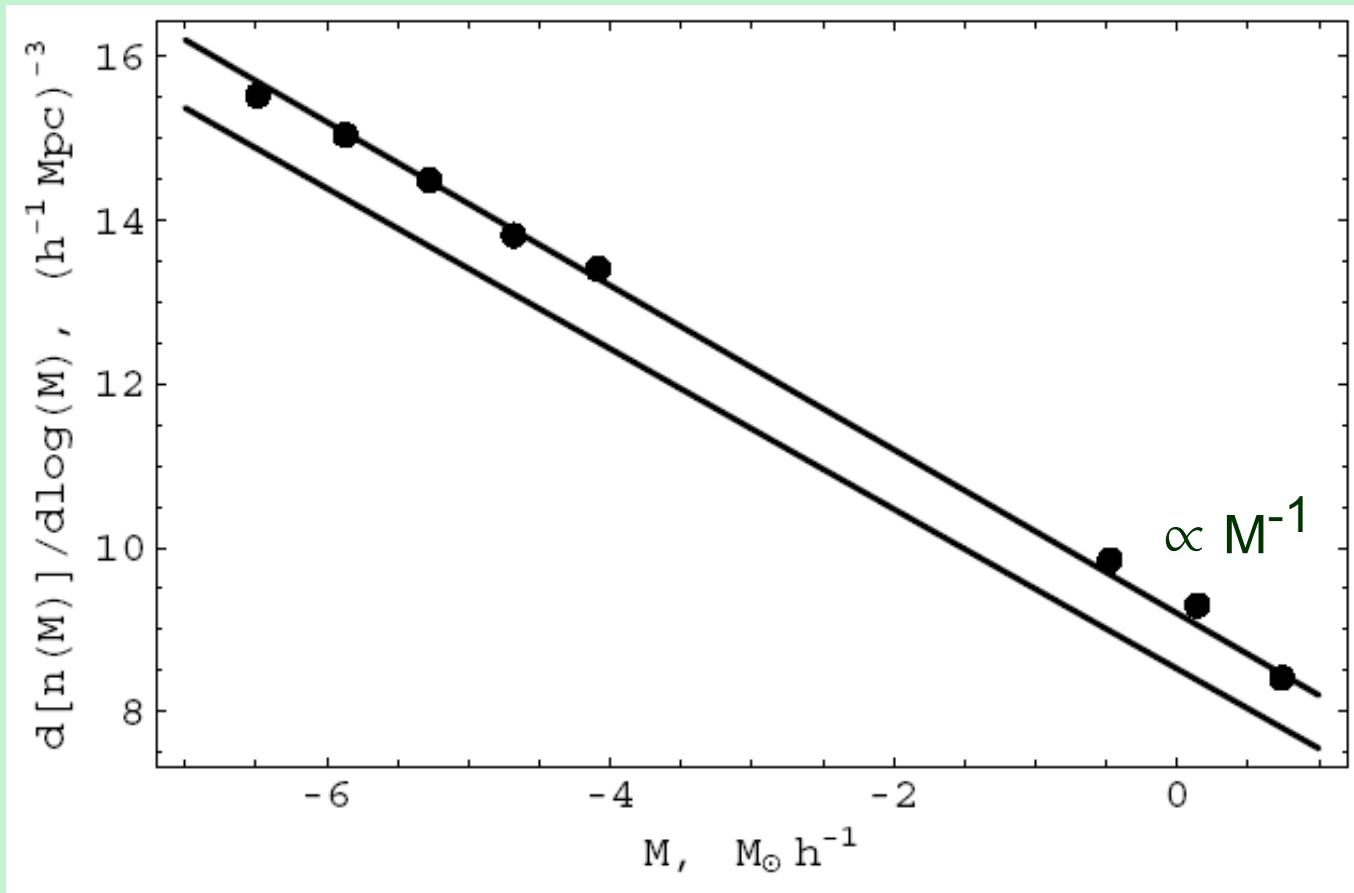
$$\phi_{PS} = \left(\frac{2}{\pi}\right)^{1/2} \frac{\rho}{M} \frac{\delta_c}{D_g(t) \sigma_{eq}^2} \frac{d\sigma_{eq}}{dM} \exp\left[\frac{-\delta_c^2}{2D_g(t)^2 \sigma_{eq}^2}\right] dM$$

Press, Shechter, 1974

Energy increase during one fly-by

$$\Delta E = \frac{1}{2} \int d^3r \rho(r) (v_x - \tilde{v}_x)^2$$

Mass function of small-scale DM clumps



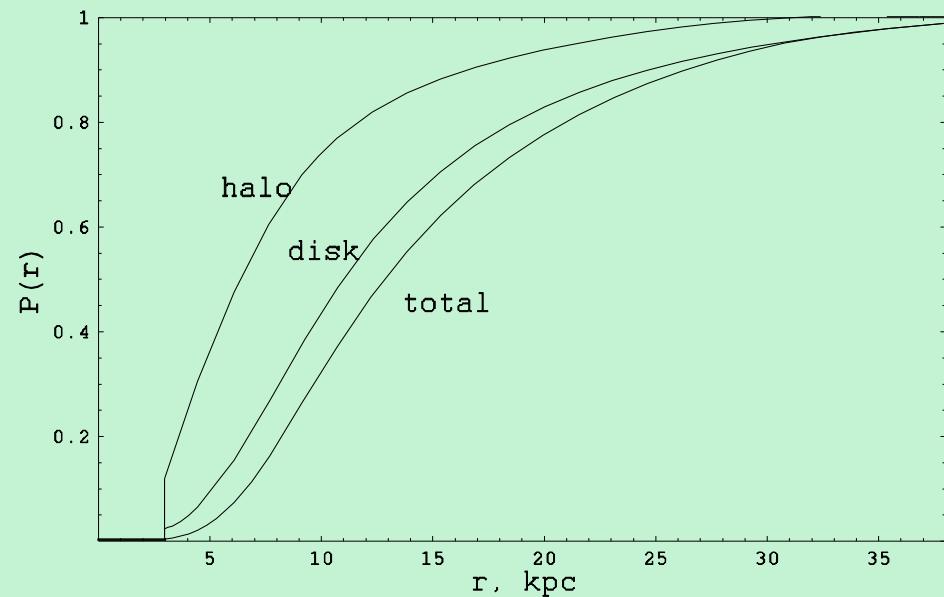
Lower line – model calculation

Berezinsky, Dokuchaev & Eroshenko 2003

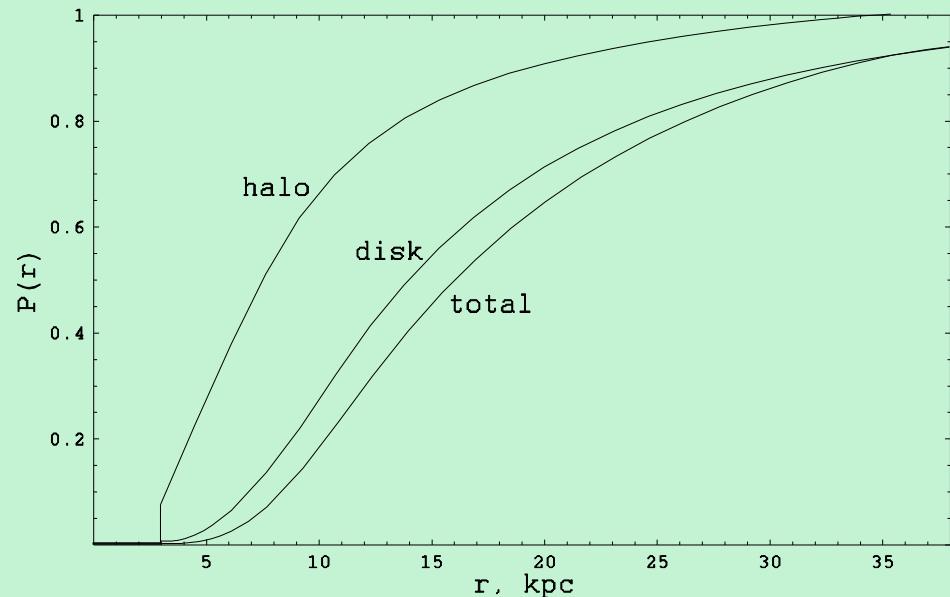
Dots - numerical simulations

Diemand, Moore & Stadel 2005

Fraction of survived clumps



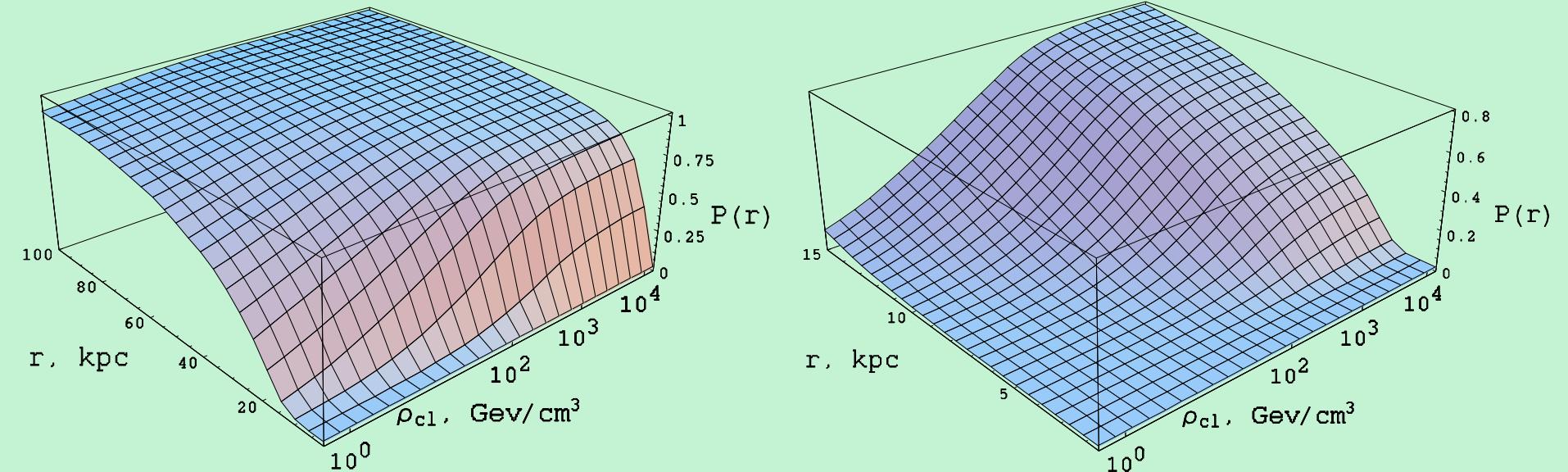
$$M_{\text{cl}} = 10^{-8} M_{\odot}$$



$$M_{\text{cl}} = 10^{-3} M_{\odot}$$

Tidal destruction of clumps in the Galactic bulge, disc, halo and the resulting total fraction of survived clumps with $M_{\text{cl}}=10^{-8}M_{\odot}$ and $10^{-3}M_{\odot}$

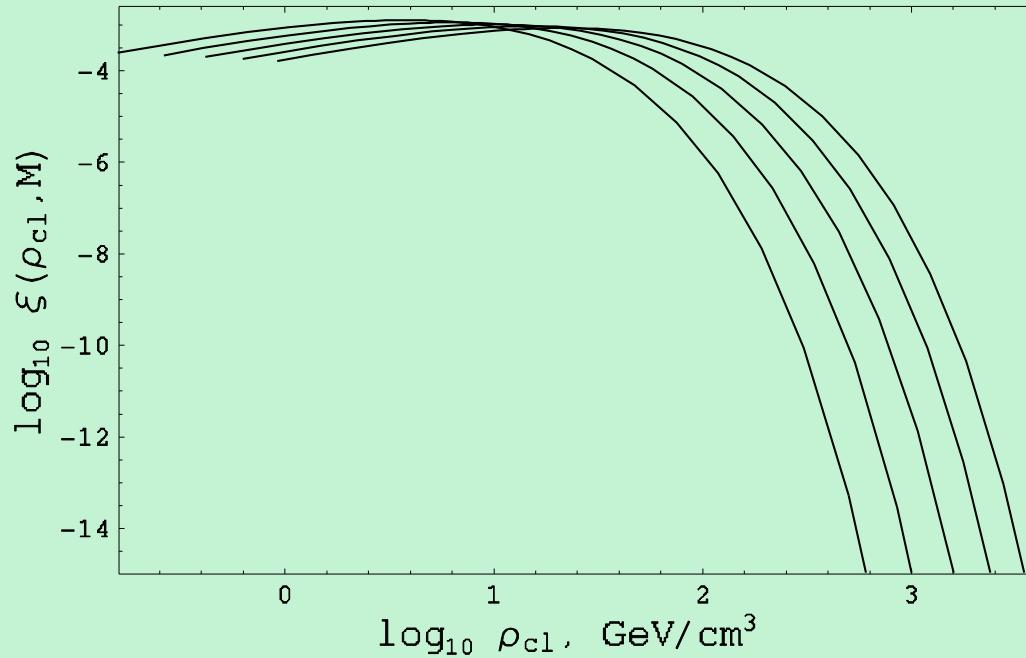
Survived fraction of clumps $P(r)$ in the Galaxy



Radial galactocentric distance r in kpc

Mean internal density of clump ρ_{cl} in GeV/cm^3

Clump distribution function



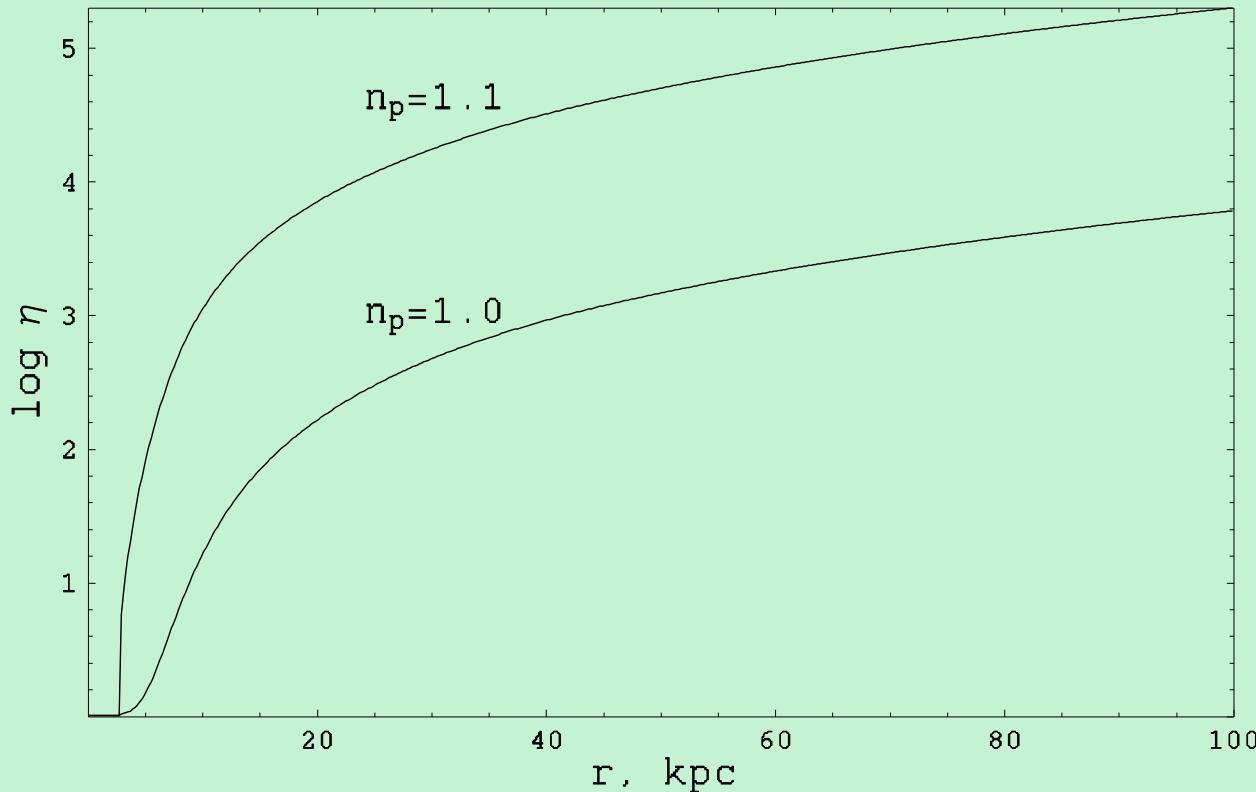
$$\xi(M, \nu) \frac{dM}{M} d\nu \simeq \frac{y(\nu)}{(2\pi)^{1/2}} e^{-\nu^2/2} \frac{d \log \sigma(M)}{dM} dM d\nu$$

Mean internal density of clump ρ_{cl} in GeV/cm^3

Right sides of curves correspond to

$M_{\text{cl}} / M_{\odot} = 10^{-8}, 10^{-6}, 10^{-4}, 10^{-2}, 1$ from up to down

Annihilation enhancement in clumps



**Local annihilation enhancement factor (boosting) $n(r)$
in isothermal spherical symmetric halo**

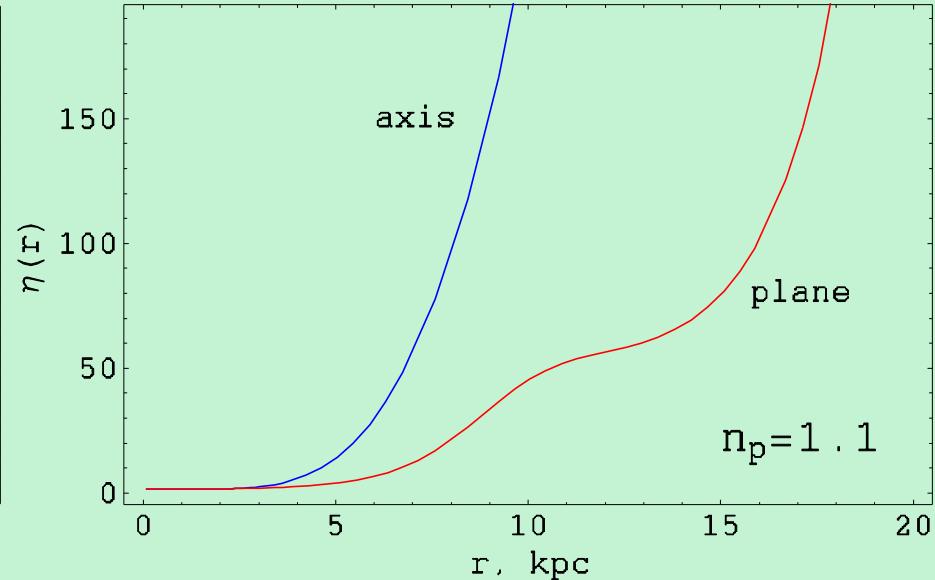
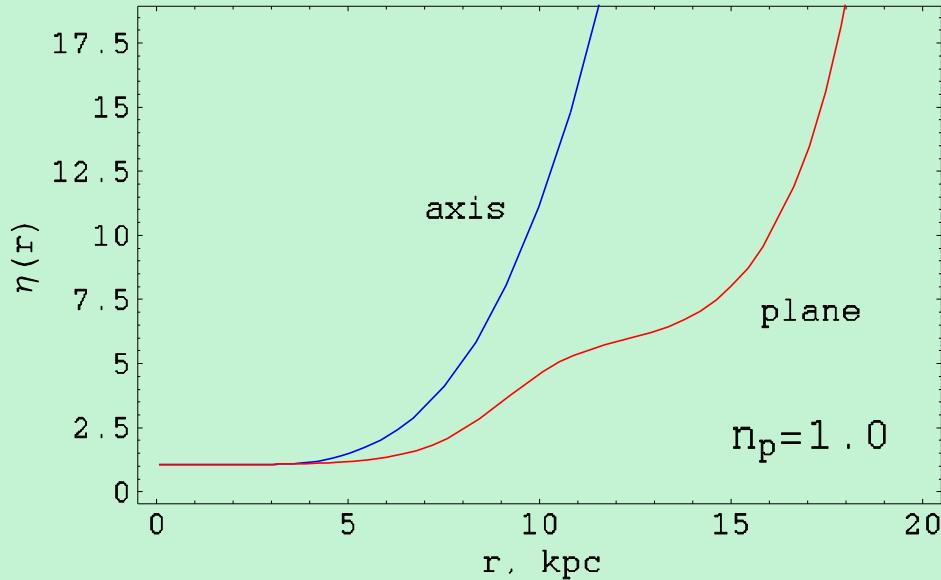
$$dI_{\text{dif}} = \frac{\langle \sigma_{\text{ann}} v \rangle}{4\pi l^2} \frac{\rho_{\text{DM}}^2(r)}{m_\chi^2} d^3 r$$

$$\bar{\rho}_{\text{cl},\text{P}}(r, n_p) = \frac{S(x_c, \beta)}{\xi_{\text{P}}(r, n_p)} \int P(r, \rho_{\text{cl}}) \rho_{\text{cl}} \xi(\rho_{\text{cl}}) d\rho_{\text{cl}}$$

$$\eta(r) = \frac{I_{\text{cl}} + I_{\text{dif}}}{I_{\text{dif}}} = 1 + \xi_{\text{P}}(r, n_p) \frac{\bar{\rho}_{\text{cl},\text{P}}(r, n_p)}{\rho_\chi(r)}$$

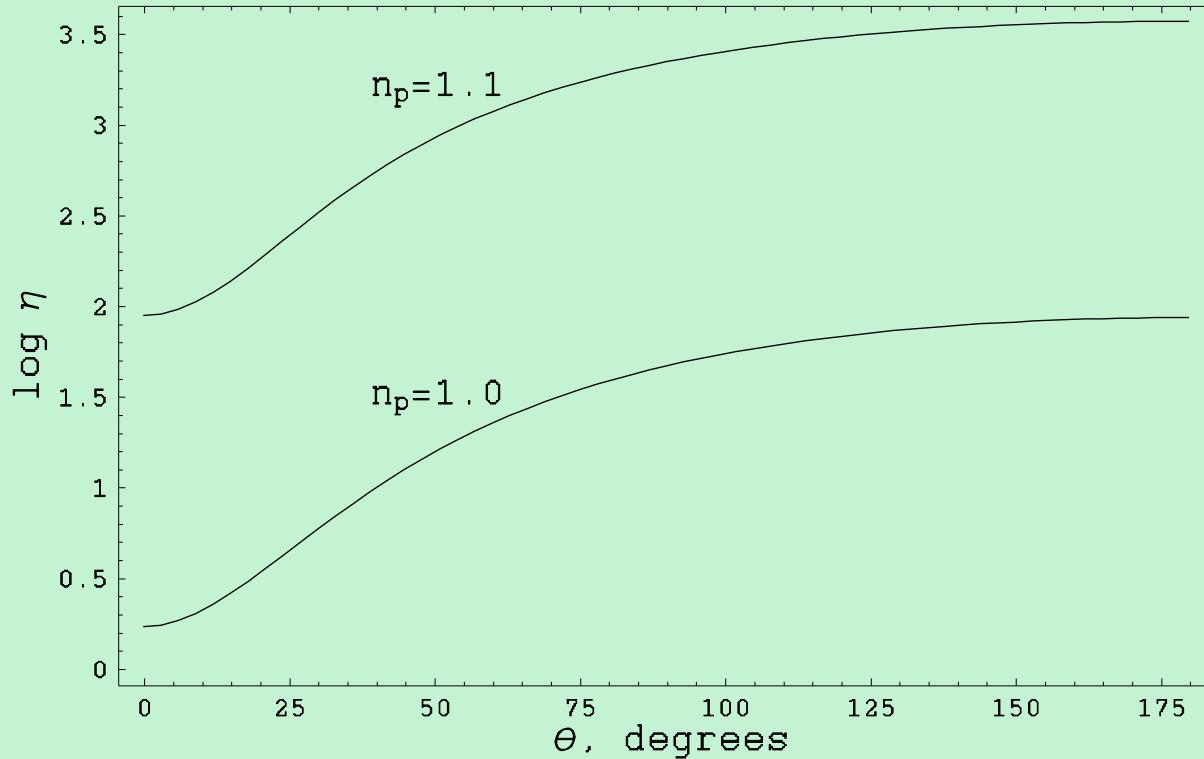
$$\xi_{\text{P}}(r, n_p) = \int P(r, \rho_{\text{cl}}) \xi(\rho_{\text{cl}}) d\rho_{\text{cl}},$$

Local boosting in the Galactic plane and along z-axis



Local annihilation enhancement factor (boosting) $\eta(r)$
in NFW halo with ring at $r=14$ kpc

Integrated along the line of sight (observed) enhancement factor $\eta(\theta)$



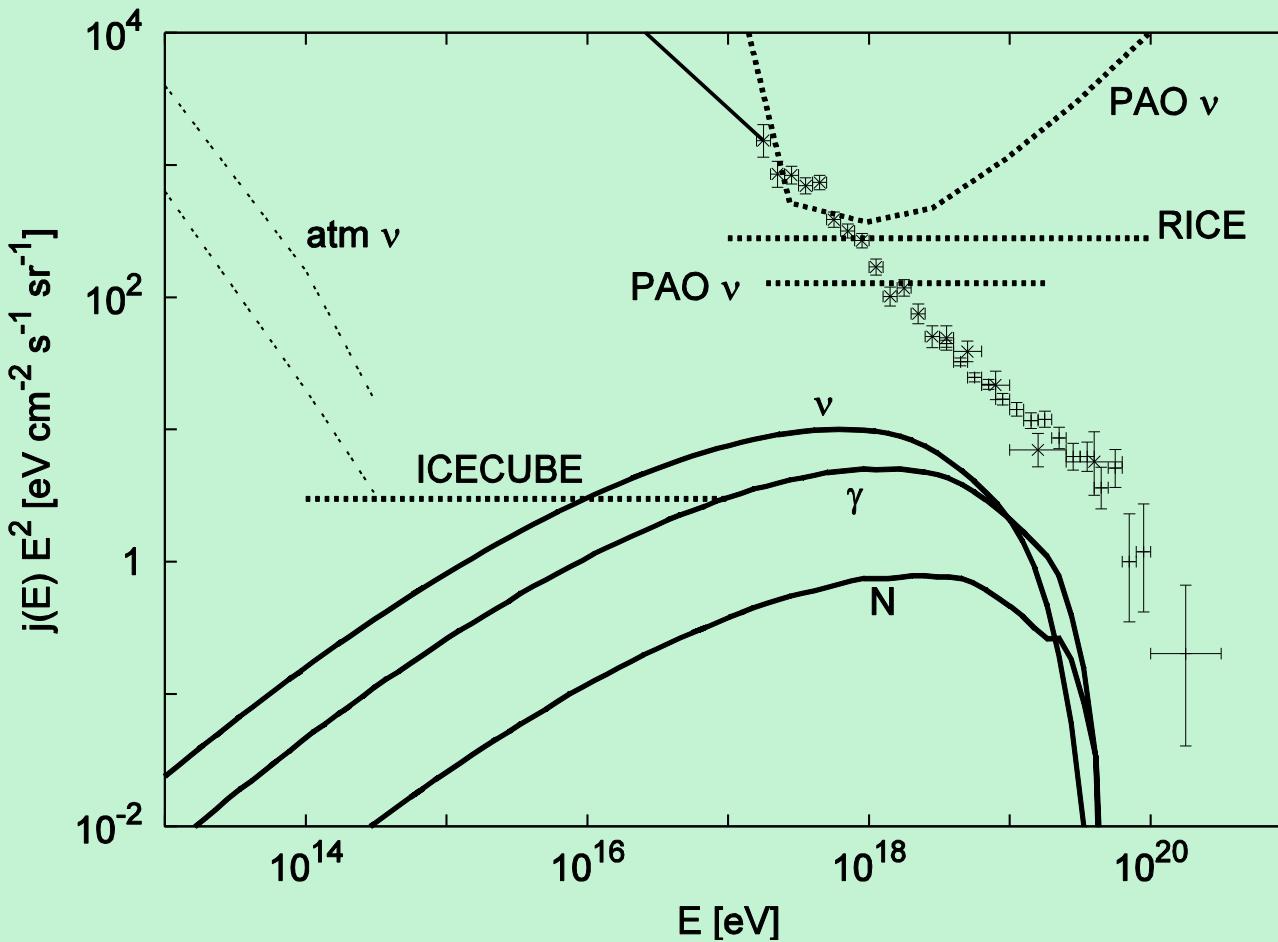
The case of isothermal spherically symmetric halo model

$$\eta(\theta, \phi) = 1 + \frac{\int \rho_\chi(r) dr \xi_P(r) \bar{\rho}_{cl,P}(r, n_p)}{\int \rho_\chi^2(r) dr}$$

$$\bar{\rho}_{cl,P}(r, n_p) = \frac{S(x_c, \beta)}{\xi_P(r, n_p)} \int P(r, \rho_{cl}) \rho_{cl} \xi(\rho_{cl}) d\rho_{cl}$$

$$\xi_P(r, n_p) = \int P(r, \rho_{cl}) \xi(\rho_{cl}) d\rho_{cl},$$

Annihilations of superheavy DM in superdense clumps



Fluxes $I_i(E)$ of photons, nucleons and neutrinos from neutralino annihilations in the Galactic halo for neutralino with $m_\chi \sim 10^{11}$ GeV

- ★ Clumps with M_{\min} give the dominant contribution to DM annihilation

In the case of $n_p = 1$, $\nu \sim 2.5$

$$M_{\min} \sim 10^{-8} M_{\odot}$$

$$R \simeq 3.6 \cdot 10^{15} \text{ cm}, \quad R_c \simeq 1.8 \cdot 10^{14} \text{ cm}$$

$$\bar{\rho}_{\text{int}} \simeq 2.5 \cdot 10^{-22} \text{ g cm}^{-3}$$

Halo mass fraction in these clumps

$$\xi_{\text{int}} \sim 0.002$$

Mean number density in the halo

$$n_{\text{cl}} \sim 25 \text{ pc}^{-3}$$

DM annihilation enhancement

$$\eta \sim 10 - 10^2$$