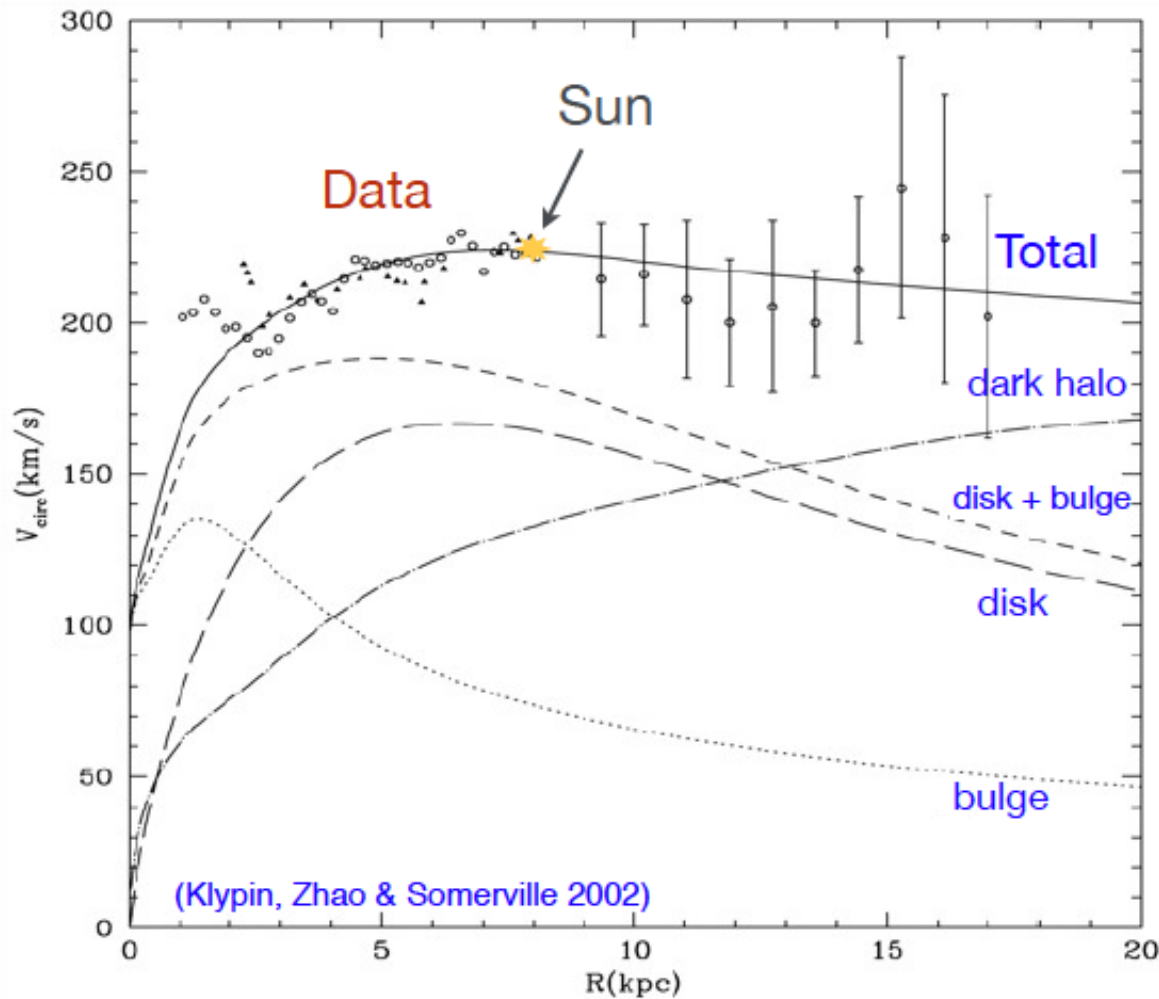


Direct detection of DM particles

Introduction to lab visit
8.5.2013

Christian Regenfus
(regenfus@cern.ch)

Density of WIMPs in the Milky Way



PDG:

$$\rho_{\text{halo}} = 0.1 - 0.7 \text{ GeVcm}^{-3}$$

$$\rho_{\text{disk}} = 2 - 7 \text{ GeVcm}^{-3}$$

Standard value:

$$\rho_{\chi} \sim 0.3 \text{ GeVcm}^{-3}$$

$$\rho_{\chi} \sim 3000 \text{ WIMPs m}^{-3}$$

$$(M_{\text{WIMP}} = 100 \text{ GeV})$$

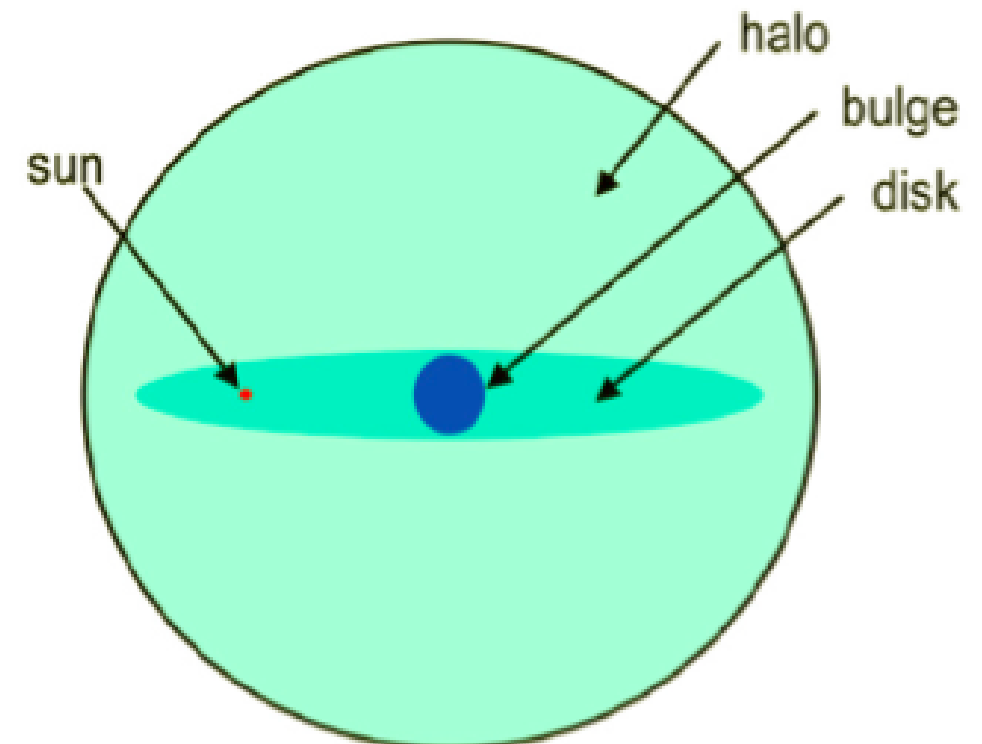
$$\text{WIMP flux on earth} \sim 10^5 \text{ cm}^{-2}\text{s}^{-1}$$

Standard halo assumptions:

- Sun moves around the galactic centre with $v = 220 \text{ km/s}$

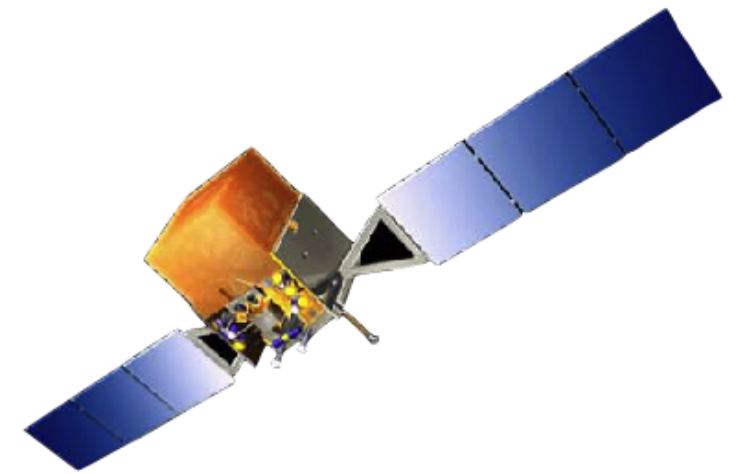
Isothermal WIMP Halo:

- On average halo is stationary, no Co-rotation around the galactic centre
- WIMP velocities are Maxwell-Boltzmann distributed
- $v_0 = 220 \text{ km/s}$
- $v_{\text{max}} = v_{\text{esc}} = 650 \text{ km/s}$

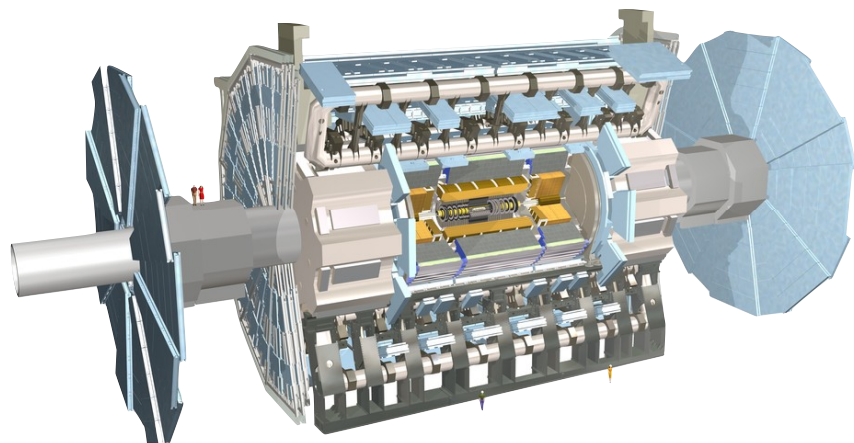


Strategies for WIMP Detection

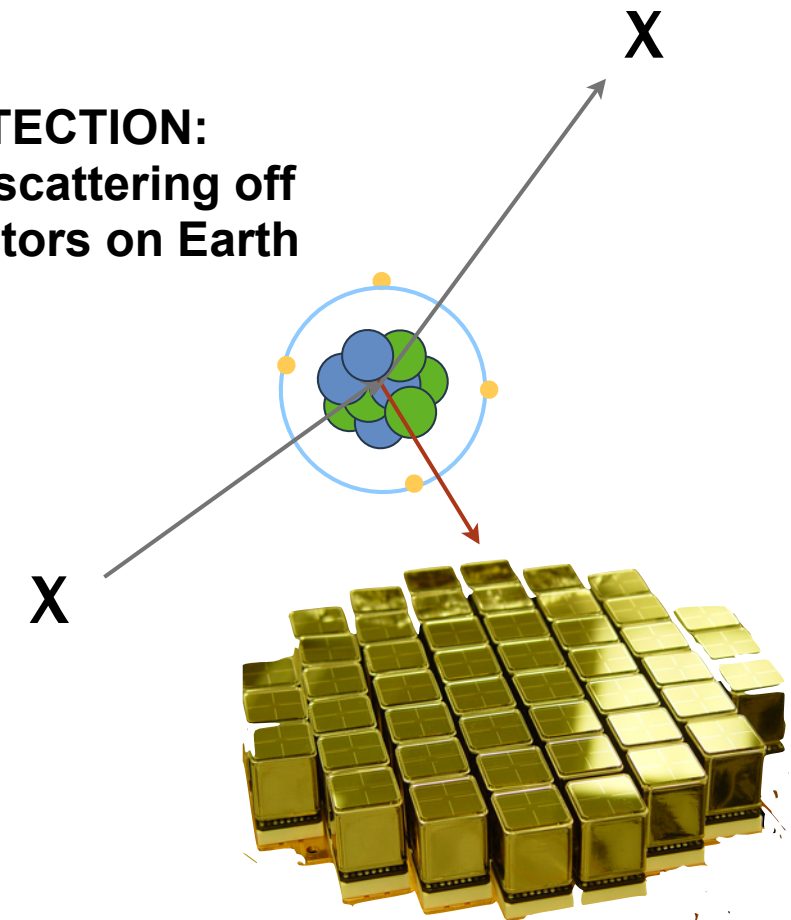
INDIRECT DETECTION: measure gamma rays, neutrinos, positrons, antiprotons, anti-deuterons, etc. from WIMP annihilation in GC, in Sun, in MW



PARTICLE COLLIDERS:
Produce (and detect) WIMPs



DIRECT DETECTION:
measure WIMP scattering off targets in detectors on Earth

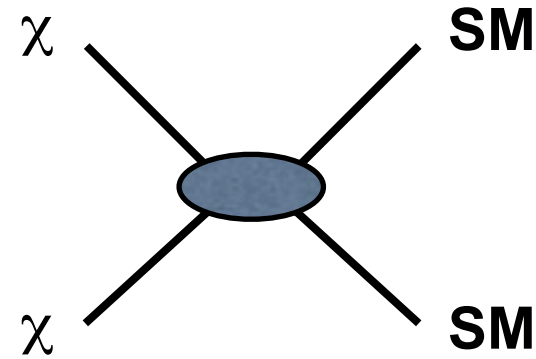
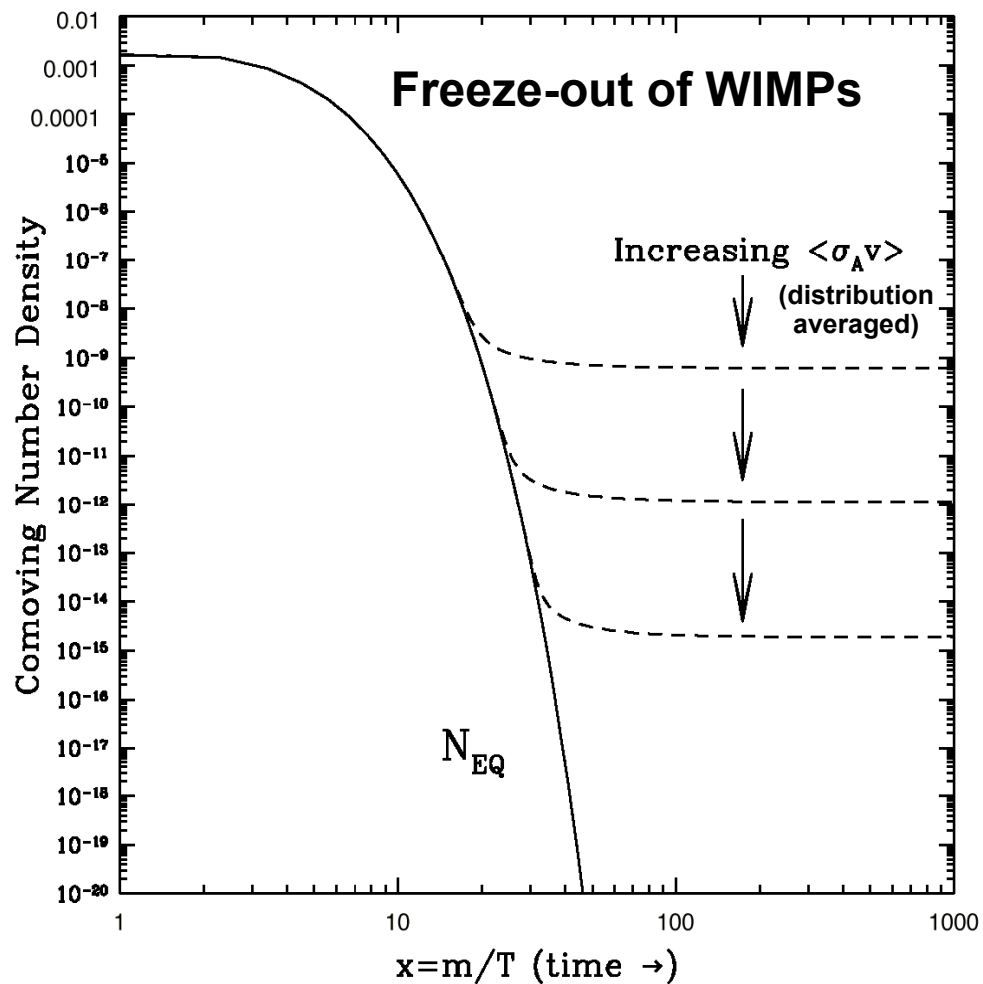


WIMP models are stringently probed by one or more of these methods

We look for direct interactions of WIMPS

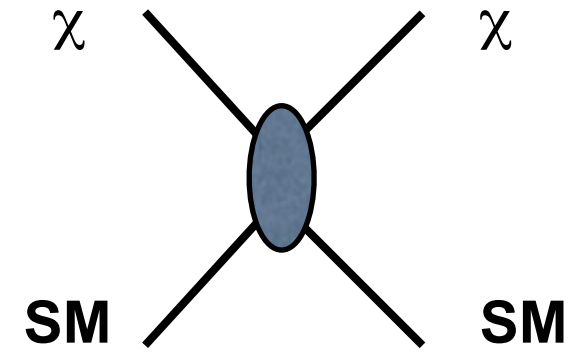
WIMP - most prominent DM candidate

Thermal (BB) relic's number density



annihilation

s-channel



scattering

t-channel

coherent weak or spin dependent IA

The observed relic density points to the **weak scale!**

Well motivated from SUSY

LSP neutralino: super partner of the gauge bosons
(need exact R parity conservation)

We don't know a massive particle which does not interact weakly !

Principle of Direct Detection

Non rel. elastic collisions with nuclei => nuclear recoils

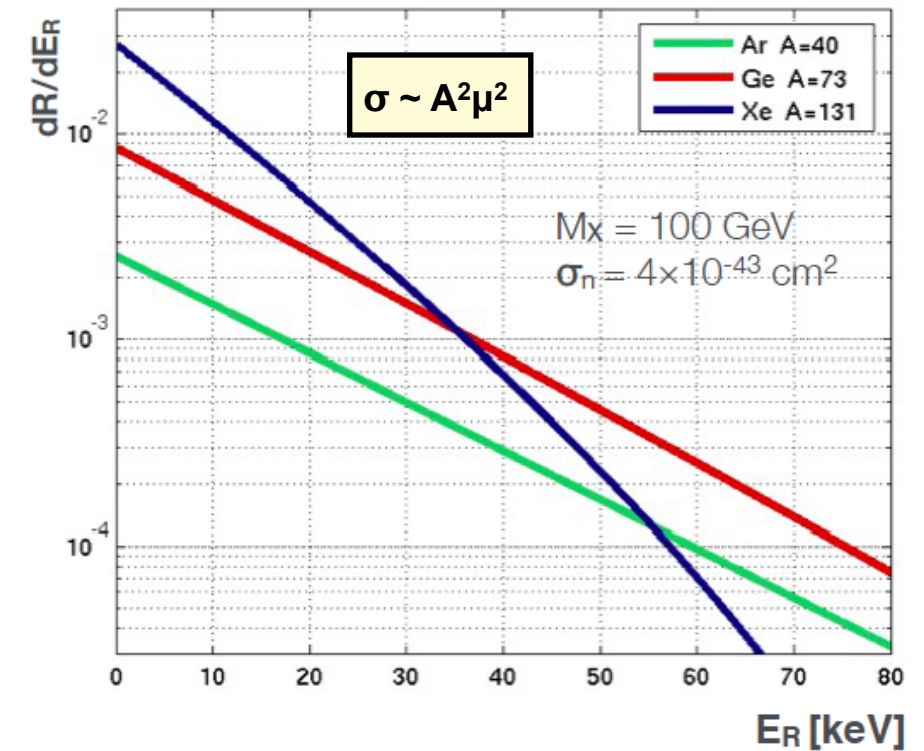
Coherent scattering of WIMPs (1985)

Goodman, M. W. and Witten, E.: 1985, Phys. Rev. D31, 3059.

Detectability of certain dark-matter candidates, weak neutrino-like couplings

- Recoil energies 0....~100 keV
- A^2 & $F(q)$ Dependence: consistency test with different targets
- Expected count rate: some events / kg year
- Time dependence of the signal:
annual: ~3 % , diurnal: ~40 %
- Combination of different A can test $A^2\mu^2$ dependence

Differential rates at different targets



Ar: ~15 events/100kg/year @Thr. = 30keV

Comparison with neutrinos:

Leptonic IA: (MeV) neutrinos (O) $\sigma \sim 10^{-45} \text{ cm}^2$

Nuclear IA: CNS (Coherent neutrino-nucleus scattering)

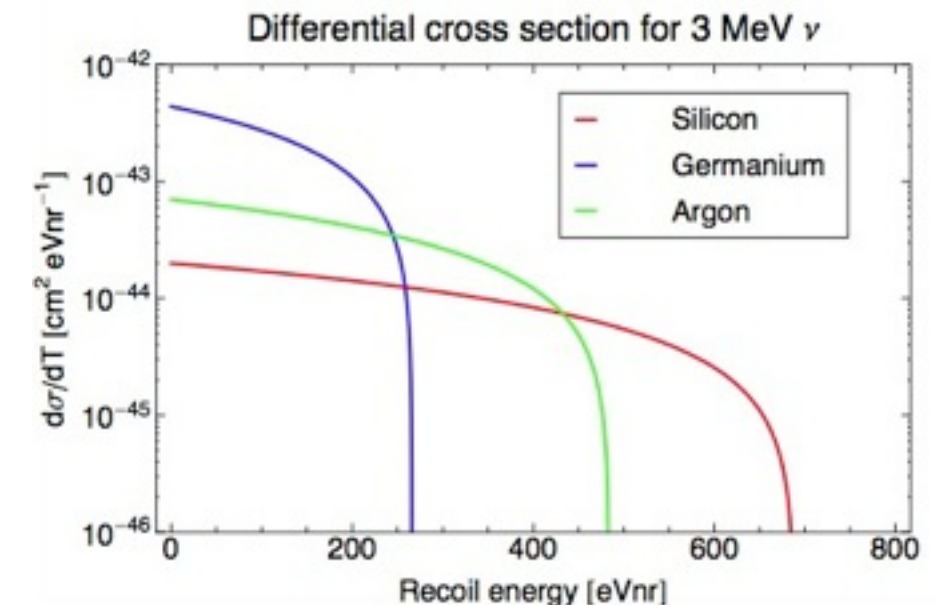
Cross sections for neutrino-nucleon scatterings are not so well known as for leptonic reactions (poor theoretical knowledge of the nucleon form factors)

=> Very small recoil energies

Need large and sensitive detectors with good background suppression

Principles and applications of a neutral-current detector for neutrino physics...,
Drukier and Stodolsky PRD 30, 11,(1984) 2295

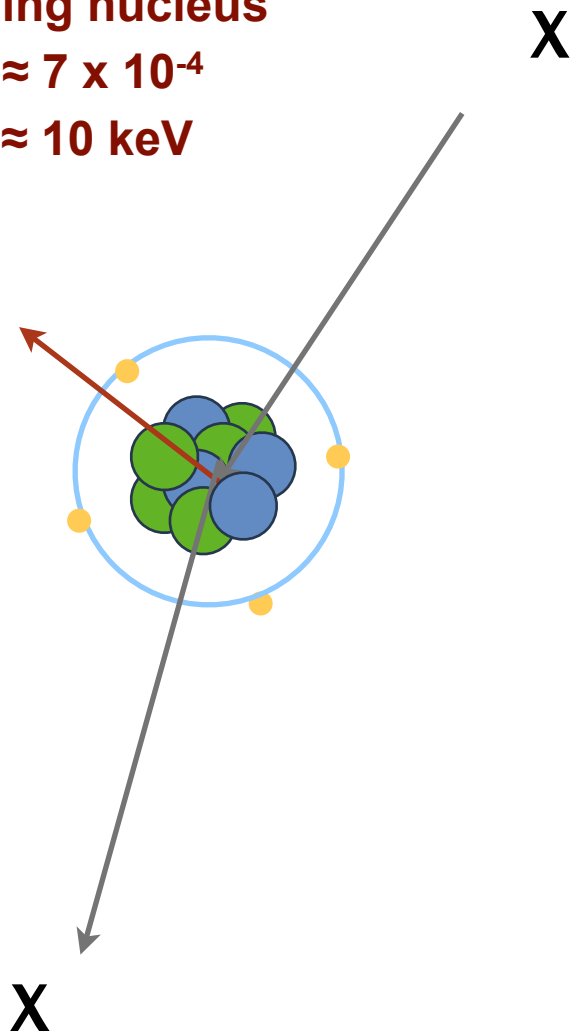
Coherent effects of a weak neutral current,
D. Freedman, PRD 9, (1974) 1389



Detection of WIMPs: Signal and Background

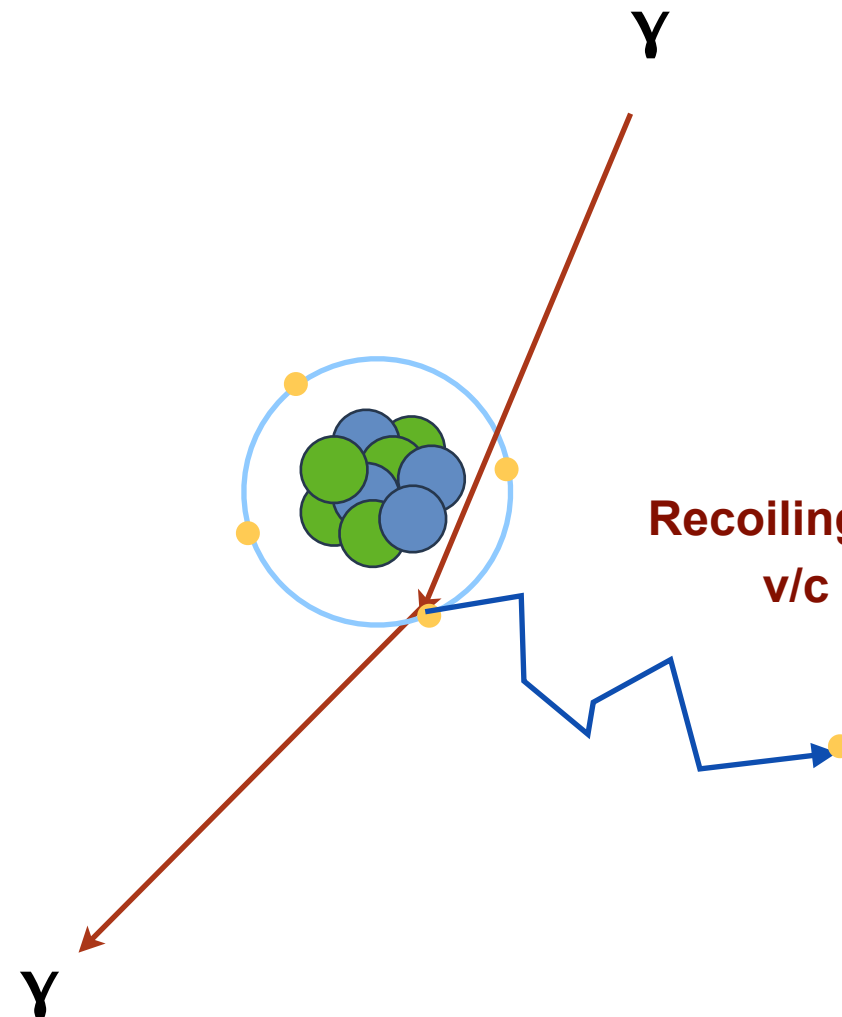
Signal (WIMPs, n)

Recoiling nucleus
 $v/c \approx 7 \times 10^{-4}$
 $E_R \approx 10 \text{ keV}$



Background (γ, β)

Recoiling electron
 $v/c \approx 0.3$



The detector material (target) converts this energy into something measurable
scintillation light - free charge (electrons) - phonons - gas bubbles - ionisation grains

Experiments are distinguished by their ability to detect

- interaction type (integral vs. event-by-event)
- single or multiple channel read out (light, charge, heat ...)

Direct detection techniques

Example listing not complete!
Lots of activities around the world

Bolometers

Targets: Ge, Si, Al₂O₃, LiF, TeO₂
E_{phonon} ~meV, NR energy col. eff. ~100%
Sensitivity (TES) << 1 keV
FWHM 4.5 eV @ 6 keV γ
CRESST-I (0.6 keV), CUORICINO, CUORE (5 keV)

Heat

Single channel techniques

Hybrid techniques

Heat & Ionisation

Bolometers

ZIP/NTD for Q & H channels
Targets: Ge, Si
CDMS, EDELWEISS, SCDMS, EURECA cryogenic (<50 mK)

Scintillation & Heat Bolometers

Targets: CaWO₄, BGO, Al₂O₃
TES/NTD for L & H channels **CRESST, ROSEBUD** even more cryogenic (~10 mK)

Combination of more detection principles allows for better identification of the interaction type.

Helps enormously to suppress background.

Event-by-event interaction type ID present in all modern DM detectors.

Ionisation Detectors

Targets: Ge, Si, CF₃I, CS₂
Energy per e/h pair 1-5 eV
NR energy collection eff. 10-30%
Thresh. ~3 keV
IGEX, HDMS, COUPP, GENIUS, DRIFT, CoGeNT

Ionisation

Scintillation

Scintillators

Targets: NaI, Xe, Ar, Ne E. per photon ~15 eV, NR energy collection eff. 1-3%,
Light sens. 2-8 phe/keV, Thresh. ~2-5 keV
ZEPLIN I (2 keV), NAIAD (4 keV) DAMA (2 keV), DEAP, CLEAN, XMASS (5 keV)

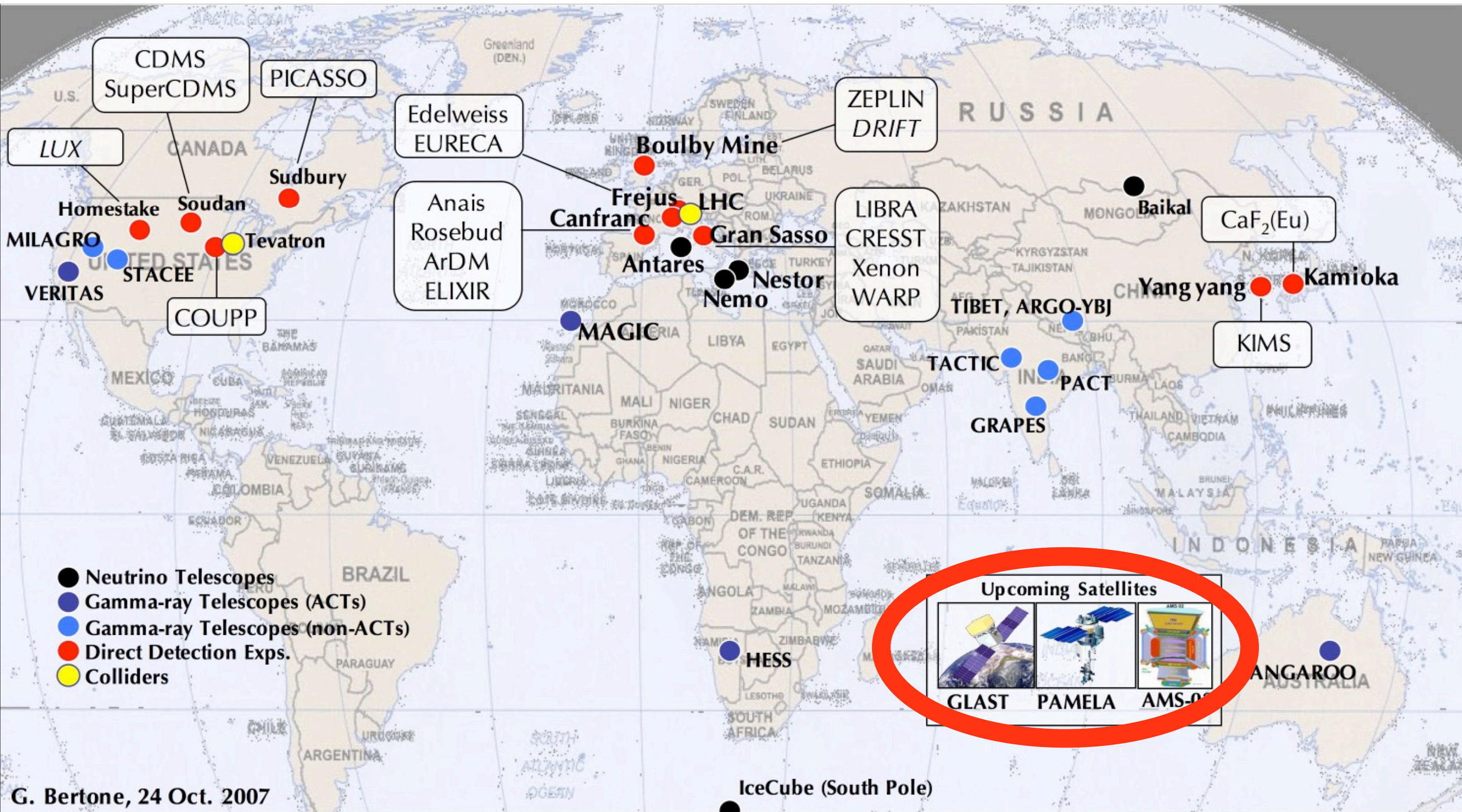
Scintillation & Ionisation Detectors

Noble liquid TPCs, PMT readout (both channels), Thresh. ~5 - 30 keV
ZEPLIN, XENON, LUX, DarkSide, WARP, ArDM, SIGN
(Best scalability, presently best limits)

In addition:

Scintillation Timing (DEAP/CLEAN)
Signal Modulation (DAMA/LIBRA, DRIFT ...)
Nucl.-recoil-only trigger (COUPP...)
Self-Shielding (XMASS) ...

DM search around the world

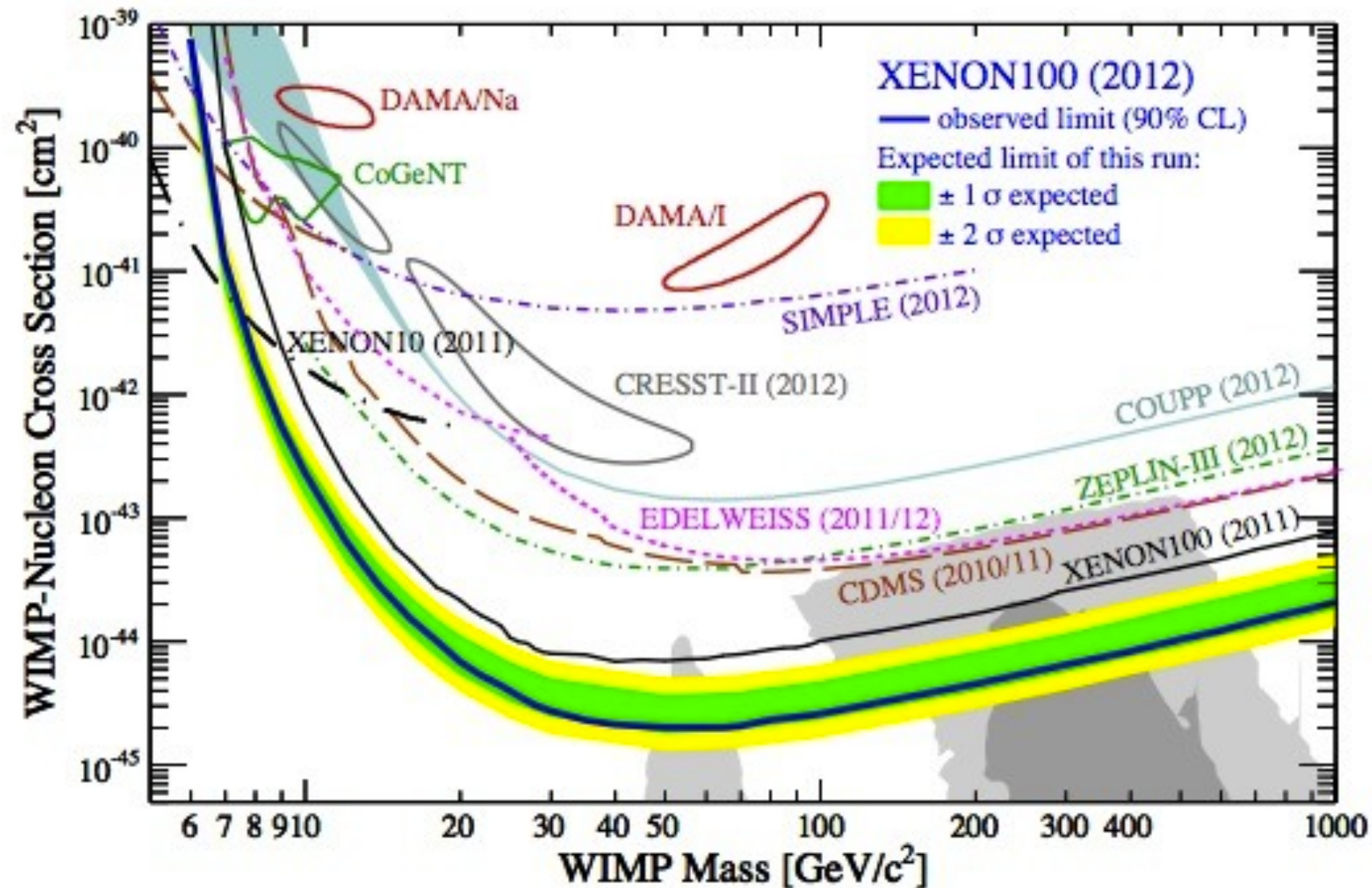


Best limits up to date

Example: XENON100

arXiv:1207.5988v1 [astro-ph.CO] 25 Jul 2012

results from 225 live days



Latest spin-independent WIMP-nucleon scattering limits from XENON100: The expected sensitivity of this data is shown by the green/yellow band (1σ/2σ) and the resulting exclusion limit (90% CL) in blue.

How do we build such a detector? (developments towards a ton scale LAr detector)

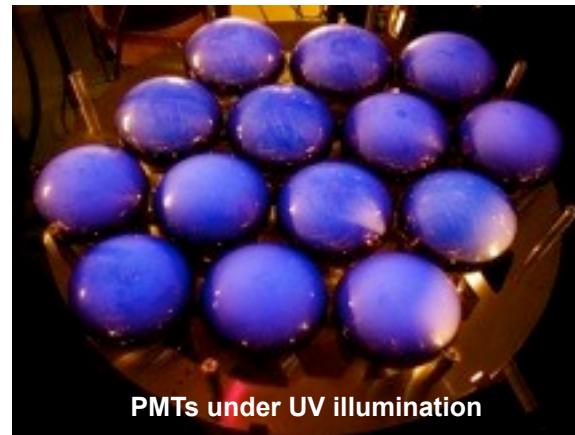
Large volume – double phase – charge/light – 3-D imaging – energy threshold ~ 30 keV – event-by-event interaction type identification – ionisation density discrimination



Developments at CERN

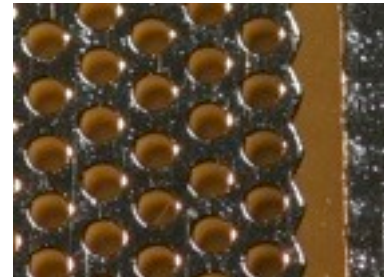
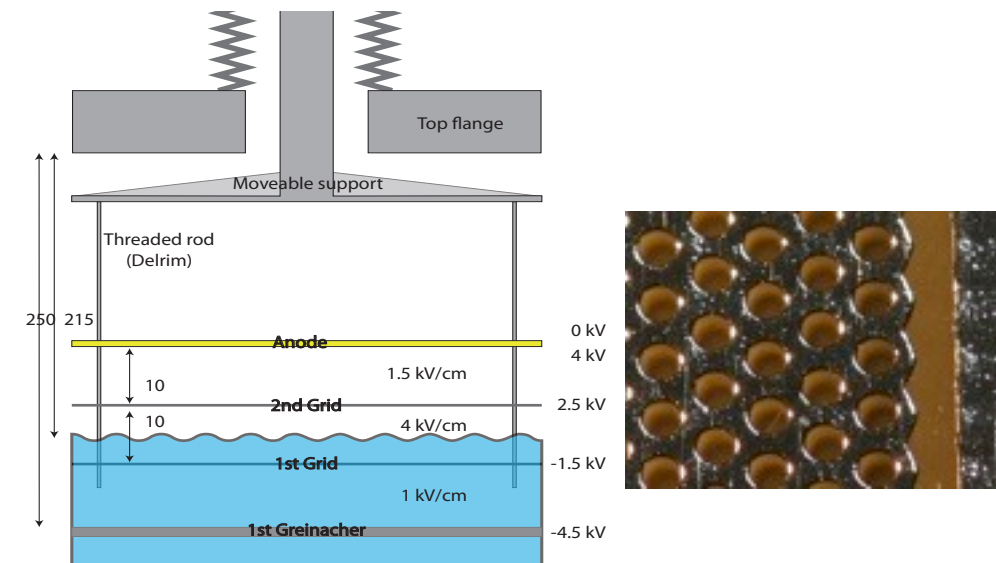
Test runs

Light readout with PMTs and TPB as WLS

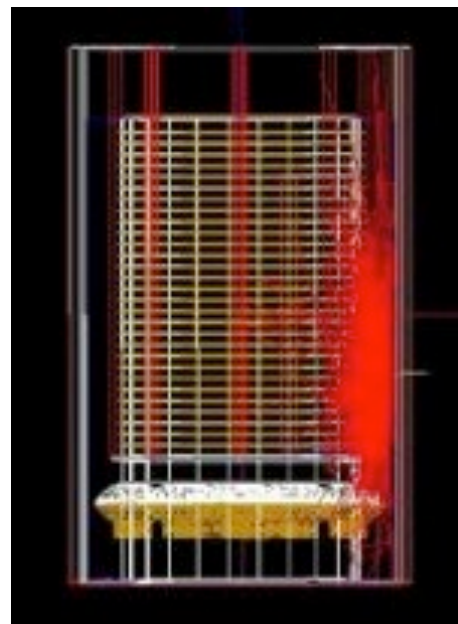


PMTs under UV illumination

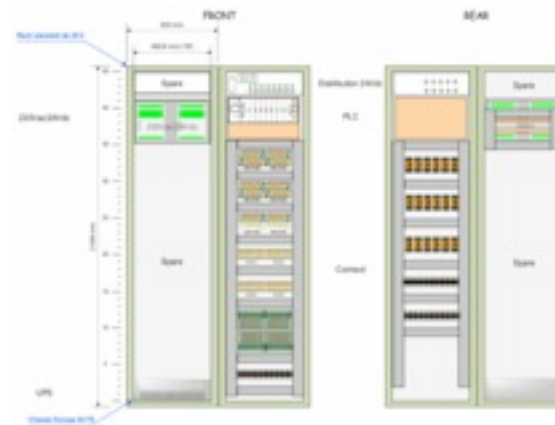
Charge extraction, charge read out



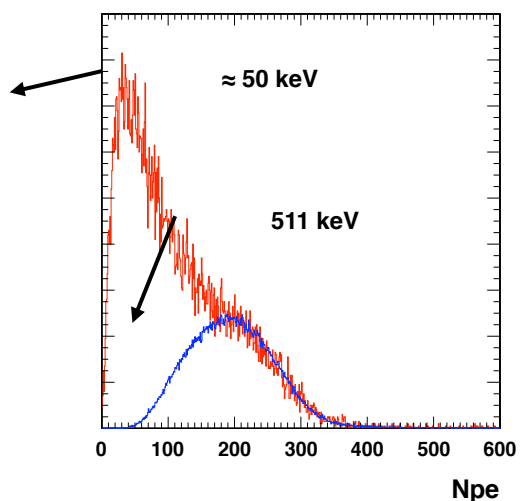
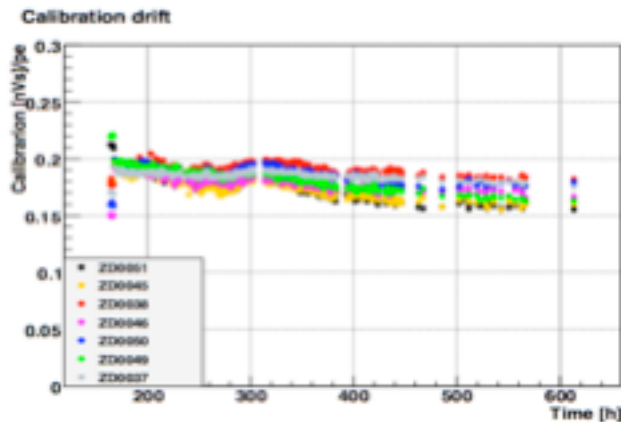
Full GEANT4 detector MC



General control system: PLC



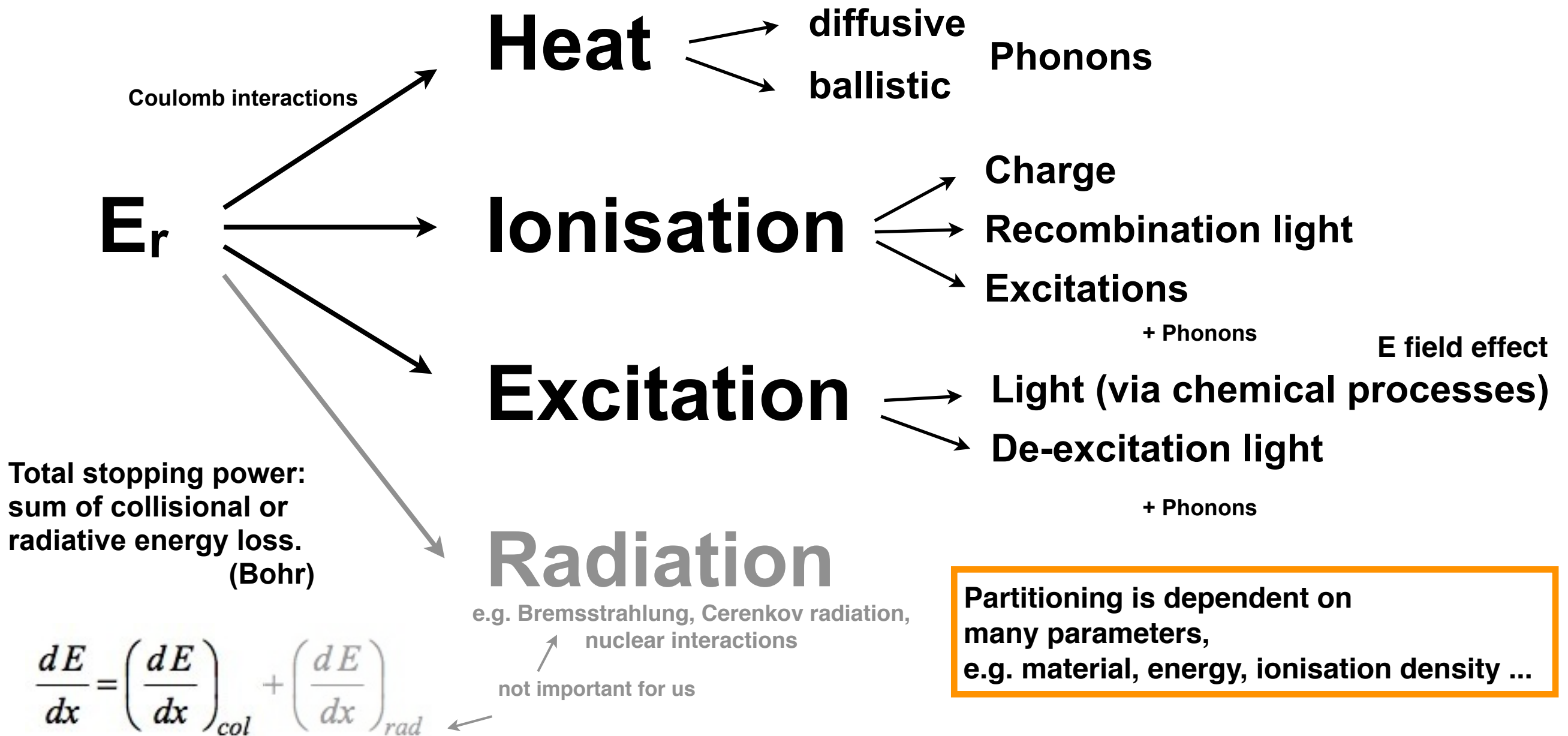
Neutrons from a dd generator for calibration and tests



Cryogenic operation

What has to be measured?

Where does the recoil energy go to ?



- $(dE/dx)_{col}$: electronic energy loss due to Coulomb interactions (i.e., the ionisation and excitation)
- $(dE/dx)_{rad}$: nuclear energy loss (emission of Bremsstrahlung, Cerenkov radiation, nuclear interactions)
- Excitation raises an electron to a higher energy shell
- Ionisation completely removes the electron from the atom
- Energetic electrons are sometimes called delta rays
- $(dE/dx)_{col}$ is also called Linear Energy transfer (LET), the rate of energy loss (dE/dx) of a charged particle due to ionisation and excitation.

We observe different energy scales (yields)

WIMPs (and neutrons) scatter off nuclei, most background sources (gammas, electrons) scatter off electrons

- Detectors have different response to nuclear and electron recoils (@ same energy)

keV_{ee} = measured signal from an electron recoil

keV_r = measured signal from a nuclear recoil

- Often expressed by a **Quench factor** (QF)

For nuclear recoil events:

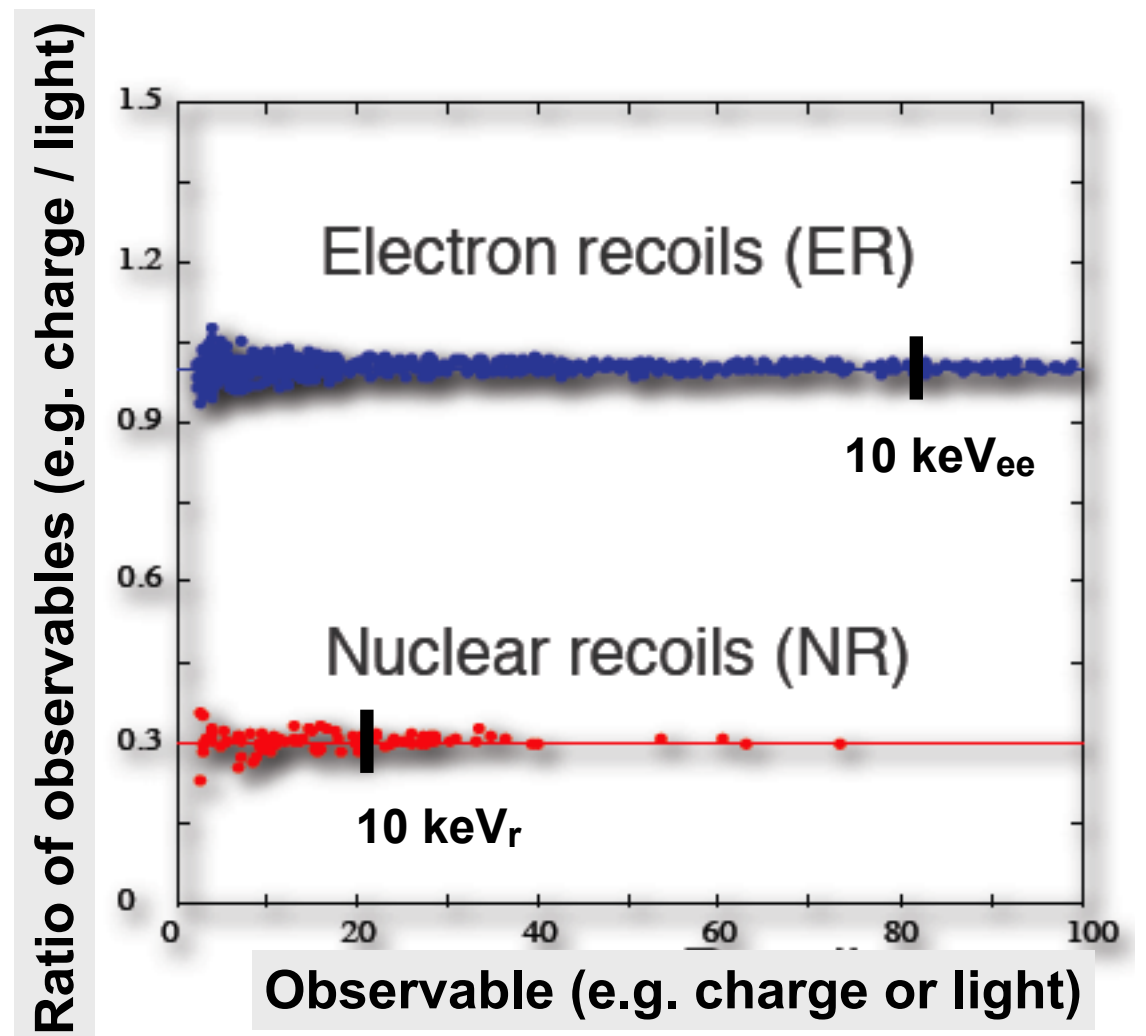
$$E_{\text{visible}} [\text{keV}_{ee}] = QF \cdot E_r [\text{keV}_r]$$

Example ->

$$QF = 25\%$$

$QF = QF(E)$ can be non-linear

The two energy scales can be calibrated with gamma (^{57}Co , ^{133}Ba , ^{137}Cs , ^{60}Co , etc) and neutron (AmBe, ^{252}Cf , n-generator, etc) sources.

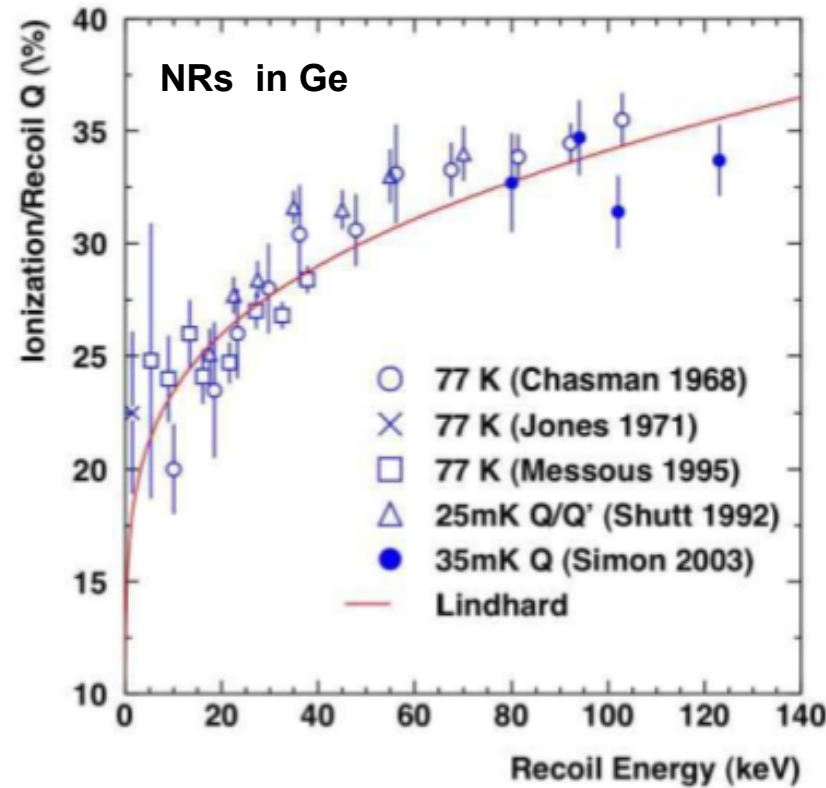


Lindhard theory (nuclear losses = phonons/heat)

Lindhard effect (nuclear quenching) in semiconductors

Ionisation charge quenching in Ge and Si

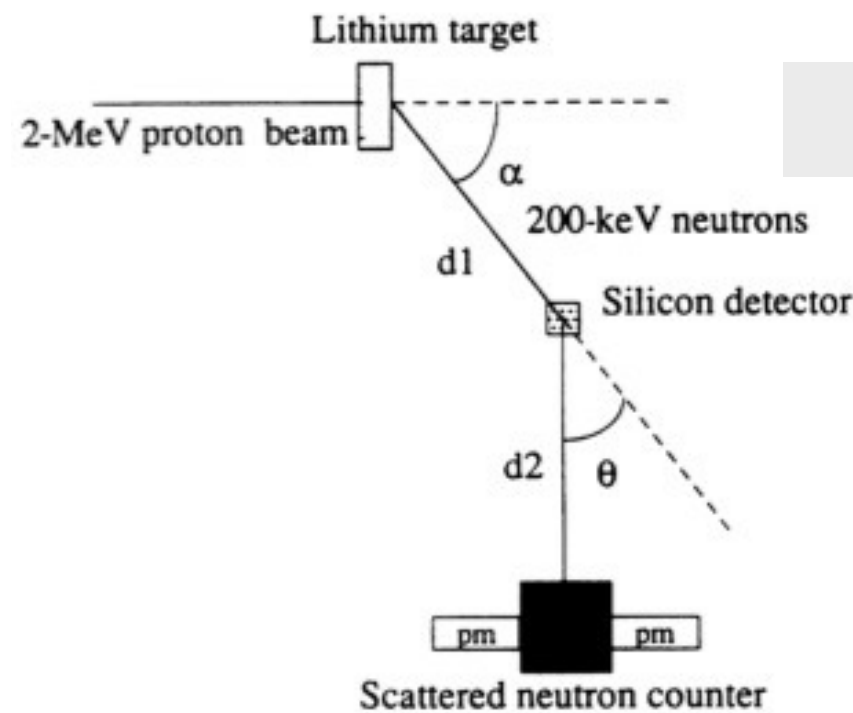
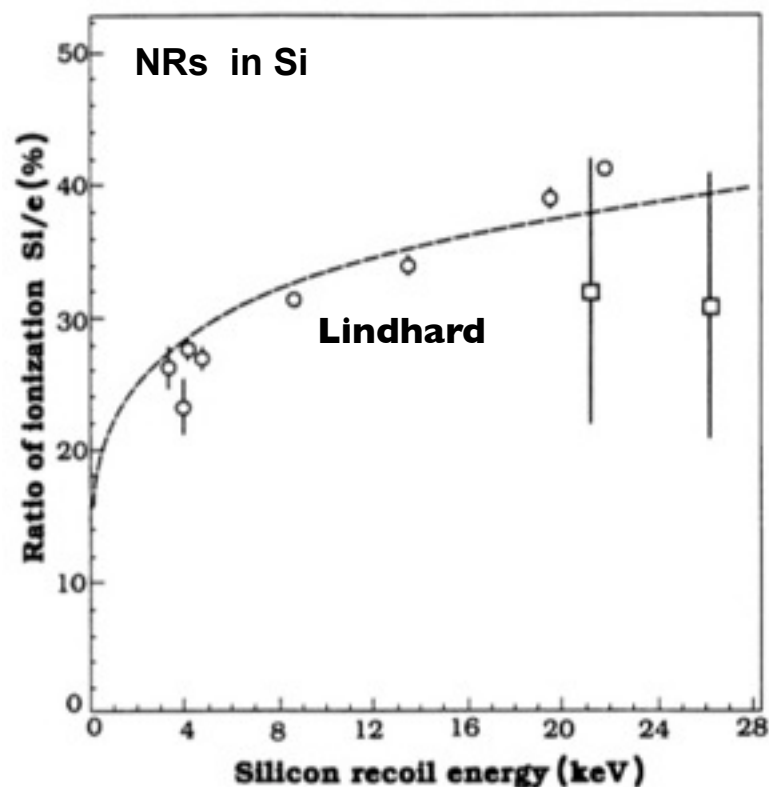
$$\left(\frac{dE}{dx}\right)_{col} = \left(\frac{dE}{dx}\right)_{elec} + \left(\frac{dE}{dx}\right)_{nuc} \quad \begin{array}{l} \text{Nuclear quenching} \\ \text{Energy dependent!} \end{array}$$



Recoiling atoms are less effective in producing primary ionisation or scintillation than electrons of the same energy.

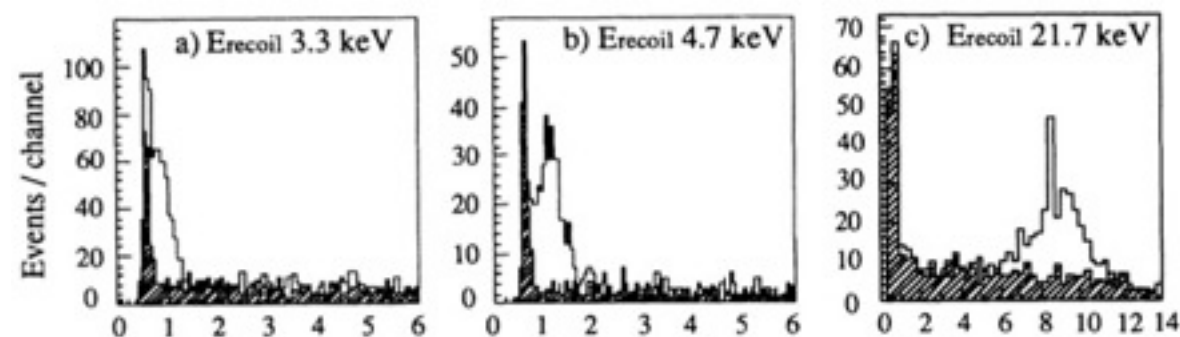
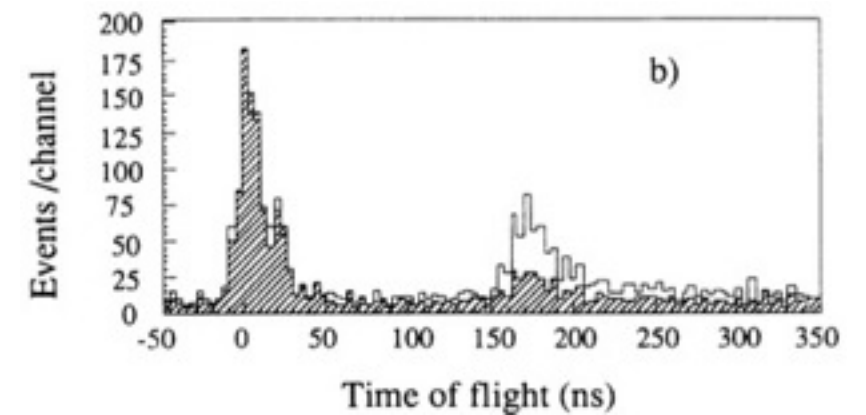
The ratio of the ionisation (and/or scintillation) yield from atomic projectiles to that from electrons, referred to as the **quench factor**

Generally decreases with energy and is material dependent.



Neutron scattering experiment (tagged neutron beam)

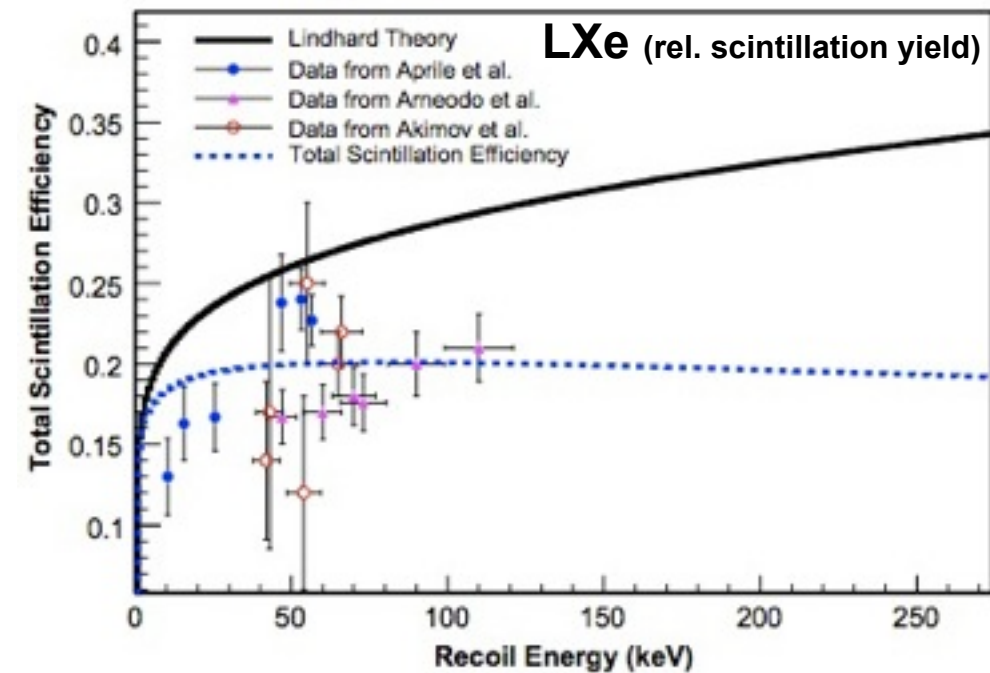
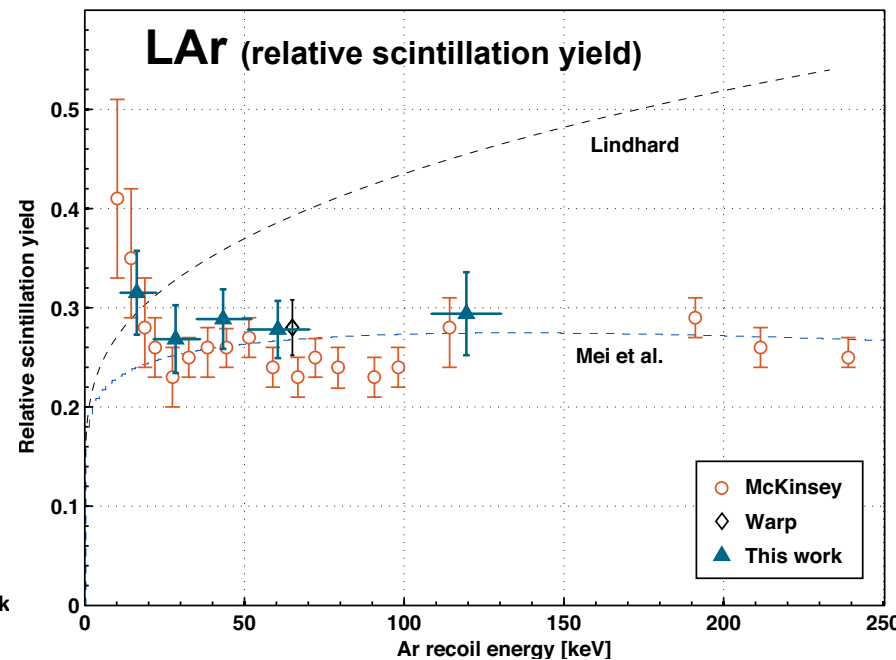
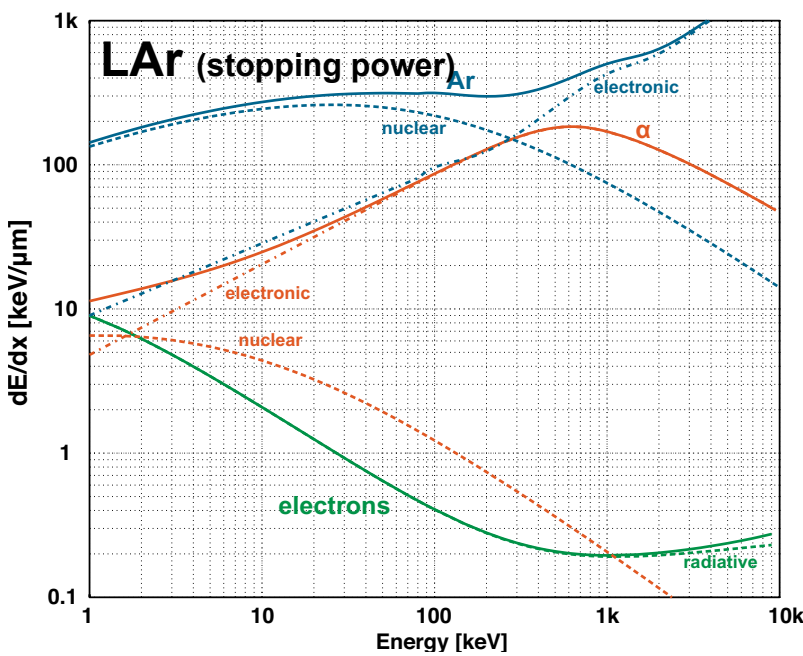
Measurement of the ionisation of slow silicon nuclei in Si
G. Gerbier et al., PRD 42, 3211 (1990)



More complicate: Quenching in noble liquids

On the example of the scintillation light (similar for charge)

$$f \left(\frac{dE}{dx} \right)_{elec} \Rightarrow \text{signal}$$



Semi-empiric approach:

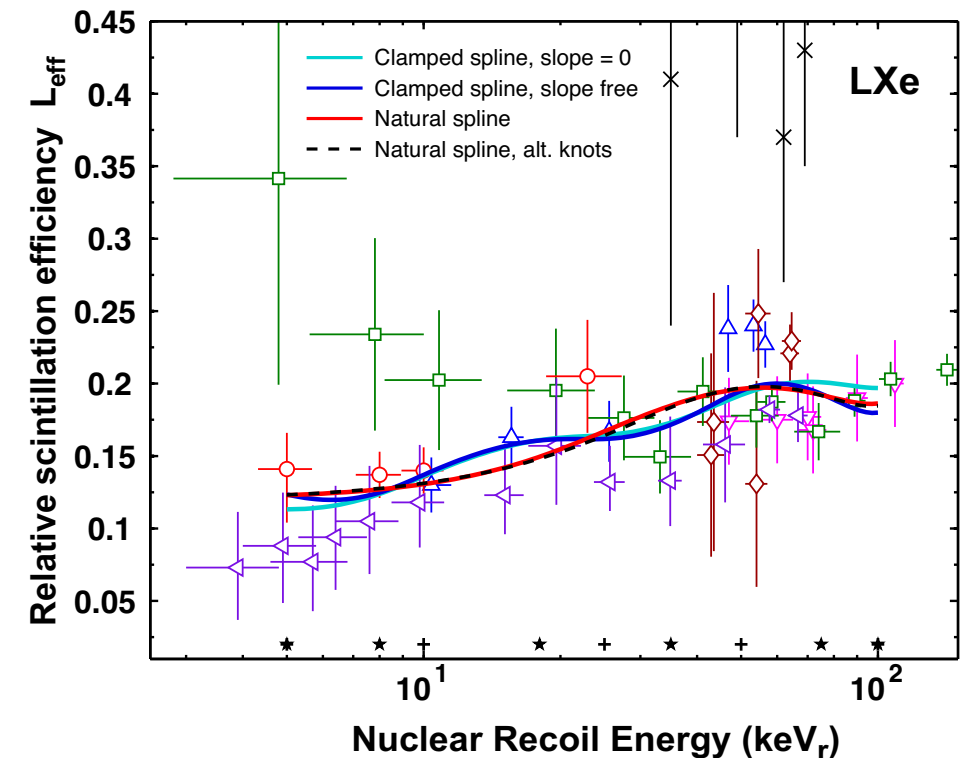
(A model of nuc. recoil scint. eff. in noble liquids; D. Mei et al., Astropart. Phys. 30 (2008) 12)

- In addition to the Lindhard factor we have an efficiency for the production of scintillation light or charge
- Latter is described as luminescence quenching (charge)
- Total yield (quenching) is factorized in a nuclear and a signal term

$$q_f = f_n \times f_l$$

Ionisation density dependent

- Only the latter leads to the production of charges and light
- Need additional theoretical assumptions, **not yet well understood**

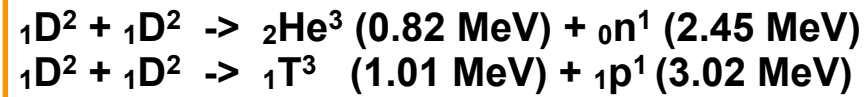


Possible explanation by collisions in signal production:
Bi-excitonic, bi-excimer (Pening) processes, Supereleastic...

See: A. Hitachi and T. Doke, Luminescence quenching in liquid argon under charged particle impact: Relative scintillation yield at different linear energy transfer. Phys.Rev., B46(18), 1992. doi:10.1103/PhysRevB.46.11463.

Setup for scattering of monochromatic neutrons

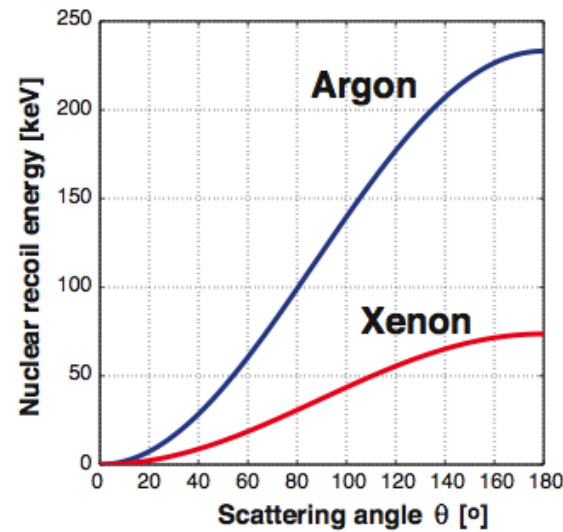
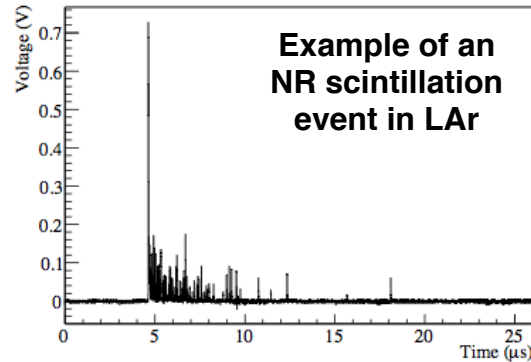
Fusion generator



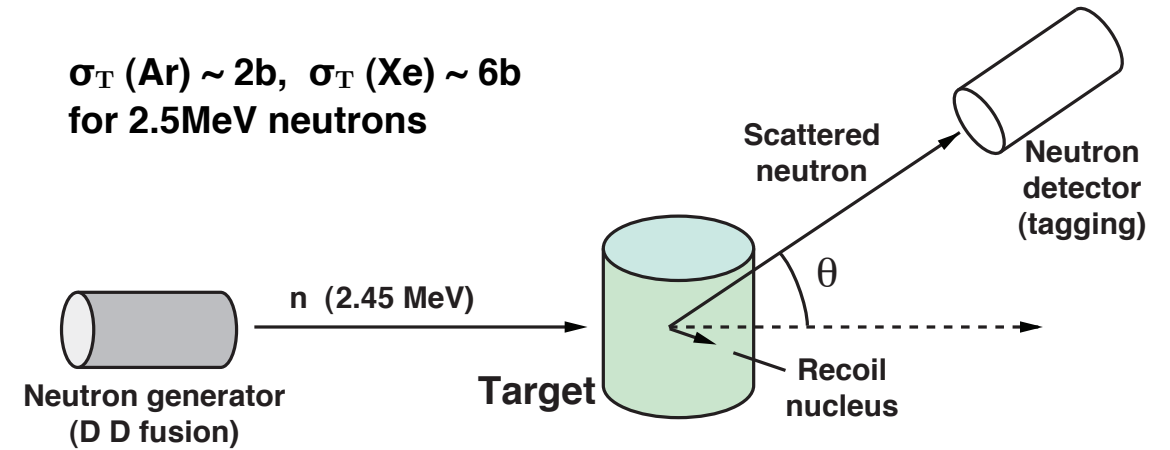
$$E_{\text{recoil}} = \frac{2E_n}{(1+A)^2} \left[1 + A - \cos^2\theta - \cos\theta\sqrt{A^2 + \cos^2\theta - 1} \right] \cong \frac{2E_n A}{(1+A)^2} (1 - \cos\theta)$$

$$R = 1 - \exp\left(-\frac{N_A \rho L A r \sigma_T}{A} l\right)$$

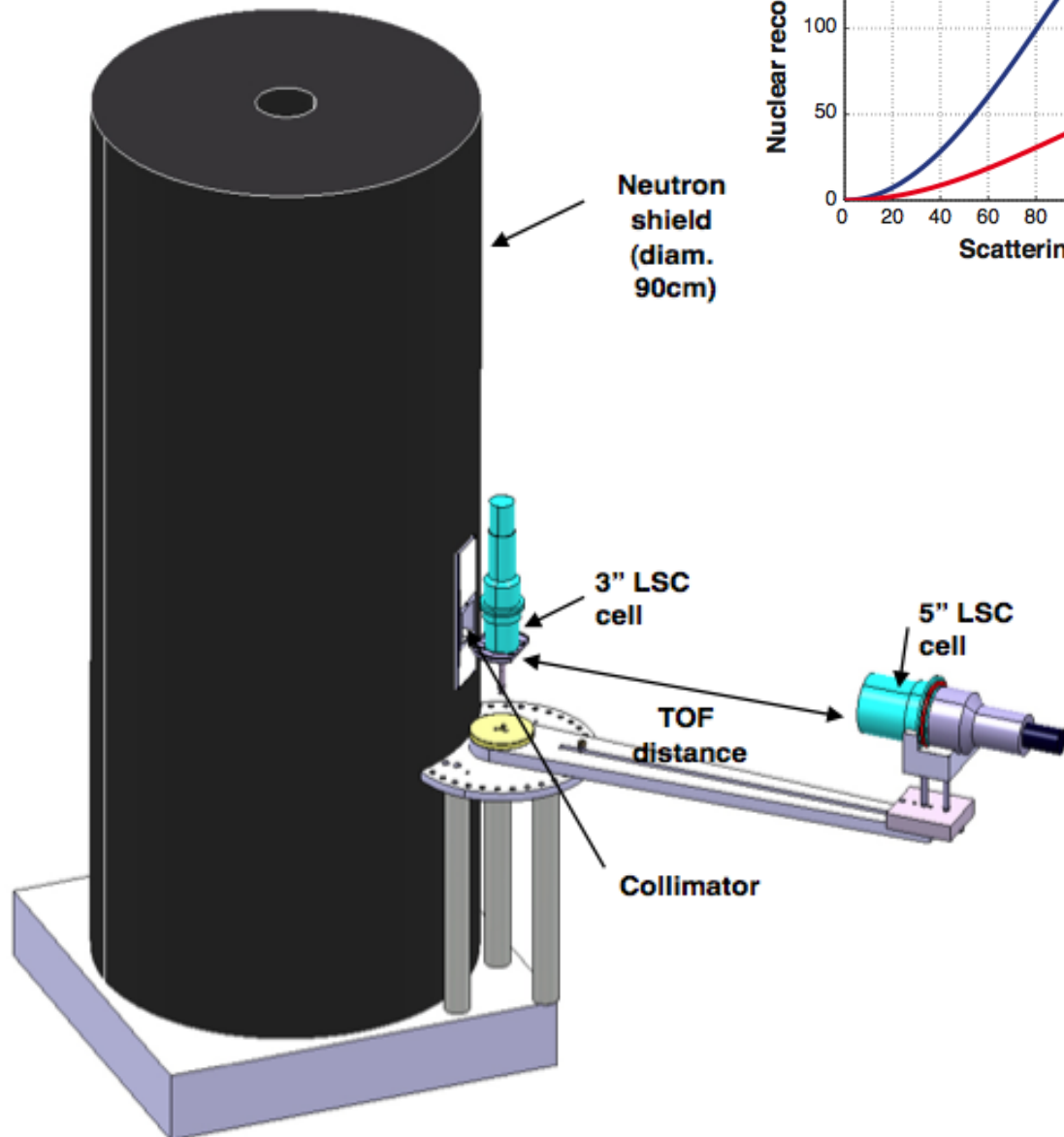
Interaction fraction
 (l length, ρ density, $N_A = 6 \cdot 10^{23}$)



$\sigma_T (\text{Ar}) \sim 2\text{b}$, $\sigma_T (\text{Xe}) \sim 6\text{b}$
 for 2.5MeV neutrons



Setup in our lab at CERN



Let's visit the experimental setup in our laboratory

Christian Regenfus (regenfus@cern.ch)