Probing Beyond the Standard Model Physics with Neutrinoless Double Beta Decay

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### Neutrinos

- Oscillations
- Absolute Mass

### Neutrinoless Double Beta Decay

- Light Neutrino Exchange
- Effective Mass and Seesaw
- New Physics Mechanisms
- Distinguishing Mechanisms
- LNV at the LHC
- Conclusion

## Neutrino Oscillations

### Neutrino interaction states different from mass eigenstates

Neutrino flavour can change through propagation

$$\nu_{i} = \sum_{\alpha} U_{i\alpha} \nu_{\alpha}, \quad \nu_{i}(t) = e^{-i(E_{i}t - p_{i}x)} \nu_{i}$$
$$\Rightarrow P_{\alpha \to \beta} = \sin^{2}(2\theta) \sin^{2} \left( 1.27 \frac{\Delta m^{2}}{eV^{2}} \frac{L/km}{E/GeV} \right)$$

- Solar neutrino oscillations Large mixing
- Atmospheric oscillations
  Maximal mixing
- Reactor and accelerator neutrinos

 $\sin^2(2\theta_{13}) = 0.092 \pm 0.021$ 

- normal hierarchy inverted hierarchy  $(m_3)^2$   $(m_2)^2$   $(\Delta m^2)_{sol}$   $(m_1)^2$   $(\Delta m^2)_{atm}$   $v_e$   $(\Delta m^2)_{atm}$   $v_\mu$   $(\Delta m^2)_{atm}$   $v_\tau$   $(\Delta m^2)_{atm}$
- **Experimental unknowns and anomalies** CP violation? Sign of  $\Delta m_{23}$ ? Sterile Neutrinos?

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### **Absolute Neutrino Mass**

**Energy endpoint in Beta decay** 

 $m_{\beta}^2 = \sum_i |U_{ei}|^2 m_i^2 < (2.2 \,\text{eV})^2$  Katrin:  $m_{\beta} \approx 0.2 \,\text{eV}$ 

Impact on Large Scale Structure

$$\Sigma = \sum_{i} m_i < 0.4 - 1 \,\mathrm{eV}$$

**Neutrinoless Double Beta Decay** 

$$m_{\beta\beta} = \left| \sum_{i} U_{ei}^2 m_{\nu_i} \right| < 0.2 - 2.0 \,\mathrm{eV}$$
 Future Experiments:  
 $m_{\beta\beta} \approx 0.01 \,\mathrm{eV}$ 

eV keV MeV meV GeV TeV v<sub>3</sub> 3 V2 2 V1  $10^2 \ 10^3 \ 10^4 \ 10^5 \ 10^6 \ 10^7 \ 10^8 \ 10^9 \ 10$  $10^{10}_{10}^{11}_{10}^{12}_{10}^{12}$ 10 De Gouvea '04 mass (eV)







generation

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### Neutrinoless Double Beta Decay

- Process:  $(A, Z) \rightarrow (A, Z+2) + 2e^{-1}$
- Uncontroversial detection of 0vββ
  of utmost importance
  - Prove lepton number to be broken
  - Prove neutrinos to be Majorana particles (Schechter, Valle '82)



### • Which mechanism triggers the decay?





General Effective Operator

## Light Neutrino Exchange



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### **Experimental Situation**



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### **Probing Majorana CP Violation?**

### Effective 0vββ Observable

$$m_{\beta\beta} \equiv (m_{\nu})_{ee} = \left| \sum_{i} U_{ei}^{2} m_{\nu_{i}} \right| = \left| m_{\nu_{1}} |V_{el}|^{2} + m_{\nu_{2}} |V_{e2}|^{2} e^{2i\phi_{12}} + m_{\nu_{3}} |V_{e3}|^{2} e^{2i\phi_{23}} \right|$$
  
$$\approx m_{1}^{2} \sqrt{1 - \sin^{2}(2\theta_{12}) \sin^{2}(\phi_{12})} \quad \text{quasi-deg. neutrinos}$$

#### Interplay of mass probes

(J. Auger, FFD, O. Lahav, D. Waters, J. Auger, work in progress)





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## Nuclear Matrix Elements

Decay Rate

$$T = T_{1/2}^{-1} = \frac{m_{\beta\beta}^2}{m_e^2} G^{0\nu} |M^{0\nu}|^2$$

- Requires calculation of matrix element of the nuclear transition via intermediate state
- Many-body problem not solvable from first principle
- Factor 2 3 uncertainty between different nuclear models
  - Shell Models
  - QRPA
  - IBM

# Important to search for 0vββ in several isotopes



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0v2β (exotic, not observed)

Up to

 5 to 6 hbar

0<sup>+</sup> g.s.

100 MeV

 Probes all states

T=6

### New Physics Contributions to 0vßß

Plethora of New Physics Scenarios



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## Effective Mass and Seesaw Mechanism

#### Effective operator for Majorana neutrino mass

$$L = \frac{1}{2} \frac{h_{ij}}{\Lambda_{\text{LNV}}} (\overline{L_i^c} \cdot \tilde{H}) (\tilde{H}^T \cdot L_j) \rightarrow \frac{1}{2} (m_v)_{ij} \overline{v_i^c} v_j$$

Unique dim-5 Operator

### Seesaw Mechanism

Add right-handed neutrinos to the Standard Model particle content,  $M \approx 10^{14} \text{ GeV}$ 

$$L = L_{\rm SM} - \frac{1}{2} \bar{\nu}_R M \nu_R^c + \bar{\nu}_R Y_{\nu} L \cdot H_u$$

### Light neutrino mass matrix at low energies

$$m_{\nu} = m_D^T M^{-1} m_D$$
 for  $m_D \ll M_R$   $m_{\nu} \approx 0.1 \text{eV} \left(\frac{m_D}{100 \text{ GeV}}\right)^2 \left(\frac{M}{10^{14} \text{ GeV}}\right)^{-1}$ 

## Minimal Left-Right Symmetrical Model

Based on

$$SU(3) \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

Pati & Salam '74 Mohapatra & Senjanovic '75

 Higgs Sector: Bidoublet (EW Breaking)
 + Left-handed Triplet + Right-handed Triplet (Breaking Lepton Number + Parity + SU(2)<sub>R</sub>)

### Generating r.h. Neutrino + WR + ZR masses

$$M_{N_i} \approx M_{W_R} \approx M_{Z_R} \approx < \Delta_R >$$

Charged current weak interactions

$$J_W^{\mu-} = \frac{g_L}{\sqrt{2}} \left( \bar{\nu} U_{LL} + \bar{N}^c U_{LR} \right) \gamma^{\mu} e_L + \frac{g_R}{\sqrt{2}} \sin \zeta_W \left( \bar{\nu} U_{RL} + \bar{N} U_{RR} \right) \gamma^{\mu} e_R,$$
  
$$J_{W'}^{\mu-} = -\frac{g_L}{\sqrt{2}} \sin \zeta_W \left( \bar{\nu} U_{LL} + \bar{N} U_{LR} \right) \gamma^{\mu} e_L + \frac{g_R}{\sqrt{2}} \left( \bar{N} U_{RR} + \bar{\nu}^c U_{RL} \right) \gamma^{\mu} e_R,$$

### Neutrinoless Double Beta Decay in the LRSM



## **Disentangling New Physics Scenarios**

### Comparison of 0vββ in multiple isotopes

(FFD, Päs PRL 2007)

$$\frac{T_{1/2}({}^{A}\mathbf{X})}{T_{1/2}({}^{B}\mathbf{Y})} = \frac{G({}^{B}\mathbf{Y}) |M({}^{B}\mathbf{Y})|^{2}}{G({}^{A}\mathbf{X}) |M({}^{A}\mathbf{X})|^{2}}$$

- Depends on  $0\nu\beta\beta$  mechanism
- Independent of details of new physics (if one mechanism dominates)
- Important to search for 0vββ
  in several isotopes



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## **Disentangling New Physics Scenarios**

### Angular and Energy distribution of emitted electrons

(Doi et al. '83; Ali et al. '06; Arnold et al. '10; FFD, Jackson, Nasteva, Söldner-Rembold '10)

$$\frac{d\Gamma}{dE_1dE_2d\cos\theta} = \frac{\Gamma}{2}(1 - k(E_1, E_2)\cos\theta) \qquad -1 < k < 2$$

- Linear in  $\cos\theta$
- $k(E_1, E_2)$  depends on  $0\nu\beta\beta$  mechanism
- Requires measurement of individual electron energies and tracking
- Testable at SuperNEMO
  - Calorimetry and Tracking: Individual electron tracks, vertices and energies
  - Expected sensitivity

$$T_{1/2}$$
 = 2×10<sup>26</sup> years

$$\langle m_v \rangle$$
 = 50-90 meV





- Constraints on the parameter space
  - Half-life measurement



- Statistical uncertainties from simulation
- Theoretical nuclear matrix element uncertainty: 30%

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#### Probing BSM Physics with Neutrinoless DBD

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### SNO+

- Search for low energy solar, geo-, reactor and Supernova neutrinos
- Search for 0vββ in <sup>150</sup>Nd
  - Loaded in liquid scintillator (0.1%, ~ 50 kg <sup>150</sup>Nd)
  - Good isotope for 0vββ discovery (Large phase space, large Q value, good natural abundance)
  - Large nuclear deformation (NME uncertainty?)

### Sensitivity to New Physics?

- Work in progress
  J. Hartnell, A. Back (Sussex),
  F. di Lodovico (QMUL), FFD (UCL)
- Majoron modes
- Possible upgrades and other isotopes







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## Lepton Number Violation at the LHC

- Inverse or "Bent" Seesaw mechanism
- R-Parity Violating SUSY
- Left-Right symmetric models
  Right-handed neutrinos couple with gauge strength to charged leptons



Probing right-handed neutrino mixing, Large LFV and LNV rates at LHC





Das, FFD, Kittel, Valle (2012)

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Probing BSM Physics with Neutrinoless DBD

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 Strong experimental program to probe absolute mass

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### 0vββ is crucial probe for BSM physics

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  New LNV physics at the EW scale
- Prepared for the worst
  Only 5-dim operator from LNV
  at the GUT scale



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- Neutrino impact on cosmology Baryogenesis via Leptogenesis?
- Rich phenomenology in models of neutrino mass generation
  - Charged lepton flavour violation
  - LFV and LNV processes at the LHC

