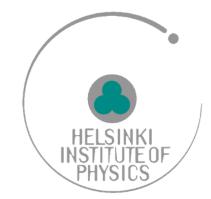
### LHC data and aspects of new physics

#### Kimmo Tuominen



#### Jyväskylä University & Helsinki Institute of Physics





### Outline

# 1. Need for new physics

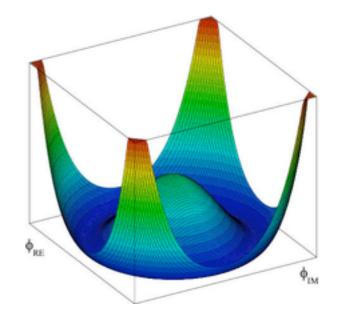
# 2. LHC data and a new physics model example (T. Alanne et al. 1303.3615)

# 3. Paradigm shift?

(M. Heikinheimo et al. arXiv: next Monday...)

# 1. The need for new physics

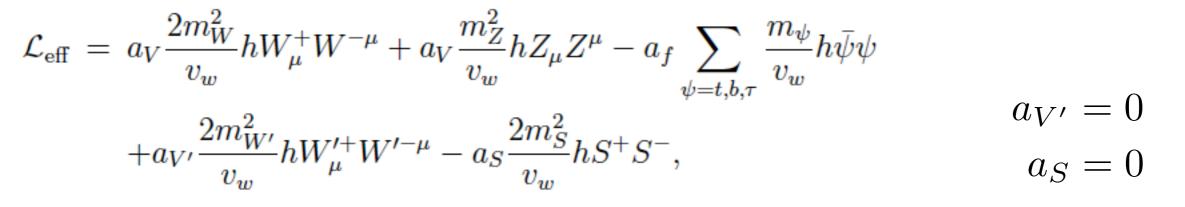
$$\mathcal{L}_{H} = (D_{\mu}\phi)^{\dagger}D^{\mu}\phi - V(\phi)$$
$$V(\phi) = \mu^{2}\phi^{\dagger}\phi + \frac{1}{4}\lambda \left(\phi^{\dagger}\phi\right)^{2}, \quad \mu^{2} < 0.$$



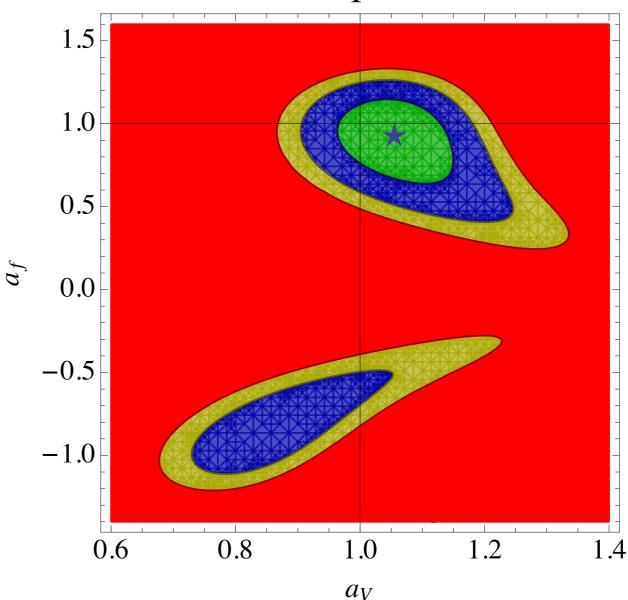
### The discovery: A SM Higgs boson exists

$$\begin{aligned} \mathcal{L}_{\text{eff}} &= a_V \frac{2m_W^2}{v_w} h W_{\mu}^+ W^{-\mu} + a_V \frac{m_Z^2}{v_w} h Z_{\mu} Z^{\mu} - a_f \sum_{\psi=t,b,\tau} \frac{m_{\psi}}{v_w} h \bar{\psi} \psi \\ &+ a_{V'} \frac{2m_{W'}^2}{v_w} h W_{\mu}'^+ W'^{-\mu} - a_S \frac{2m_S^2}{v_w} h S^+ S^-, \end{aligned}$$

### The discovery: A SM Higgs boson exists



Data: post-Moriond '13



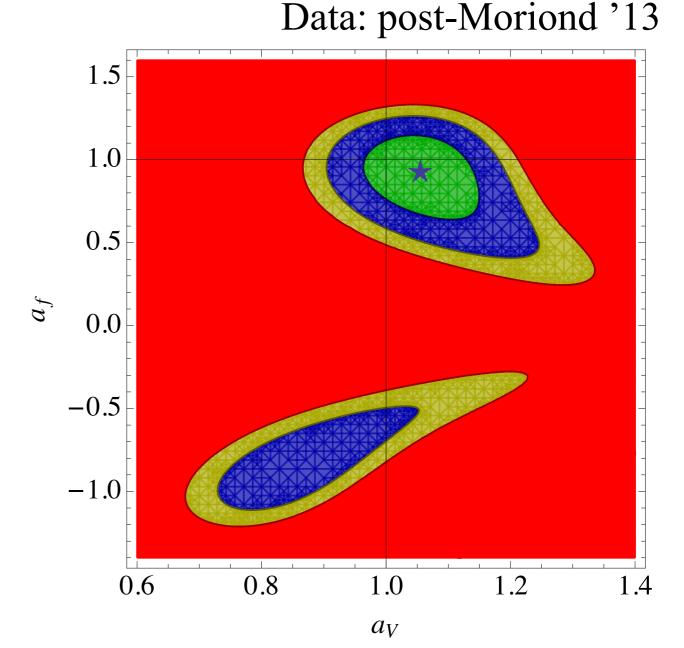
### The discovery: A SM Higgs boson exists

$$\begin{aligned} \mathcal{L}_{\text{eff}} &= a_V \frac{2m_W^2}{v_w} h W_{\mu}^+ W^{-\mu} + a_V \frac{m_Z^2}{v_w} h Z_{\mu} Z^{\mu} - a_f \sum_{\psi=t,b,\tau} \frac{m_{\psi}}{v_w} h \bar{\psi} \psi \\ &+ a_{V'} \frac{2m_{W'}^2}{v_w} h W_{\mu'}^{\prime+} W^{\prime-\mu} - a_S \frac{2m_S^2}{v_w} h S^+ S^-, \end{aligned} \qquad \begin{aligned} a_{V'} &= 0 \\ a_S &= 0 \end{aligned}$$

SM corresponds to

 $a_V = a_f = 1$  $a_{V'} = a_S = 0$ 

Giardino et al., (2013), Ellis & You (2013), Djouadi & Moreau (2013), Falkowski et al. (2013), Alanne et al. (2013),



# Even if we know the mass of the Higgs boson, we do not know what is the origin of it!

# Even if we know the mass of the Higgs boson, we do not know what is the origin of it!

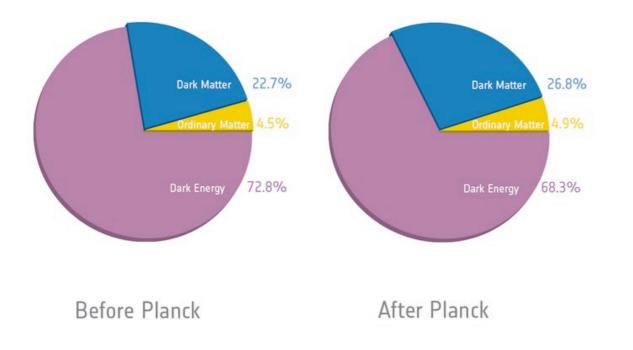
#### Further puzzles:

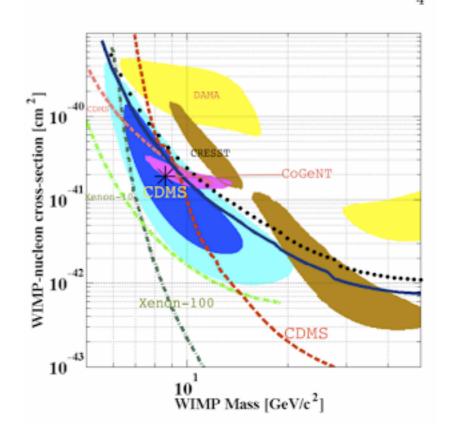
Dark matter?

Origin of matter over antimatter?

Neutrino masses?

Flavor physics?





# Even if we know the mass of the Higgs boson, we do not know what is the origin of it!

$$v_{\rm weak} \sim F_{\pi}$$

### 2. Example of a viable model

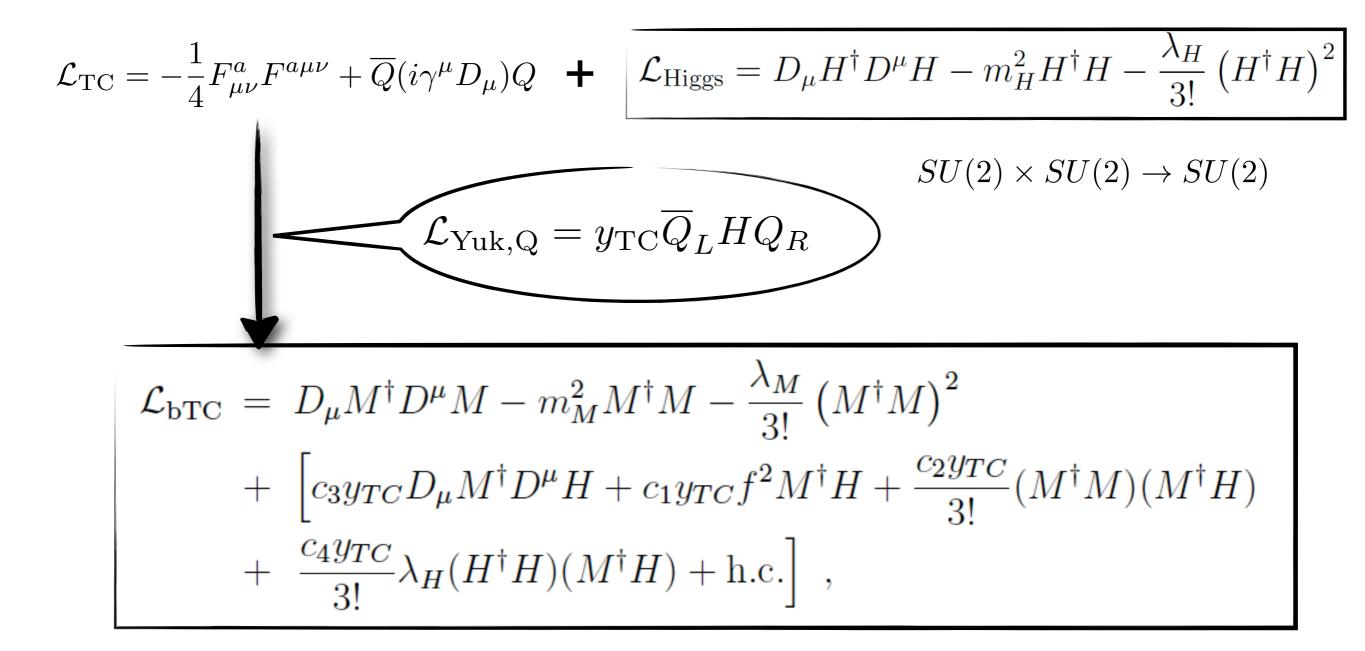
Antola et al. (2009), Fukano et al. (2011), Alanne et al. (2013)

$$\mathcal{L}_{\mathrm{TC}} = -\frac{1}{4} F^a_{\mu\nu} F^{a\mu\nu} + \overline{Q} (i\gamma^{\mu} D_{\mu}) Q + \mathcal{L}_{\mathrm{Higgs}} = D_{\mu} H^{\dagger} D^{\mu} H - m_H^2 H^{\dagger} H - \frac{\lambda_H}{3!} \left( H^{\dagger} H \right)^2$$

 $SU(2) \times SU(2) \rightarrow SU(2)$ 

 $\mathcal{L}_{\mathrm{Yuk},\mathrm{Q}} = y_{\mathrm{TC}} \overline{Q}_L H Q_R$ 

Antola et al. (2009), Fukano et al. (2011), Alanne et al. (2013)



Antola et al. (2009), Fukano et al. (2011), Alanne et al. (2013)

$$\mathcal{L}_{\mathrm{TC}} = -\frac{1}{4} F^{a}_{\mu\nu} F^{a\mu\nu} + \overline{Q} (i\gamma^{\mu} D_{\mu}) Q + \mathcal{L}_{\mathrm{Higgs}} = D_{\mu} H^{\dagger} D^{\mu} H - m_{H}^{2} H^{\dagger} H - \frac{\lambda_{H}}{3!} (H^{\dagger} H)^{2}$$

$$SU(2) \times SU(2) \rightarrow SU(2)$$

$$\mathcal{L}_{\mathrm{Yuk}, Q} = y_{\mathrm{TC}} \overline{Q}_{L} H Q_{R}$$

$$SU(2) \times SU(2) \rightarrow SU(2)$$

$$\mathcal{L}_{\mathrm{bTC}} = D_{\mu} M^{\dagger} D^{\mu} M - m_{M}^{2} M^{\dagger} M - \frac{\lambda_{M}}{3!} (M^{\dagger} M)^{2}$$

$$+ \left[ c_{3} y_{TC} D_{\mu} M^{\dagger} D^{\mu} H + c_{1} y_{TC} f^{2} M^{\dagger} H + \frac{c_{2} y_{TC}}{3!} (M^{\dagger} M) (M^{\dagger} H) \right]$$

$$+ \frac{c_{4} y_{TC}}{3!} \lambda_{H} (H^{\dagger} H) (M^{\dagger} H) + \mathrm{h.c.} \right] ,$$

$$\mathcal{L}_{\text{Yuk}} = (y_u)_{ij} H \bar{Q}_i U_j + (y_d)_{ij} H^{\dagger} \bar{Q}_i D_j + (y_\ell)_{ij} H^{\dagger} \bar{L}_i E_j + \text{h.c.}$$

Antola et al. (2009), Fukano et al. (2011), Alanne et al. (2013)

$$\mathcal{L}_{\mathrm{TC}} = -\frac{1}{4} F^{a}_{\mu\nu} F^{a\mu\nu} + \overline{Q} (i\gamma^{\mu} D_{\mu}) Q + \mathcal{L}_{\mathrm{Higgs}} = D_{\mu} H^{\dagger} D^{\mu} H - m_{H}^{2} H^{\dagger} H - \frac{\lambda_{H}}{3!} (H^{\dagger} H)^{2}$$

$$\mathcal{L}_{\mathrm{Yuk}, Q} = y_{\mathrm{TC}} \overline{Q}_{L} H Q_{R}$$

$$SU(2) \times SU(2) \rightarrow SU(2)$$

$$\mathcal{L}_{\mathrm{bTC}} = D_{\mu} M^{\dagger} D^{\mu} M - m_{M}^{2} M^{\dagger} M - \frac{\lambda_{M}}{3!} (M^{\dagger} M)^{2}$$

$$+ \left[ c_{3} y_{TC} D_{\mu} M^{\dagger} D^{\mu} H + c_{1} y_{TC} f^{2} M^{\dagger} H + \frac{c_{2} y_{TC}}{3!} (M^{\dagger} M) (M^{\dagger} H) \right]$$

$$+ \frac{c_{4} y_{TC}}{3!} \lambda_{H} (H^{\dagger} H) (M^{\dagger} H) + \mathrm{h.c.} \right] ,$$

$$\mathcal{L}_{\text{Yuk}} = (y_u)_{ij} H \bar{Q}_i U_j + (y_d)_{ij} H^{\dagger} \bar{Q}_i D_j + (y_\ell)_{ij} H^{\dagger} \bar{L}_i E_j + \text{h.c.}$$

A simple model for EW symmetry breaking and fermion masses

$$v_w^2 = v^2 + f^2 + 2c_3 y_{TC} f v = (246 \text{ GeV})^2 , \quad \langle M \rangle = \frac{f}{\sqrt{2}} , \quad \langle H \rangle = \frac{v}{\sqrt{2}}$$

8 parameters:  $\lambda_H, \lambda_M, c_1, c_2, c_3, c_4, y_{\text{TC}}, v$ 

$$v_w^2 = v^2 + f^2 + 2c_3 y_{TC} f v = (246 \text{ GeV})^2 , \quad \langle M \rangle = \frac{f}{\sqrt{2}} , \quad \langle H \rangle = \frac{v}{\sqrt{2}}$$

8 parameters:  $\lambda_H, \lambda_M, c_1, c_2, c_3, c_4, y_{\text{TC}}, v$ 

Theory constraints:  $\lambda_H, \lambda_M > 0$ ;  $\lambda_H + \lambda_M > 2(c_2 + c_4\lambda_H) y_{TC}$ .

#### Pheno constraints:

 $m_{h^0} = 125 \pm 1 \text{ GeV}$  ,  $m_{H^\pm} = m_{A^0} > 100 \text{ GeV}$  ,  $m_{H^0} > 600 \text{ GeV}$  ,

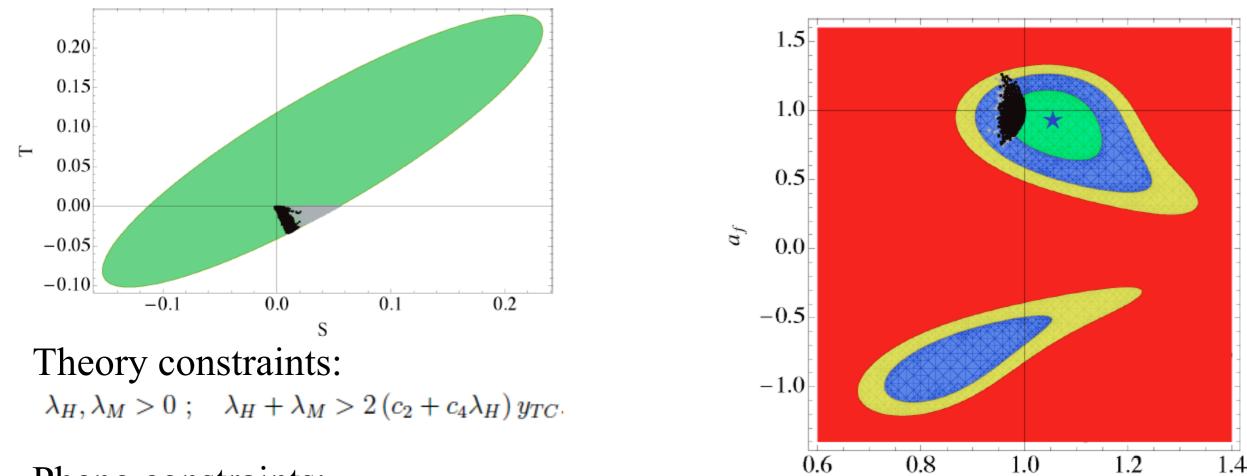
 $S = 0.04 \pm 0.09$ ,  $T = 0.07 \pm 0.08$ , r(S,T) = 88%,  $m_{A^0}, m_{H^0} < 5\Lambda_{TC}$ .

#### Random scan:

$$0 < \lambda_H, \lambda_M < (4\pi)^2, \ 2\pi < |c_1|, |c_2|, |c_3^{-1}|, |c_4^{-1}| < 8\pi, \ |c_3y_{TC}| < 1$$
  
$$|y_t| < 4\pi, \ f = \pm \sqrt{v_w^2 - v^2 \left(1 - c_3^2 y_{TC}^2\right)} - v c_3 y_{TC}, |v| < v_w \left(1 - c_3^2 y_{TC}^2\right)^{-1/2}$$

$$v_w^2 = v^2 + f^2 + 2c_3 y_{TC} f v = (246 \text{ GeV})^2 , \quad \langle M \rangle = \frac{f}{\sqrt{2}} , \quad \langle H \rangle = \frac{v}{\sqrt{2}}$$

#### 8 parameters: $\lambda_H, \lambda_M, c_1, c_2, c_3, c_4, y_{\text{TC}}, v$



 $a_V$ 

#### Pheno constraints:

 $m_{h^0} = 125 \pm 1 \text{ GeV}$  ,  $m_{H^\pm} = m_{A^0} > 100 \text{ GeV}$  ,  $m_{H^0} > 600 \text{ GeV}$  ,

 $S = 0.04 \pm 0.09$ ,  $T = 0.07 \pm 0.08$ , r(S,T) = 88%,  $m_{A^0}, m_{H^0} < 5\Lambda_{TC}$ .

#### Random scan:

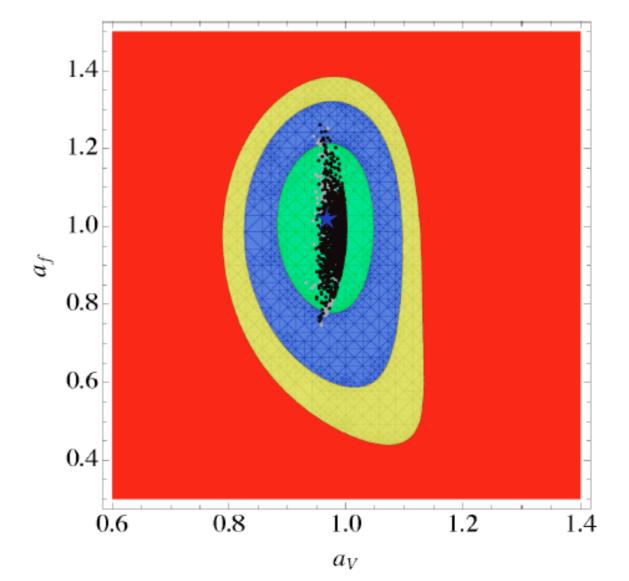
$$0 < \lambda_H, \lambda_M < (4\pi)^2, \ 2\pi < |c_1|, |c_2|, |c_3^{-1}|, |c_4^{-1}| < 8\pi, \ |c_3y_{TC}| < 1$$
  
$$|y_t| < 4\pi, \ f = \pm \sqrt{v_w^2 - v^2 \left(1 - c_3^2 y_{TC}^2\right)} - v c_3 y_{TC}, |v| < v_w \left(1 - c_3^2 y_{TC}^2\right)^{-1/2}$$

## To distinguish BTC from 2HDM: Vector mesons

$$+ m^{2} \operatorname{Tr} \left[ C_{L\mu}^{2} + C_{R\mu}^{2} \right] + \frac{1}{2} \operatorname{Tr} \left[ D_{\mu} M D^{\mu} M^{\dagger} \right] - \tilde{g^{2}} r_{2} \operatorname{Tr} \left[ C_{L\mu} M C_{R}^{\mu} M^{\dagger} \right]$$
$$- \frac{i \tilde{g} r_{3}}{4} \operatorname{Tr} \left[ C_{L\mu} \left( M D^{\mu} M^{\dagger} - D^{\mu} M M^{\dagger} \right) + C_{R\mu} \left( M^{\dagger} D^{\mu} M - D^{\mu} M^{\dagger} M \right) \right] + \frac{\tilde{g}^{2} s}{4} \operatorname{Tr} \left[ C_{L\mu}^{2} + C_{R\mu}^{2} \right] \operatorname{Tr} \left[ M M^{\dagger} \right]$$
$$C_{L\mu} \equiv A_{L\mu} - \frac{g}{\tilde{g}} \widetilde{W_{\mu}} , \quad C_{R\mu} \equiv A_{R\mu} - \frac{g'}{\tilde{g}} \widetilde{B_{\mu}}$$
(Belyaev et al. 2008)

### To distinguish BTC from 2HDM: Vector mesons

$$+ m^{2} \operatorname{Tr} \left[ C_{L\mu}^{2} + C_{R\mu}^{2} \right] + \frac{1}{2} \operatorname{Tr} \left[ D_{\mu} M D^{\mu} M^{\dagger} \right] - \tilde{g^{2}} r_{2} \operatorname{Tr} \left[ C_{L\mu} M C_{R}^{\mu} M^{\dagger} \right] \\ - \frac{i \tilde{g} r_{3}}{4} \operatorname{Tr} \left[ C_{L\mu} \left( M D^{\mu} M^{\dagger} - D^{\mu} M M^{\dagger} \right) + C_{R\mu} \left( M^{\dagger} D^{\mu} M - D^{\mu} M^{\dagger} M \right) \right] + \frac{\tilde{g}^{2} s}{4} \operatorname{Tr} \left[ C_{L\mu}^{2} + C_{R\mu}^{2} \right] \operatorname{Tr} \left[ M M^{\dagger} \right] \\ C_{L\mu} \equiv A_{L\mu} - \frac{g}{\tilde{g}} \widetilde{W_{\mu}} , \quad C_{R\mu} \equiv A_{R\mu} - \frac{g'}{\tilde{g}} \widetilde{B_{\mu}}$$
(Belyaev et al. 2008)



 $m_{W'} \sim 1 \,\mathrm{TeV}$ 

OK with ATLAS search on sequential W'

Life of walking technicolor

Life of walking technicolor

The birth: Absence of FCNCs.

Life of walking technicolor

The birth: Absence of FCNCs.

The youth: S parameter can be small.

Life of walking technicolor

The birth: Absence of FCNCs.

The youth: S parameter can be small.

Adulthood: The Higgs can be light.

Life of walking technicolor

The birth: Absence of FCNCs.

The youth: S parameter can be small.

Adulthood: The Higgs can be light.



# The naturality paradigm (simplified):

- New physics needed to explain the origin of  $v_{\rm weak} \sim F_{\pi}$
- Typically implies a rich spectrum around  $\Lambda \sim 4\pi F_\pi$
- SM is an eff. theory below  $\Lambda$  No hierarchy problem

# The naturality paradigm (simplified):

- New physics needed to explain the origin of  $v_{\rm weak} \sim F_{\pi}$
- Typically implies a rich spectrum around  $\Lambda \sim 4\pi F_{\pi}$
- SM is an eff. theory below  $\Lambda$  No hierarchy problem

But SM never had the hierarchy problem!

(-- all radiative corrections are logarithmic if there is no physical scale above EW scale.)

# 3.A paradigm shift?

Can we explain the origin of the EW scale, weakly coupled dark matter,

and have no obvious new dofs (beyond Higgs)?

**Proposal: everything is scale invariant at tree level.** 

# **Proposal: everything is scale invariant at tree level.**

All scales generated quantum mechanically.

# **Proposal: everything is scale invariant at tree level.**

All scales generated quantum mechanically.

Coleman--Weinberg '73, Bardeen 95, C.T. Hill 05,...

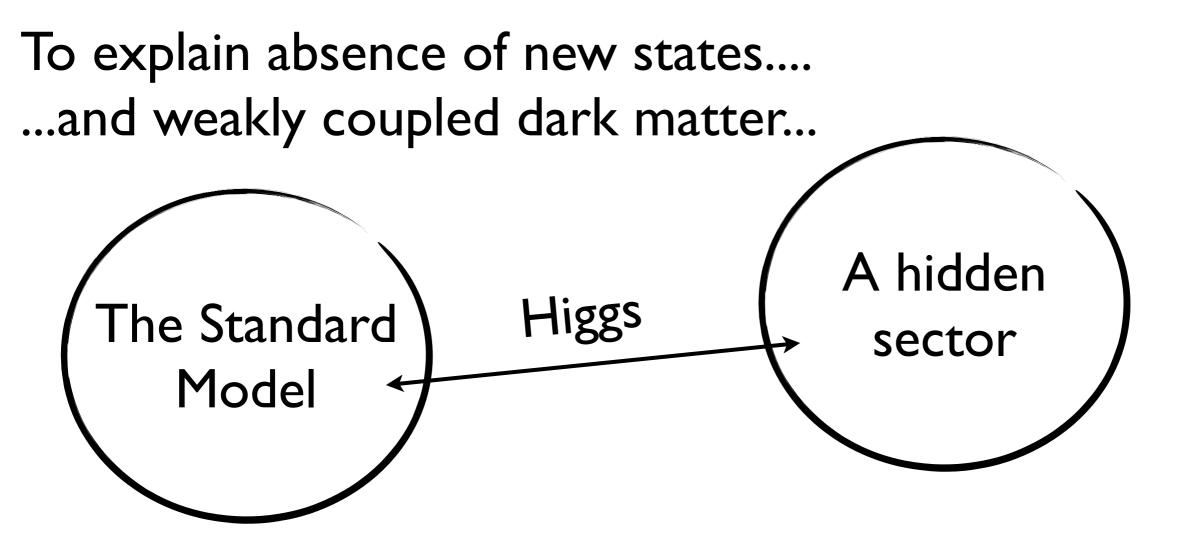
# **Proposal: everything is scale invariant at tree level.**

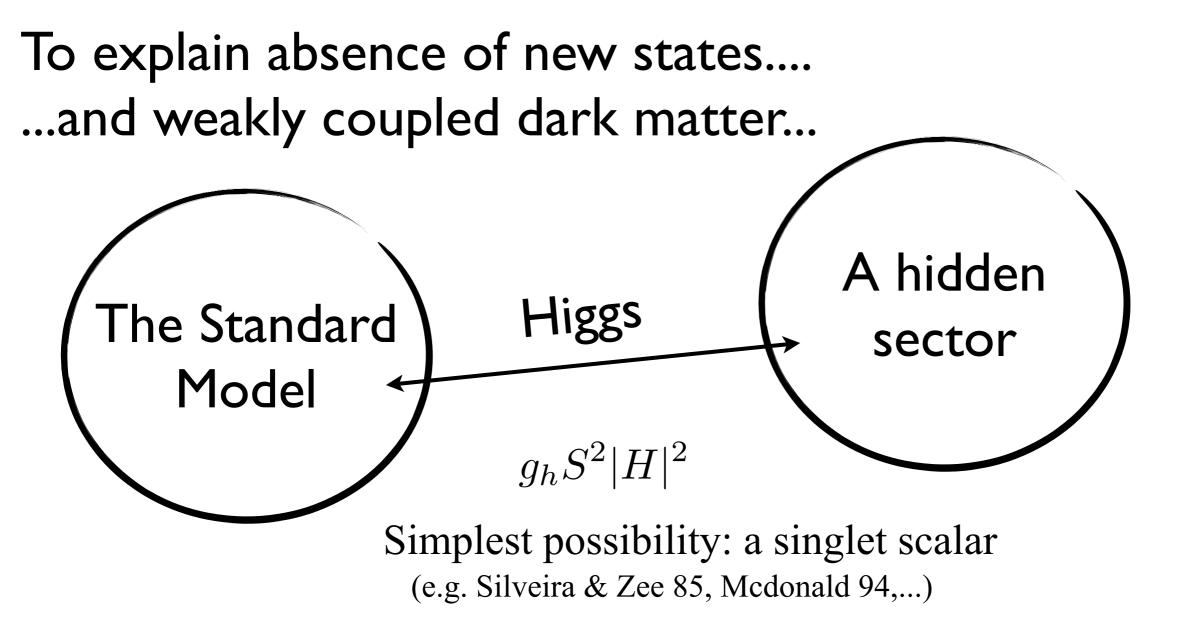
All scales generated quantum mechanically.

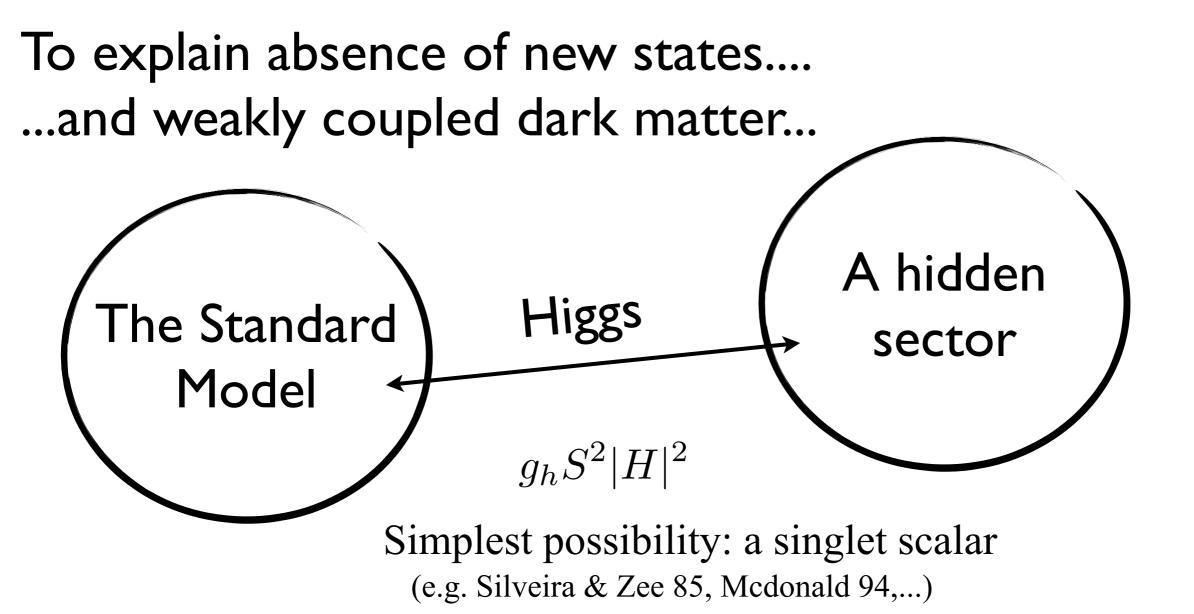
Coleman--Weinberg '73, Bardeen 95, C.T. Hill 05,...

(Compare with the copernican principle in cosmology)

To explain absence of new states.... ...and weakly coupled dark matter...

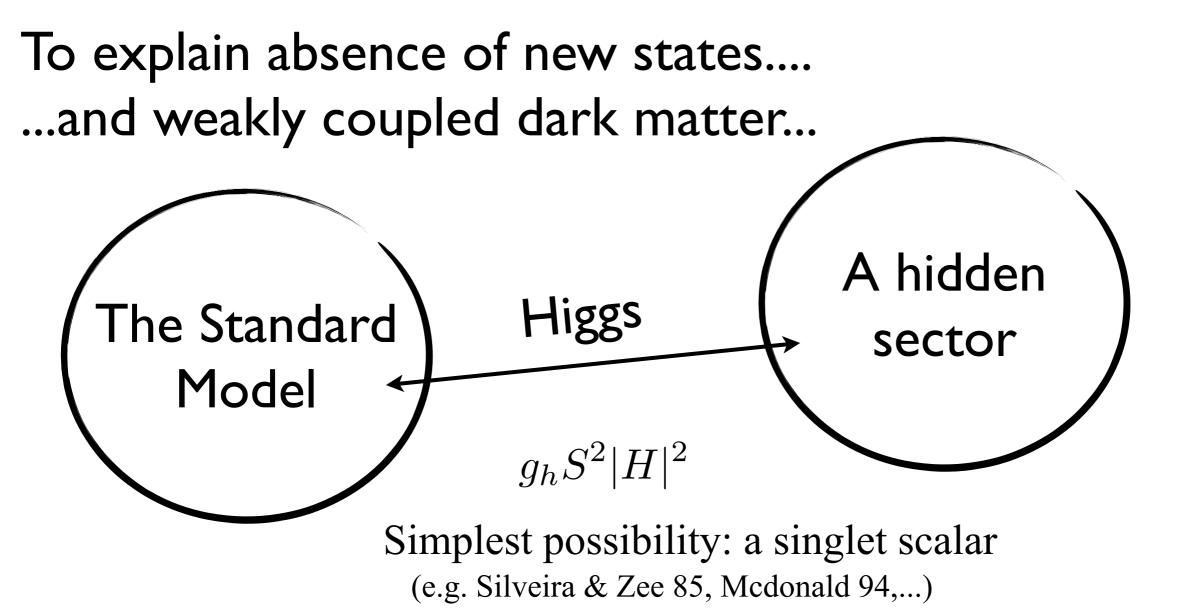






A novel possibility: complex scalar with U(1) charge & BEC. Leads to EW scale and protects it from radiative corrections.

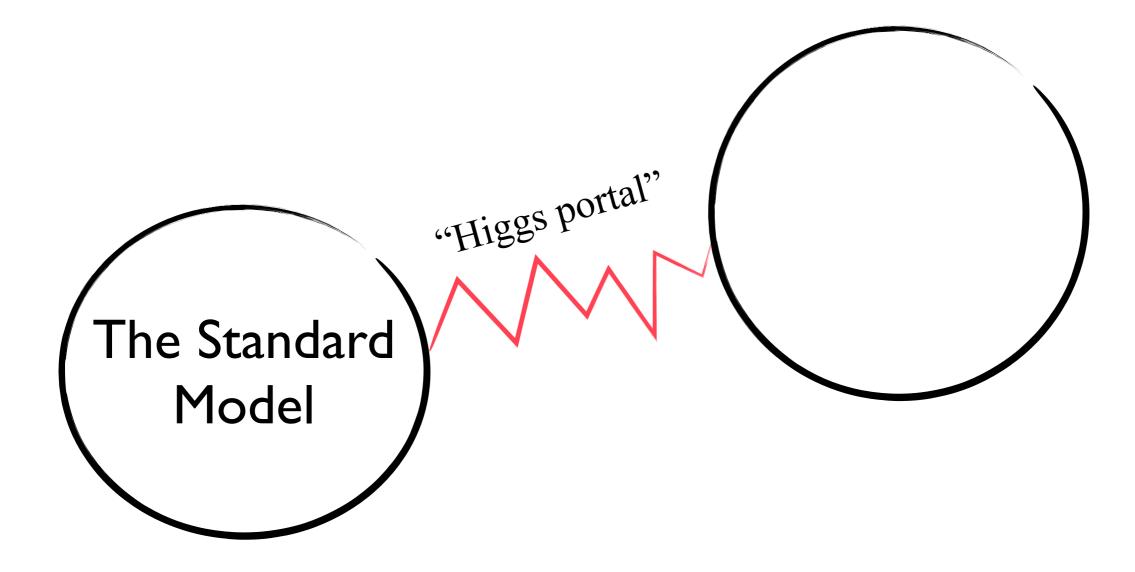
(Sannino, Tuominen 2003)

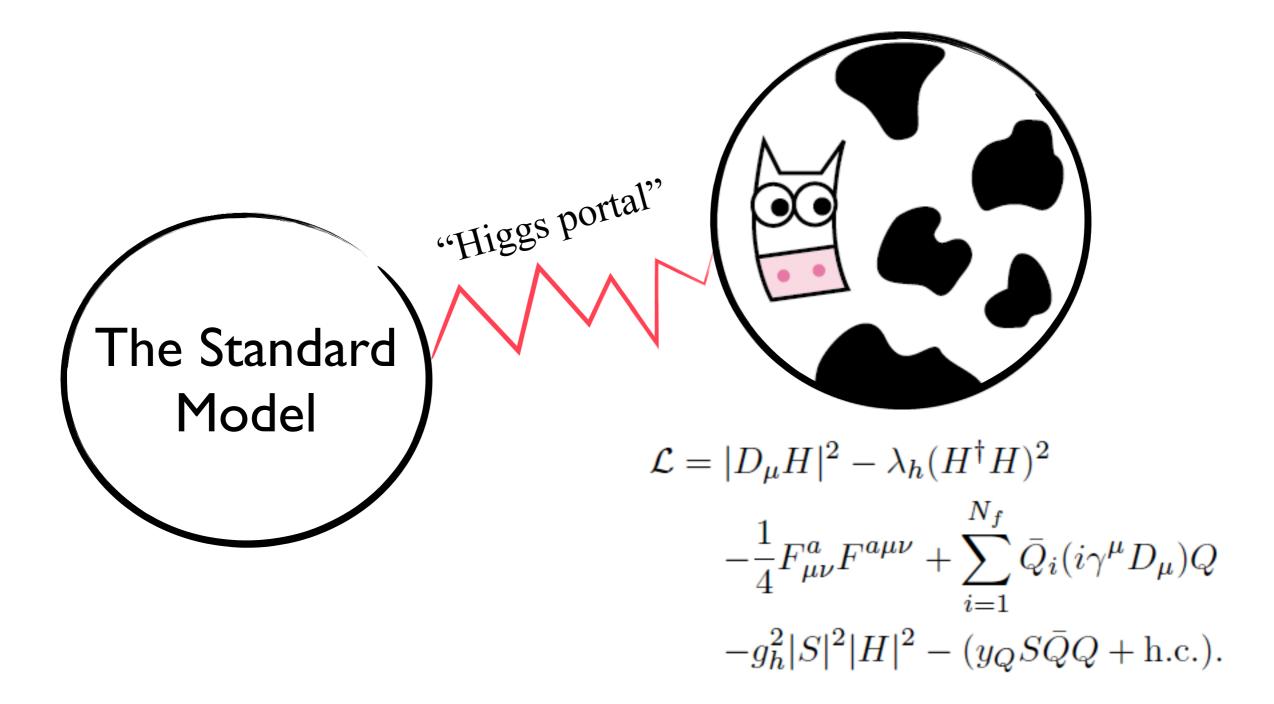


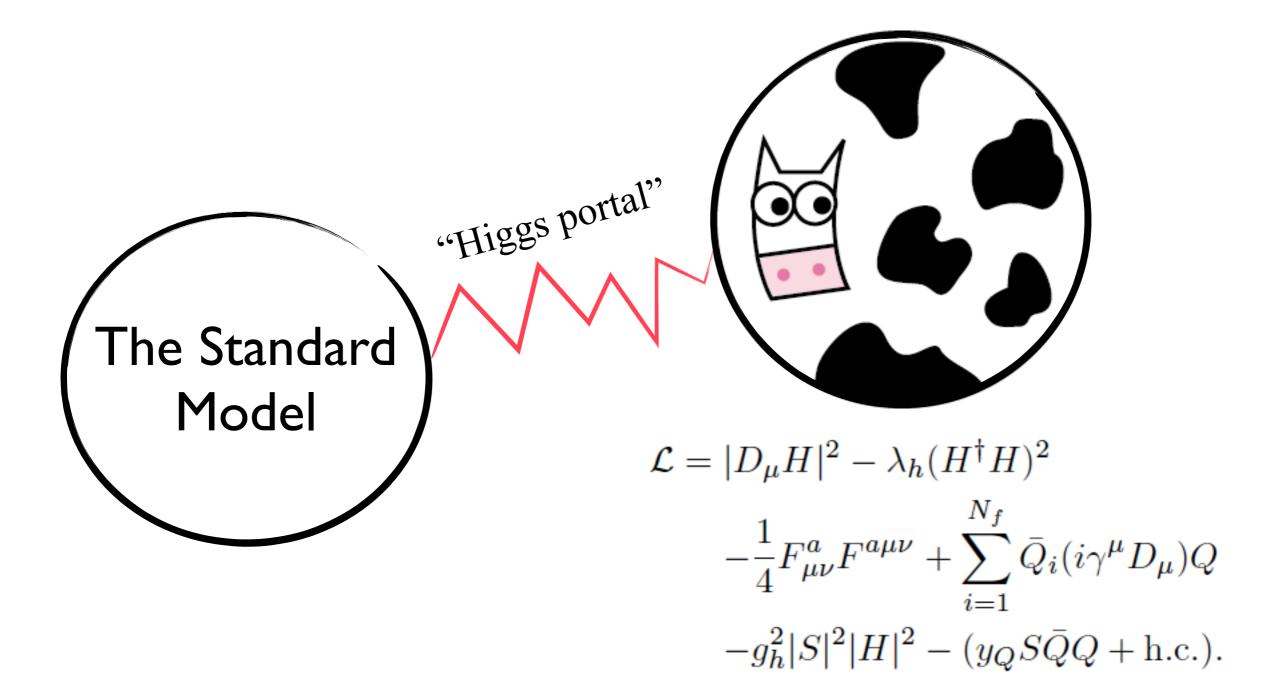
A novel possibility: complex scalar with U(1) charge & BEC. Leads to EW scale and protects it from radiative corrections.

(Sannino, Tuominen 2003)

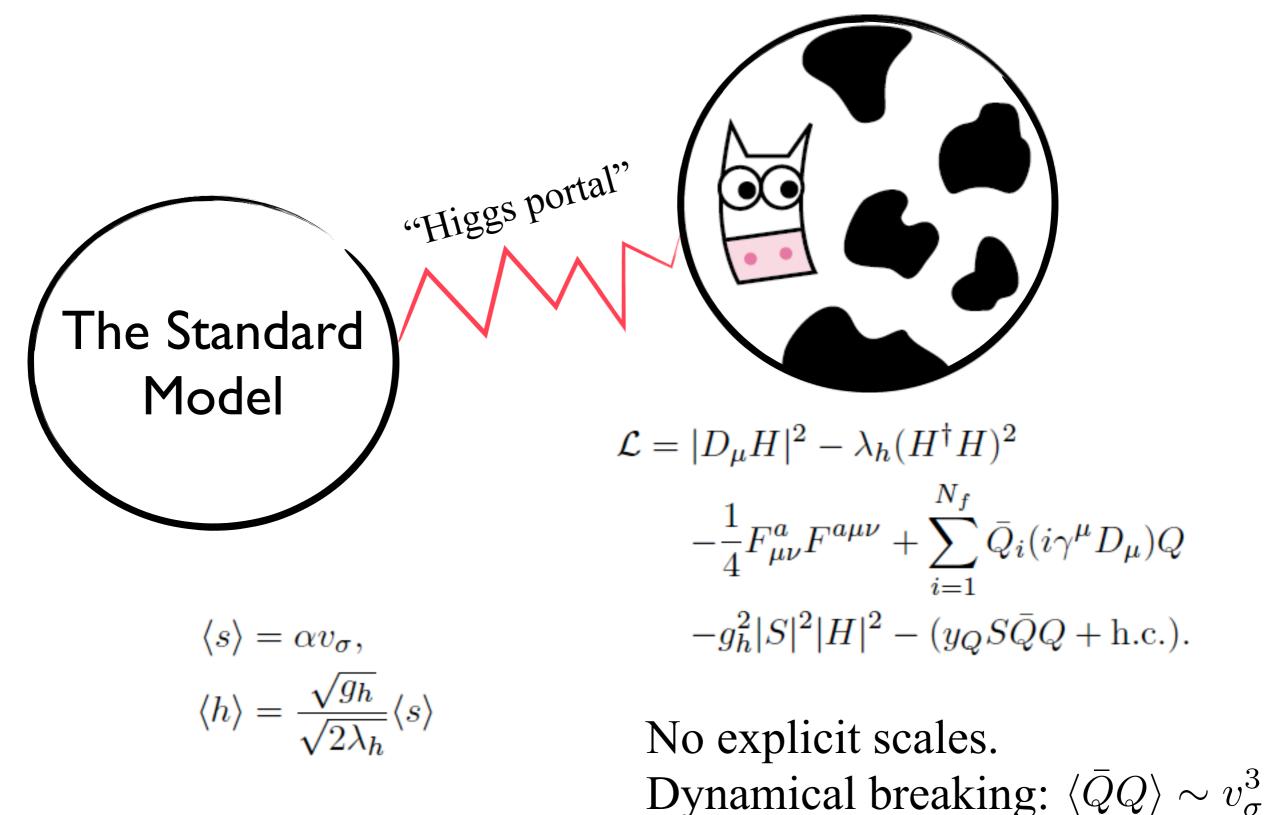
Combine this picture with scale invariance.







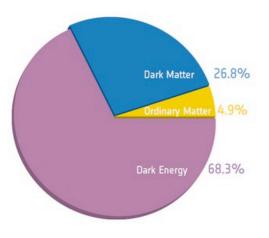
No explicit scales. Dynamical breaking:  $\langle \bar{Q}Q \rangle \sim v_{\sigma}^{3}$ (Hur & Ko, 2011; Heikinheimo et al. 2013)



(Hur & Ko, 2011; Heikinheimo et al. 2013)

Several novel dark matter candidates:

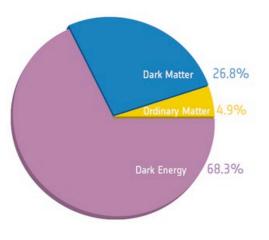
- Lightest (pseudo) goldstone boson
- Lightest dark baryon



After Planck

Several novel dark matter candidates:

- Lightest (pseudo) goldstone boson
- Lightest dark baryon

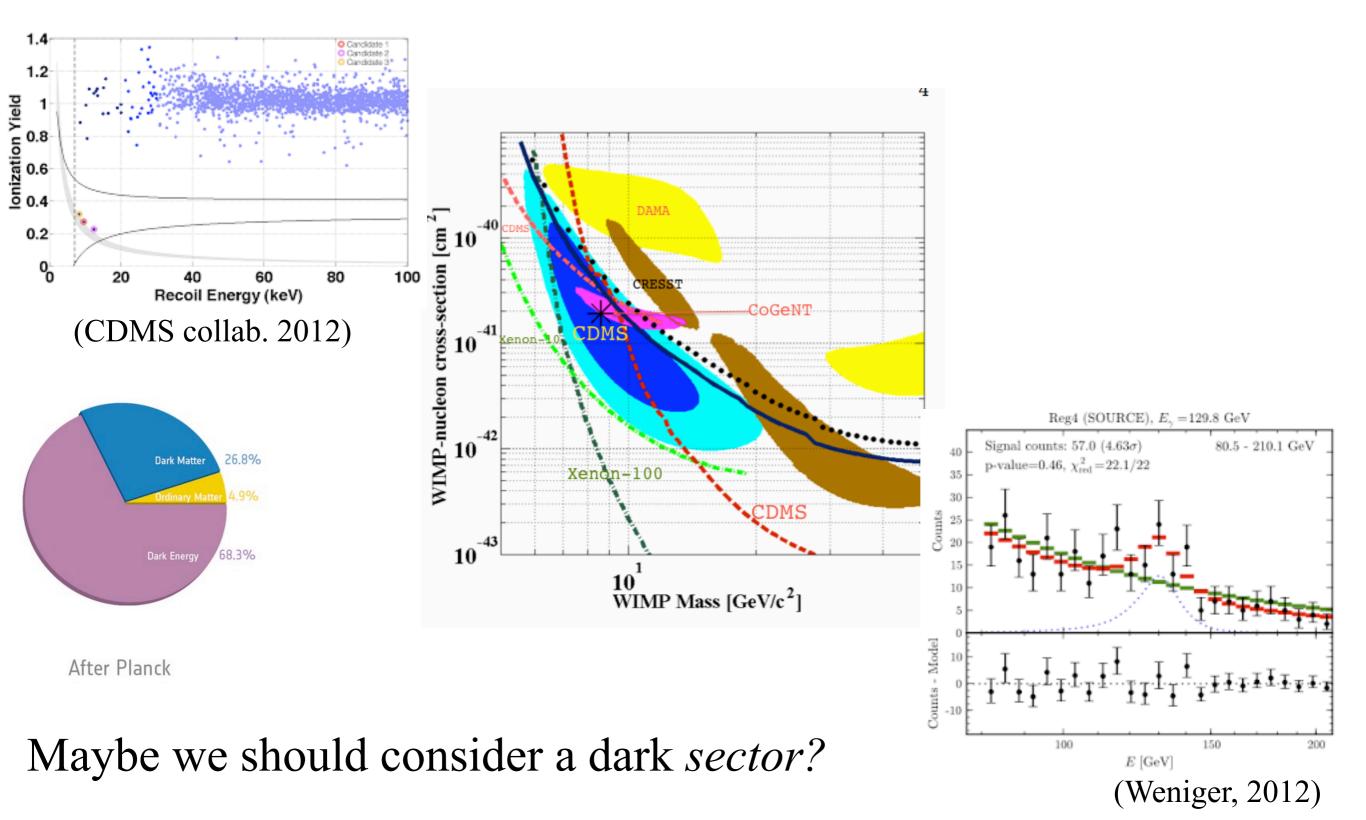


After Planck

### Maybe we should consider a dark *sector*?

Several novel dark matter candidates:

- Lightest (pseudo) goldstone boson
- Lightest dark baryon



### Conclusions

Discovery of only SM-like Higgs leads to stringent cuts in theory space.

Origin of Fermi scale remains unexplained; need new physics.

Some traditional models of DEWSB remain viable and predict states accessible to LHC.

New directions needed?; revised paradigm of naturality?; a sector of dark matter?

(Quasi) Conformal (gauge) theories expected to play a major role.