

**Higgs workshop  
Edinburgh 2013**

# **The sextet BSM model and the Higgs impostor**

**Lattice Higgs Collaboration (LHC) ->**

**with Zoltan Fodor, Kieran Holland, Daniel Negradi,  
Chris Schroeder, Chik Him Wong**

**Julius Kuti**

University of California, San Diego

**Strongly interacting dynamics beyond the Standard Model and the  
Higgs boson**

**Higgs Center workshop, University of Edinburgh, April 24-26, 2013**

# Outline

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- why the sextet model?

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- why the sextet model?
- chiral symmetry breaking sextet model
- running (walking?) gauge coupling
- dilaton as Higgs impostor?
- spectroscopy and confining force
- dark matter, finite temperature
- summary and outlook

## for more details of the discussions:

**Can the nearly conformal sextet gauge model hide the Higgs impostor?**

Zoltan Fodor, Kieran Holland, Julius Kuti, Daniel Negradi, Chris Schroeder, Chik Him Wong

Published in **Phys.Lett. B718 (2012) 657-666**

**The sextet gauge model, light Higgs, and the dilaton**

Zoltan Fodor, Kieran Holland, Julius Kuti, Daniel Negradi, Chris Schroeder, Chik Him Wong

Published in **PoS LATTICE2012 (2012) 024**

**Confining force and running coupling with twelve fundamental and two sextet fermions**

Zoltan Fodor, Kieran Holland, Julius Kuti, Daniel Negradi, Chris Schroeder, Chik Him Wong

Published in **PoS LATTICE2012 (2012) 025**

Zoltan Fodor, Kieran Holland, Julius Kuti, Daniel Negradi, Chik Him Wong

**preliminary new results, published soon**

**why the sextet model?**

# LATTICE GAUGE THEORIES AT THE ENERGY FRONTIER

Thomas Appelquist, Richard Brower, Simon Catterall, George Fleming,  
Joel Giedt, Anna Hasenfratz, Julius Kuti, Ethan Neil, and David Schaich

(USQCD Collaboration)

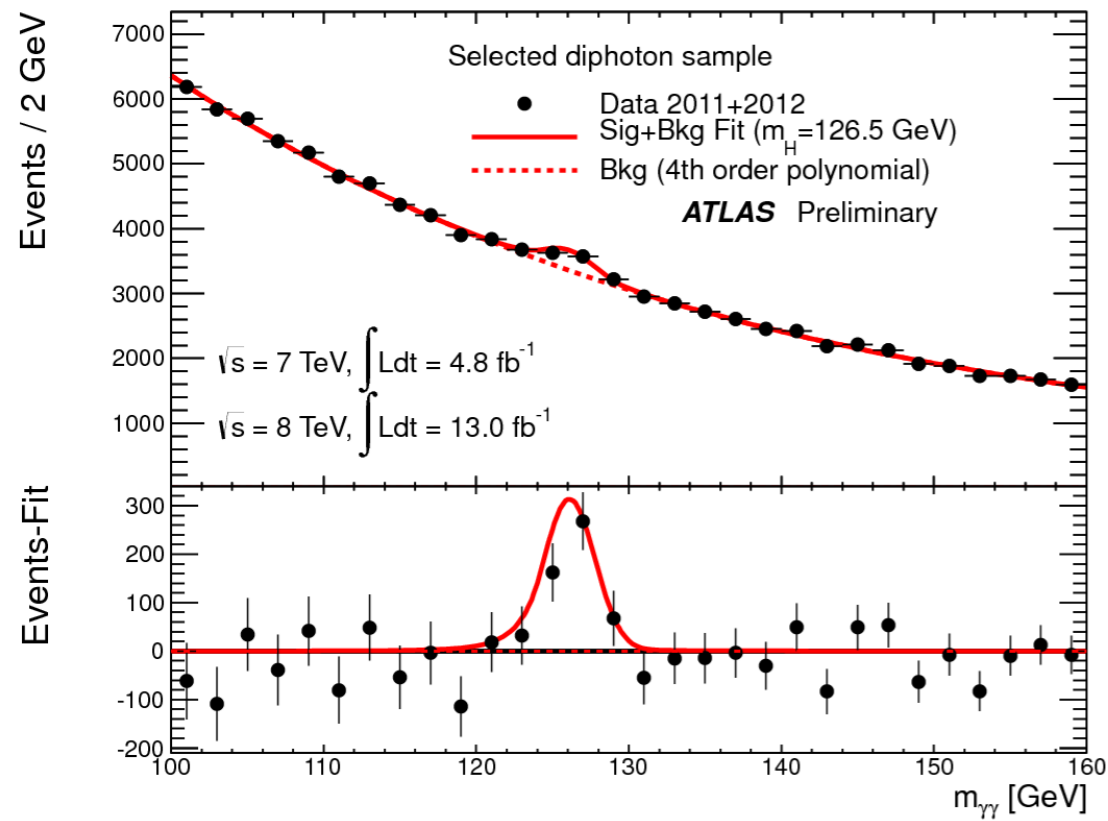
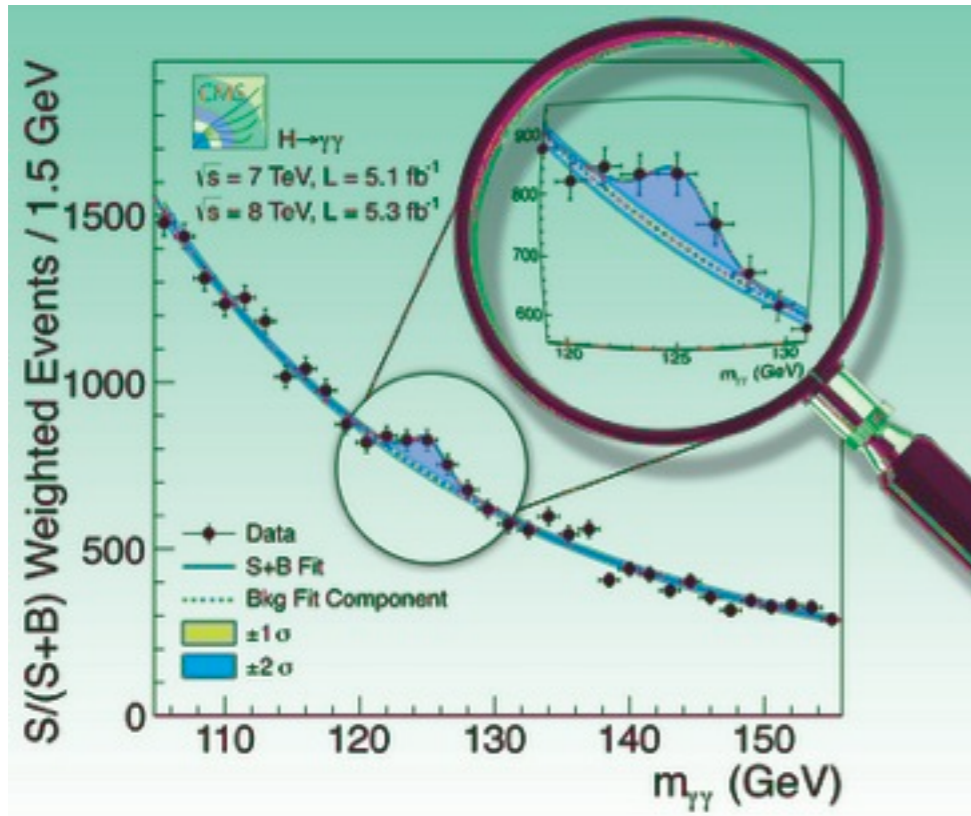
(Dated: March 10, 2013)

## White paper - USQCD BSM community based effort :

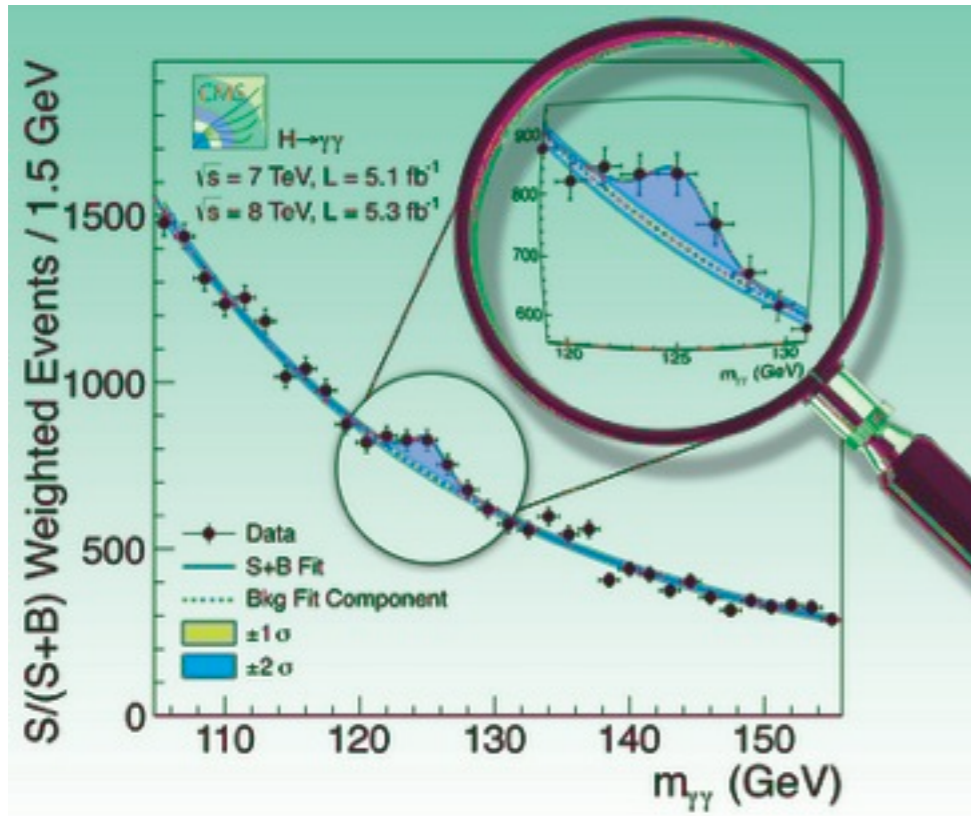
- identify most significant accomplishments of last few years
- identify our three major research directions for planning
- describe the toolset and its phenomenological applications
- estimate resources needed for the plan

New hardware proposal of USQCD submitted to DOE  
with three main directions

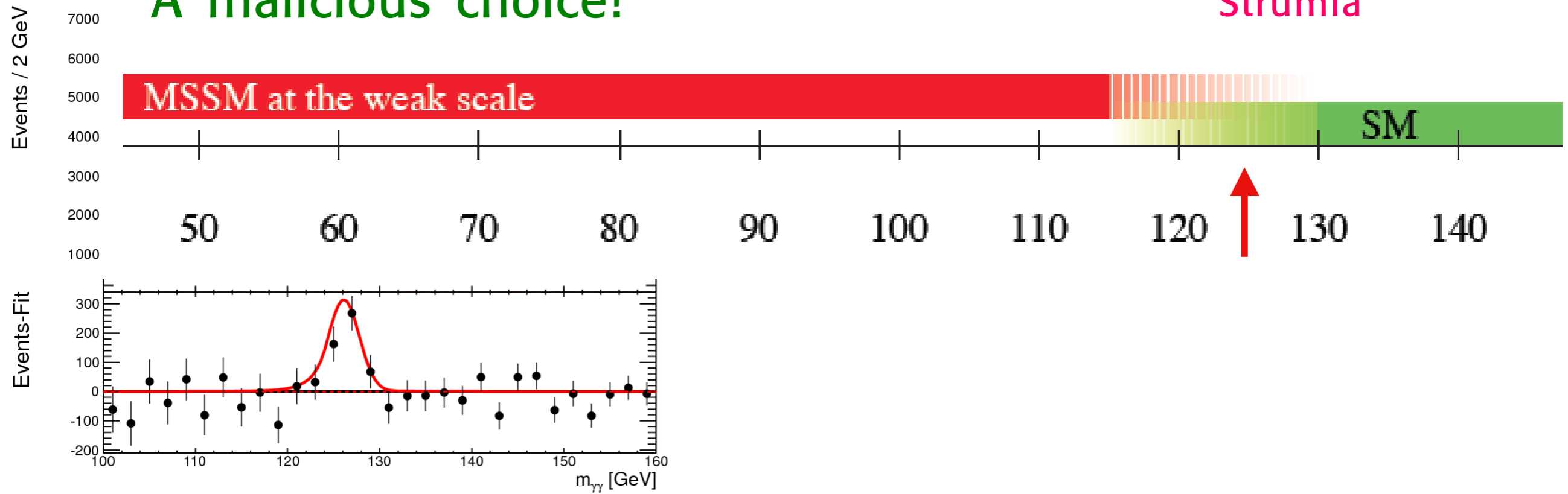
- two plots effected our planning no kidding :)
- where is the Higgs? no more asked  
what is it made of? asked now all the time
- “Mexican hat” solution parametrization rather than dynamical explanation?
- has fine tuning and hierarchy problems



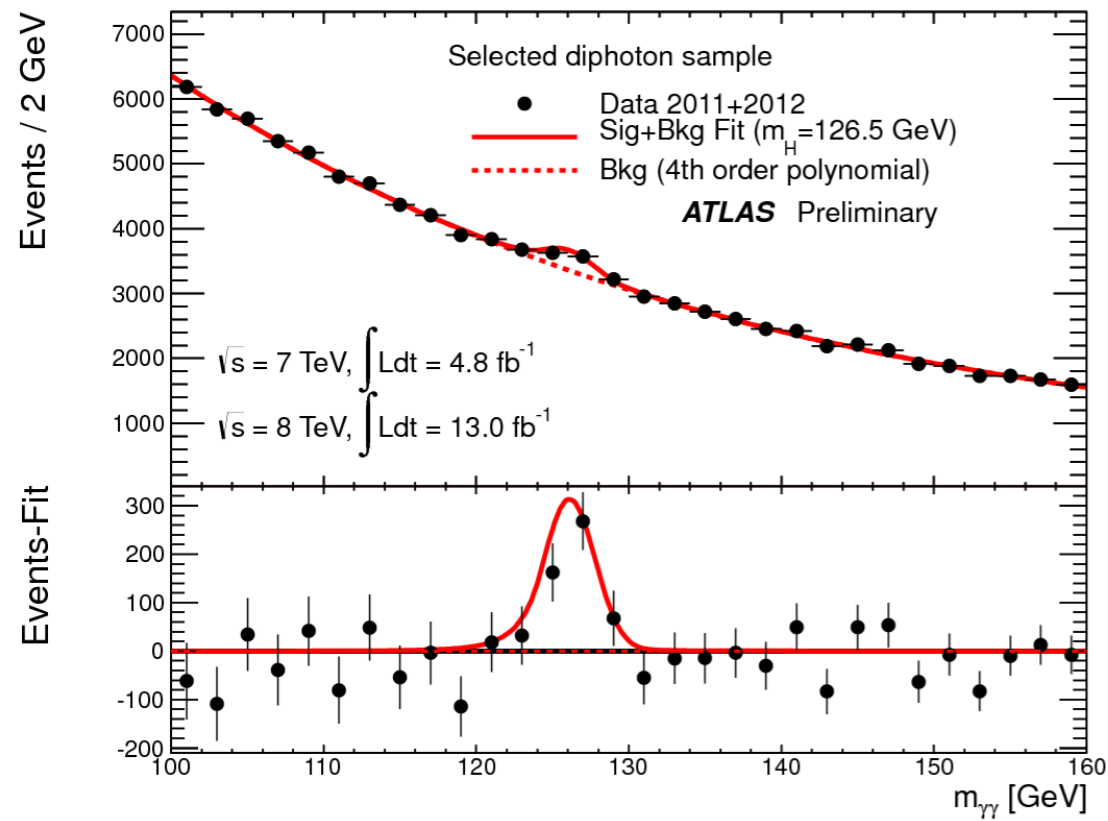
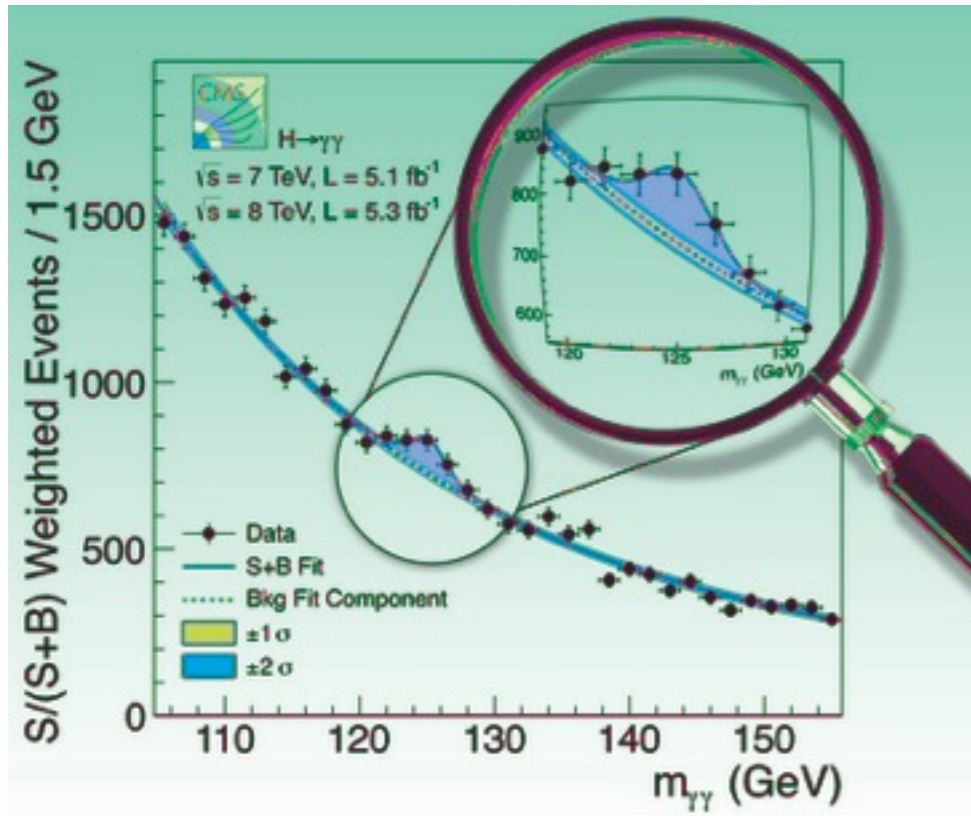
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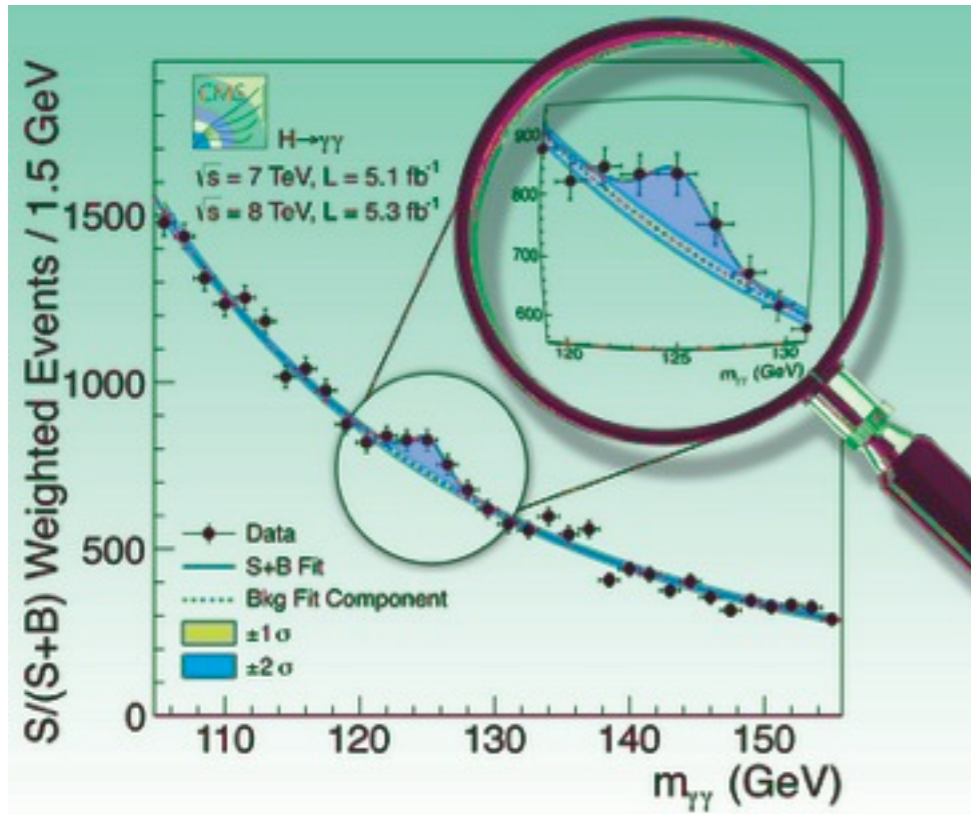
A malicious choice!



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- has fine tuning and hierarchy problems

- three BSM directions to do better:

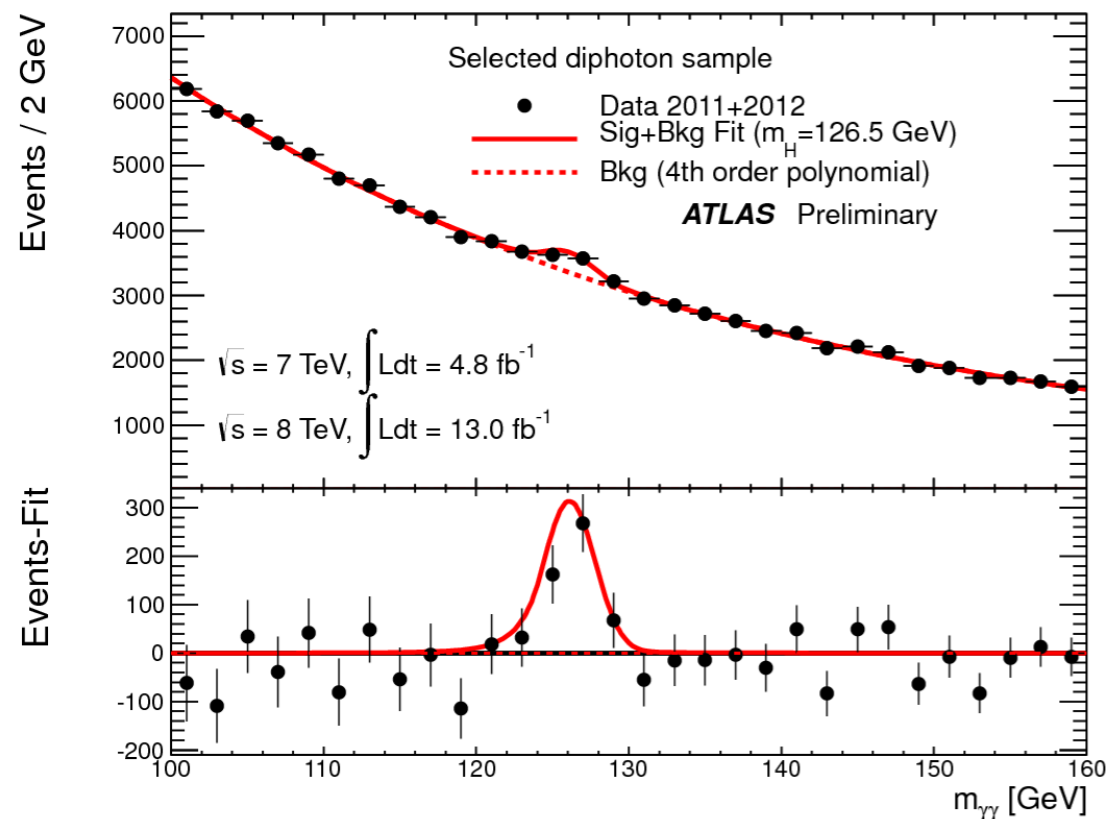
sextet model

- strongly coupled near-conformal gauge theories

- light pseudo-Goldstone boson (like little Higgs)

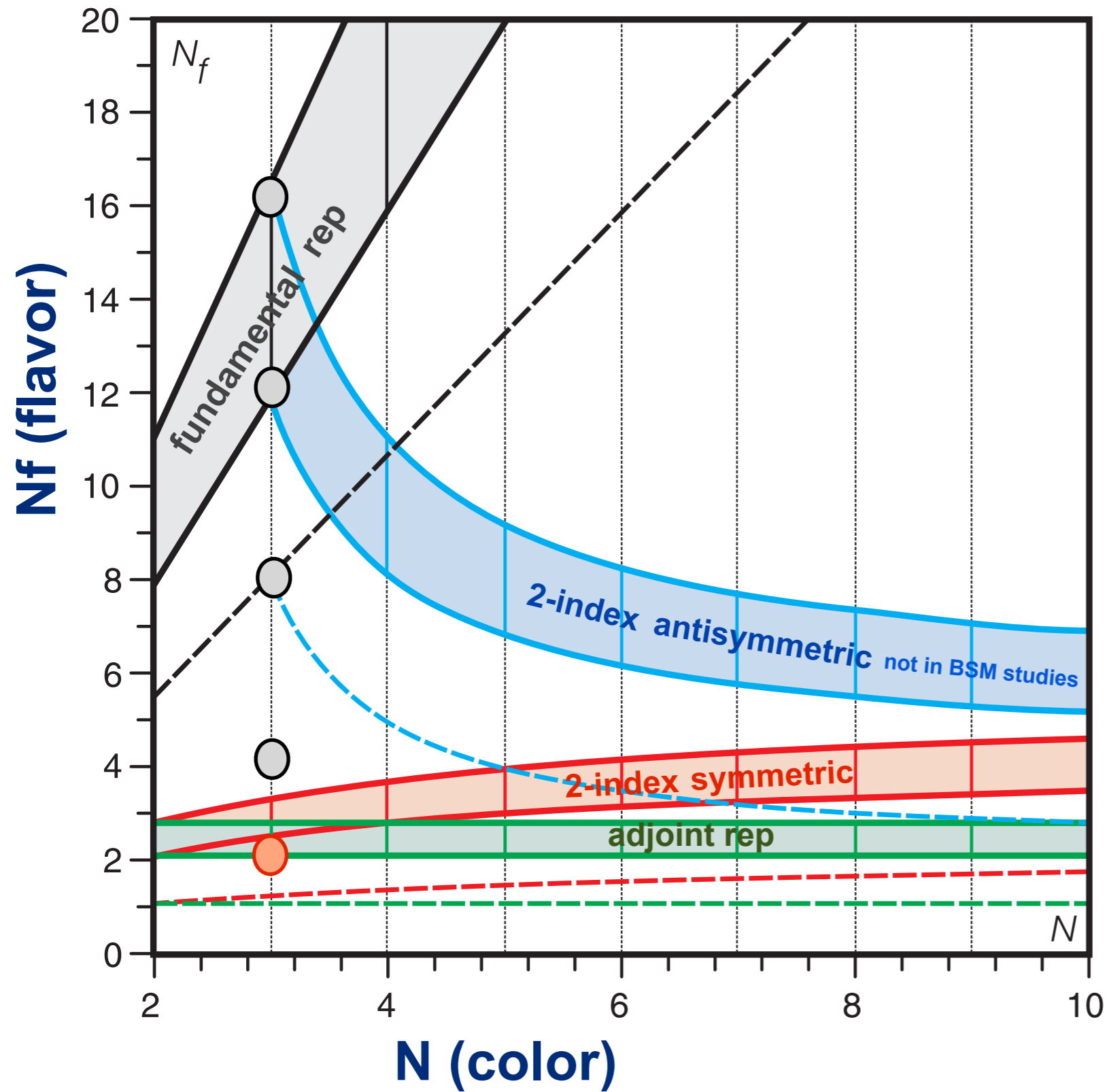
- SUSY

- new physics with little or no tuning within LHC reach, or hiding just above LHC reach?

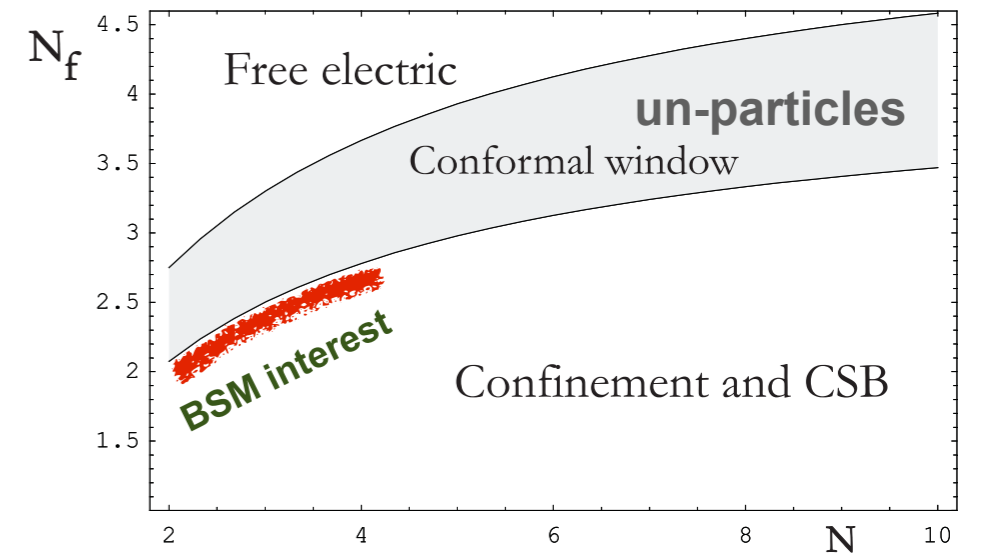




**theory space and conformal window**  
 color, flavor, and massless fermion representation

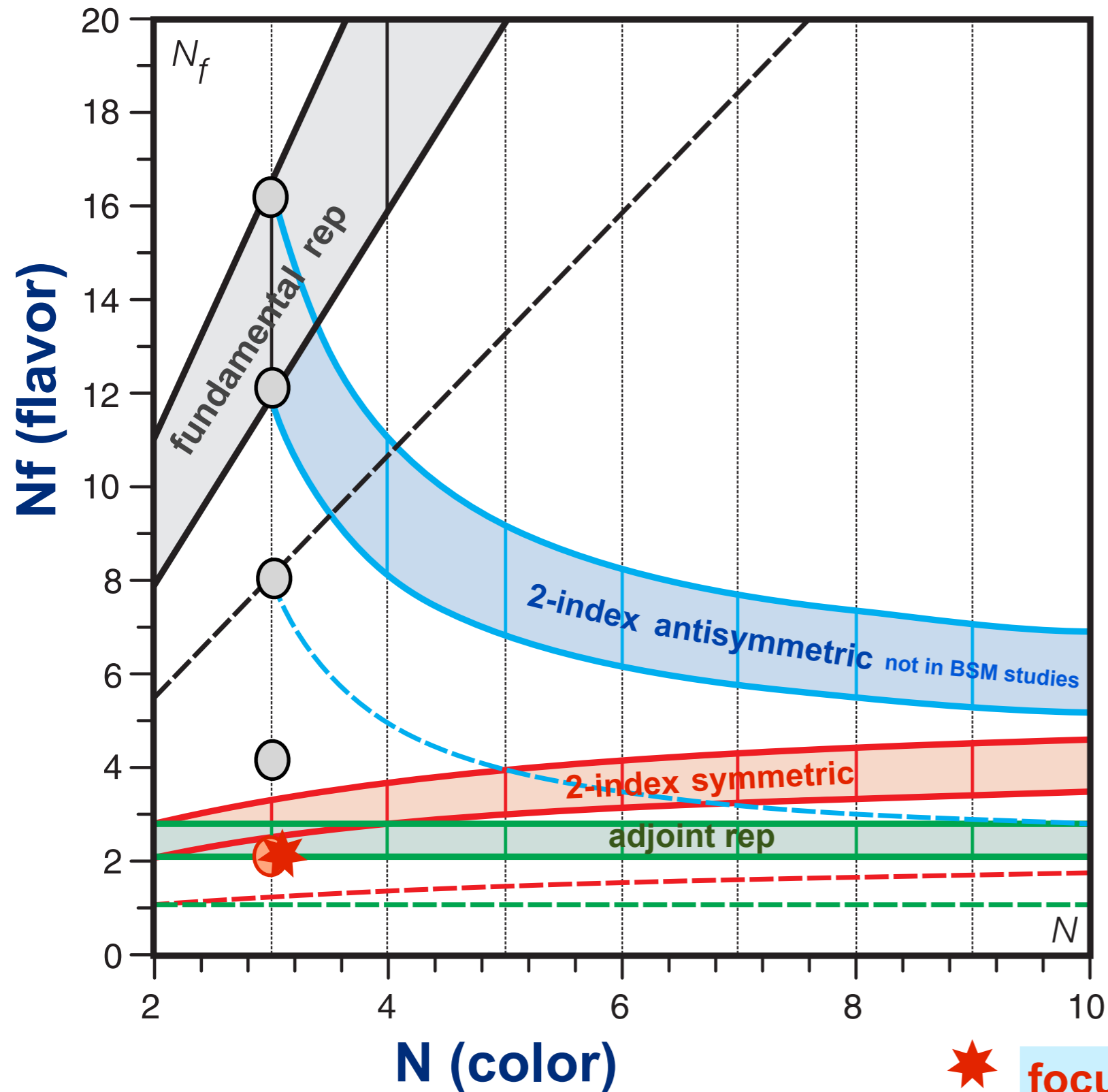


for each rep BSM interest is below conformal window but close to it:



physics prospects of sextet look promising ?

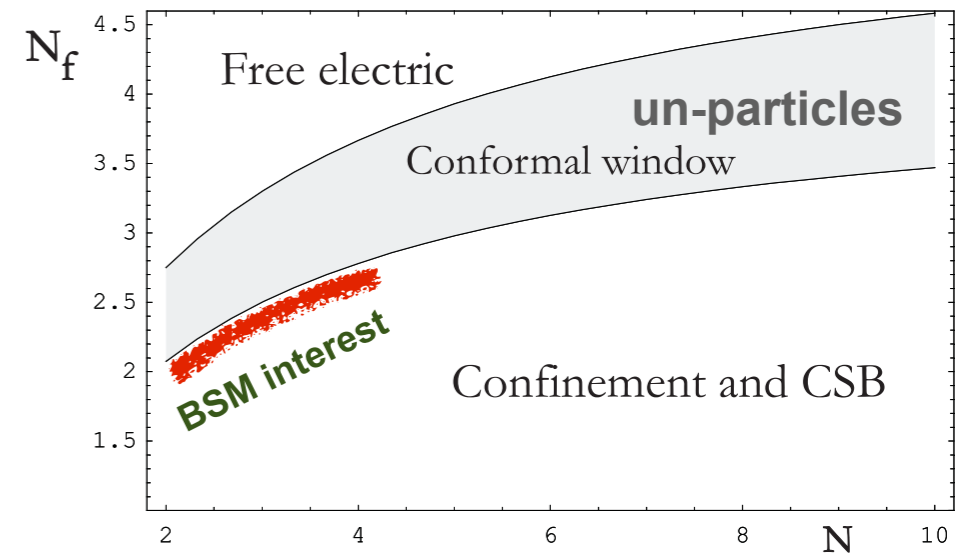
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★ focus of talk

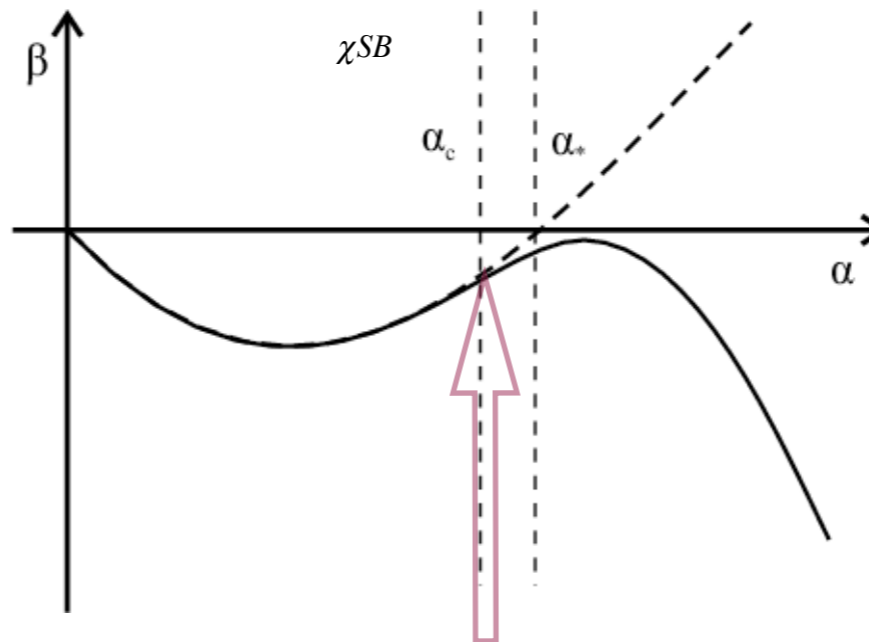
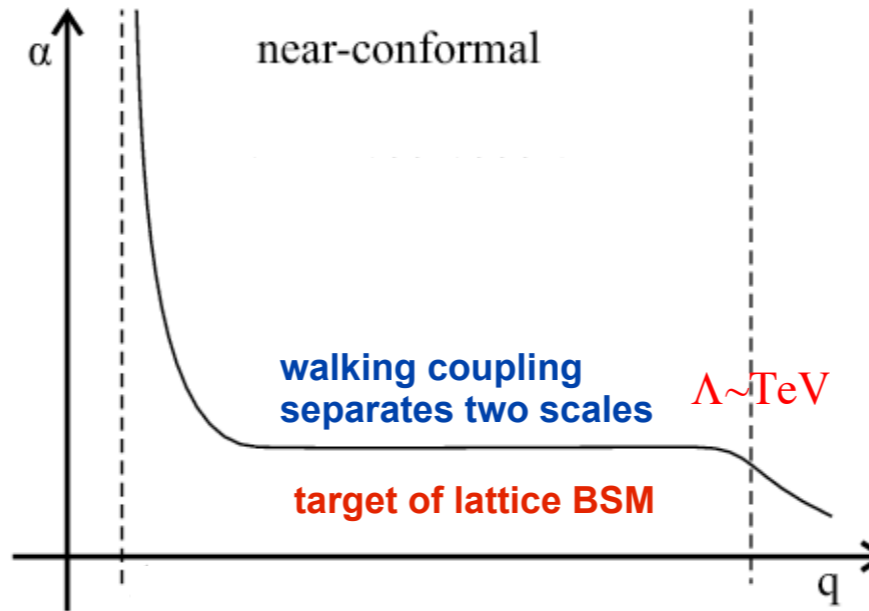
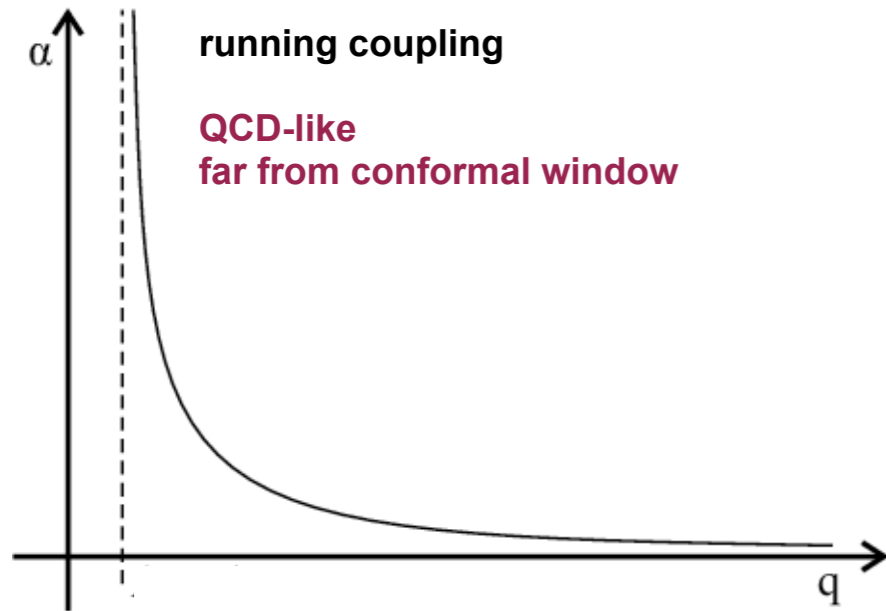
New extended data set and analysis

for each rep BSM interest is below conformal window but close to it:



physics prospects of sextet look promising ?

# simplest realization of composite Higgs: $N_f=2$ $SU(3)$ sextet representation



$\chi_{SB}$  on  $\Lambda \sim \text{TeV}$  scale

walking gauge coupling

fermion mass generation not addressed

what composite Higgs mechanism ?

broken scale invariance (dilaton) and/or light non-SM composite Higgs particle?

Early work using sextet rep:

Marciano (QCD paradigm, 1980)

Kogut, Shigemitsu, Sinclair (quenched, 1984)

recent work:

DeGrand, Shamir, Svetitsky IRFP or walking gauge coupling

Lattice Higgs Collaboration  $\chi_{SB}$

Kogut, Sinclair finite temperature

original Technicolor paradigm replaced with sextet  $SU(3)$  color rep:

one massless fermion doublet  $\chi_{SB}$  on  $\Lambda \sim \text{TeV}$  scale

$$\begin{bmatrix} u \\ d \end{bmatrix}$$

three Goldstone pions become longitudinal components of weak bosons

composite Higgs mechanism scale of Higgs condensate  $\sim F=250$  GeV

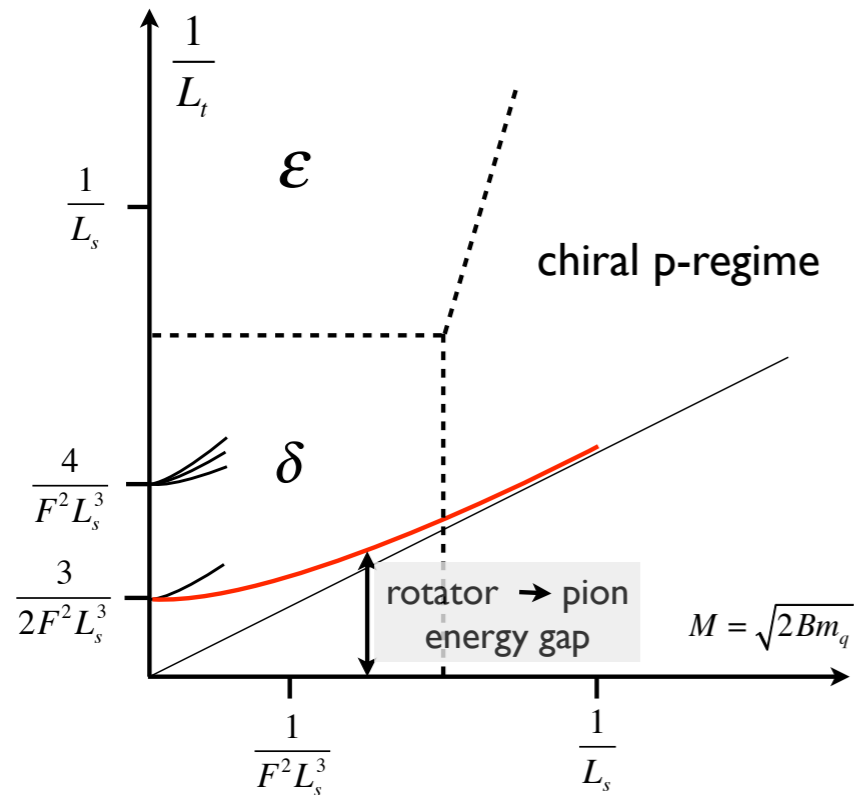
conflicts with EW precision constraints?

sextet model with two critical requirements: (1)  $\chi_{SB}$  (2) walking gauge coupling?

role of the third fermion flavor?

**chiral symmetry breaking in the sextet model**

# mass deformed chiral SB in finite volume below conformal window:

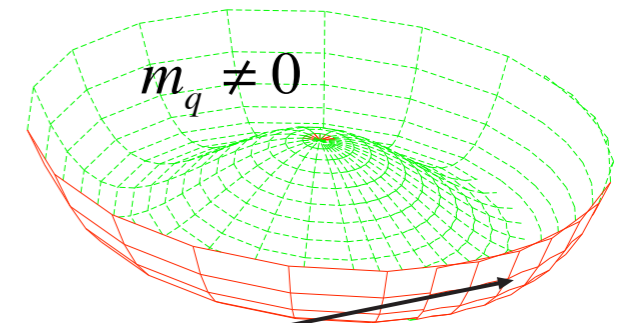
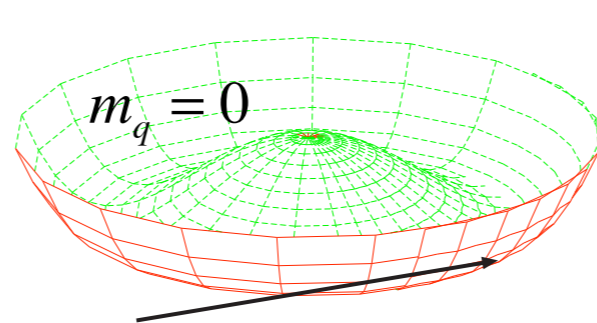


Goldstone dynamics is different in each regime

We study  $\delta$  and  $\epsilon$ -regimes (RMT) and p-regime (probing chiral loops)

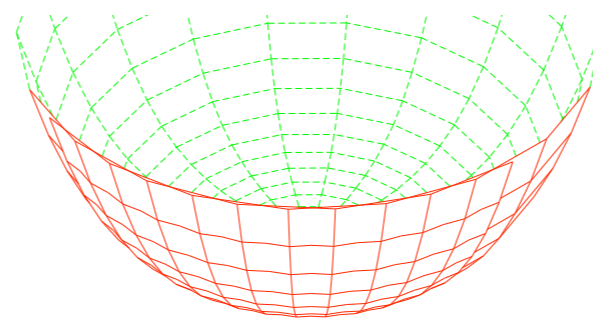
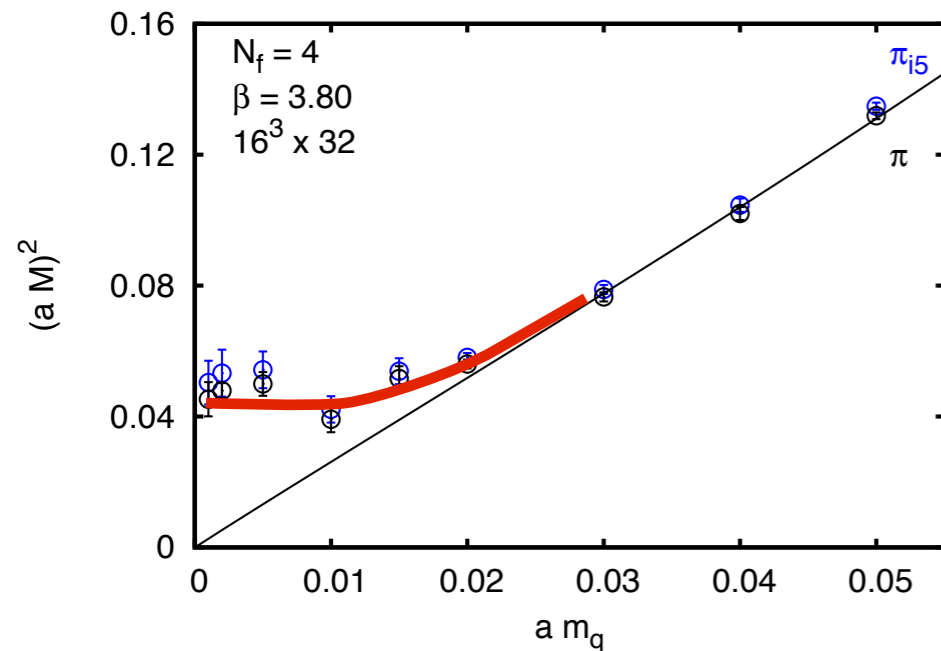
complement each other

interpretation of rotator levels in  $m_q \rightarrow 0$  limit:



$V_{\text{eff}}$ : chiral condensate in flavor space  
arbitrary orientation of condensate

tilted condensate



Not to misidentify rotator gaps as evidence of chirally symmetric phase !

Our sextet simulations are in the p-regime  $\beta = 3.2$  and  $\beta = 3.25$

## simulation details:

tree level improved Symanzik gauge action;  $\beta = 6/g^2$  normalization

smearing in staggered fermions: 2 stout steps

rooting with two flavors (follow-up work without rooting if model will pass first tests)

## RHMC

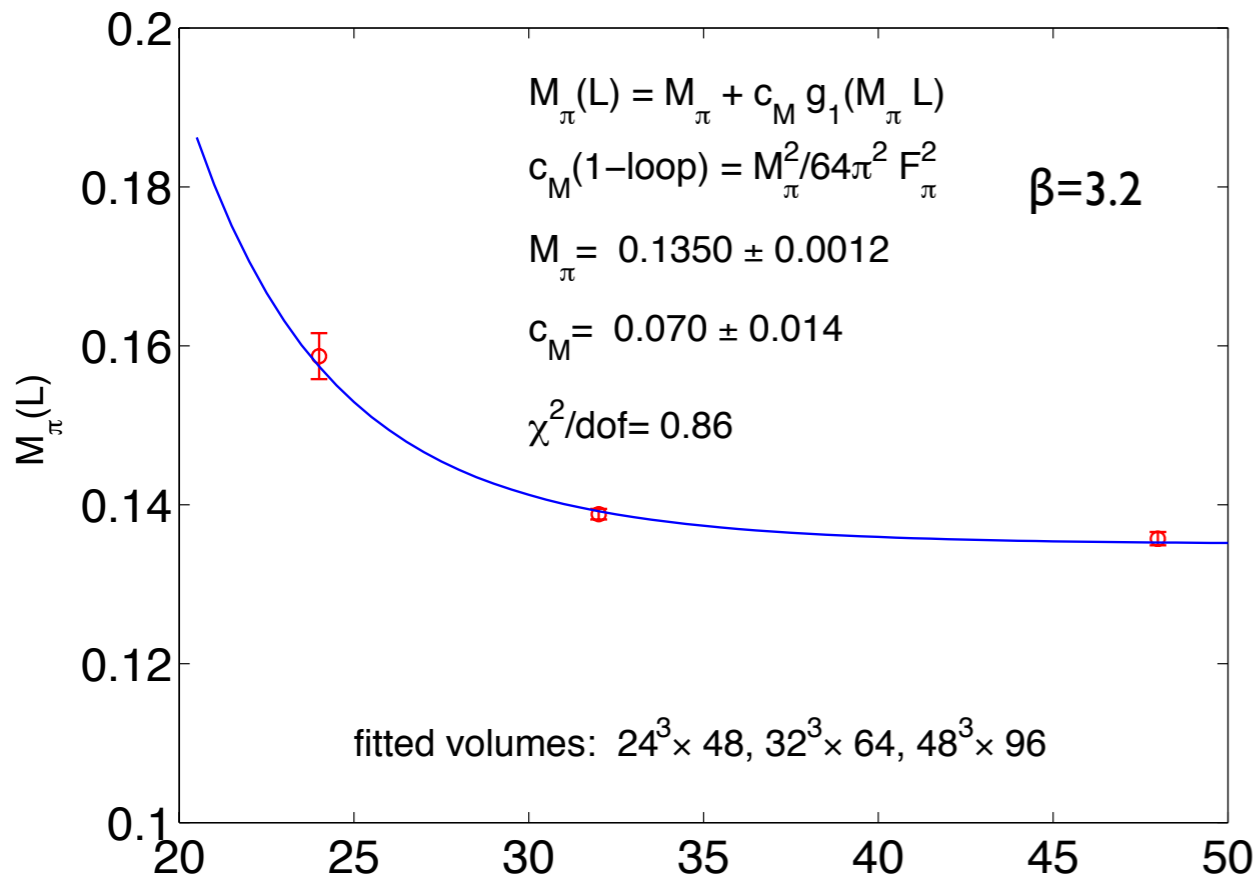
multiple time scales and Omelyan integrator

$\beta = 3.20$   $m = 0.003 - 0.010$  mass range  $24^3 \times 48$ ,  $28^3 \times 56$ ,  $32^3 \times 64$ ,  $48^3 \times 96$  lattices

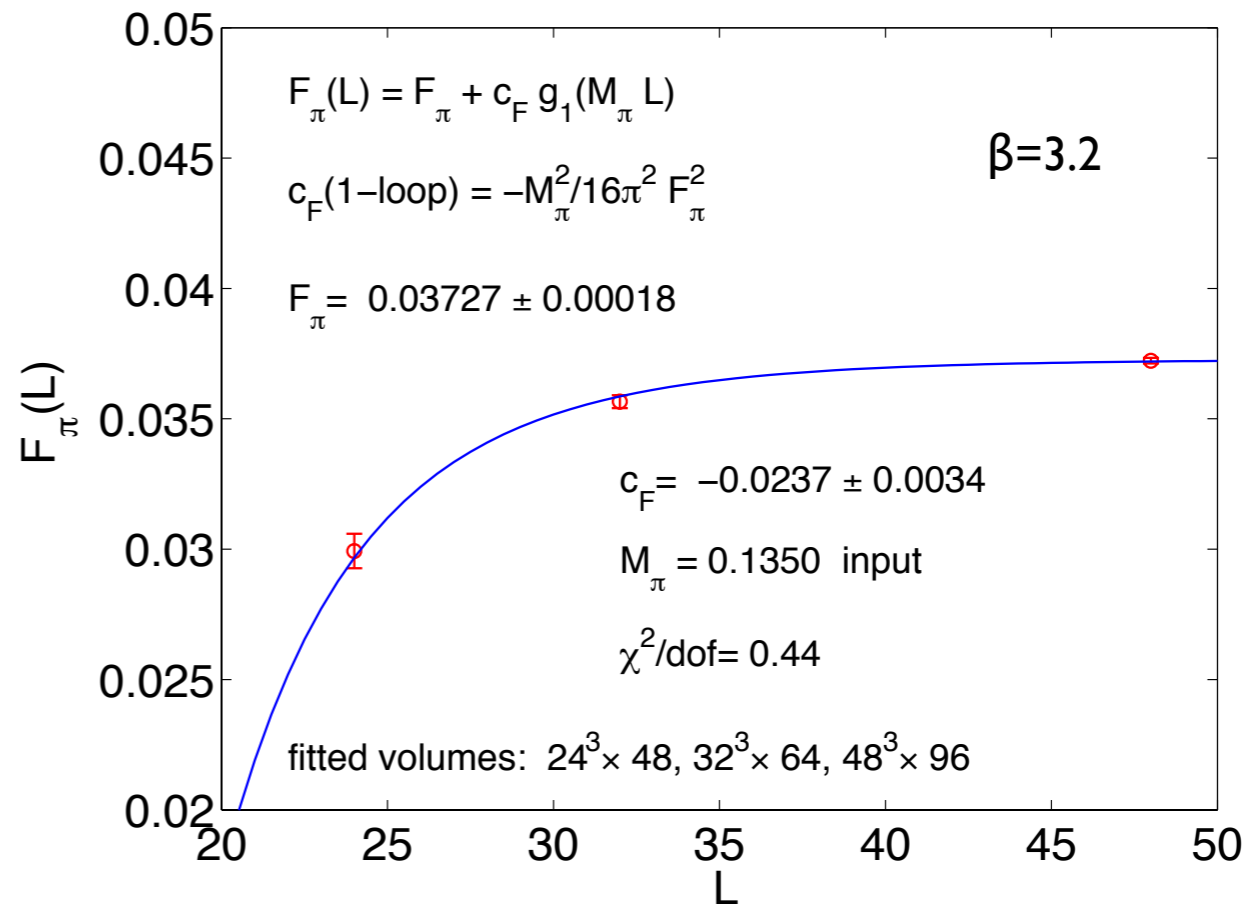
$\beta = 3.25$   $m = 0.004 - 0.008$  mass range  $24^3 \times 48$ ,  $28^3 \times 56$ ,  $32^3 \times 64$  lattices

error analysis: mass fits with double Jackknife procedure on covariance matrices

$M_\pi$   $\beta = 3.2$   $m=0.003$



$F_\pi$   $\beta = 3.2$   $m=0.003$



**strategy:  $L \rightarrow \infty$  extrapolation first**

**mass-deformed theory**

**close to  $m=0$  critical surface**

**$L \rightarrow \infty$  extrapolated chiral and conformal**

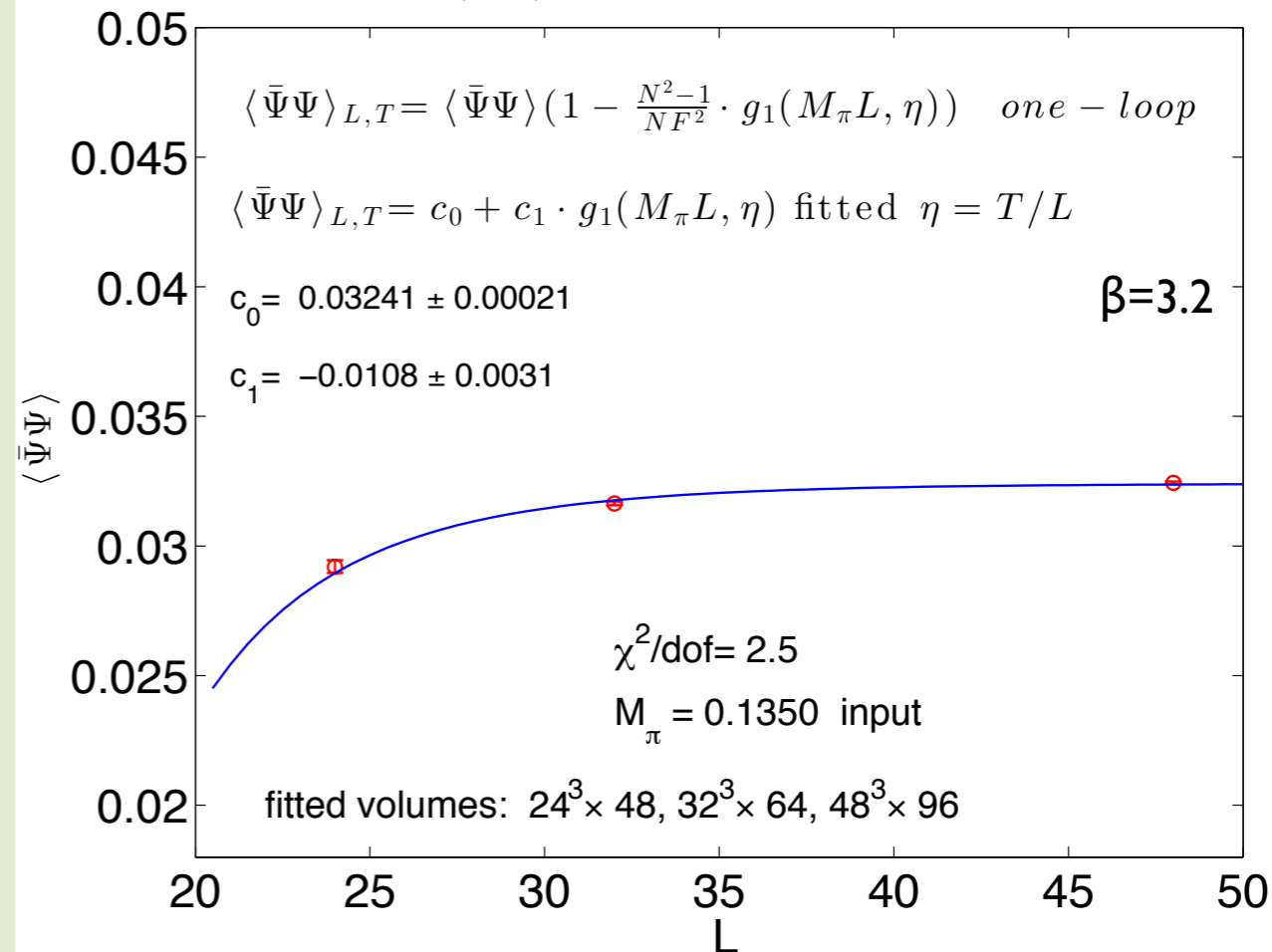
**scaling tests in sextet model**

**for  $L \cdot M_\pi > 5$  less than one percent**

**$L$  correction left**

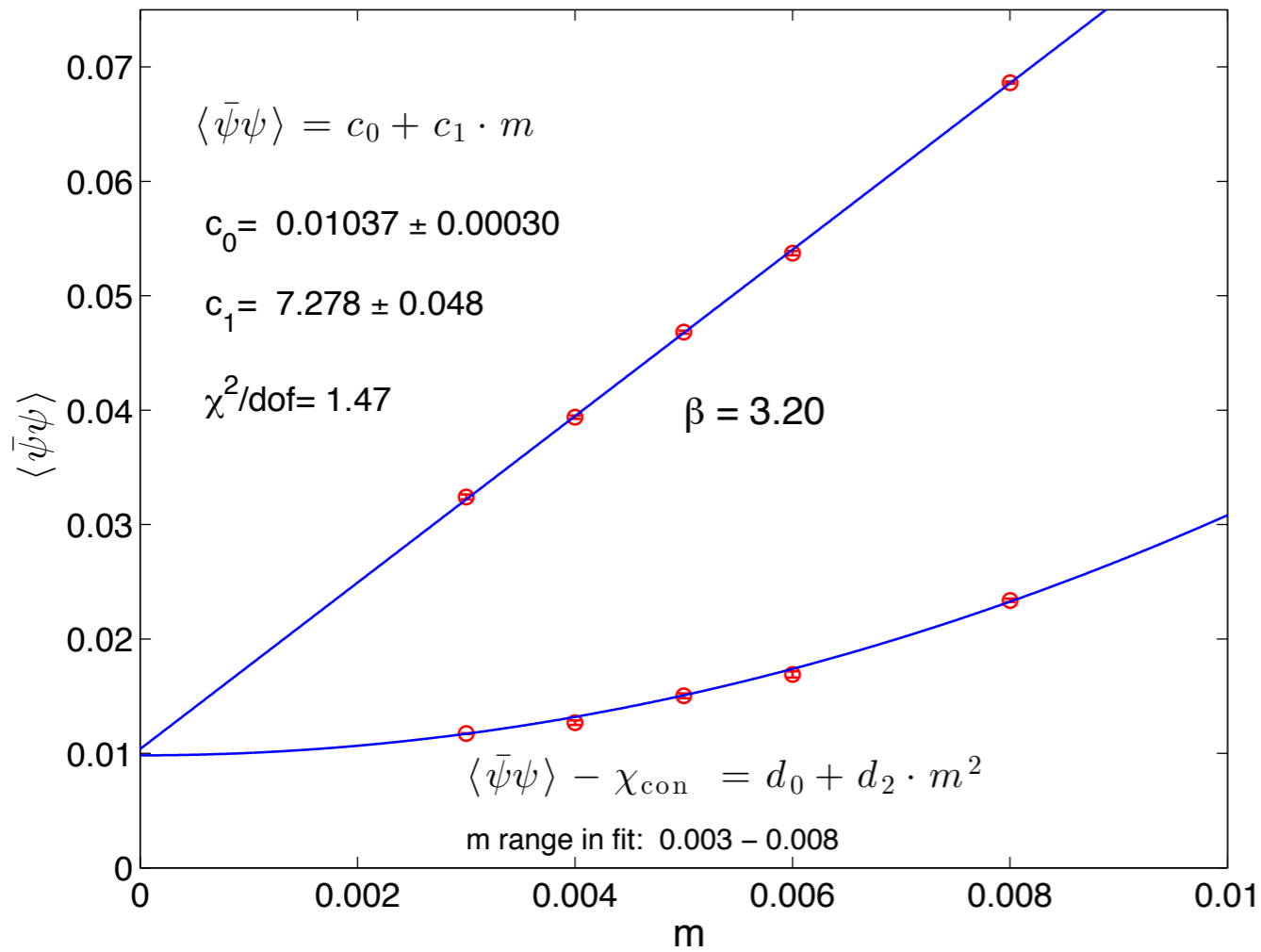
**$C_M$  and  $C_F$  signs correct, numerically off**

$\langle \bar{\Psi}\Psi \rangle$   $m=0.003$  finite volume fit

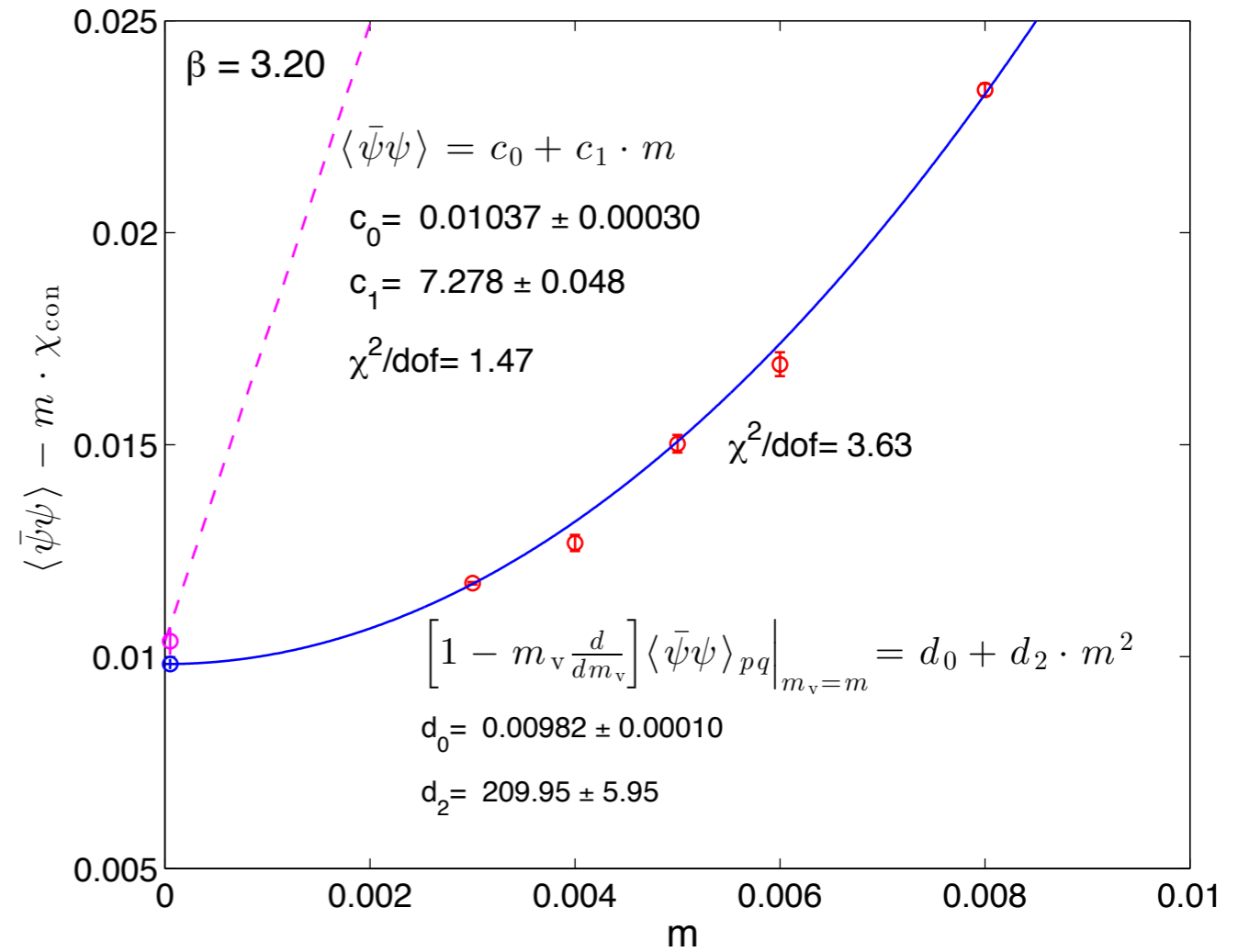


# Nf=2 SU(3) sextet chiral condensate

chiral condensate and its subtracted form



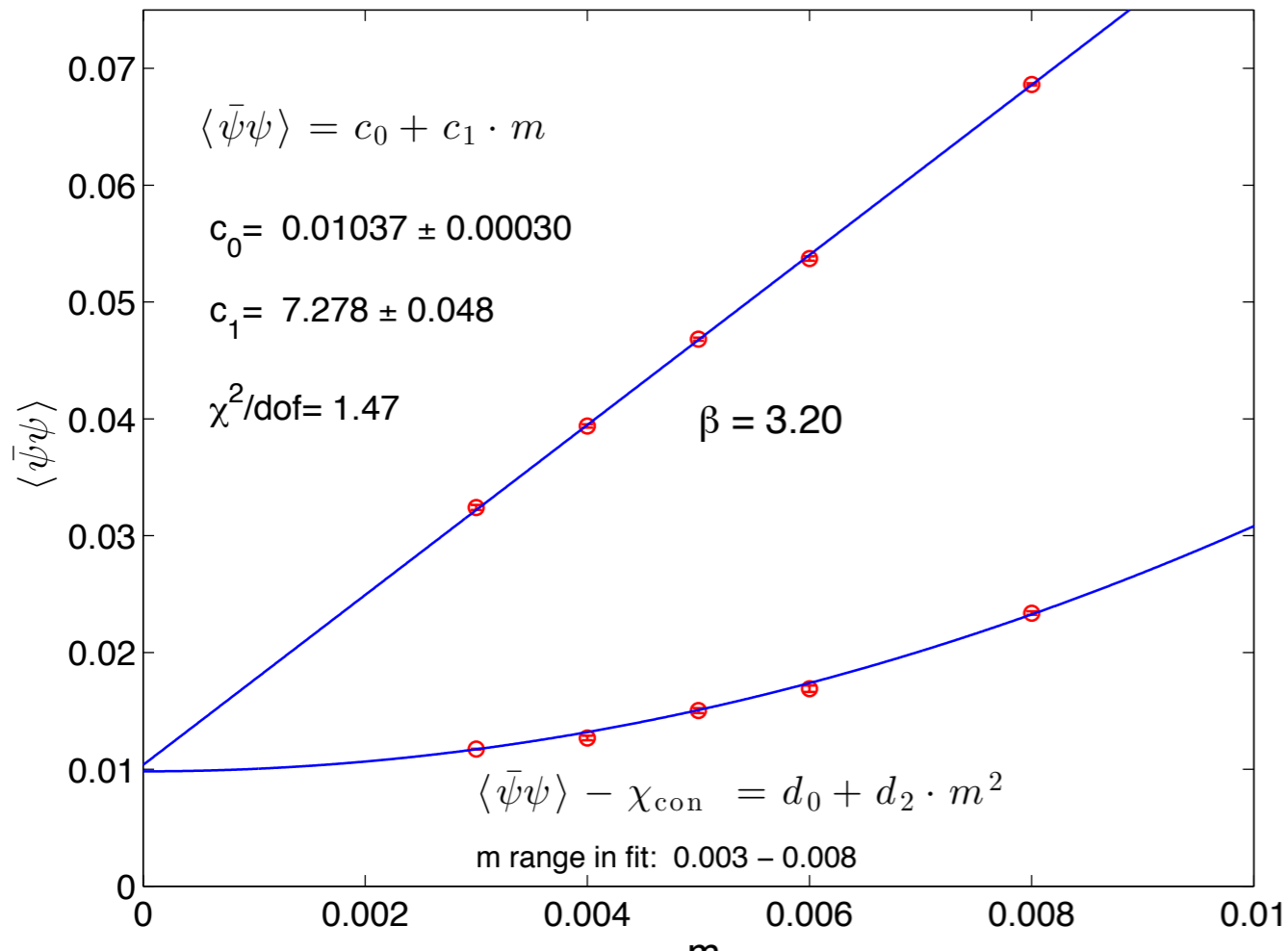
subtracted chiral condensate



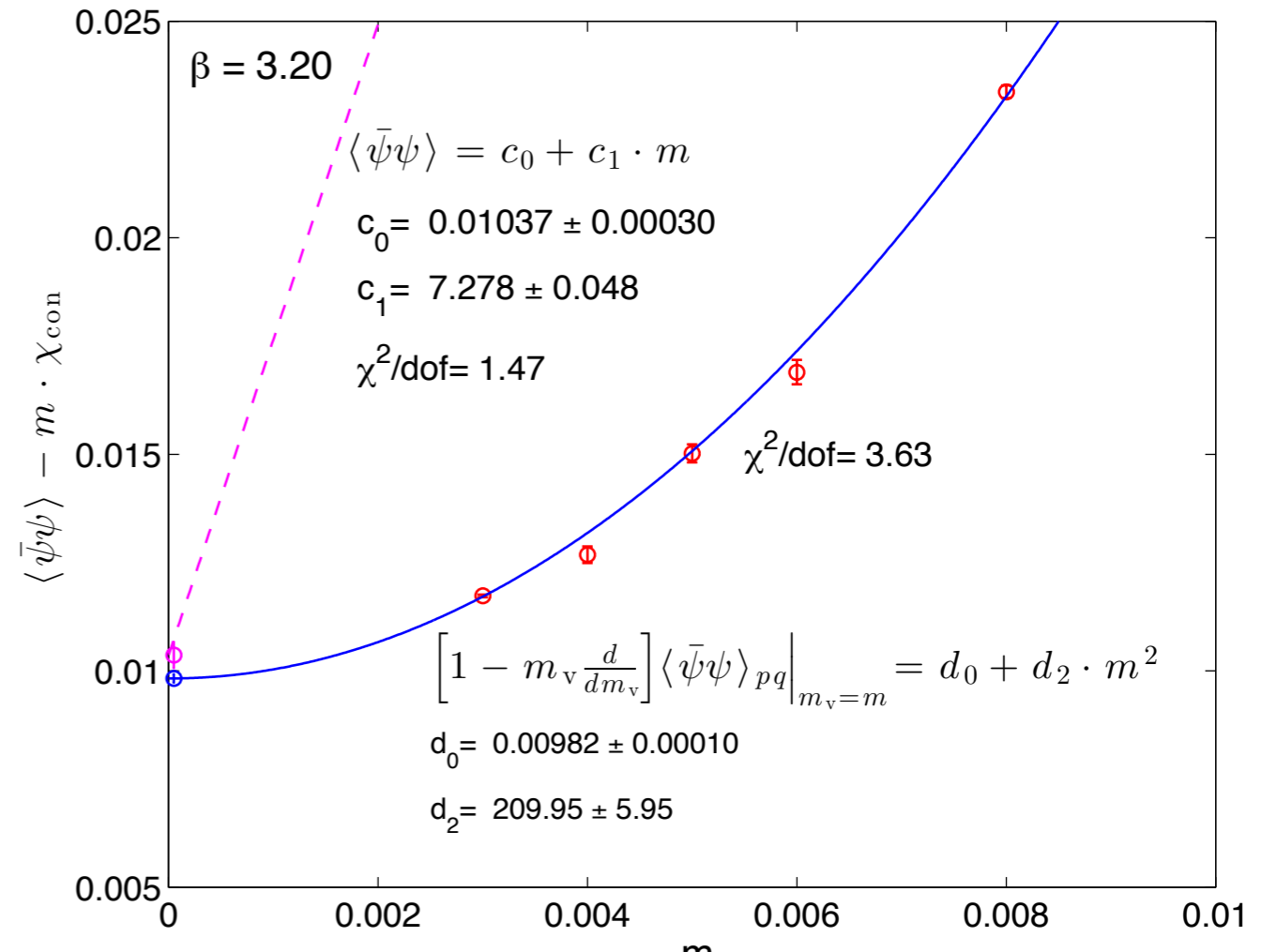


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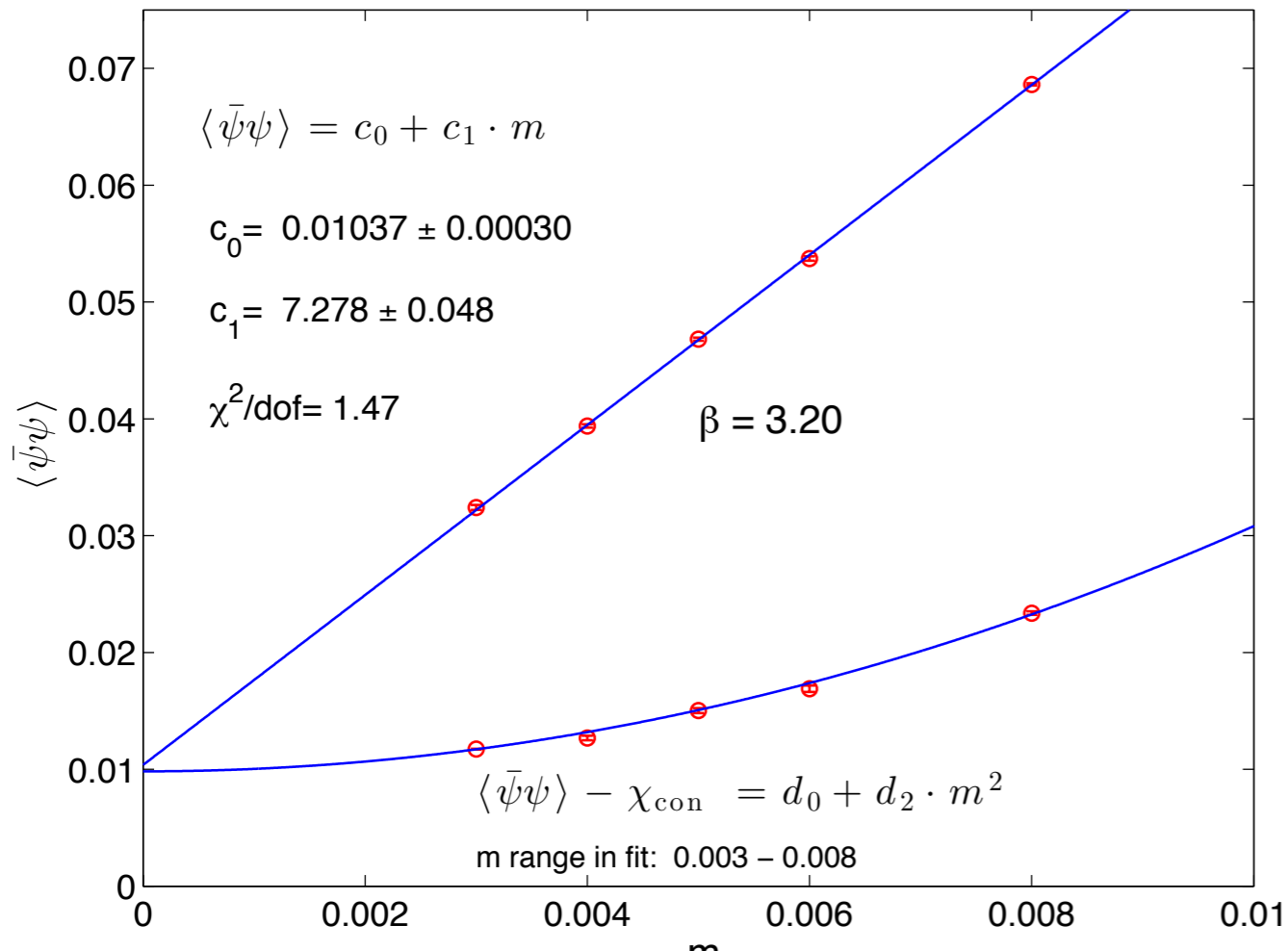
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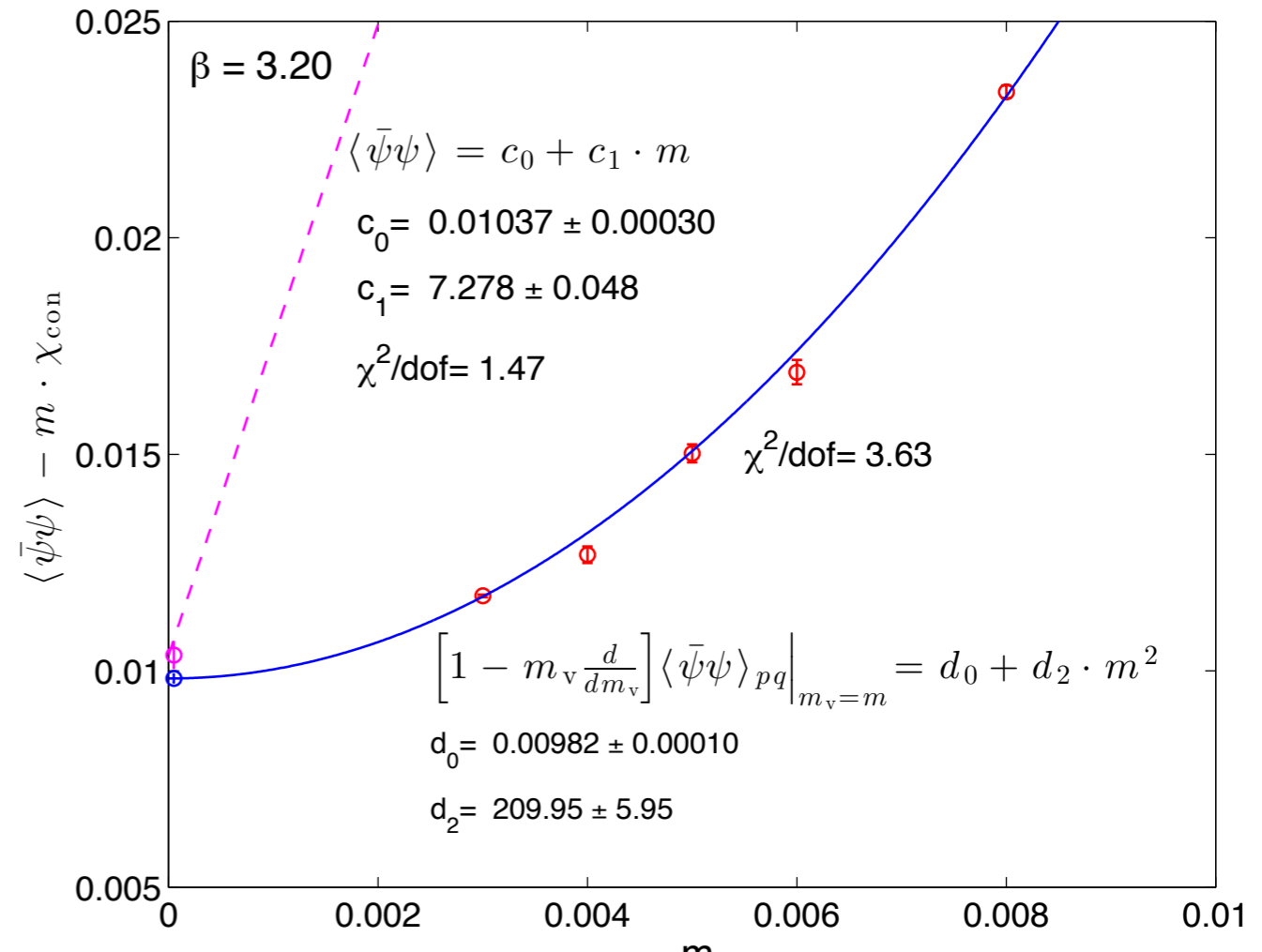
two independent determinations of the chiral condensate

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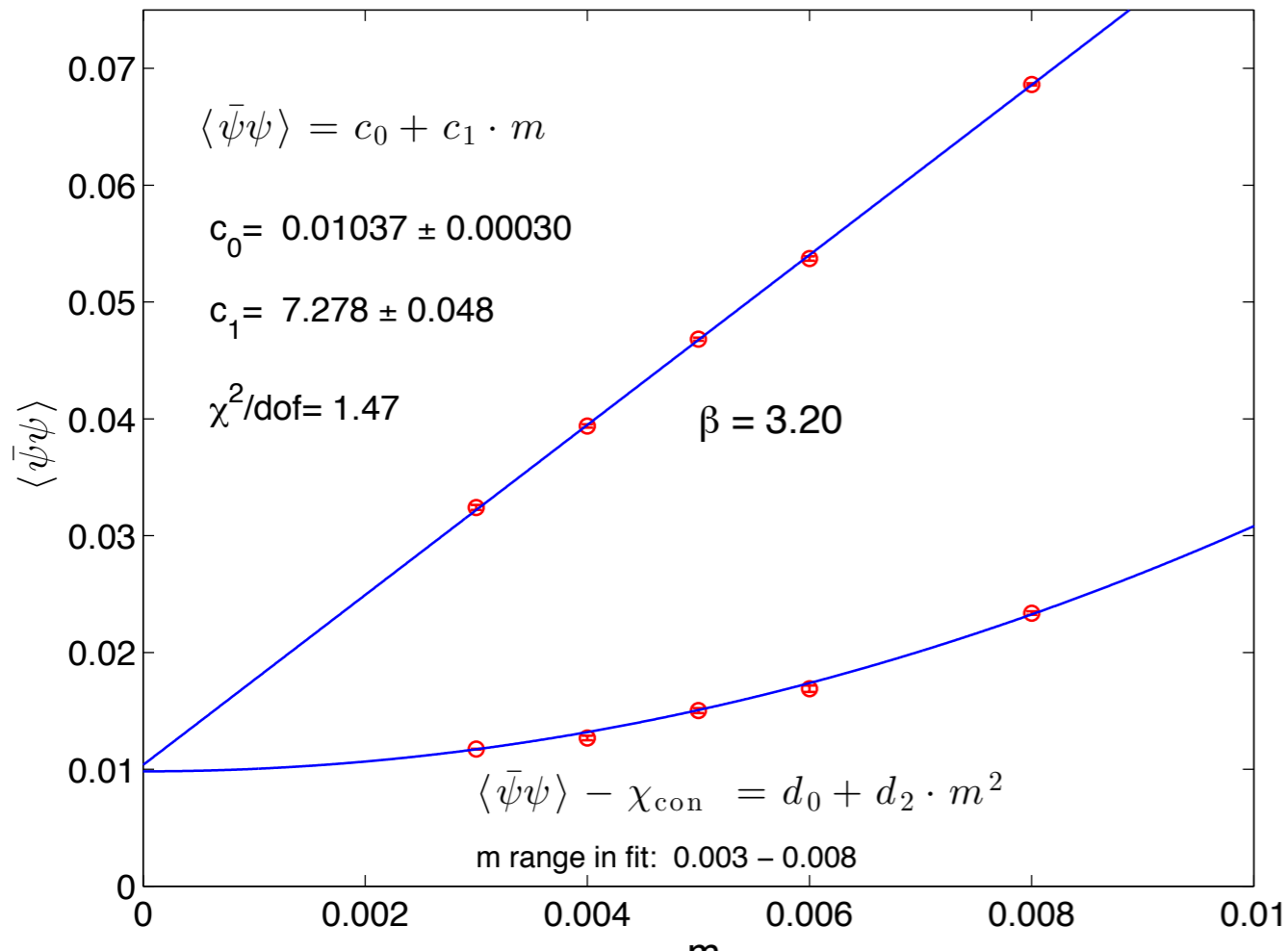
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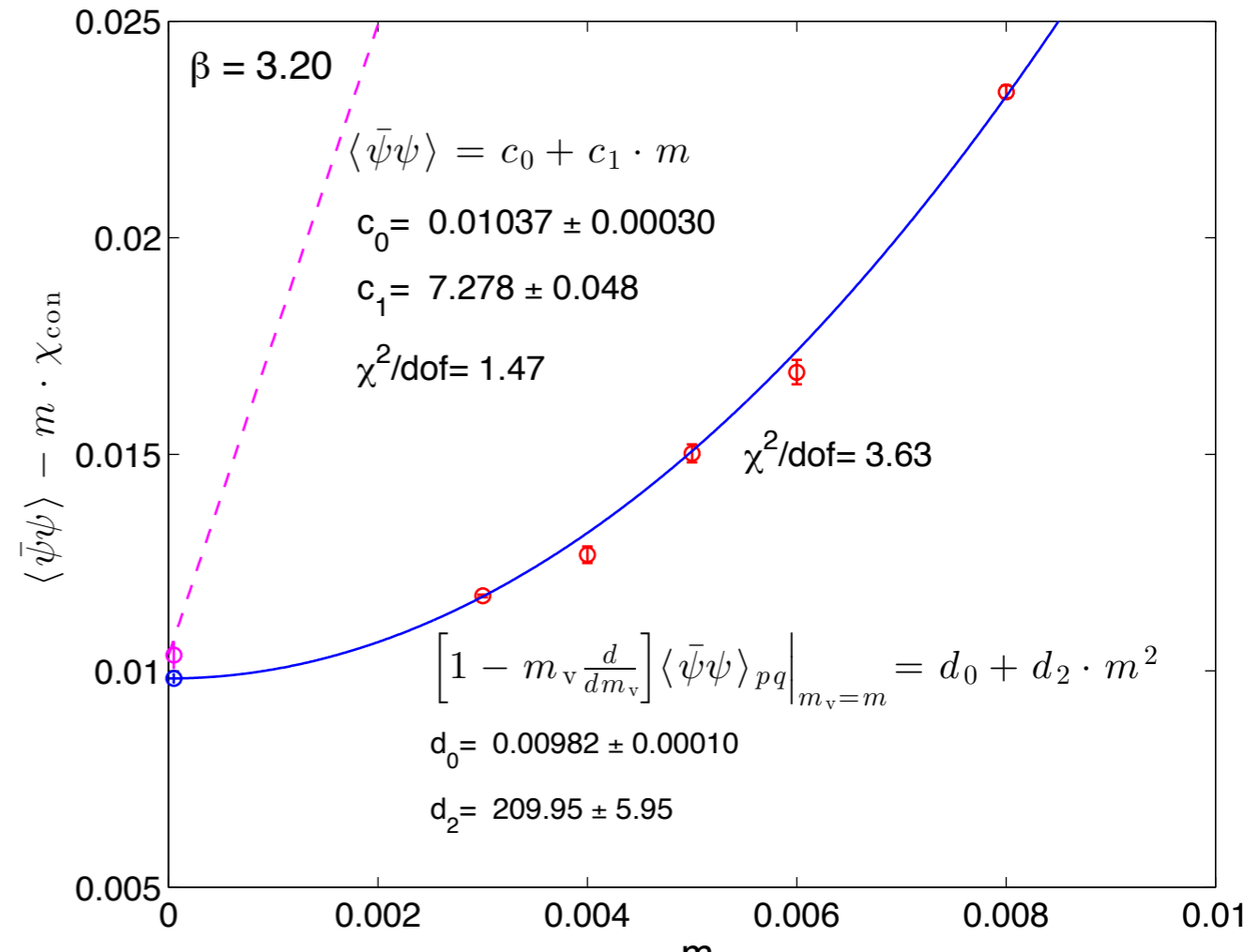
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 (partially cancelled UV divergences in subtracted form)

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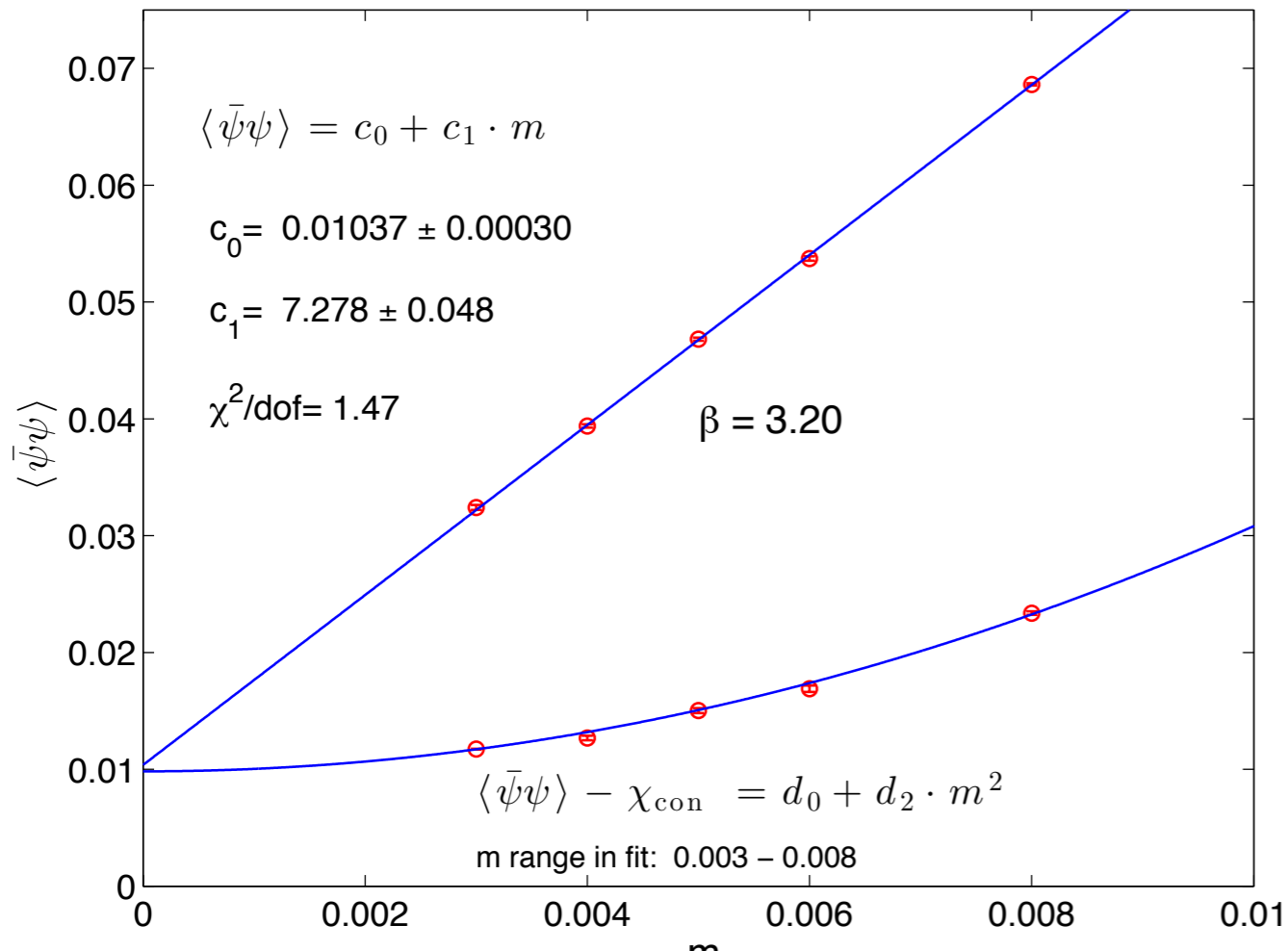
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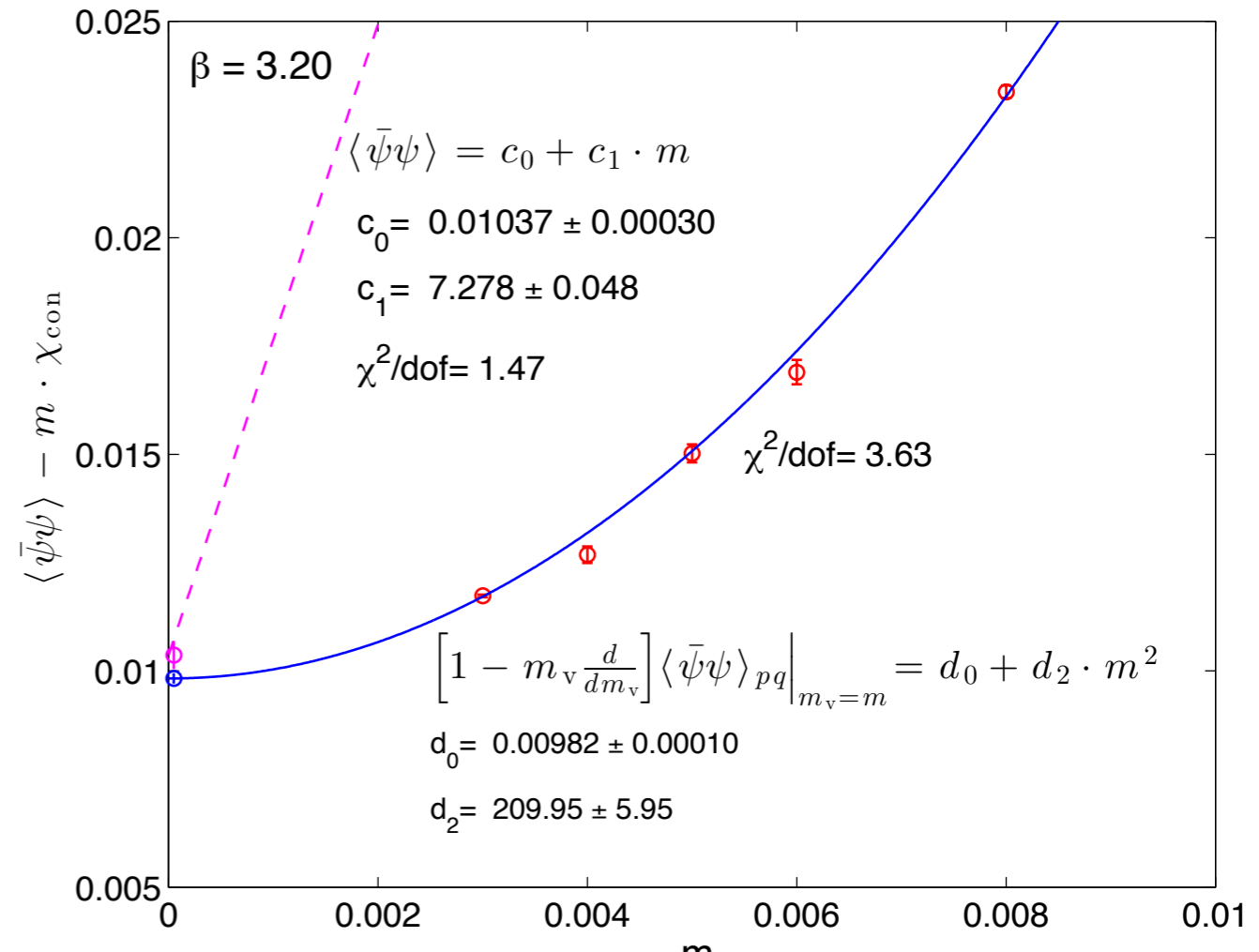
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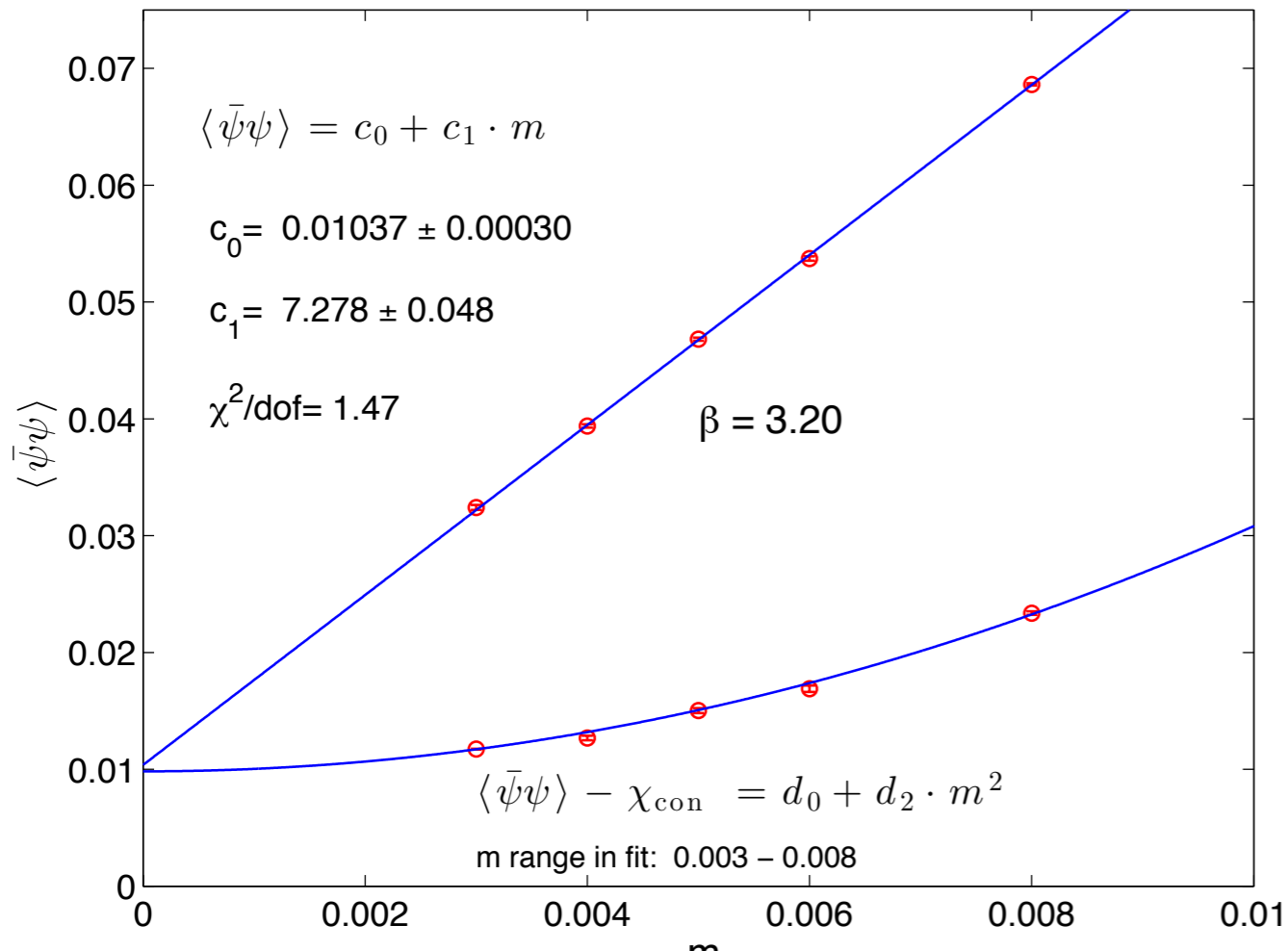
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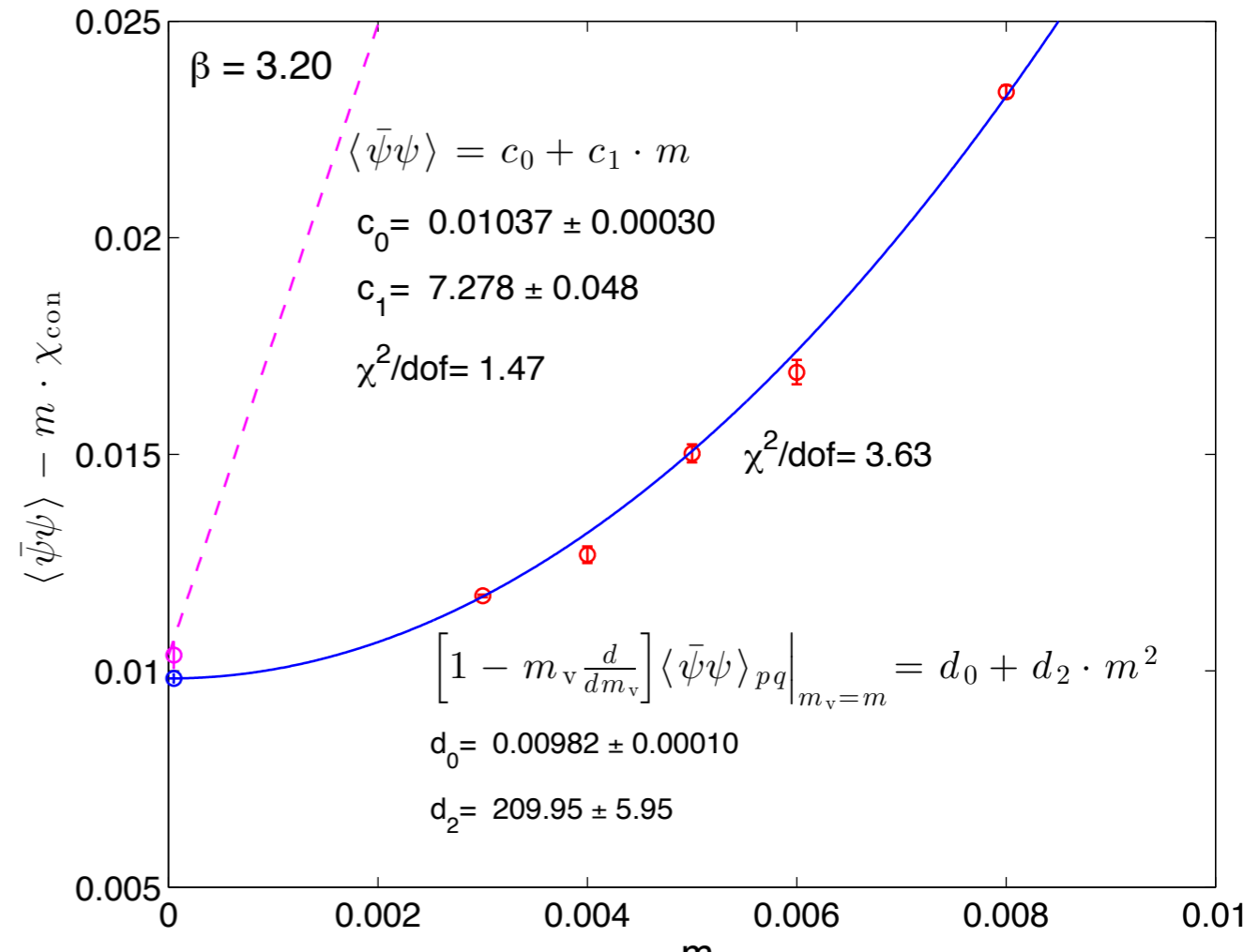
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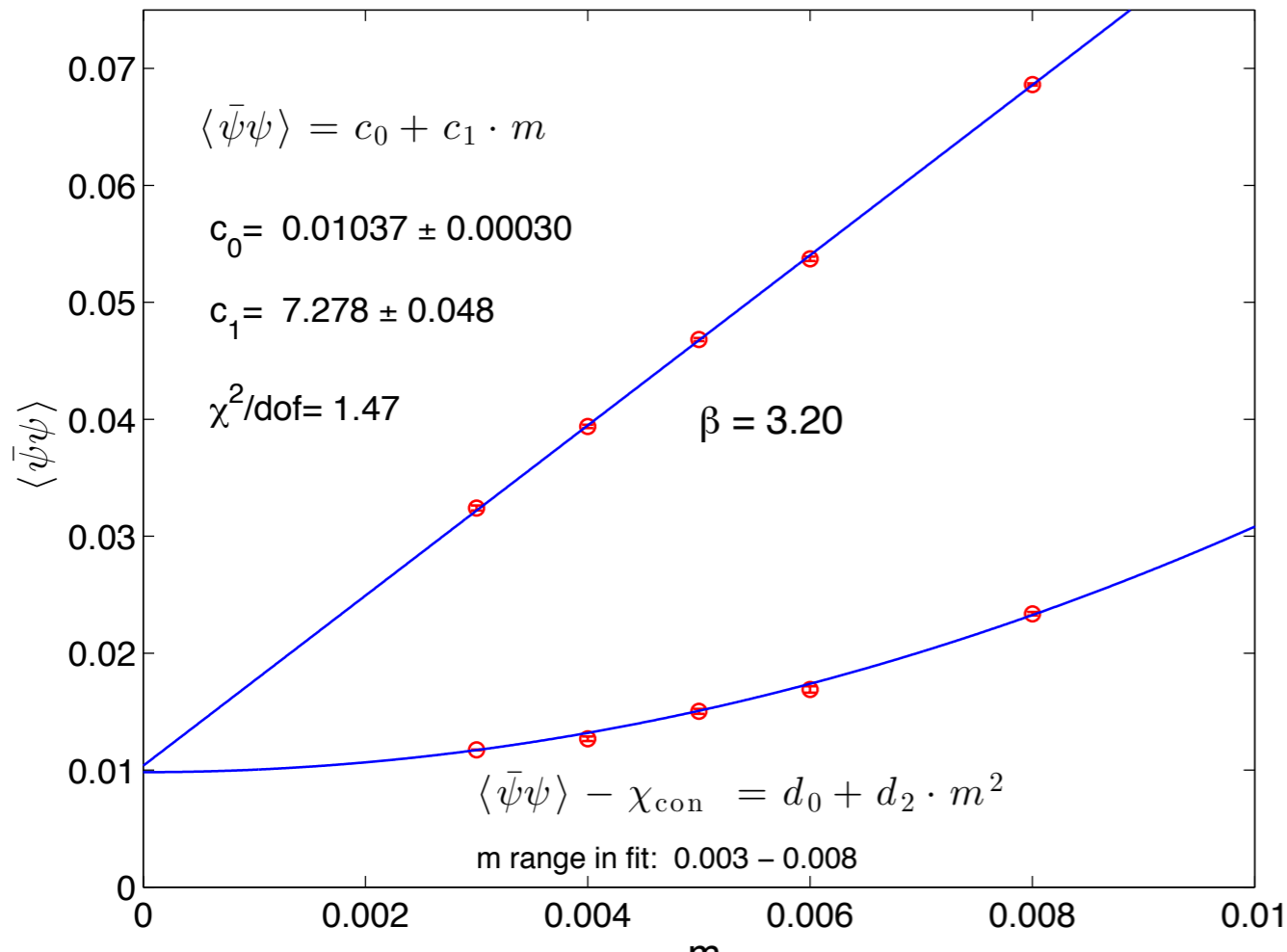
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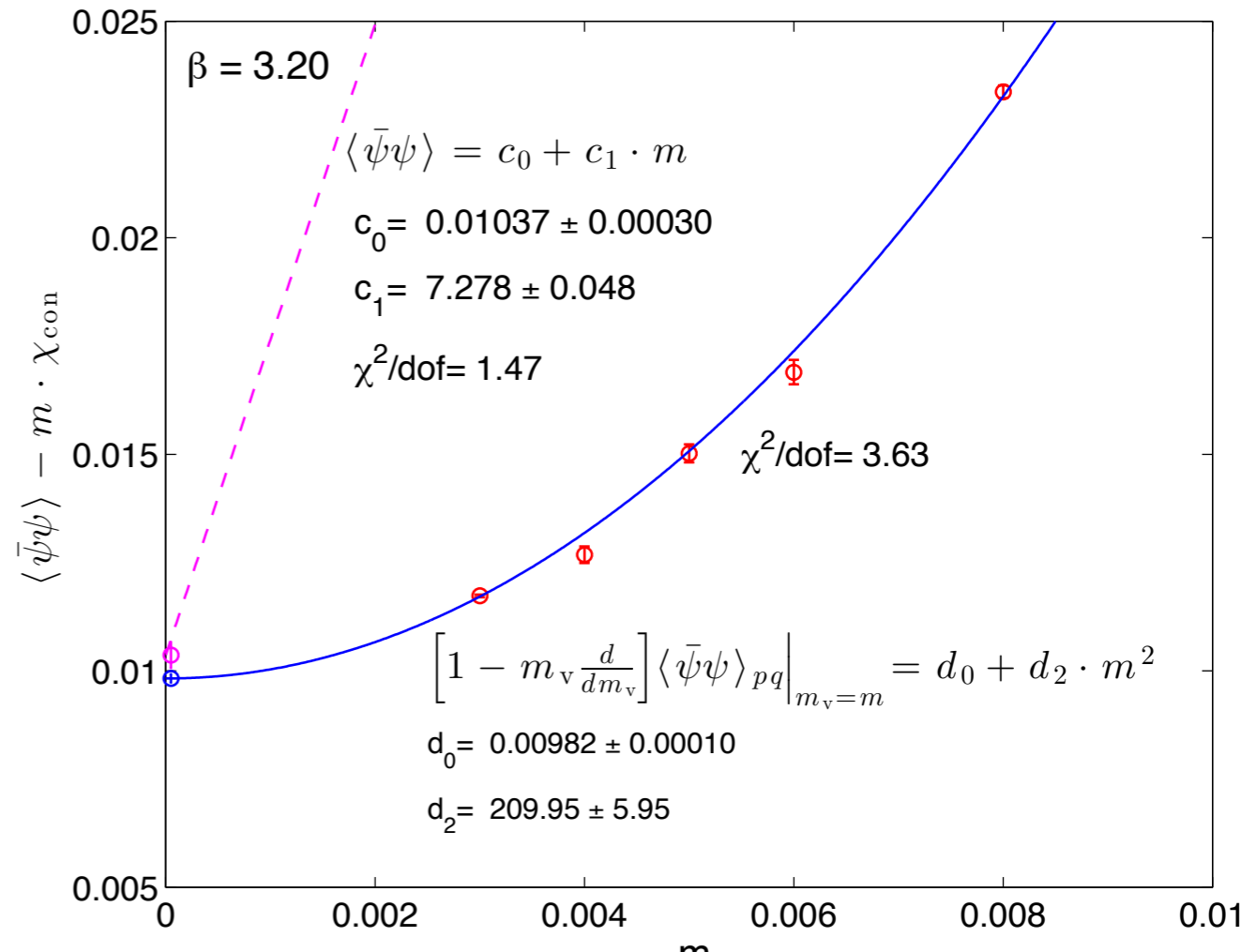
relying on  $L \cdot M\pi > 5$  (less than one percent L correction)

# Nf=2 SU(3) sextet chiral condensate

chiral condensate and its subtracted form



subtracted chiral condensate



two independent determinations of the chiral condensate

(partially cancelled UV divergences in subtracted form)

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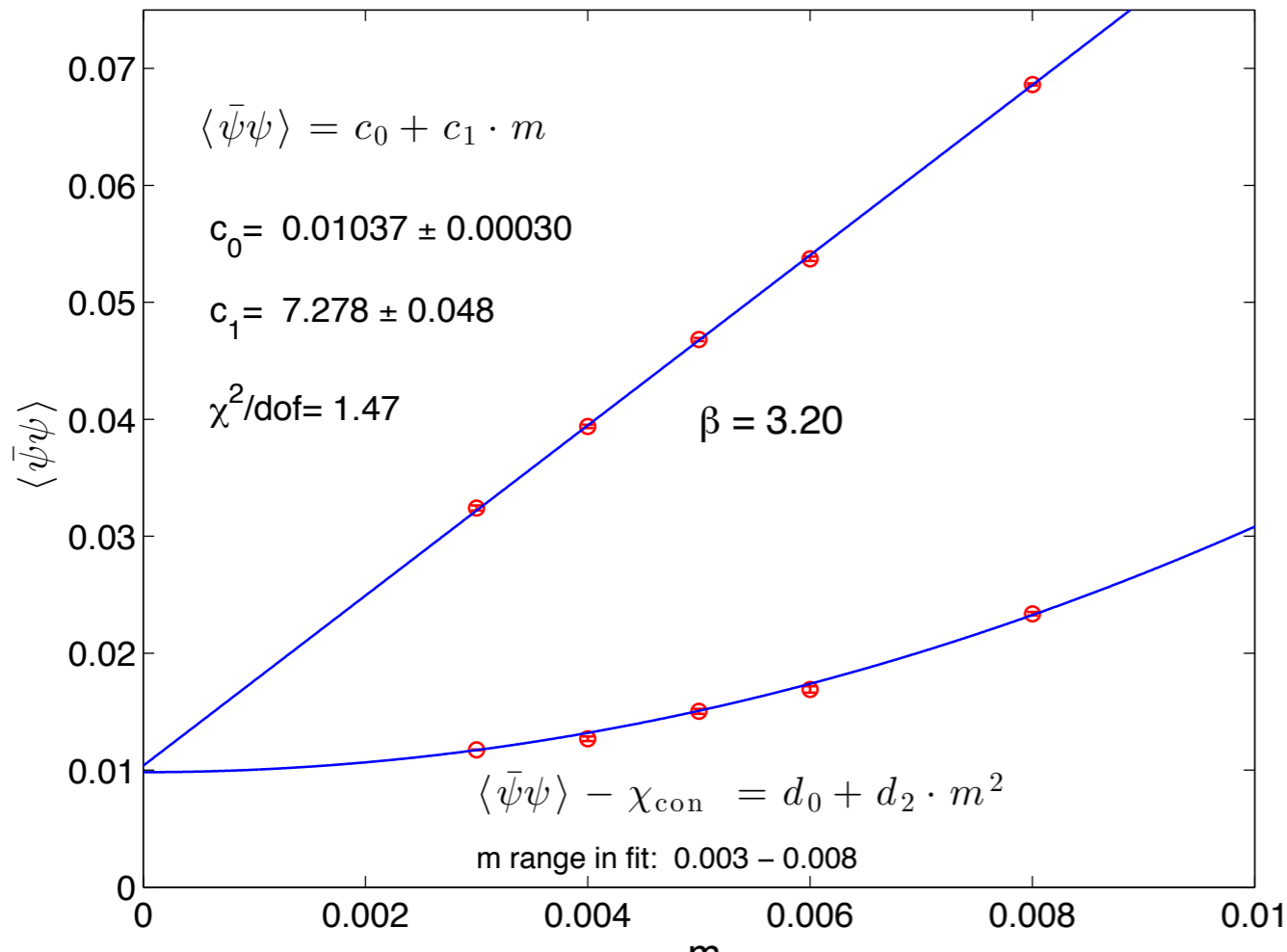
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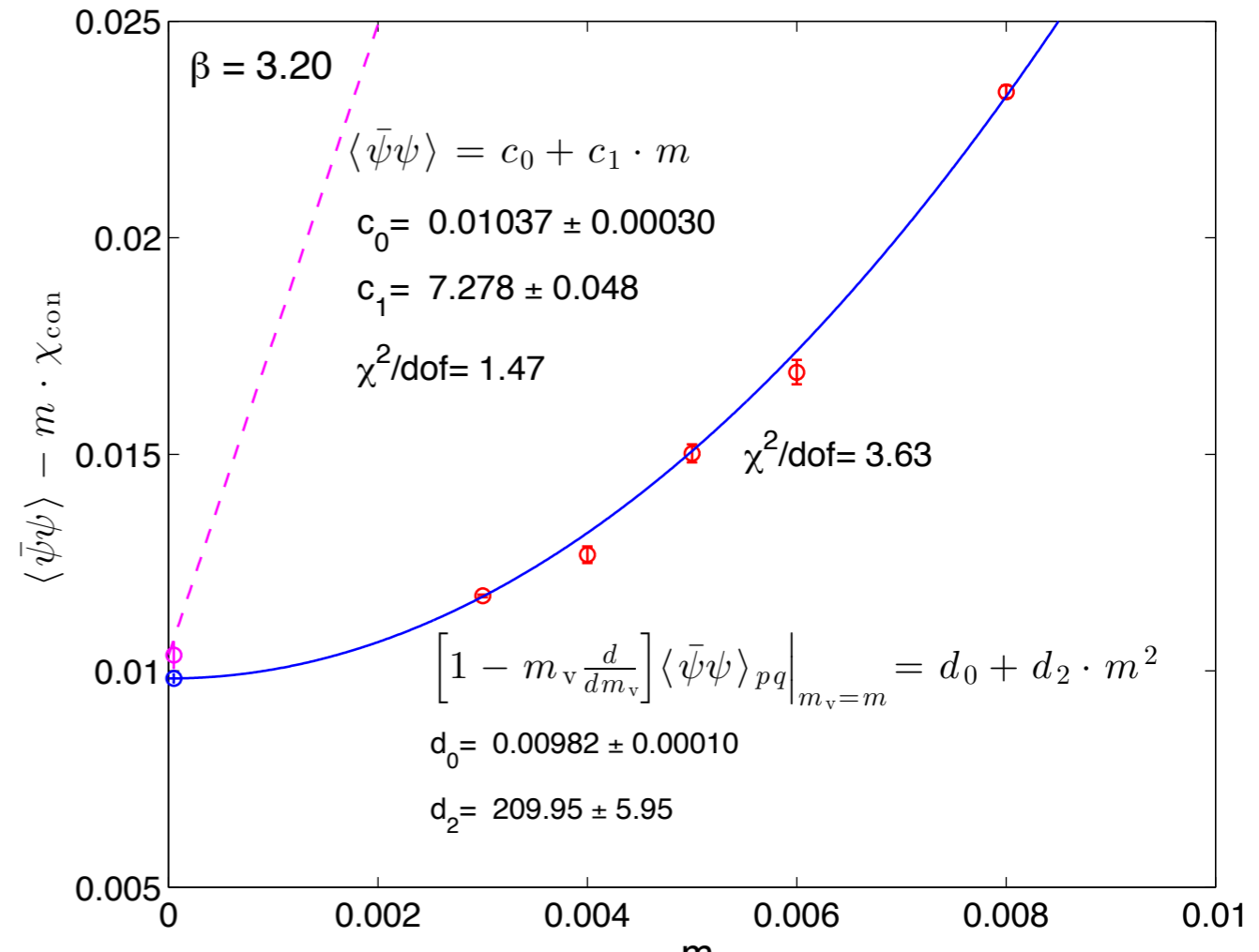
spectral density analysis more powerful

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spectral density analysis more powerful

(Del Debbio et al., Giusti and Luscher, Boulder group, Patella ...)

complete control on UV divergences:

node number density of chiral condensate

$$\rho(\lambda, m) = \frac{1}{V} \sum_{k=1}^{\infty} \langle \delta(\lambda - \lambda_k) \rangle$$

$$\lim_{\lambda \rightarrow 0} \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \rho(\lambda, m) = \frac{\Sigma}{\pi}$$

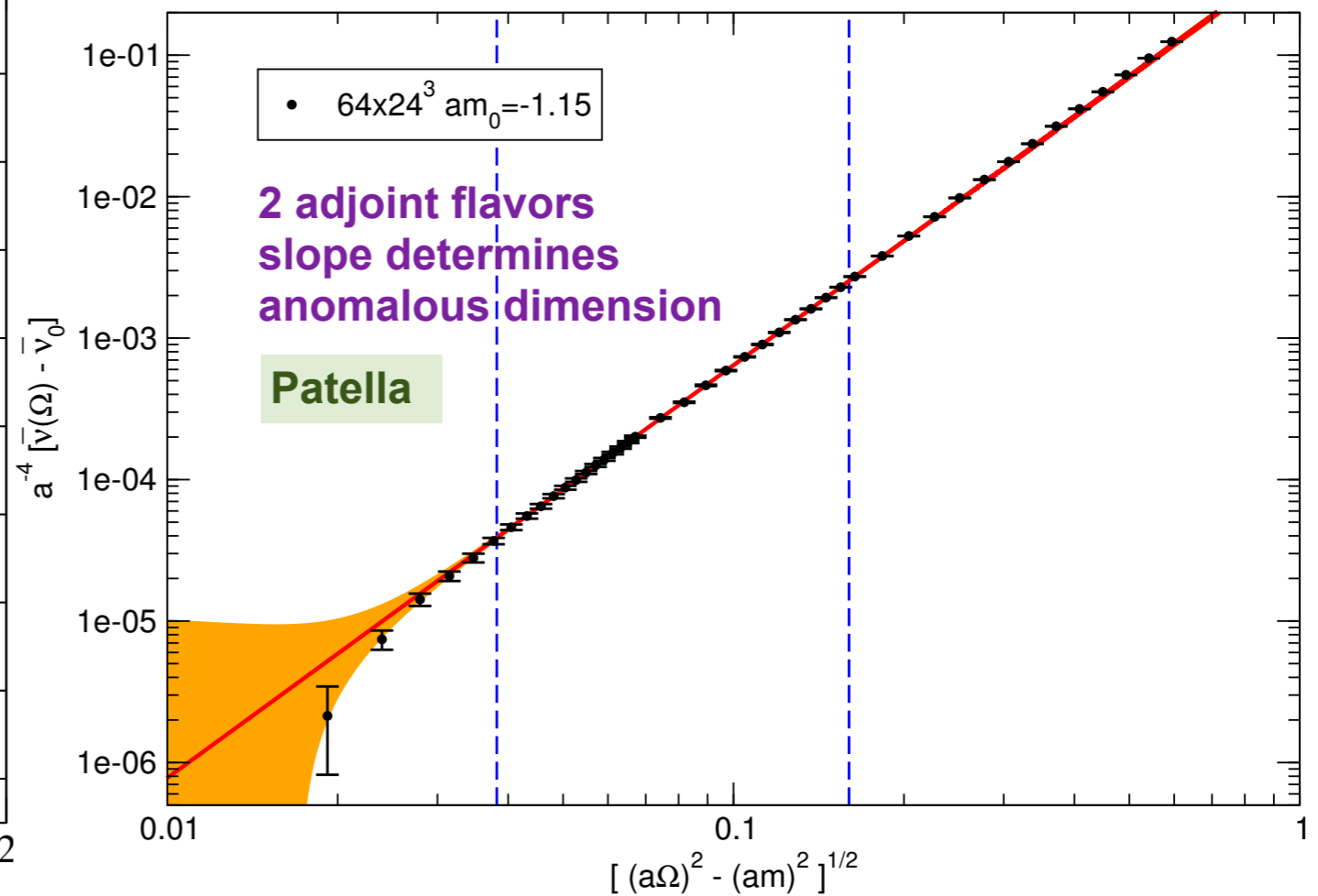
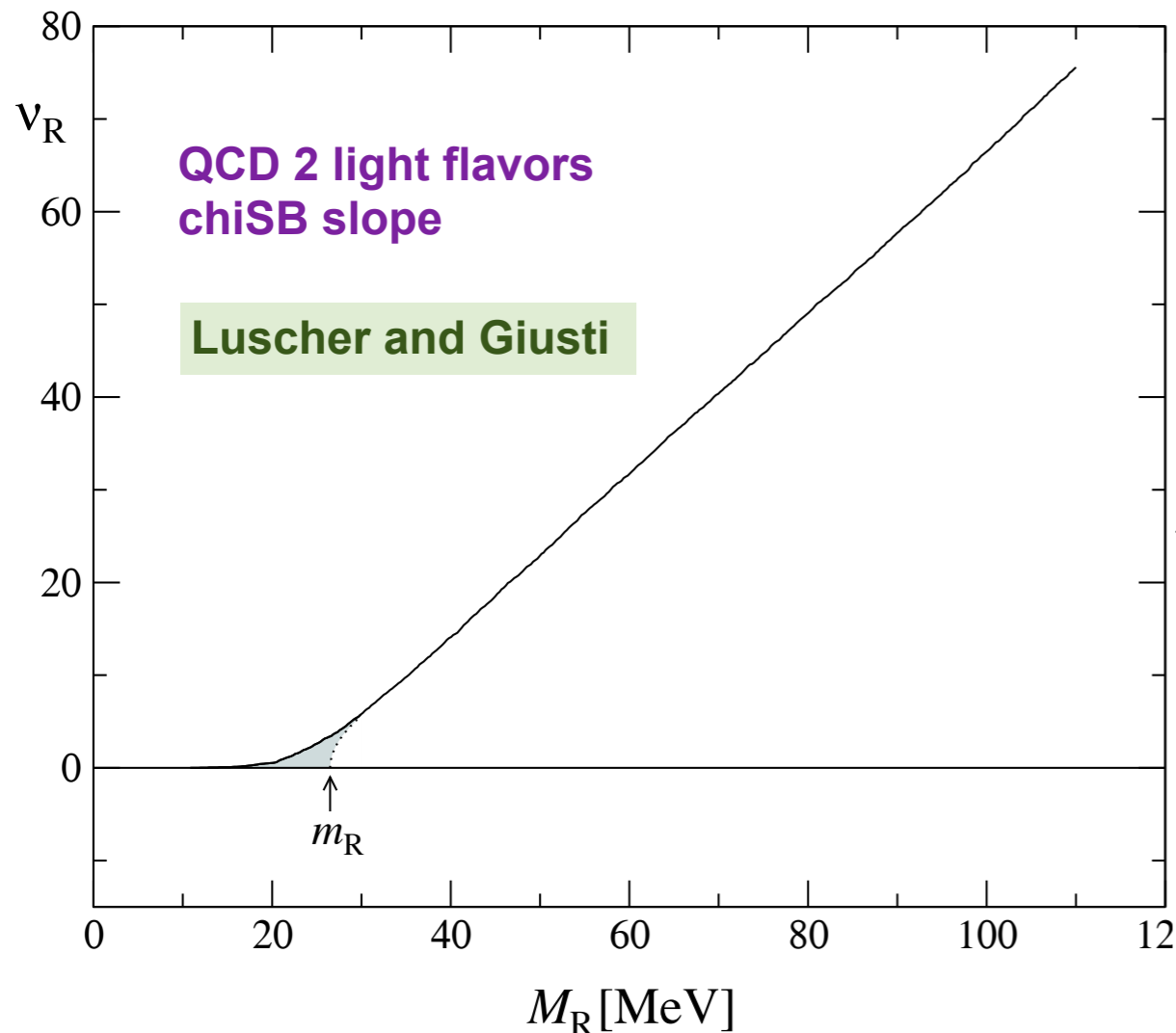
spectral density

$$\nu(M, m) = V \int_{-\Lambda}^{\Lambda} d\lambda \rho(\lambda, m), \quad \Lambda = \sqrt{M^2 - m^2}$$

node number density

$$\nu_R(M_R, m_R) = \nu(M, m_q) \text{ renormalized and RG invariant}$$

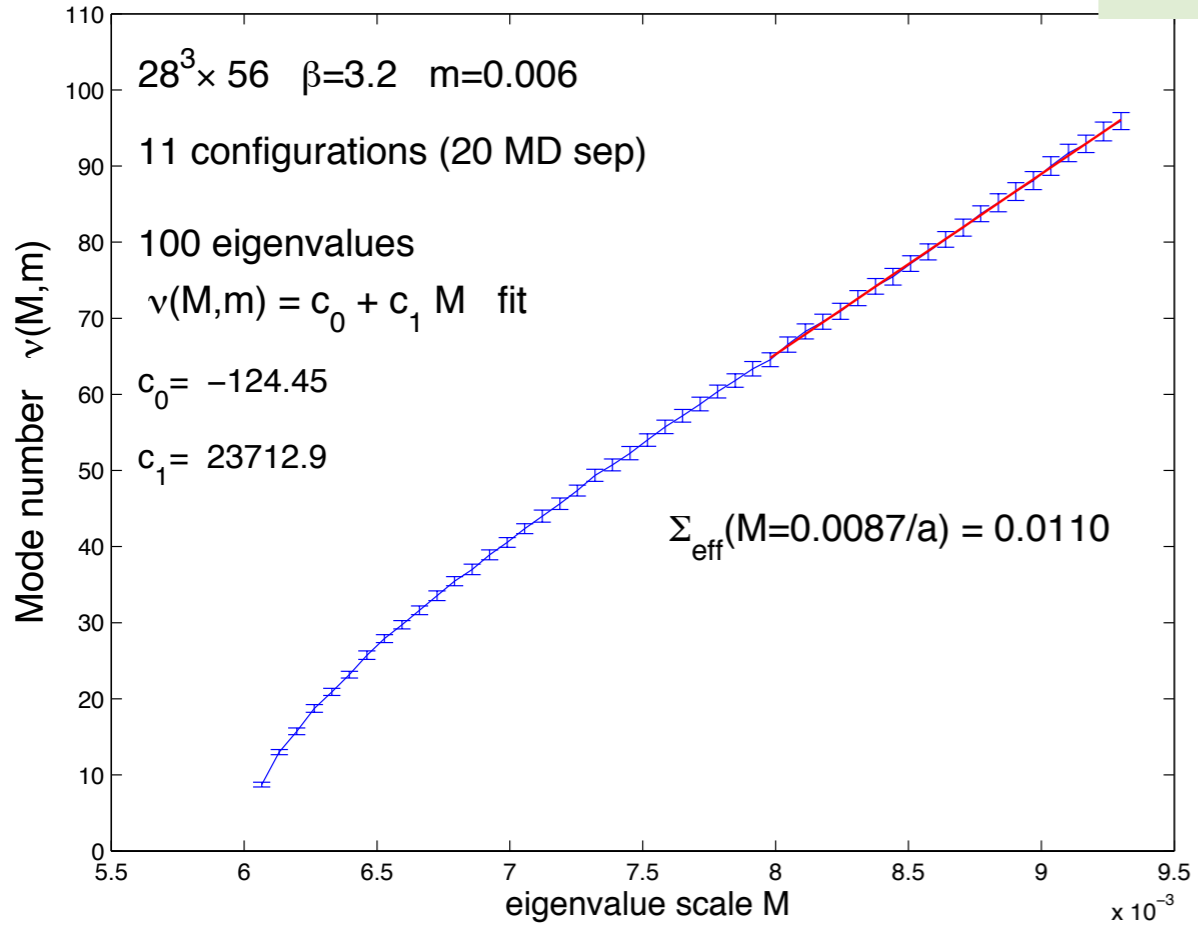
Del Debbio et al.  
Luscher and Giusti  
Patella  
Boulder group



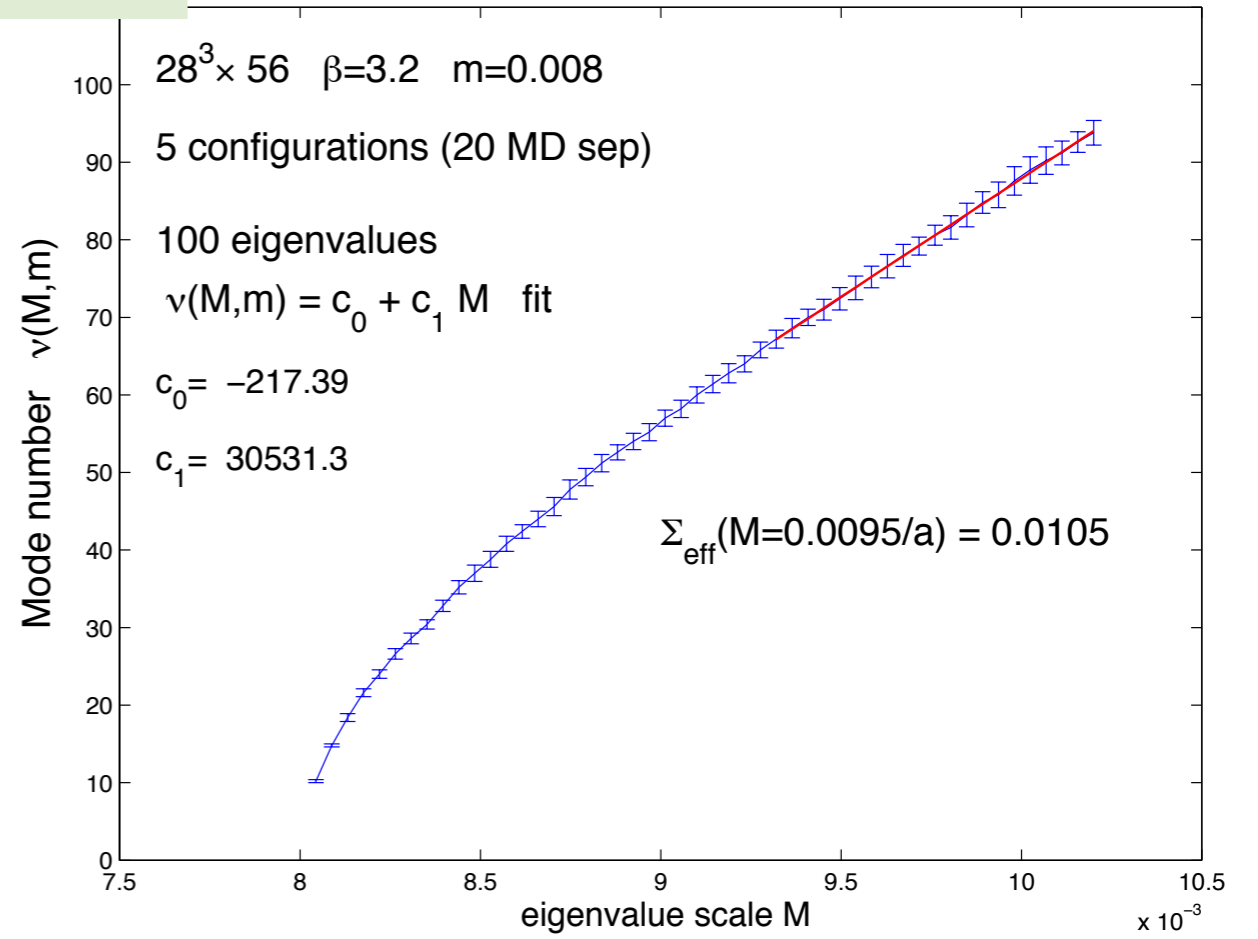


# sextet model

Scale dependence of Mode number distribution  $\nu(M,n)$

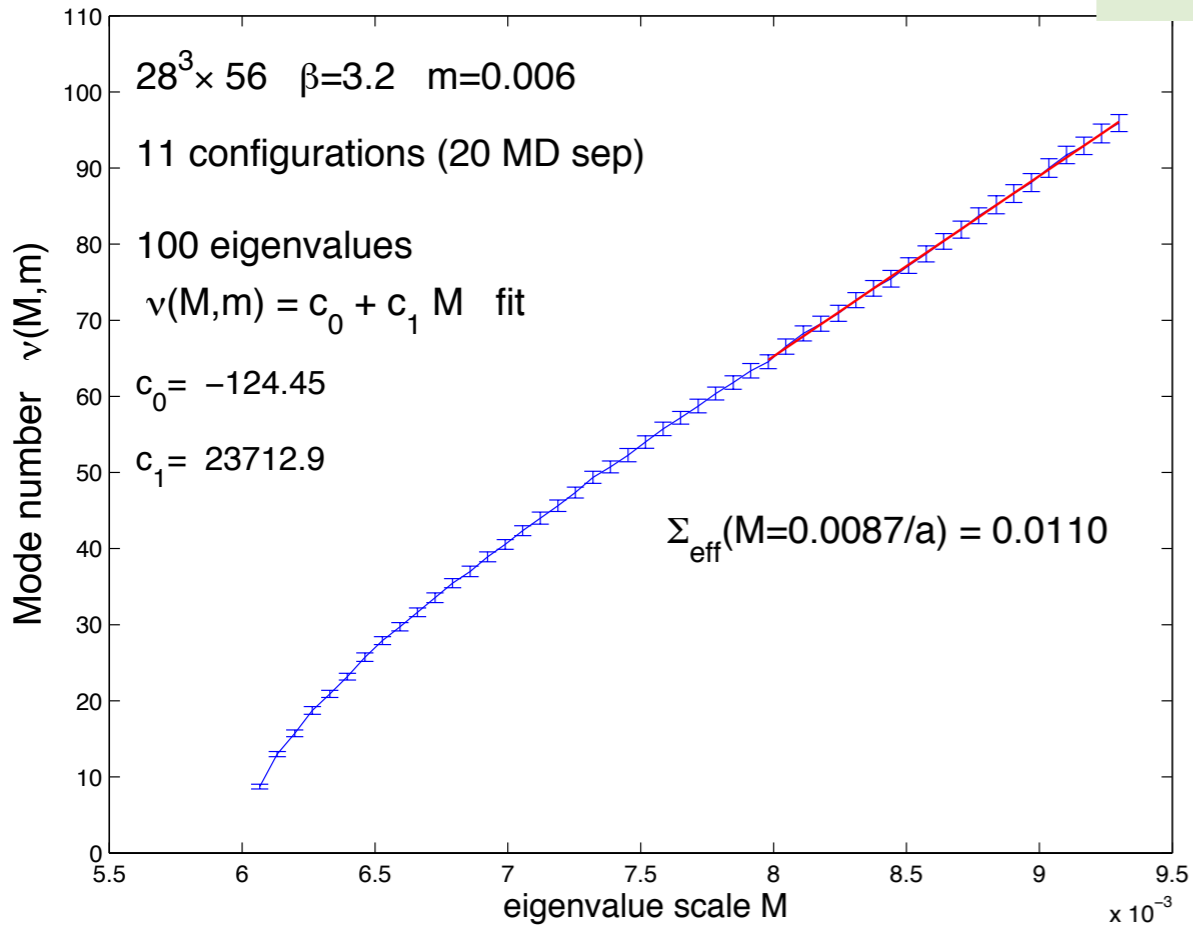


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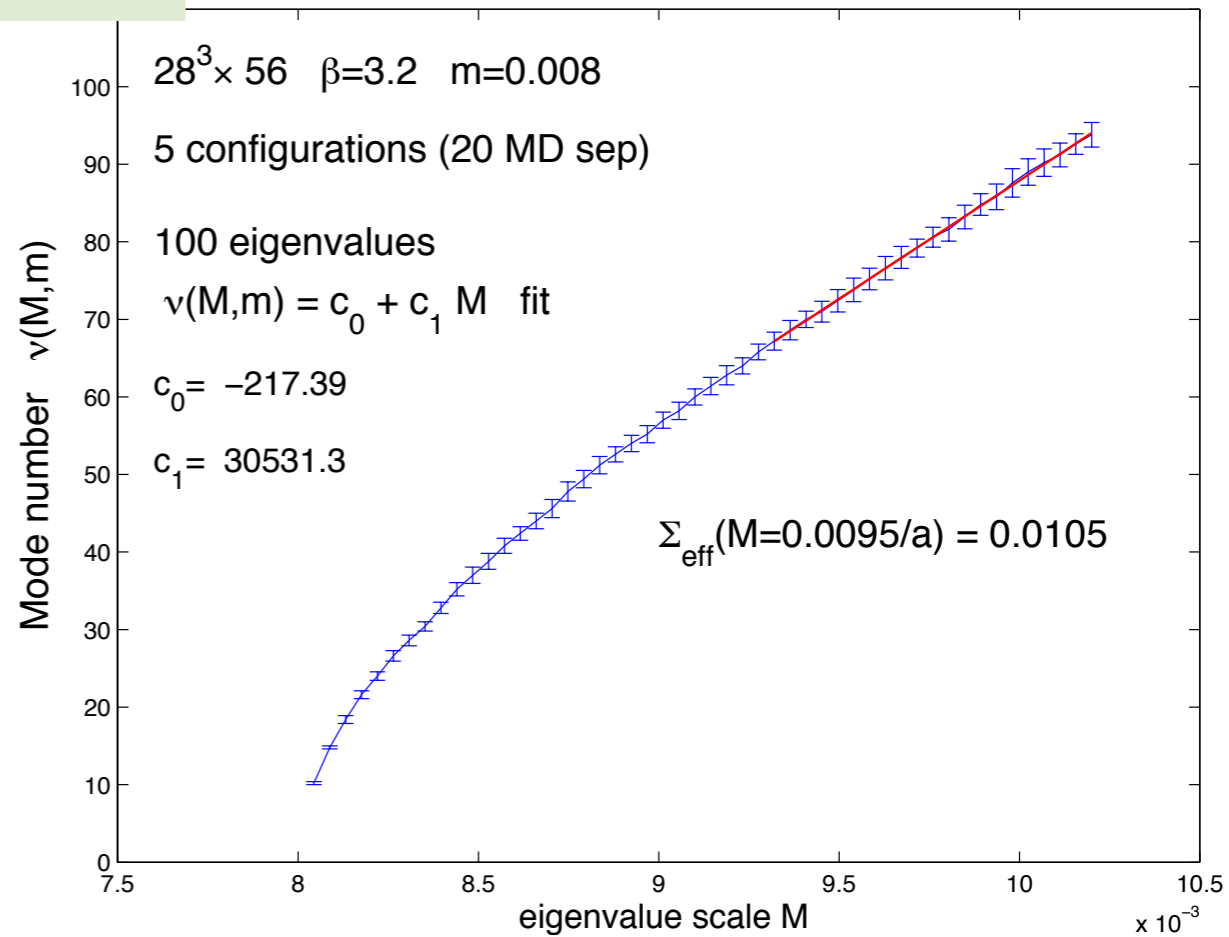


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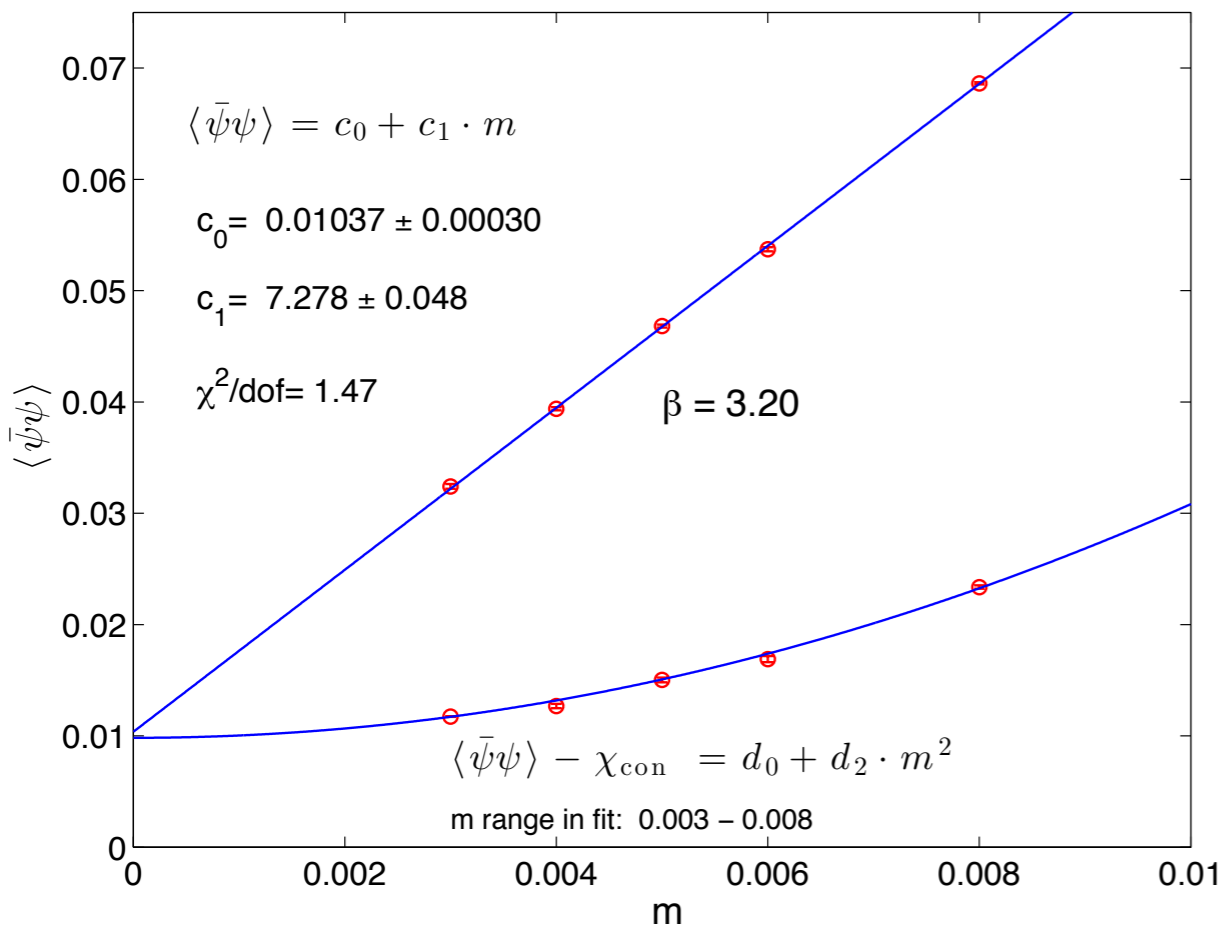
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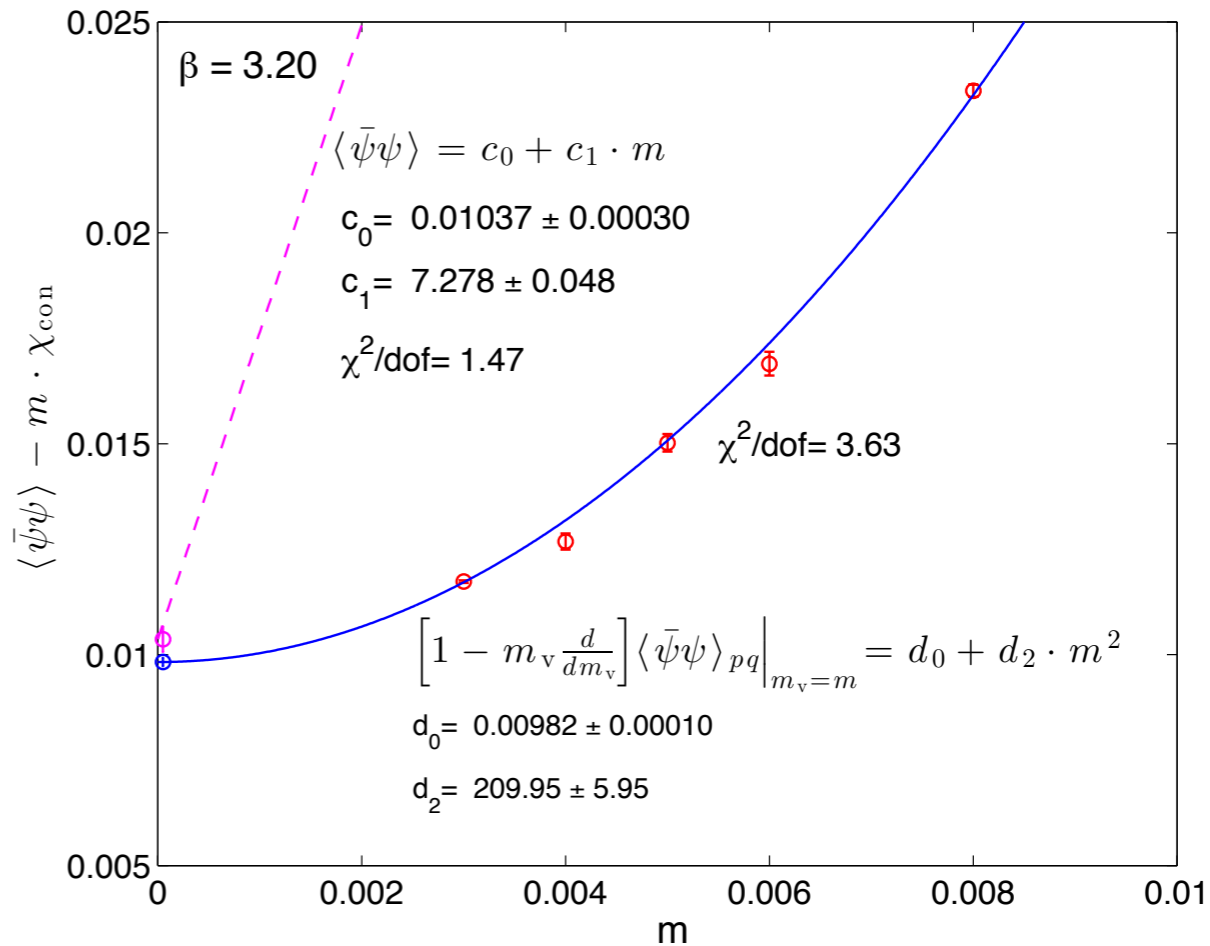
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chiral condensate and its subtracted form

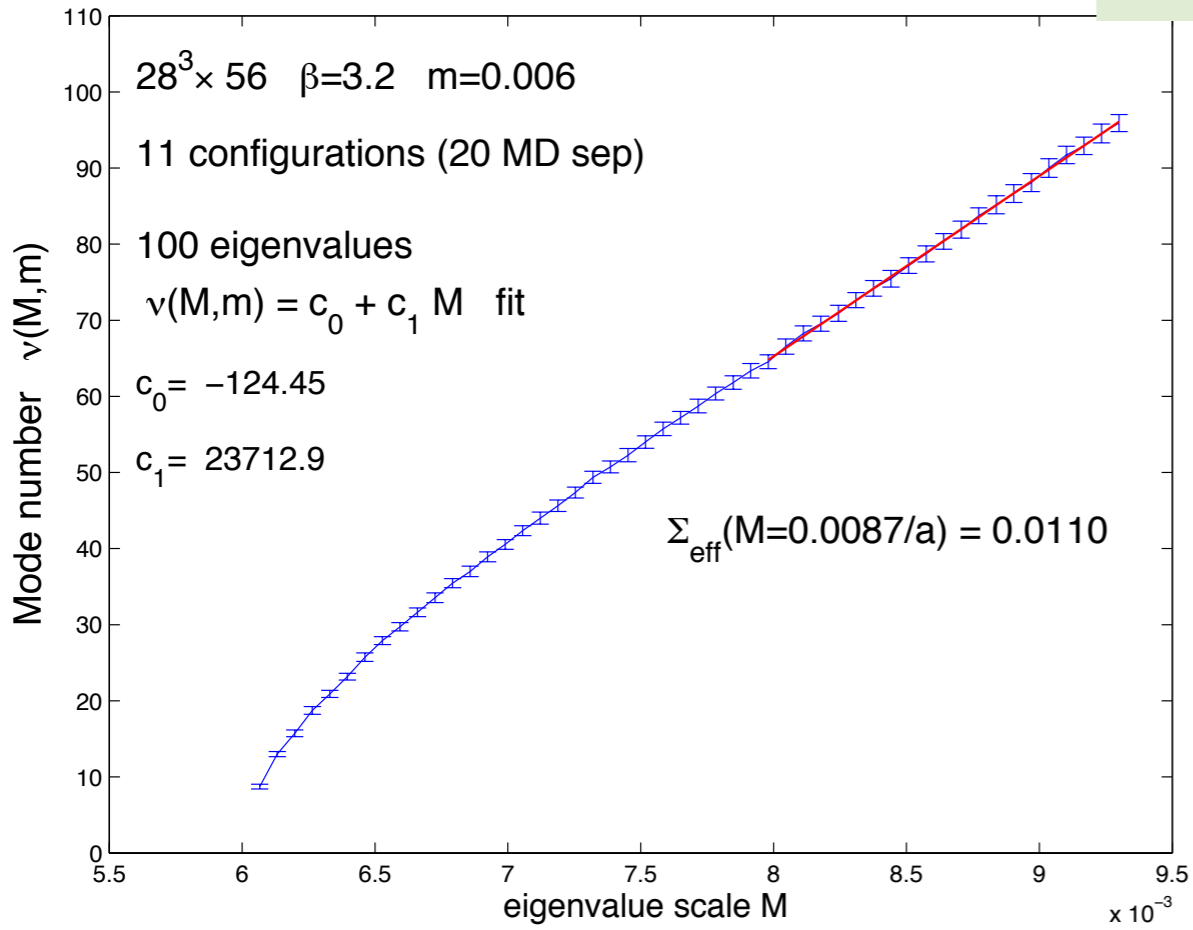


subtracted chiral condensate

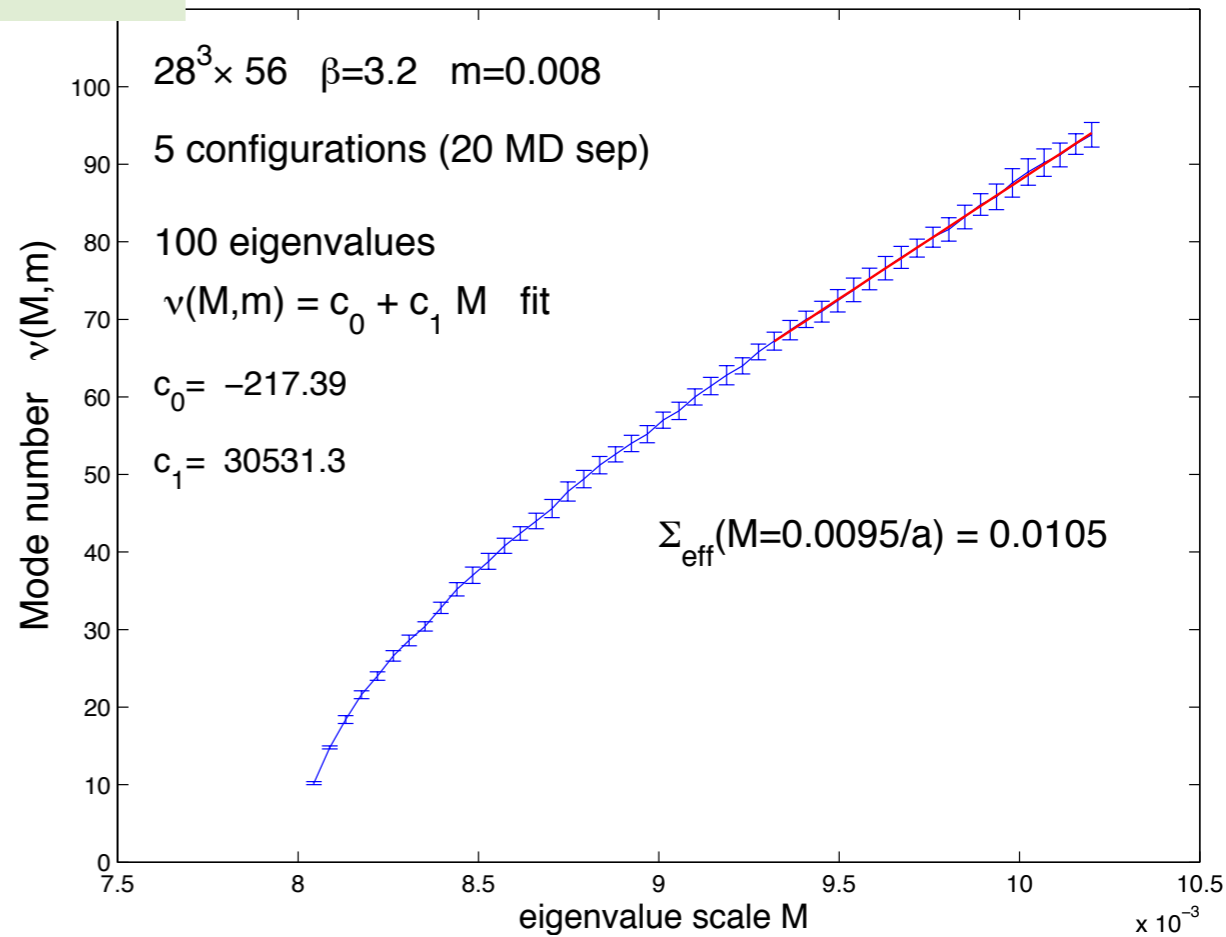


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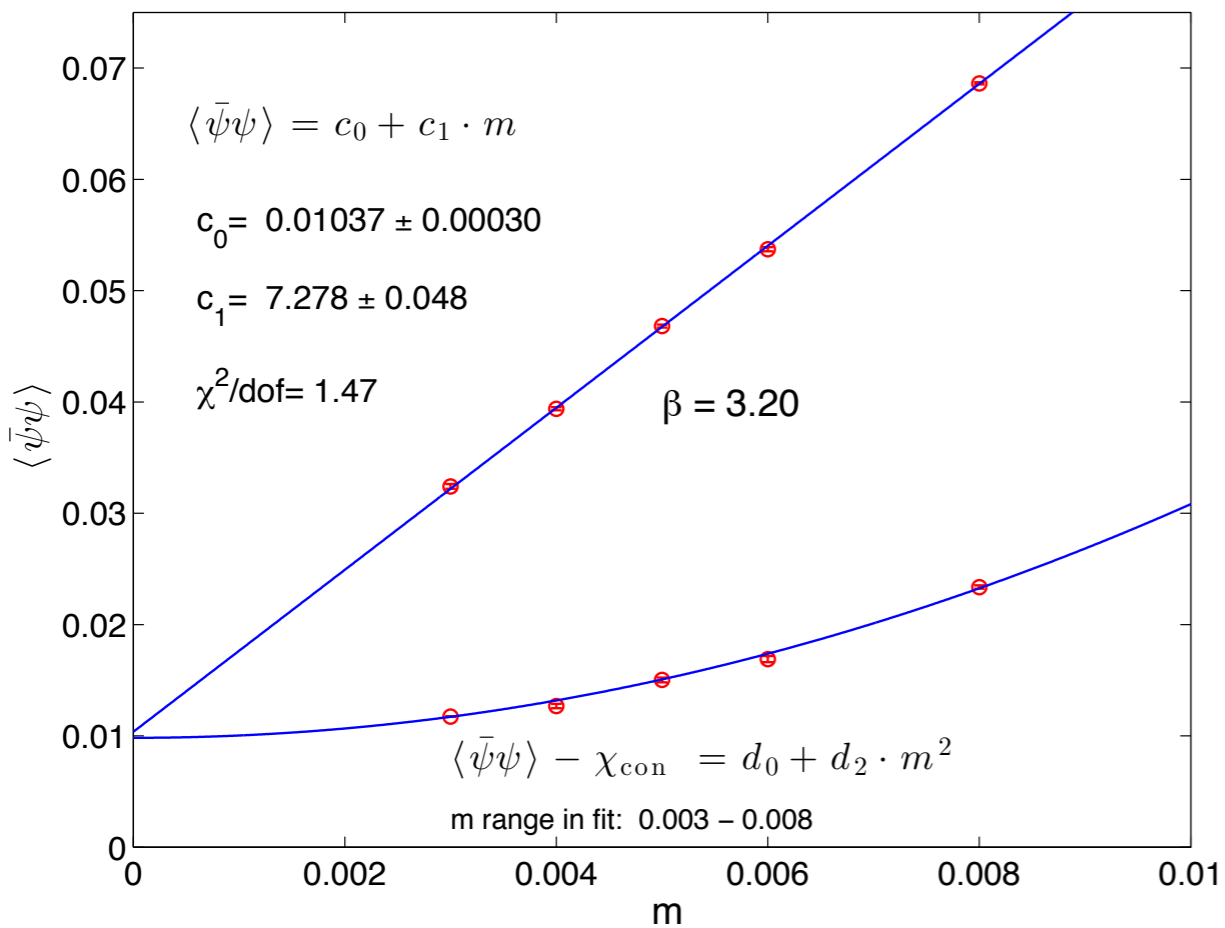
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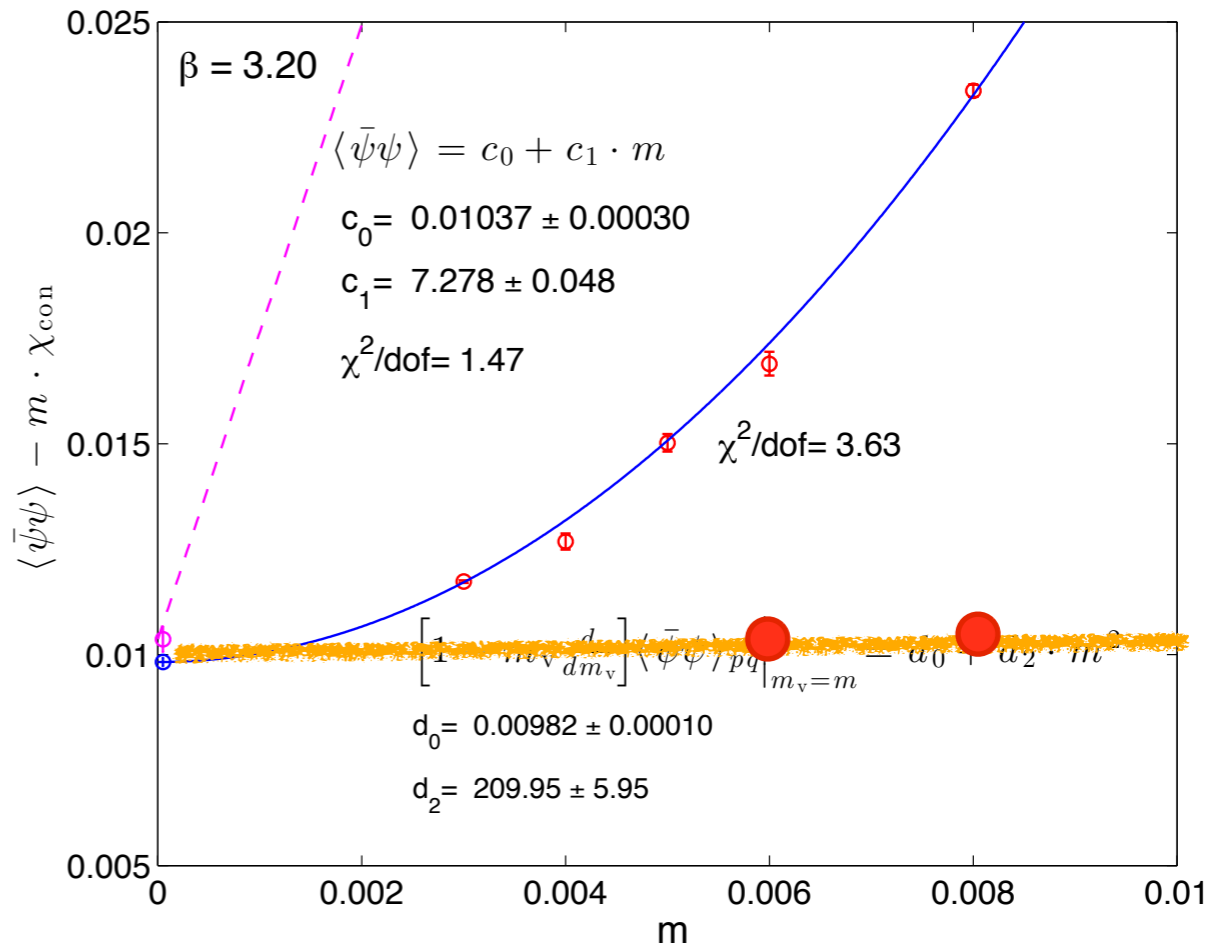
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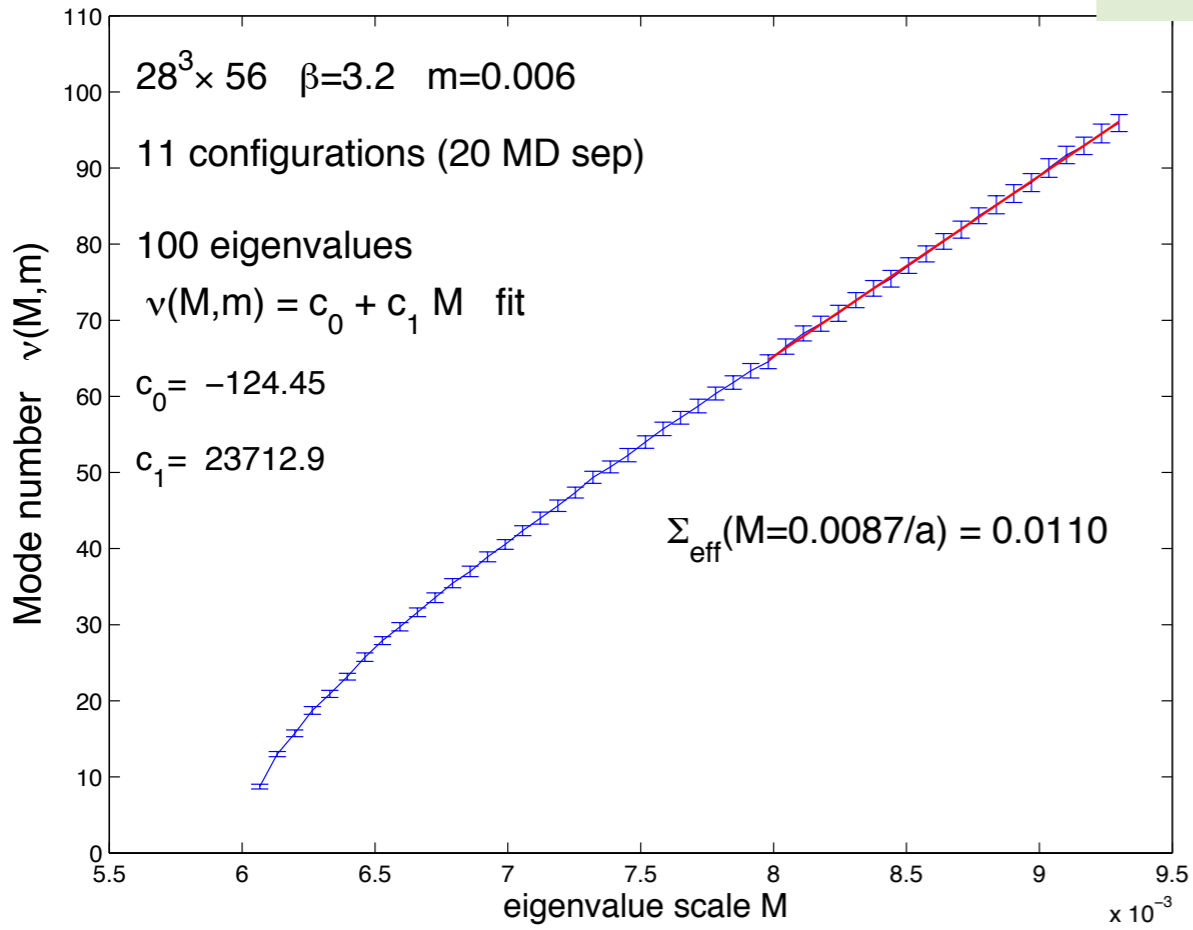


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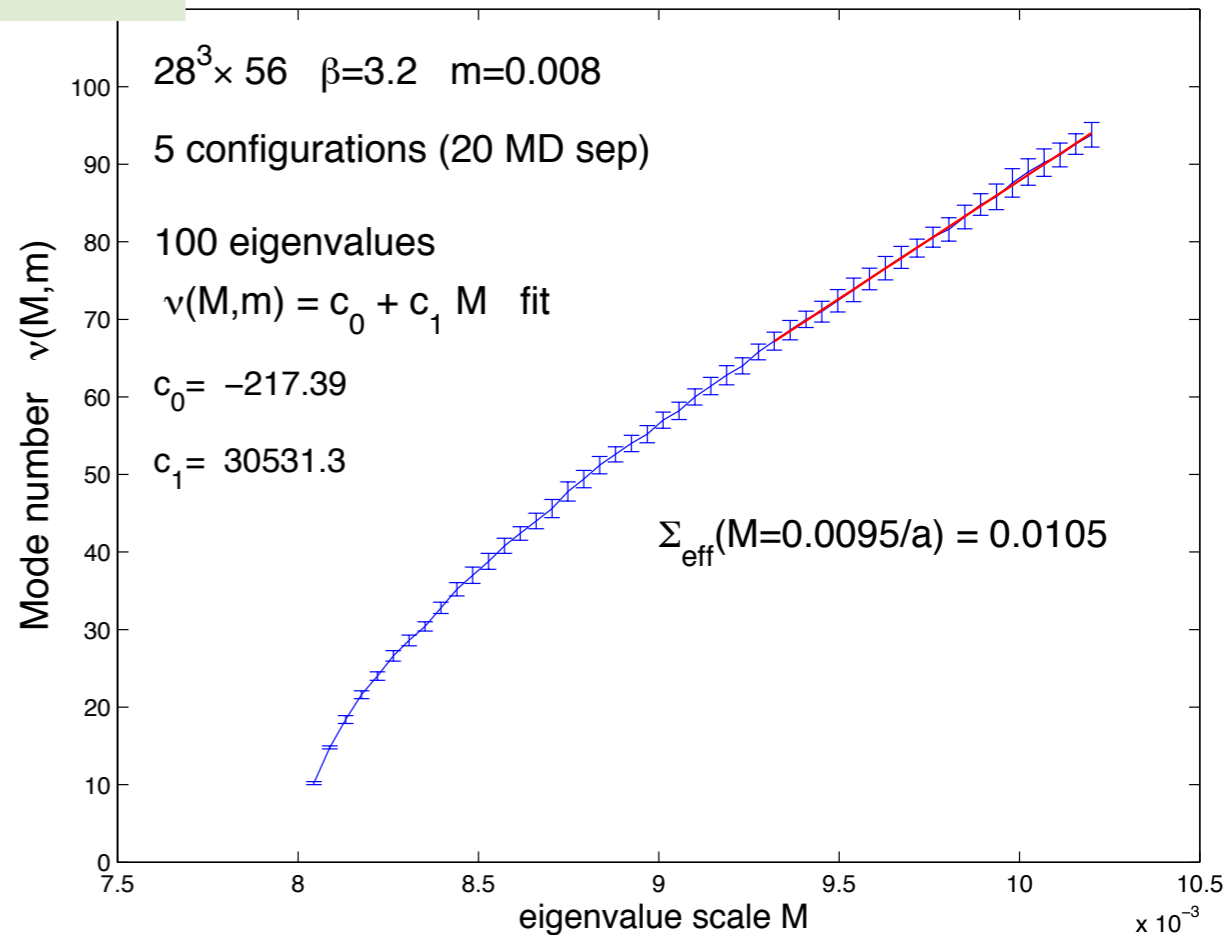


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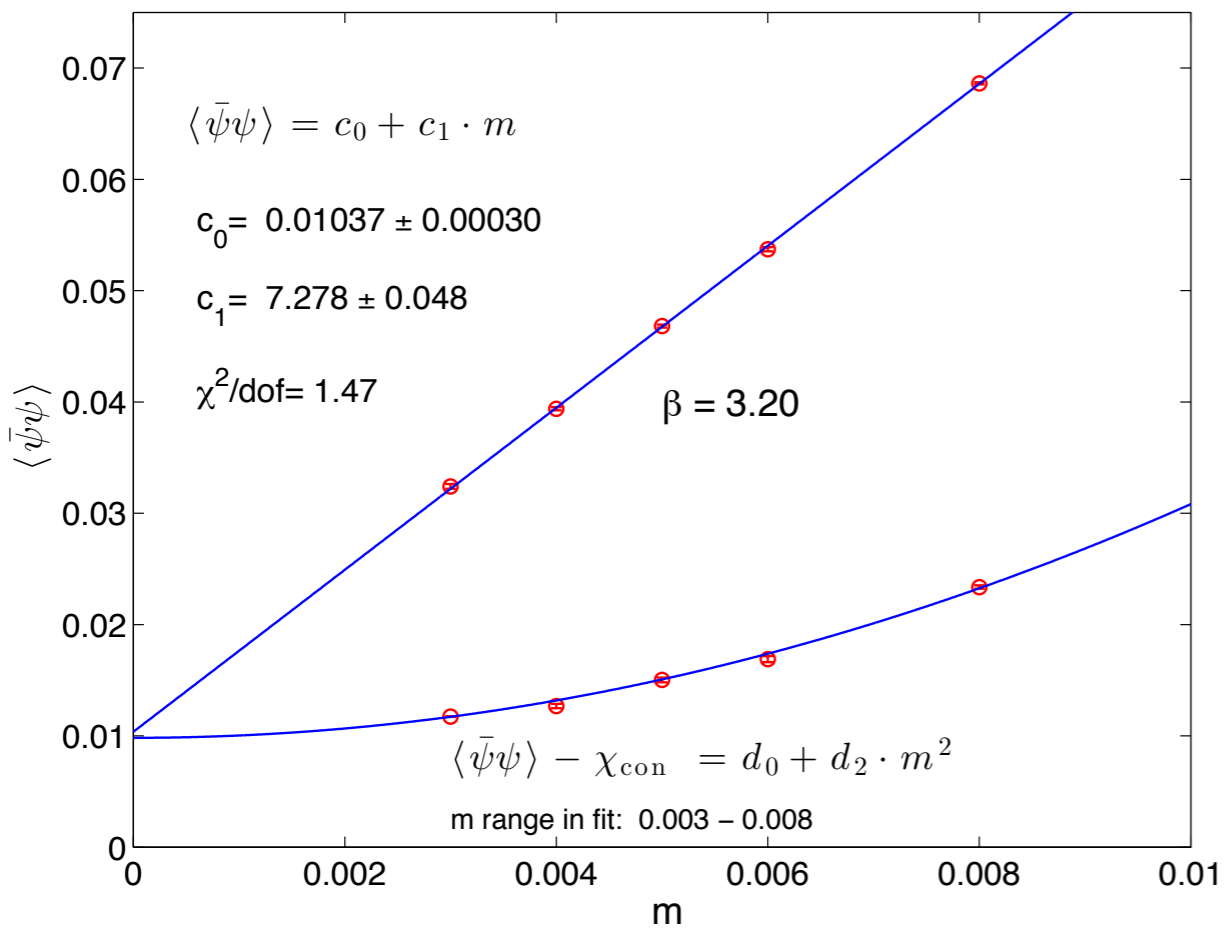
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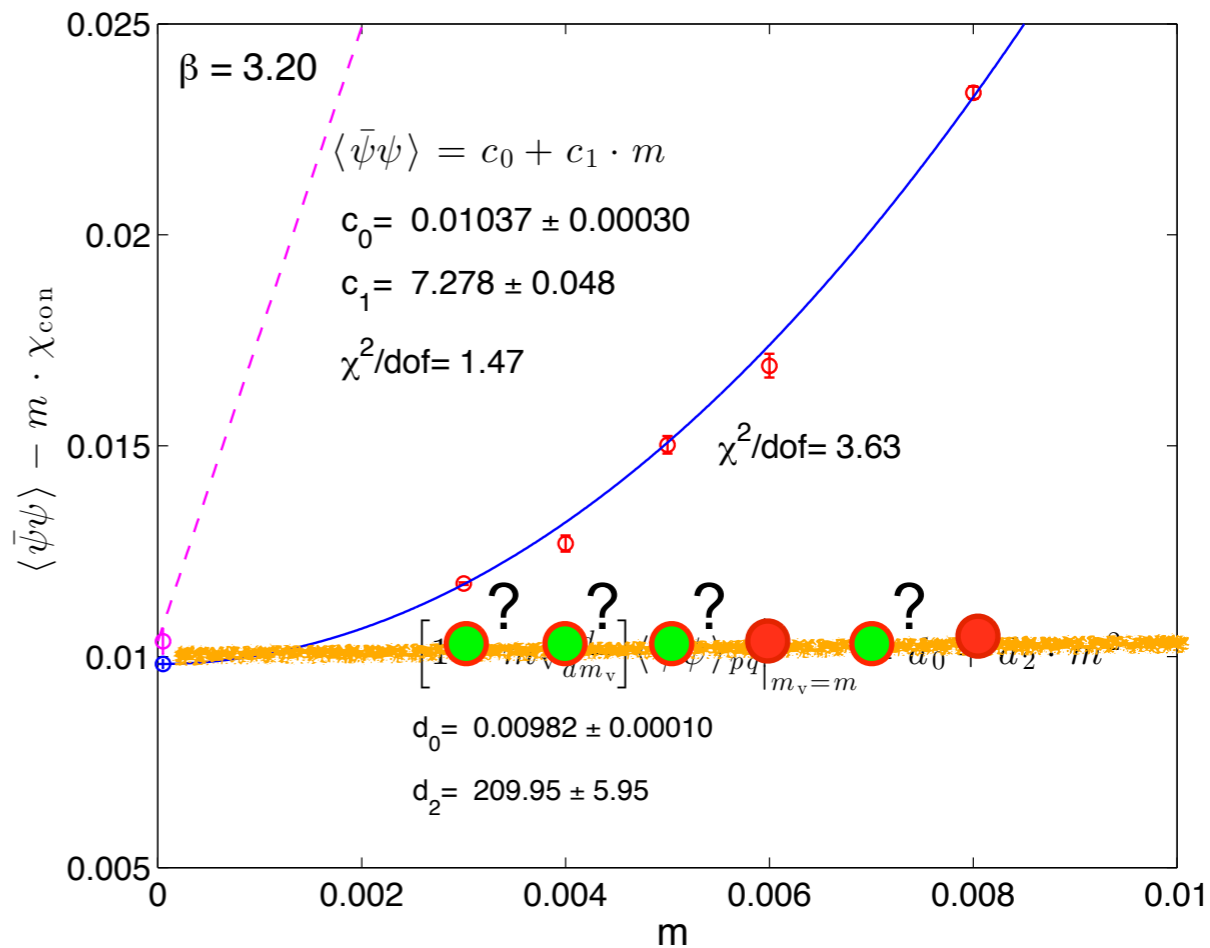
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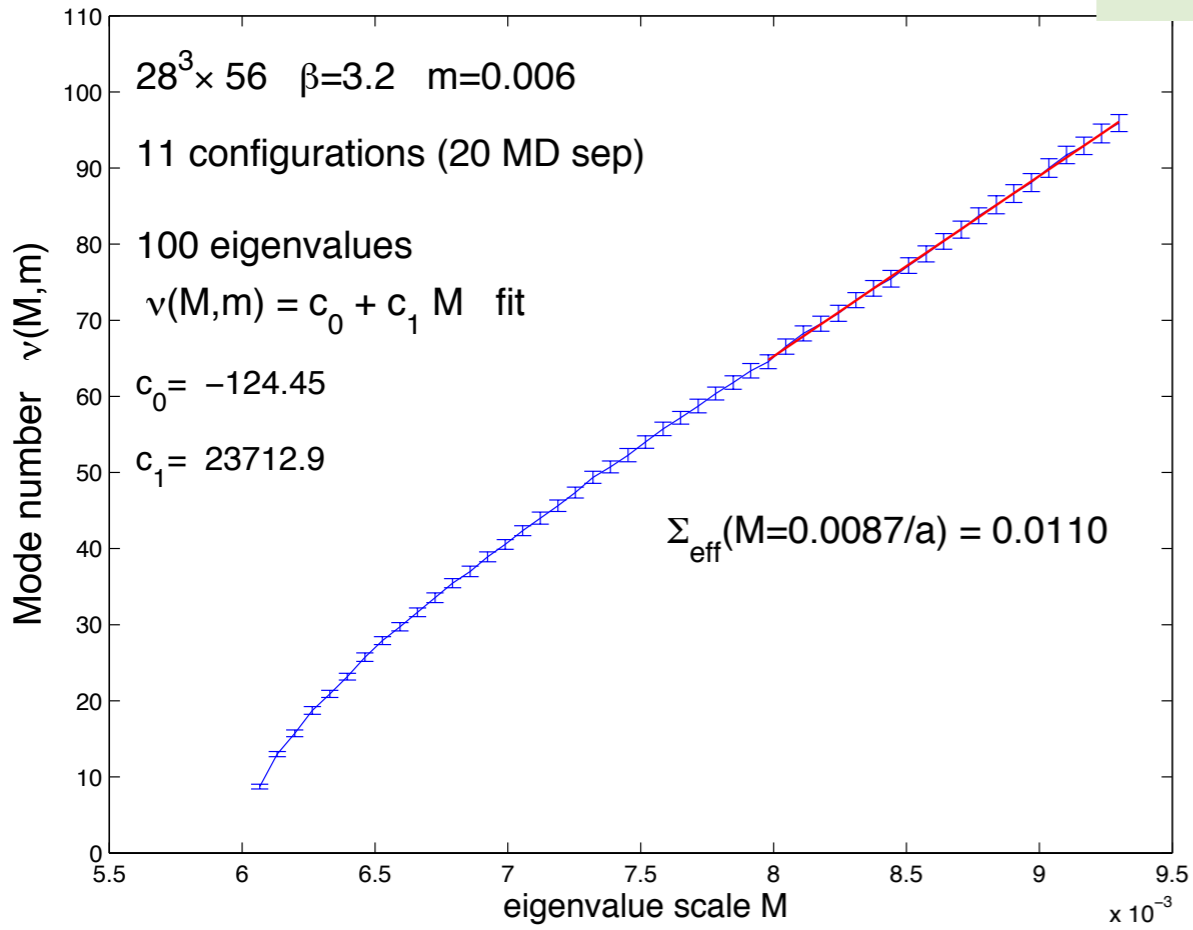


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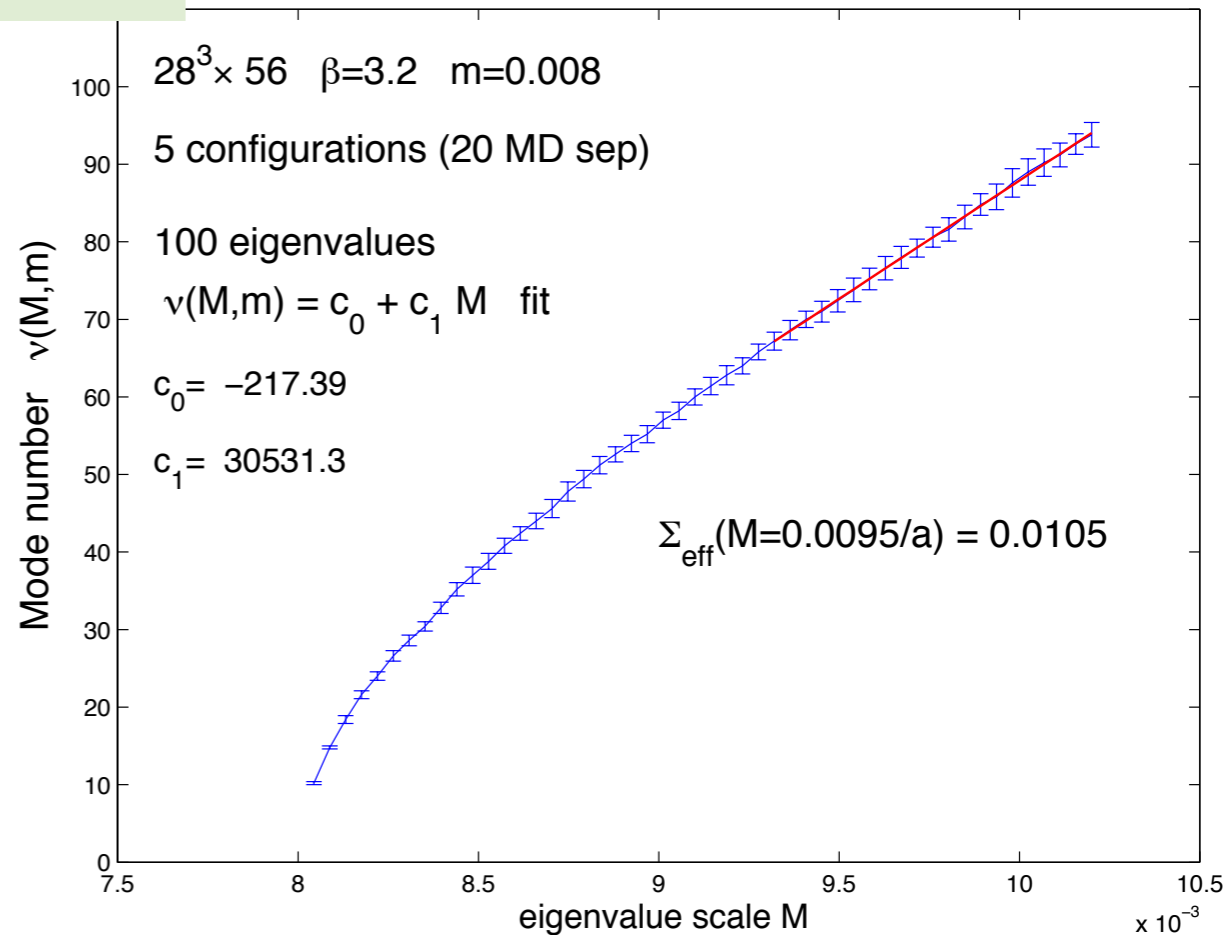


# sextet model

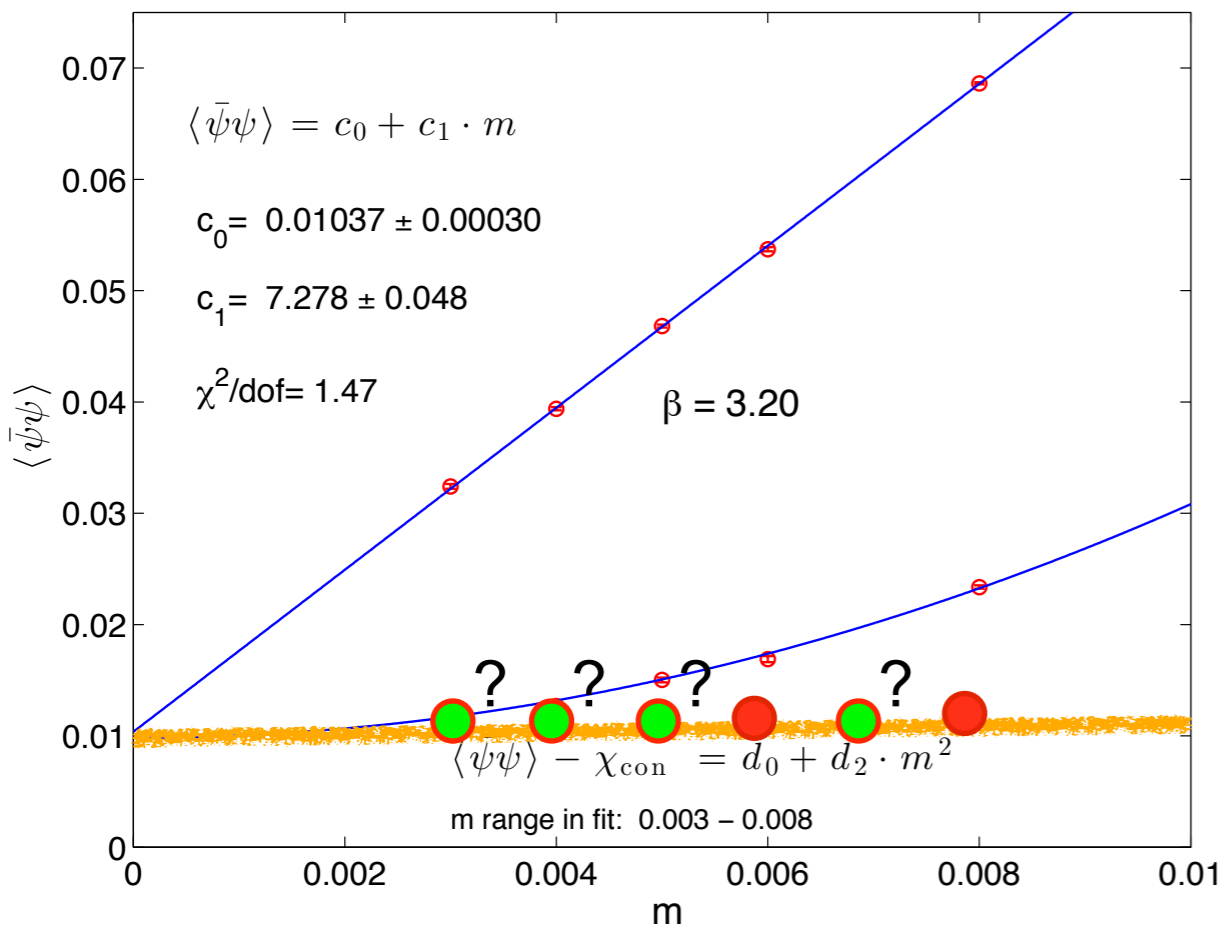
Scale dependence of Mode number distribution  $\nu(M,n)$



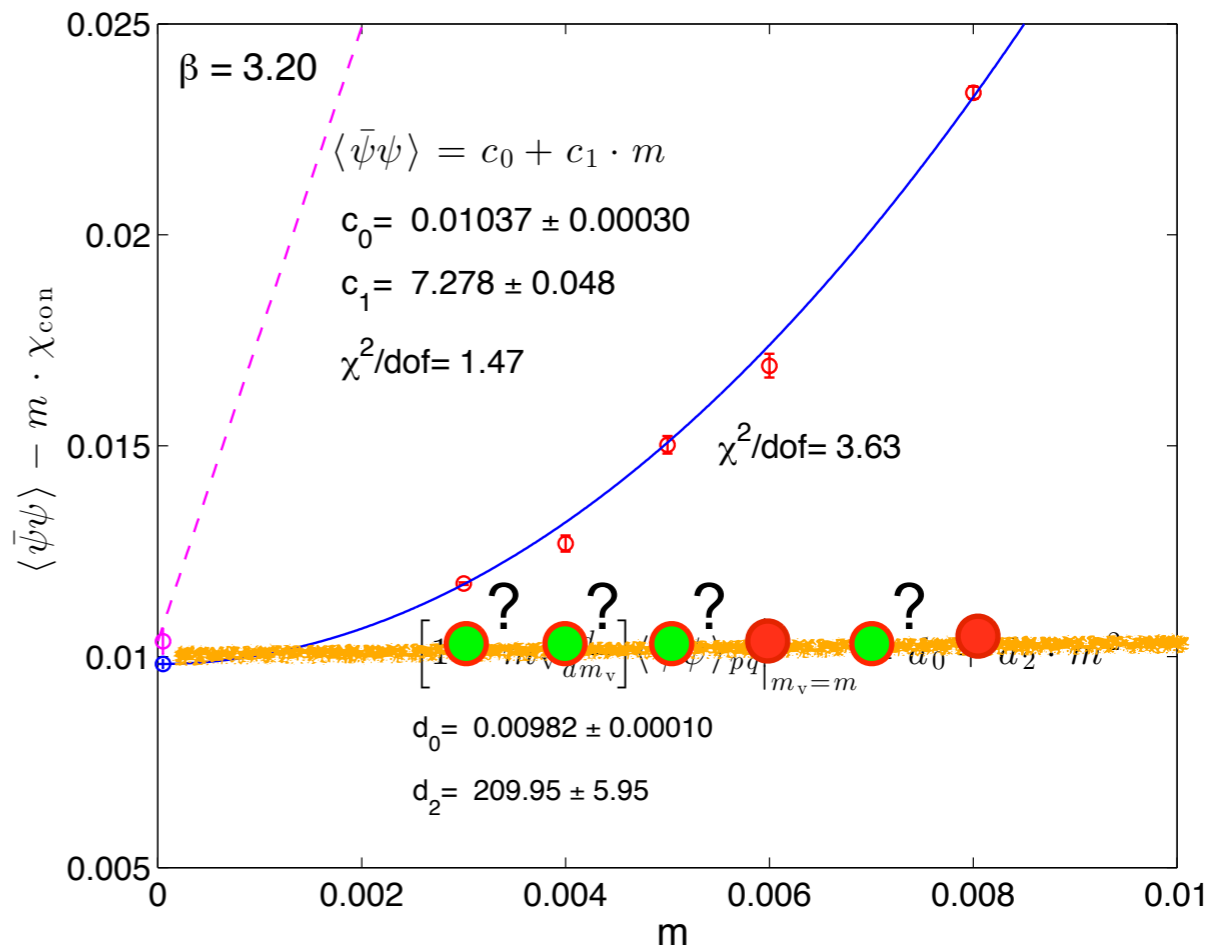
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chiral condensate and its subtracted form



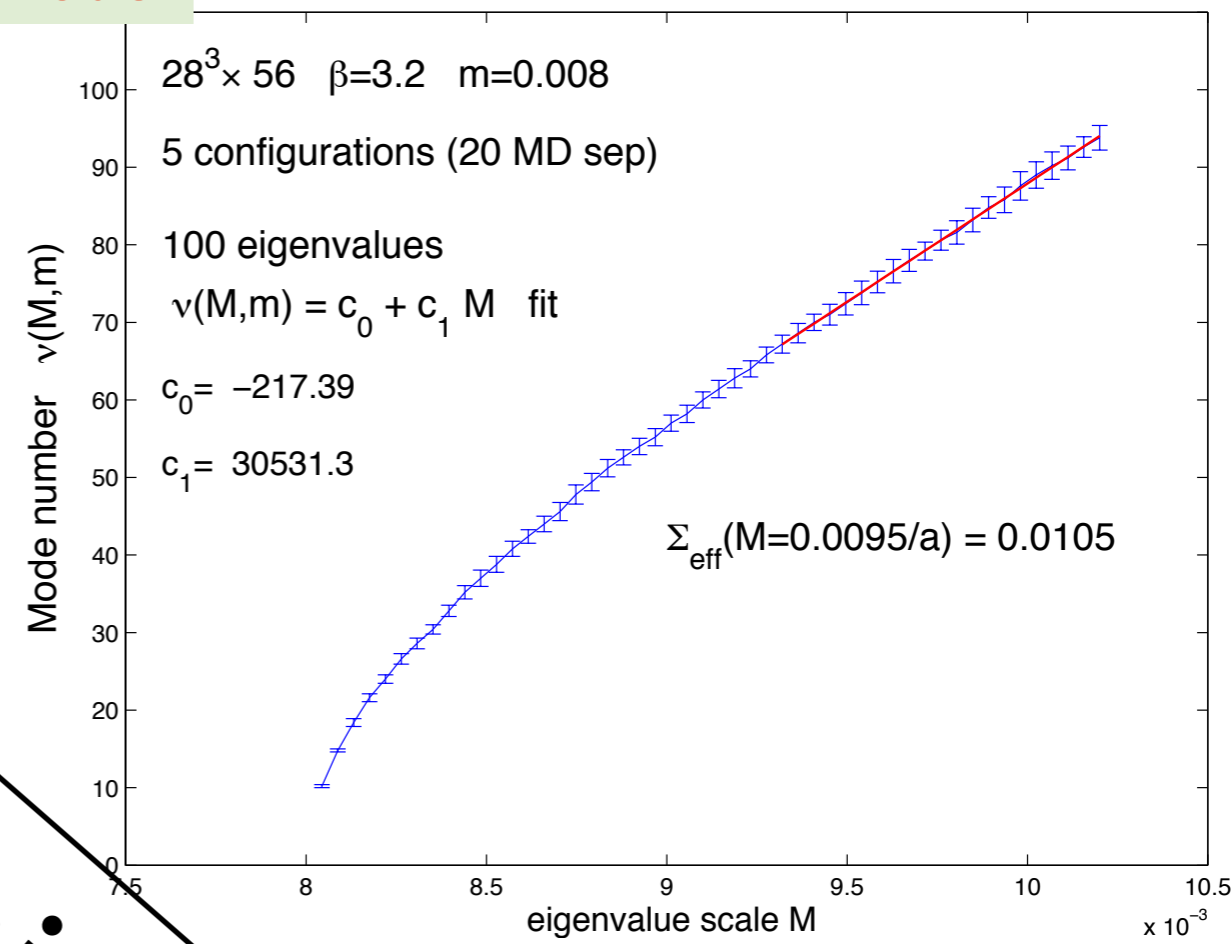
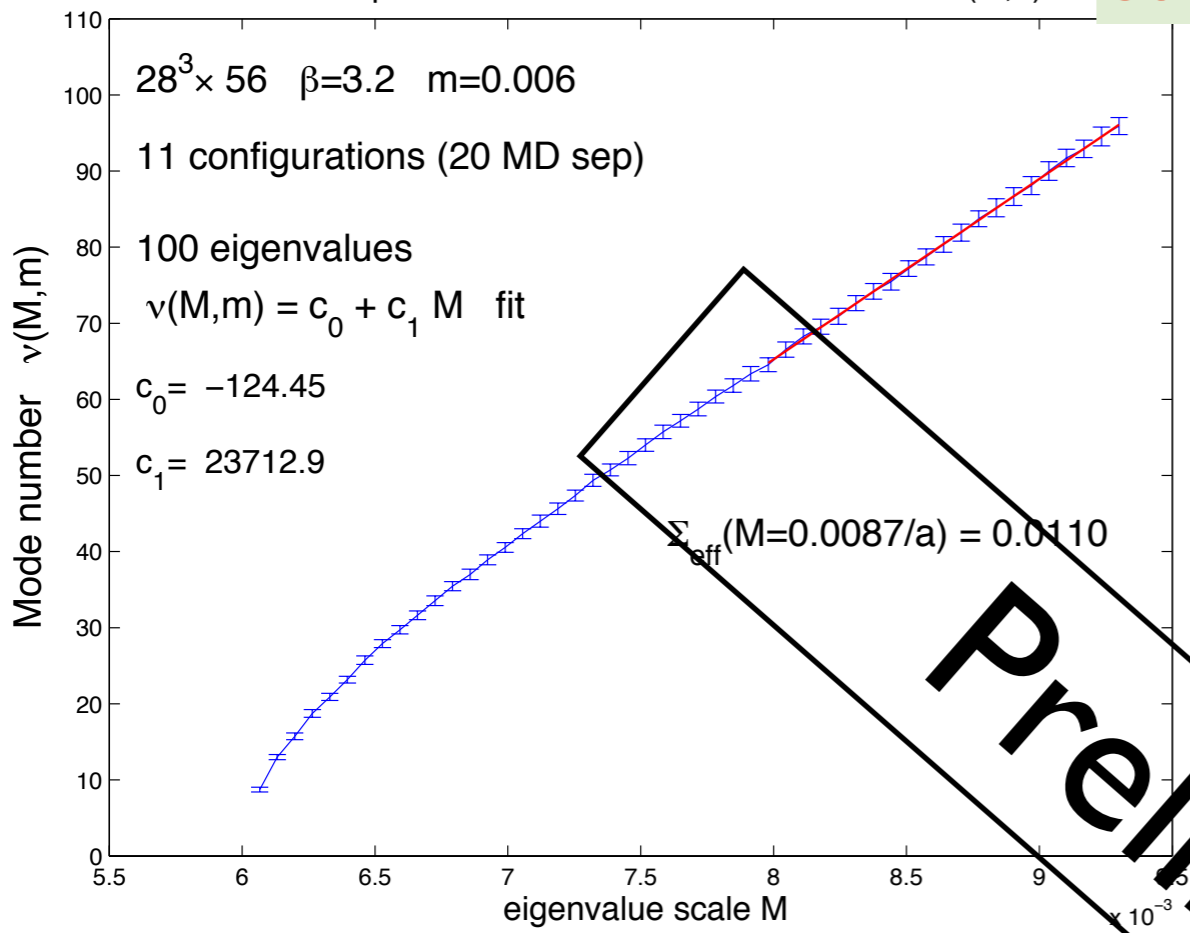
subtracted chiral condensate



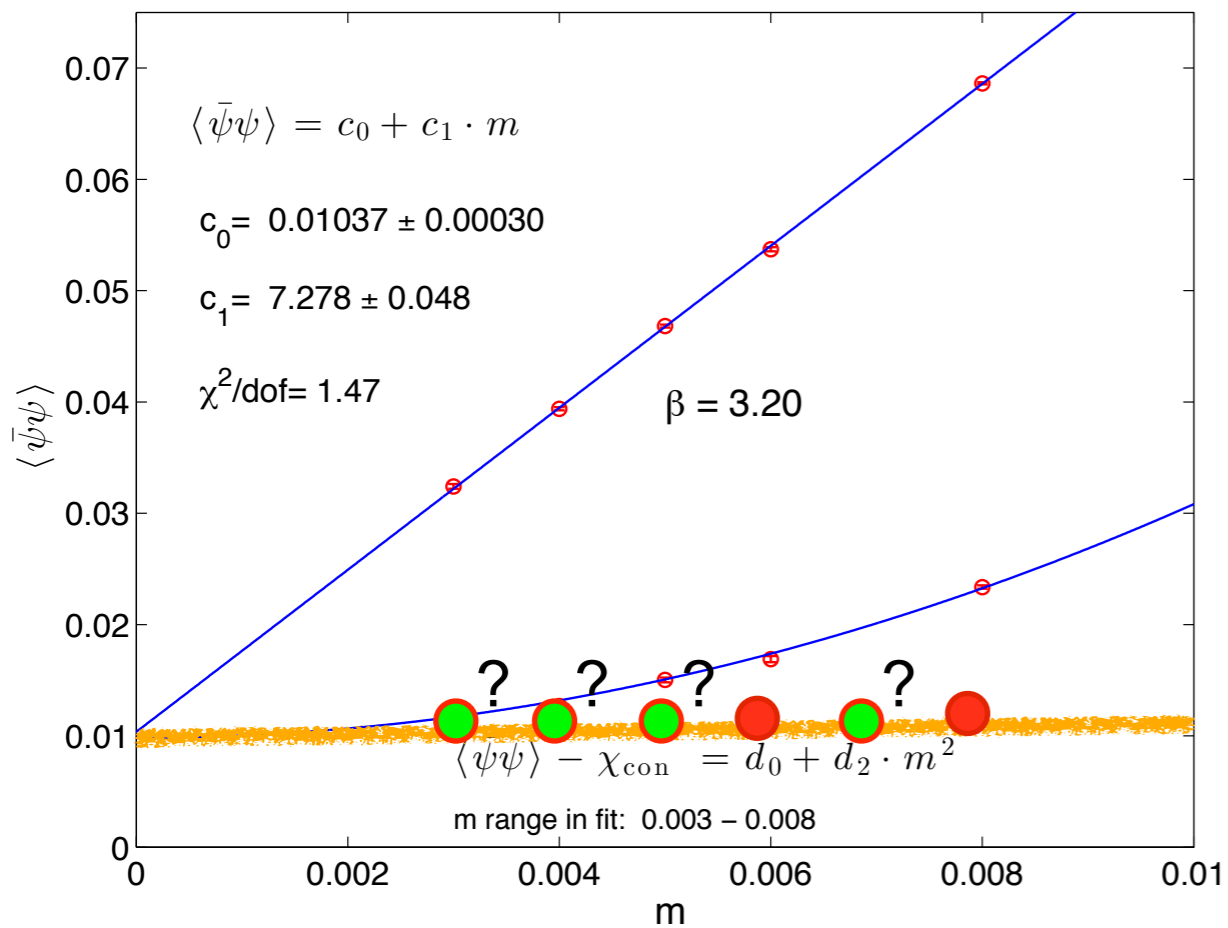
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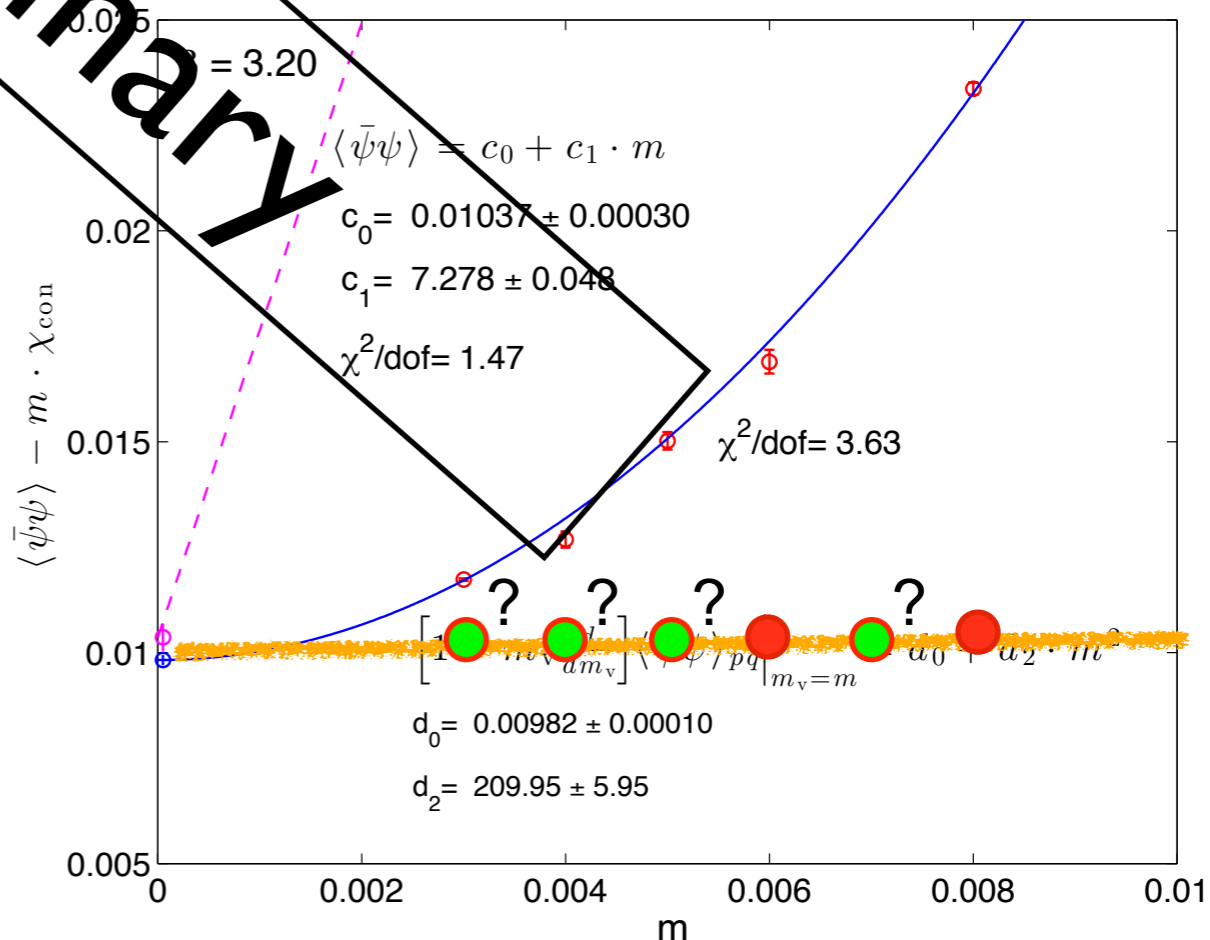
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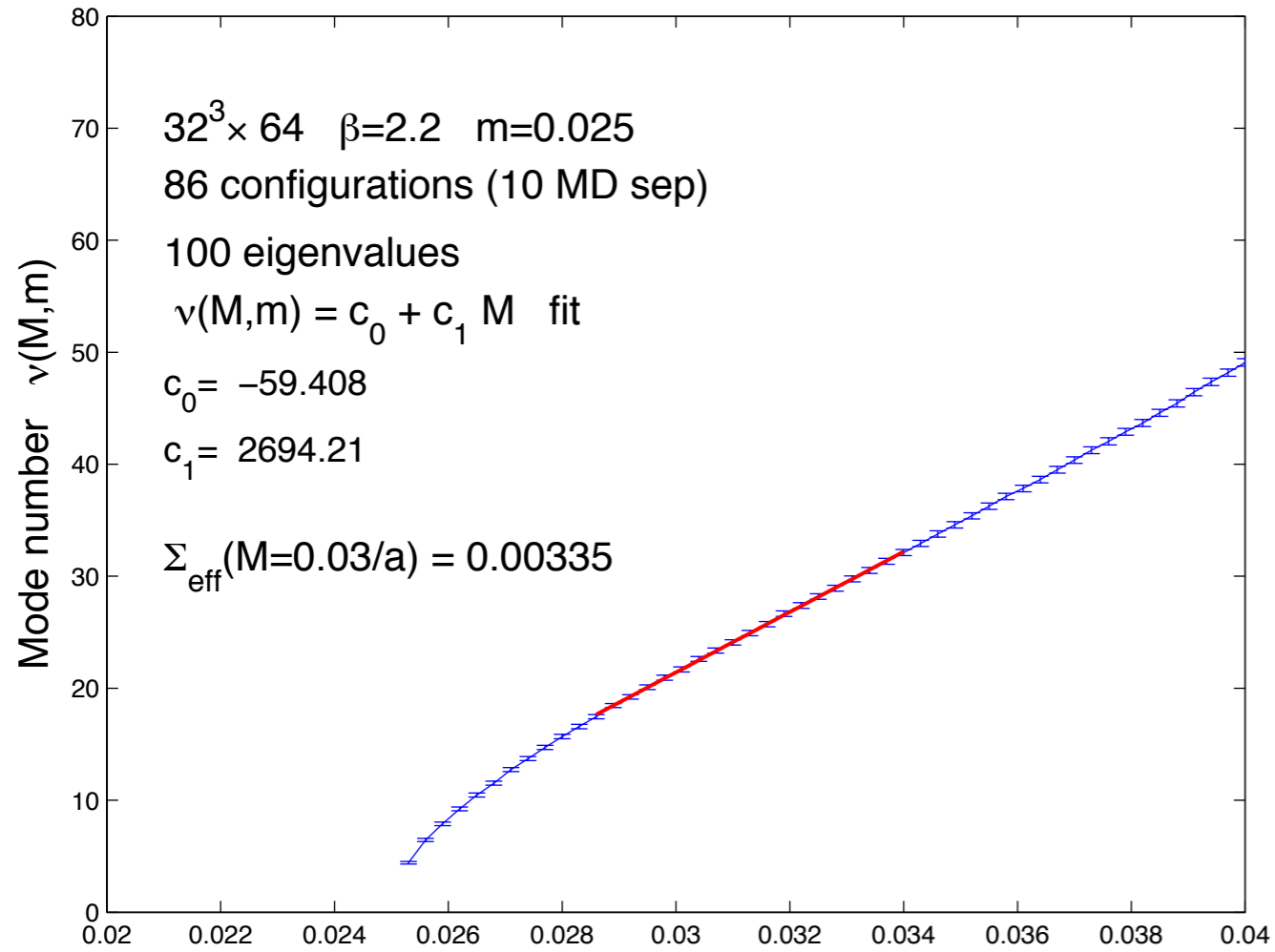
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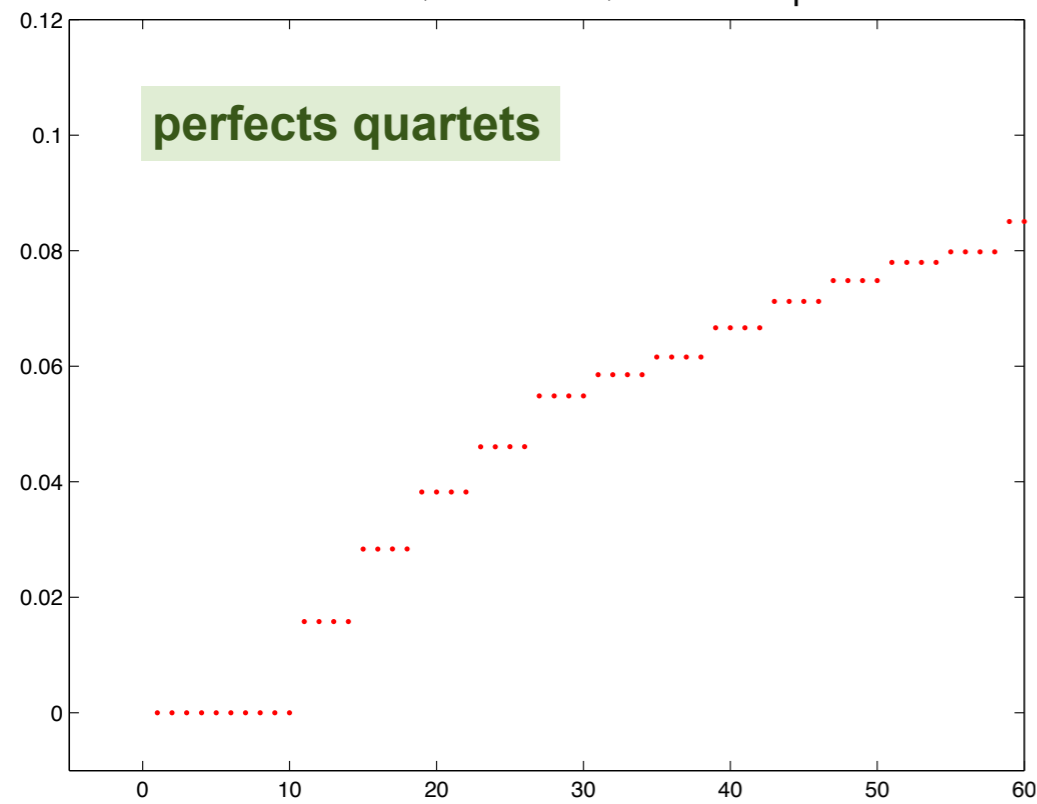
Preliminary

# Nf=12 fundamental rep:

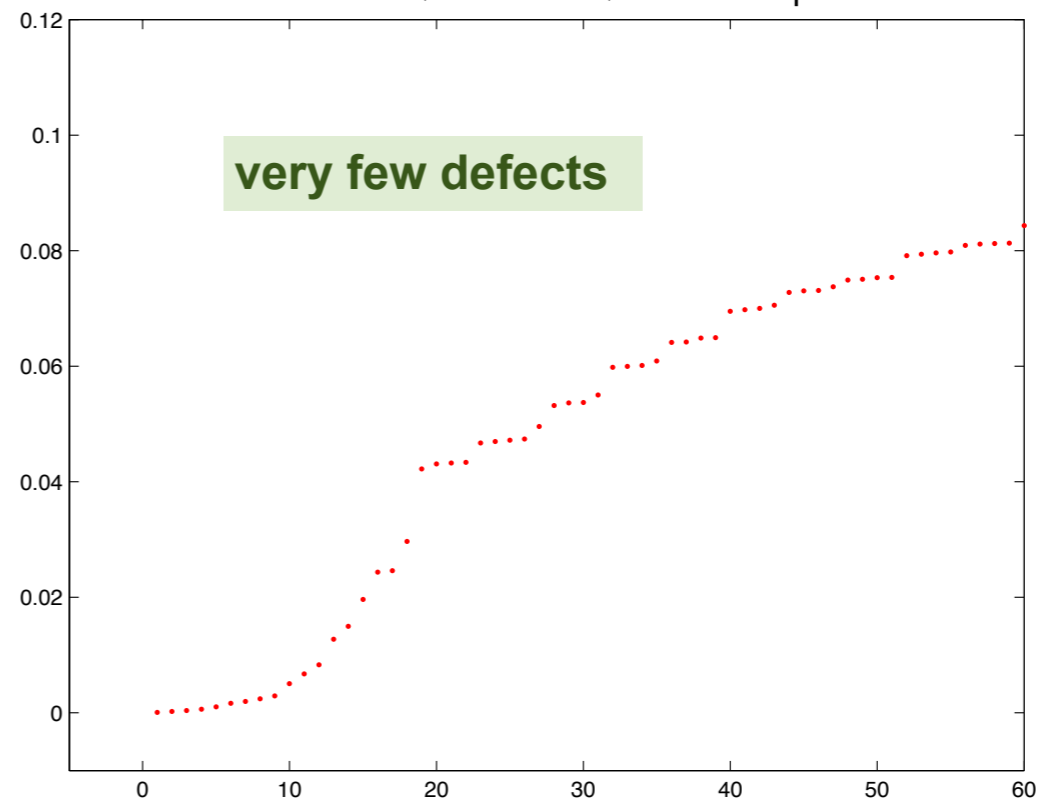
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sextet model Quartets of  $Q=-1$  Dirac spectrum

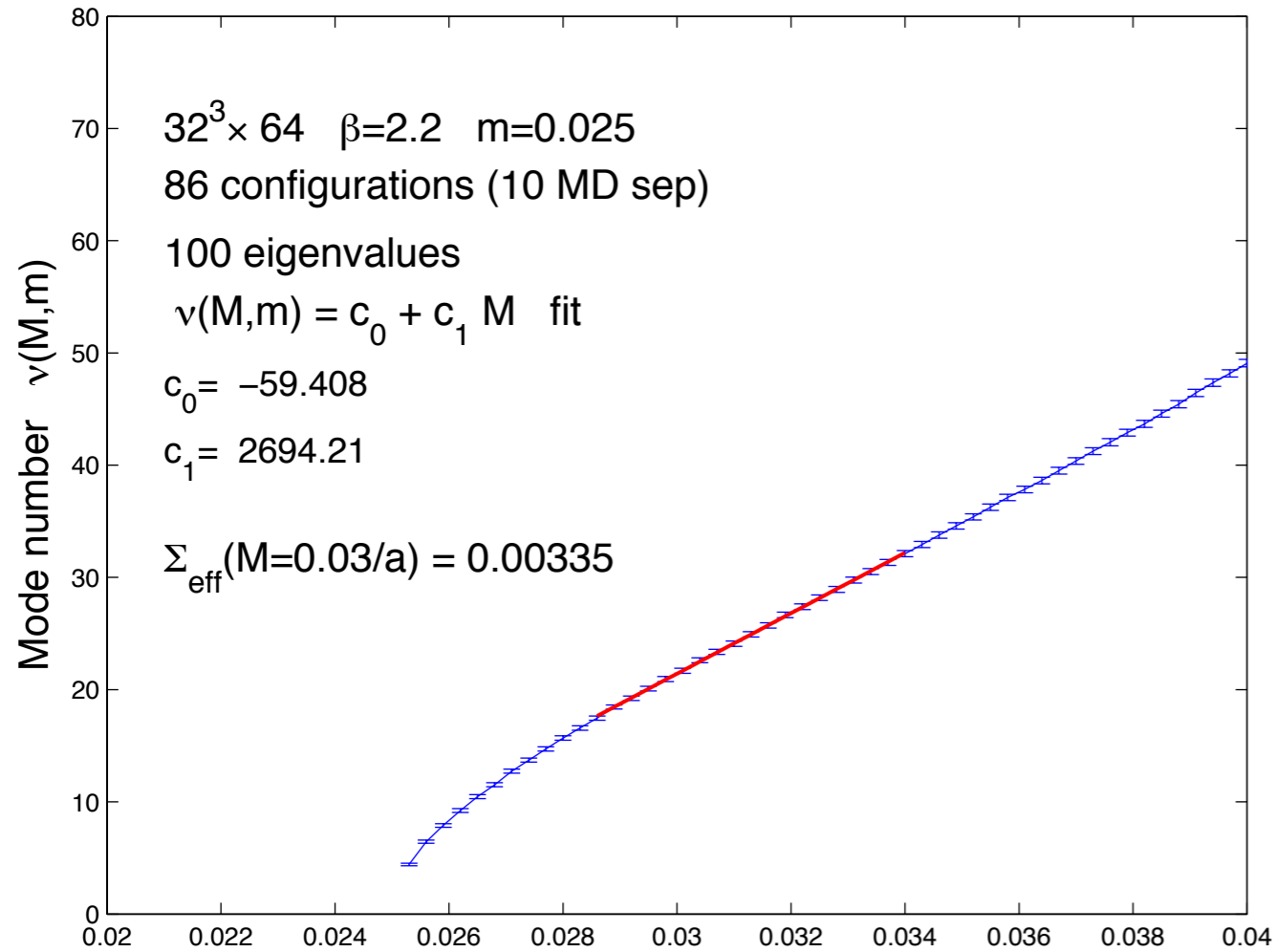


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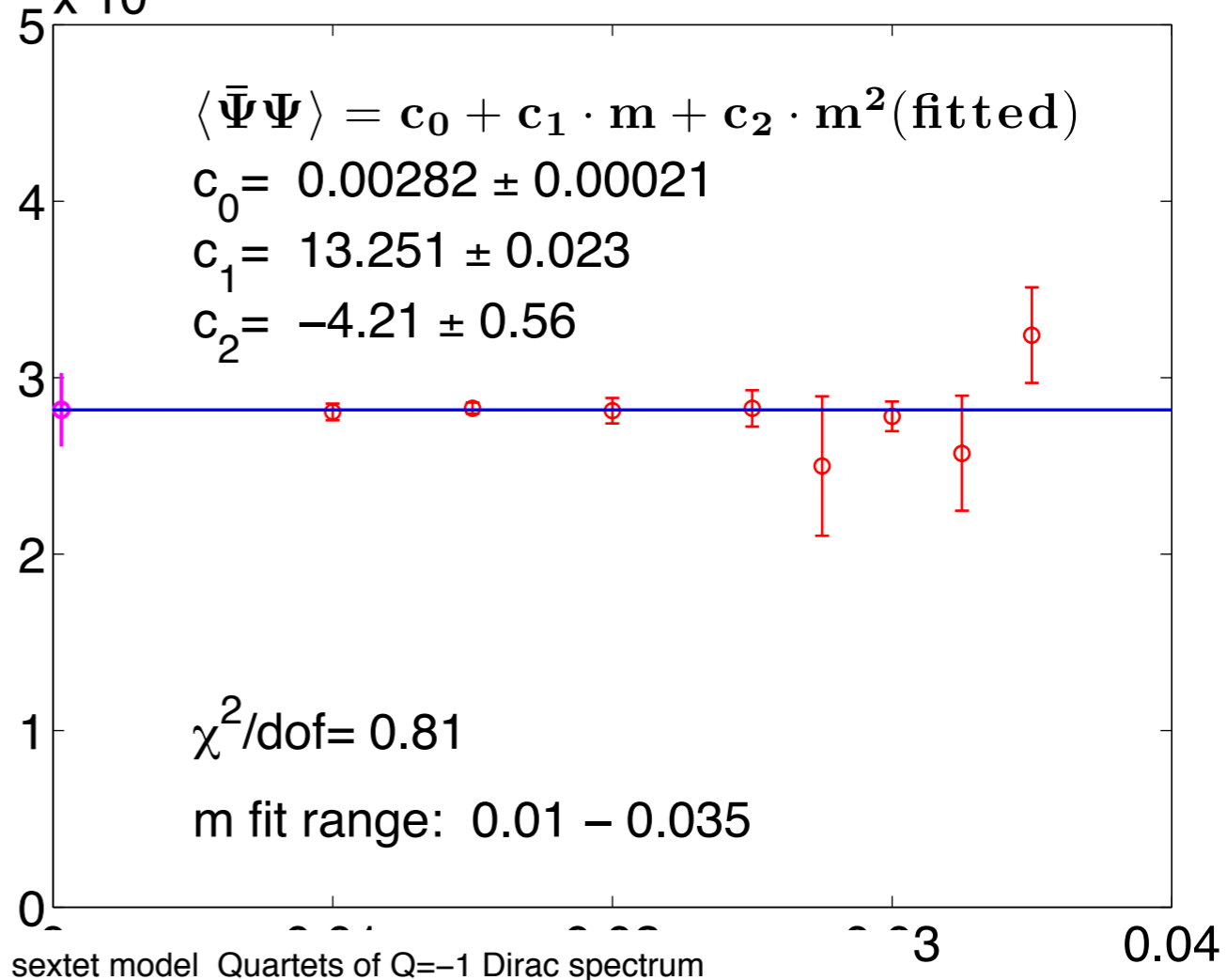


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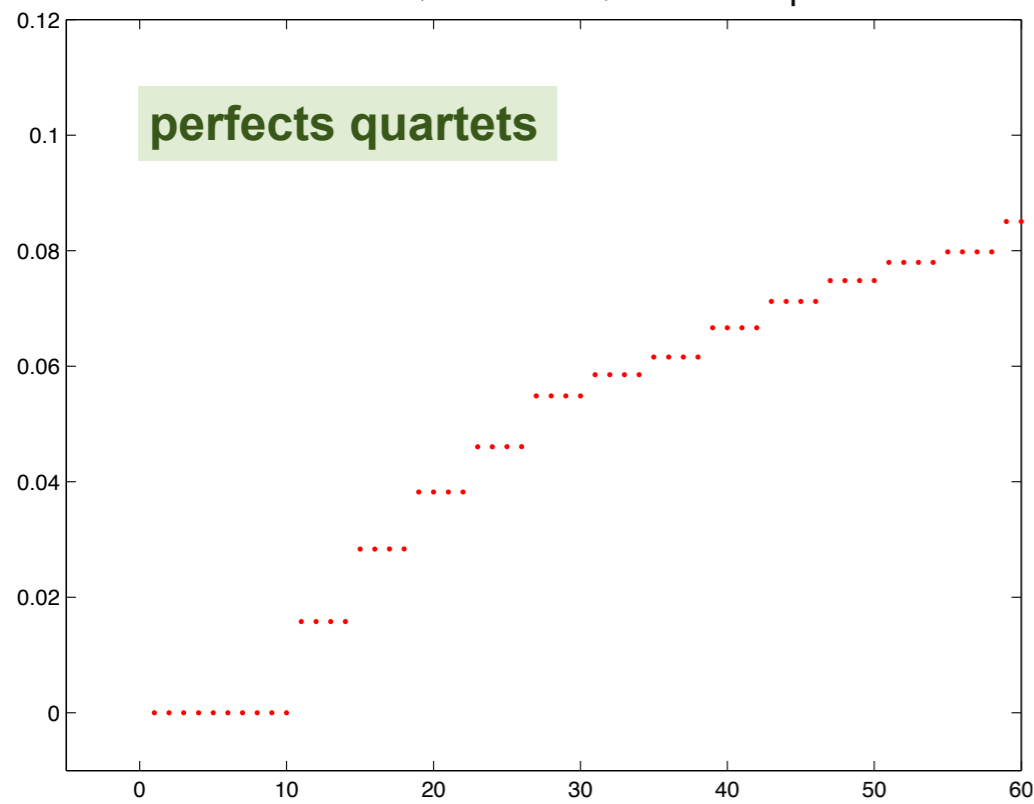
Scale dependence of Mode number distribution  $v(M,n)$



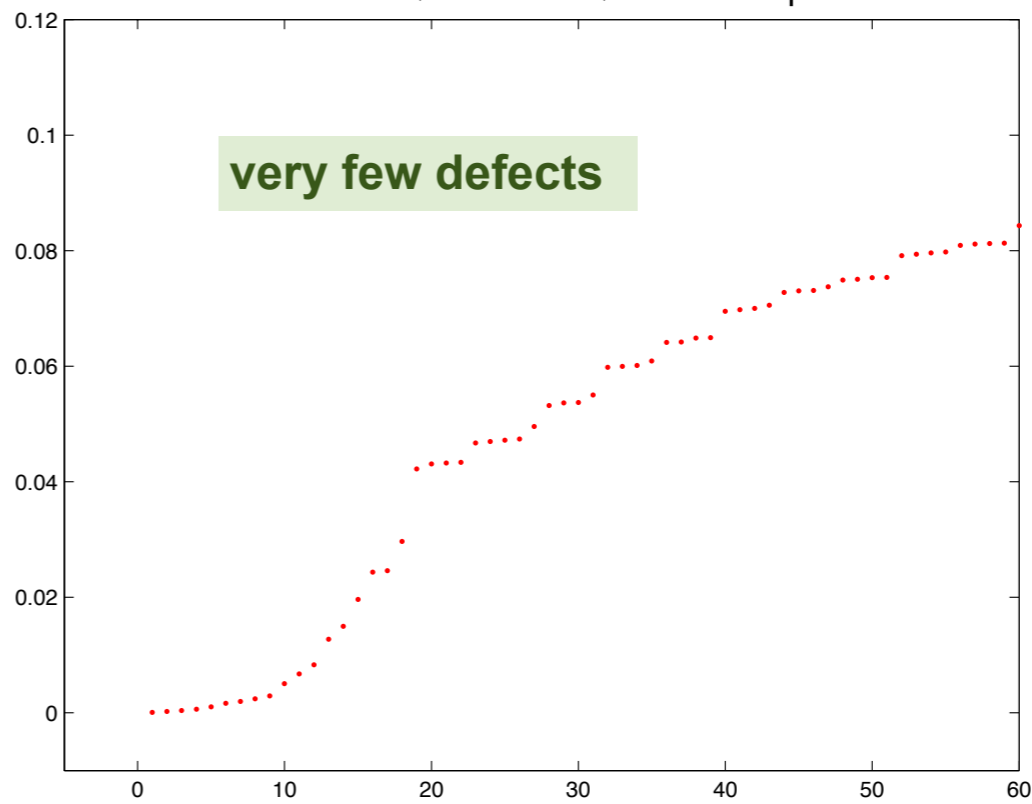
$\langle \bar{\Psi}\Psi \rangle - c_1 m - c_2 m^2$  plotted



sextet model Quartets of Q=-1 Dirac spectrum



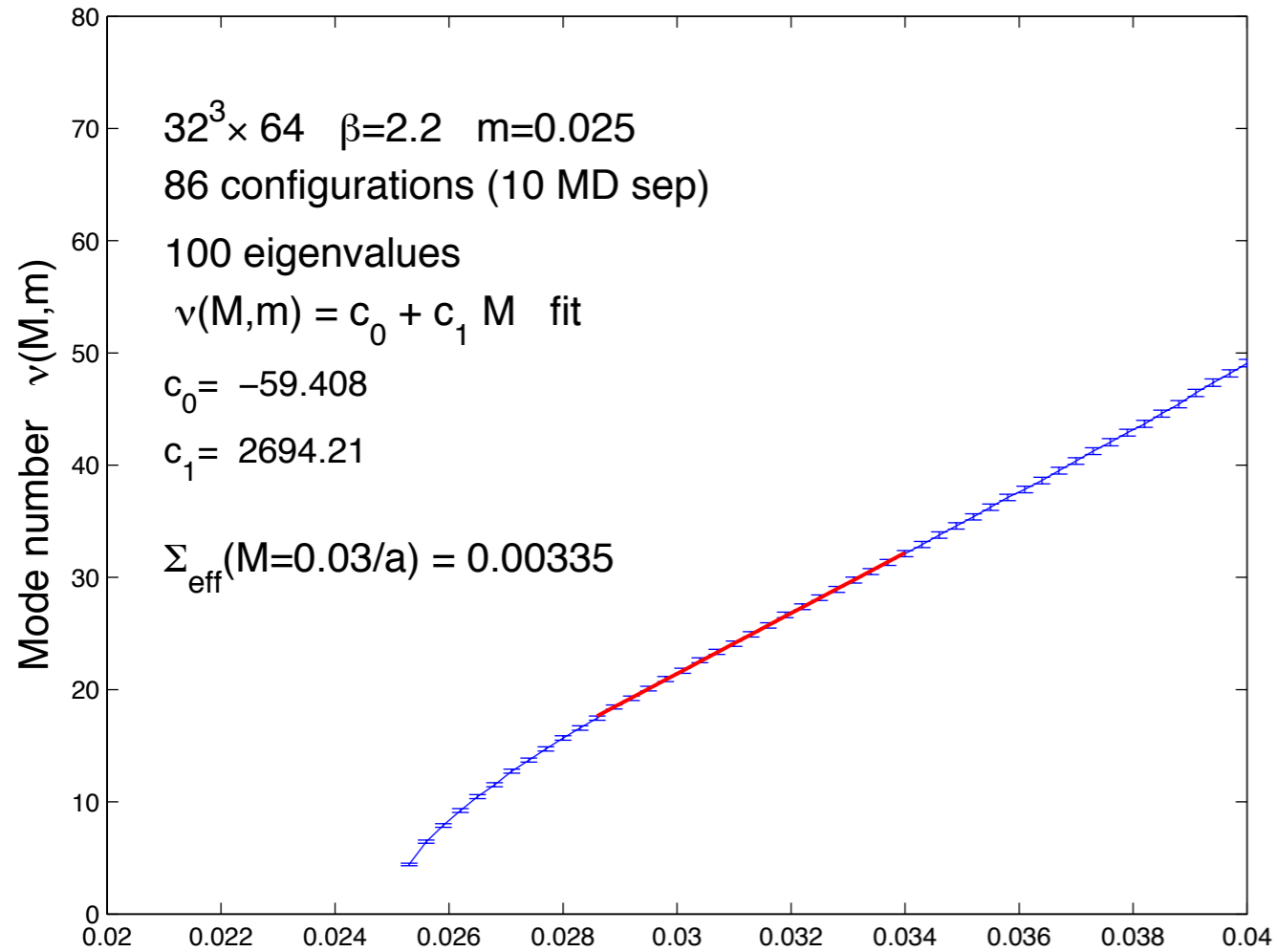
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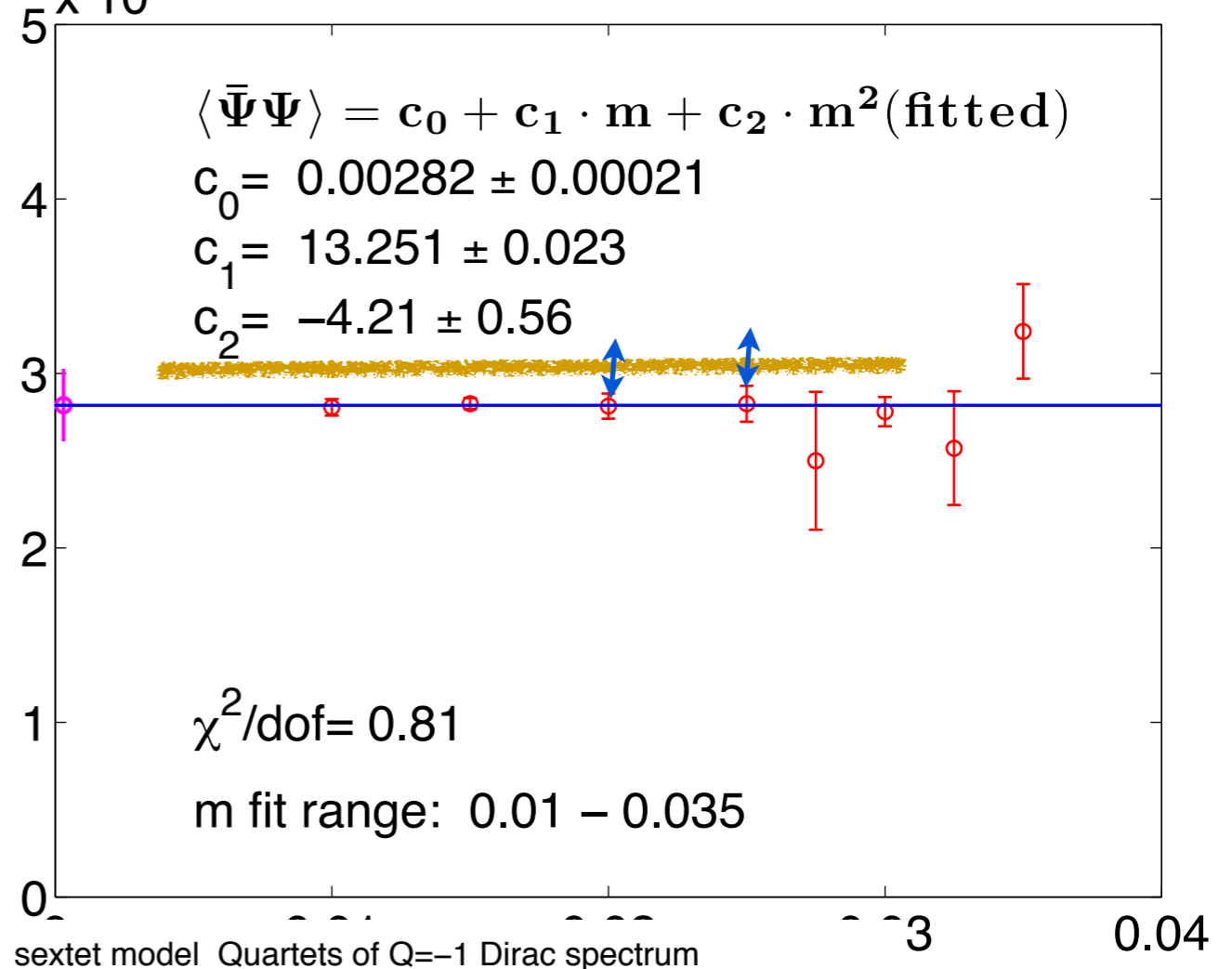


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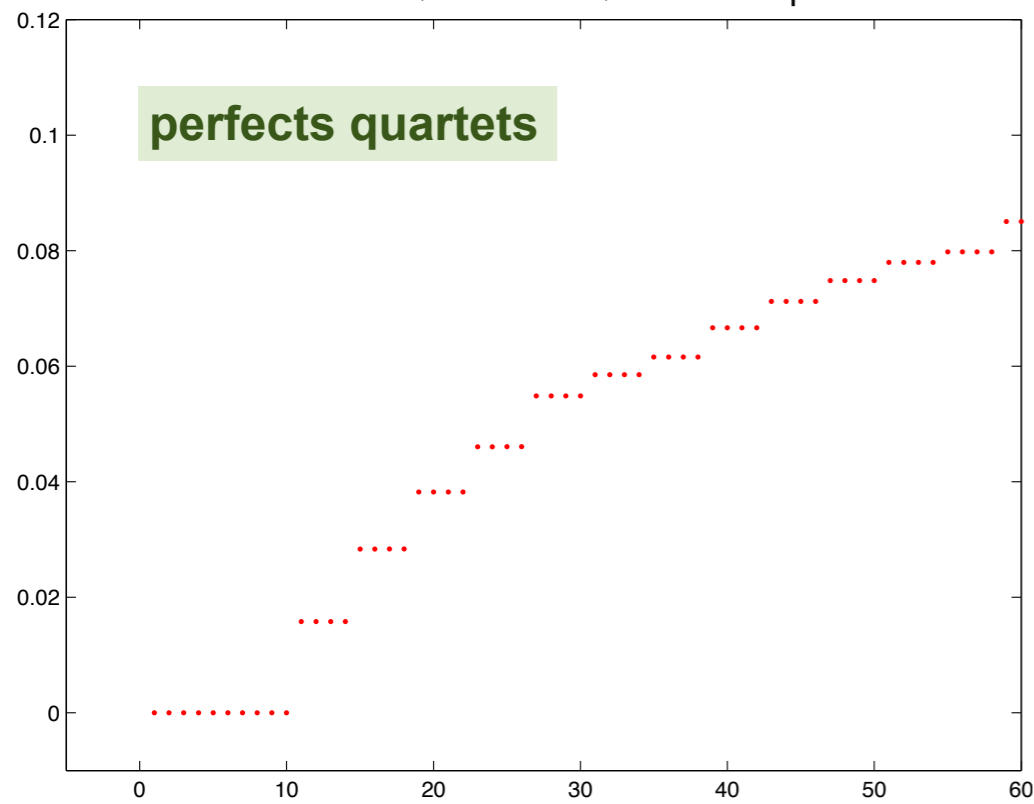
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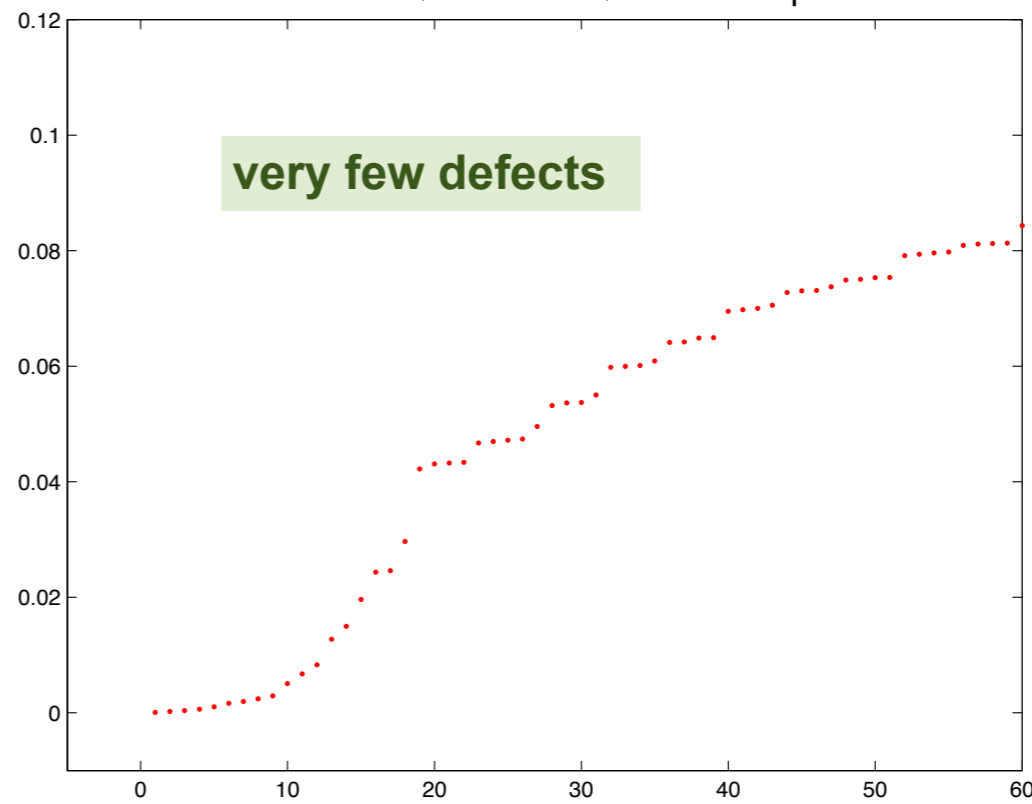
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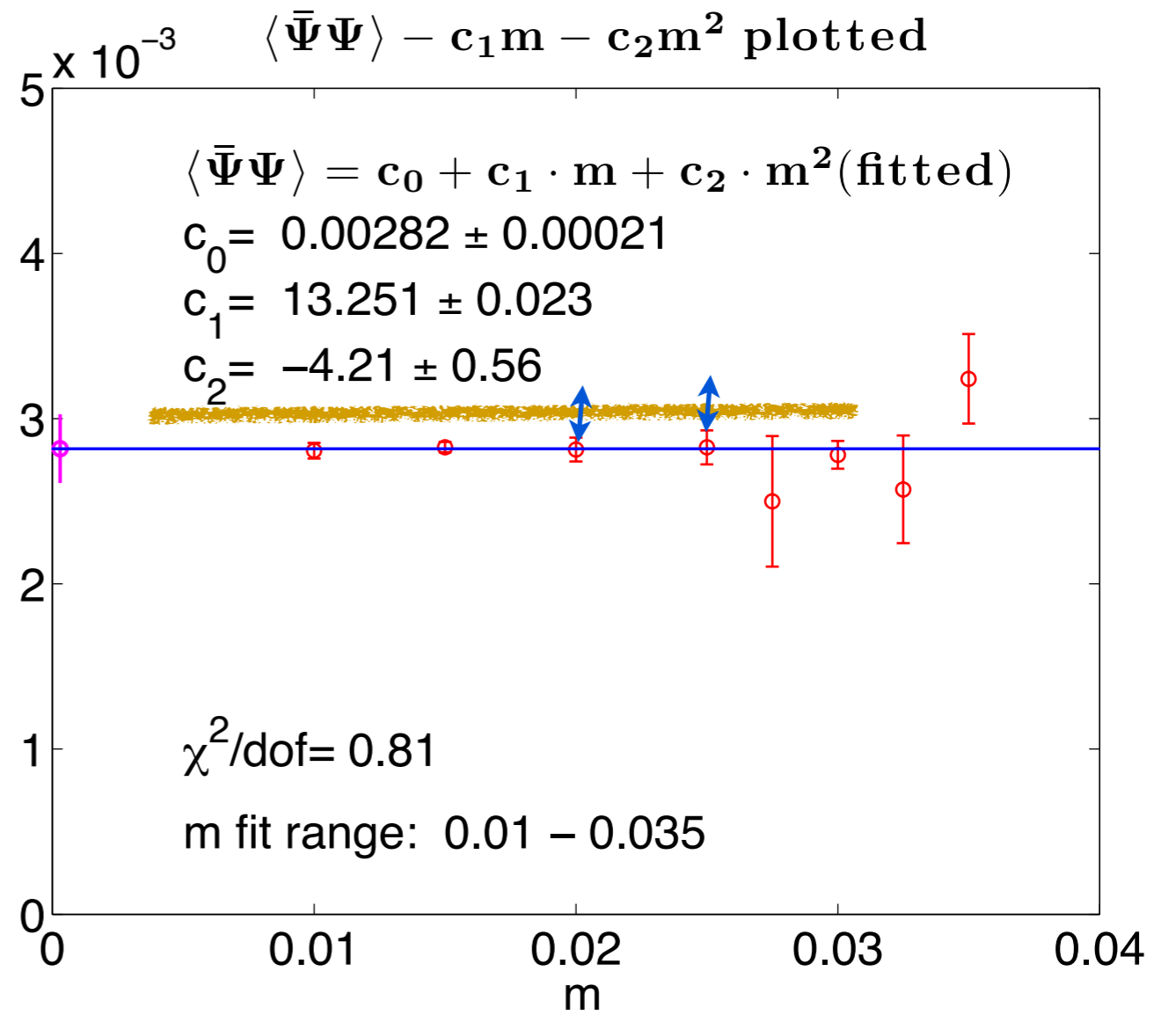
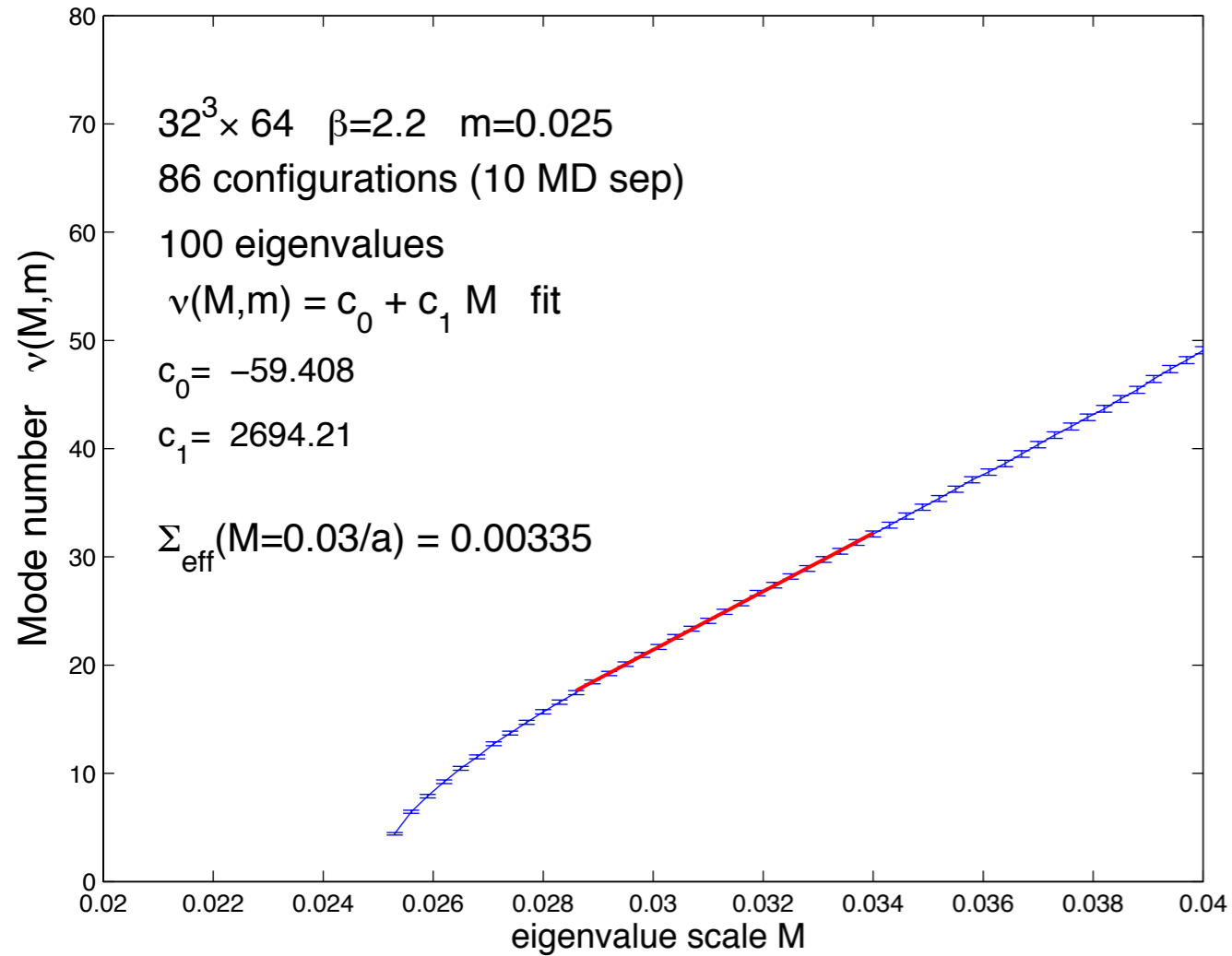


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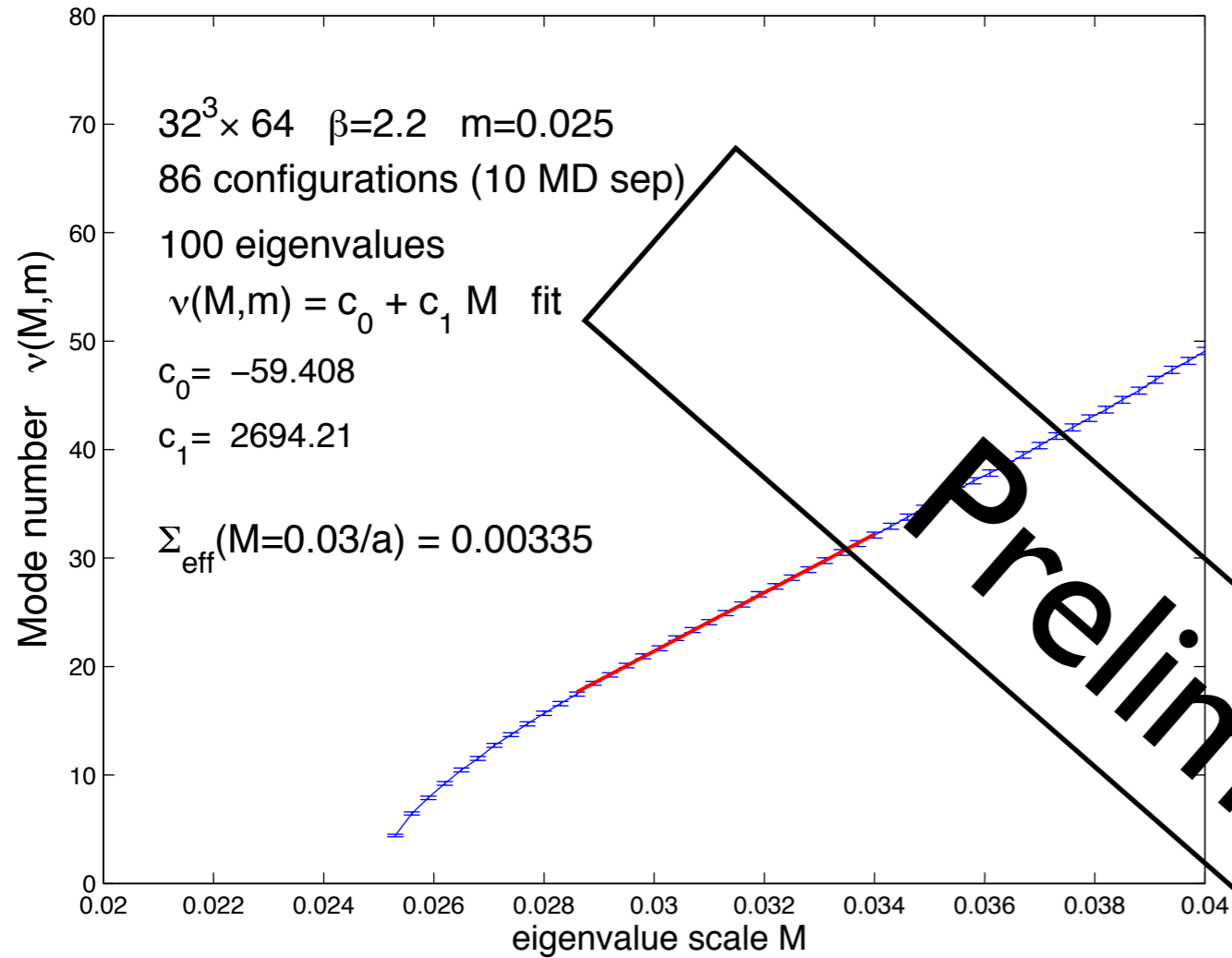
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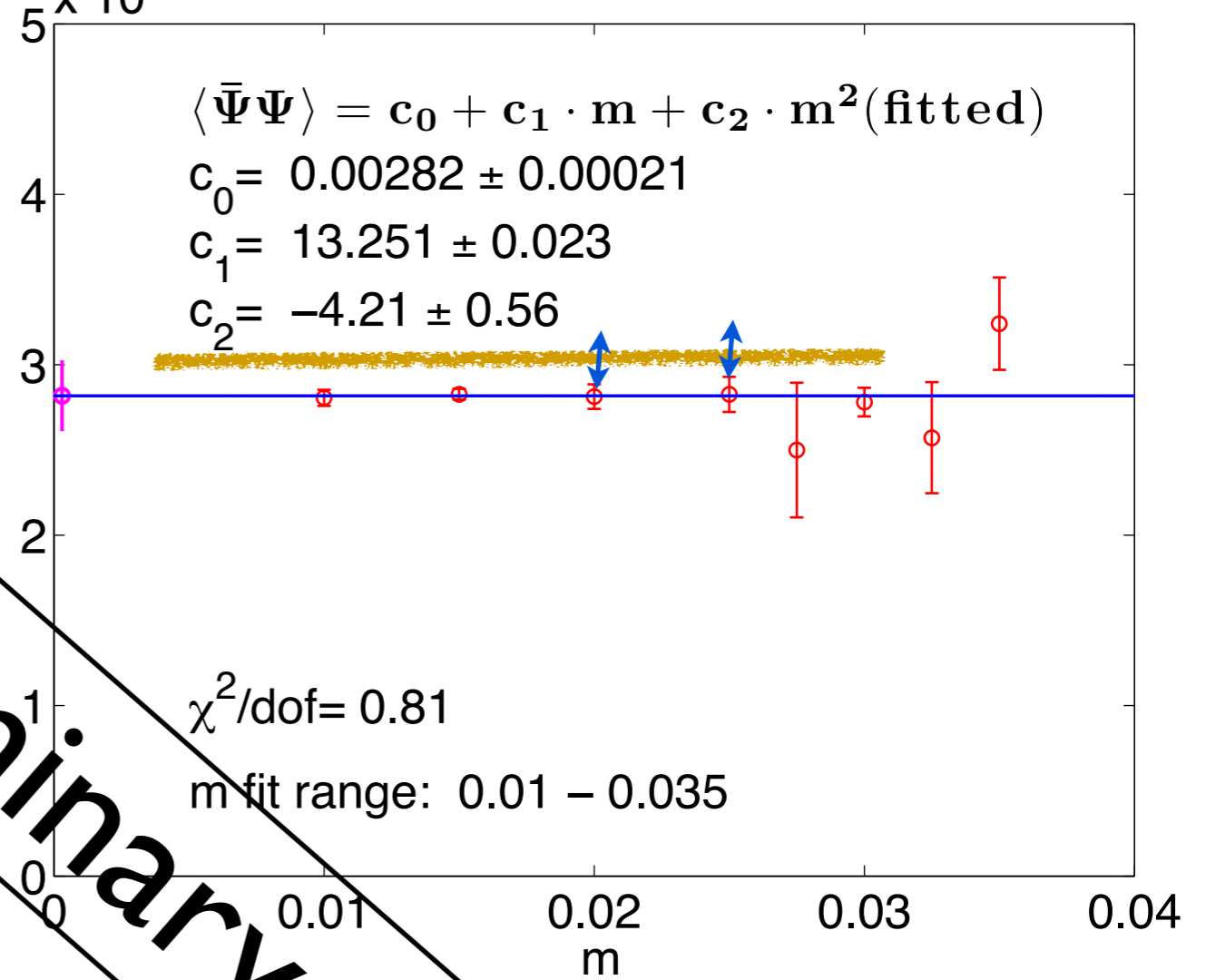


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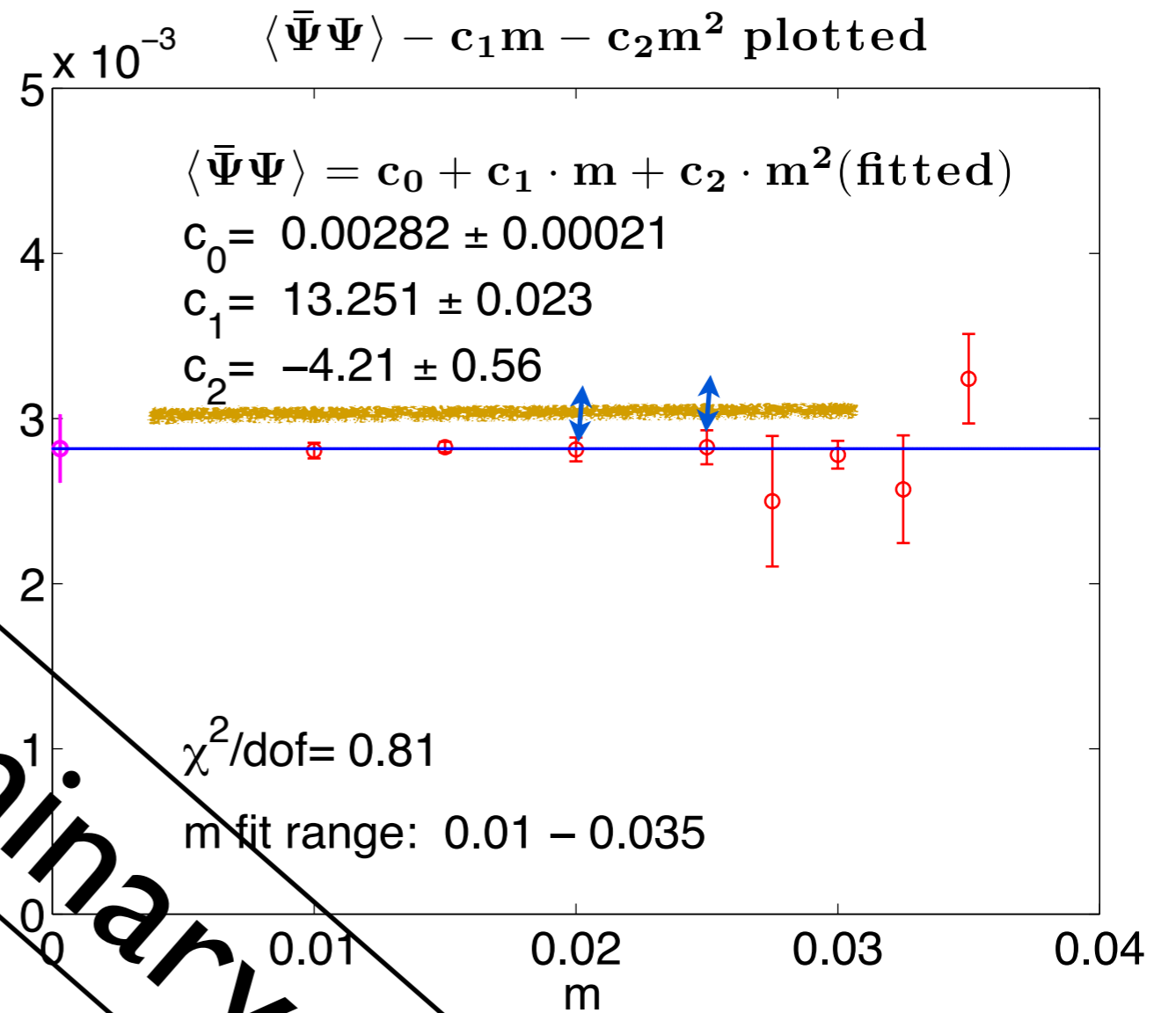
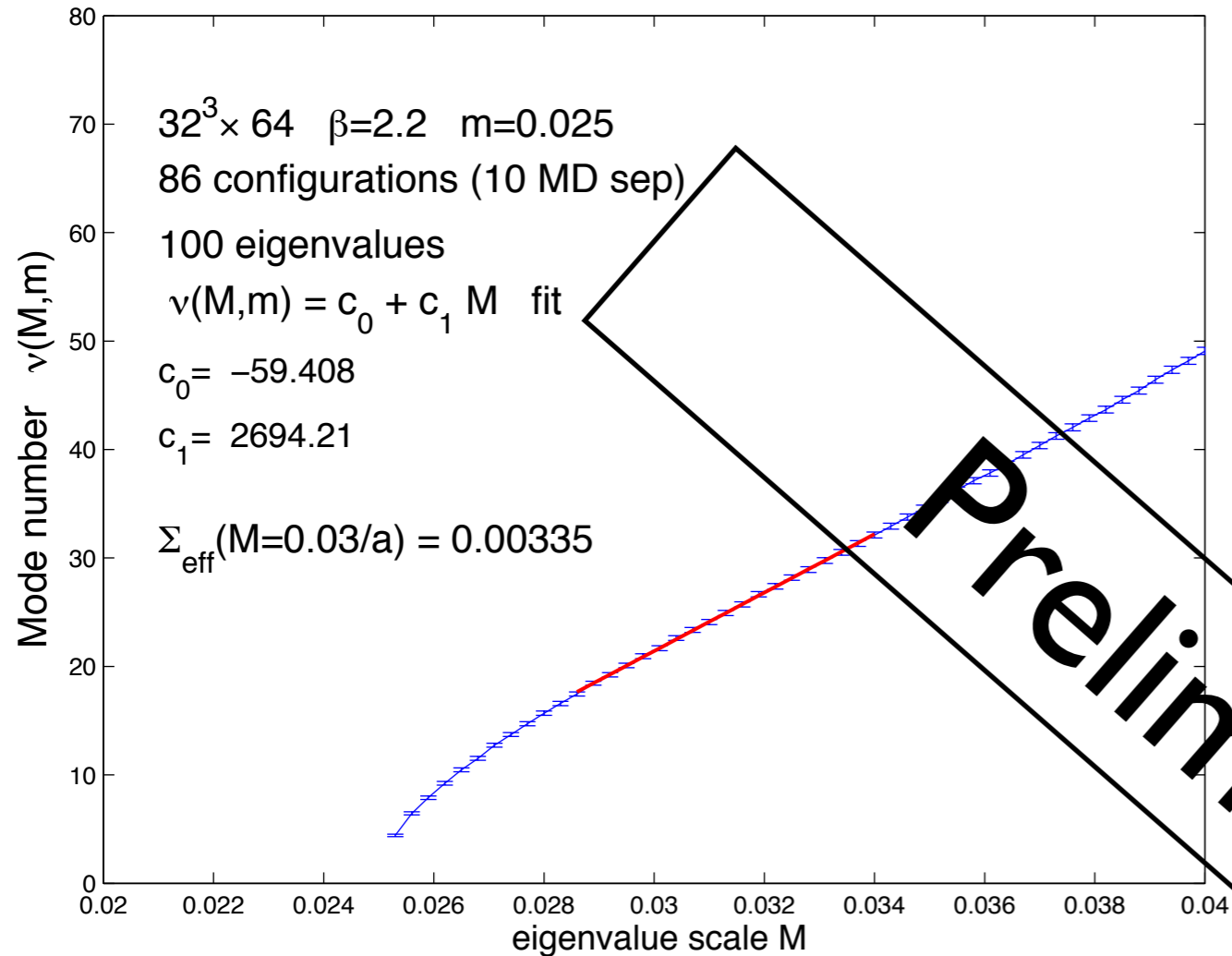
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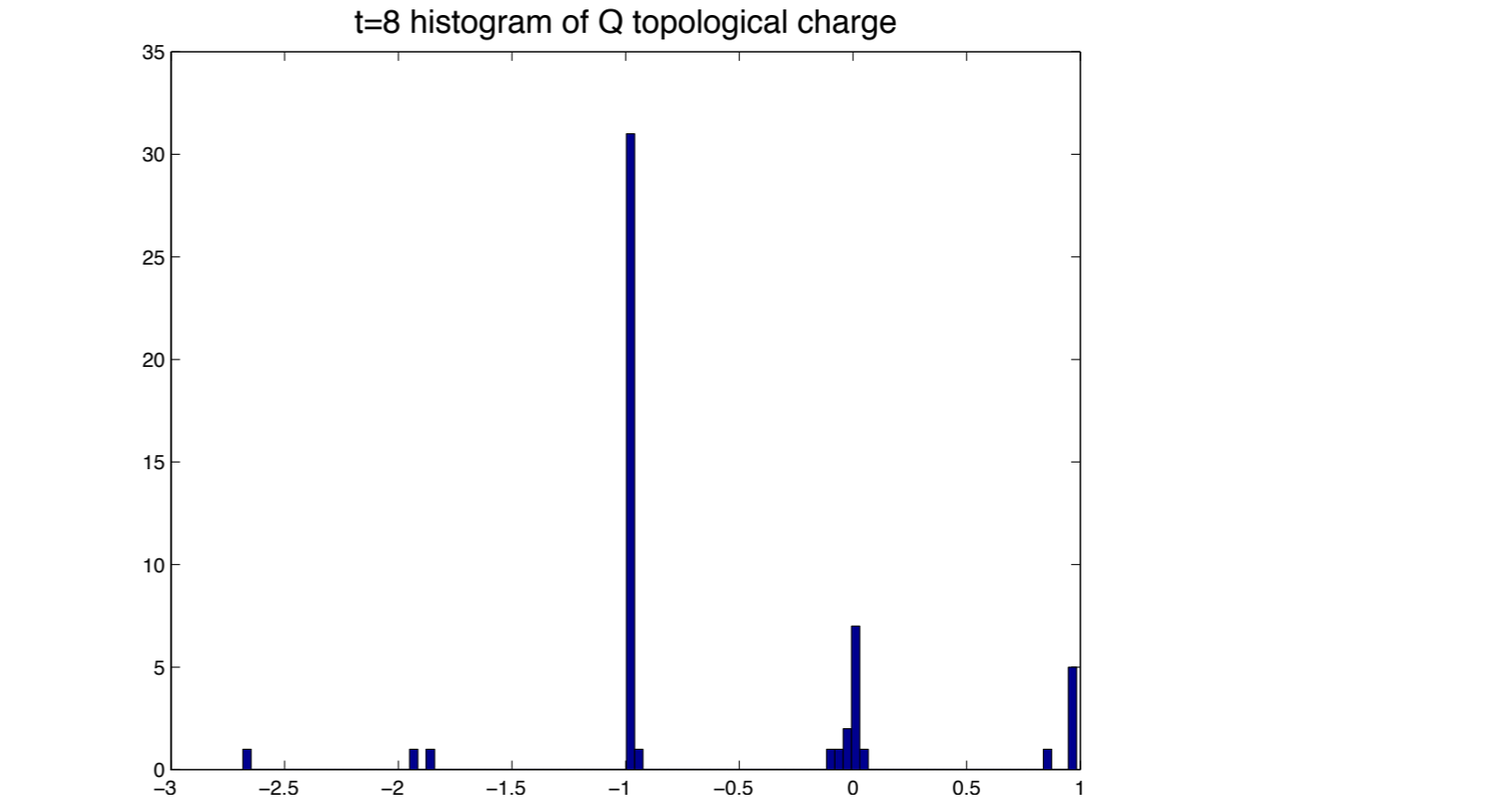
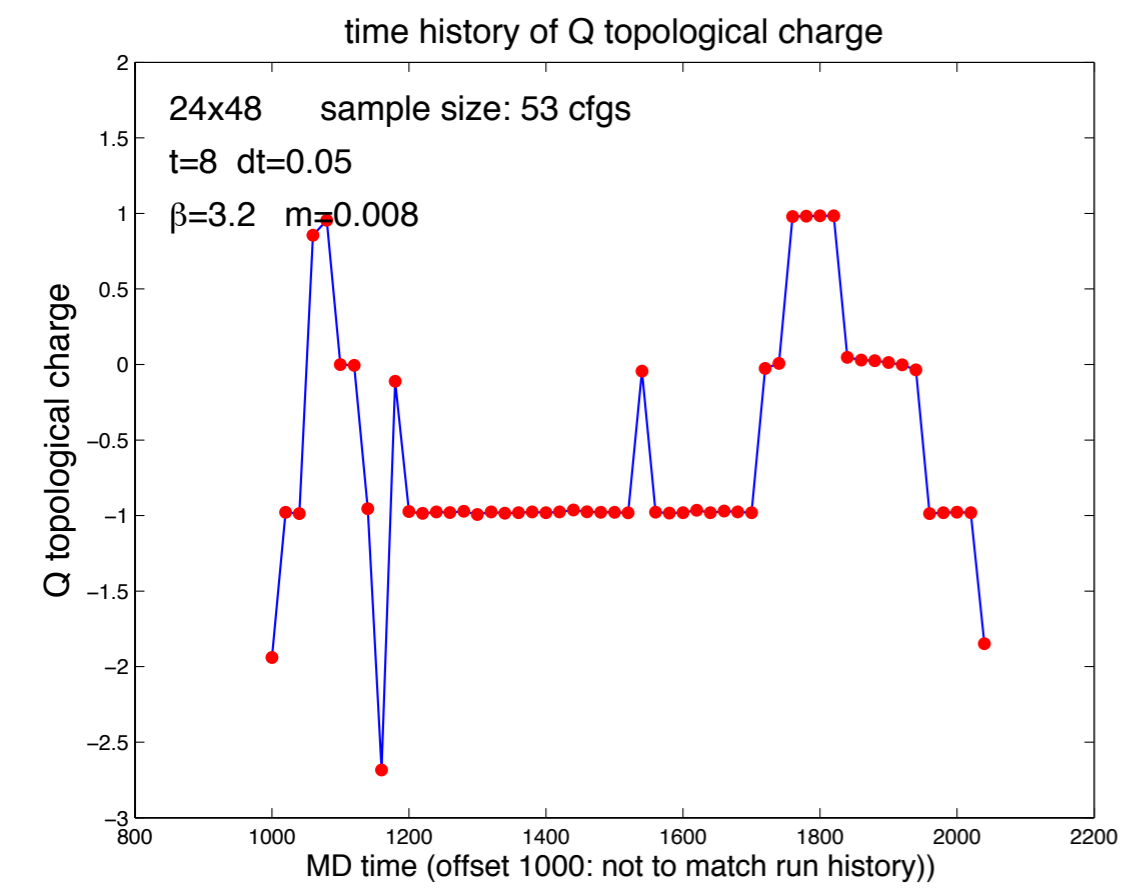
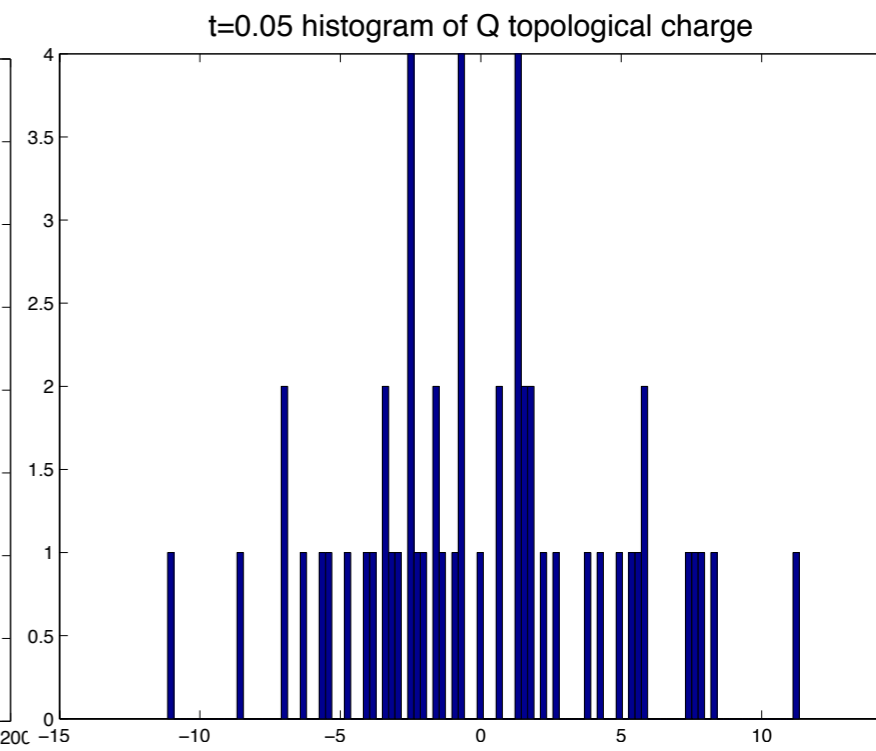
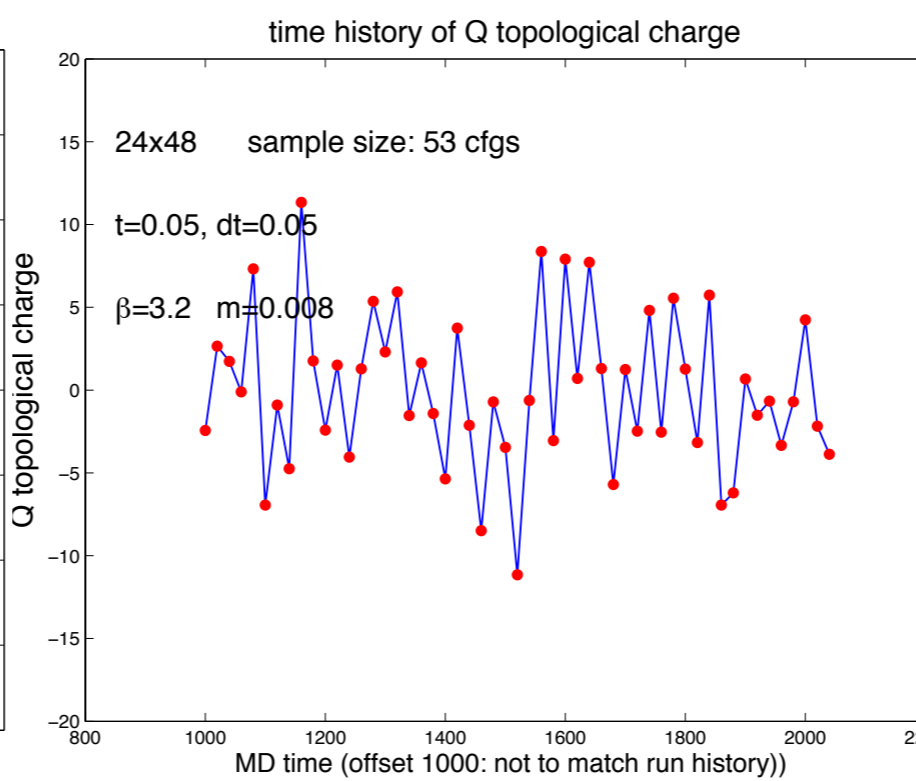
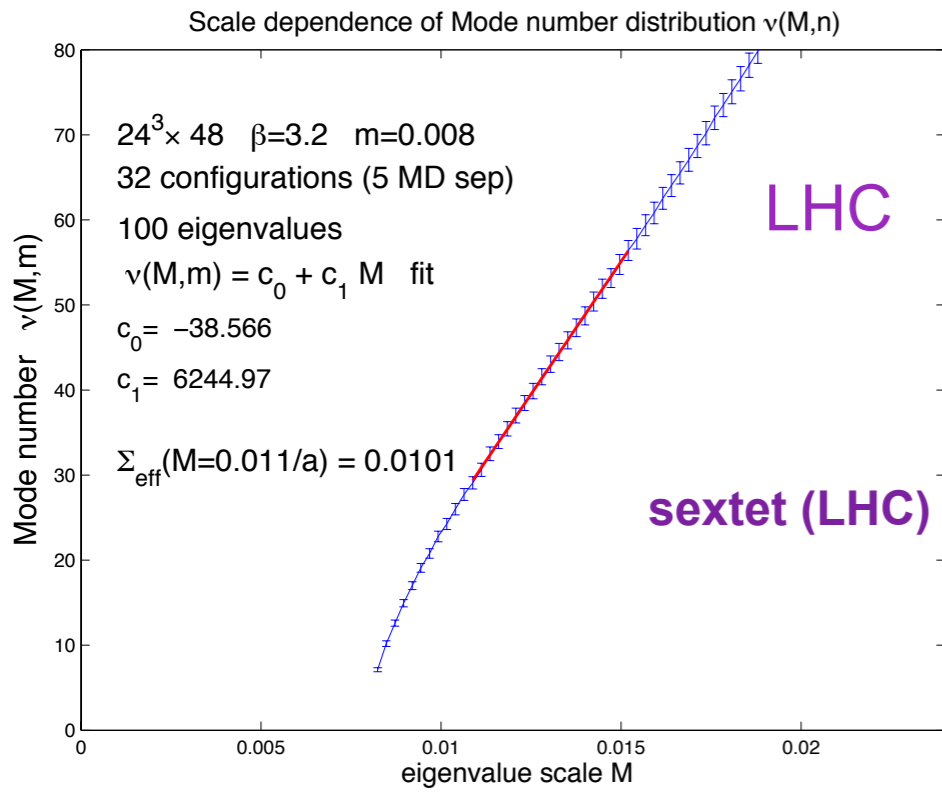
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is Nf=12 chirally broken after all? :)

# gradient flow reveals problems with slowly developing topological distribution:



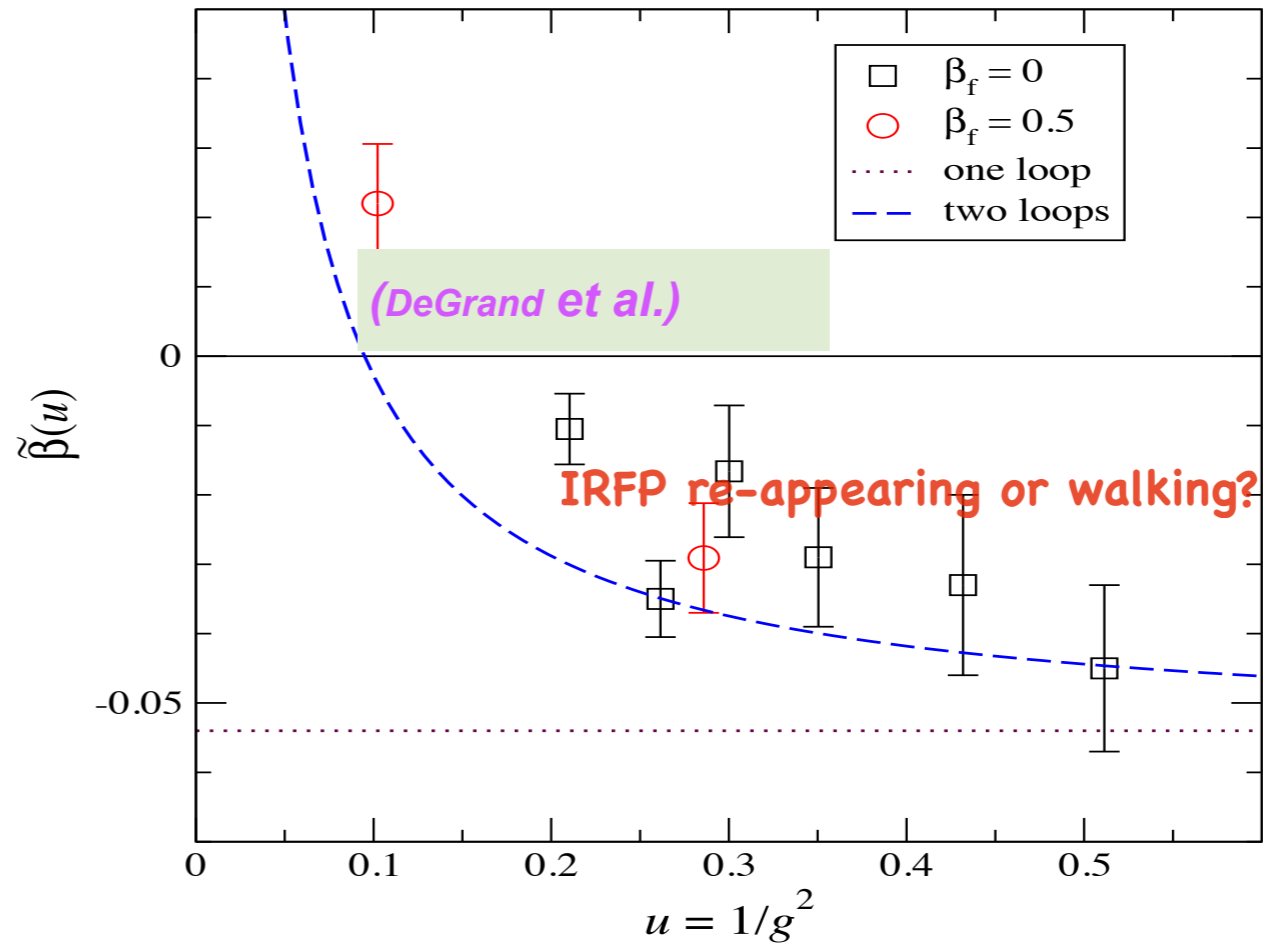
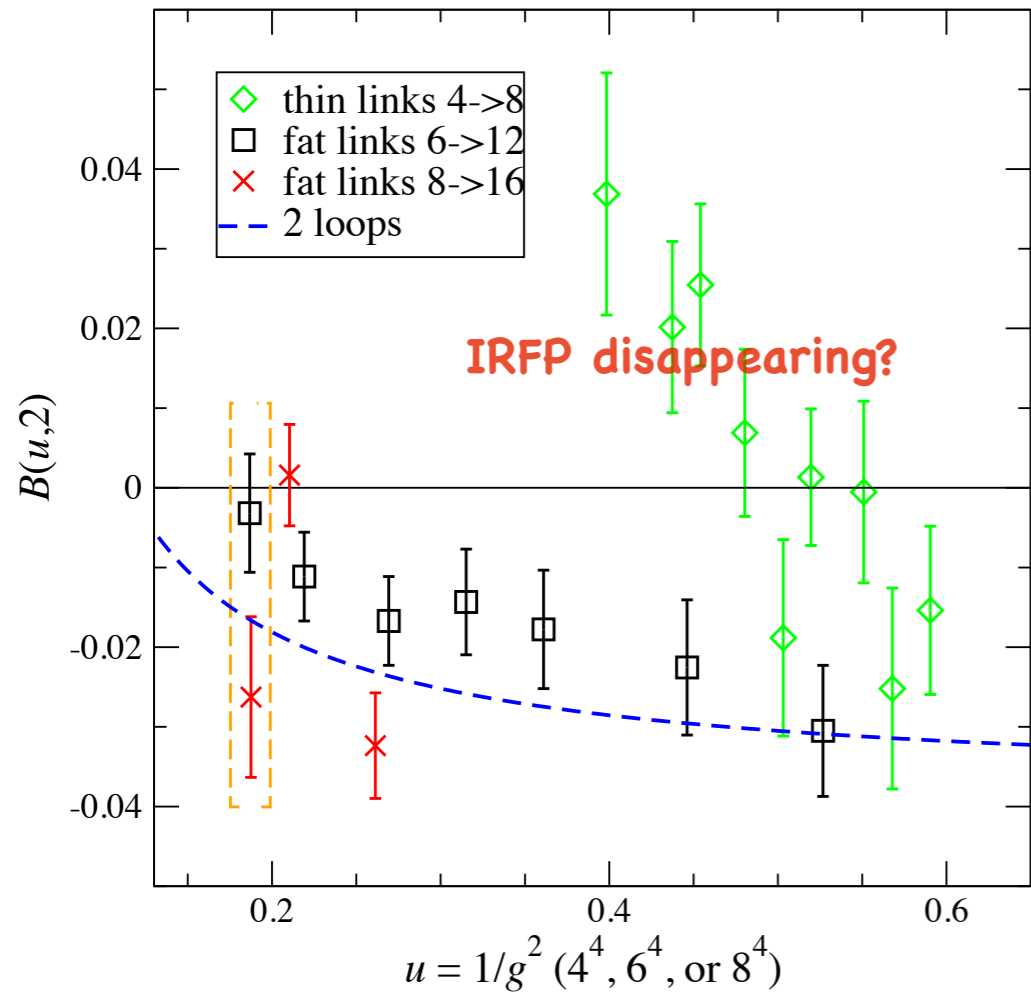
**running (walking?) gauge coupling**

DeGrand et al. find: very small  $N_f=2$  sextet beta function

IRFP zero or walking

$\chi^{SB}$  is inconsistent with IRFP slow running (walking?) what is  $\gamma(\mu)$ ?

$\gamma(\mu) < 0.45$  cannot happen with  $\chi^{SB}$

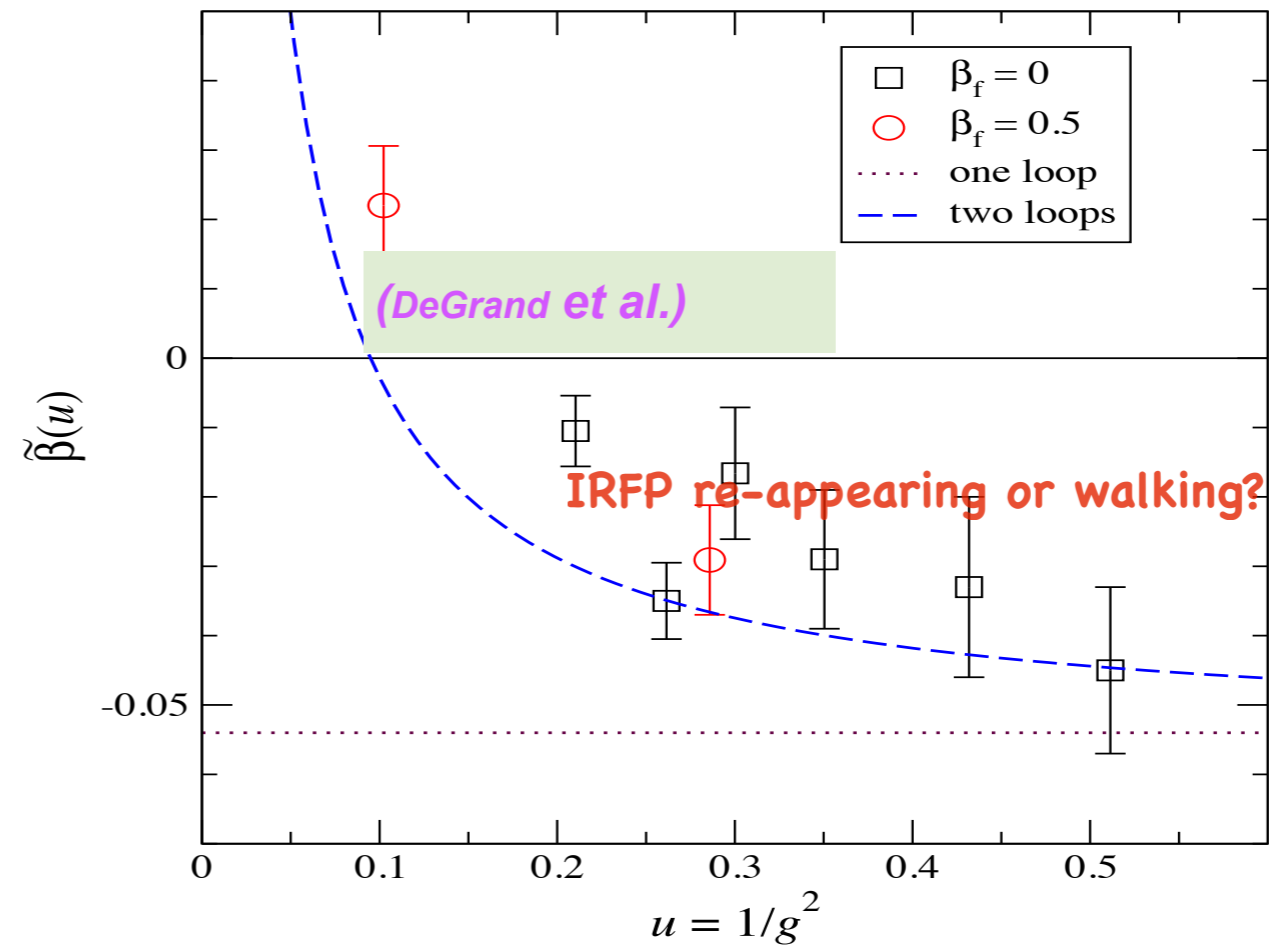
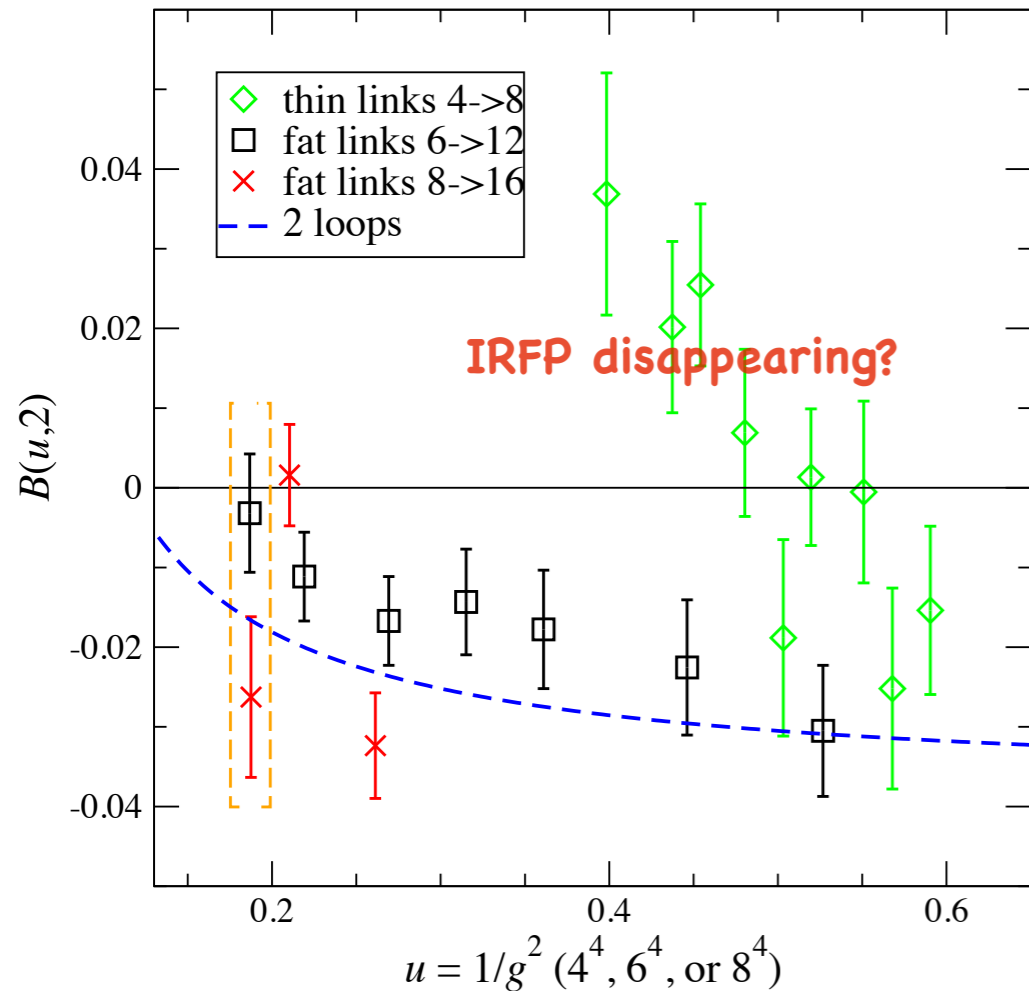


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sextet model becomes legitimate target  
for dilaton-like scenario with Higgs impostor



**dilaton as Higgs impostor?**

# The light Higgs and the dilaton near conformality

Partially Conserved Dilatation Current (PCDC)

$$m_\sigma^2 \simeq -\frac{4}{f_\sigma^2} \langle 0 | [\Theta_\mu^\mu(0)]_{NP} | 0 \rangle$$

n.p. stress-energy from gradient flow?

there are two different expectations when conformal window is approached:  $g(\mu = \Lambda) = g_c$

1. dilaton mass parametrically vanishes  $m_\sigma^2 \simeq (N_f^c - N_f) \cdot \Lambda^2 \quad \frac{m_\sigma}{f_\sigma} \rightarrow 0$

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important role of  $\frac{f_\pi}{f_\sigma}$  in electroweak phenomenology

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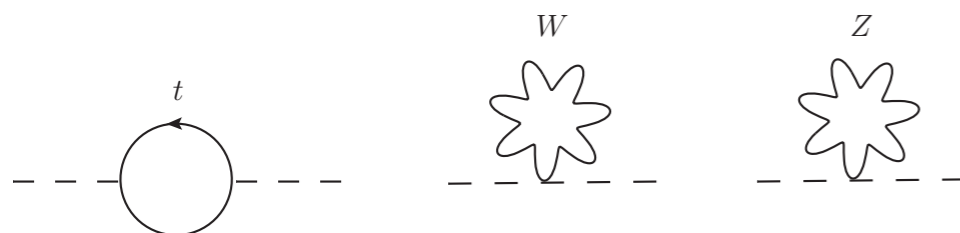
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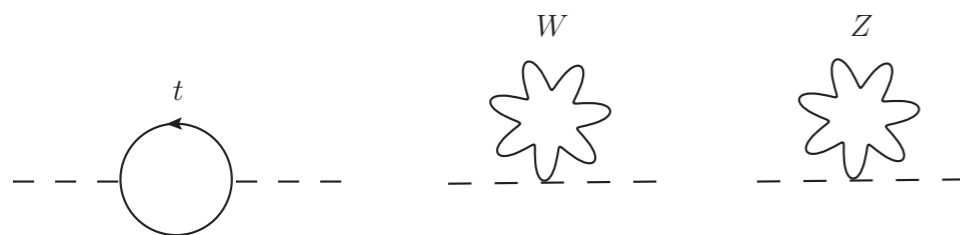
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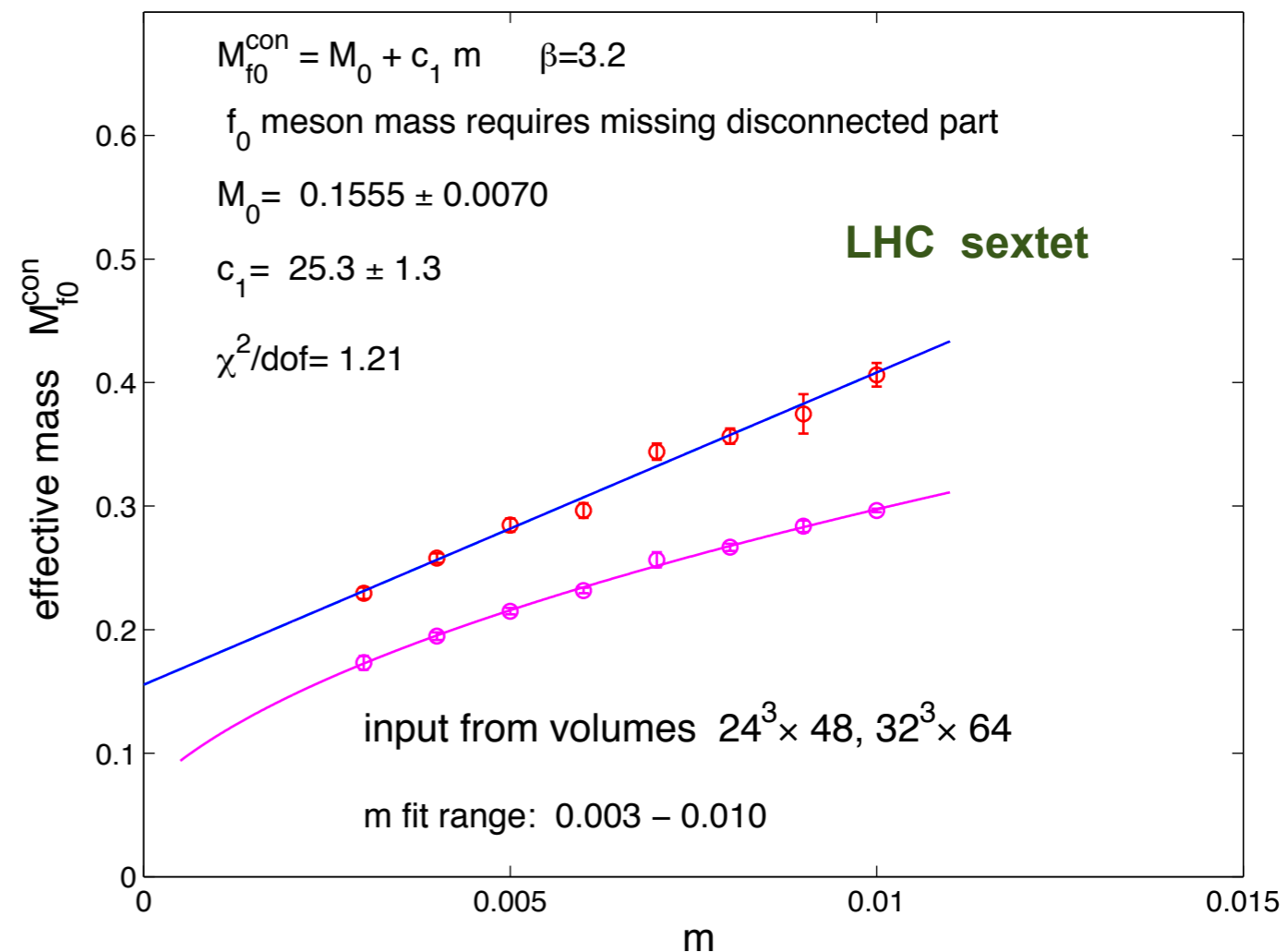
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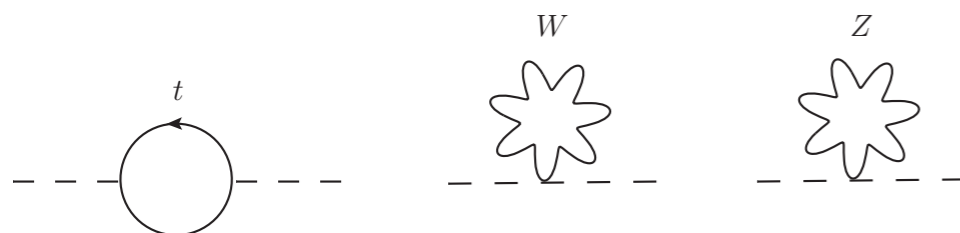
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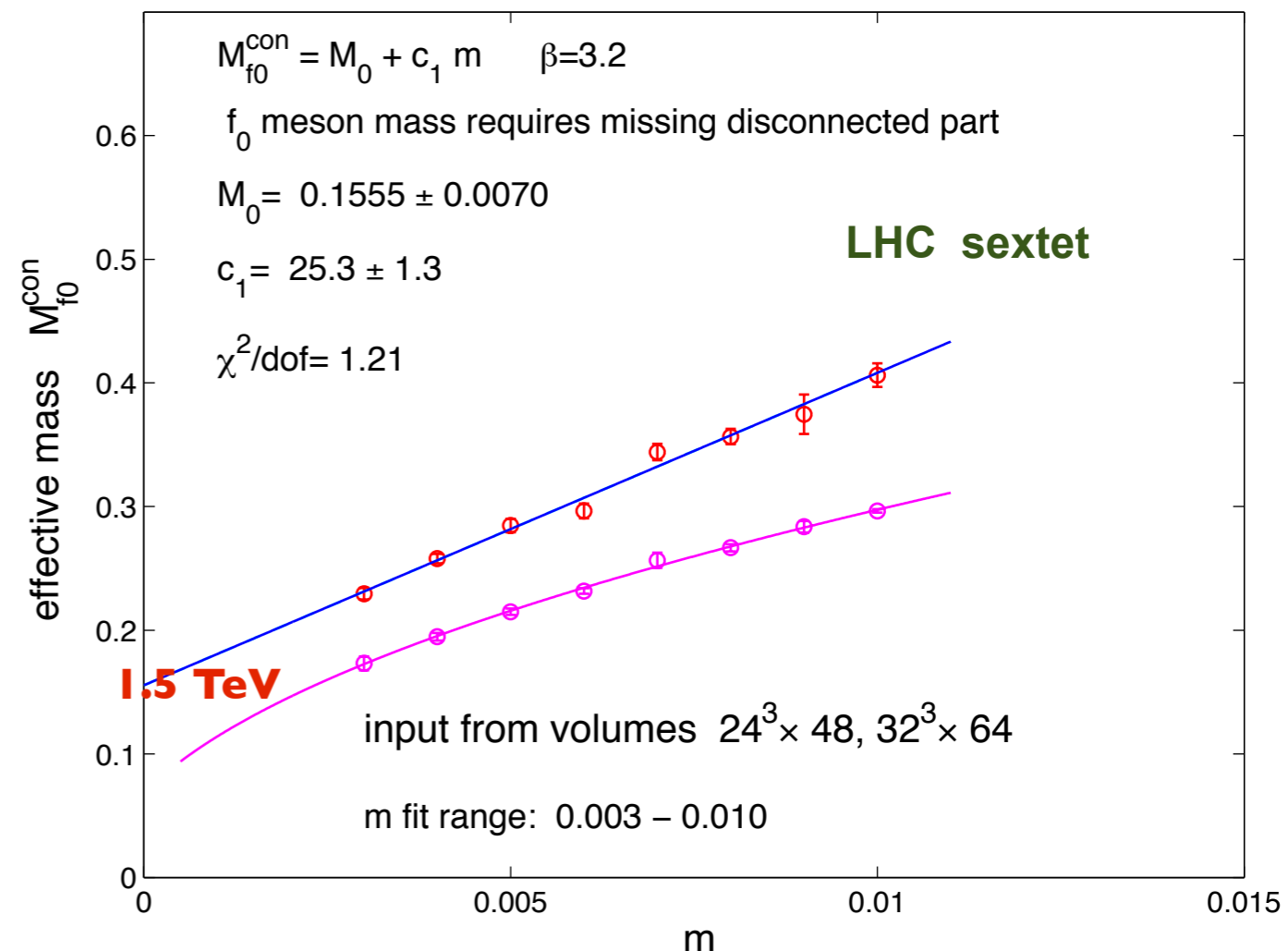
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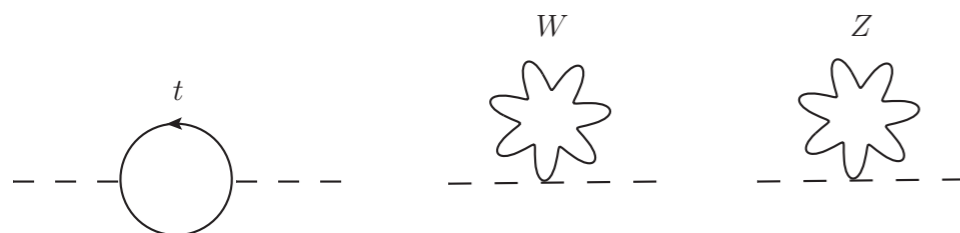
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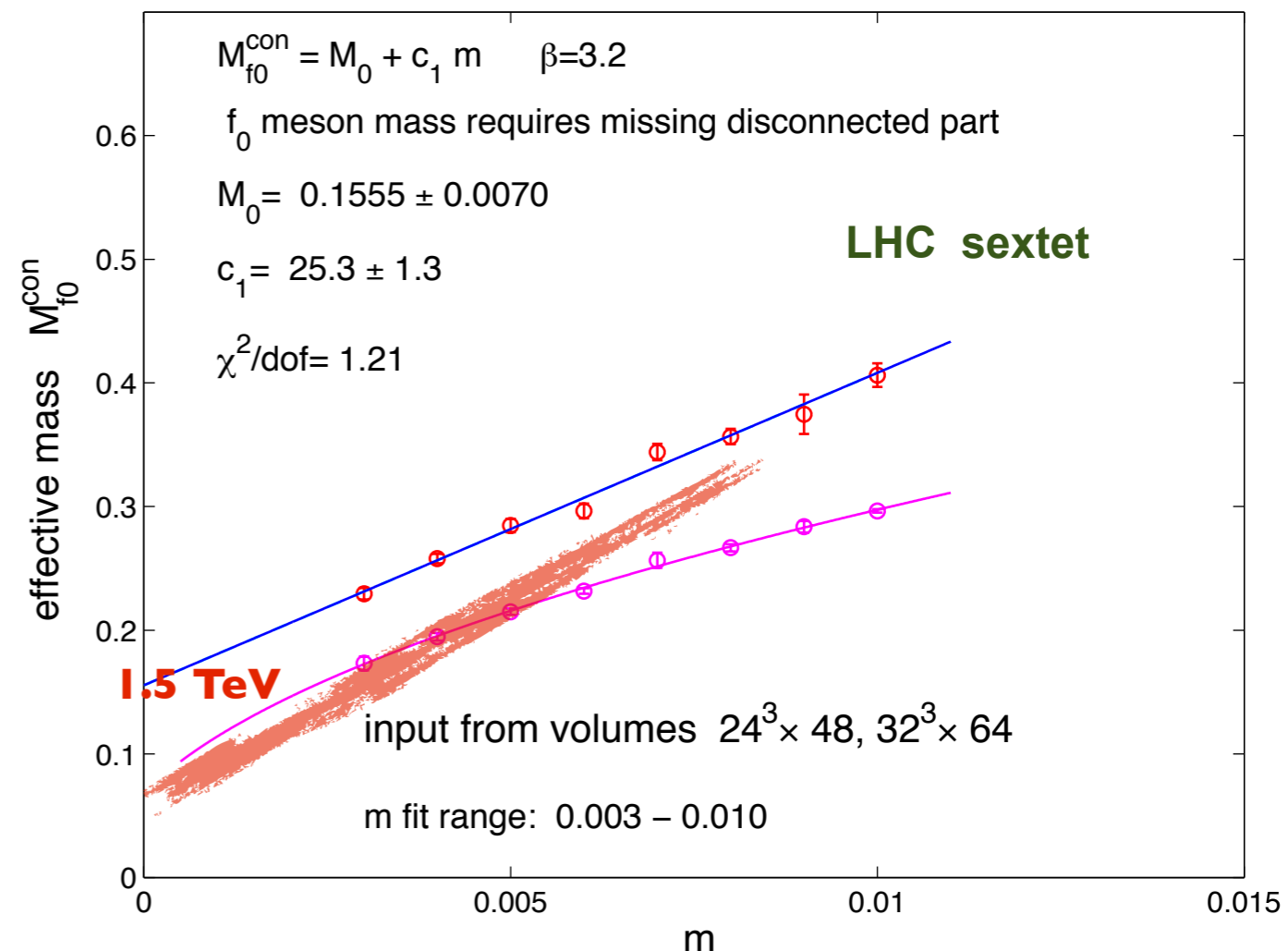
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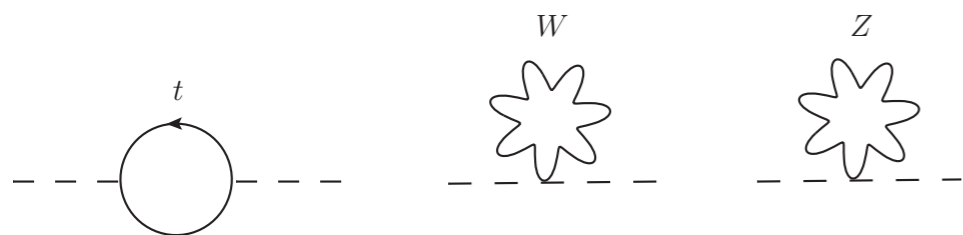
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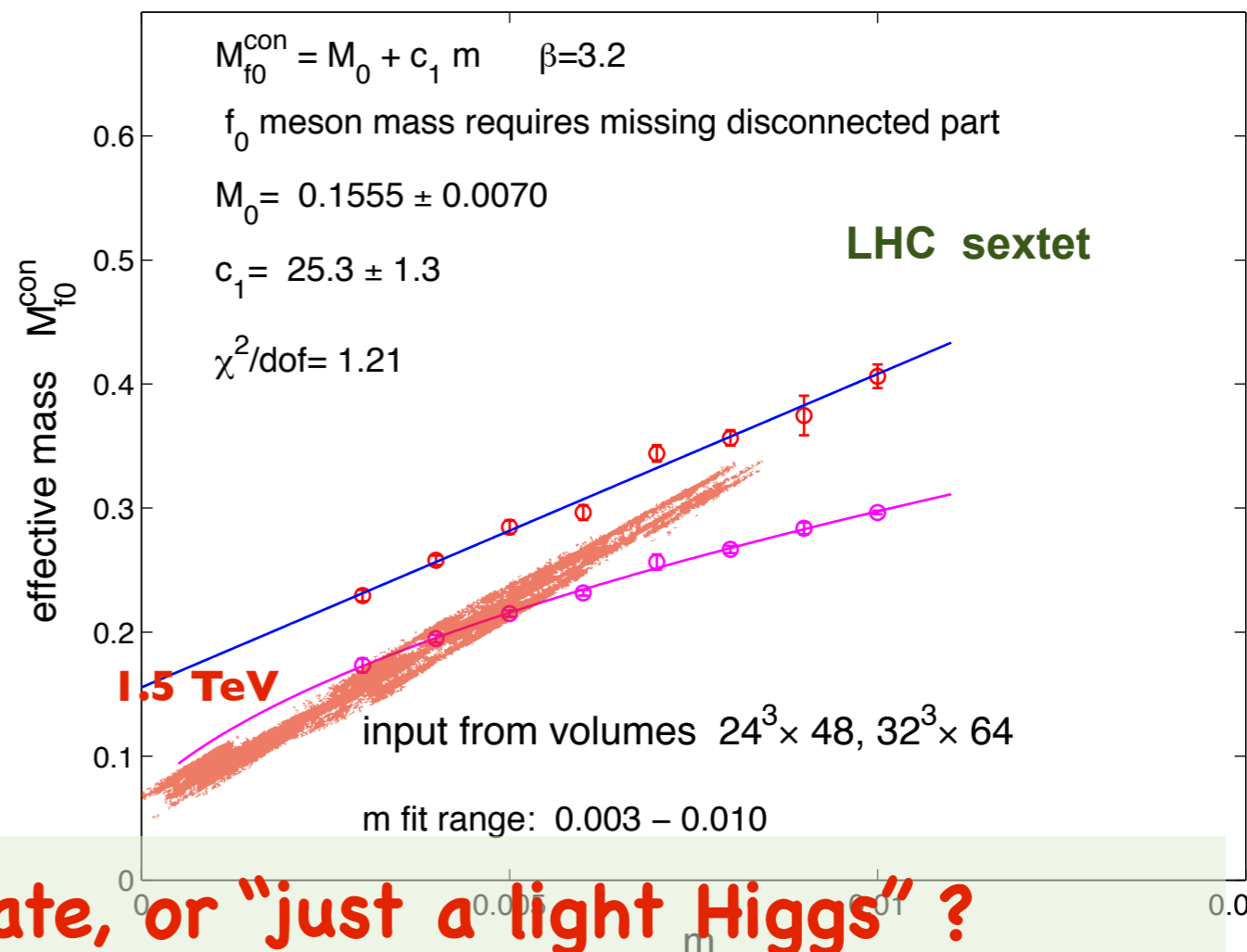
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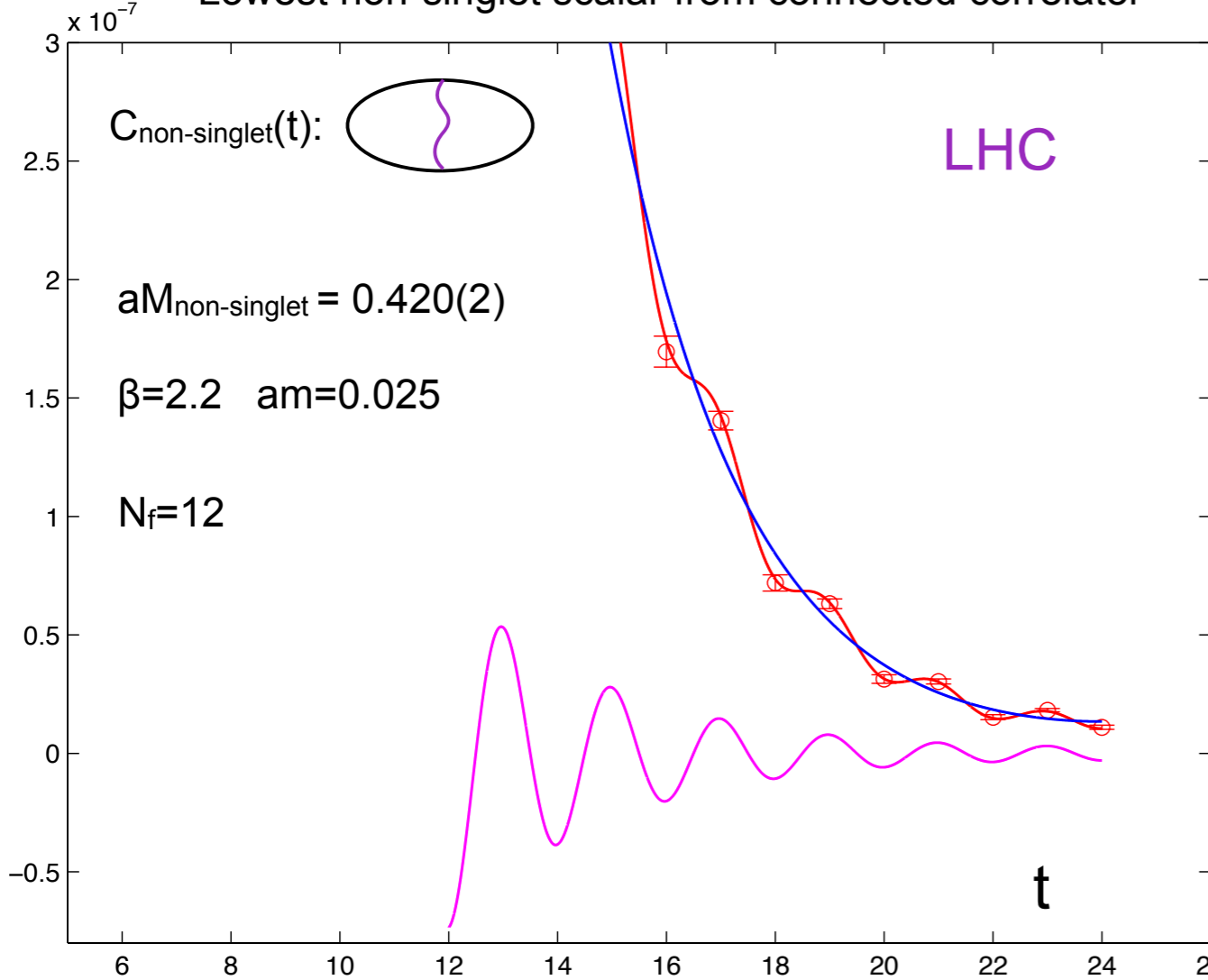
is this SCGT dilaton-like scalar state, or "just a light Higgs"?



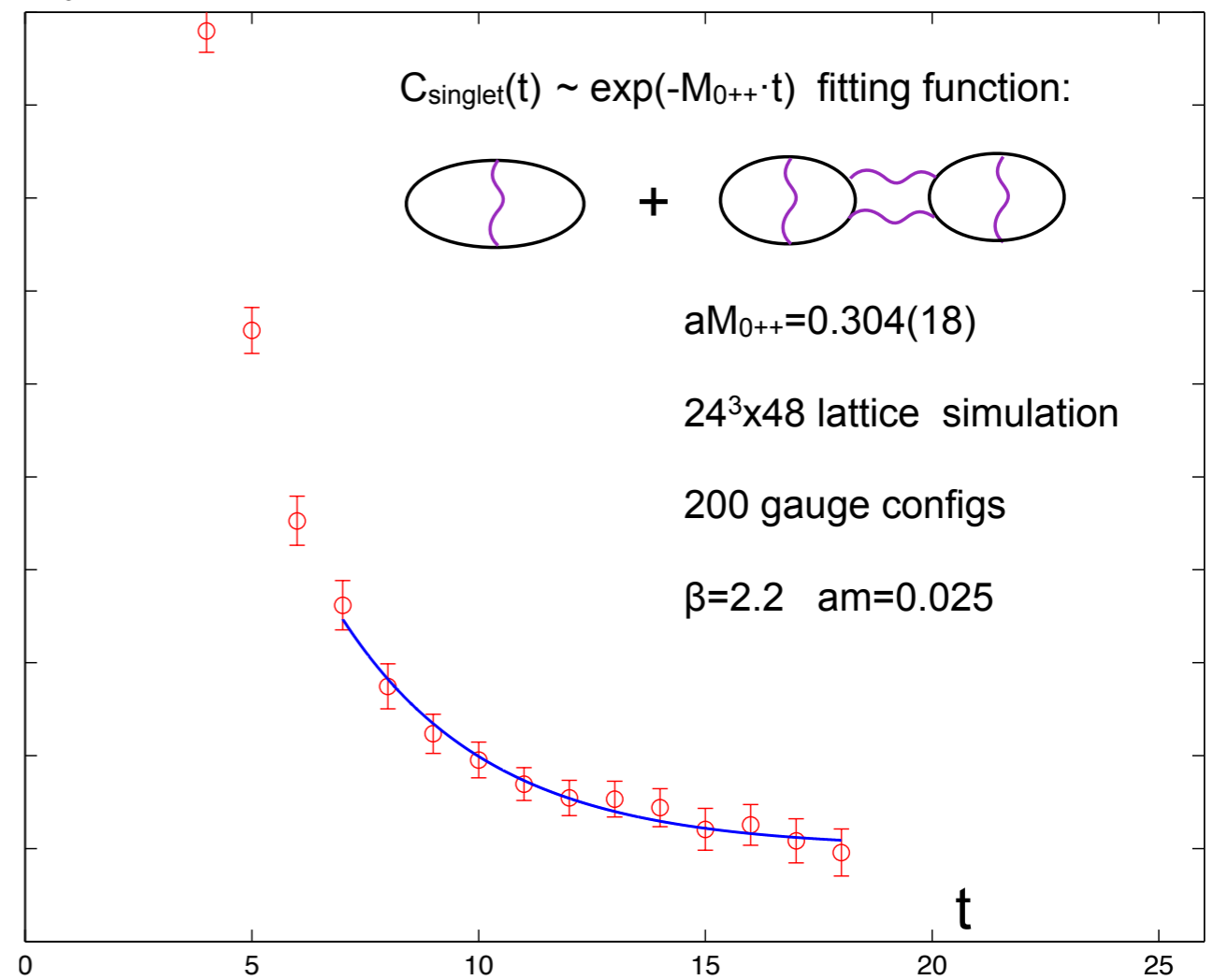
# The light Higgs and the dilaton near conformality

## proof of life:

Lowest non-singlet scalar from connected correlator



$N_f=12$  Lowest  $0^{++}$  scalar state from singlet correlator



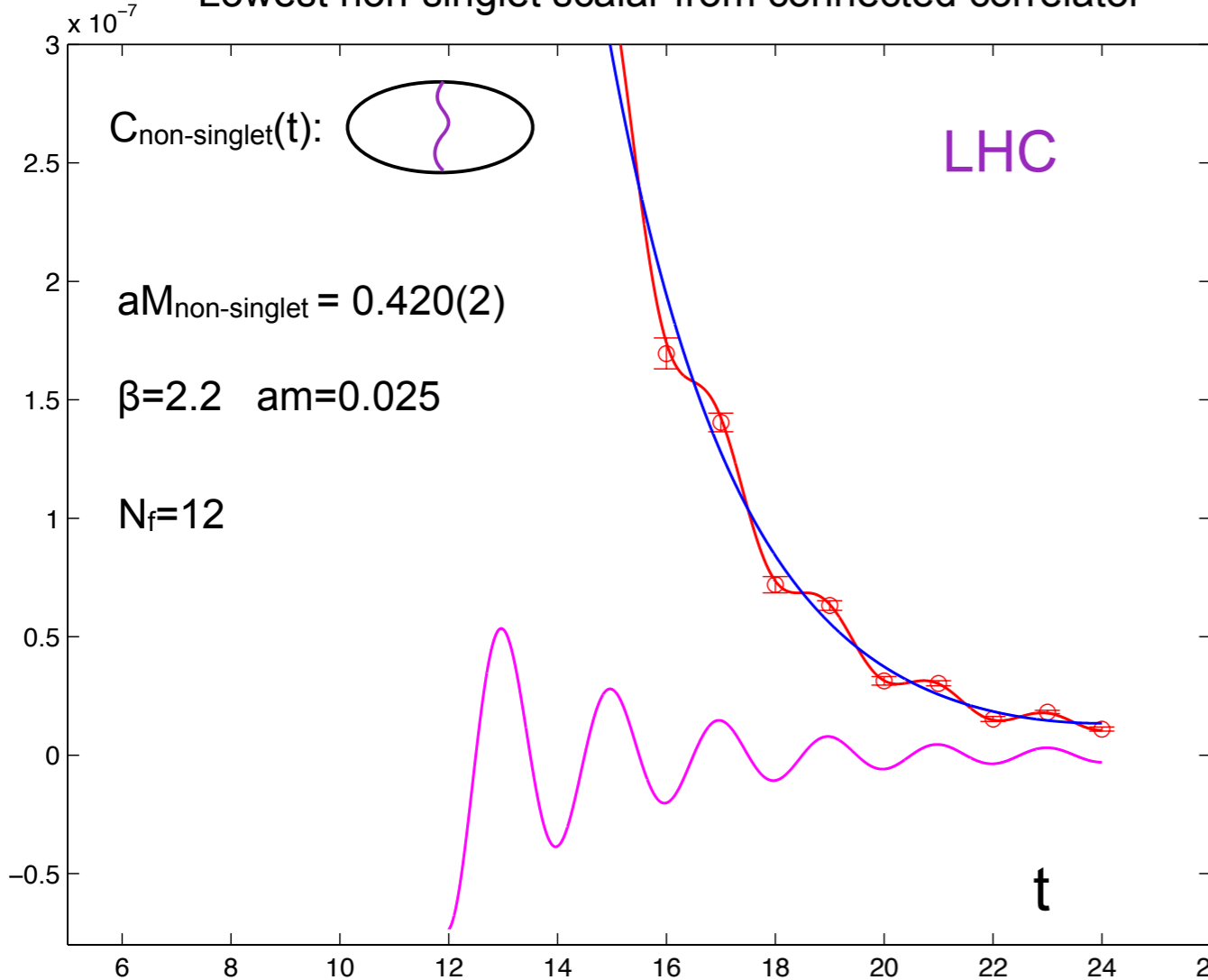
$$C(t) = \sum_n \left[ A_n e^{-m_n(\Gamma_S \otimes \Gamma_T)t} + (-1)^t B_n e^{-m_n(\gamma_4 \gamma_5 \Gamma_S \otimes \gamma_4 \gamma_5 \Gamma_T)t} \right]$$

staggered correlator

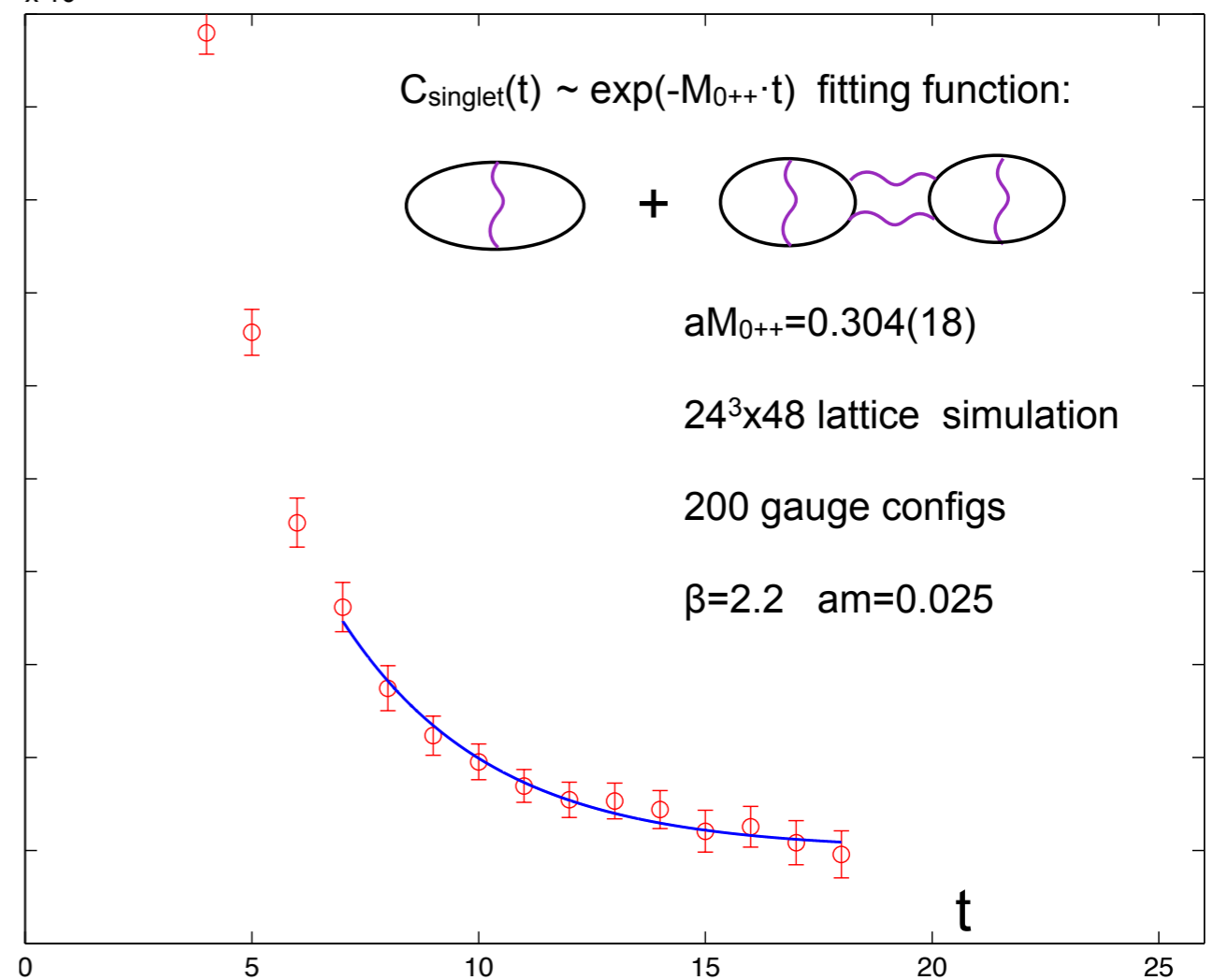
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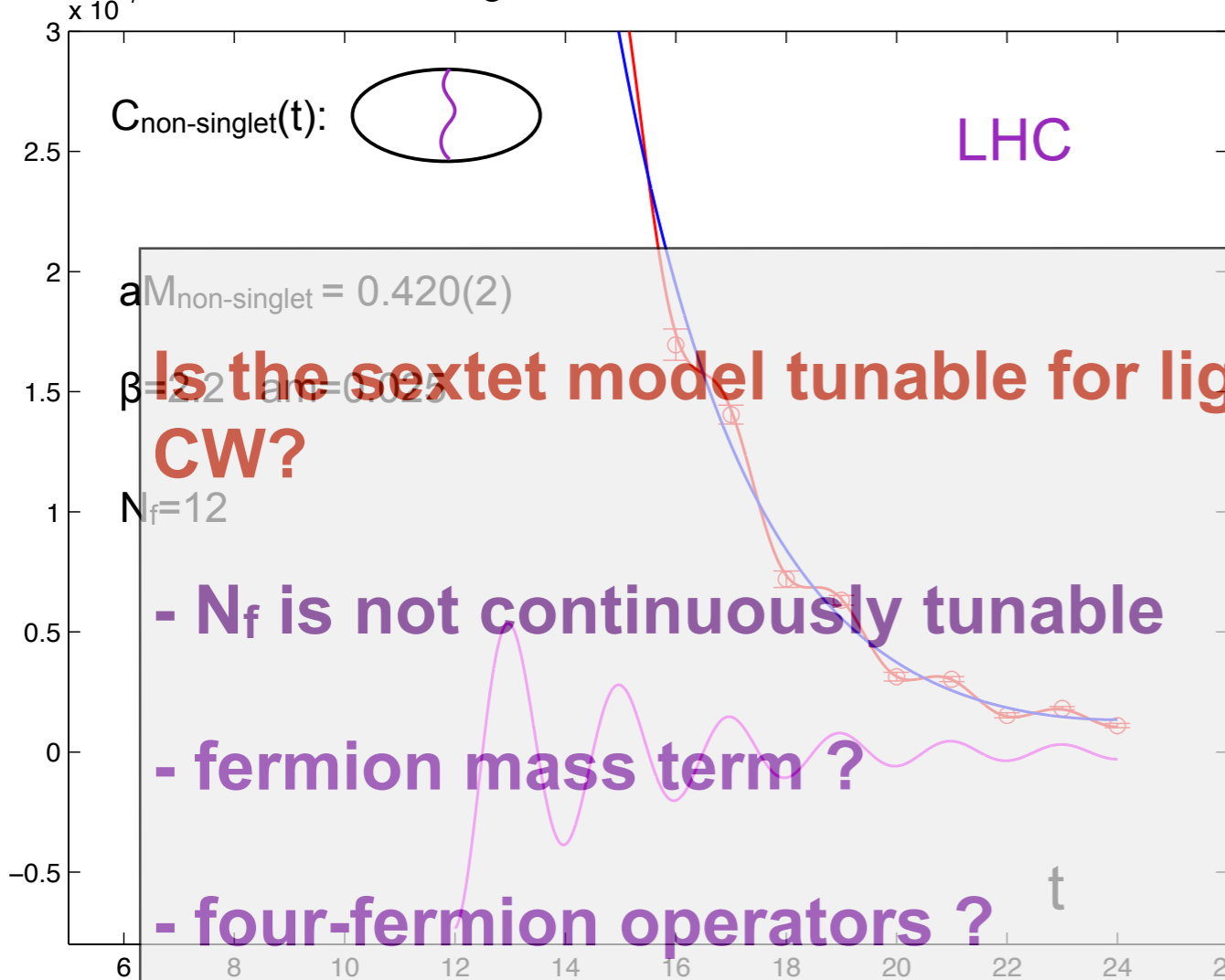
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**similar analysis is being done in sextet model with  $N_f=2$**

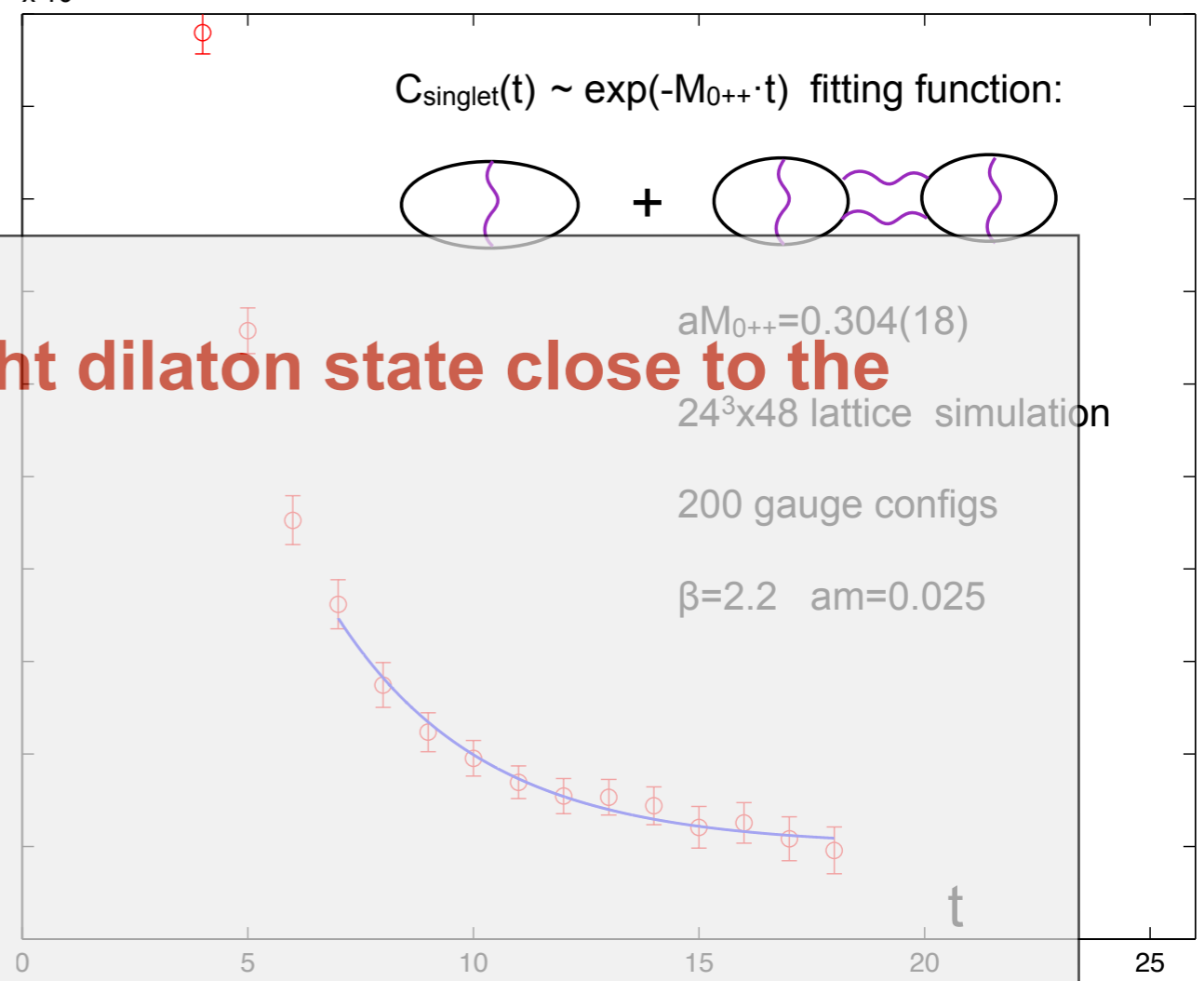
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Is the sextet model tunable for light dilaton state close to the CW?

-  $N_f$  is not continuously tunable

- fermion mass term ?

- four-fermion operators ?

$$C(t) = \sum_n \left[ A_n e^{-m_n(\Gamma_S \otimes \Gamma_T)t} + (-1)^t B_n e^{-m_n(\gamma_4 \gamma_5 \Gamma_S \otimes \gamma_4 \gamma_5 \Gamma_T)t} \right] \quad \text{staggered correlator}$$

similar analysis is being done in sextet model with  $N_f = 2$

# mass tuning, like partially gauged (conformal) Technicolor?

Sannino, Dietrich, Luty, ...

freeze-out from  $N_f=3$  to  $N_f=2$  inside conformal window

mass-dependent beta-function of sextet model

MOM scheme (Yoshino and Hagiwara)

- position of IRFP inside CW not tunable with  $m$
- plateau length tunable, its position is not (in or out)

$$\mu \frac{d}{d\mu} \alpha(\mu) = -\alpha(\mu)^2 \pi^{-1} [\beta_0 + \beta_1 \cdot \alpha(\mu)]$$

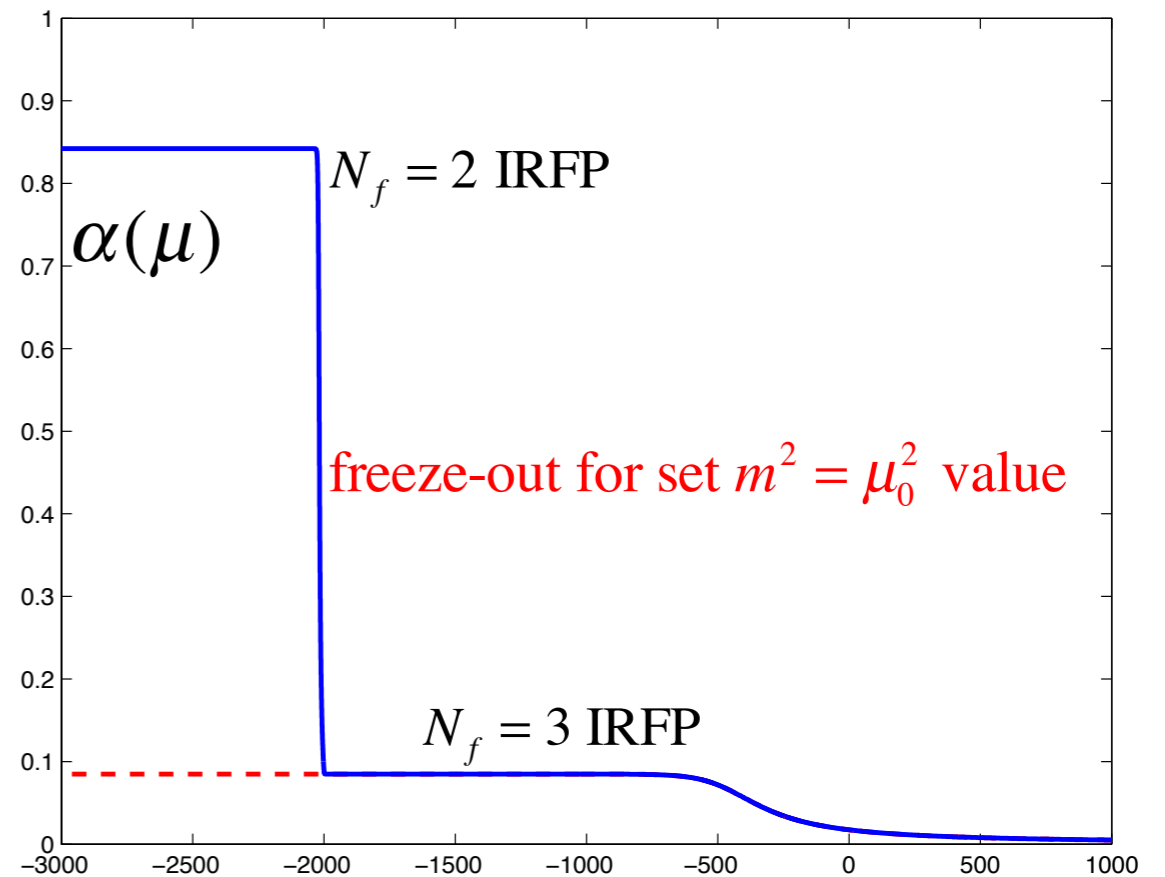
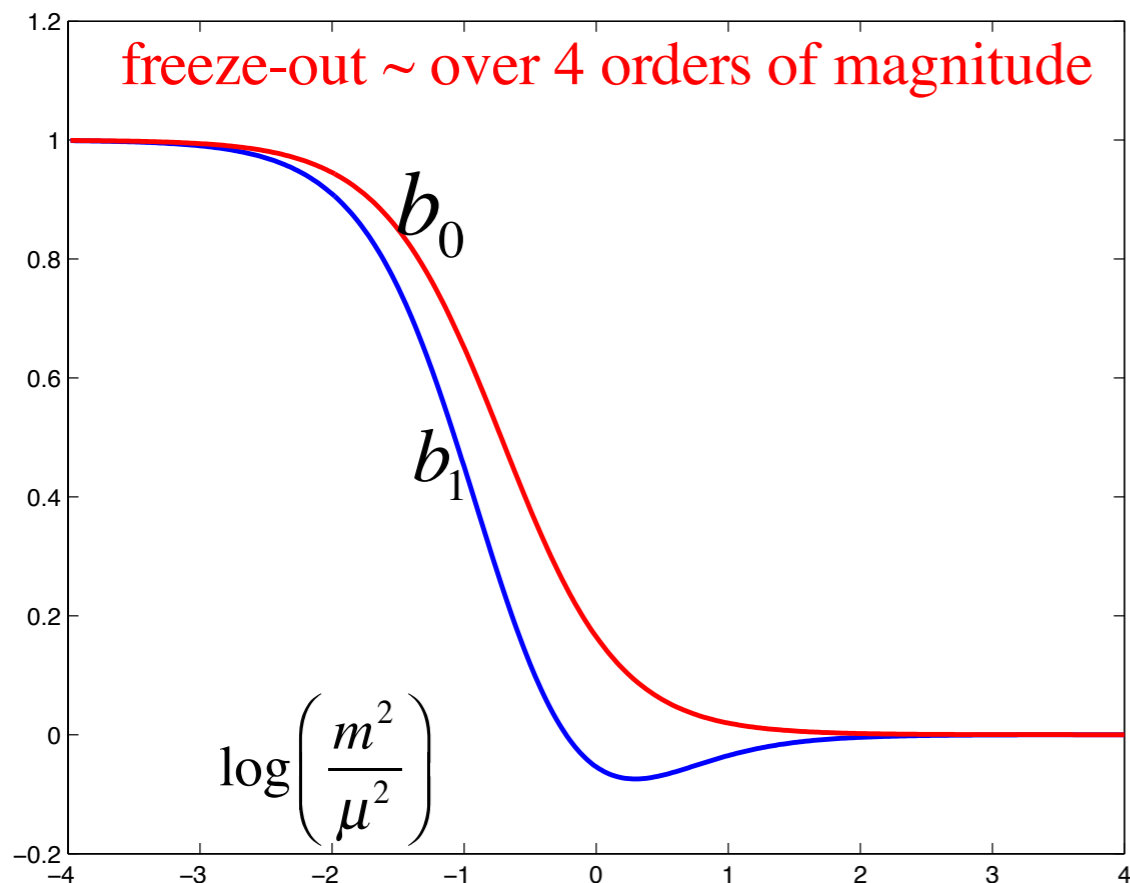
$$\beta_0 = \frac{11}{6} C_A - \frac{2}{3} T_F \sum_f b_0(x), \quad x = \frac{m_f^2}{\mu^2}$$

$$b_0(x) = 1 - 6x + \frac{12x^2}{\sqrt{1+4x}} \ln \frac{\sqrt{1+4x} + 1}{\sqrt{1+4x} - 1}$$

$$\beta_1(x) = \frac{17}{12} C_A^2 - \left( \frac{5}{6} C_A + \frac{1}{2} C_F \right) \cdot T_F \sum_f b_1(x)$$

$$b_1(x) = \frac{-0.45577x + 0.26995}{x^2 + 2.1742x + 0.26995}$$

sextet parameters:  $C_A = 3$ ;  $C_F = \frac{10}{3}$ ;  $T_F = \frac{5}{2}$



four-fermion operator

tunable deformation of IRFP?

large-N double trace limit

Witten, Rastelli, Vecchi ...

which scenario is realized will depend on the scaling dimension  $\Delta$  of  $\langle \bar{\psi}\psi \rangle$  and on other intrinsic properties of IRFP

for any model choice things are set

$$\mathcal{L}_{\text{CFT}} + \frac{f}{2} \mathcal{O}_{ij}^\dagger \mathcal{O}^{ij}$$

$$\mathcal{L}_{\text{def}} = f(q_R^\dagger q_L)(q_L^\dagger q_R) \quad \langle e^{i \int \mathcal{L}_{\text{def}}} \rangle_{\text{CFT}}$$

$$f(\Lambda) - \frac{v}{2\Delta - d} f^2(\Lambda) (\Lambda^{2\Delta-d} - \Lambda'^{2\Delta-d}) = f(\Lambda')$$

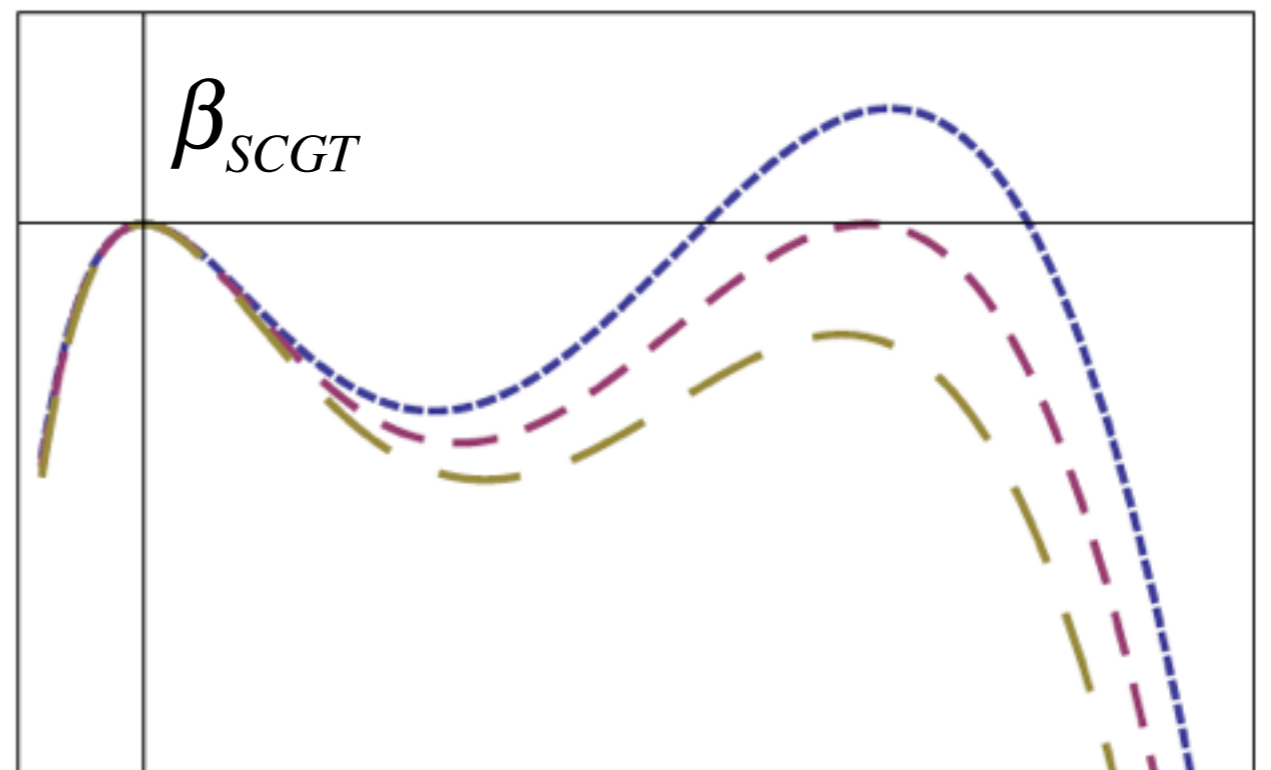
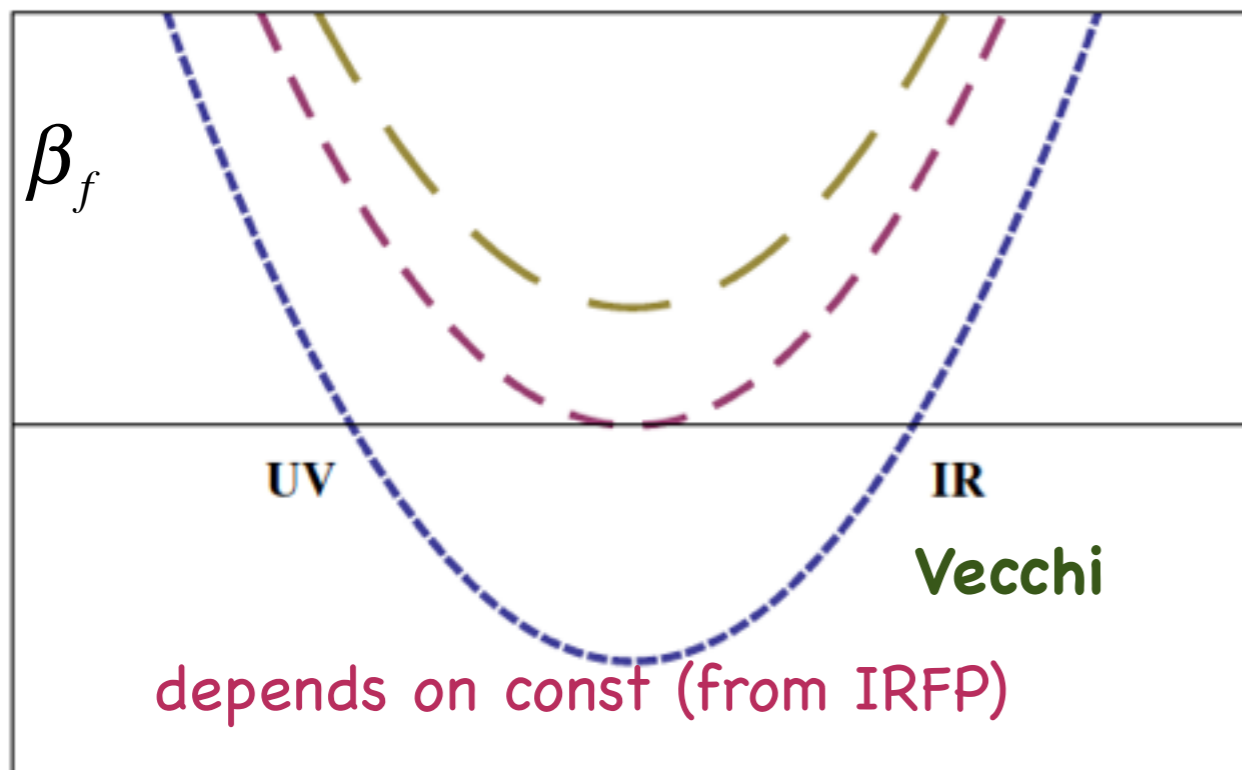
$$\langle \mathcal{O}(x)\mathcal{O}(y) \rangle \equiv \frac{v}{|x-y|^{2\Delta}} \frac{\Gamma(d/2)}{2\pi^{d/2}}$$

$$\Lambda \frac{d\bar{f}}{d\Lambda} = v\bar{f}^2 + (2\Delta - d)\bar{f}$$

+ const (from conform IRFP)

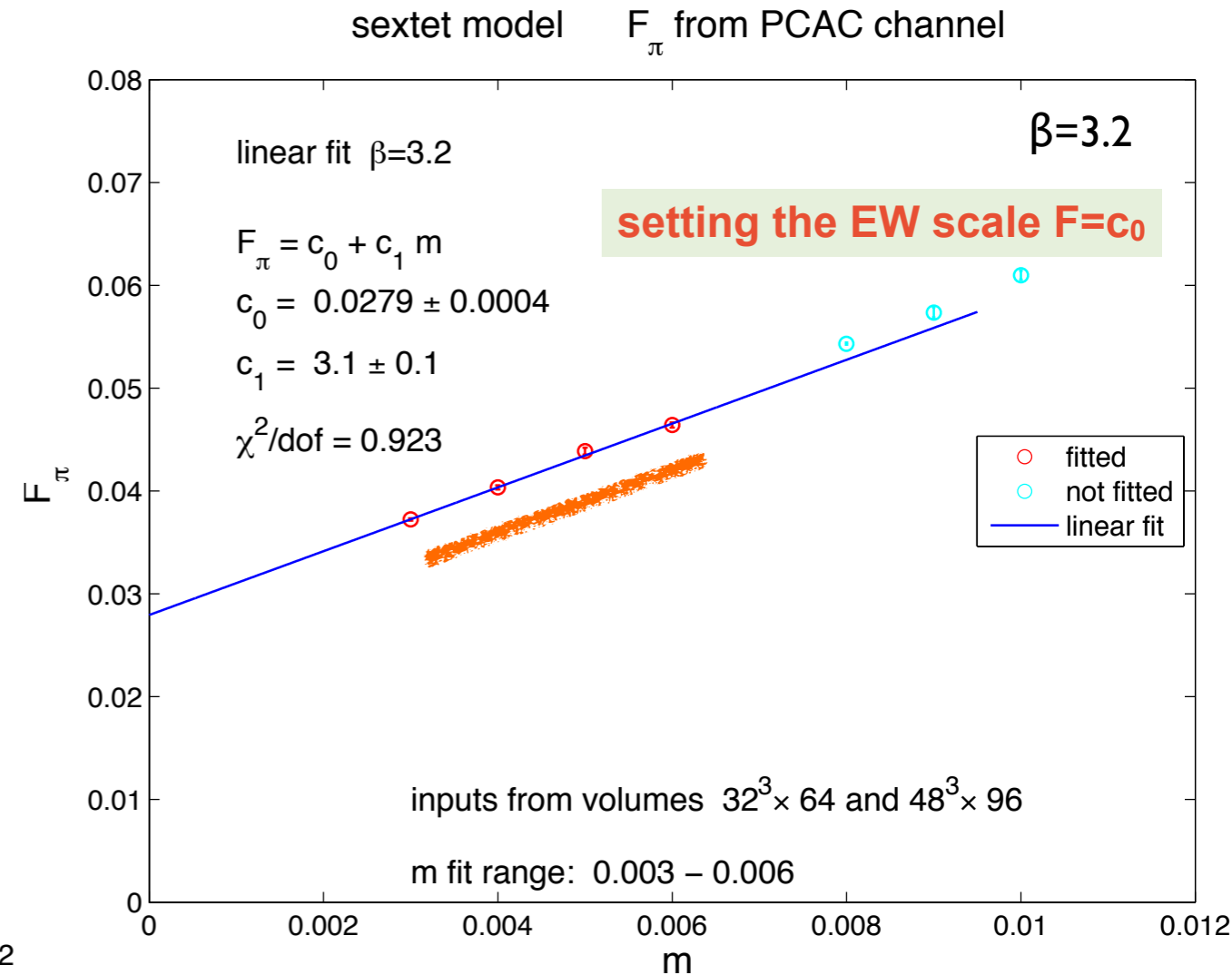
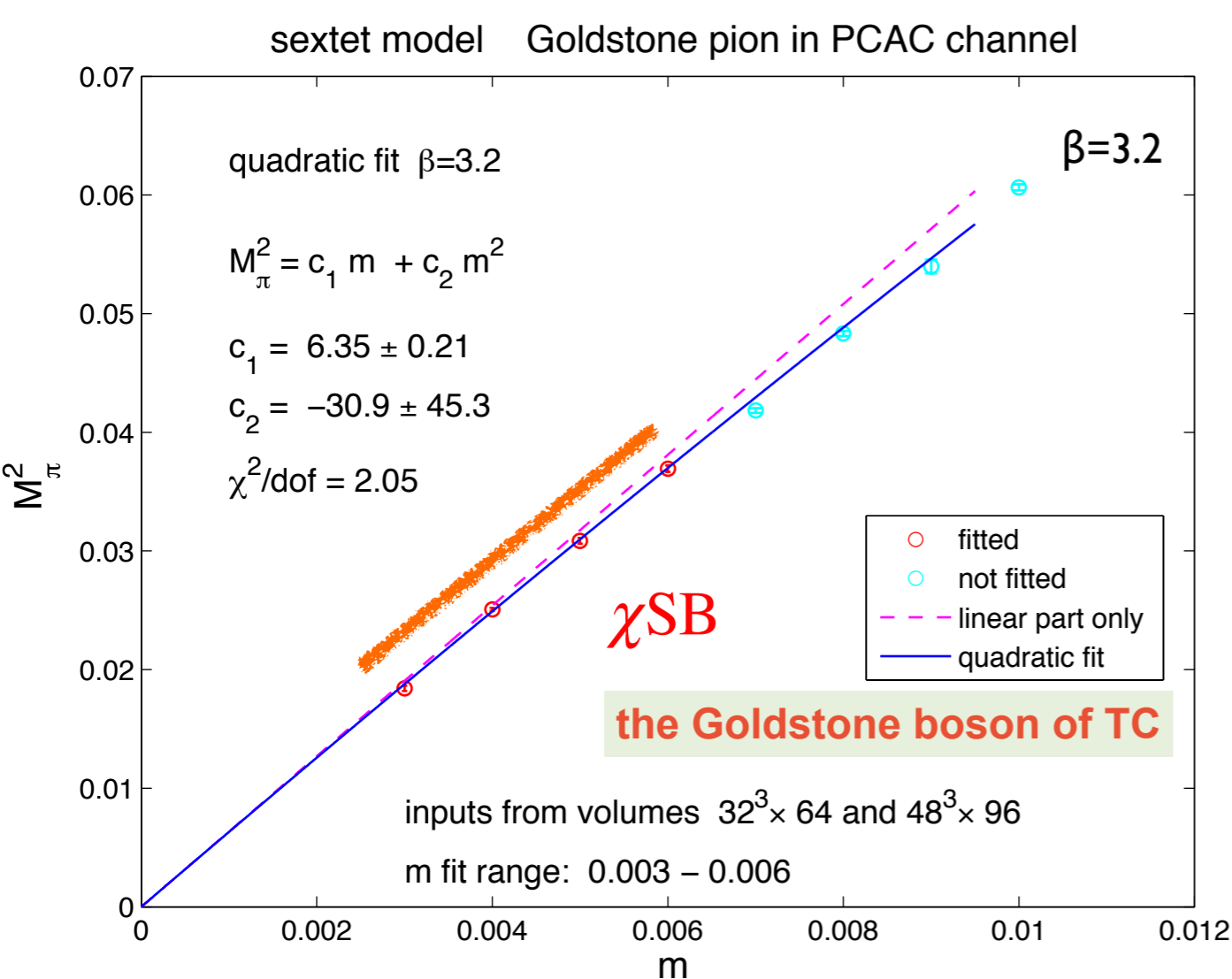
f flows to fixed points

not continuously tunable, no guarantee



**spectroscopy and confining force**

# Nf=2 SU(3) sextet chiral fits of $M_\pi$ and $F_\pi$



$m=0.003-0.006$  range close to chiral log regime?

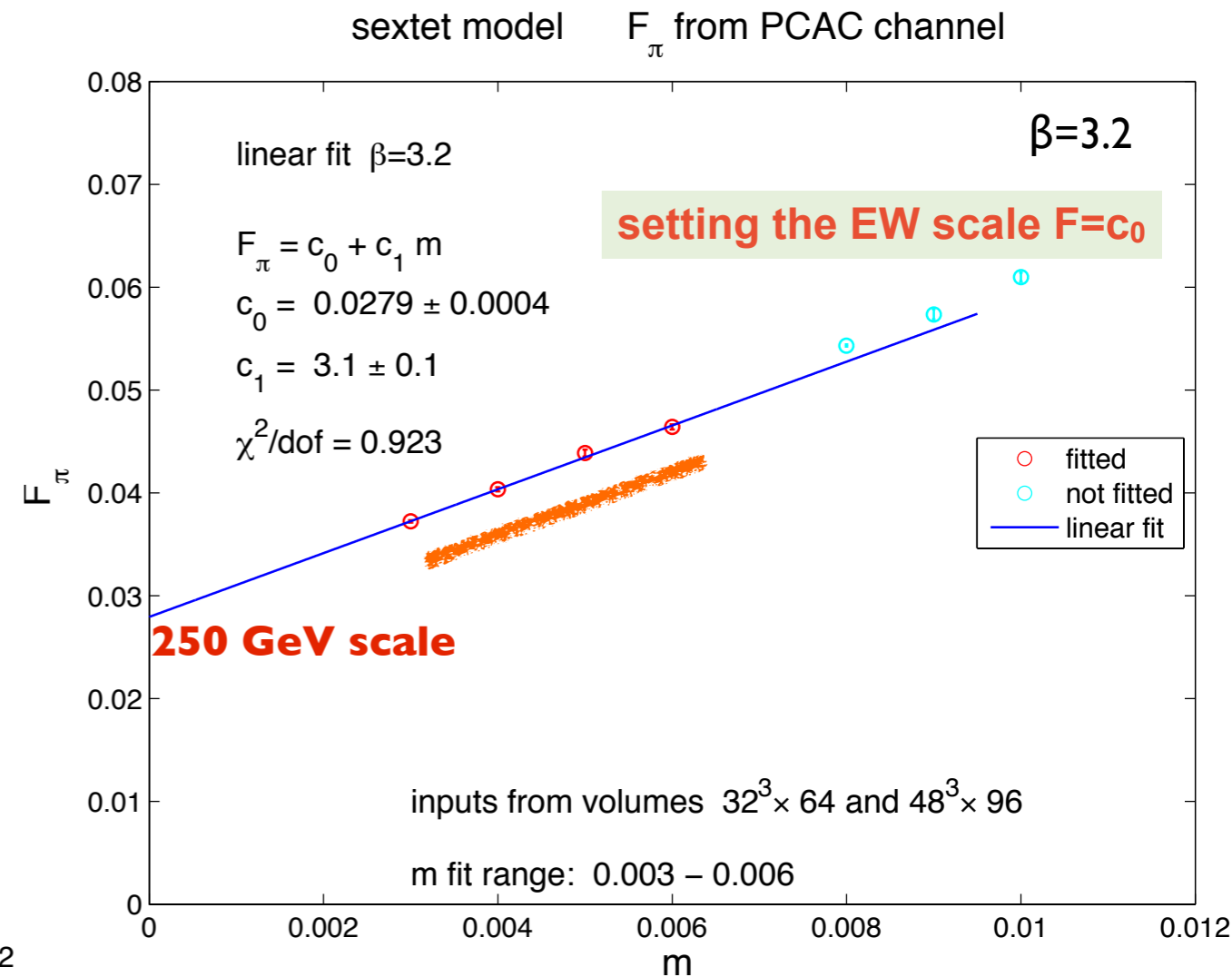
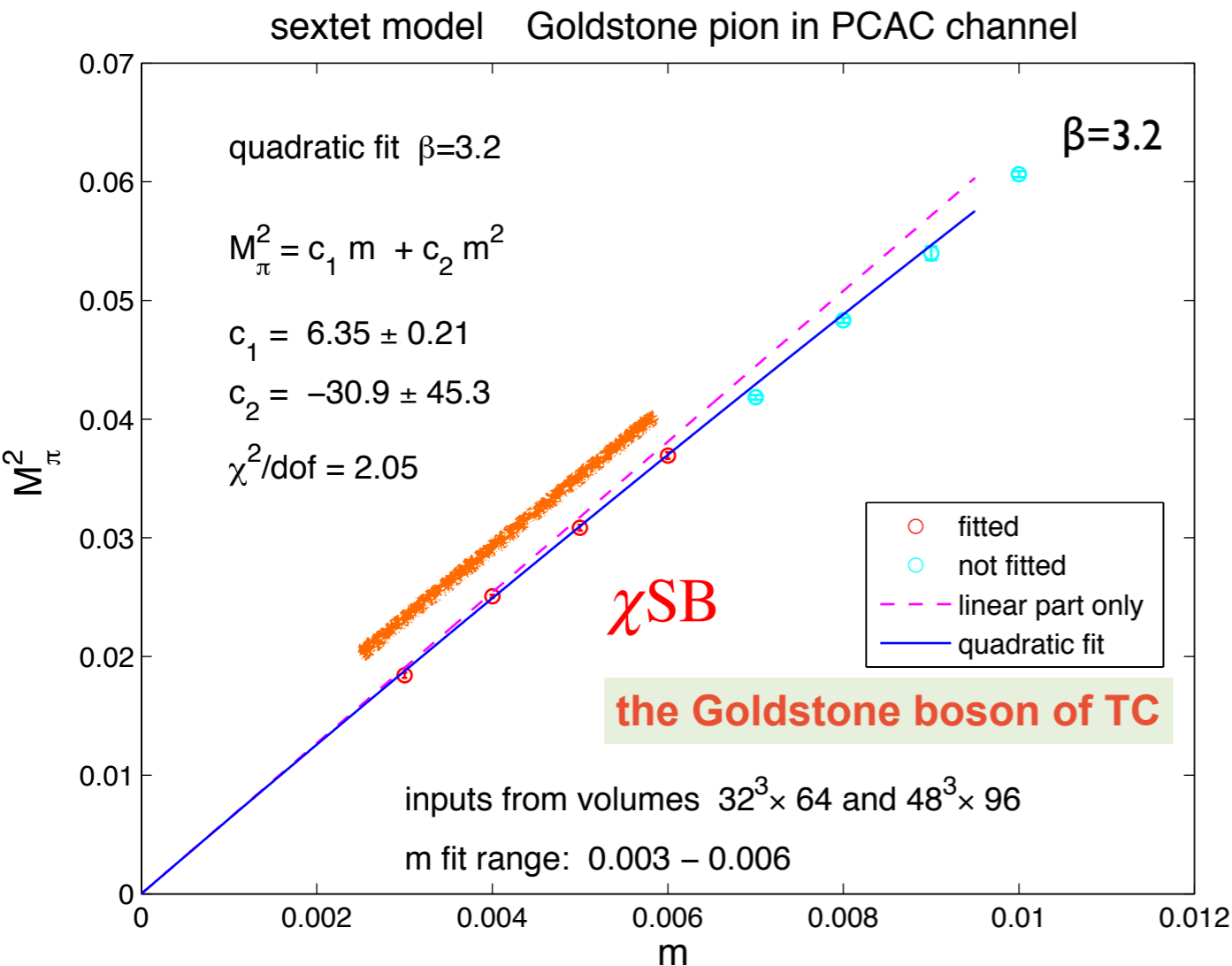
Nf=2 helps, more QCD-like

log detection will require even more precise data

consistency with partially quenched staggered chiral perturbation theory?

To better understand GMOR relation

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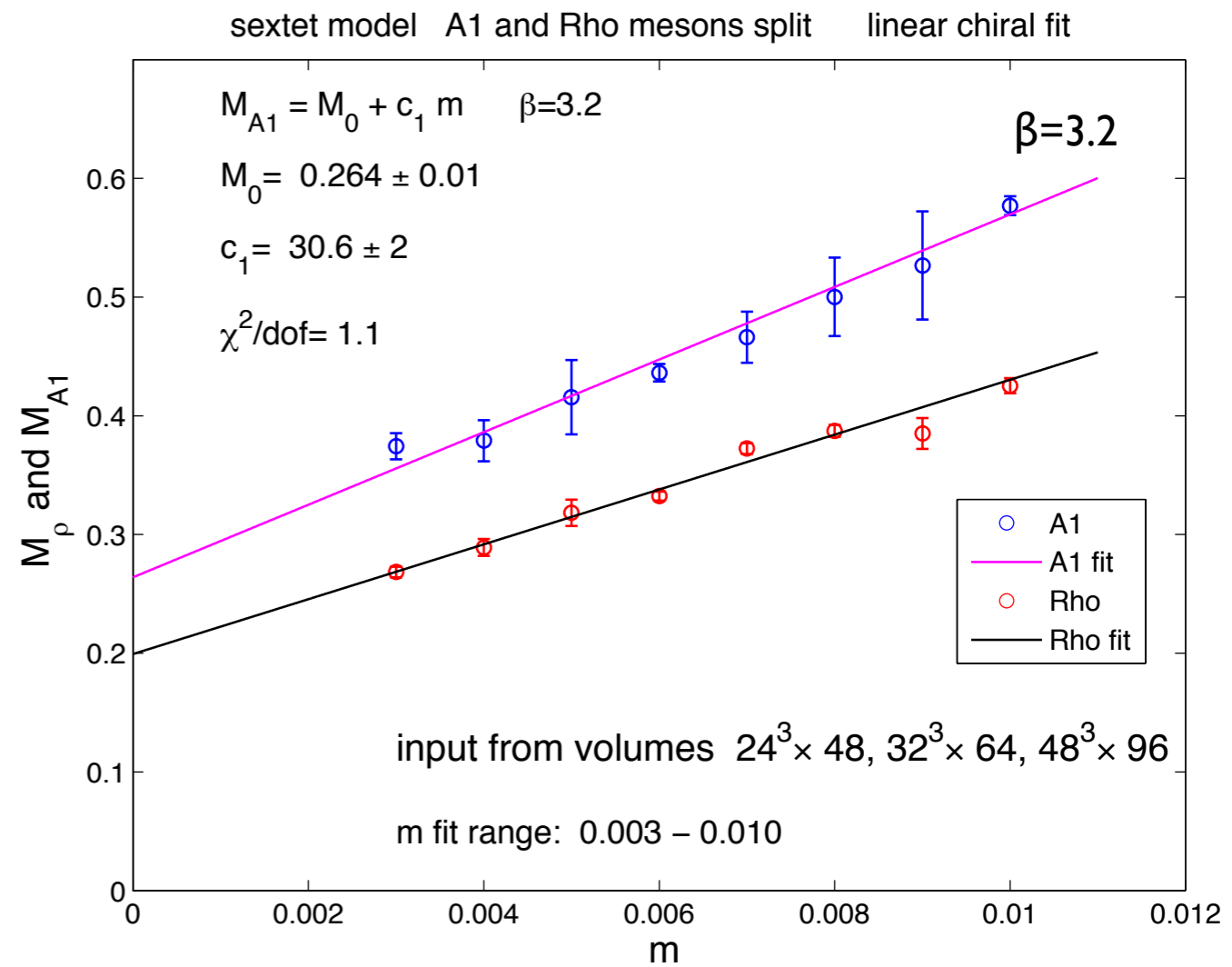
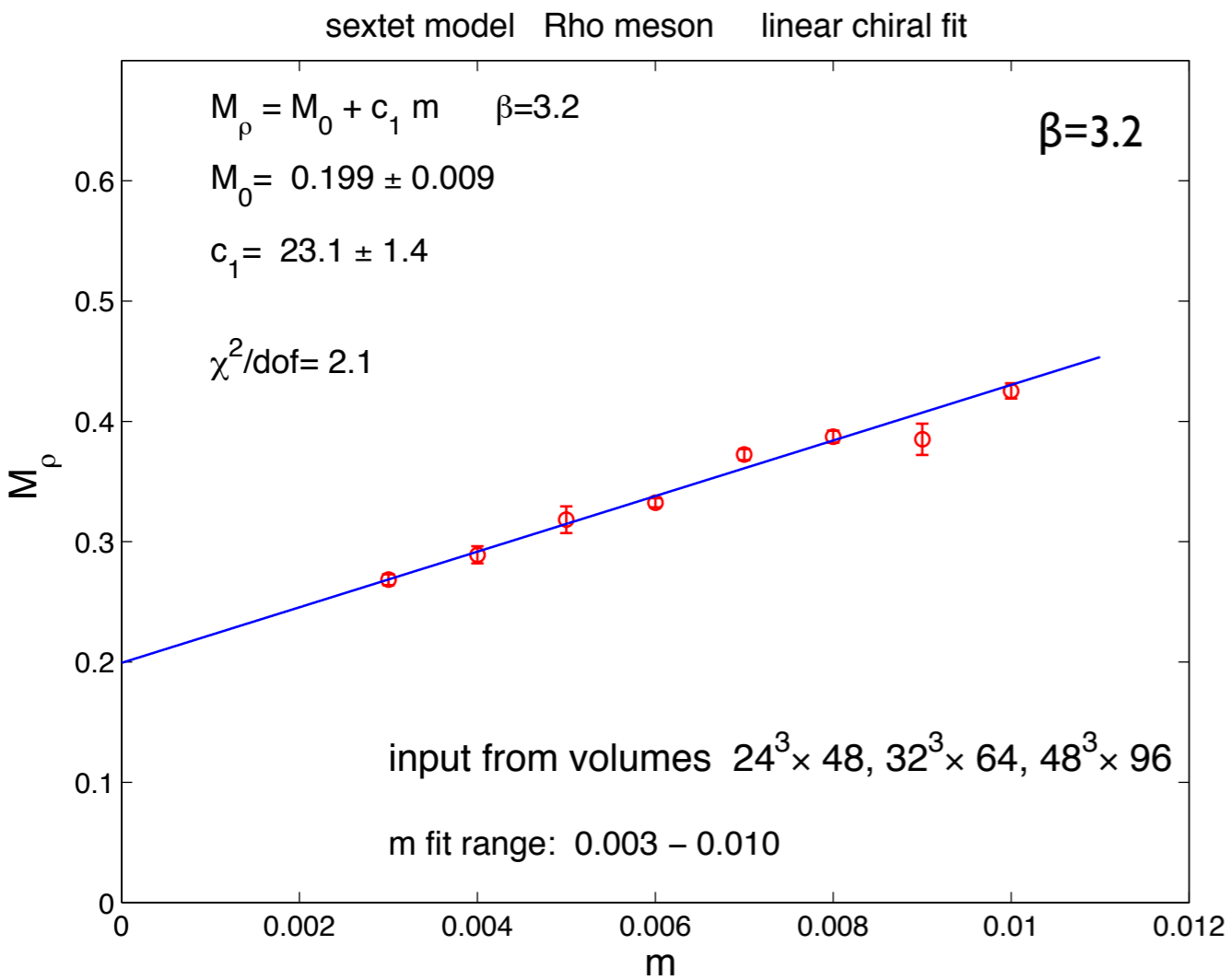
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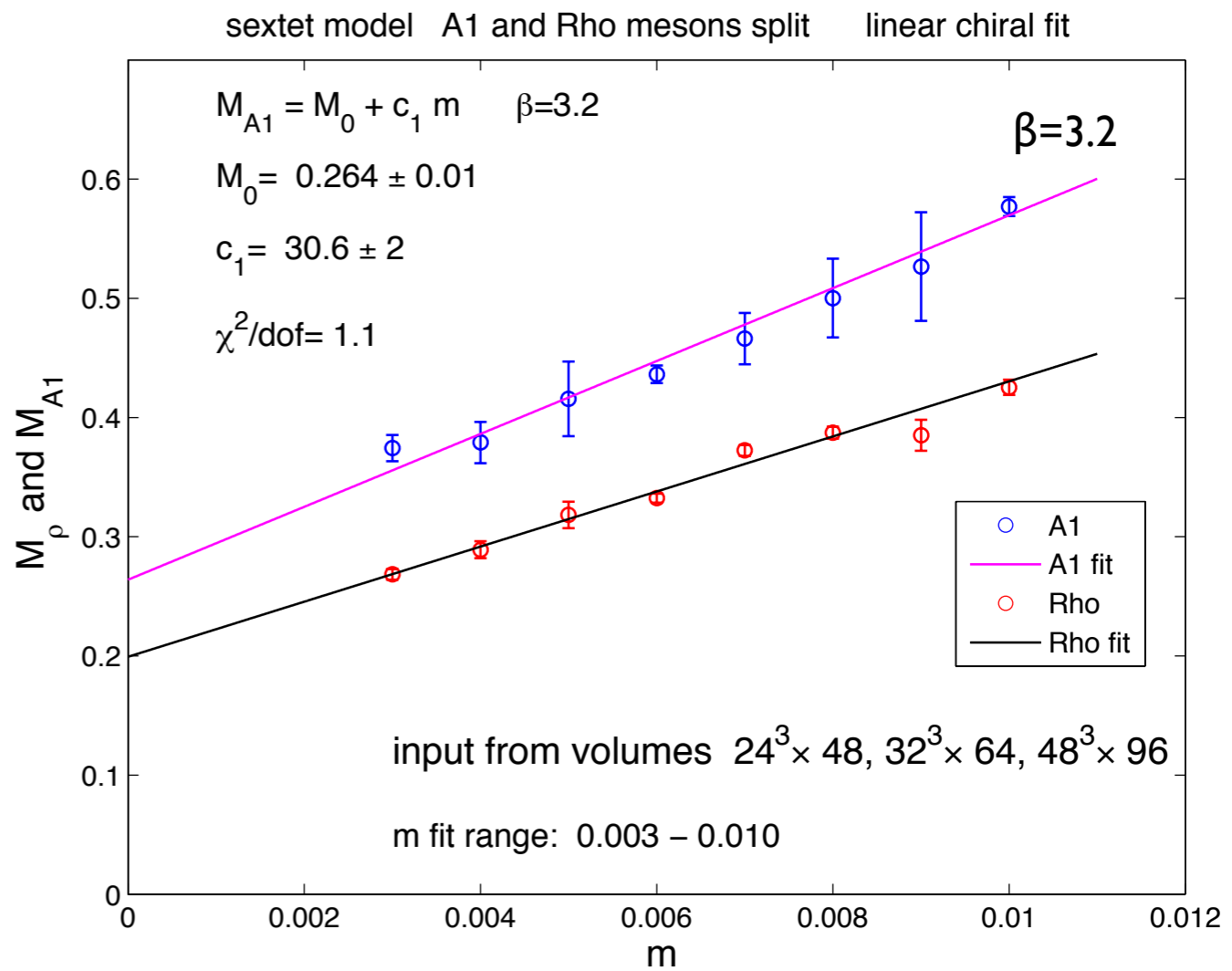
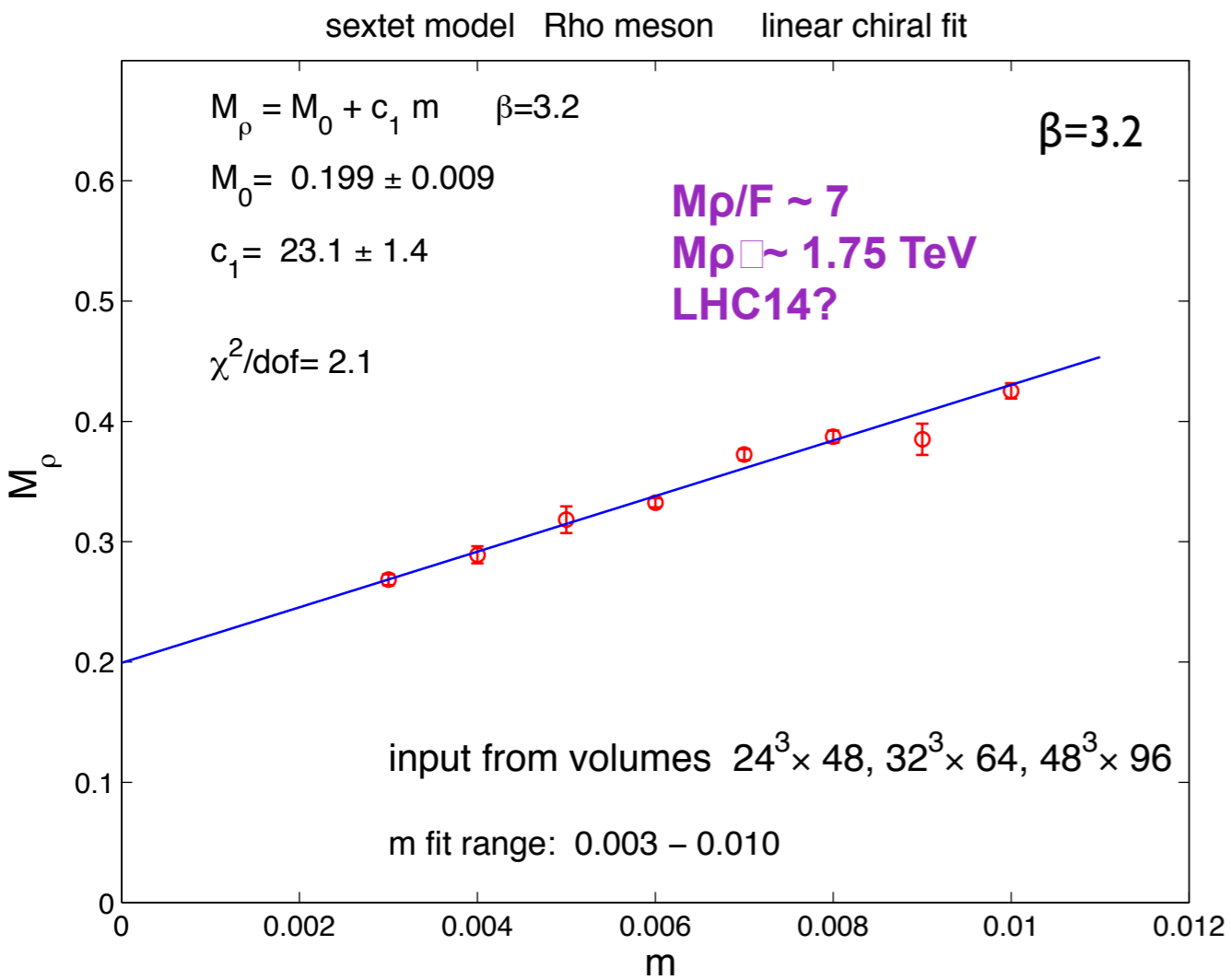
To better understand GMOR relation



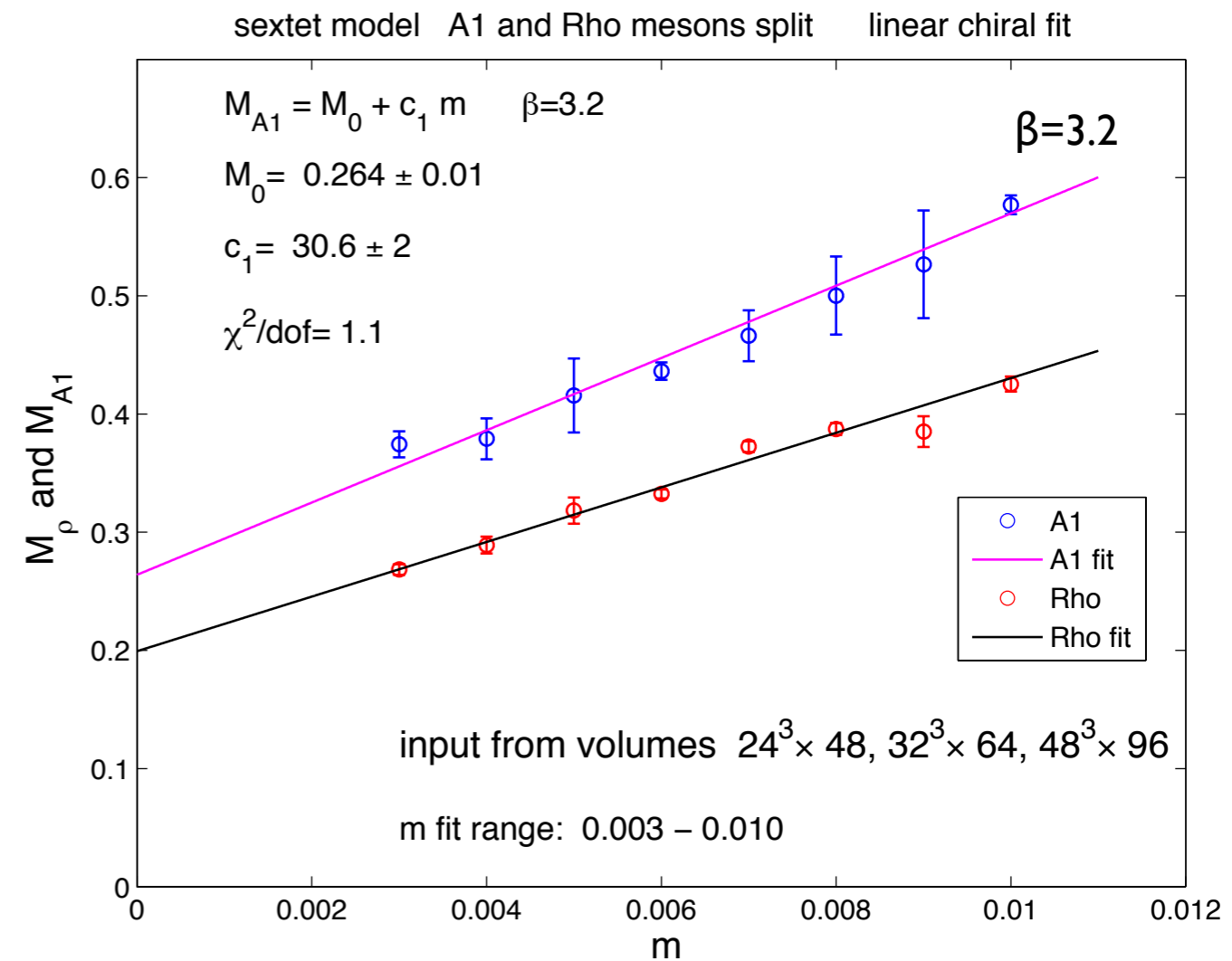
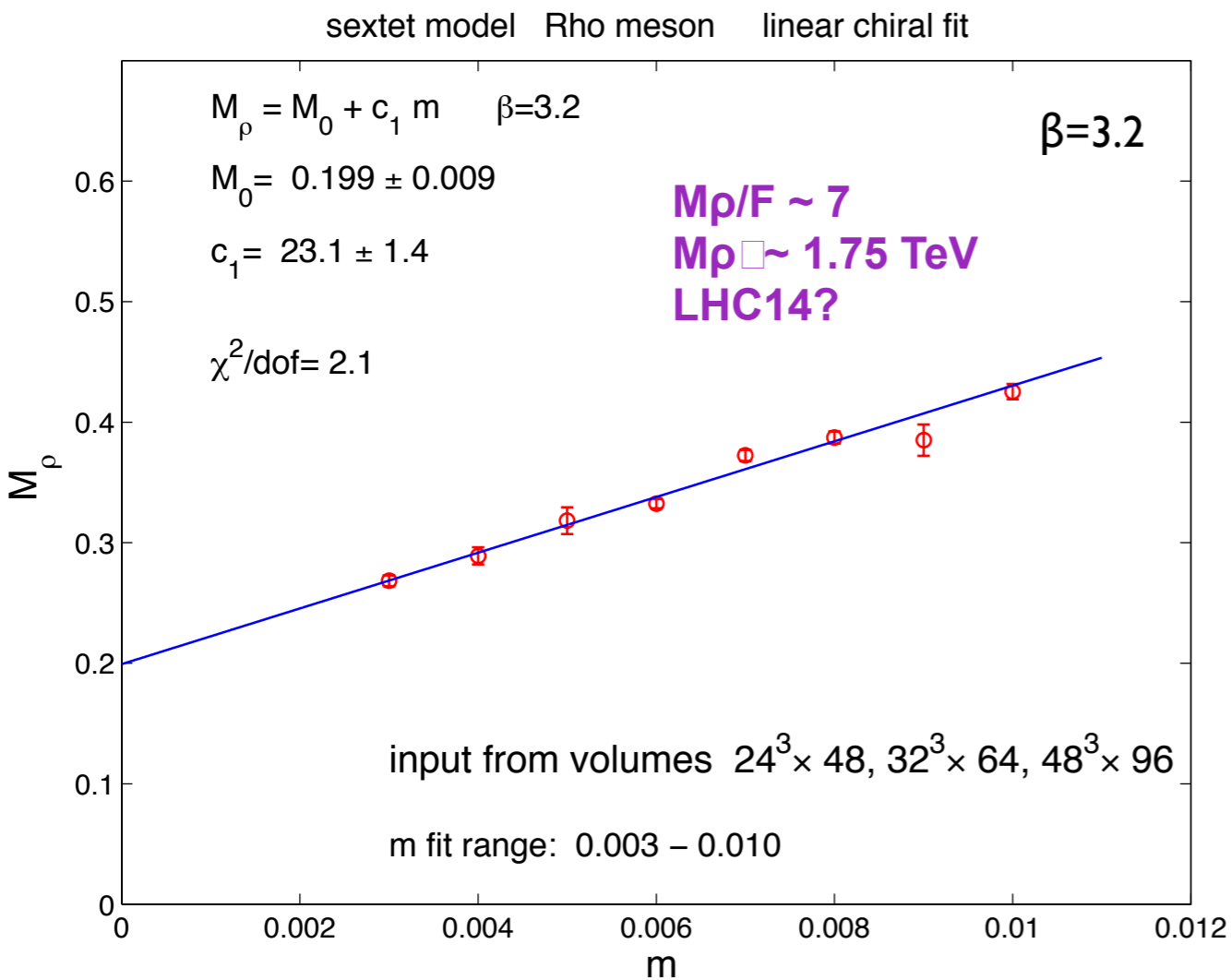
# Nf=2 SU(3) sextet chiral fits $M_\rho$ and $M(A_1)$



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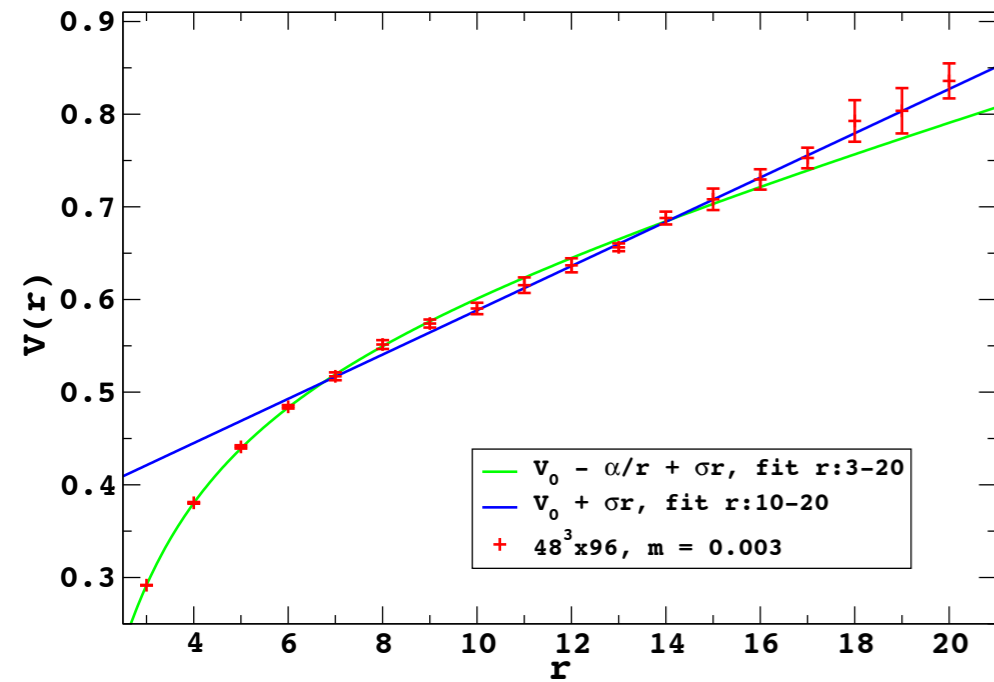
**$M_\rho$  remains heavy in massless fermion limit**

**perhaps within LHC reach?**

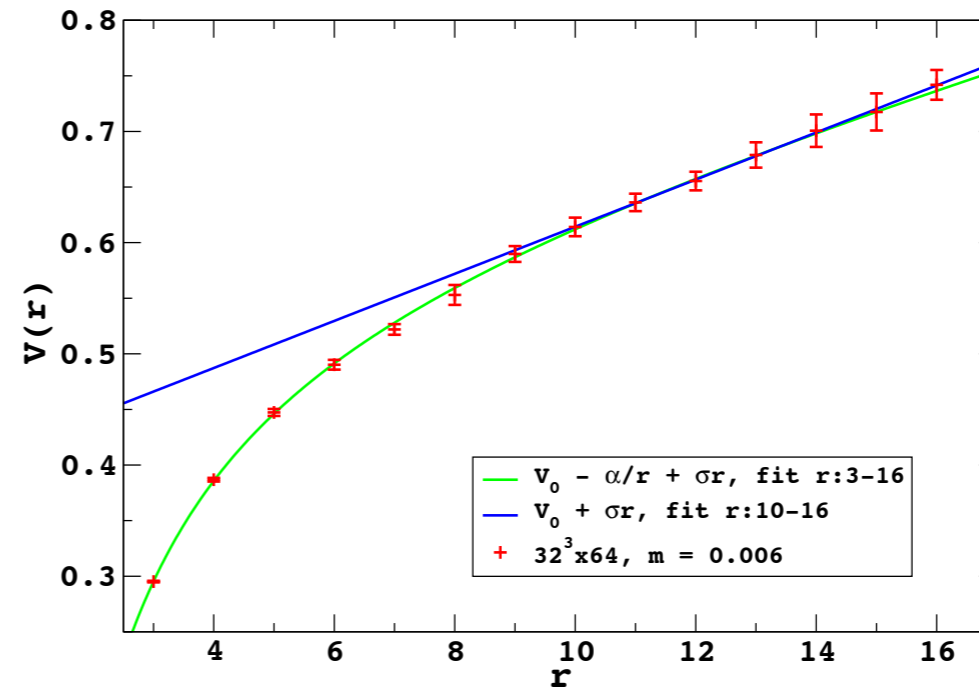
**parity partners remain split in massless fermion limit**

# sextet simulations **confining force at finite m** (LHC group)

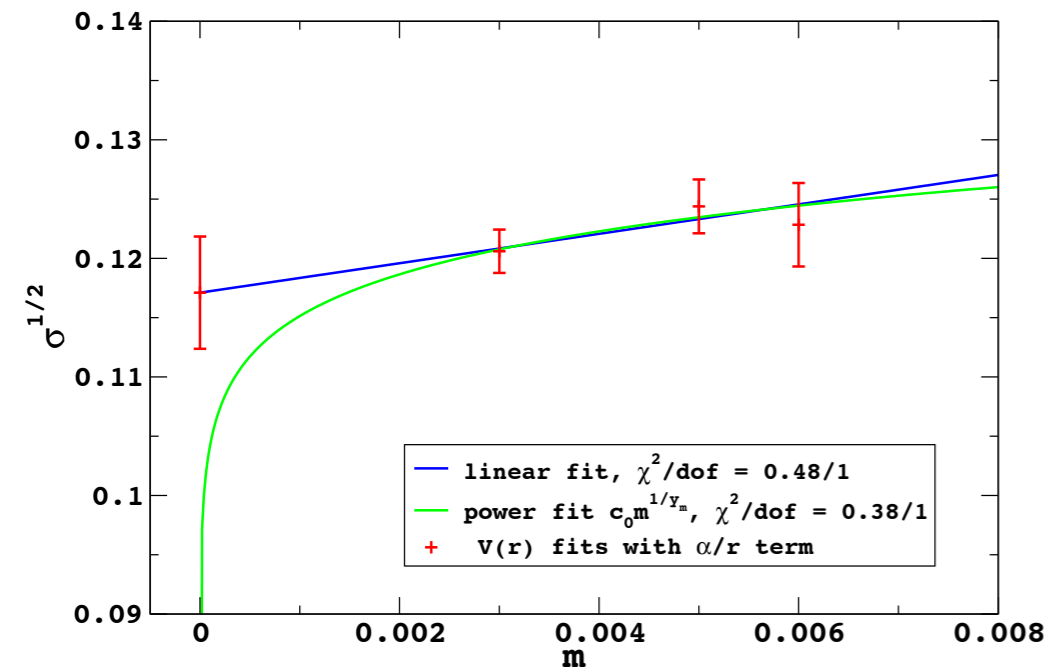
sextet  $N_f = 2, \beta = 3.20$



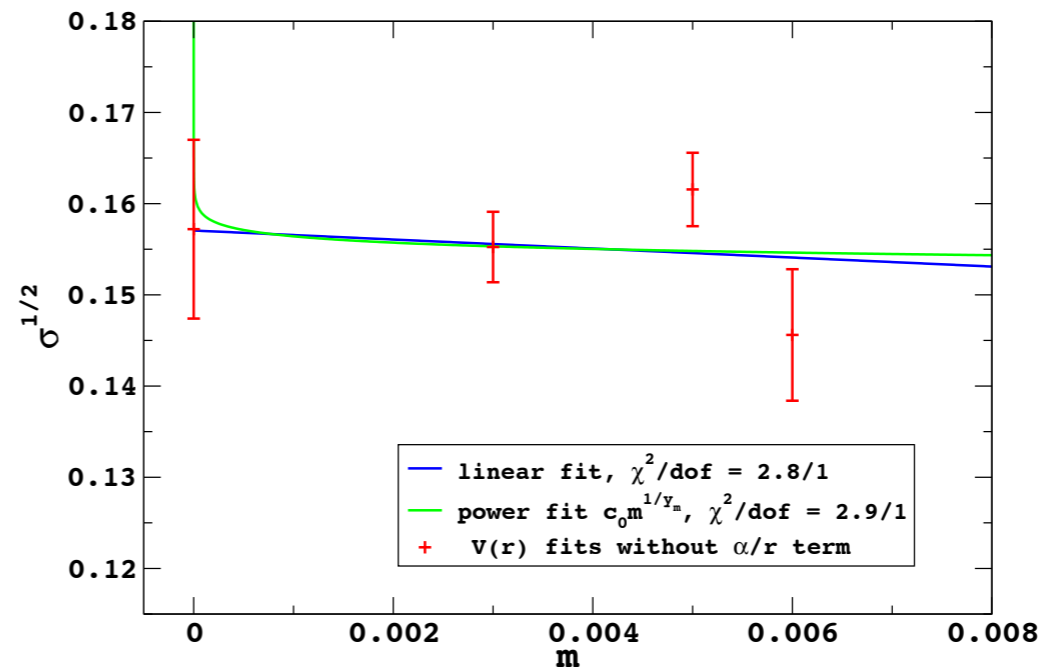
sextet  $N_f = 2, \beta = 3.20$



sextet  $N_f = 2, \beta = 3.20$



sextet  $N_f = 2, \beta = 3.20$



**$1/1+\gamma \sim 0.04(4)$  ? conformal  $\gamma \sim \text{infinite}$  would be needed**

**finite temperature**

# EW phase transition in sextet model - early universe

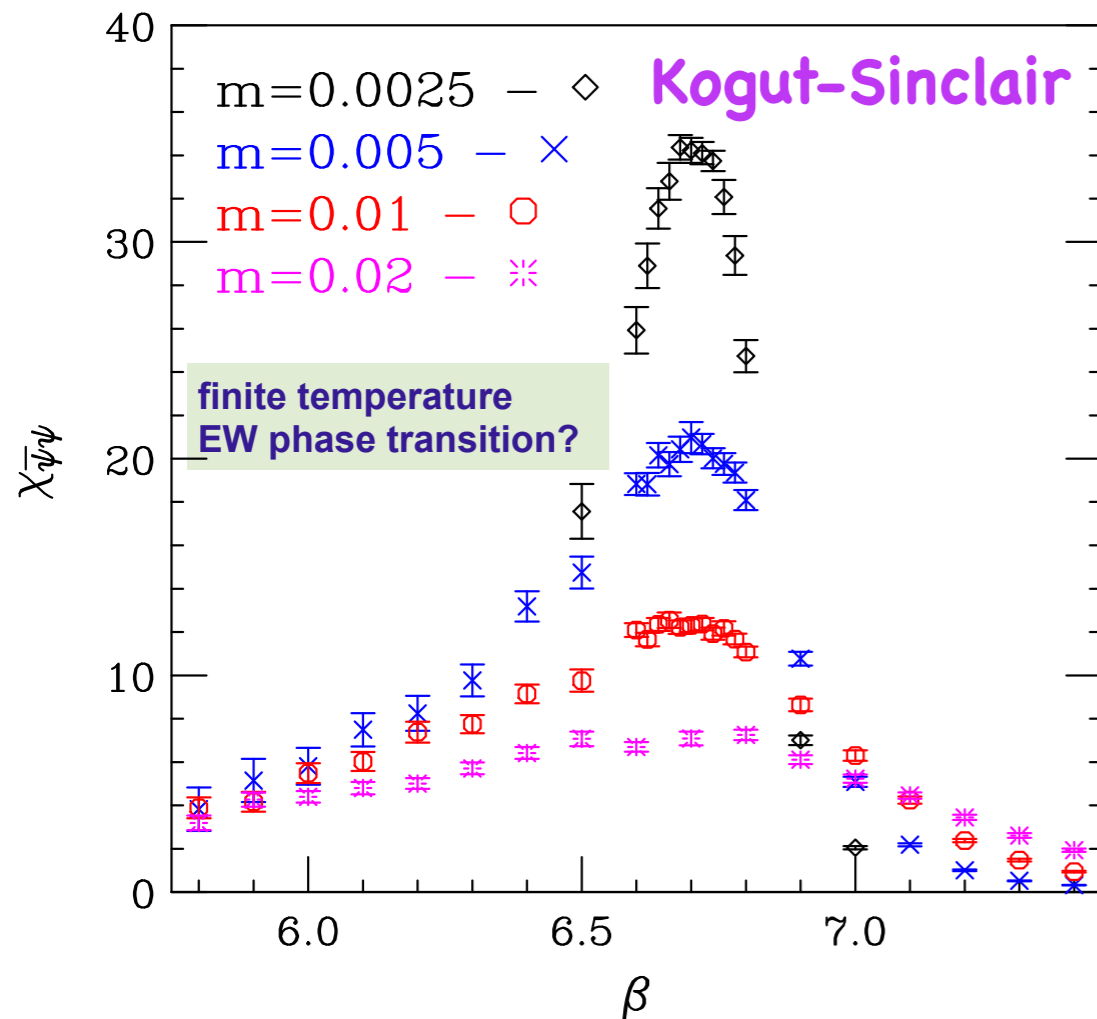
Kogut-Sinclair consistent with  $\chi$  SB phase at T=0

relevance in early cosmology

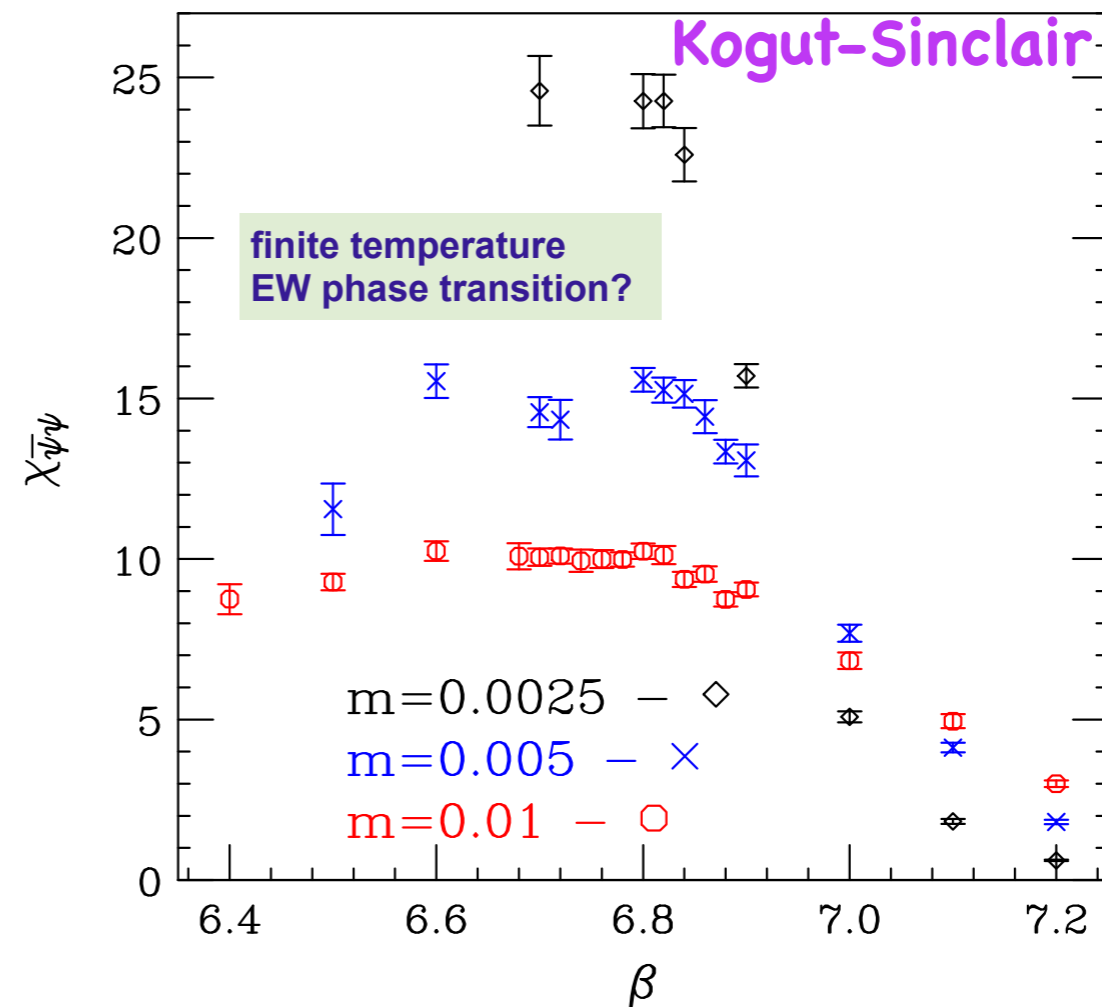
We are planning to run sextet thermo after model passed other tests

Third massive fermion flavor (electroweak singlet) dark matter?

$16^3 \times 8$  lattice



$24^3 \times 12$  lattice



# Summary and Outlook

Consistency with  $\chi_{SB}$  in  $N_f=2$   $SU(3)$  sextet model

Results of DeGrand et al. reconciled if walking or nearly walking coupling

Scalar spectrum from disconnected correlator remains highest priority  
dilaton-like?

Scale-dependent anomalous dimension of condensate from Dirac spectrum?

S-parameter and WW scattering if model remains standing? LHC14 phenomenology

Dark matter and electroweak phase transition ?

Do we have an impostor?





