



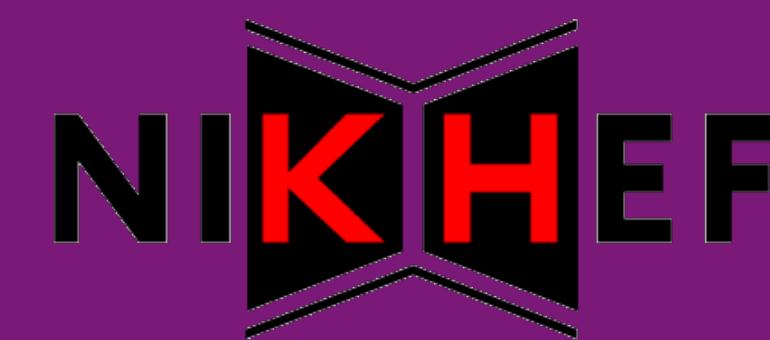
# First Evidence for the Rare Decay

$$B_s^0 \rightarrow \mu^+ \mu^-$$



Siim Tolk<sup>1</sup> on behalf of the LHCb collaboration

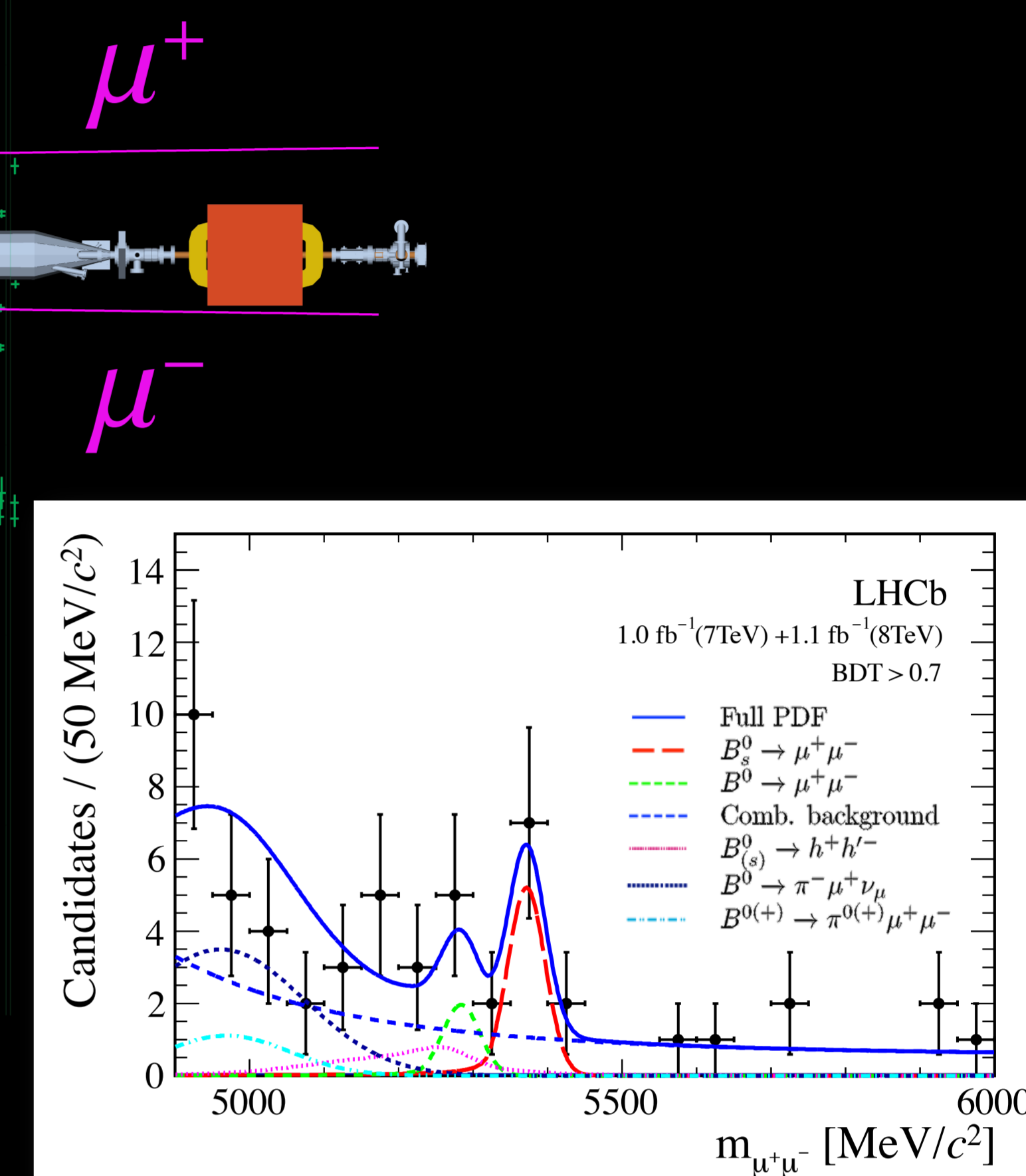
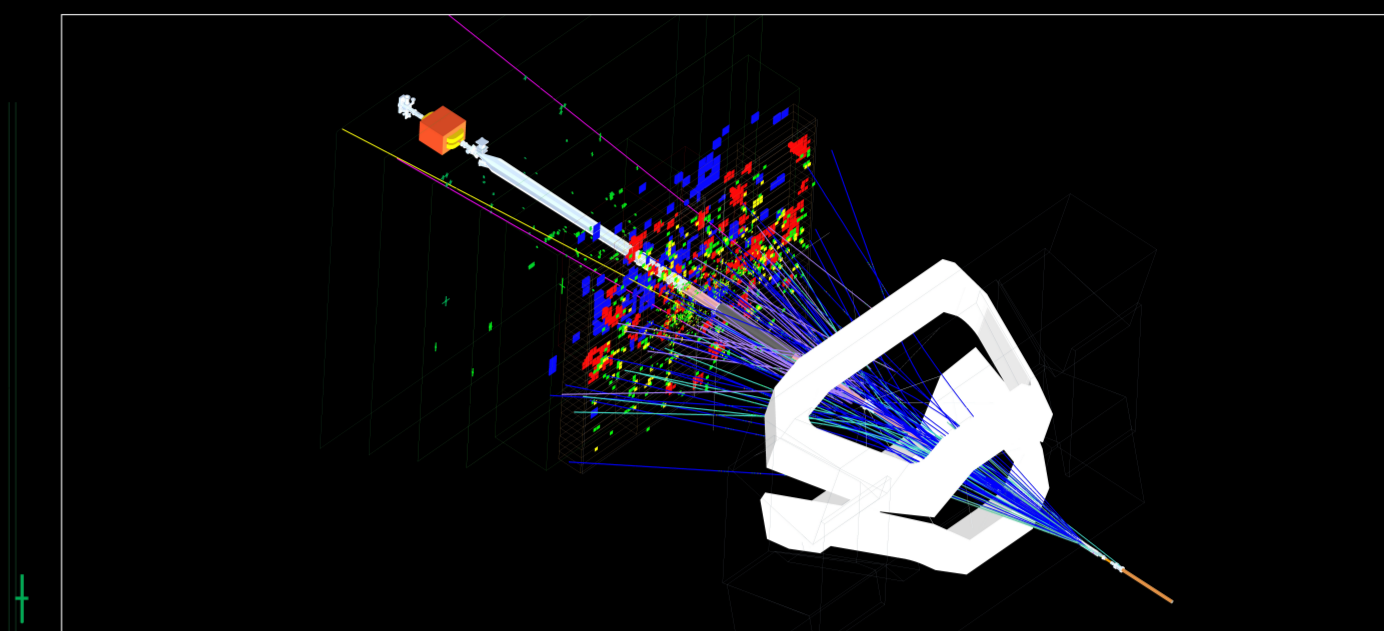
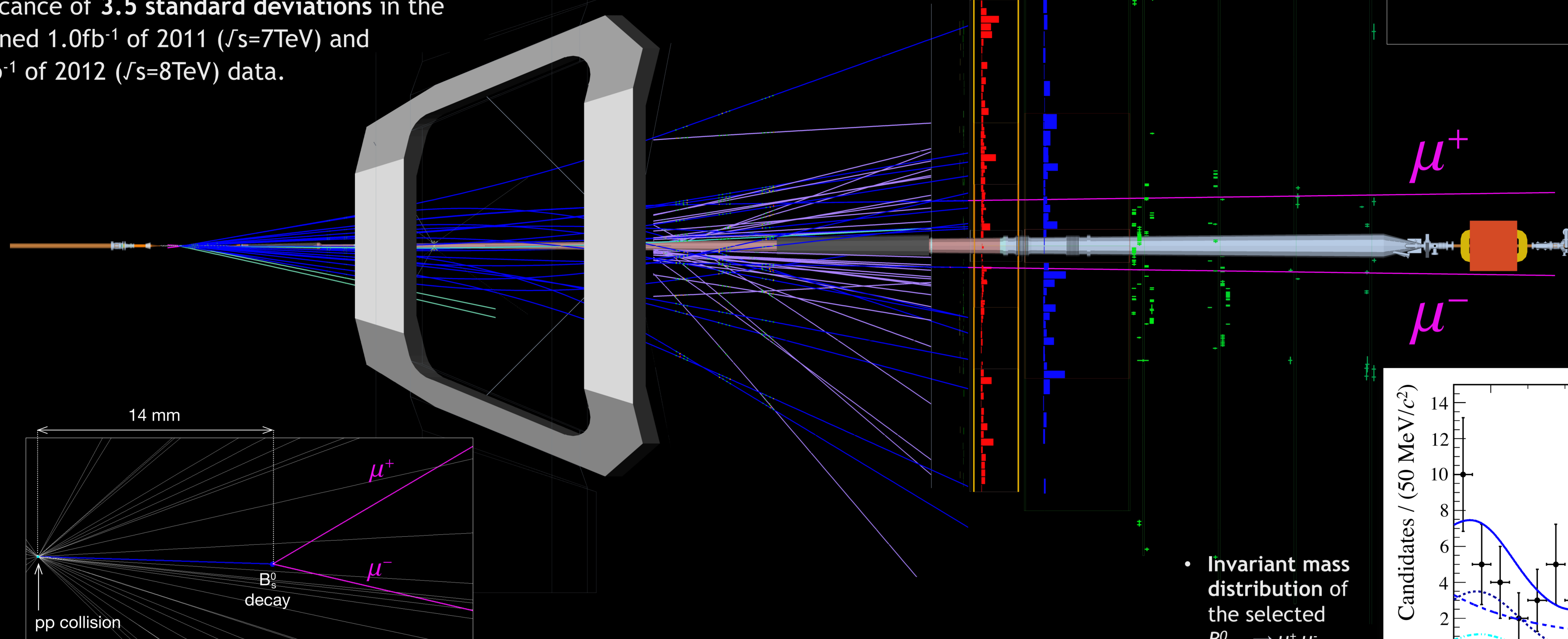
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The LHCb Collaboration  
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## ONE OF THE RAREST DECAYS EVER SEEN

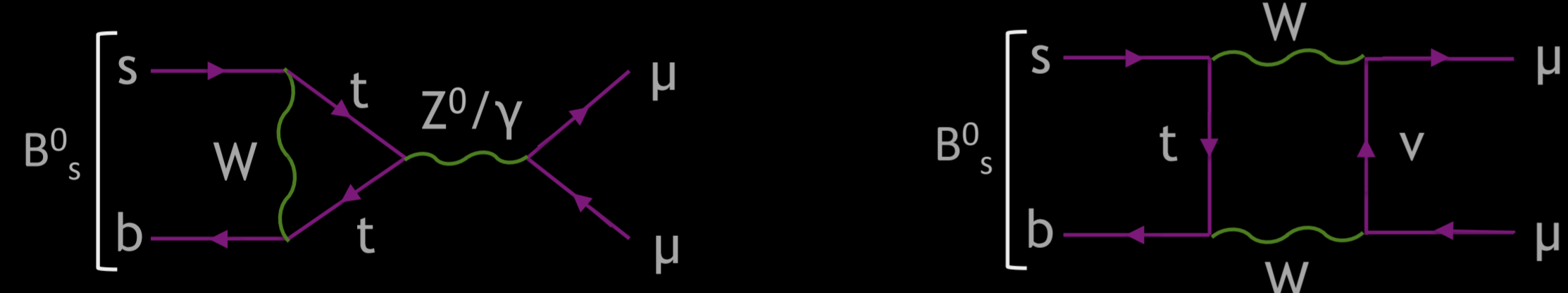
LHCb observed evidence for a very rare decay of a neutral  $B_s$  meson into a pair of muons[1]. The probability that the background could produce such an excess or larger is mere  $5.3 \times 10^{-4}$ . This corresponds to signal significance of **3.5 standard deviations** in the combined  $1.0 \text{ fb}^{-1}$  of 2011 ( $\sqrt{s}=7 \text{ TeV}$ ) and  $1.11 \text{ fb}^{-1}$  of 2012 ( $\sqrt{s}=8 \text{ TeV}$ ) data.



Invariant mass distribution of the selected  $B_{s,d}^0 \rightarrow \mu^+ \mu^-$  candidates in high BDT region

## Why so rare?

New physics (NP) is easiest to spot where according to the Standard Model (SM) an event is not likely to happen. In SM,  $B_{s,d}^0 \rightarrow \mu^+ \mu^-$  decays are so rare because (i) only decays through loop diagrams are allowed (flavor-changing neutral-current), and (ii) strong helicity suppression. The main SM diagrams contributing to  $BR(B_s^0 \rightarrow \mu^+ \mu^-)$  are



LHCb measures a flavor-averaged and time-integrated branching ratio. For the same quantity, precise SM predictions read[2]:

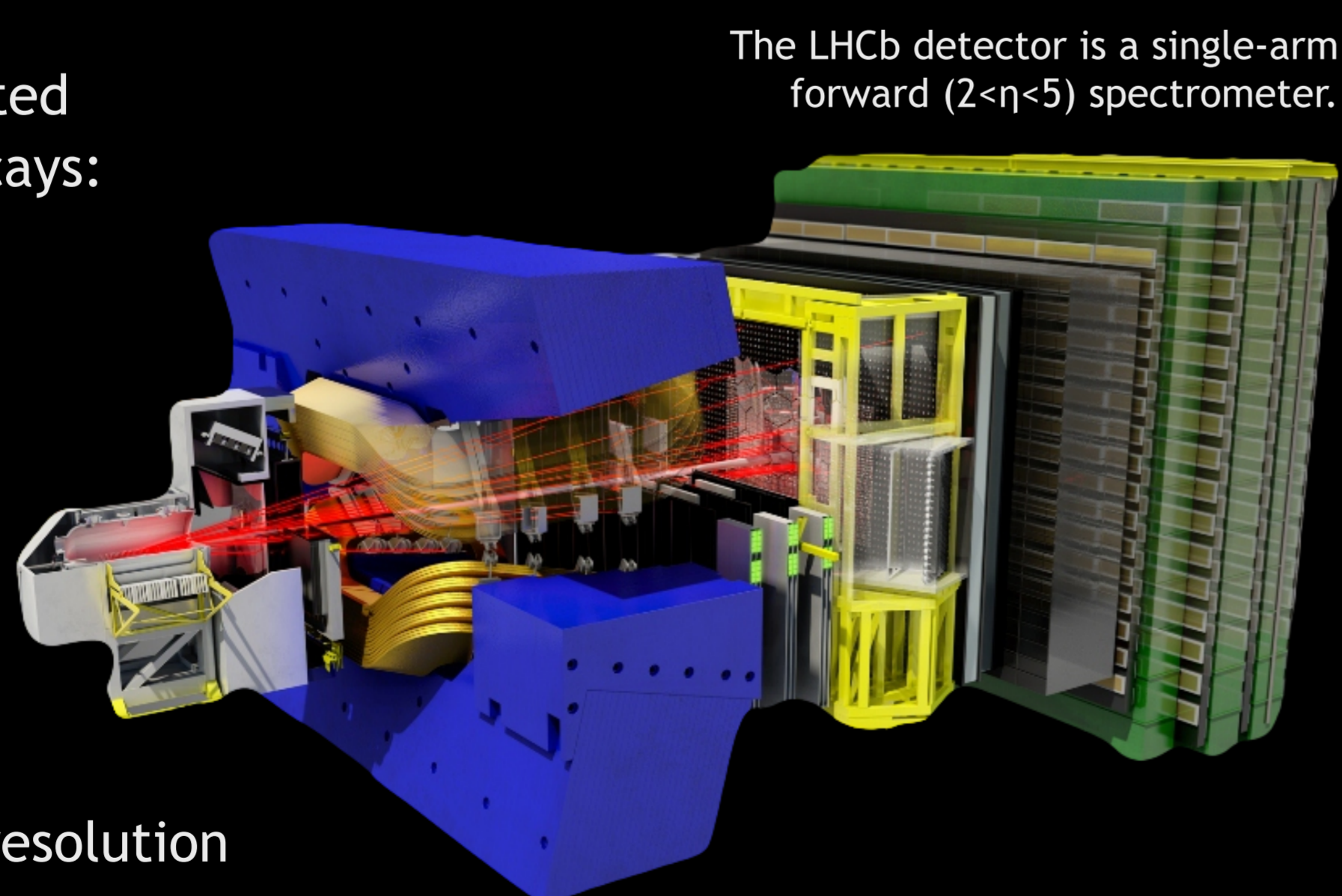
$$BR(B_s^0 \rightarrow \mu^+ \mu^-)^{(t) \chi_{CP}} = (3.54 \pm 0.30) \times 10^{-9}$$

$$BR(B_d^0 \rightarrow \mu^+ \mu^-)^{(t) \chi_{CP}} = (1.07 \pm 0.10) \times 10^{-10}$$

## Why LHCb?

The LHCb detector is well suited for measuring  $B_{s,d}^0 \rightarrow \mu^+ \mu^-$  decays:

- Abundant  $bb_{bar}$  pairs  
 $\sigma(pp \rightarrow bb_{bar} X) \sim 300 \mu\text{b}$  at 7 TeV [2]
- Outstanding muon ID performance  
 $\varepsilon(\mu \rightarrow \mu) \sim 97\%$  ( $p > 10 \text{ GeV}/c$ ),  
 $\varepsilon(h \rightarrow \mu) < 1\%$  ( $p > 10 \text{ GeV}/c$ ).
- Excellent mass and impact parameter resolution  
 $\sigma(B_{s,d}^0 \rightarrow \mu \mu) \sim 25 \text{ MeV}/c^2$   
 $\sigma(IP) \sim 25 \mu\text{m}$  ( $p_T = 2 \text{ GeV}/c^2$ )
- Efficient trigger on  $low-p_T$  muons



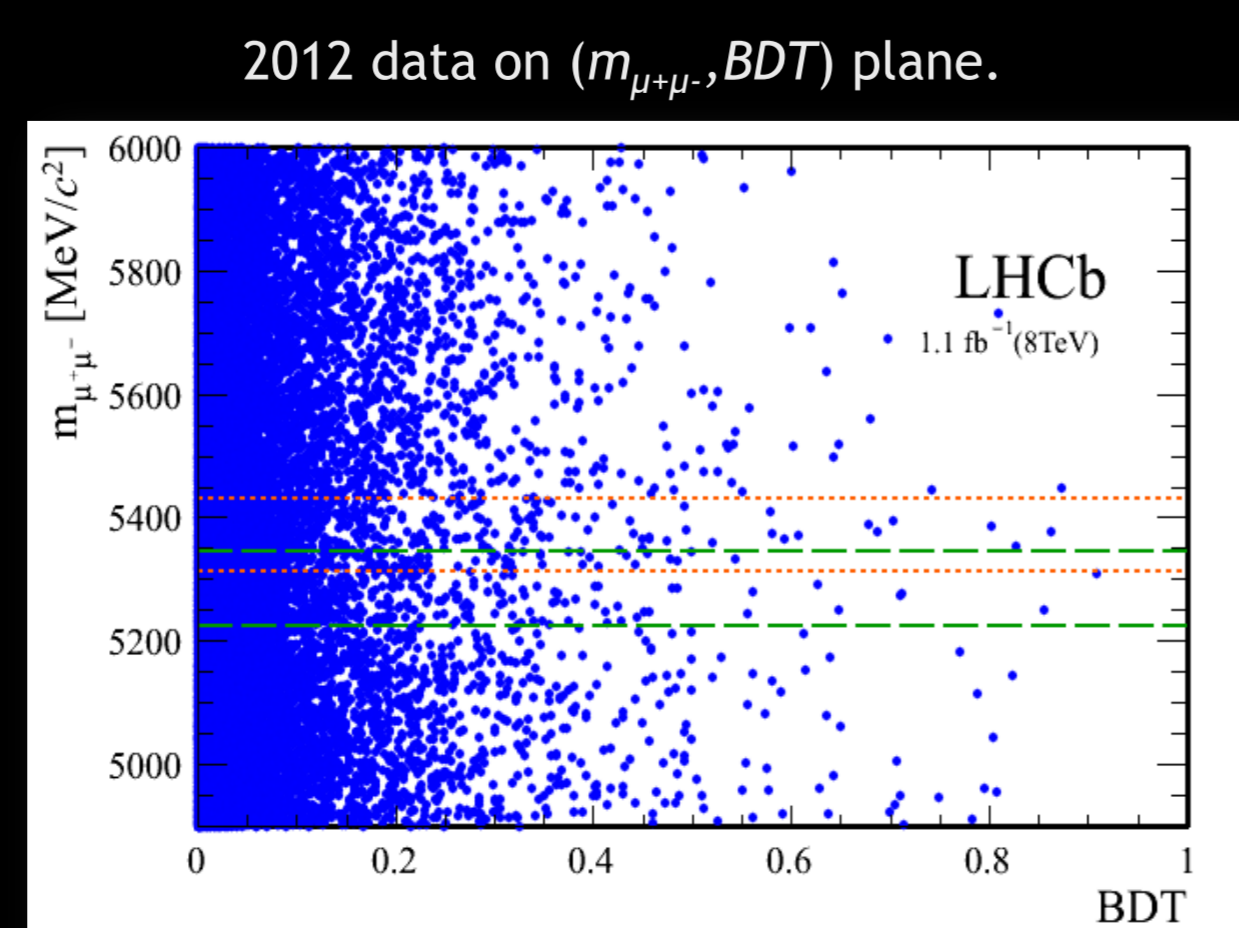
The LHCb detector is a single-arm forward ( $2 < \eta < 5$ ) spectrometer.

## How do we separate the signal candidates?

After a loose pre-selection, we classify the events in two independent variables:

- Di-muon invariant mass ( $m_{\mu^+ \mu^-}$ )
- Boosted Decision Tree output (BDT)

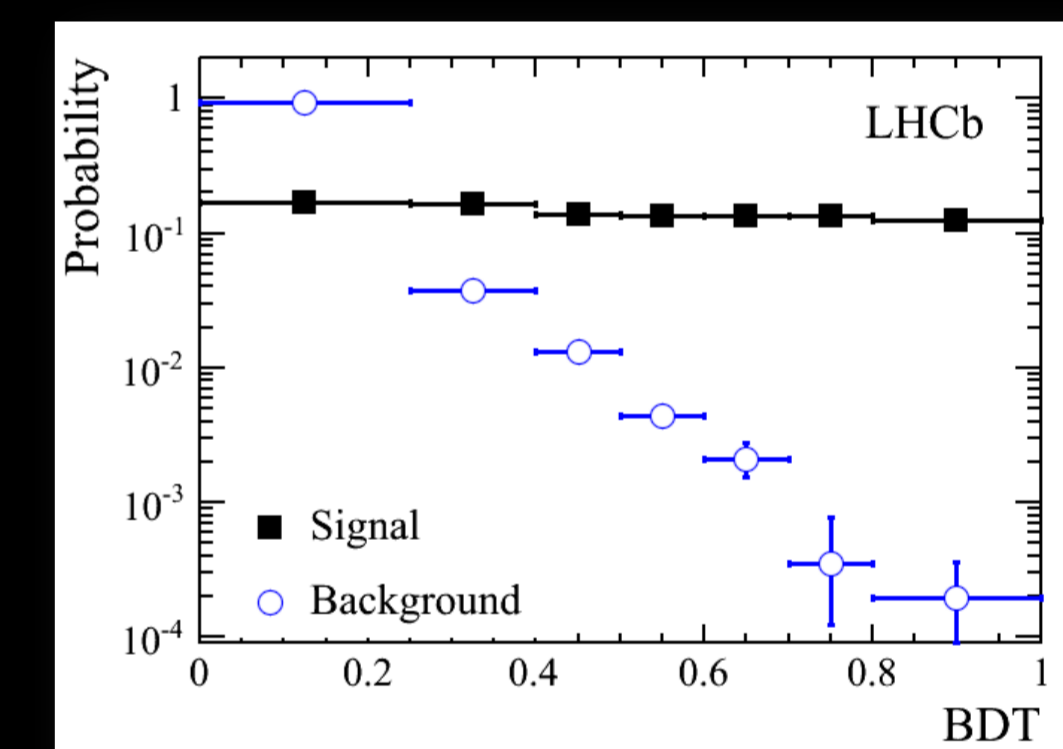
We divide the BDT range into 8 bins. In each bin, we compare the number of observed events to the number of expected signal and background events.



## Boosted Decision Tree (BDT)

- BDT combines the separation power of 9 carefully chosen topological and kinematic variables.
- We train the BDT with simulated  $B_{s,d}^0 \rightarrow \mu^+ \mu^-$  events for the signal and  $bb_{bar} \rightarrow \mu^+ \mu^- X$  for the background.
- We use  $B_s^0 \rightarrow K^+ \pi^-$  data to describe the signal distribution in BDT.

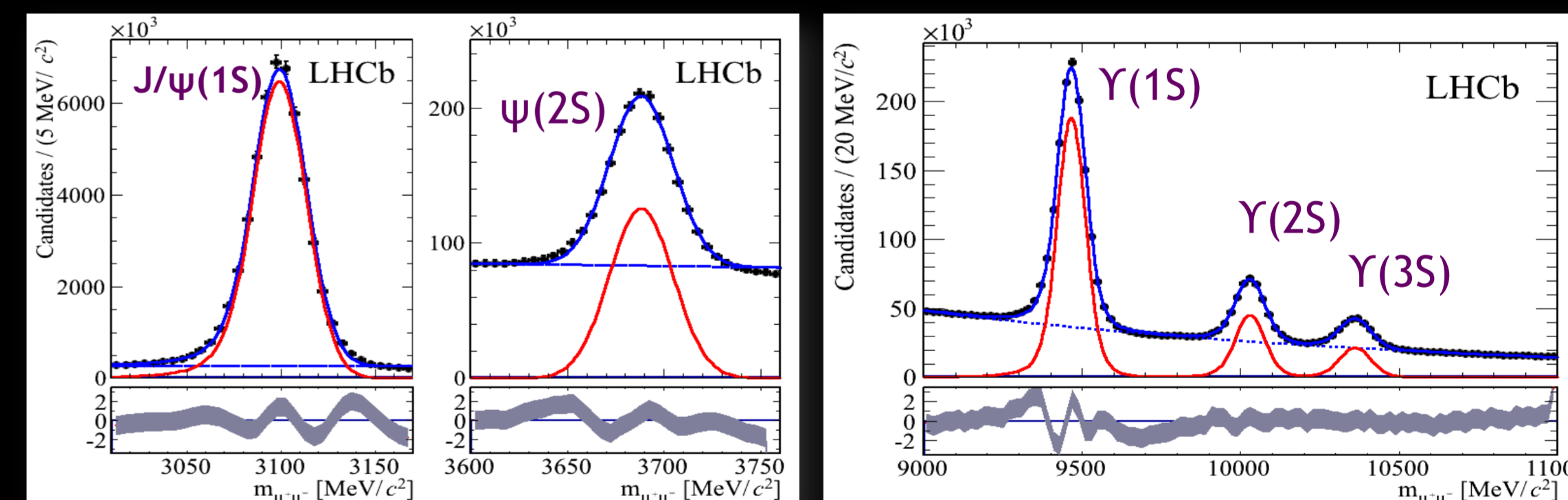
BDT distribution for signal (black) and combinatorial background (blue).



## Signal distribution in mass

We obtain the  $B_{s,d}^0$  mass central values from the  $B_{s,d}^0 \rightarrow h^+ h^-$  data (where  $h$  is a kaon or pion). The mass resolutions are determined by combining the results from the mass central value fits with the power-law interpolation between  $\mu^+ \mu^-$  resonances:

- $B_s^0$  and  $B_d^0$  mass peaks are separated by more than three times the mass resolution ( $\Delta m \sim 87 \text{ MeV}/c^2$ ).



## Normalization

The observed  $B_{s,d}^0 \rightarrow \mu^+ \mu^-$  yields are translated into absolute branching ratios by normalizing to well known  $B_{u,c}^0 \rightarrow J/\psi K^+$  and  $B_d^0 \rightarrow K^+ \pi^-$  branching ratios.

- Single event sensitivity:  
 $\sim 2.5 \times 10^{-10}$  for  $B_s^0 \rightarrow \mu^+ \mu^-$   
 $\sim 6.5 \times 10^{-11}$  for  $B_d^0 \rightarrow \mu^+ \mu^-$

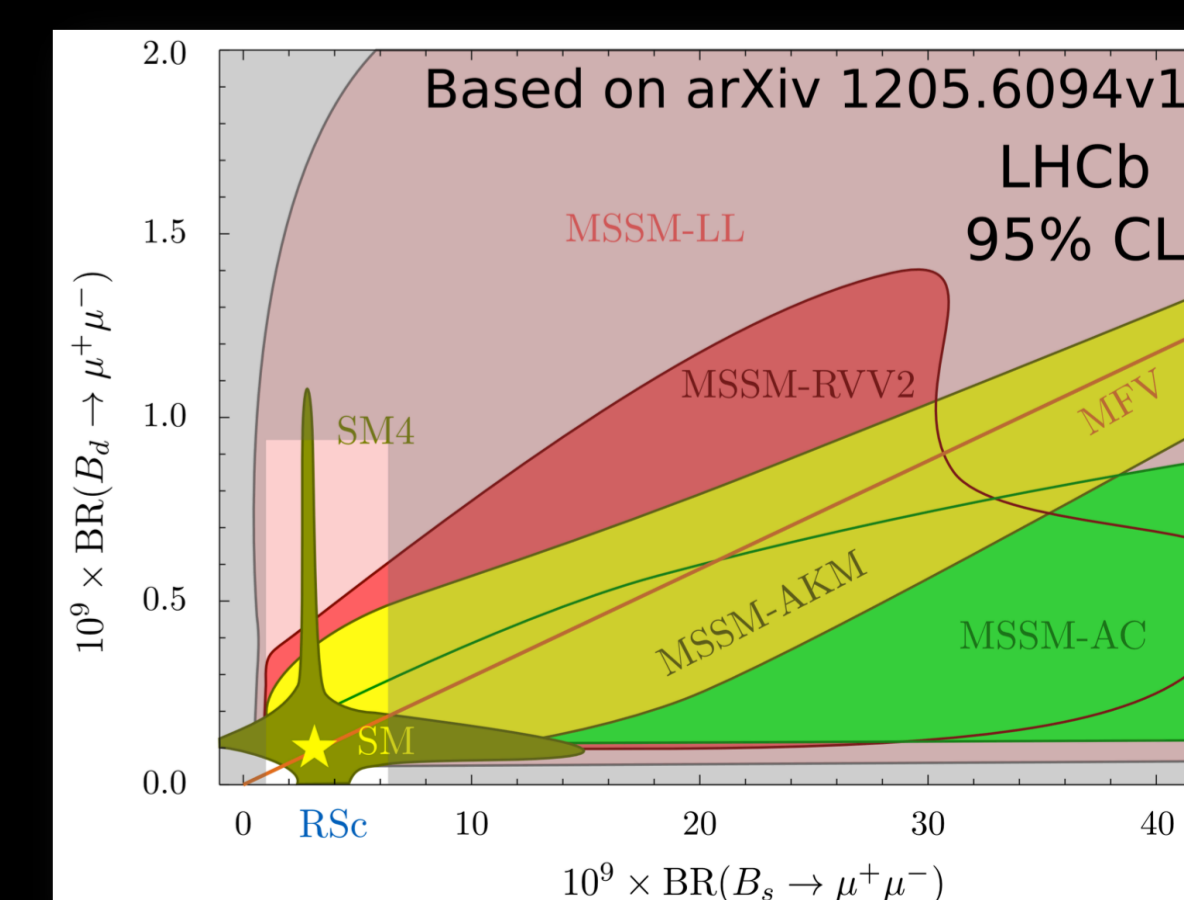
$$BR(B_{s,d}^0 \rightarrow \mu^+ \mu^-) = \frac{N_{B_{s,d}^0 \rightarrow \mu^+ \mu^-} \varepsilon_{sig}^{norm} f_{sig}^{norm}}{N_{norm} \varepsilon_{sig} f_{sig}} \times BR^{norm}$$

## Impact on new physics

Contributions from NP models to the  $BR(B_s^0 \rightarrow \mu^+ \mu^-)$  and  $BR(B_d^0 \rightarrow \mu^+ \mu^-)$  are strongly limited by the LHCb results[1]:

$$BR(B_s^0 \rightarrow \mu^+ \mu^-) = [3.2^{+1.4}_{-1.2} (stat)^{+0.5}_{-0.3} (syst)] \times 10^{-9}$$

$$BR(B_d^0 \rightarrow \mu^+ \mu^-) < 9.4 \times 10^{-10} @ 95\% CL$$



Constraints on new physics models [4]  
Shaded area is excluded by LHCb @ 95% CL.

## References

- The LHCb Collaboration, Phys. Rev. Lett. 110, 021801 (2013)
- A.J. Buras *et al.*, Eur. Phys. J. C 72, 2172 (2012)
- The LHCb Collaboration, Phys. Lett. B 694 (2010) 209
- D. M. Straub, arXiv:1205.6094v1 (2012)