

Review of Synchrotron Radiation effects in the FFS

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February 27, 2013

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Beamsize

We are interested in the horizontal beamsize at the IP.

Horizontal plane

$$\sigma^2 = \sigma_0^2 + \sigma_i^2 + \sigma_{rad}^2$$

$\sigma_0 \equiv$ zeroth order approx.

$\sigma_i \equiv$ result from aberrations

$\sigma_{rad} \equiv$ interaction with magnets

Evaluated by:

- ▶ tracking of particles
- ▶ mathematical approximations

Beam Radiation Model

x describes the displacement of a particle at ($s = L$, e.g. IP) due to radiation, where all other effects are ignored.

$$x = \sum_{i=1}^{N(T)} \Delta x_{i,total} - x_0 \quad (1)$$

- ▶ $\Delta x_{i,total}$: is the total deviation due to the i^{th} photon radiated
- ▶ x_0 : is $\langle \sum_{i=1}^{N(T)} \Delta x_{i,total} \rangle$,
in order to make $\langle x \rangle = 0$, and $\sigma_{rad}^2 = \langle x^2 \rangle$
- ▶ N : is the number of photons radiated
- ▶ T : time to cross the bending magnet

Finding Δx_i

Δx_i is the effect at ($s = L$) due to a photon of energy u radiated at $s = s_i$, therefore it has to be propagated from s_i to L .

$$\Delta x_i = (u/E)R_{16}(s_i, L) \quad (2)$$

It becomes:

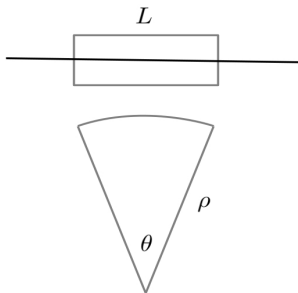
$$\sigma_{rad}^2 \approx C_2 \int \frac{E^5}{\rho^3} \left\{ \sqrt{\frac{\beta_L}{\beta_s}} [\eta \cos \Delta\phi(s, L) + (\alpha\eta + \beta\eta') \sin \Delta\phi(s, L)] - \eta_L \right\}^2 ds$$

- ▶ $C_2 = 4.13 \times 10^{-11} [\text{m}^2 \cdot \text{GeV}^{-5}]$
- ▶ E : is the beam energy

included in MAPCLASS2.

One dipole (theoretical expression)

Using $R_{16} = \rho(1 - \cos \theta)$



$$\begin{aligned}\sigma_{rad}^2 &= C_2 \int_0^L \frac{E^5}{\rho^3} R_{16}^2(s, L) ds \\ &= C_2 \int_0^\theta \frac{E^5}{\rho^3} [\rho(1 - \cos(\theta - \chi))]^2 \rho d\chi \\ &= C_2 E^5 \left[\frac{1}{4} (6\theta - 8 \sin \theta + \sin(2\theta)) \right]^1 \\ &= C_2 E^5 \left(\frac{\theta^5}{20} - \frac{\theta^7}{168} + \frac{\theta^9}{2880} - \frac{17\theta^{11}}{1330560} + O(\theta^{13}) \right)\end{aligned}$$

Theoretical expression for a drift has been also derived. Now, the **theoretical expression**, the **approximated result** and **tracking with PLACET** could be compared.

Some care should be taken when using **the expression above** due to numerical precision.

MAPCLASS2 uses on MAD-X twiss table results.

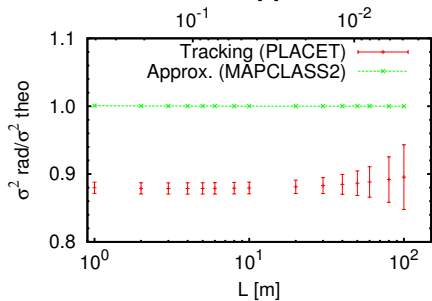
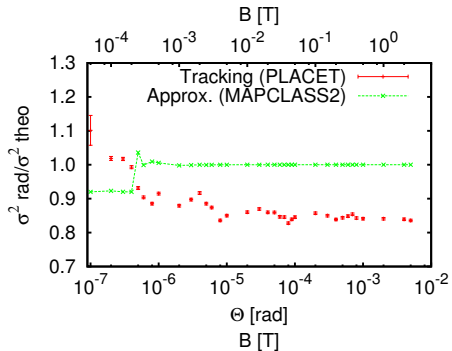
PLACET has two different radiation methods "Default" and "six_dim".

Radiation beamsizes has been normalized to the **theoretical value**.

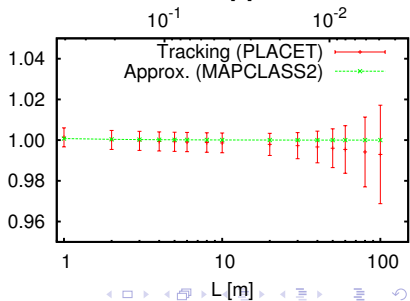
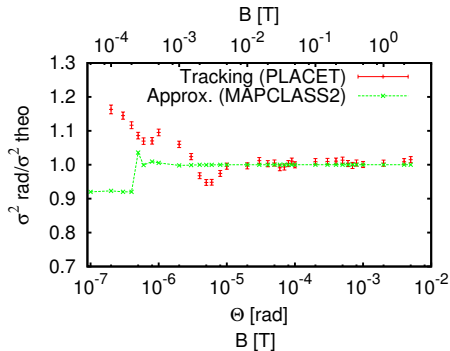
1

¹If E is considered constant.

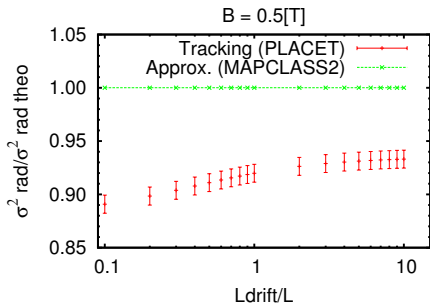
Default Synrad



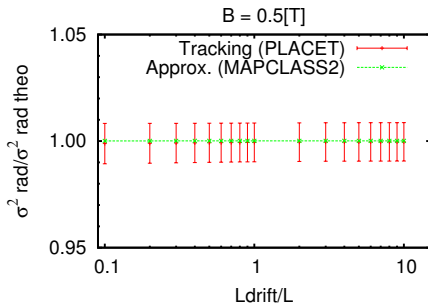
Flag "-six_dim 1"

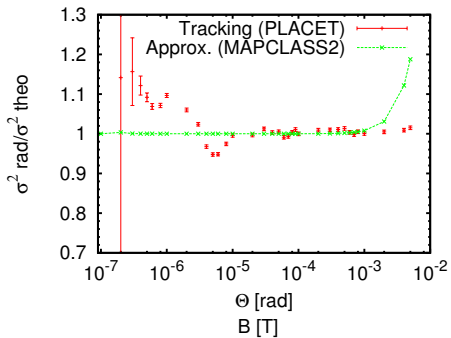


Default Synrad

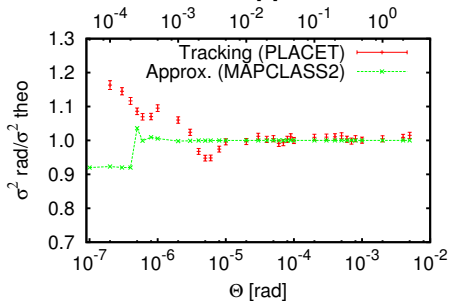


Flag "-six_dim 1"





Twiss



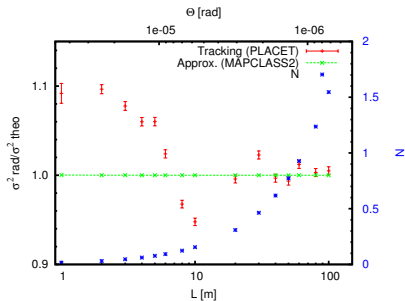
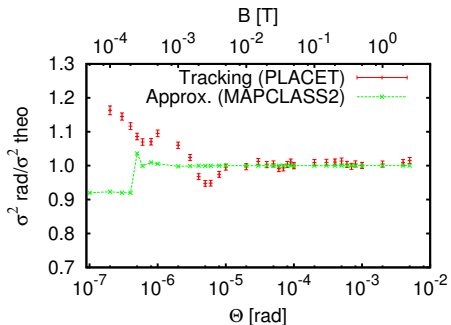
ptc_twiss

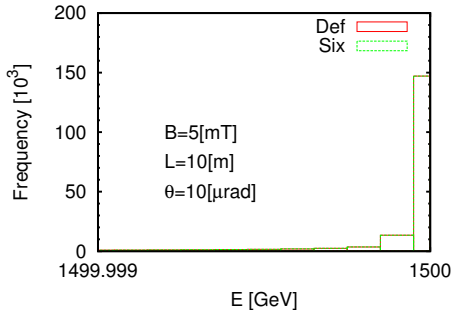
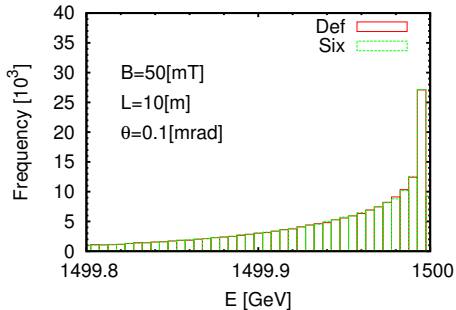
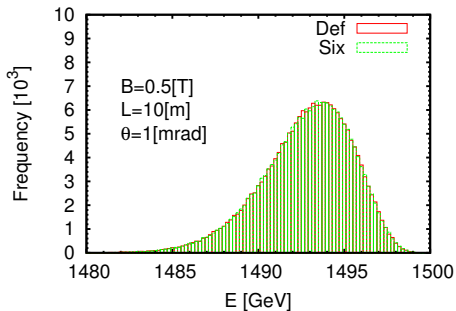
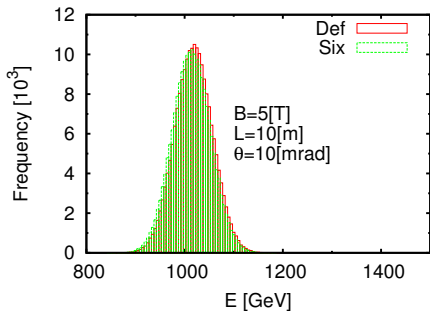
Low number of photons

Average number of photons emitted is

$$\langle N \rangle = \frac{1}{c} \int_0^L ds \int_0^\infty du n(u, s) \quad (3)$$

$$\approx C_1 E \theta \quad ; C_1 = 20.61 [\text{GeV}]^{-1} \quad (4)$$







Conclusions

- ▶ Mapclass2 agrees with theory, some care should be taken with twiss file.
- ▶ “six_dim” in PLACET radiation model is more accurate than default radiation calculation. Both shows differences with theoretical model for bends causing low photon emission.
- ▶ Theoretical model for low photon emission stills in check.

References

-  Sands, Matthew. Emittance growth from radiation fluctuations. SLAC/AP – 47. December, 1985.
-  Renier, Ives. Implementation and validation of the linear collider final focus prototype : ATF2 at KEK (Japan). Doctoral Thesis, LAL10-91. June 2010.

Additional slides

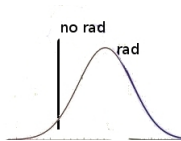
Dispersion function

$$\begin{pmatrix} \eta(s_2) \\ \eta'(s_2) \\ 1 \end{pmatrix} = \begin{pmatrix} C(s_1, s_2) & S(s_1, s_2) & R_{16}(s_1, s_2) \\ C'(s_1, s_2) & S'(s_1, s_2) & R_{26}(s_1, s_2) \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \eta(s_1) \\ \eta'(s_1) \\ 1 \end{pmatrix}$$

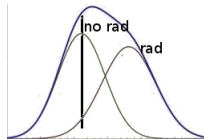
PLACET error bars

$$f = \frac{x_{rad}^2 - x_{norad}^2}{x_0^2}$$
$$\delta f = \frac{2}{\sqrt{N_{part}}} \frac{(x_{rad}^2 + x_{norad}^2)}{x_0^2}$$

Higher angle: most particles radiate.



Lower angle: some particles radiate, others don't.



$$un(u)du \equiv P(u/\hbar)du/\hbar = P(\omega)d\omega \quad (5)$$