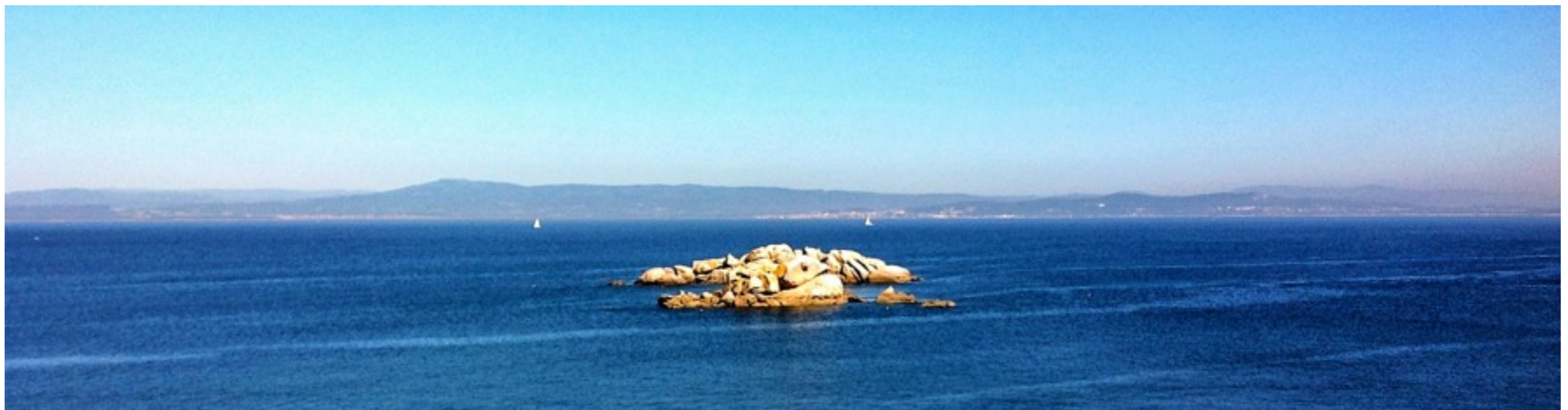
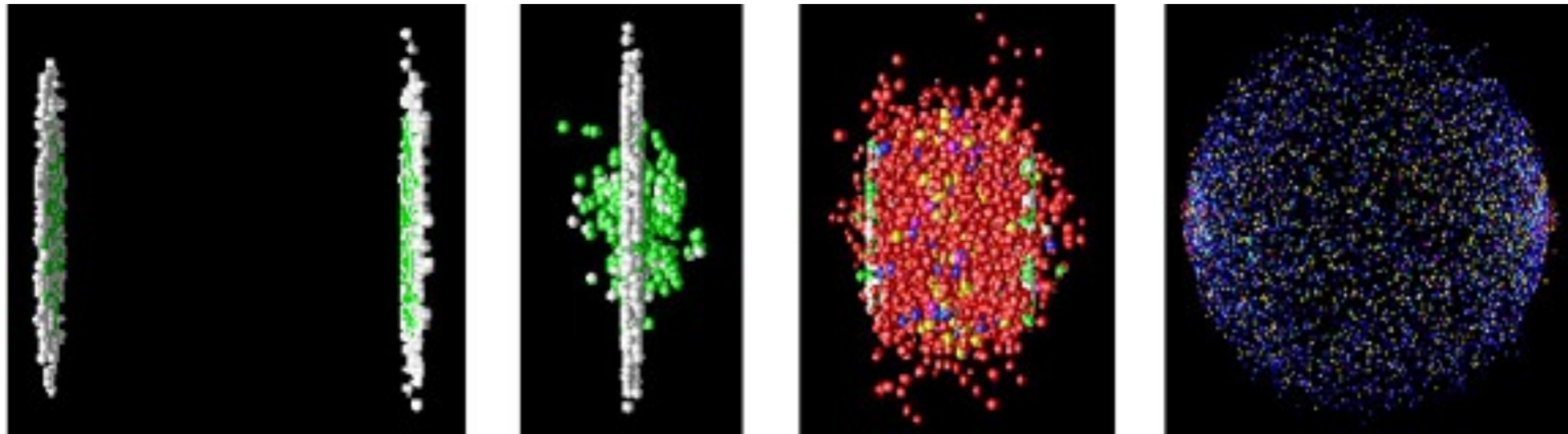


Harmonics in harmony

Jean-Yves Ollitrault (IPhT, Saclay)
Illa da Toxa, Sept.9, 2013



Nuclear collisions at the LHC



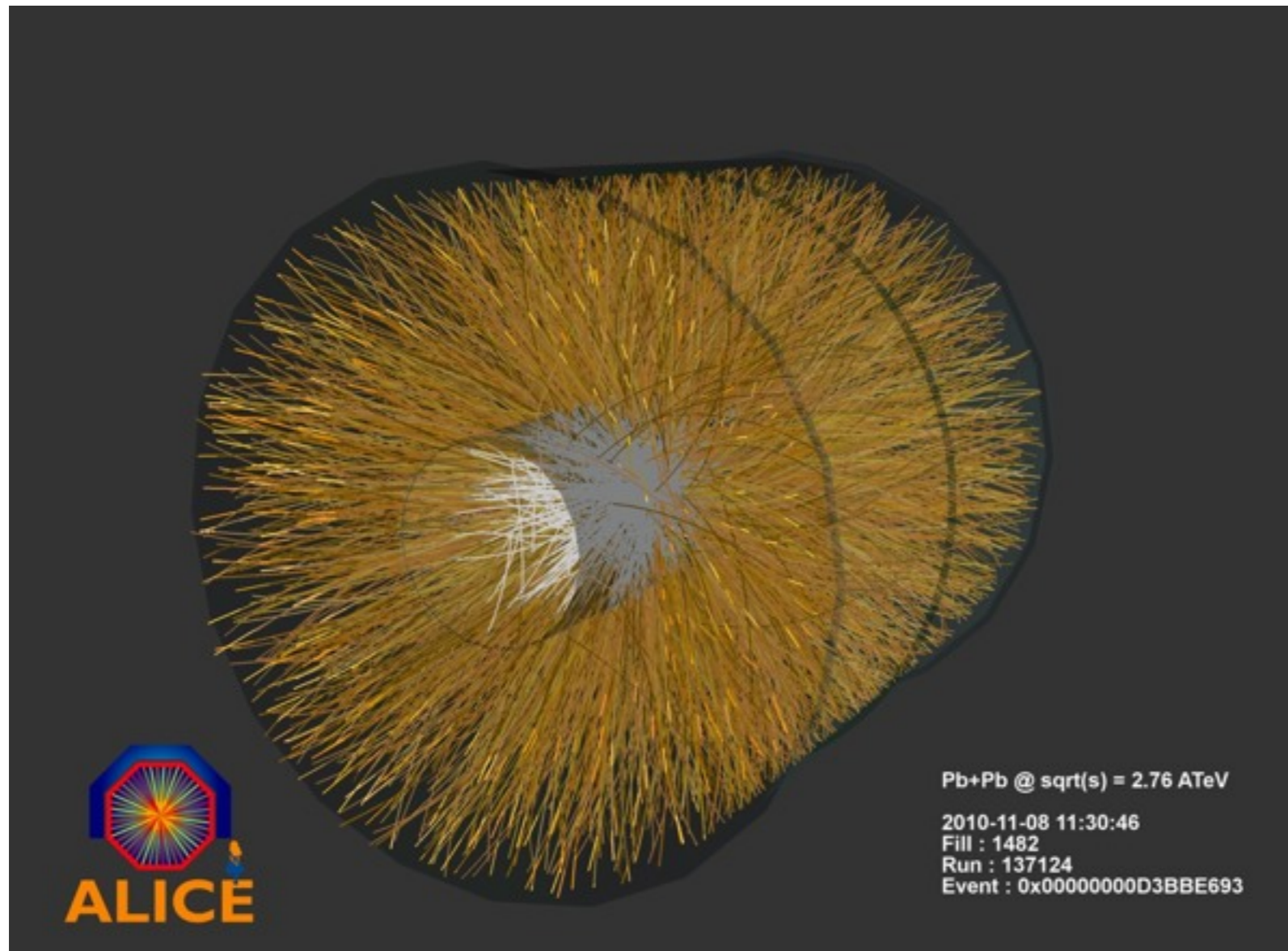
Initial stage: energy deposited in central rapidity region, thermalization

Kurkela Heller Epelbaum

Strongly-interacting expanding plasma: best described as a (macroscopic) small lump of **fluid**

*Luzum Strickland Schenke
Niemi Flörchinger Denicol*

What we see



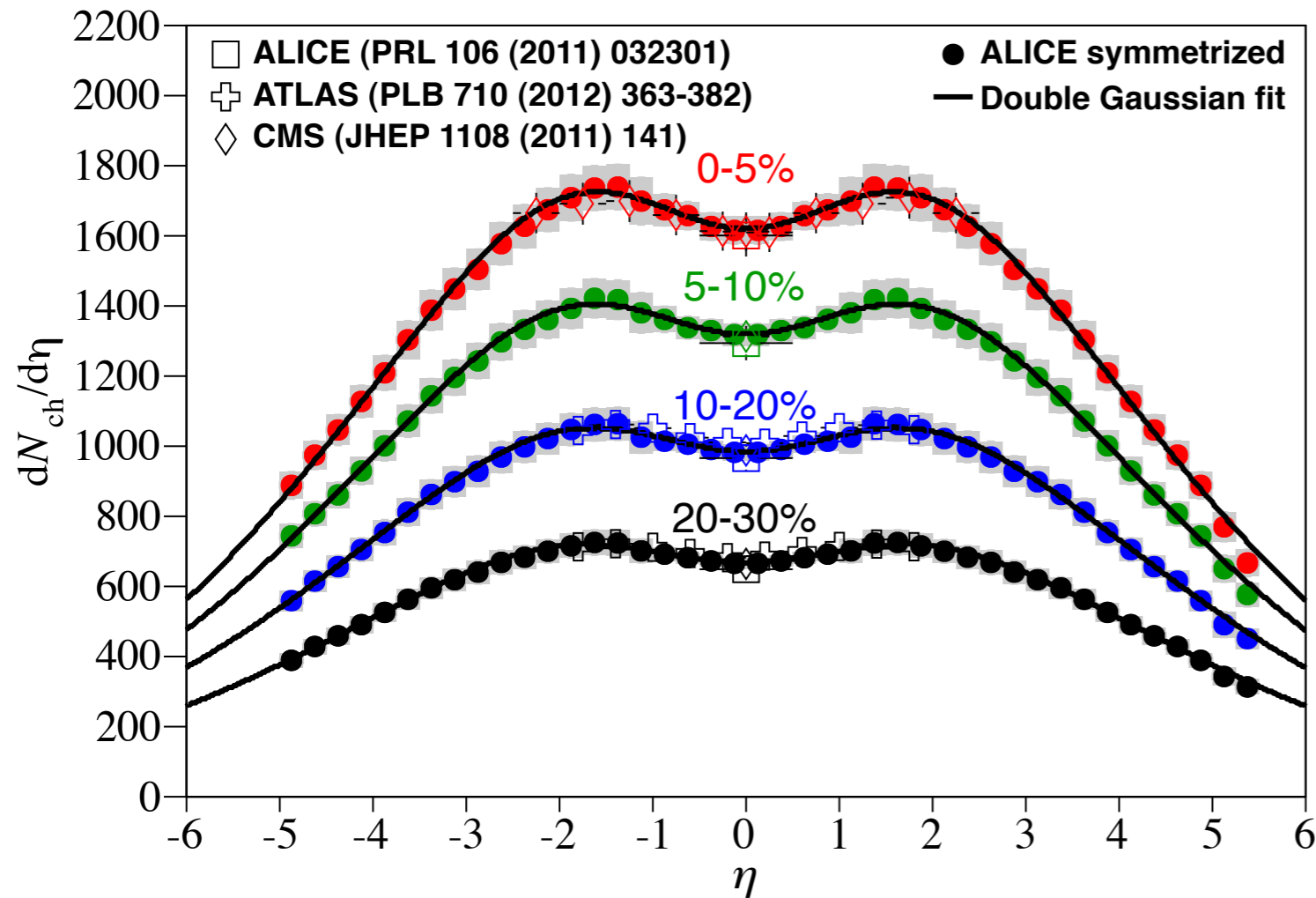
Trajectories of
charged particles:

polar angle θ
(or pseudorapidity
 $\eta = -\ln \tan \theta/2$)

azimuthal angle ϕ

Counting charged particles

distribution in pseudorapidity (\sim polar angle)



ALICE arXiv:1304.0347

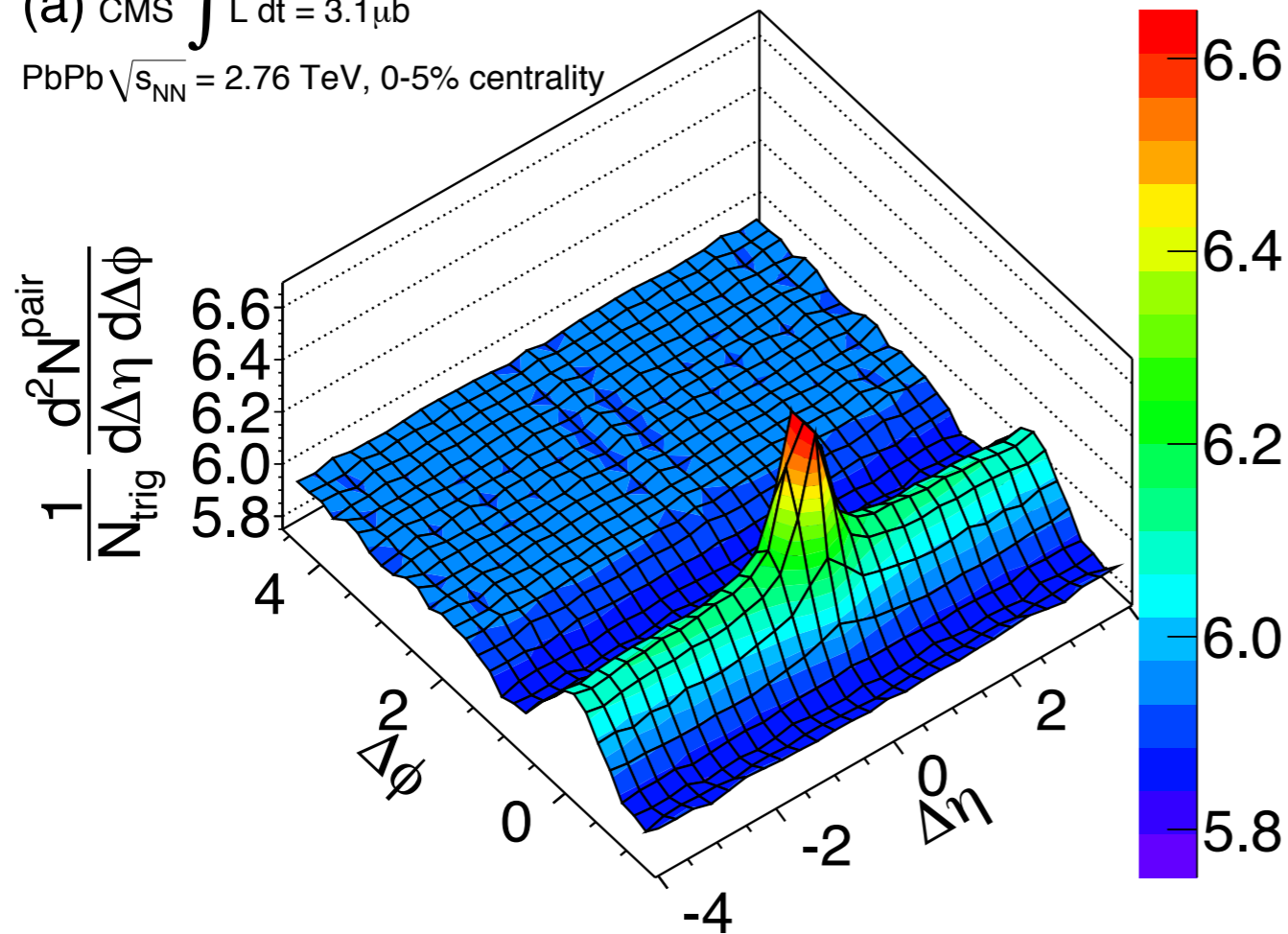
Collisions classified
from more central to
less central in 5% bins

More central creates
more particles

A central collision
typically produces
25000 particles.

Observing the small fluid: counting pairs

(a) CMS $\int L dt = 3.1 \mu\text{b}^{-1}$
PbPb $\sqrt{s_{\text{NN}}} = 2.76 \text{ TeV}$, 0-5% centrality

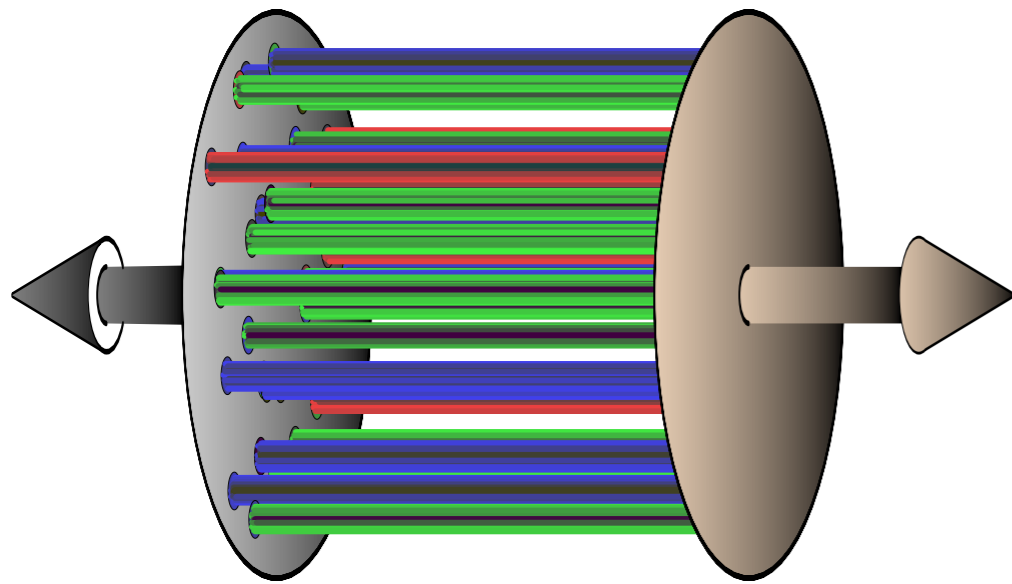


Number of pairs of particles versus relative azimuthal angle and pseudorapidity (\sim polar angle) in central Pb-Pb collisions

CMS arXiv:1105.2438

- *Ripple in a pond: cleanest signature of fluid behavior*

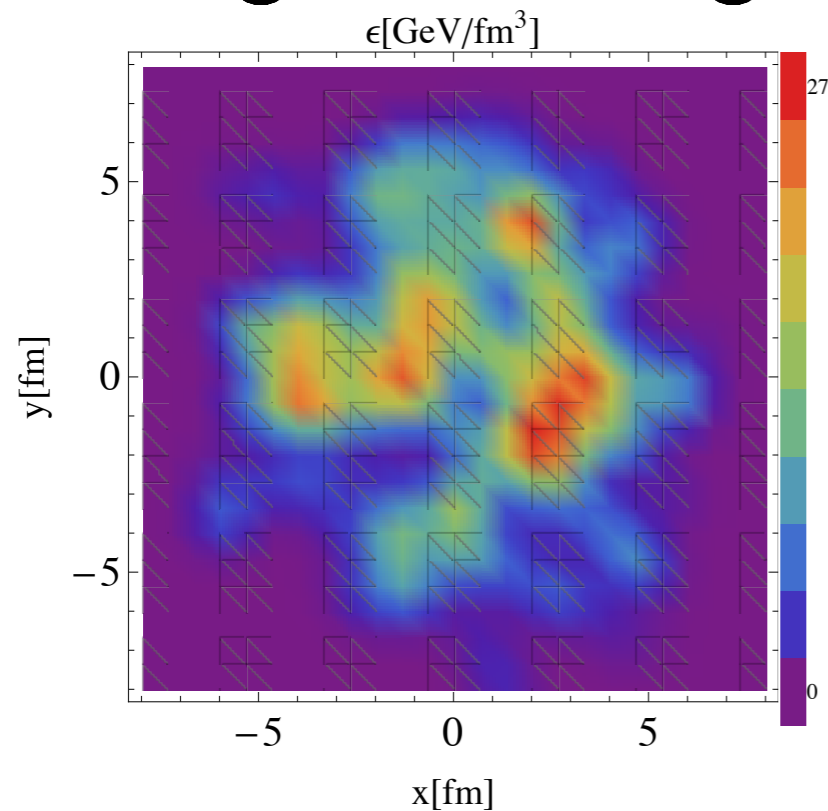
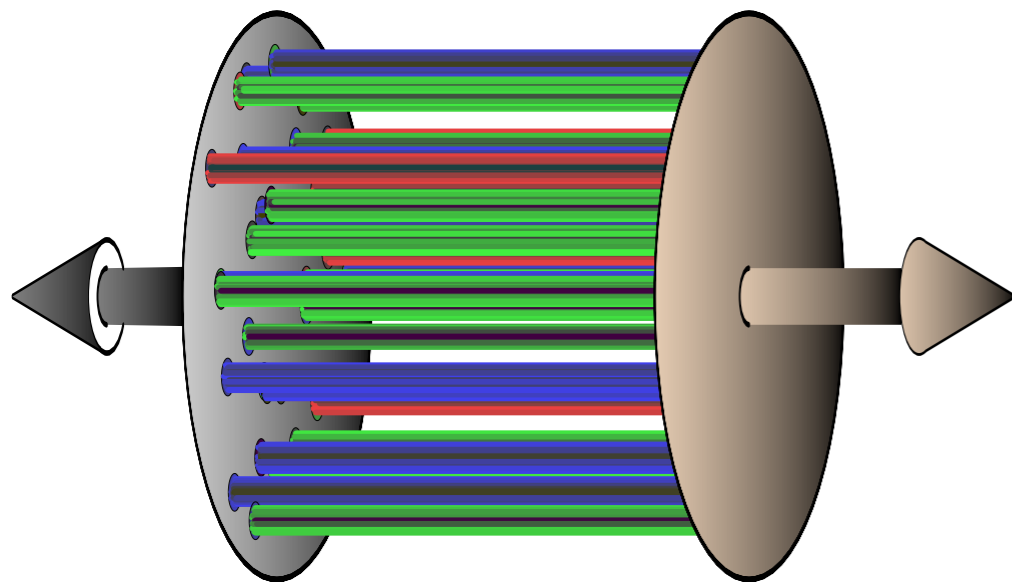
Lorentz contraction produces longitudinally-extended fluctuations



Gelis 1110.1544

- Pb nucleus=208 nucleons
- Each nucleon-nucleon collision deposits energy **at the transverse location** of the nucleons
- Thus the initial density profile is typically **uniform longitudinally**, but with a **bumpy transverse profile**

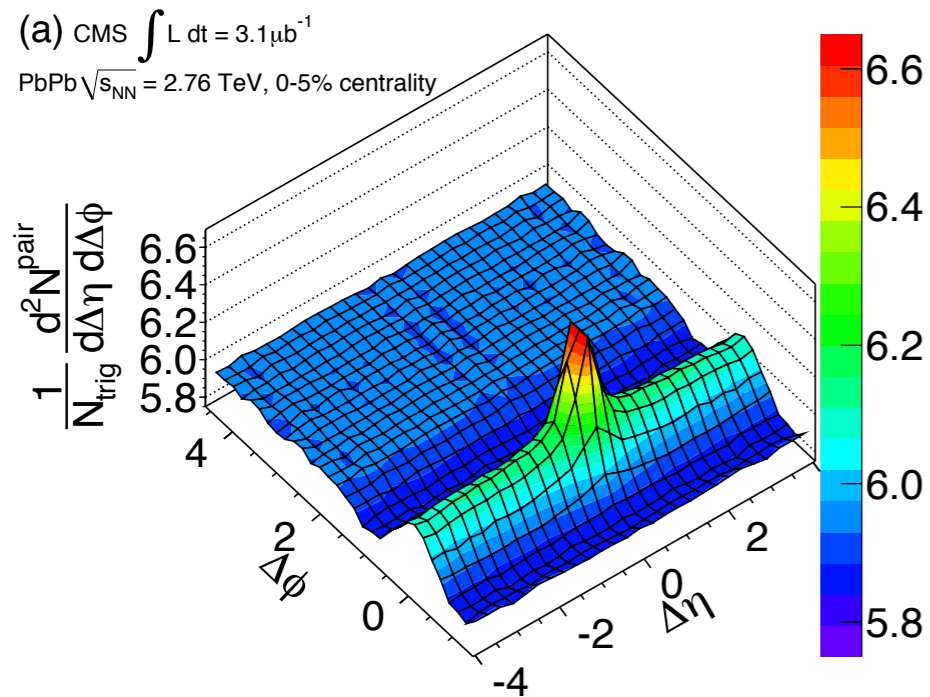
The fluid and its symmetries



- The fluid has the **same symmetry** as the initial state: **longitudinally invariant**, with **large transverse fluctuations**.
- Transverse velocity is generated by **expansion into the vacuum**
- The transverse velocity of the fluid **depends on the azimuthal direction φ** due to transverse fluctuations

Fluid to particles

- Eventually the fluid freezes and transforms into **independent** individual particles
- Particle velocity = fluid velocity, plus (small) thermal motion
- Some azimuthal directions φ have more/faster particles than others due to fluctuations.
- More particles \Rightarrow more pairs
- **Explains the observed peak near $\Delta\varphi=0$, independent of longitudinal separation $\Delta\eta$**



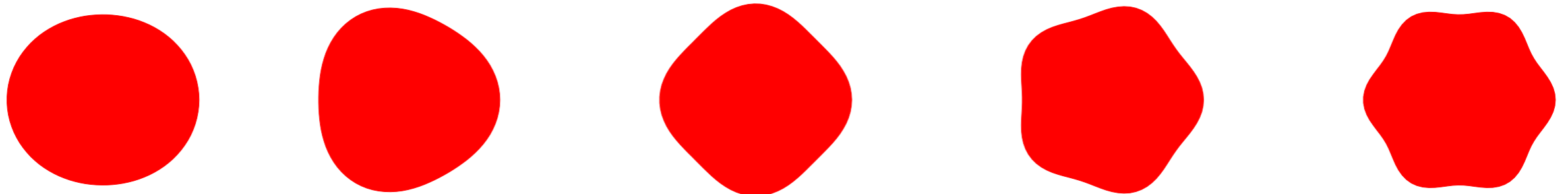
Understanding the *ripple in a pond*

- **Independent** particle emission from a thermal fluid
- + **fluctuations** in the initial state
- \Rightarrow **Correlations** between outgoing particles

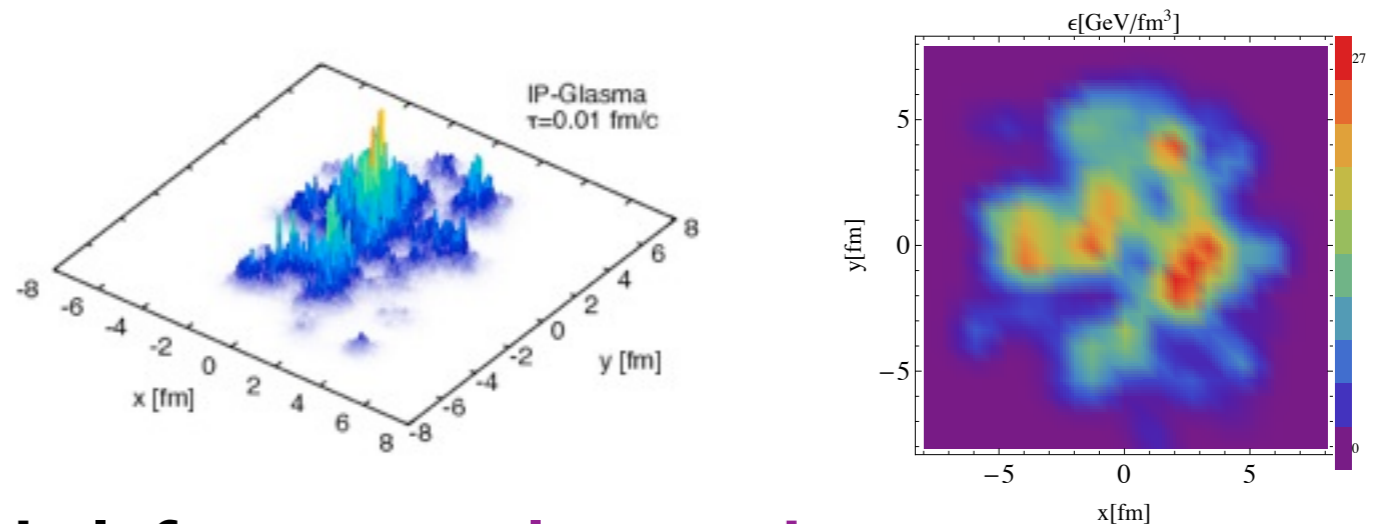
Flow harmonics

Hydrodynamics: information about flow is captured by the **single**-particle azimuthal (ϕ) distribution and its *complex* Fourier **harmonics**

$$dN/d\phi \sim \sum_n V_n e^{-in\phi}$$



From initial stage to flow



First strategy:

- take a particular model for **initial conditions**
- solve viscous hydro with these **initial conditions**

Difficulties:

- Numerically heavy (solve hydro **event by event**)
- Hydro=theory for long-wavelength modes, **not reliable for short-scale structures**

From initial stage to flow

*Alternative strategy: **symmetry** arguments*

- V_n is determined by a **property** of initial conditions which is invariant under **translation** and $2\pi/n$ **rotation**
- Construct smooth initial conditions with the same **property**
- Solve hydro with these smooth initial conditions

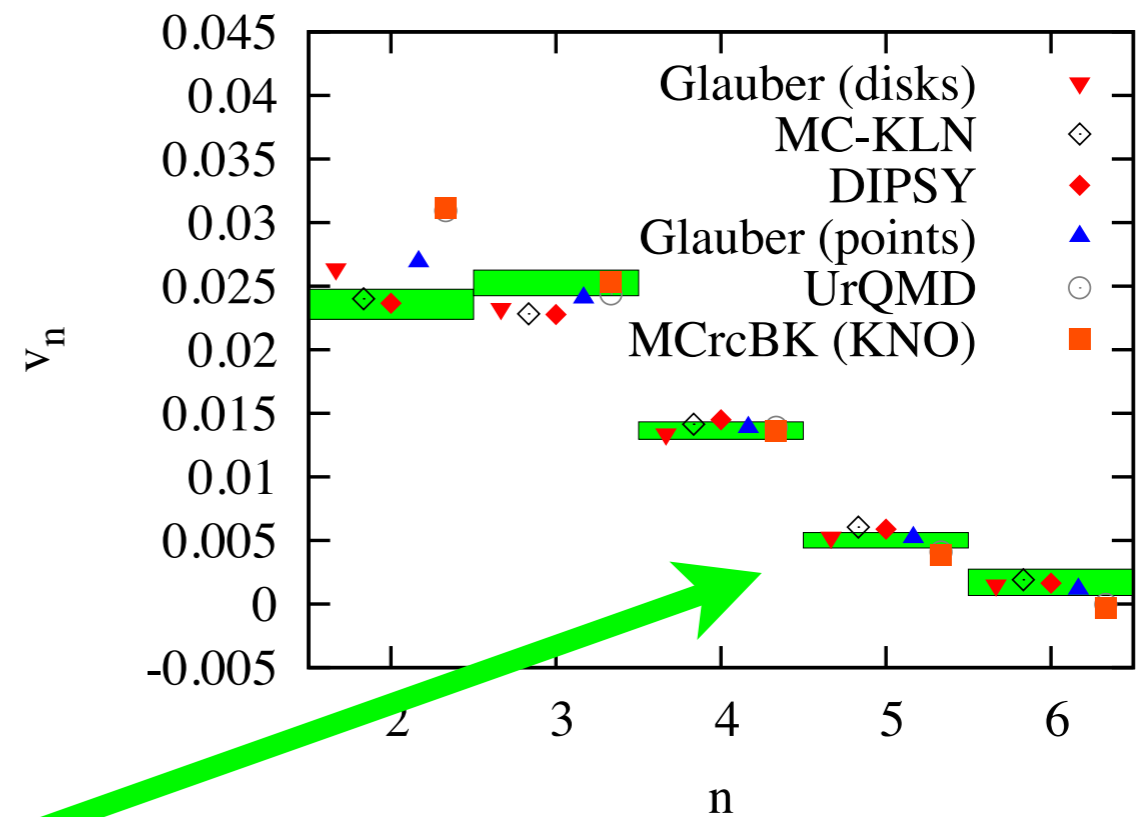
Example : Teaney and Yan's **cumulants** [1010.1876]

$\varepsilon_n \sim \langle r^n e^{in\varphi} \rangle$ (averaged over initial profile)

Central collisions: linear response

$V_n \propto \epsilon_n$ (magnitude & phase)
= **good** approximation to
event-by-event hydro,
better as viscosity increases
[talk by H. Niemi]

Within this approximation,
initial models+viscous hydro
describe ATLAS data



Luzum JYO 1210.6010

Flow observables

- Methods used to analyze flow were devised before the importance of flow fluctuations was fully recognized (2005)
- V_n fluctuates event-to-event
- All observables are averaged over events
- Which properties of the distribution of V_n can we measure?

Good flow observables

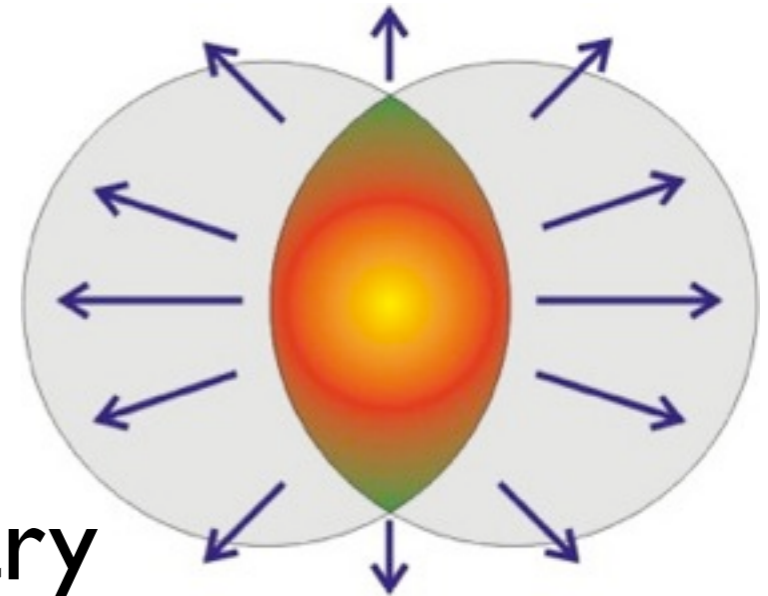
- Perform Fourier decomposition on an **event-by-event** basis $Q_n = (1/N) \sum e^{in\phi}$
- **Independent particle emission** \Rightarrow
average of **product** = product of averages
for a fixed underlying flow.
- Good observables are **products** of Q_n vectors.

Good flow observables

- $\langle Q_n Q_{-n} \rangle = \langle |V_n|^2 \rangle \rightarrow \text{rms } v_n$
- More generally, all **even** moments of the v_n distribution
- Mixed harmonics $\langle (Q_2)^2 Q_{-4} \rangle = \langle (V_2)^2 V_{-4} \rangle$
involve angles *and magnitudes* of V_2 and V_4 .
- **Not directly measurable: mean v_n , correlations between event plane angles.**

Luzum JYO 1209.2323

Non-central collisions



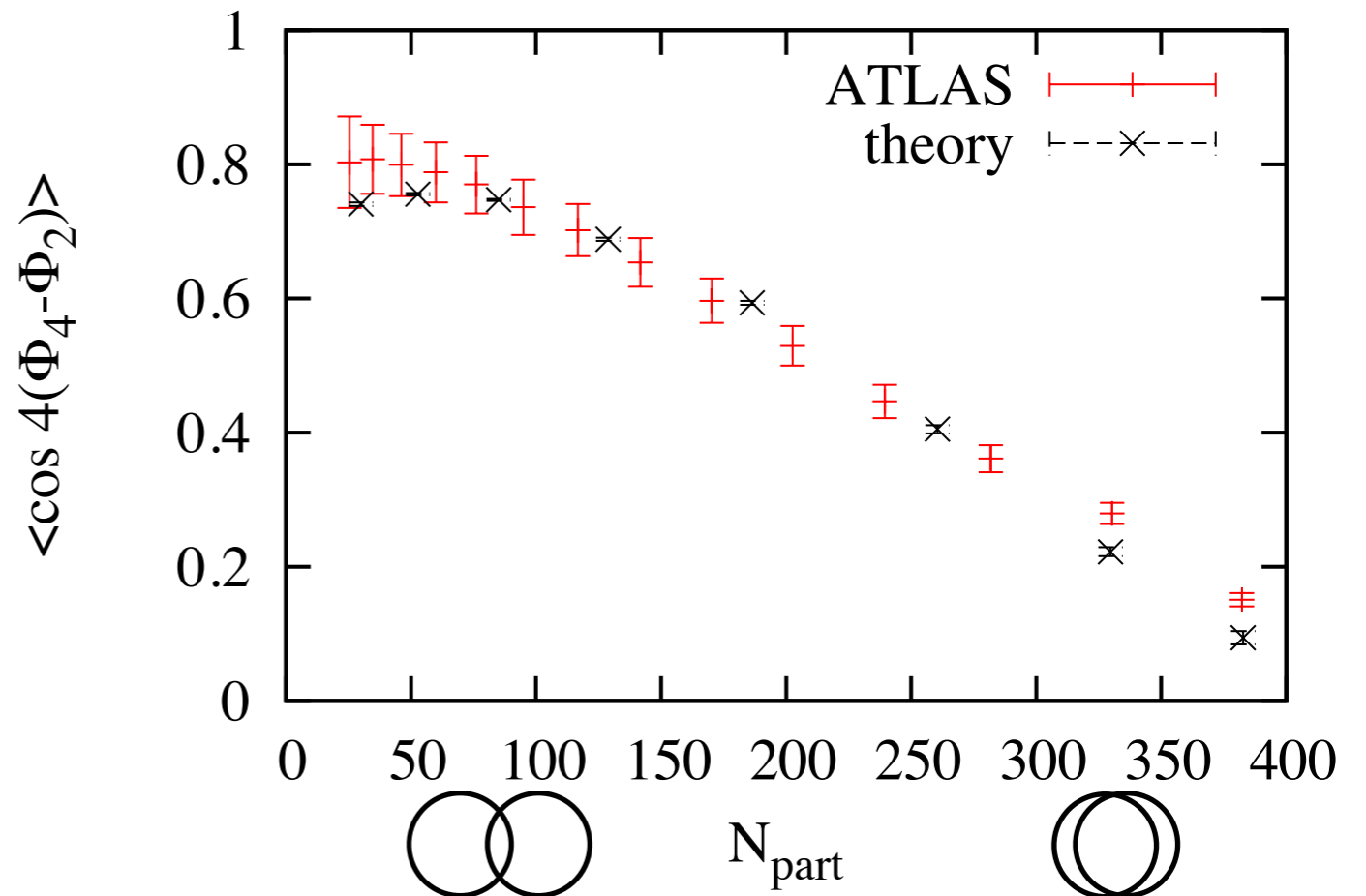
- Large **elliptic** anisotropy ϵ_2 due to average collision geometry
- Since ϵ_2 is large, nonlinear effects involving ϵ_2 are important
- $V_4 \propto (\epsilon_2)^2$ is a better rule than $V_4 \propto \epsilon_4$

Event plane correlators

ATLAS has measured
14 mixed correlations.

e.g., linear correlation
between V_4 and $(V_2)^2$

=Large when non-linear
response dominates



ATLAS 1208.1427

Bhalerao JYO Pal 1307.0980

Summary

- Flow fluctuations: a rich, new subfield of heavy-ion physics
- Hydrodynamic response to initial fluctuations not fully understood: effects of small-scale fluctuations + thermalization?
- Flow methods must be revisited to take flow fluctuations into account: new analyses needed.