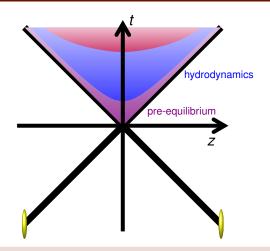
WHAT CAN WE LEARN ABOUT THE INITIAL STAGES FROM FLOW MEASUREMENTS?

Matthew Luzum

McGill University / Lawrence Berkeley National Laboratory

International Conference on the Initial Stages in High-Energy Nuclear Collisions 11 September, 2013

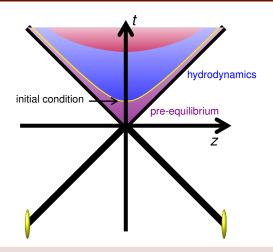
HYDRODYNAMIC EVOLUTION



• Standard picture of A-A collision: system evolves as a fluid

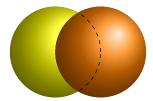
• Depends on initial $T^{\mu
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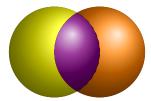
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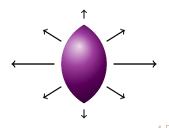


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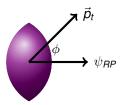






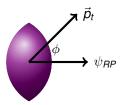
FLOW

$$rac{2\pi}{N}rac{dN}{d\phi}=\sum_n V_n e^{-in\phi}$$



FLOW

$$\frac{2\pi}{N}\frac{dN}{d\phi} = \sum_{n} V_{n}e^{-in\phi} = \sum_{n} v_{n}e^{in\Psi_{n}}e^{-in\phi}$$



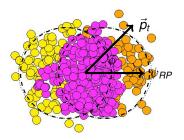
FLOW

Asymmetric pressure gradients \rightarrow anisotropic momentum distribution

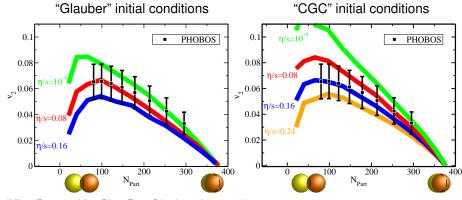
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FLUCTUATIONS ARE IMPORTANT!

- Not symmetric
- Varies from event to event



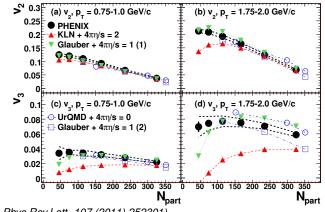
HISTORY: ELLIPTIC FLOW



(ML & Romatschke, Phys.Rev. C78 (2008) 034915)

- Depends on both medium properties (η/s) and initial conditions
- Need more information about one to constrain the other

NEW OBSERVABLES



(PHENIX, Phys.Rev.Lett. 107 (2011) 252301)

- New observables more constraints
- E.g., some initial conditions incompatible with (v₂, v₃)

- Usual procedure:
 - Choose initial state model
 - 2 Evolve with hydro for many events
 - Ompare to data.
 - Goto step 1
- Better: find constraints that can be checked without running hydro
- Note: in each event:

$$\frac{dN}{d\phi} = \mathcal{F}\left(T^{\mu\nu}(\mathbf{x})\right)$$

• Need to understand hydro response ${\cal F}$

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Idea: (Teaney, Yan; Phys.Rev. C83 (2011) 064904)

Characterize density by moments (or cumulants) of 2-D Fourier transform

$$ho(\mathbf{k}) = \int d^2 x
ho(\mathbf{x}) e^{i \mathbf{k} \cdot \mathbf{x}}$$

- (Small *m* = small power of *k* in Taylor series)
- Can write complete hydro response as set of functions

$$v_n e^{in\Psi_n} = f(W_{m,n})$$

• Can use other bases (e.g., Bessel expansion of k modes) (Coleman-Smith, Petersen, Wolpert; J.Phys. G40 (2013) 095103

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- Must make quantities with correct symmetries out of $W_{m,n} \sim \langle r^m e^{in\phi} \rangle$
- If anisotropies are small, can arrange in Taylor series. E.g., to first order:

$$V_n \equiv \mathbf{v}_n \mathrm{e}^{in \Psi_n} = \sum_{p=0}^{\infty} C_{n+2p,n} W_{n+2,p}$$

- If hydro is sensitive to large scale structure, first terms are most important (smaller powers of *k* in Fourier transform).
- If single term sufficient:

$$V_n e^{in\Psi_n} = -C \frac{\langle r^n e^{in\phi} \rangle}{\langle r^n \rangle} \equiv C \varepsilon_n e^{in\Phi_n}$$

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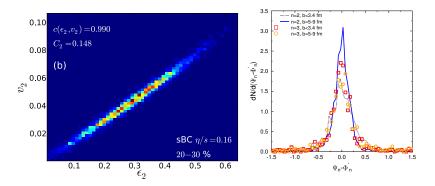
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 $V_n \propto \varepsilon_n$

Relation works well for $n \leq 3$:

$$v_n e^{in\Psi_n} = C \varepsilon_n e^{in\Phi_n}$$



(Niemi, Denicol, Holopainen, Huovinen; arXiv:1212.1008)

(Petersen, Qin, Bass, Muller, Phys. Rev. C 82, 041901 (2010))

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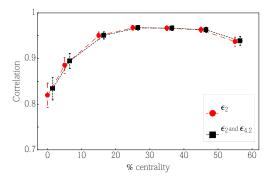
HYDRO RESPONSE

• Hydro insensitive to higher cumulants: including several terms,

$$v_2 e^{i2\Psi_2} = C \langle r^2 e^{i2\Phi_2} \rangle + C' \langle r^4 e^{i2\Phi_2} \rangle + \dots$$

is no better than just lowest term

• Nonlinear terms are important for $n \ge 4$



(Gardim, Grassi, ML, Ollitrault; Phys.Rev. C85 (2012) 024908)

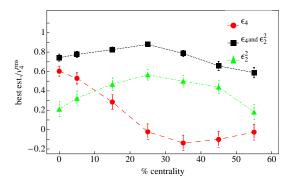
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• Nonlinear terms are important for $n \ge 4$

To a good approximation:

$$v_{1}e^{i\Psi_{1}} = C_{1} \varepsilon_{1}e^{i\Phi_{1}}$$

$$v_{2}e^{i2\Psi_{2}} = C_{2} \varepsilon_{2}e^{i2\Phi_{2}}$$

$$v_{3}e^{i3\Psi_{3}} = C_{3} \varepsilon_{3}e^{i3\Phi_{3}}$$

$$v_{4}e^{i4\Psi_{4}} = C_{4,1} \varepsilon_{4}e^{i4\Phi_{4}} + C_{4,2}\varepsilon_{2}^{2}e^{i4\Phi_{2}}$$

$$v_{5}e^{i5\Psi_{5}} = C_{5,1} \varepsilon_{5}e^{i5\Phi_{5}} + C_{5,2}\varepsilon_{2}\varepsilon_{3}e^{i(2\Phi_{2}+3\Phi_{3})}$$

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$$\begin{aligned} v_1 e^{i\Psi_1} &= C_1 \varepsilon_1 e^{i\Phi_1} \\ v_2 e^{i2\Psi_2} &= C_2 \varepsilon_2 e^{i2\Phi_2} \\ v_3 e^{i3\Psi_3} &= C_3 \varepsilon_3 e^{i3\Phi_3} \\ v_4 e^{i4\Psi_4} &= C_{4,1} \varepsilon_4 e^{i4\Phi_4} + C_{4,2} \varepsilon_2^2 e^{i4\Phi_2} \\ v_5 e^{i5\Psi_5} &= C_{5,1} \varepsilon_5 e^{i5\Phi_5} + C_{5,2} \varepsilon_2 \varepsilon_3 e^{i(2\Phi_2 + 3\Phi_3)} \end{aligned}$$

Can use these relations to derive constraints for initial state:

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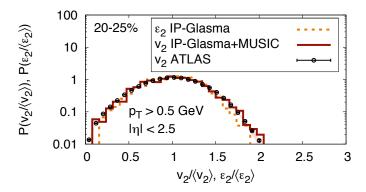
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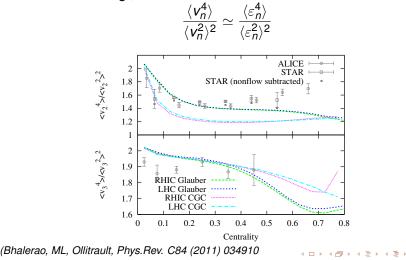


(Gale, Jeon, Schenke, Tribedy, Venugopalan, Nucl.Phys.A904-905 2013 (2013) 409c-412c)

- For $n \leq 3$, normalized v_n distribution same as ε_n distribution
- ullet \Longrightarrow can compare initial state model directly to data

Ratios of moments of v_n , Ψ_n distributions

Can take ratios to generate single numbers that can be compared to measurements. E.g.,



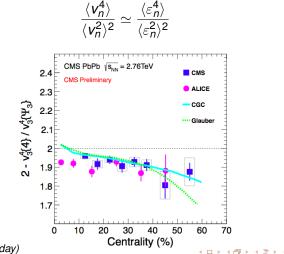
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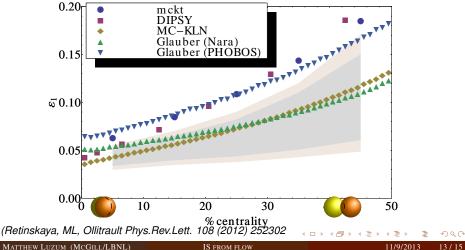
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(Sanders talk on Monday) MATTHEW LUZUM (MCGILL/LBNL)

PRECISION CONSTRAINTS: ε_1

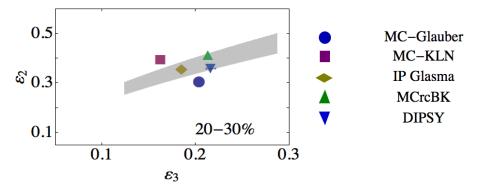
Can also constrain magnitude of lowest moments of ε_n distribution. E.g., allowed values of $\sqrt{\langle \varepsilon_1^2 \rangle}$:



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PRECISION CONSTRAINTS: $\varepsilon_2, \varepsilon_3$

Can also constrain magnitude of lowest moments of ε_n distribution. E.g., allowed values of $\langle \varepsilon_2^2 \rangle / \langle \varepsilon_3^2 \rangle^{0.6}$:



(Retinskaya, ML, Ollitrault — See talk on Friday)

- With the many new (and forthcoming) flow measurements, we can directly constrain theory — both medium properties and the initial state
- Eccentricities ε_n represent the lowest k mode of the Fourier transformed initial density
- The set of ε_n accurately predict flow in each event (not always linearly)
- medium response effectively codified in a few coefficients
- ⇒ can provide simple and direct quantitative constraints that must be satisfied for any model of the initial stages

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