

A state of the art I-QCD + HRG EOS: Full vs Partial Chemical Equilibrium

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Outline

- Motivation
- The conceptual setup
- Results

NB Work done in collaboration with M. Bluhm, P. Alba, W.M. Alberico and C. Ratti and found in [arXiv:1306.6188 \[hep-ph\]](https://arxiv.org/abs/1306.6188).

Motivations

The **QCD Equation of State** affects soft-physics observables

- *directly*, entering into the **hydrodynamic equations** describing the medium evolution

$$\partial_\mu T^{\mu\nu} = 0 \quad \text{with} \quad P = P(\epsilon)$$

- *indirectly*, entering into the description of the **transition from fluid to particles** through the Cooper-Frey formula

$$E \frac{dN_i}{d\mathbf{p}} = \frac{g_i}{(2\pi)^3} \int_{\Sigma_{fo}} p_\mu d\Sigma^\mu f_i(x, p),$$

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- The formula fixes unambiguously *both the shape* of the spectrum *and its normalization*;
- The particle distributions must be consistent with the ones entering into the hadronic EOS.

Hadron-Resonance Gas EOS

- A system of stable hadrons interacting through the exchange of resonances (e.g. $\pi + \pi \rightarrow \rho \rightarrow \pi + \pi$) can be described as an ideal gas of hadrons (e.g. π) and resonances (e.g. ρ)
- The partition function factorizes:

$$Z_{\text{GC}} = \prod_i Z_{\text{GC}}^{(i)},$$

where i is taken to run over all hadrons/resonances up to ~ 2 GeV.

- The hadronic phase is usually described by such an EOS, where

$$P = \sum_i g_i \int \frac{d\mathbf{p}}{(2\pi)^3} \frac{1}{E} \frac{\mathbf{p}^2}{3} f_{\text{eq}}^{(i)}(\mathbf{p}),$$

the other quantities following from thermodynamics:

$$s = (\partial P / \partial T)_\mu \quad \text{and} \quad \epsilon = Ts - P + \mu n$$

HRG-EOS: chemical composition

At least three different scenarios can be conceived:

- **Full chemical equilibrium:** the multiplicity of each H/R species follows from $f_i^{\text{eq}}(\mathbf{p})$. In the absence of stopping (LHC conditions!) $\langle B \rangle = \langle S \rangle = \langle I_3 \rangle = 0$ and all chemical potentials vanish. A **unique temperature** T_{fo} would describe the *slope of the spectra* and the *particle abundances*;

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- **Early chemical freeze-out:** the **multiplicity of all H/R species is fixed at chemical freeze-out** ($T_{\text{chem}} \approx 160$ MeV). In order to keep the chemical composition fixed until *kinetic freeze-out* for each hadron/resonance an independent **chemical potential** is introduced:

$$f_i(\mathbf{p}) = \frac{1}{e^{(E_p - \mu_i)/T} \pm 1}$$

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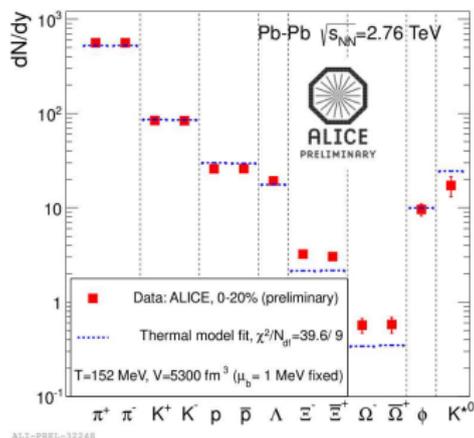
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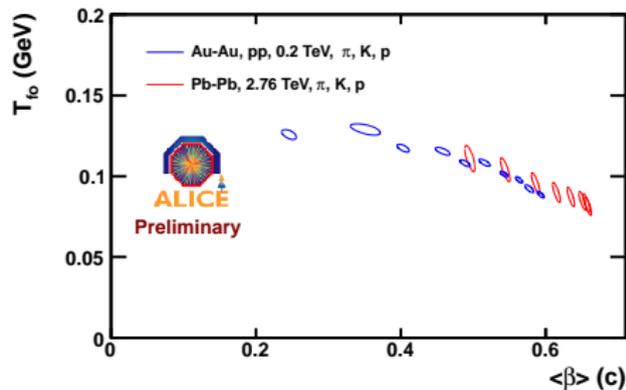
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- **Partial chemical equilibrium:** H/R abundances are initially fixed at T_{chem} , but reactions mediated by short-living resonances (e.g. $\pi + p \rightarrow \Delta \rightarrow \pi + p$) are allowed.

Chemical vs kinetic freeze-out: experimental evidence



Particle yields consistent with thermal production at $T_{\text{chem}}^{\text{fo}} \approx 155$ MeV



Blast-wave fits of the p_T -spectra require $T_{\text{kin}}^{\text{fo}} \sim 100$ MeV

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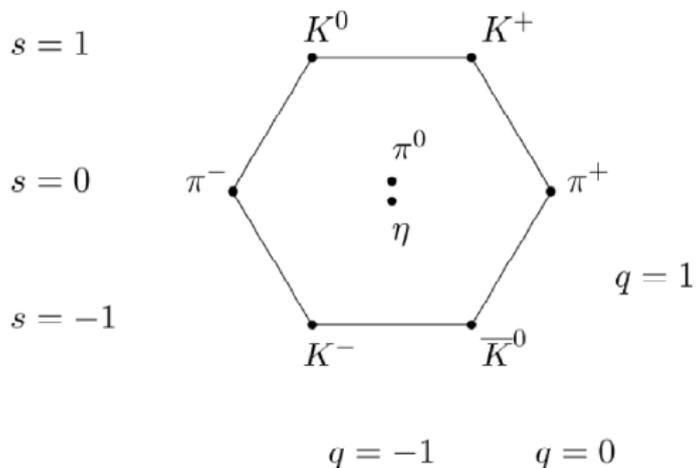
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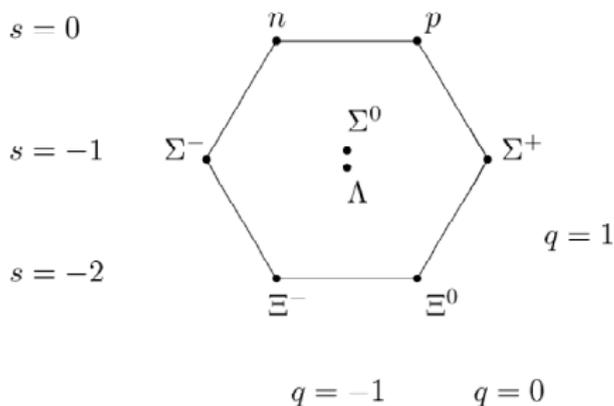
- The chemical potentials of all resonances can then be expressed in terms of the ones of the stable hadrons!

Stable hadrons



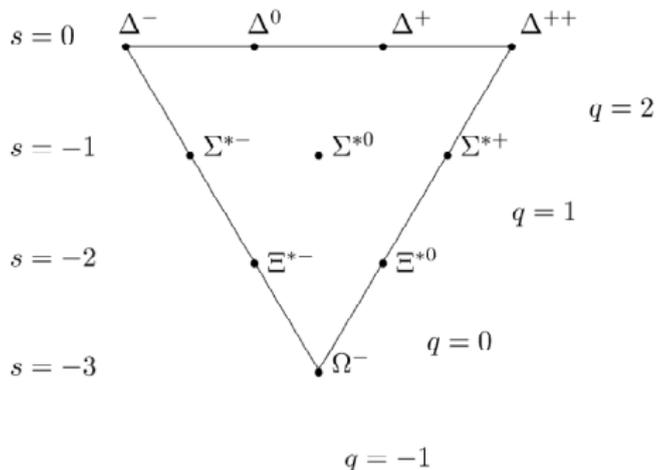
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- At equilibrium the Gibbs free-energy $\Phi \equiv E - TS + PV$ is stationary:

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- In the **decay** of P_n one has $\Delta N_n = -1 \rightarrow \Delta N_i = \sum_c \Gamma_c N_i^c$, so that

$$\mu_n = \sum_c \Gamma_c \sum_{i=1}^{n-1} N_i^c \mu_i$$

Equilibrium relations for chemical potentials: outcomes

- At the end the chemical potential of each resonance r is fixed by δ independent chemical potentials of stable hadrons

$$\mu_r = \sum_h \langle N_h^{(r)} \rangle \mu_h, \quad \text{with } h = \pi, K, \eta, N, \Lambda, \Sigma, \Xi, \Omega$$

- One obtains a table like (code, mass, spin-deg, baryon-num, $\{\mu_h\}$)

code	mass	spin-deg	baryon-num	μ_π	μ_K	μ_η	μ_N	μ_Λ	μ_Σ	μ_Ξ	μ_Ω
22	0.2	0	0	0	0	0	0	0	0	0	0
211	0.13957	1	0	1	0	0	0	0	0	0	0
111	0.13498	1	0	1	0	0	0	0	0	0	0
-211	0.13957	1	0	1	0	0	0	0	0	0	0
321	0.49368	1	0	1	0	0	0	0	0	0	0
-321	0.49368	1	0	1	0	0	0	0	0	0	0
311	0.49765	1	0	1	0	0	0	0	0	0	0
-311	0.49765	1	0	1	0	0	0	0	0	0	0
221	0.54775	1	0	0	1	0	0	0	0	0	0
213	0.7758	3	0	2	0	0	0	0	0	0	0
113	0.7758	3	0	2	0	0	0	0	0	0	0
-213	0.7758	3	0	2	0	0	0	0	0	0	0
223	0.78259	3	0	2	0	0	0	0	0	0	0
323	0.89166	3	0	1	1	0	0	0	0	0	0
-323	0.89166	3	0	1	1	0	0	0	0	0	0
313	0.8961	3	0	1	1	0	0	0	0	0	0
-313	0.8961	3	0	1	1	0	0	0	0	0	0
2212	0.93827	2	1	0	0	0	1	0	0	0	0
2112	0.93957	2	1	0	0	0	1	0	0	0	0
331	0.95778	1	0	1.98209	0	0.654	0	0	0	0	0
9010221	0.9741	1	0	1.56	0.44	0	0	0	0	0	0
9000211	0.9847	1	0	0.844	0.312	0.844	0	0	0	0	0
9000111	0.9847	1	0	0.844	0.312	0.844	0	0	0	0	0
-9000211	0.9847	1	0	0.844	0.312	0.844	0	0	0	0	0

EOS with partial chemical freeze-out (I)

- Thermodynamic relations

$$Ts = \epsilon + P - \sum_r \mu_r n_r \quad \text{and} \quad dP = s dT + \sum_r n_r d\mu_r$$

r running over *all the resonances* in the cocktail.

- Entropy and particle density: $s = (\partial P / \partial T)_{\mu_r}$ and $n_r = (\partial P / \partial \mu_r)_T$
- Chemical potentials of the resonances r expressed in terms of the ones of the “stable” hadrons h :

$$\mu_r = \sum_h \langle N_h^{(r)} \rangle \mu_h,$$

- Density of hadrons h

$$\bar{n}_h \equiv \left(\frac{\partial P}{\partial \mu_h} \right)_T = \sum_r \left(\frac{\partial P}{\partial \mu_r} \right)_T \frac{d\mu_r}{d\mu_h} = \sum_r n_r \langle N_h^{(r)} \rangle :$$

“primary” (n_h) + “feed-down” ($\sum_{r \neq h} n_r \langle N_h^{(r)} \rangle$) contributions.

EOS with PCE (II): entropy conservation

The conservation of \bar{N}_h entails the conservation of entropy! In fact...

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$$\partial_\mu T^{\mu\nu} = \partial_\mu [(\epsilon + P)u^\mu u^\nu - g^{\mu\nu}P] = 0$$

together with the thermodynamic relations

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lead to

$$\partial_\mu [su^\mu] = - \sum_h \frac{\mu_h}{T} \partial_\mu [\bar{n}_h u^\mu]$$

In spite of partial deviations from full CE, the conservation of the “effective charges” \bar{N}_h leads to the conservation of the entropy S (thanks to P. Huovinen for discussions)

EOS with PCE (III)

- ...Hence the recipe to impose **partial chemical freeze-out** at T_c :

$$\frac{\bar{n}_h(T, \{\mu_{h'}\})}{s(T, \{\mu_{h'}\})} = \frac{\bar{n}_h(T_c, \{\mu_{h'}=0\})}{s(T_c, \{\mu_{h'}=0\})}$$

expressing the conservation of entropy and of the “charges” \bar{N}_h ;

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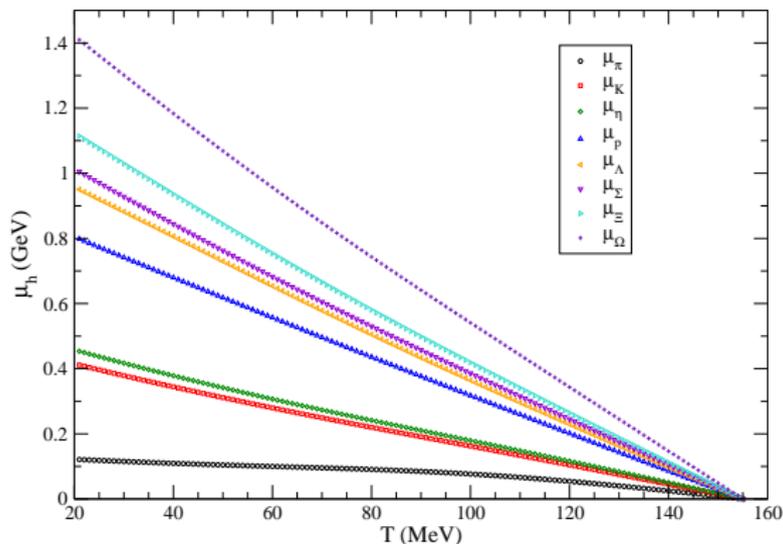
- **Particle ratios** are fixed at T_c :

$$\frac{\bar{n}_{h_1}(T, \{\mu_{h'}\})}{\bar{n}_{h_2}(T, \{\mu_{h'}\})} = \frac{\bar{n}_{h_1}(T_c, \{\mu_{h'}=0\})}{\bar{n}_{h_2}(T_c, \{\mu_{h'}=0\})}$$

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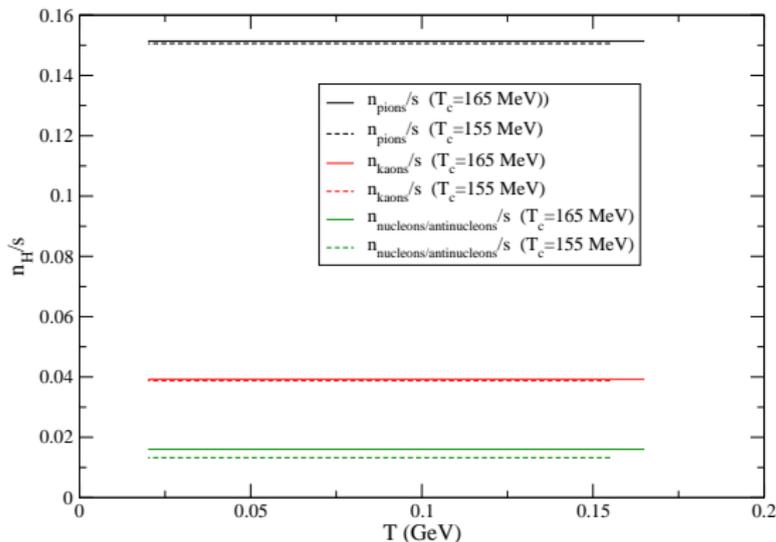
- T. Hirano and K. Tsuda, PRC 66 (2002), 054905;
- H. Bebie et al., NPB 378 (1992), 95;
- R. Rapp, PRC 96 (2002), 017901;
- P.F. Kolb and R. Rapp, PRC 67 (2003), 044903.

HRG with PCE: some results



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The full setup: I-QCD + HRG

In order to have an EOS covering the full range of temperatures of interest for the experiment and for hydro simulations...

- For high temperatures ($T > 172$ MeV) continuum-extrapolated lattice-QCD results with realistic quark masses by the Wuppertal-Budapest collaboration¹ were employed;

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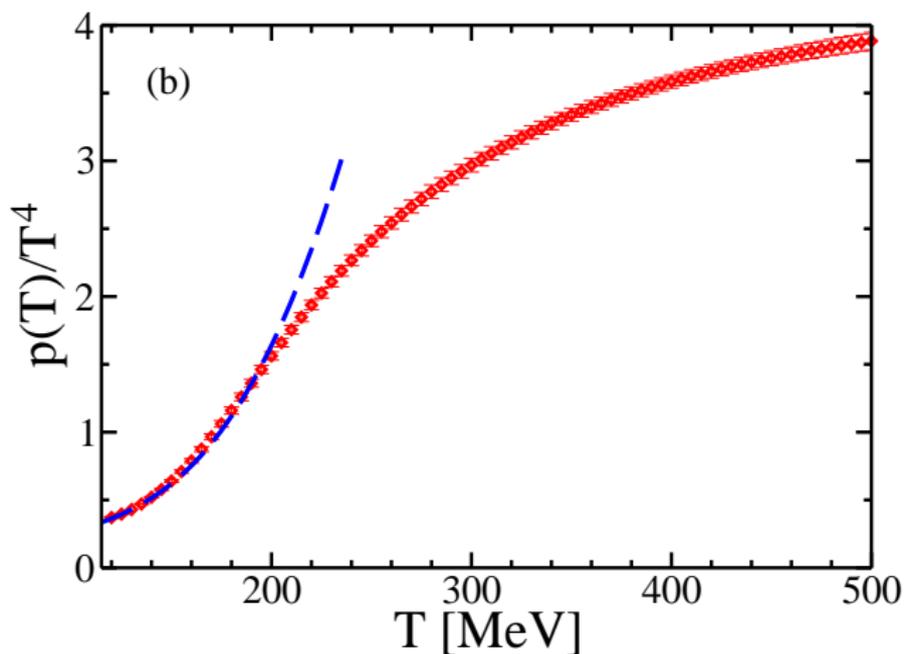
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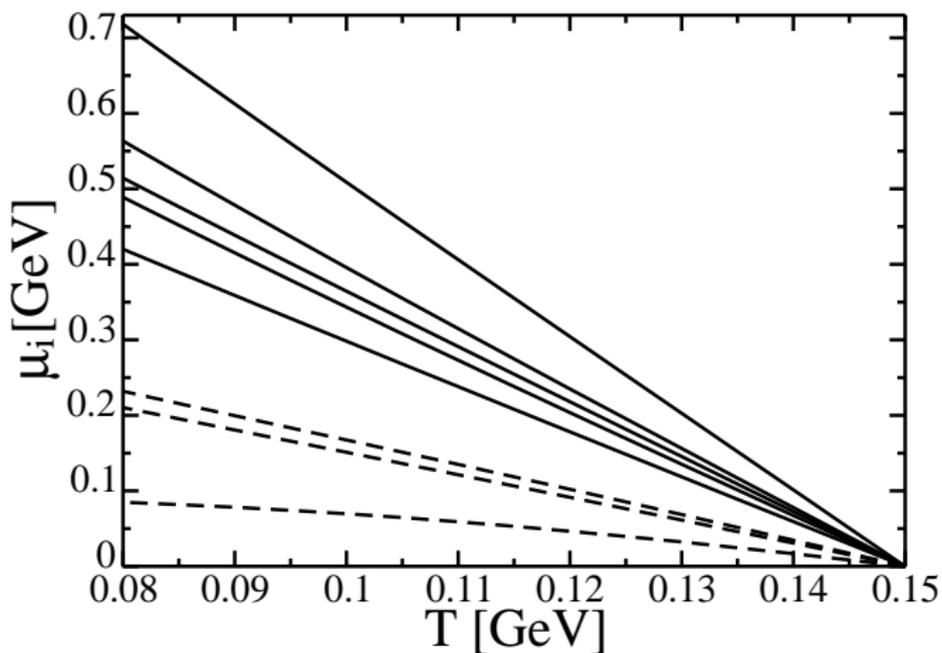
NB in view of a generalization to the finite-density case we have treated each isospin state (e.g. π^+ , π^0 , π^- ...) separately.

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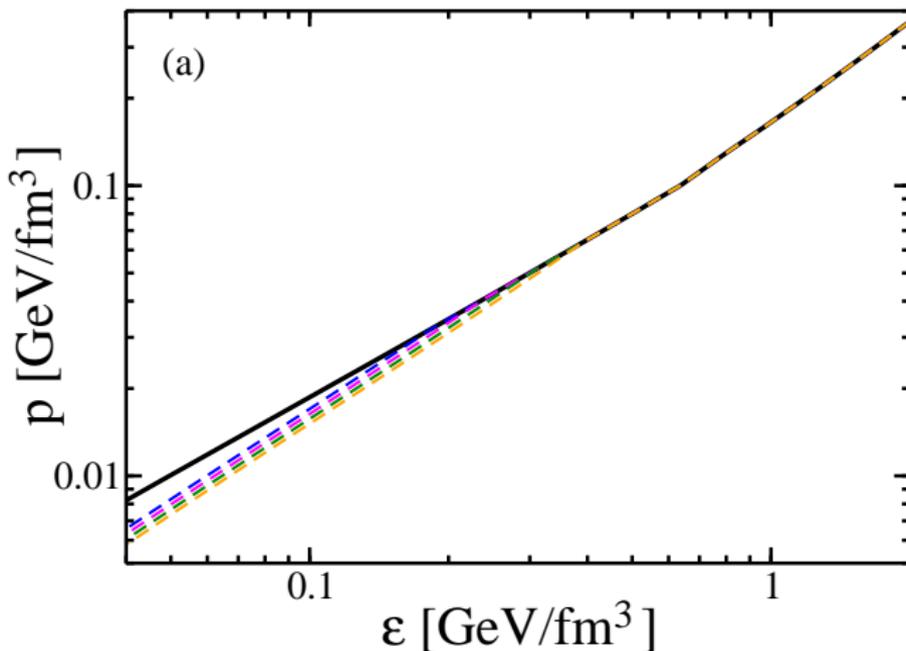
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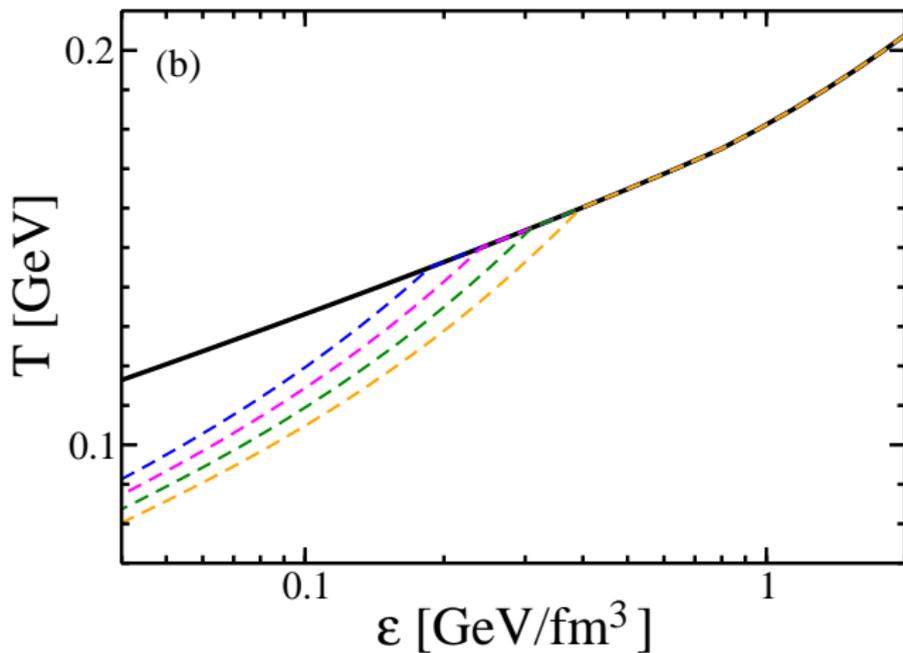
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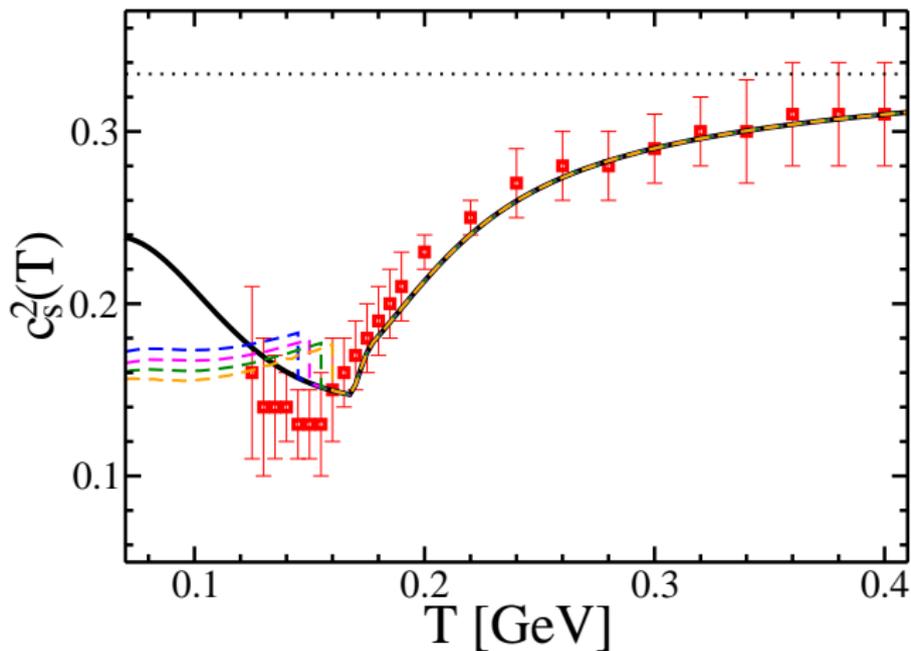
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- For the future we plan to extend our approach to the finite-density case, of interest for the low-energy scan at RHIC and for FAIR.