

three-body HEAVY meson decays:  
final state interactions

M. R. Robilotta





three-body HEAVY meson decays:

final state interactions

final state interactions

**THEORY**

weak vertex

final state  
interactions

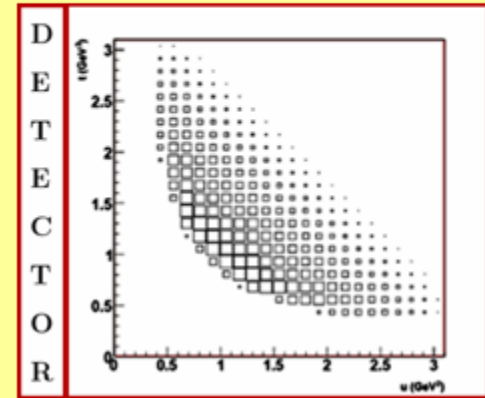
Towards three-body unitarity in  $D^+ \rightarrow K^- \pi^+ \pi^+$

research programme

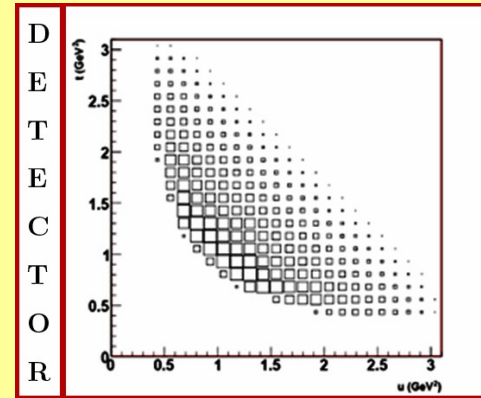
recent progress



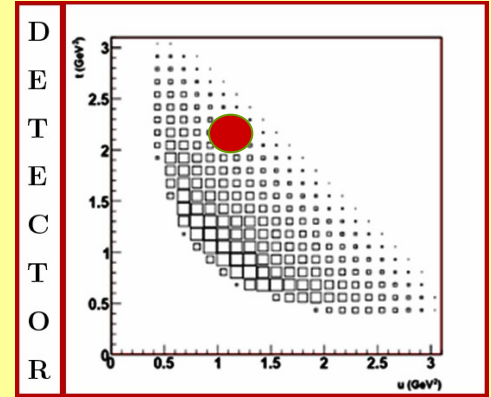
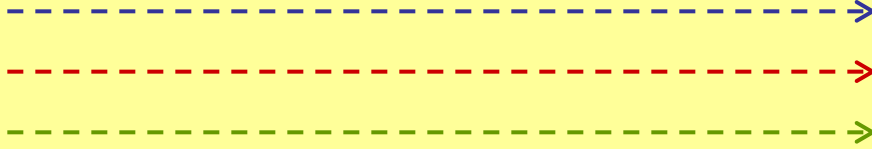
# final state interactions



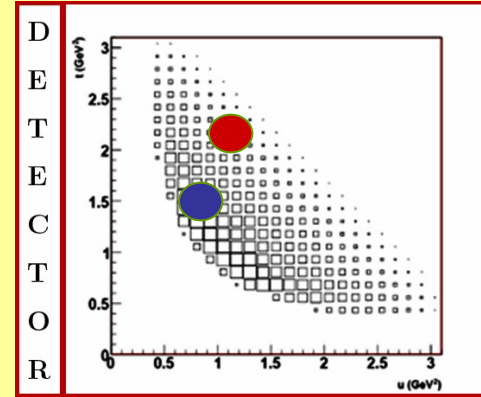
# final state interactions



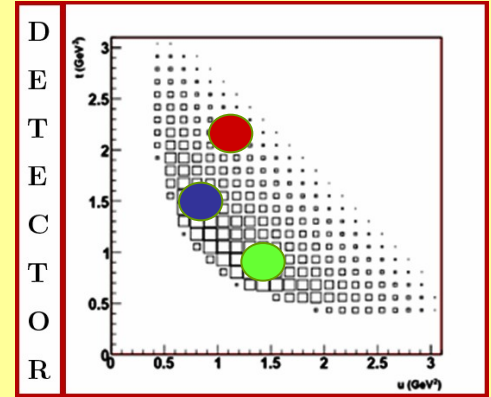
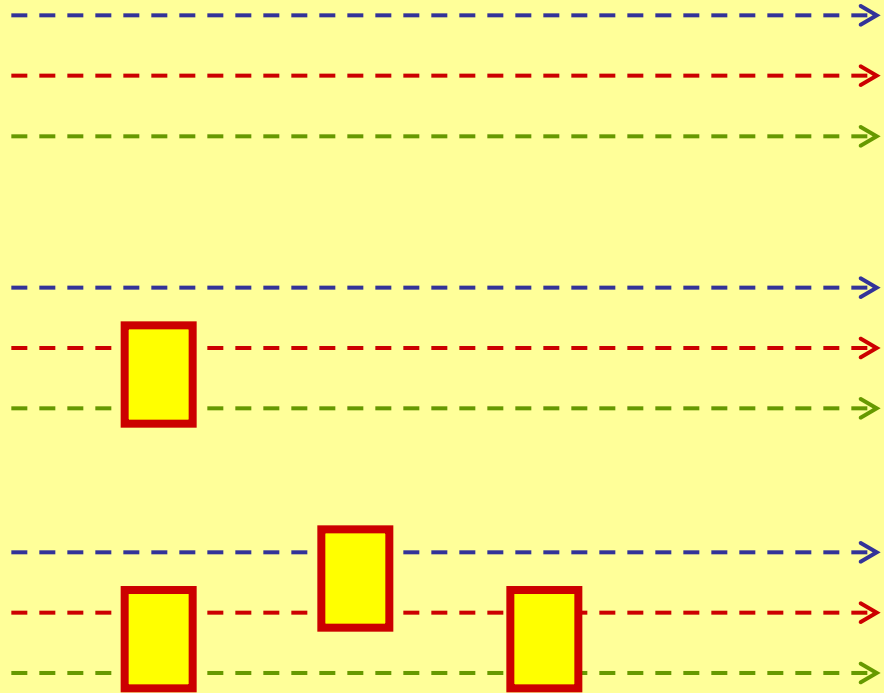
# final state interactions



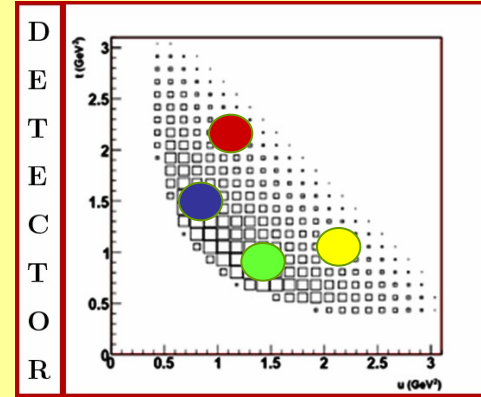
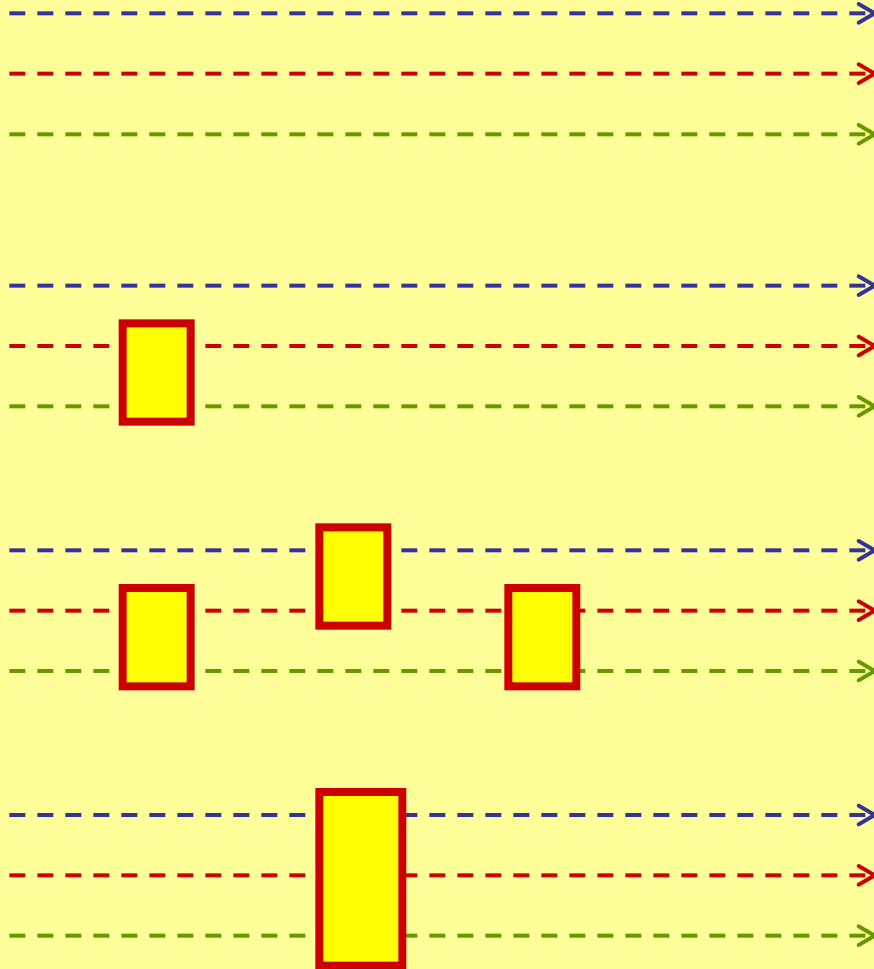
# final state interactions



# final state interactions

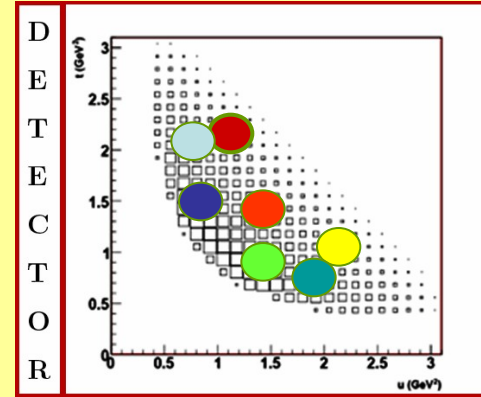
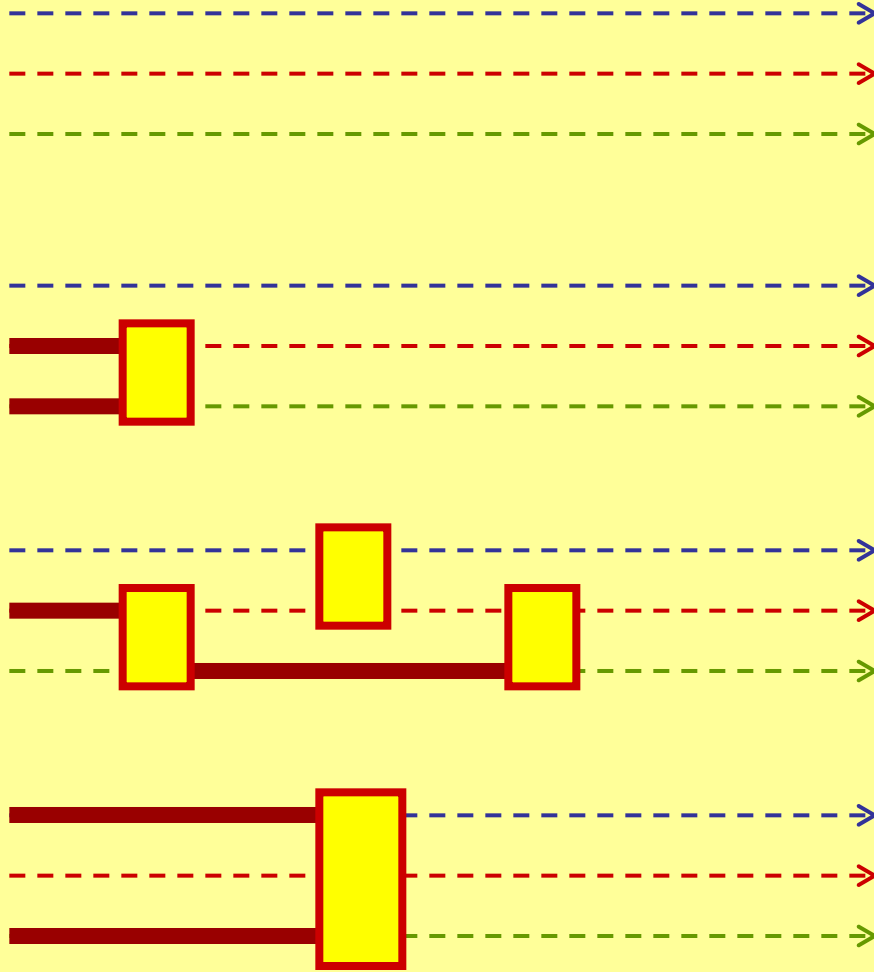


# final state interactions

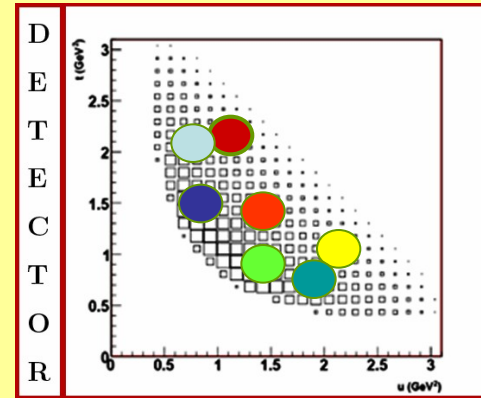
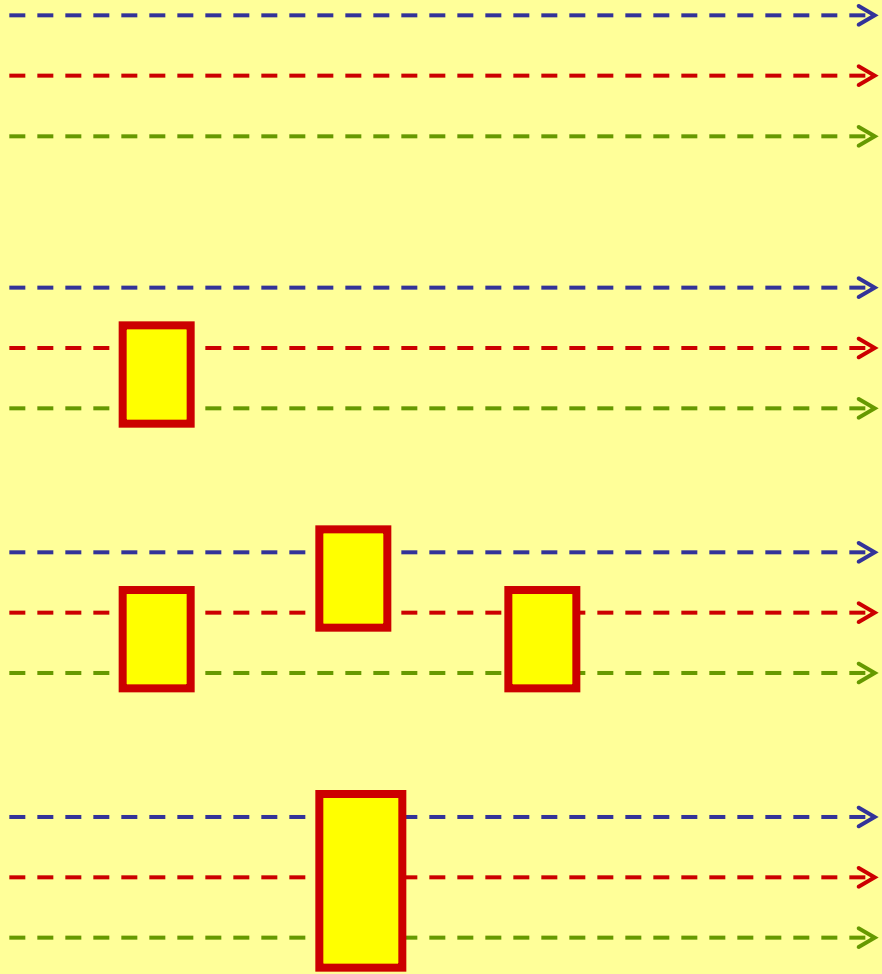




# final state interactions



# final state interactions

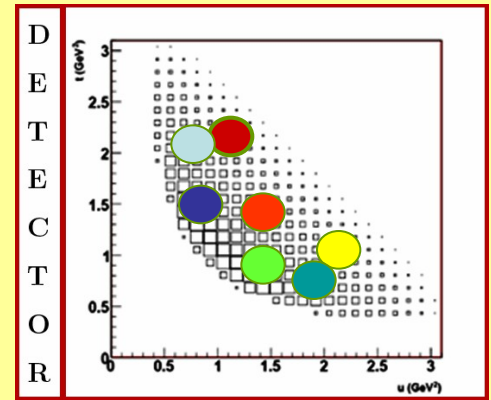
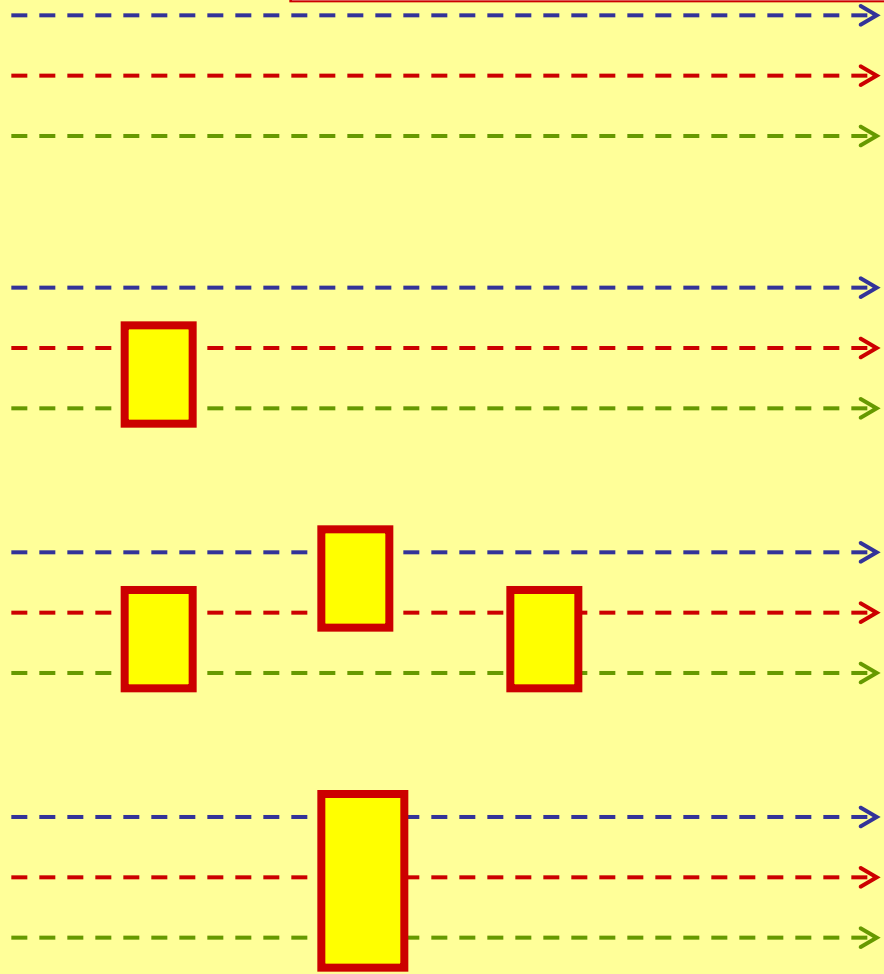


GUESS  
this  
MESS ?



# THEORY

general guides ?

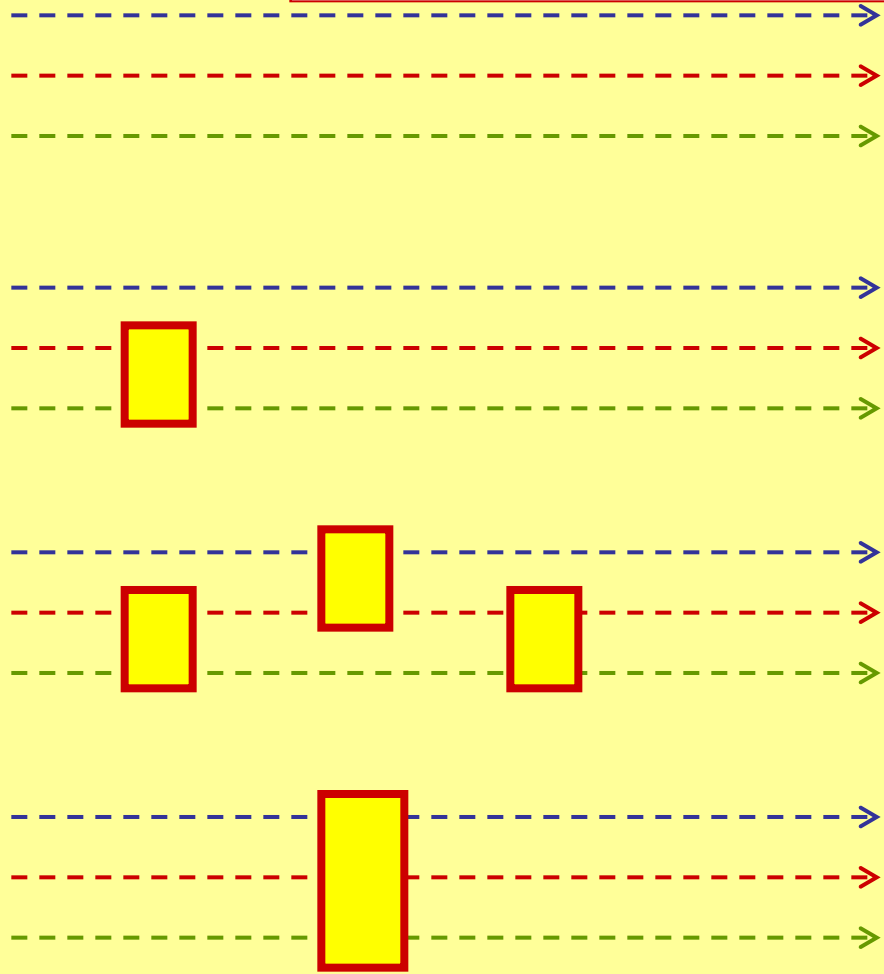


GUESS  
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# THEORY

general guides ?

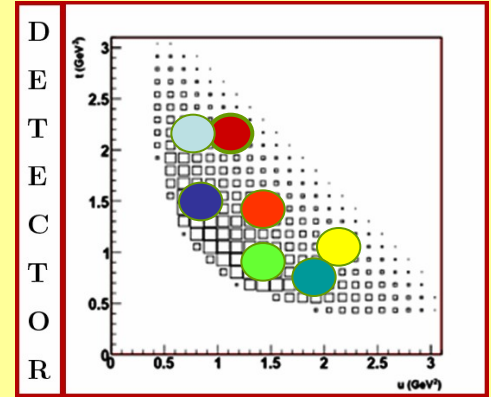


atomic  
nuclear  
physics

methods

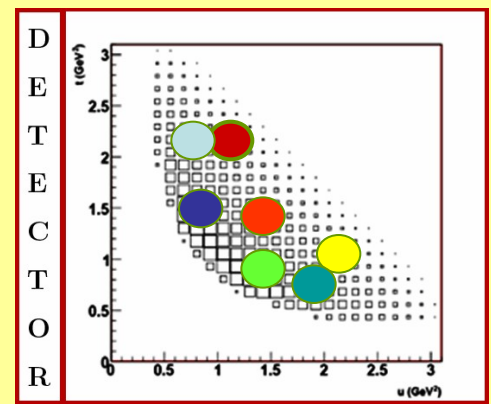
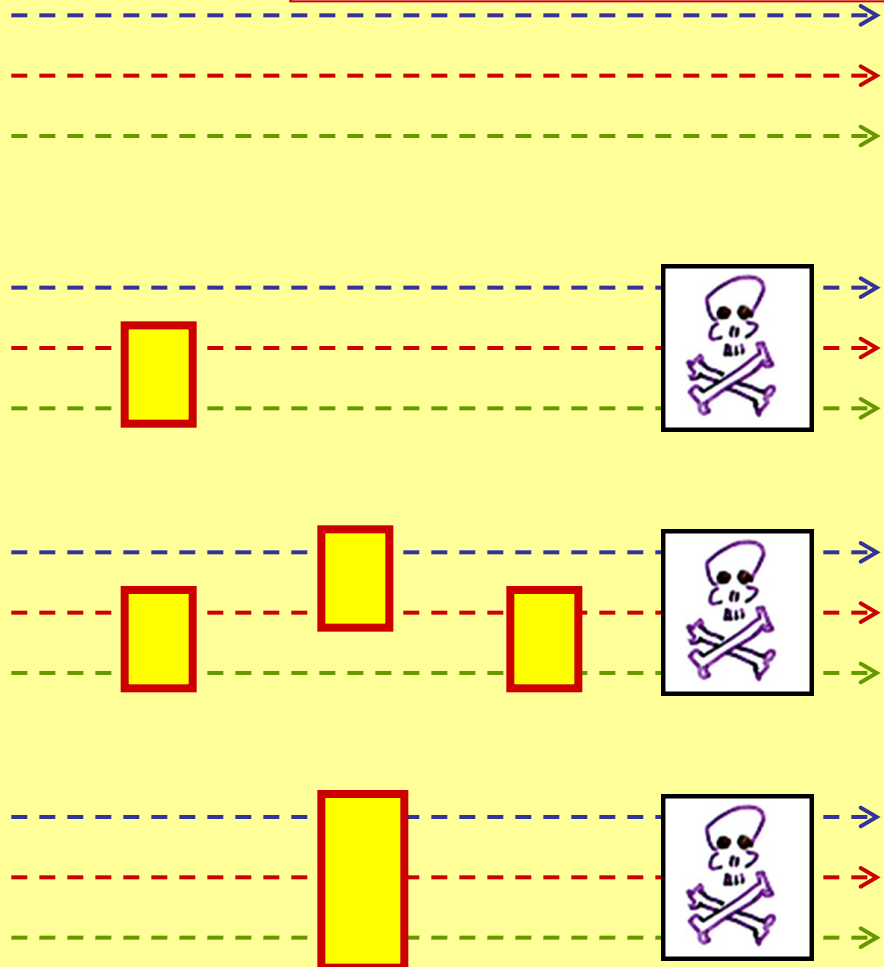
Faddeev  
series

energy



# THEORY

general guides ?



atomic  
nuclear  
physics

methods

Faddeev  
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energy

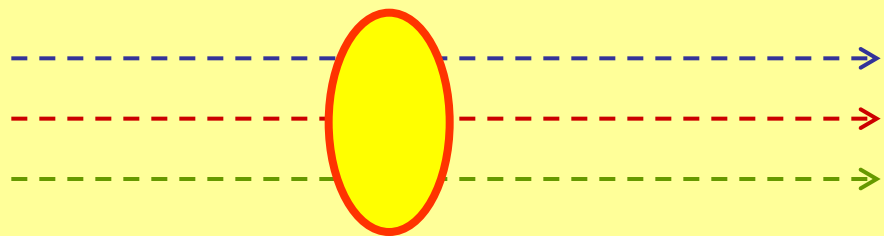
Watson's  
theorem

# THEORY

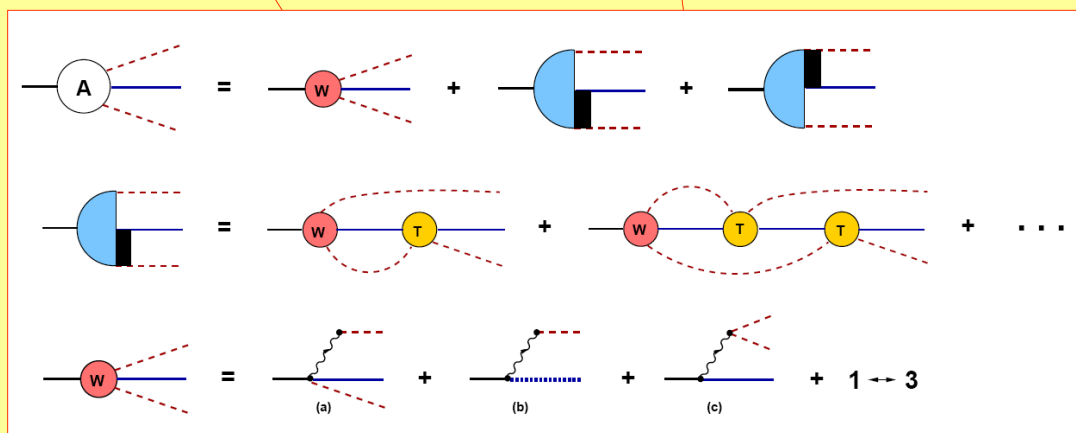
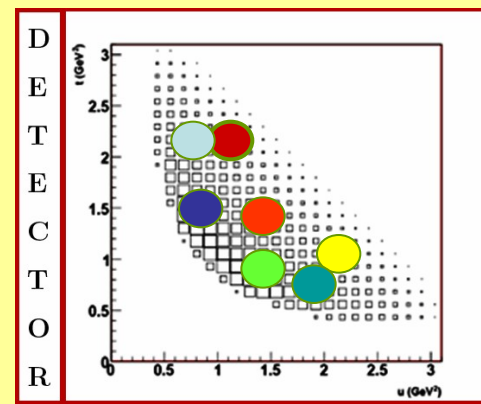
two problems



weak vertex



final state interactions



# THEORY

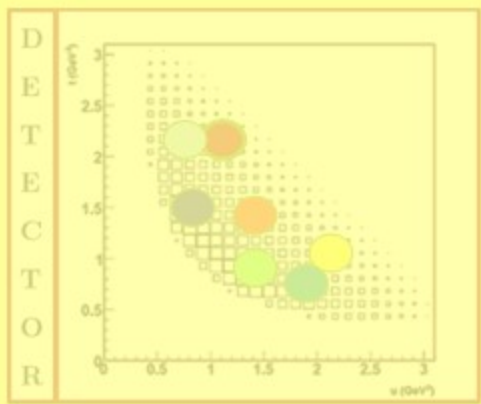
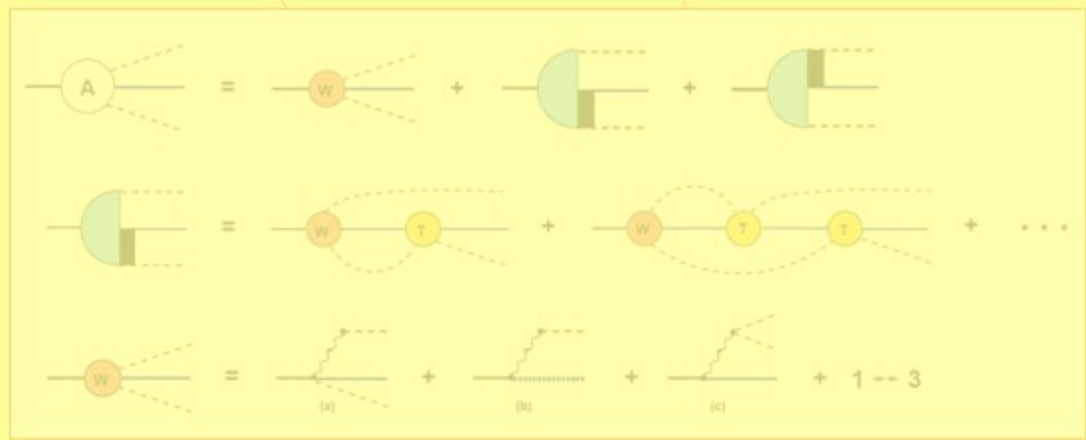
two problems



$$D^+ \rightarrow K^- \pi^+ \pi^+$$

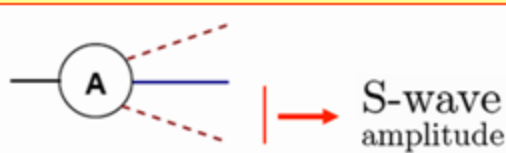
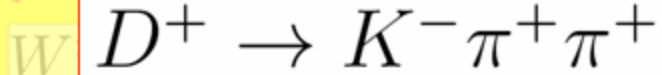
weak vertex

final state interactions



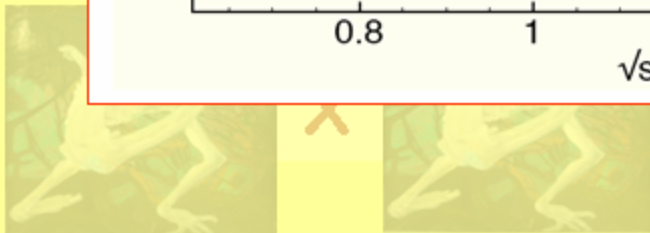
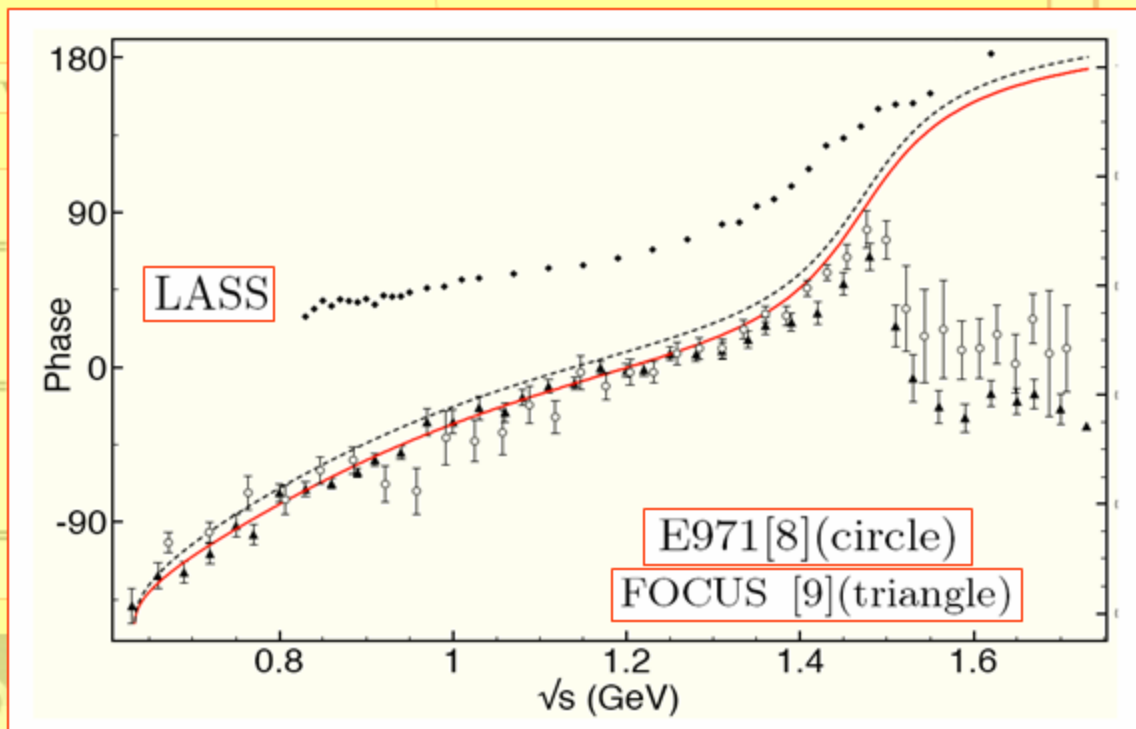
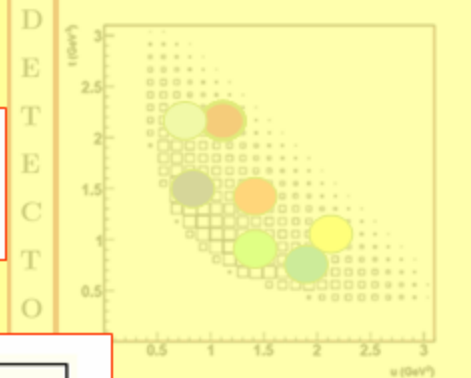
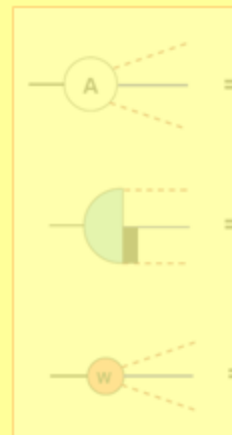
# THEORY

two problems



W

weak

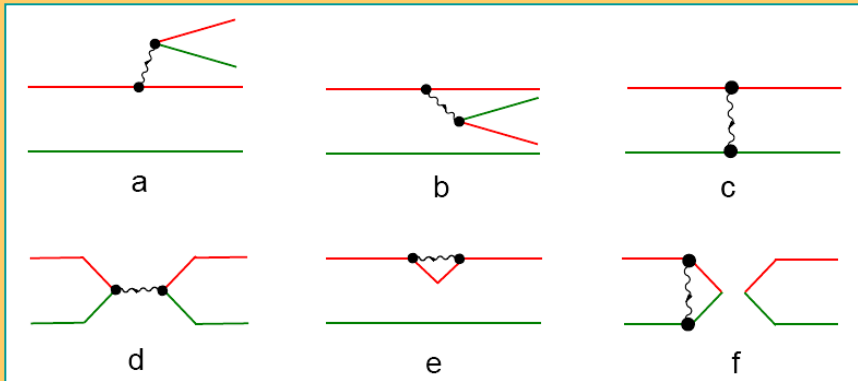




# THEORY

weak vertex

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

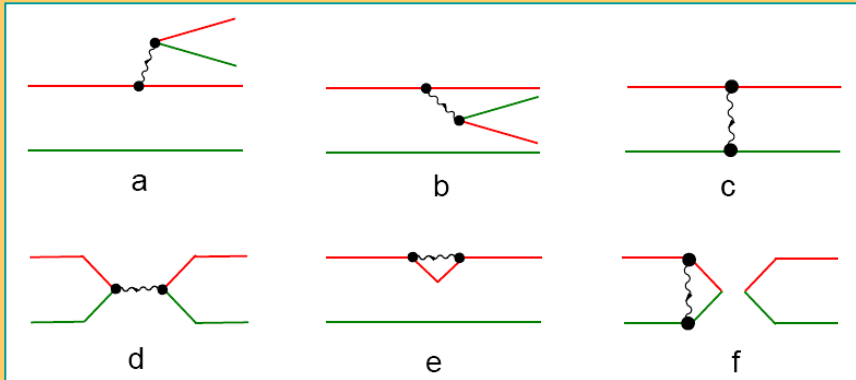


# THEORY

weak vertex

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons

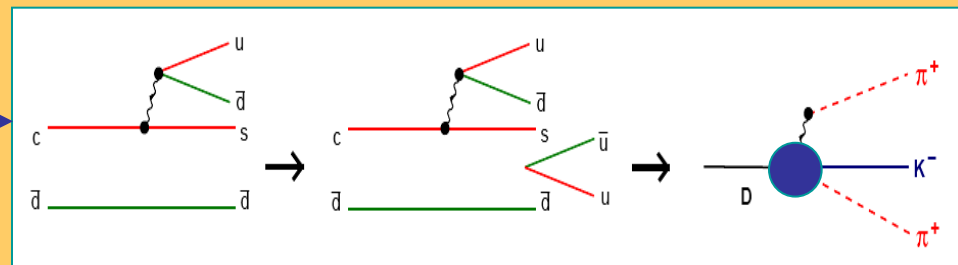
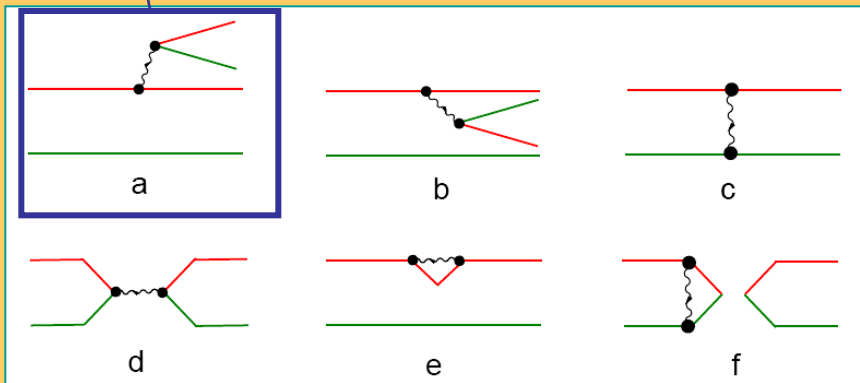


# THEORY

## weak vertex

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons

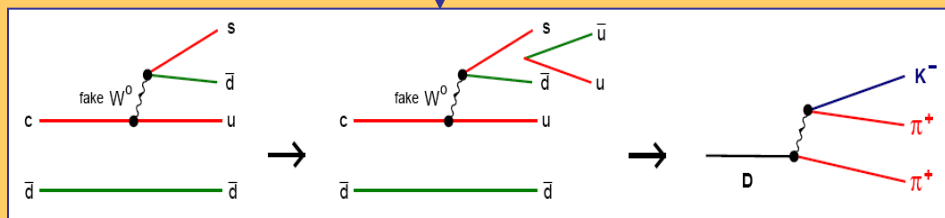
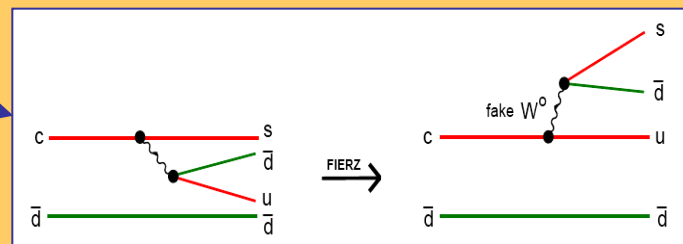
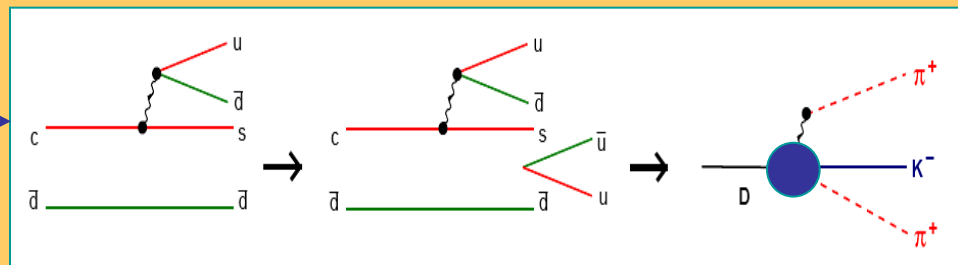
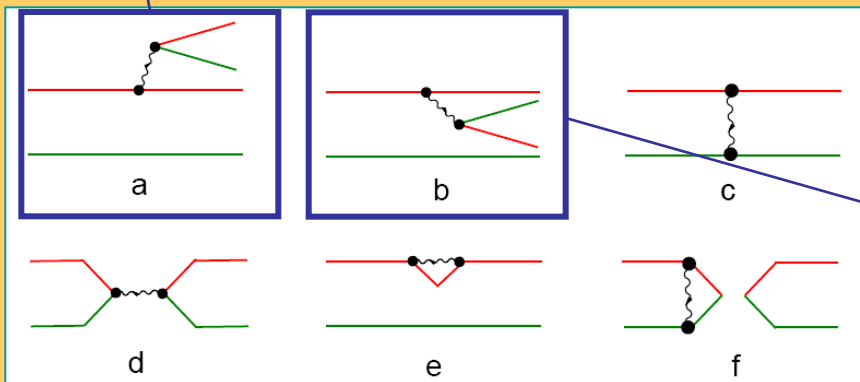


# THEORY

## weak vertex

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons



# THEORY

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons



$W^\pm$

phases in QCD ?

$\rightarrow$  unitarity


$\rightarrow$  loops

field theory &  
dispersion relations

# THEORY

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

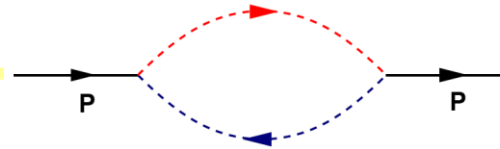
quarks  $\rightarrow$  hadrons



$W^\pm$

effective lagrangians

phases in QCD ?  
 $\rightarrow$  unitarity  
 $\rightarrow$  loops  
field theory &  
dispersion relations



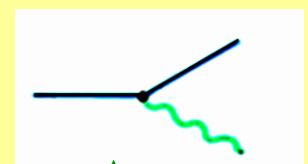
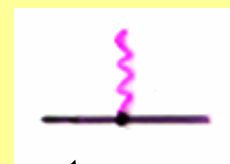
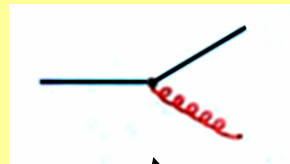
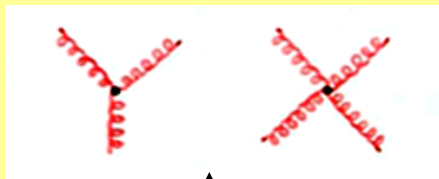
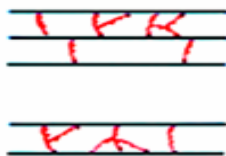
phases from  
first principles

# THEORY

# QCD



# HADRONS



$$\mathcal{L} = \frac{1}{4} \sum F_{\mu\nu}^a F_{\mu\nu}^a + \sum \sum \bar{q} \cdot [i\gamma^\mu (\partial_\mu - i g \frac{\lambda}{2} A_\mu) - m] q$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c$$

external parameter

external quantum number

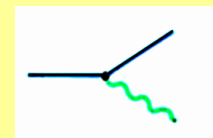
electroweak probes

# THEORY

# QCD

effective lagrangians

$$\mathcal{L} = \text{[red scribble]} + \sum \bar{q} \cdot [i\gamma^\mu \partial_\mu \text{[red scribble]} - m] q.$$



simplified version

$$\mathcal{L} = \frac{1}{4} \sum F_{\mu\nu}^a F_{\mu\nu}^a + \sum \sum \bar{q} [i\gamma^\mu (\partial_\mu - i g \frac{\lambda}{2} A_\mu) - m] q.$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c$$

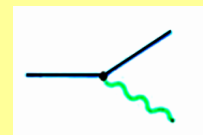


# THEORY

# QCD

effective lagrangians

$$\mathcal{L} = \text{gluon} + \sum \bar{q} \cdot [i\gamma^\mu \partial_\mu \text{gluon} - m] q.$$



**symmetries**

**hadrons**

→ baryon number

→ isospin → all quarks have the same  $m_q$

→ chiral symmetry → all quarks have  $m_q = 0$

# THEORY

# QCD

effective lagrangians

$$\mathcal{L} = \text{[gluon term]} + \sum \bar{q} \cdot [i \gamma^\mu \partial_\mu \text{[quark term]} - m \cdot] q.$$

2 or 3 light quarks

$$\begin{aligned} \mathcal{L}_M^{(2)} = & 3 F^2 B s_0 \\ & + \frac{1}{2} \partial_\mu \phi_i \partial_\mu \phi_i - B (s_0 \delta_{ij} + s_a d_{aij}) \phi_i \phi_j \\ & - \frac{1}{6 F^2} f_{ijs} f_{kls} \phi_i \partial_\mu \phi_j \phi_k \partial_\mu \phi_l + \frac{B}{24 F^2} \left[ s_0 \left( \frac{4}{3} \delta_{ij} \delta_{kl} + 2 d_{ijs} d_{kls} \right) \right. \\ & \left. + s_a \left( \frac{4}{3} \delta_{ij} d_{kla} + \frac{4}{3} d_{ija} \delta_{kl} + 2 d_{ijm} d_{kln} d_{amn} \right) \right] \phi_i \phi_j \phi_k \phi_l + \dots \end{aligned}$$

hadrons

$\phi$

pion field

$$|\phi| = F_\pi \sin \theta$$

pion  
Goldstone boson



MULTIPLETS

PARITY

BOSONS

$|V\rangle \leftrightarrow |0\rangle$

Wigner  
Weyl

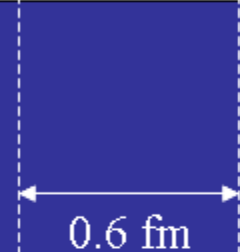
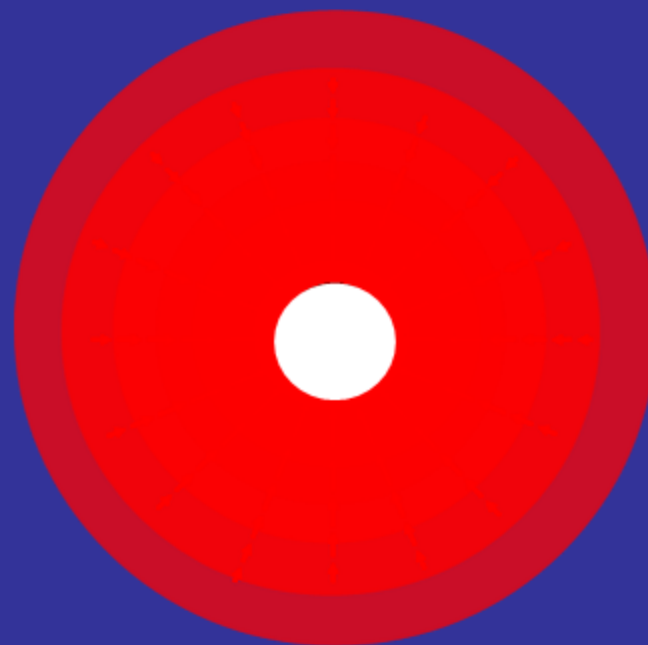
Nambu  
Goldstone

# THEORY

quark condensates



pion  
Goldstone boson




pion  
Goldstone boson

# THEORY

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons



$W^\pm$

effective lagrangians

light mesons  $\rightarrow$  chiral symmetry  $\rightarrow SU(3) \times SU(3)$

quark condensates

weak interactions  $\rightarrow$  gauge external couplings

# THEORY

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons



$W^\pm$

effective lagrangians

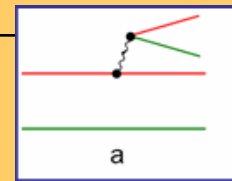
light mesons  $\rightarrow$  chiral symmetry  $\rightarrow SU(3) \times SU(3)$

quark condensates

weak interactions  $\rightarrow$  gauge external couplings

heavy mesons  $\rightarrow$  heavy mass symmetry

# THEORY



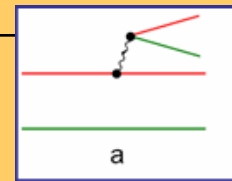
$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons

•••  $\mathcal{L}_{wD_0}^{(1)} = -G_{wD_0} \left[ \frac{g}{2\sqrt{2}} \cos\theta_C \right] W_\mu^- \left\{ \frac{i}{F} \left[ \sqrt{2} D^0 \partial^\mu K^+ + \sqrt{2} D^+ \partial^\mu K^0 + \dots \right] - \frac{1}{2F^2} \left[ \sqrt{2} D^+ \left( \sqrt{2} \pi^- \partial^\mu K^+ - \pi^0 \partial^\mu K^0 + \dots \right) \right] \right\}$

- $D^+ \rightarrow W^+ \bar{K}^0$  [FR-09|2/1]
- $D^+ \rightarrow W^+ \pi^+ K^-$  [FR-09|2/1]
- $D^+ \rightarrow W^+ \pi^0 \bar{K}^0$  [FR-09|2/1]

# THEORY

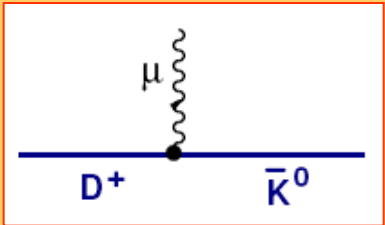


$$D^+ \rightarrow K^- \pi^+ \pi^+$$

quarks  $\rightarrow$  hadrons

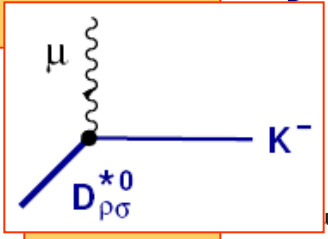
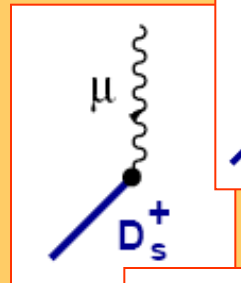


$$\dots \mathcal{L}_{wD_0}^{(1)} = -G_{wD_0} \left[ \frac{g}{2\sqrt{2}} \cos \theta_C \right] W_\mu^- \left\{ \frac{i}{F} \left[ \sqrt{2} D^0 \partial^\mu K^+ + \sqrt{2} D^+ \partial^\mu K^0 + \dots \right] - \frac{1}{2F^2} \left[ \sqrt{2} D^+ \left( \sqrt{2} \pi^- \partial^\mu K^+ - \pi^0 \partial^\mu K^0 + \dots \right) \right] \right\}$$



•  $D^+ \rightarrow W^+ \bar{K}^0$  [FR-09|2/1]

$$\rightarrow i \left[ \frac{g}{2\sqrt{2}} \cos \theta_C \right] \left[ \frac{\sqrt{2} G_{wD_0}}{F} \right] p_K^\mu$$

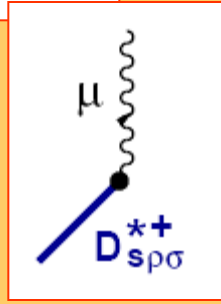


$W^+ \pi^+ K^-$  [FR-09|2/1]

$$\rightarrow - \left[ \frac{g}{2\sqrt{2}} \cos \theta_C \right] \left[ \frac{G_{wD_0}}{F^2} \right] p_K^\mu$$

•  $D^+ \rightarrow W^+ \pi^0 \bar{K}^0$  [FR-09|2/1]

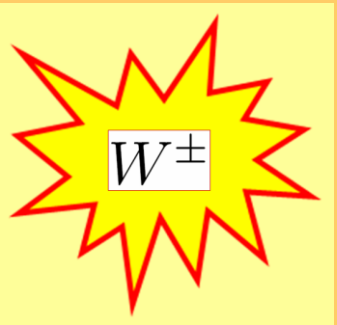
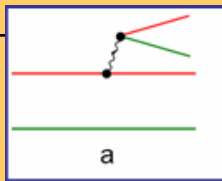
$$\rightarrow \frac{1}{\sqrt{2}} \left[ \frac{g}{2\sqrt{2}} \cos \theta_C \right] \left[ \frac{G_{wD_0}}{F^2} \right] p_K^\mu$$





# THEORY

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

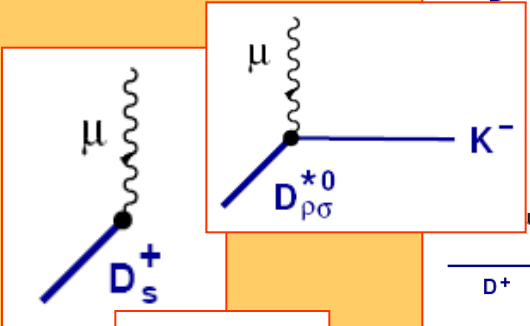
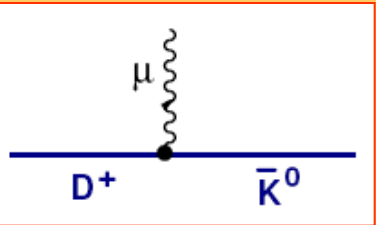
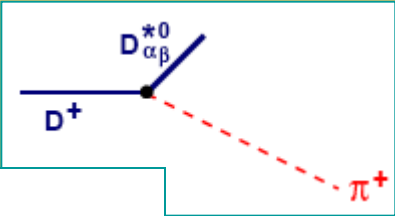
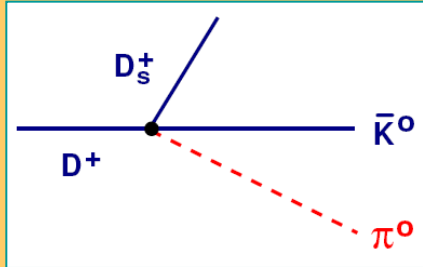


quarks  $\rightarrow$  hadrons

$$\dots \mathcal{L}_{wD_0}^{(1)} = -G_{wD_0} \left[ \frac{g}{2\sqrt{2}} \cos\theta_C \right] W_\mu^- \left\{ \frac{i}{F} \left[ \sqrt{2} D^0 \partial^\mu K^+ + \sqrt{2} D^+ \partial^\mu K^0 + \dots \right] - \frac{1}{2F^2} \left[ \sqrt{2} D^+ \left( \sqrt{2} \pi^- \partial^\mu K^+ - \pi^0 \partial^\mu K^0 + \dots \right) \right] \right\}$$

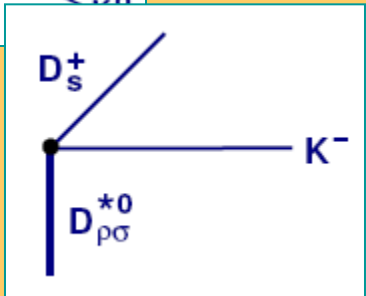
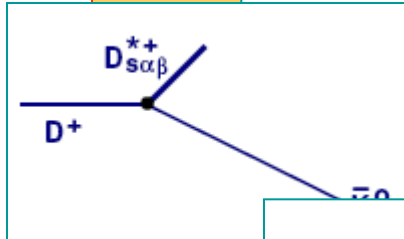
- $D^+ \rightarrow W^+ \bar{K}^0$  [FR-09|2/1]

$$\rightarrow i \left[ \frac{g}{2\sqrt{2}} \cos\theta_C \right] \left[ \frac{\sqrt{2} G_{wD_0}}{F} \right] p_K^\mu$$



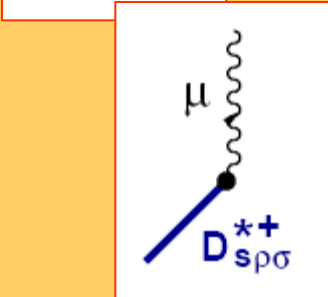
- $D^+ \rightarrow W^+ \pi^+ K^-$  [FR-09|2/1]

$$\rightarrow - \left[ \frac{g}{2\sqrt{2}} \cos\theta_C \right] \left[ \frac{G_{wD_0}}{F^2} \right] p_K^\mu$$



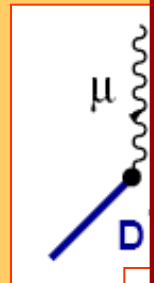
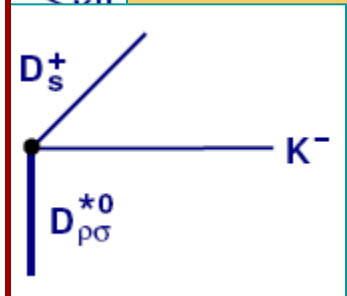
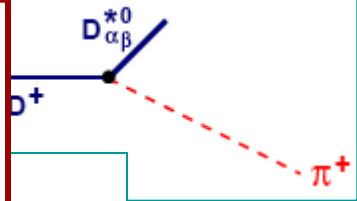
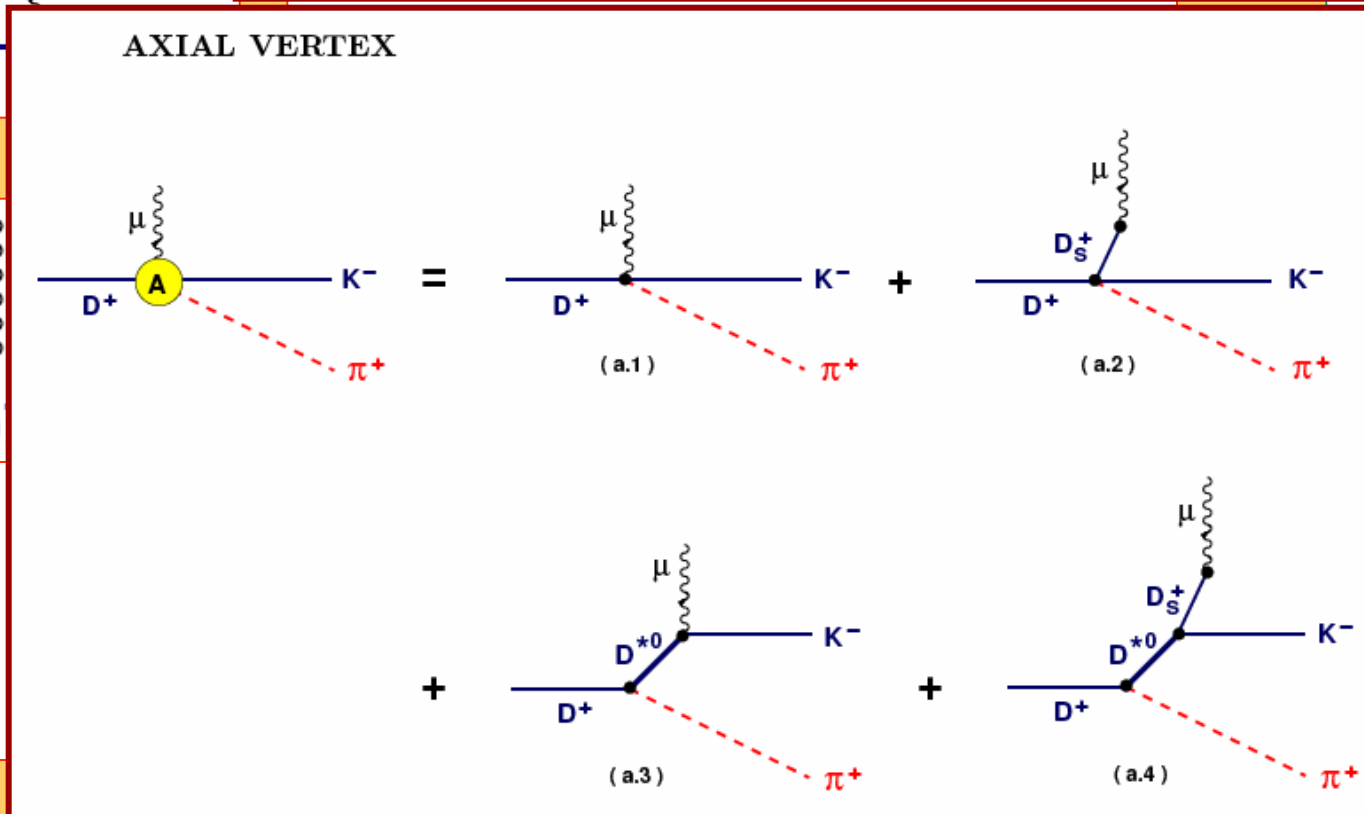
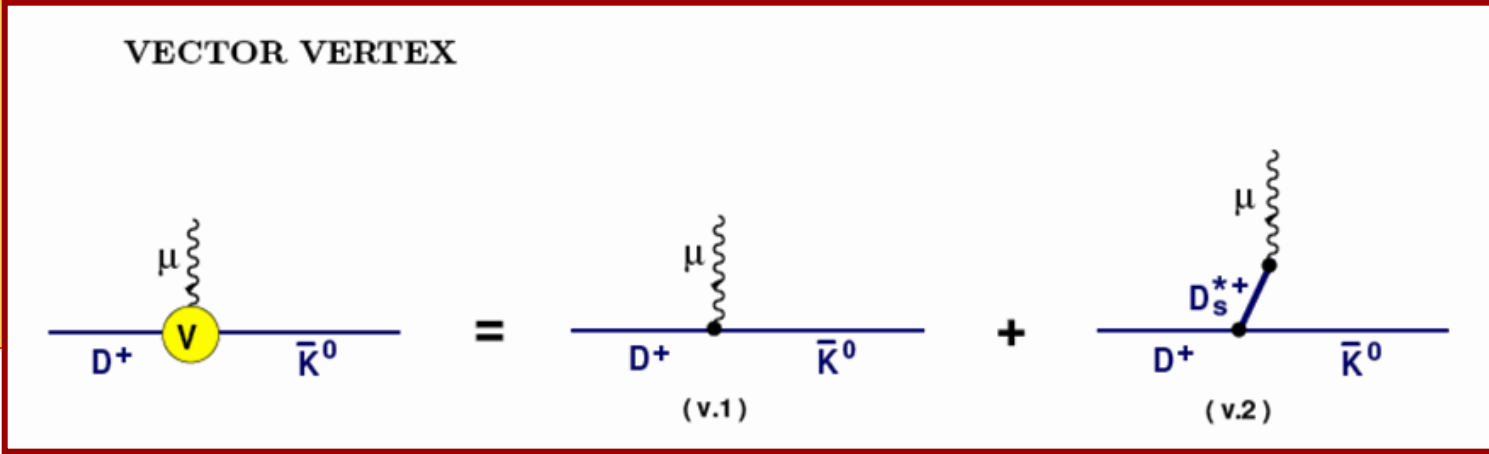
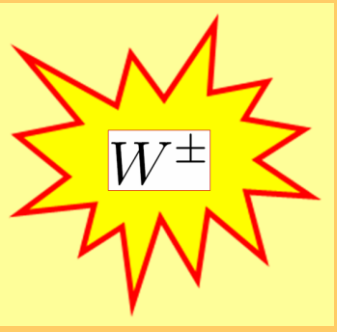
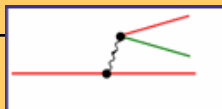
- $D^+ \rightarrow W^+ \pi^0 \bar{K}^0$  [FR-09|2/1]

$$\rightarrow \frac{1}{\sqrt{2}} \left[ \frac{g}{2\sqrt{2}} \cos\theta_C \right] \left[ \frac{G_{wD_0}}{F^2} \right] p_K^\mu$$



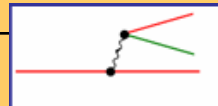
# THEORY

## weak vertex

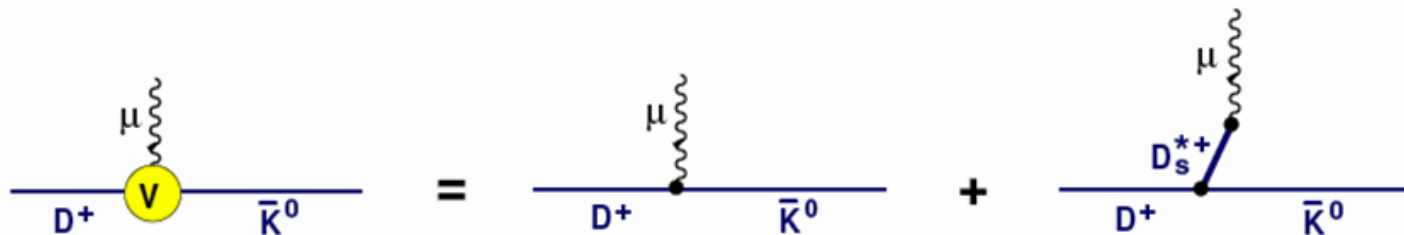


# THEORY

## weak vertex

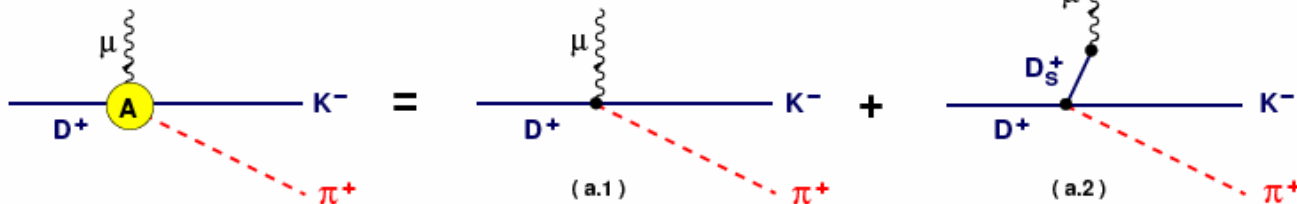
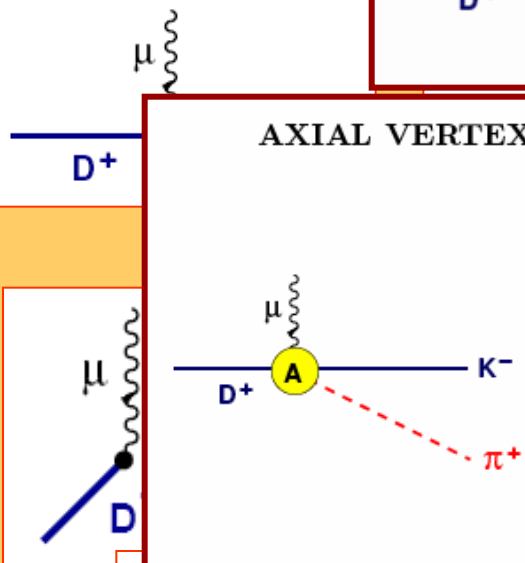


### VECTOR VERTEX

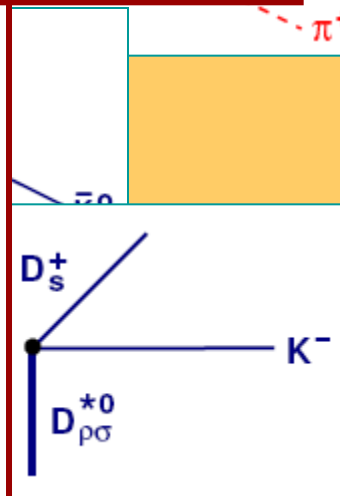
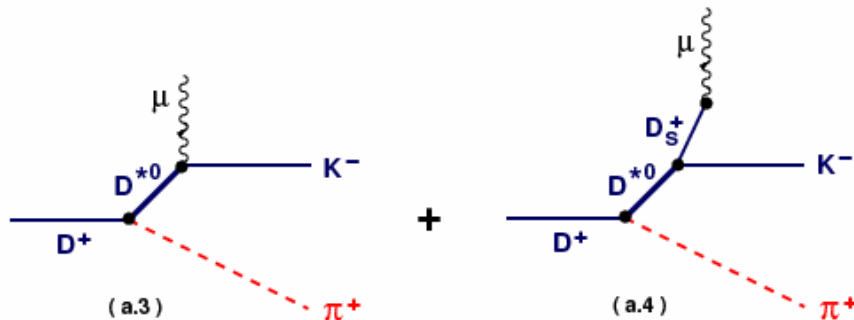


# form factors

### AXIAL VERTEX



some control  
know how to turn  
effects on and off

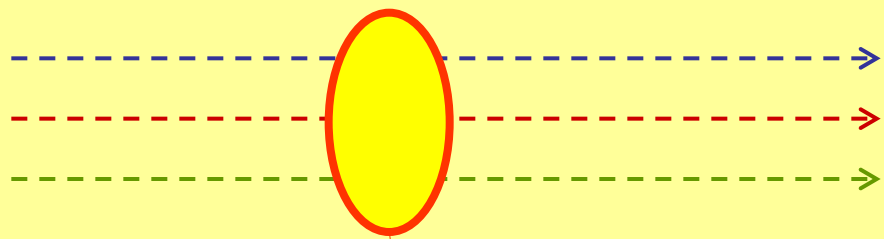


# THEORY

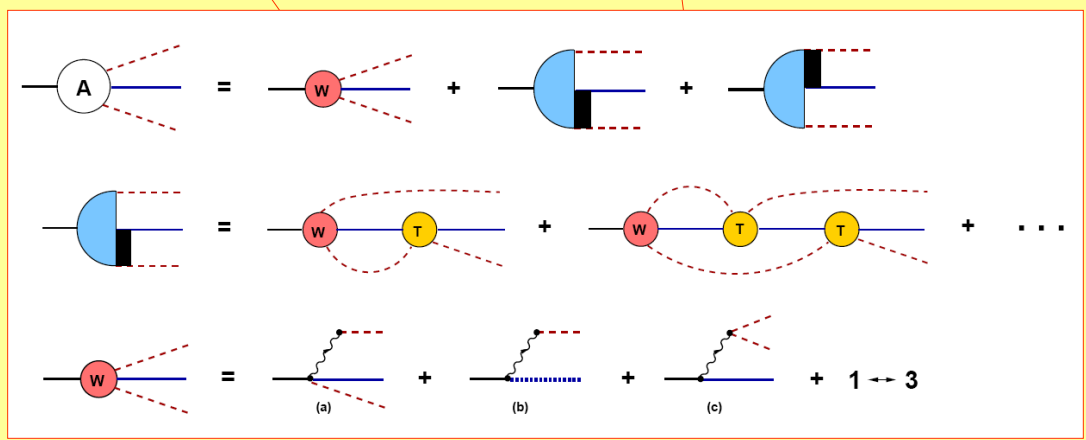
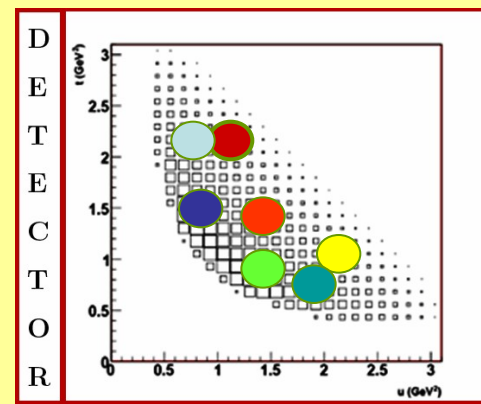
final state interactions



weak vertex



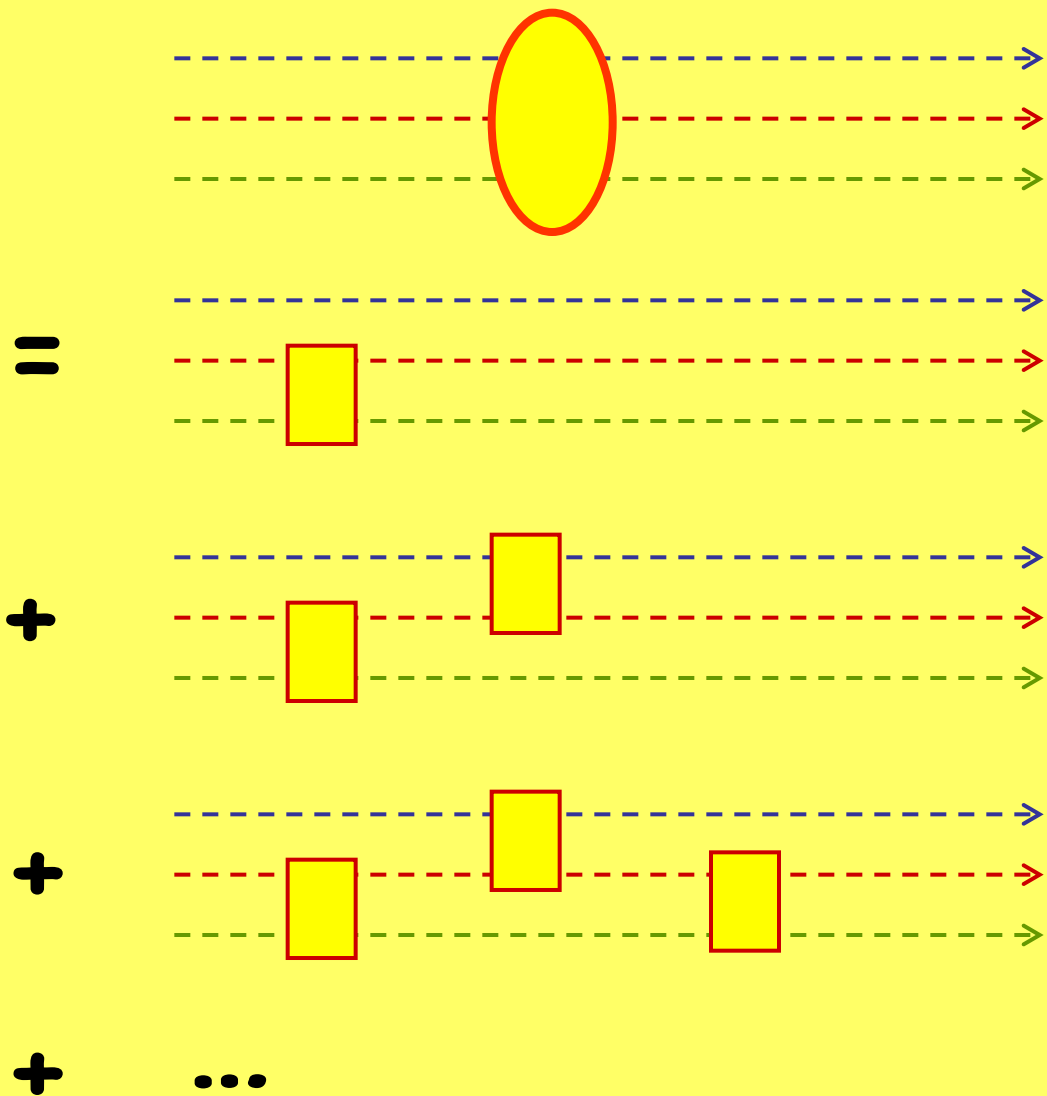
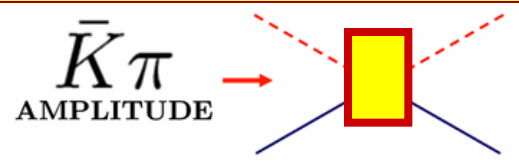
final state interactions



# THEORY

final state interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

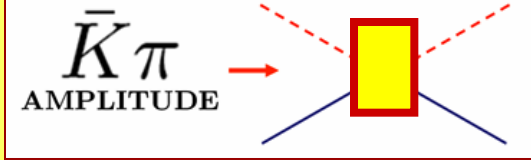


Faddeev series

# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



effective lagrangians

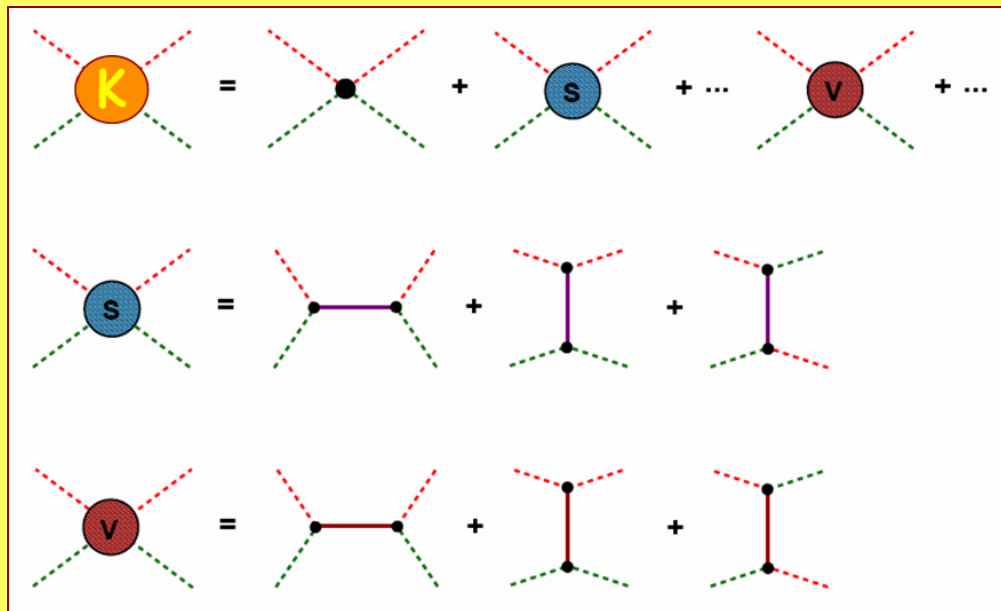
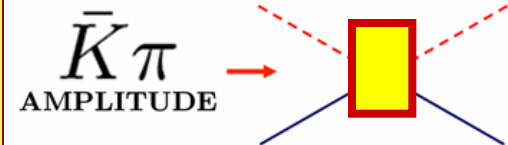
light mesons  $\rightarrow$  chiral symmetry  $\rightarrow SU(3) \times SU(3)$

quark condensates

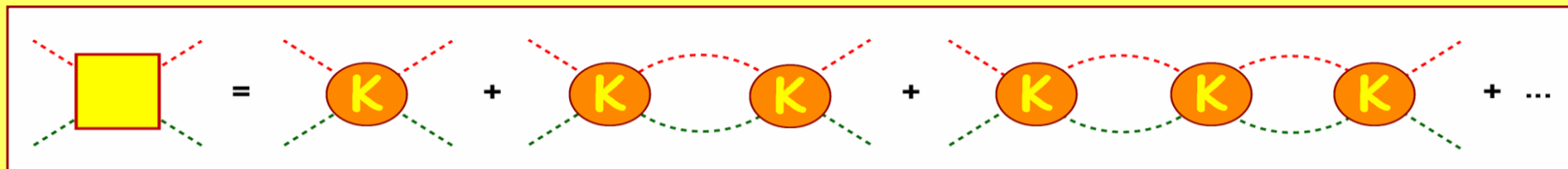
# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



Bethe-Salpeter equation



# THEORY

final state interactions

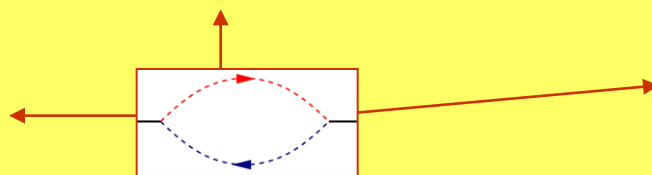
$$D^+ \rightarrow K^- \pi^+ \pi^+$$



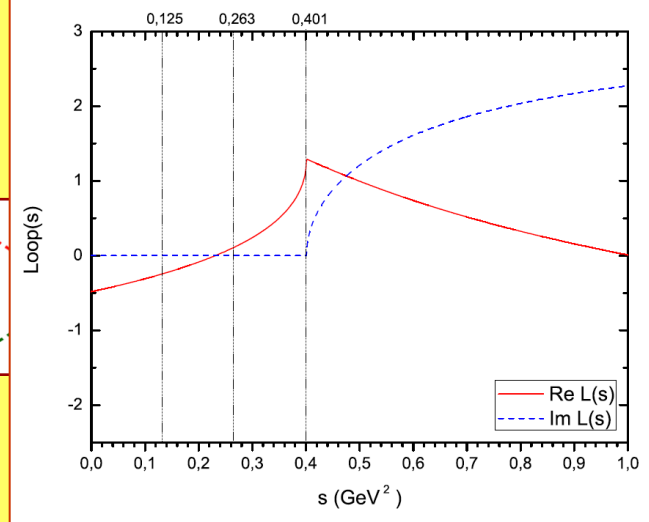
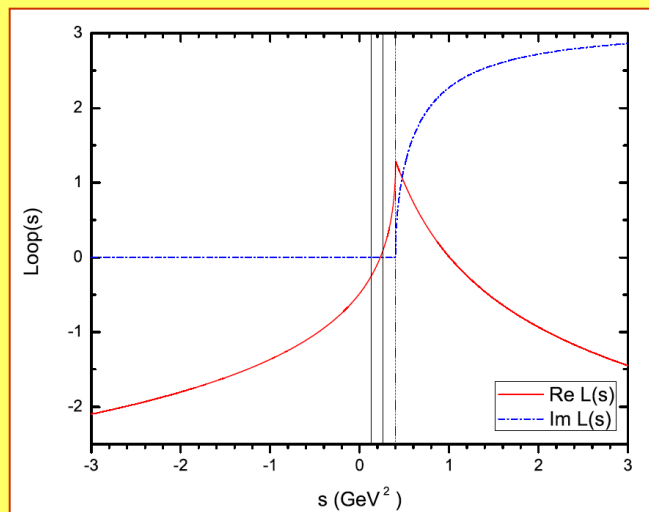
two-body unitarity

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

phase



$$= \text{K} + \text{K} + \text{K} + \dots$$

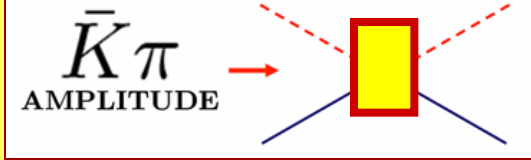




# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



two-body  
unitarity

$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$

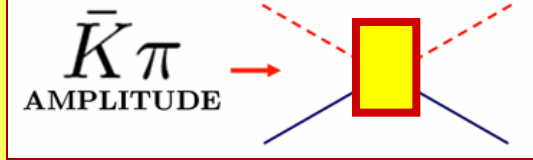
$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

Breit-Wigner  
+ chiral symmetry

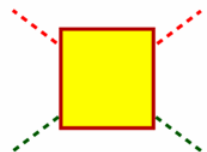
# THEORY

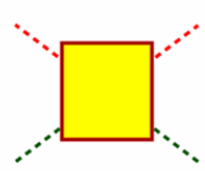
final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



two-body  
unitarity


$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$


$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

Breit-Wigner  
+ chiral symmetry

isobar model  
is crude

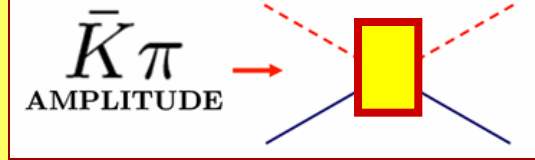
# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$

K-matrix  
approximation

$$\mathcal{M} \rightarrow m^*$$



two-body  
unitarity

A Feynman diagram for two-body unitarity, showing a yellow square box with a red border. Two red dashed lines enter from the top, and two green dashed lines enter from the bottom. Two red dashed lines exit towards the top right, and two green dashed lines exit towards the bottom right.

$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$

A Feynman diagram for the K-matrix approximation, showing a yellow square box with a red border. Two red dashed lines enter from the top, and two green dashed lines enter from the bottom. Two red dashed lines exit towards the top right, and two green dashed lines exit towards the bottom right.

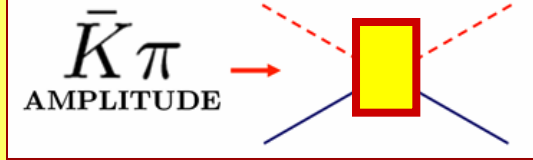
$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

Breit-Wigner  
+ chiral symmetry

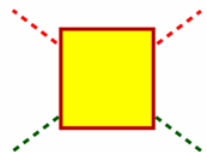
# THEORY

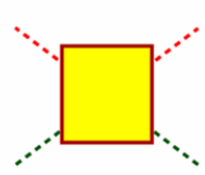
final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



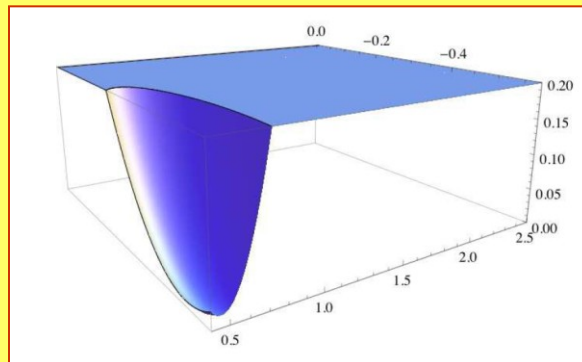
two-body  
unitarity


$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$


$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

poles

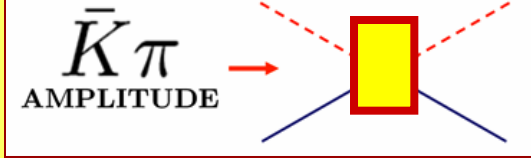
kappa  
dynamical



# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



two-body  
unitarity

A Feynman diagram for two-body unitarity, showing a yellow square box with a red border. Two red dashed lines enter from the top, and two green dashed lines enter from the bottom.

$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$

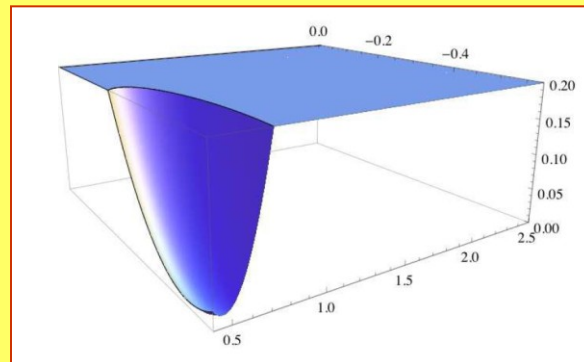
A Feynman diagram for the  $\bar{K} \pi$  amplitude, showing a yellow square box with a red border. Two red dashed lines enter from the top, and two green dashed lines enter from the bottom.

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

light sector of QCD  $\rightarrow$  chiral symmetry  
sigma & kappa  
cannot be accommodated into chiral lagrangians

poles

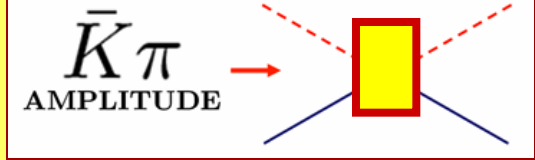
kappa  
dynamical



# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



two-body  
unitarity

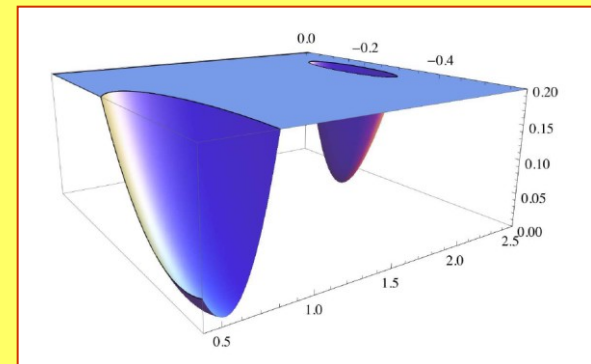
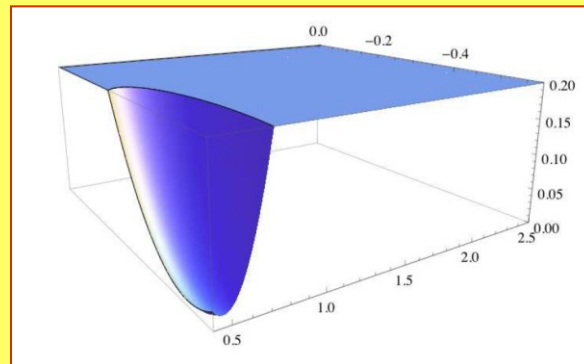
$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

light sector of QCD  $\rightarrow$  chiral symmetry  
sigma & kappa  
cannot be accommodated into chiral lagrangians

poles

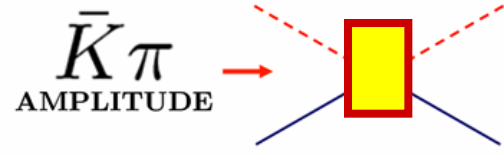
kappa  
dynamical



# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



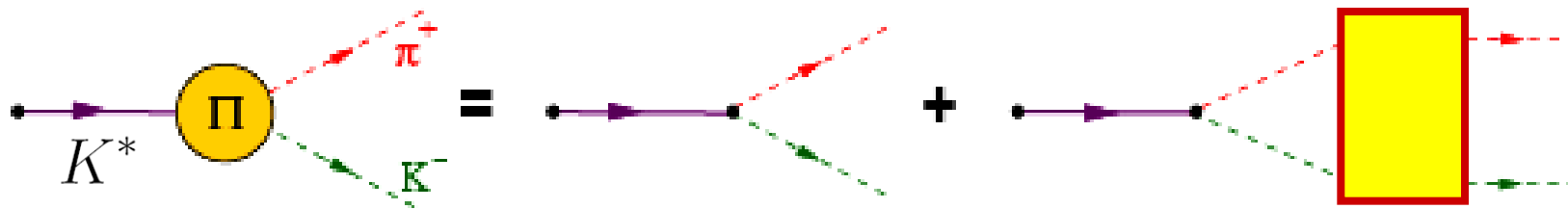
two-body  
unitarity

$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

production amplitude

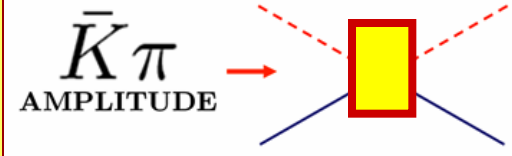
dynamical  
width



# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



two-body  
unitarity

$$= - \left[ \frac{16\pi}{\rho} \frac{m_r \Gamma}{s - \mathcal{M}^2 + i m_r \Gamma} \right]$$

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

CONSISTENT  
theoretical  
scheme



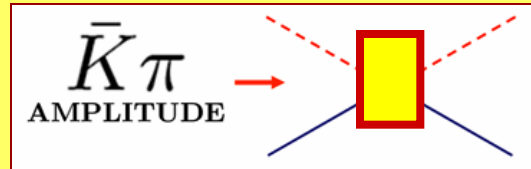
order  
in the  
MESS



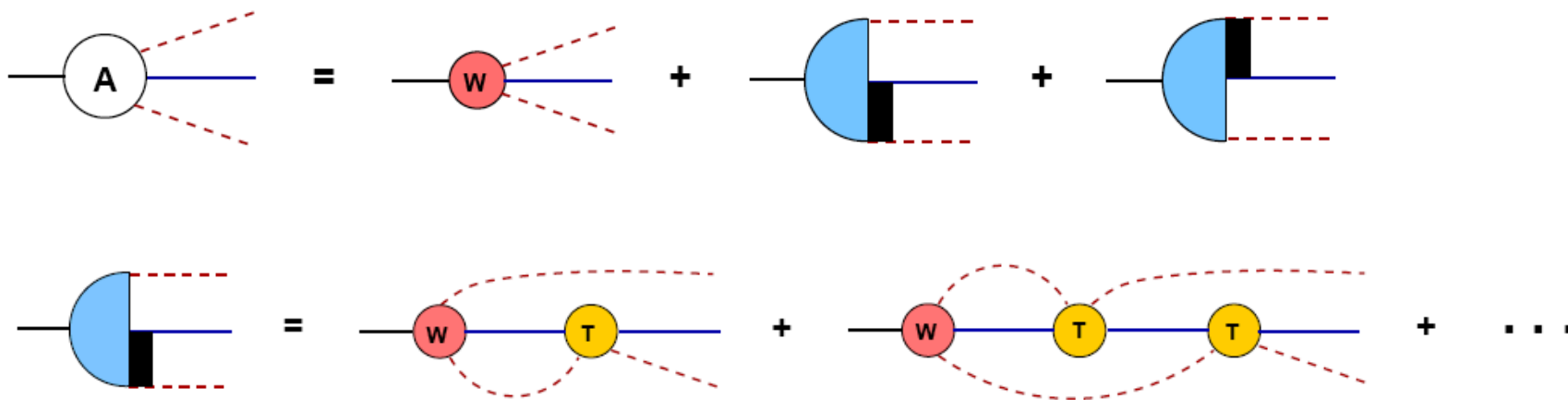
# THEORY

final state  
interactions

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



Faddeev  
series

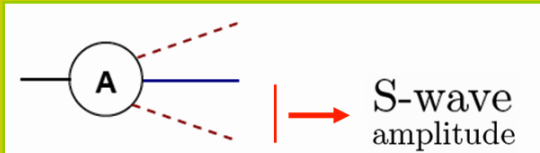


# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

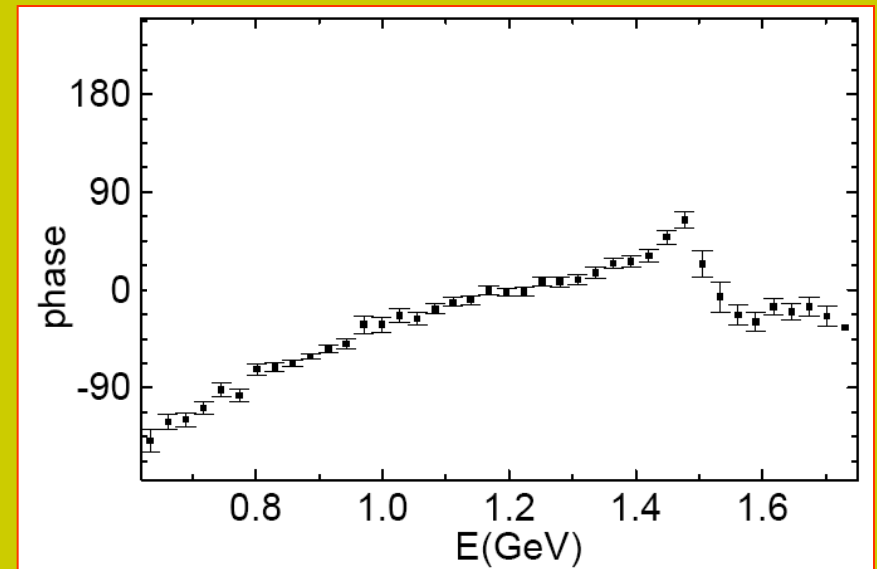
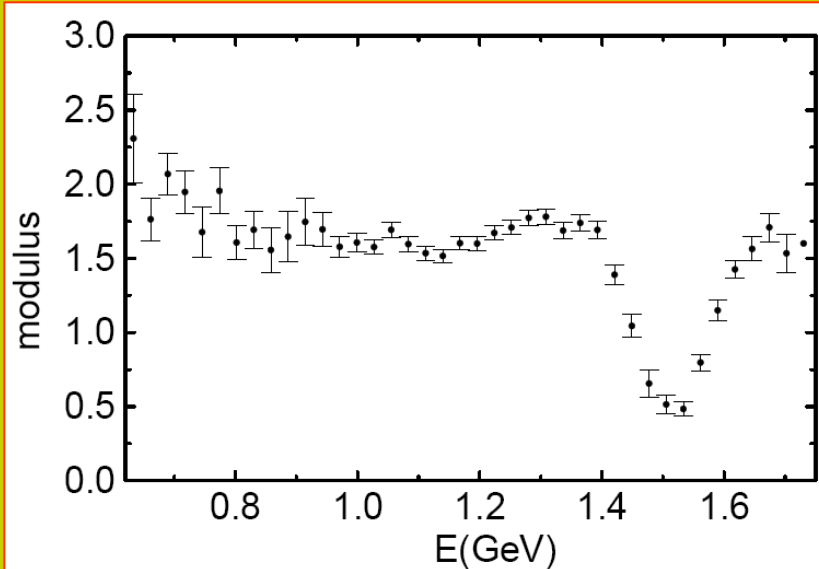
PHYSICAL REVIEW D **84**, 094001 (2011)

P.C. Magalhães, M.R. R., K.S.F.F. Guimarães, T. Frederico, W.S. de Paula,  
I. Bediaga, A.C. dos Reis, C.M. Maekawa and G.R.S. Zarnauskas

data



E.M. Aitala *et al.* [E791 Collaboration], Phys. Rev. D **73**, 032004 (2006); Erratum-ibid. D **74**, 059901 (2006)  
J.M. Link *et al.* [FOCUS Collaboration], Phys. Lett. B **681**, (2009) 14

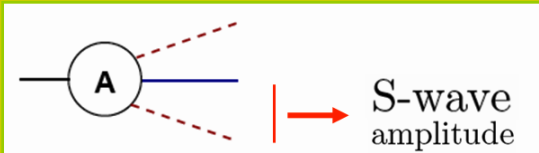


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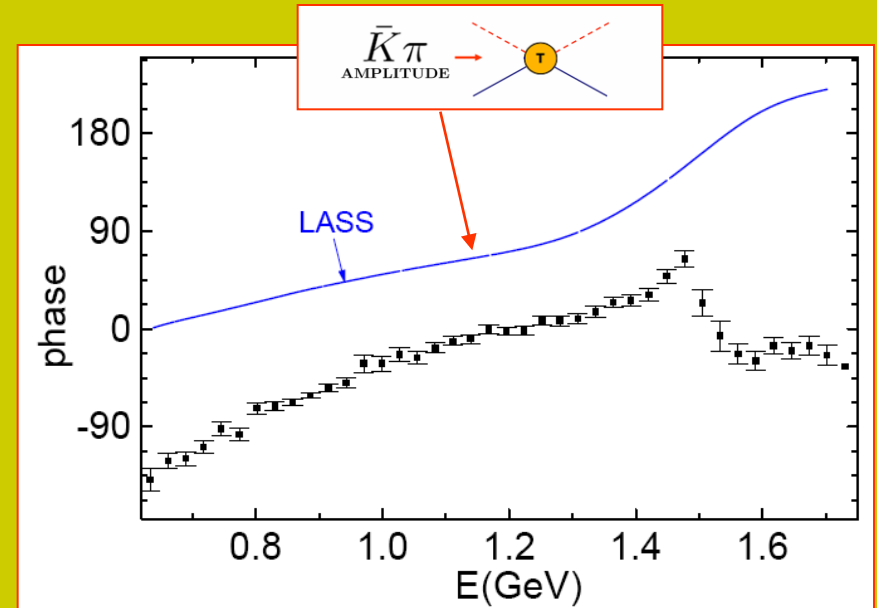
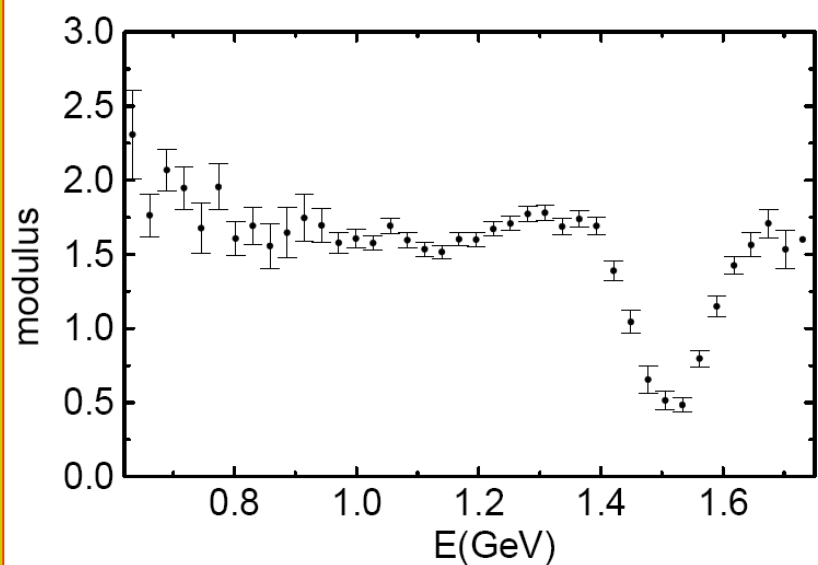
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data



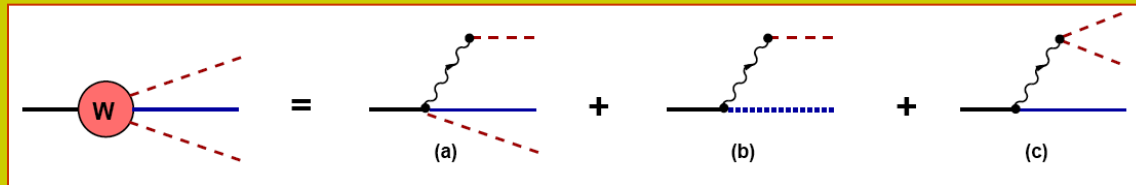
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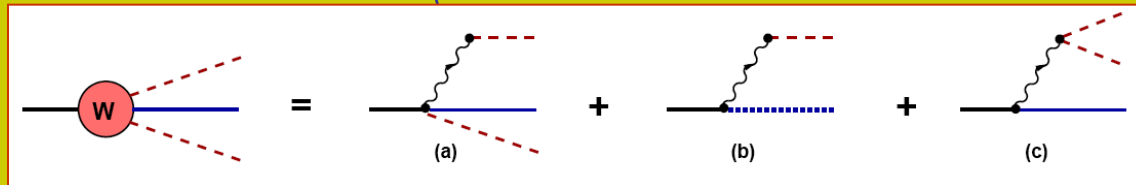
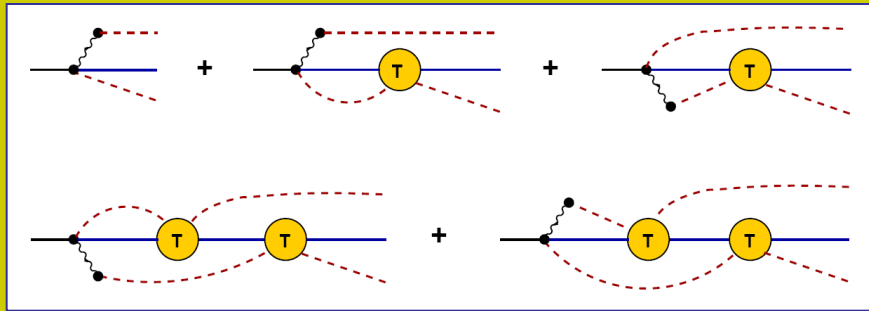


topologies without structure

# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

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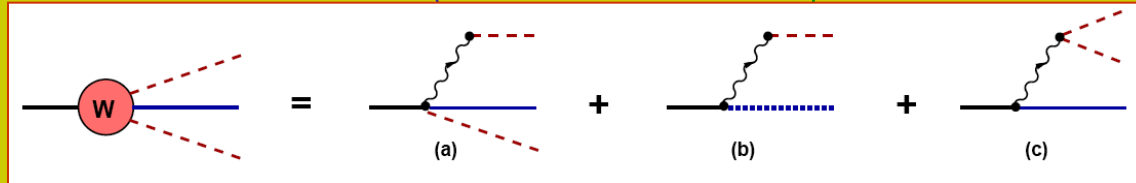
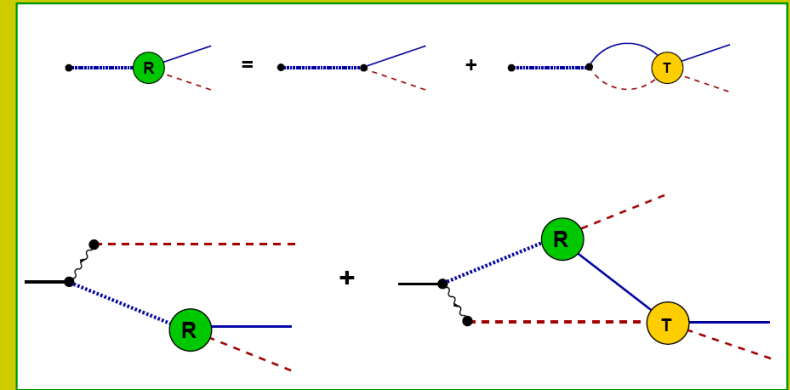
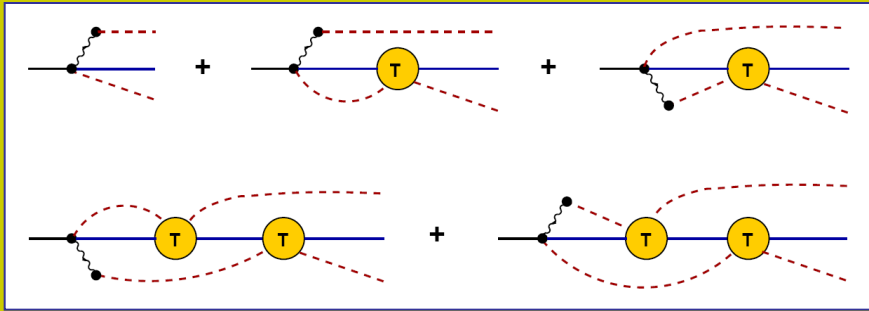


topologies without structure

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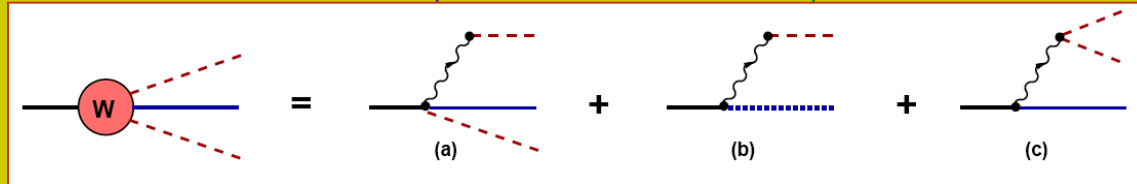
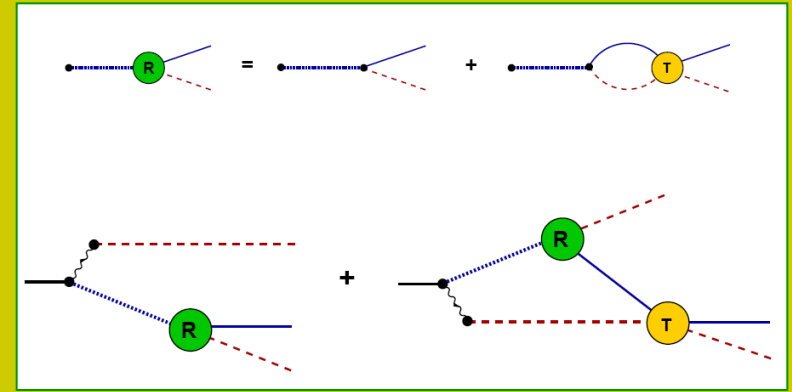
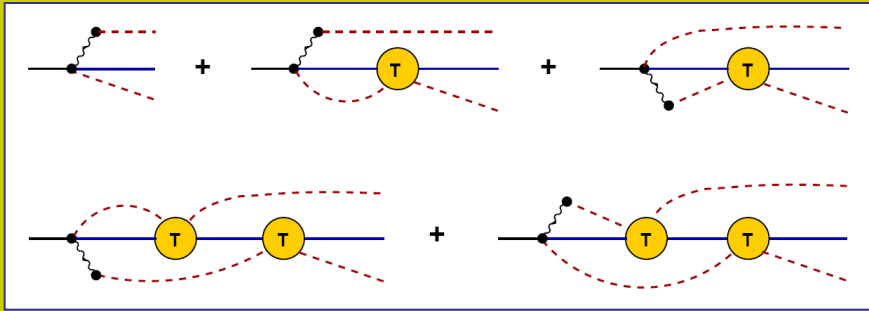


topologies without structure

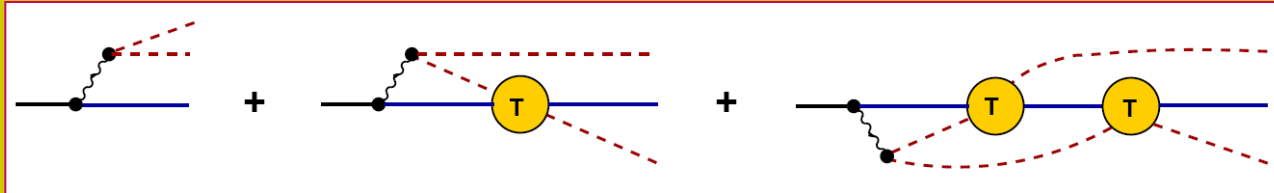
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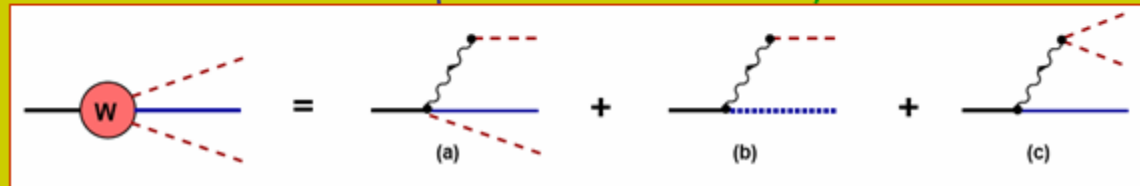
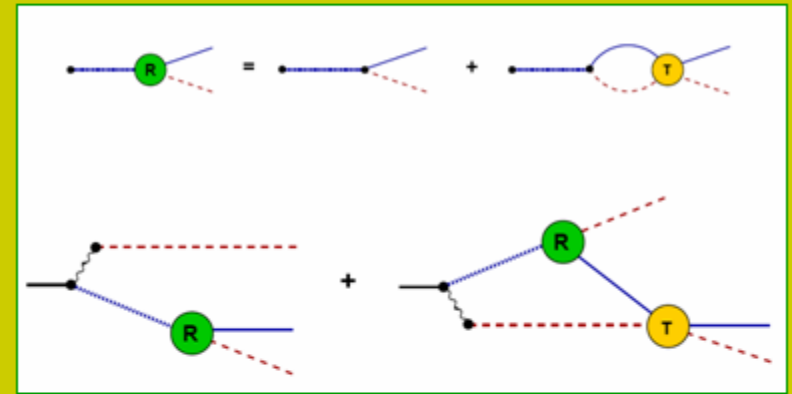
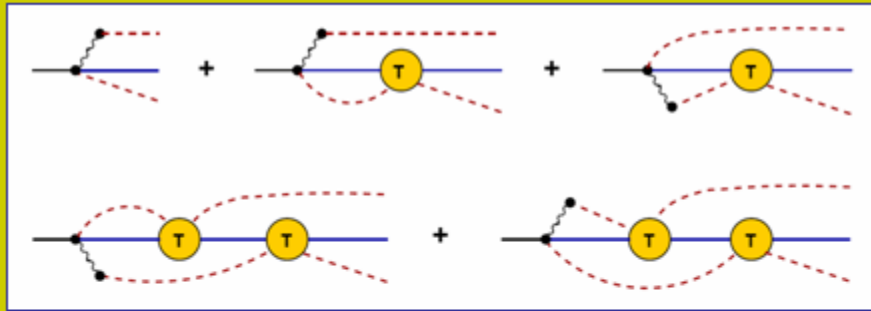
topologies without structure



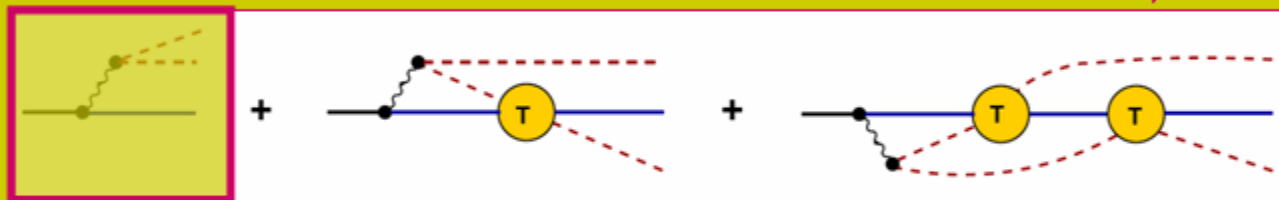
# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

PHYSICAL REVIEW D **84**, 094001 (2011)

P.C. Magalhães, M.R. R., K.S.F.F. Guimarães, T. Frederico, W.S. de Paula,  
I. Bediaga, A.C. dos Reis, C.M. Maekawa and G.R.S. Zarnauskas



topologies without structure

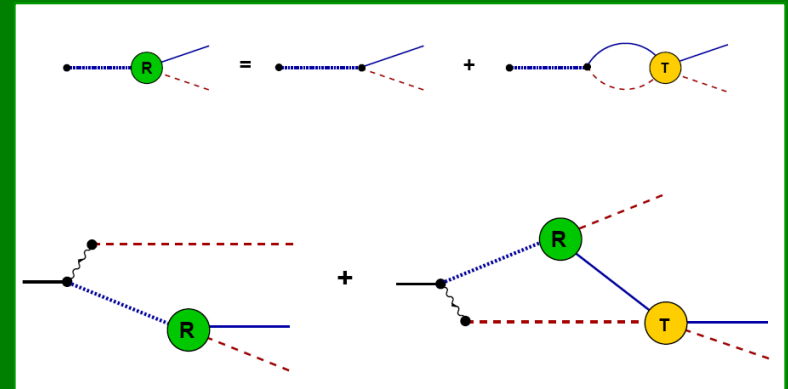
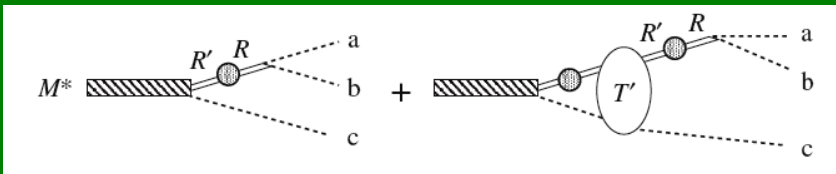
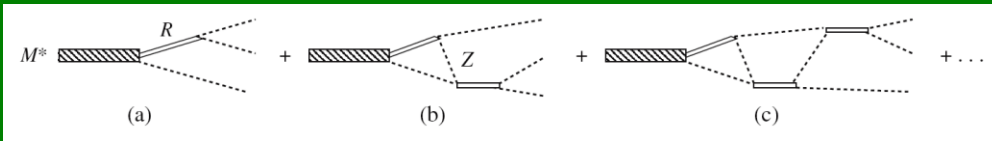




PHYSICAL REVIEW D **84**, 114019 (2011)

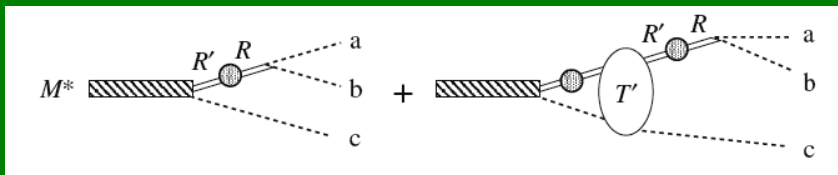
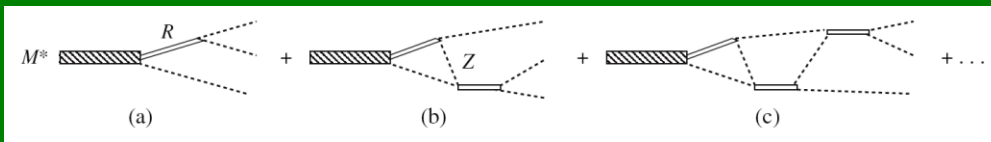
# Unitary coupled-channels model for three-mesons decays of heavy mesons

H. Kamano,<sup>1,2</sup> S. X. Nakamura,<sup>2</sup> T.-S. H. Lee,<sup>3,2</sup> and T. Sato<sup>4,2</sup>



# Unitary coupled-channels model for three-mesons decays of heavy mesons

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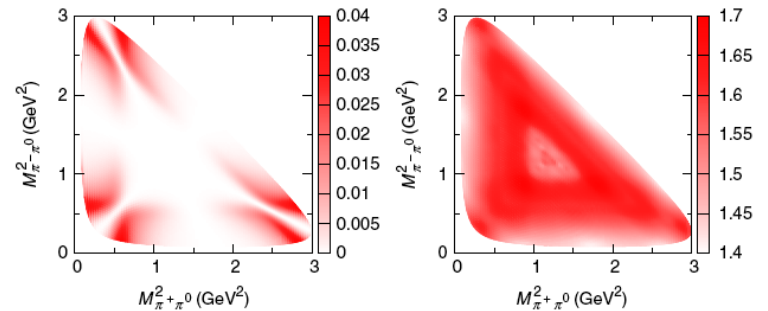


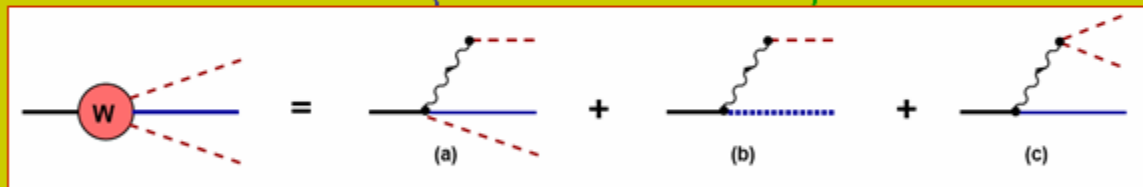
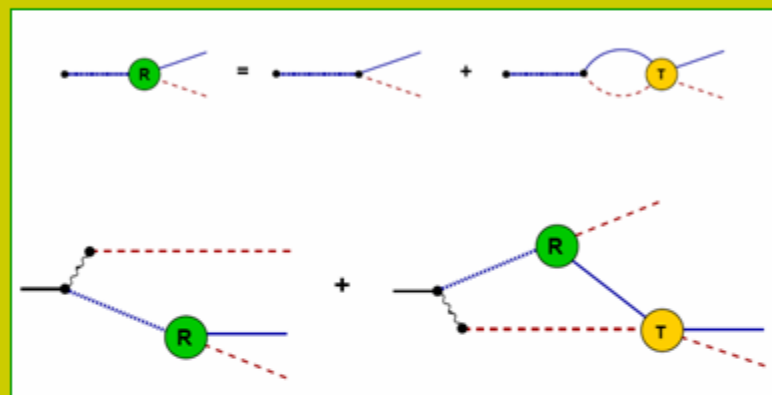
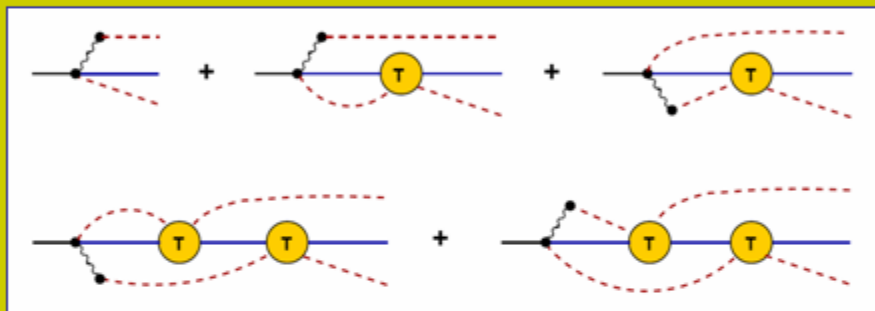
FIG. 4 (color online). (left) Dalitz plot of  $D^0 \rightarrow \pi^+ \pi^- \pi^0$  decay; (right) ratio of Dalitz plot distributions with and without the  $Z$  graphs.

Dalitz plot distributions by as much as 50% in magnitudes in the regions where  $f_0(600)$ ,  $\rho(770)$ , and  $f_2(1270)$  dominate the distributions. Also, by fitting to the same Dalitz plot distributions, we demonstrate that the decay amplitudes obtained with the unitary model and the isobar model can be rather different, particularly in the phase that plays a crucial role in extracting the Cabibbo-Kobayashi-Maskawa  $CP$ -violating phase from the data of  $B$  meson decays. Our results indicate that the commonly used isobar-model analysis must be extended to account for the final state interactions required by the three-body unitarity to reanalyze the three-mesons decays of heavy mesons, thereby exploring hybrid or exotic

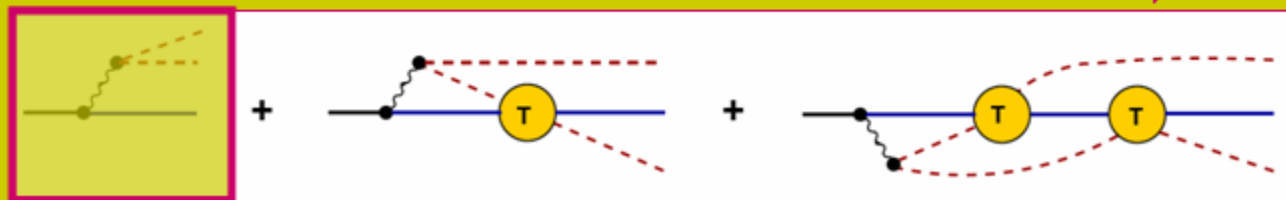
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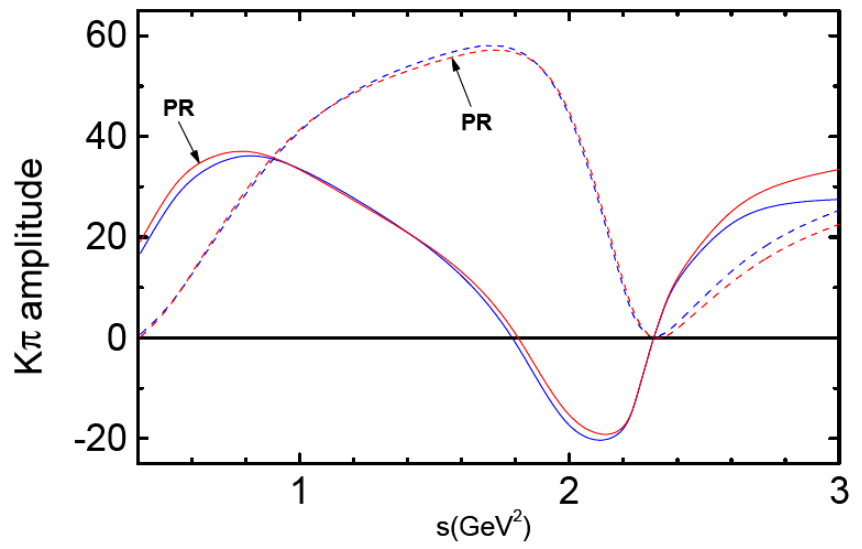
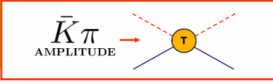
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topologies without structure

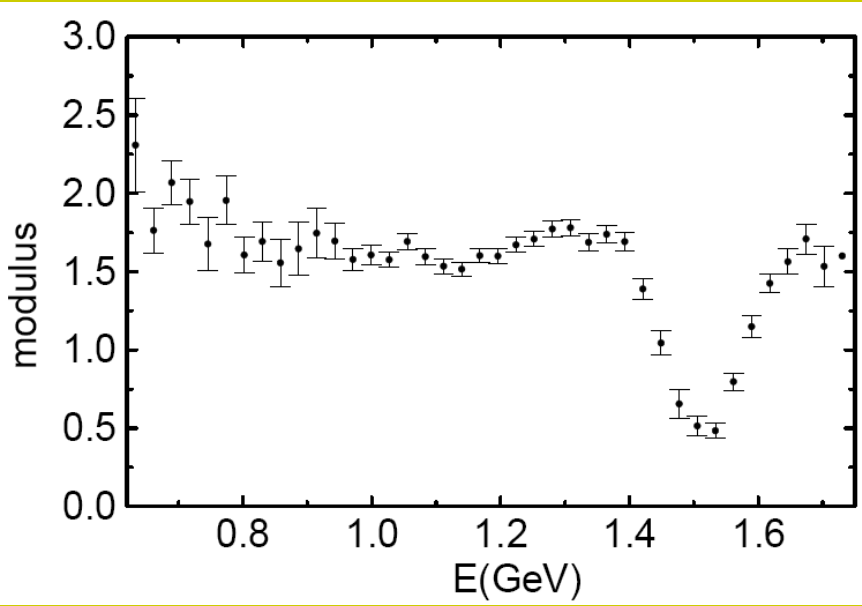


# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

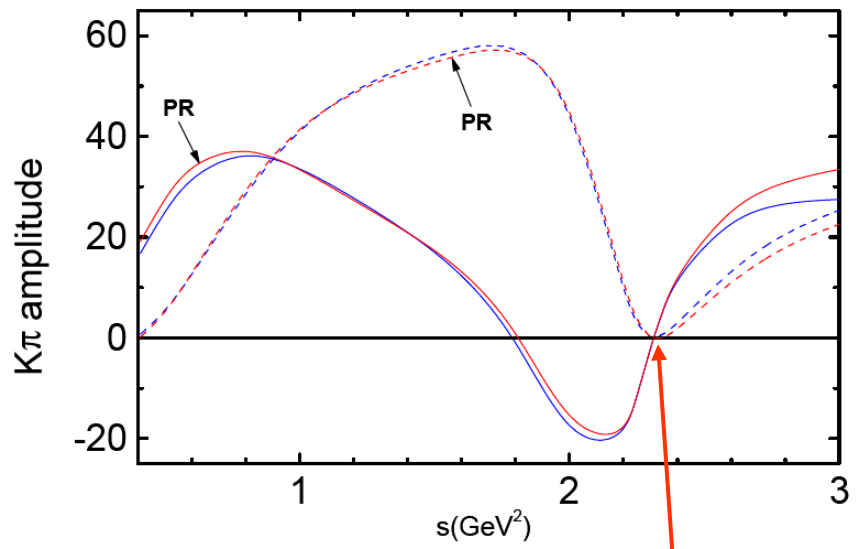
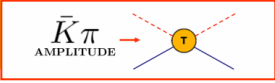


A diagram of a yellow square box with a red border. Dashed red lines enter from the top-left and top-right corners, and dashed green lines enter from the bottom-left and bottom-right corners. To the right of the box is the equation:

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$



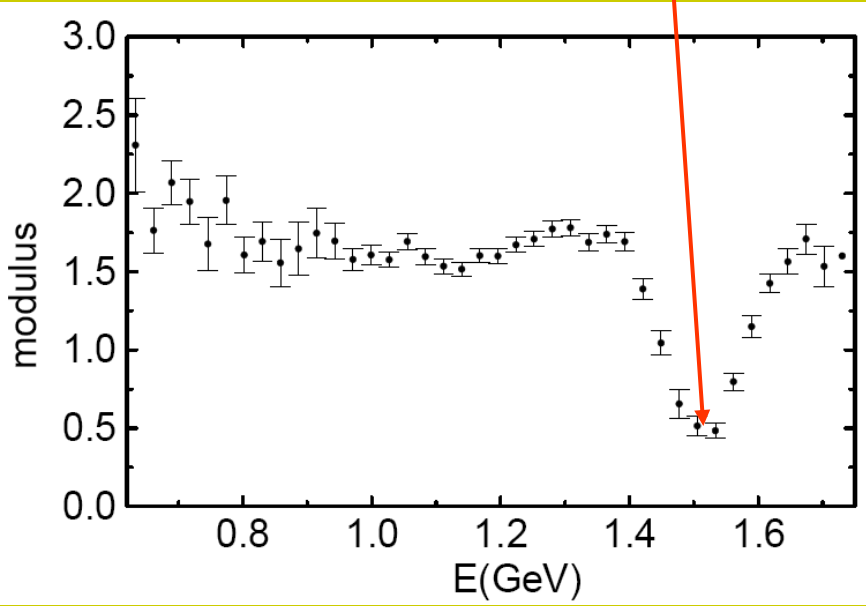
# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$



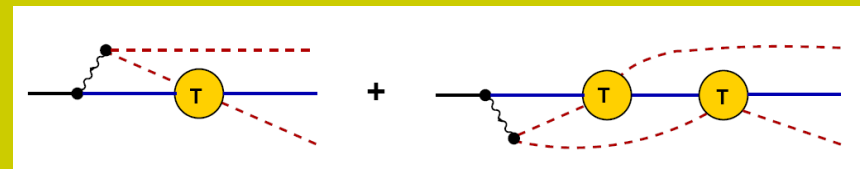
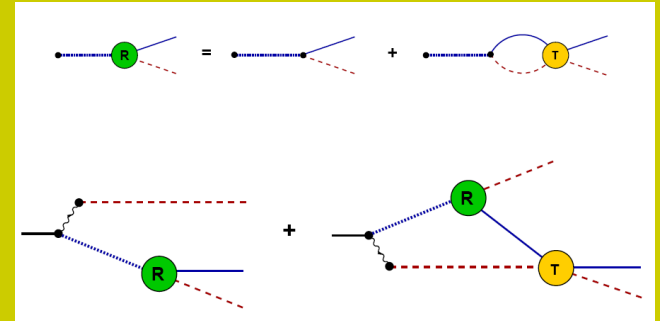
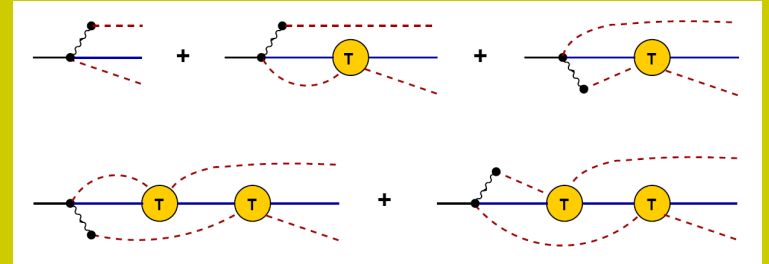
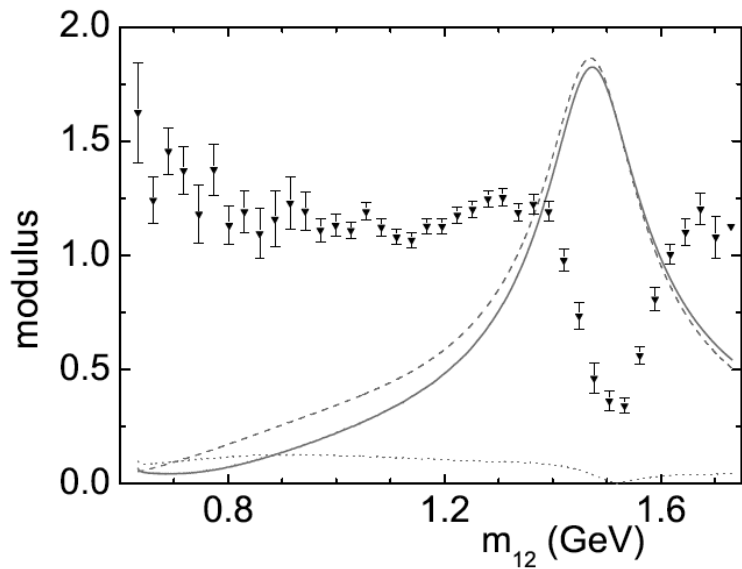
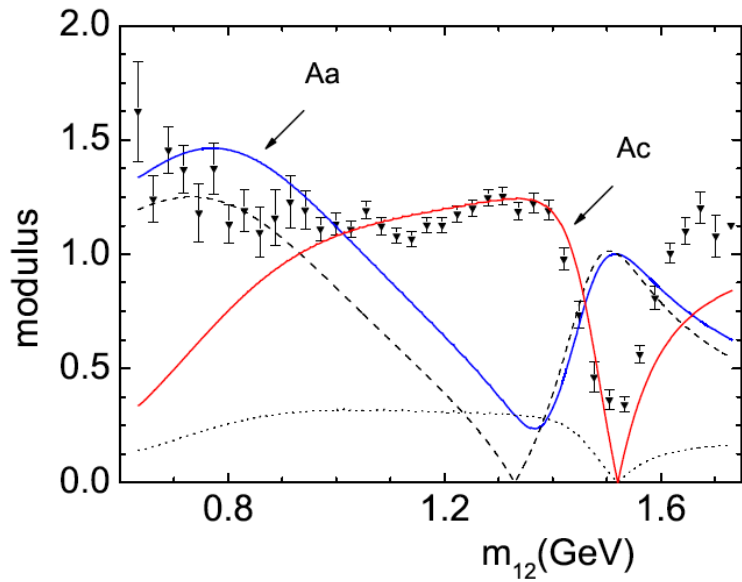
A diagram of a yellow square loop with dashed lines entering and exiting. A red arrow points from this diagram to the left. To the right of the diagram is the equation:

$$= \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$

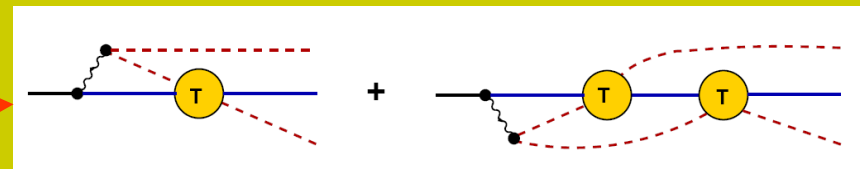
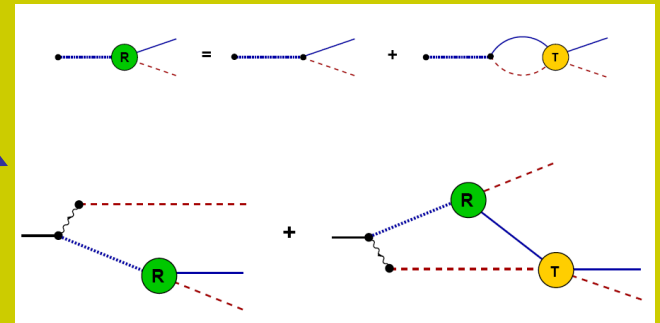
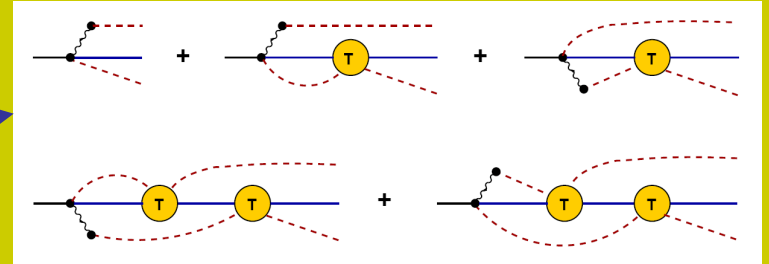
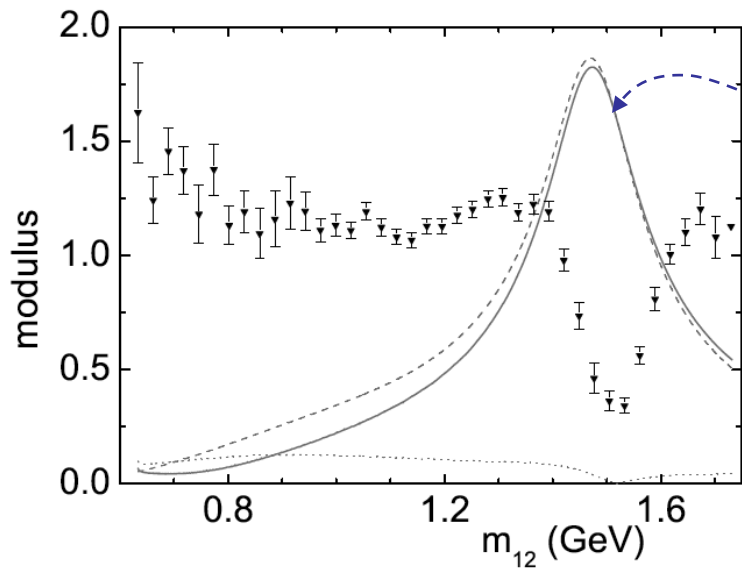
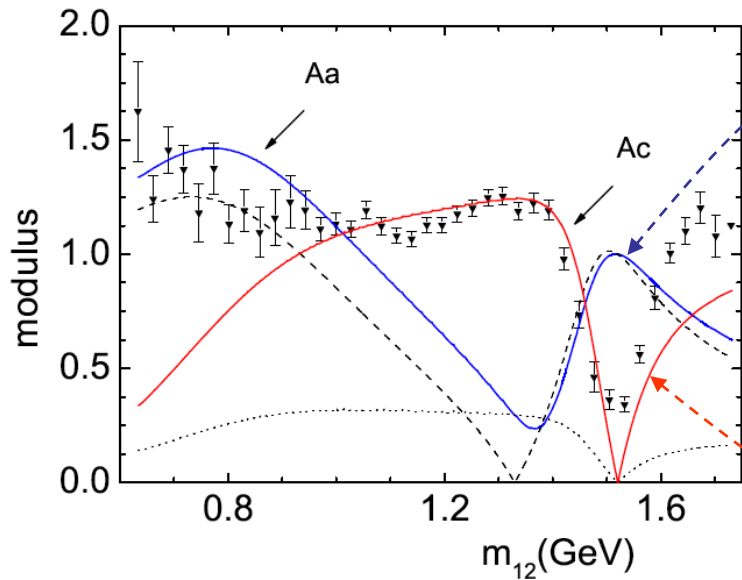
The symbol  $\mathcal{K}$  is circled in red.



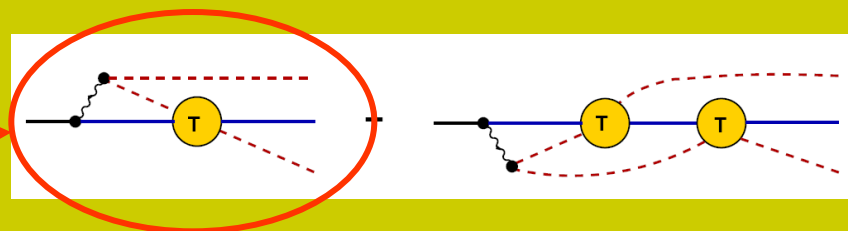
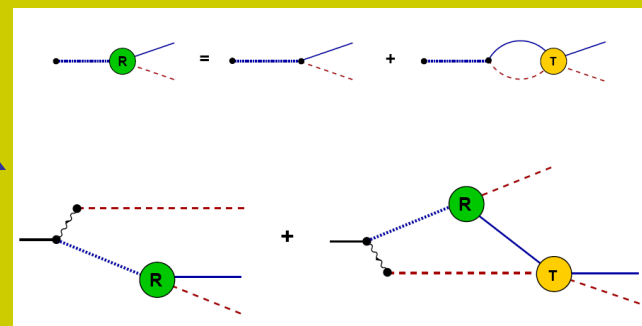
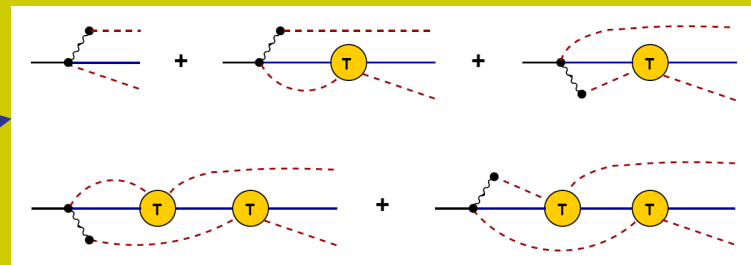
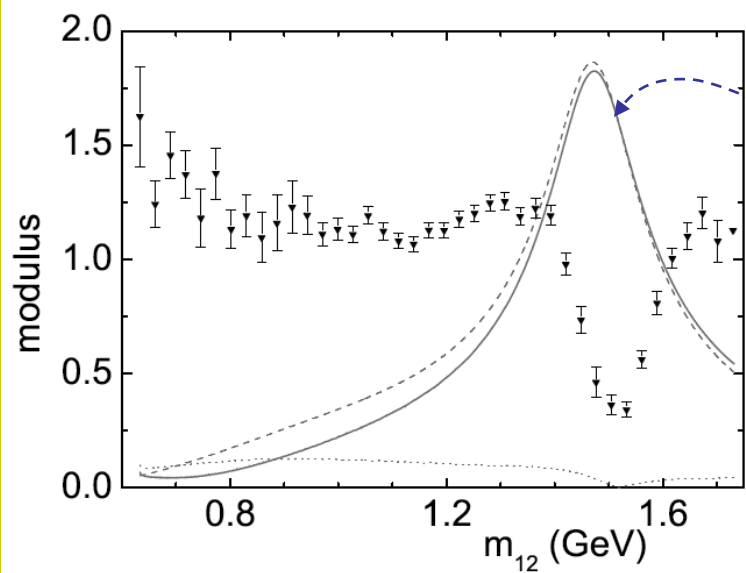
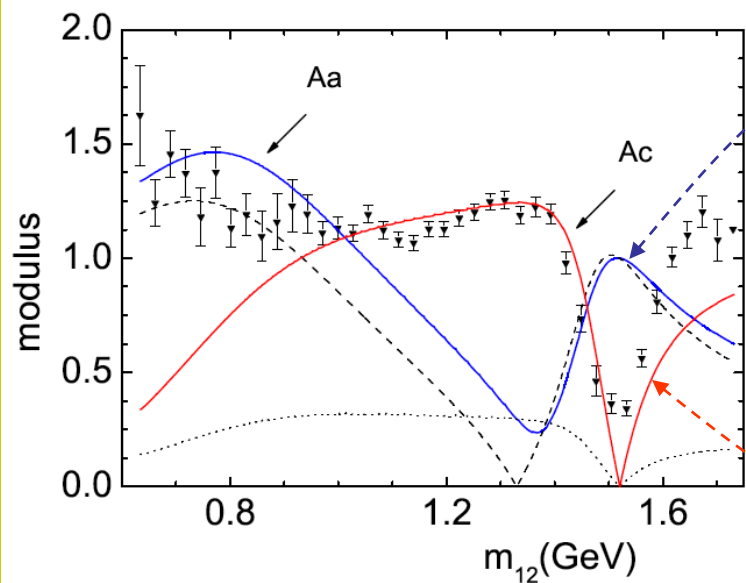
# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$



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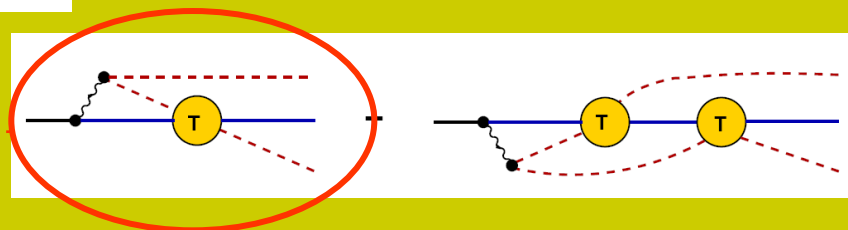
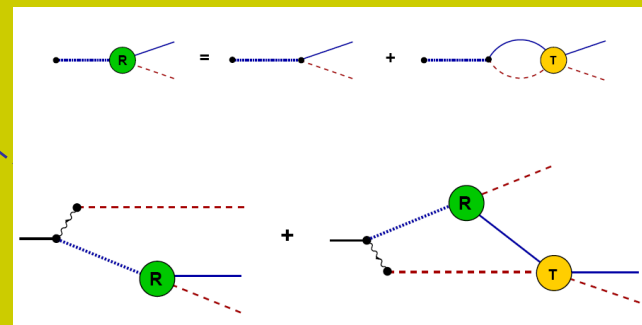
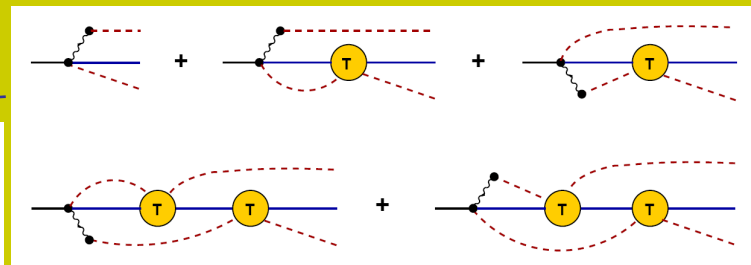
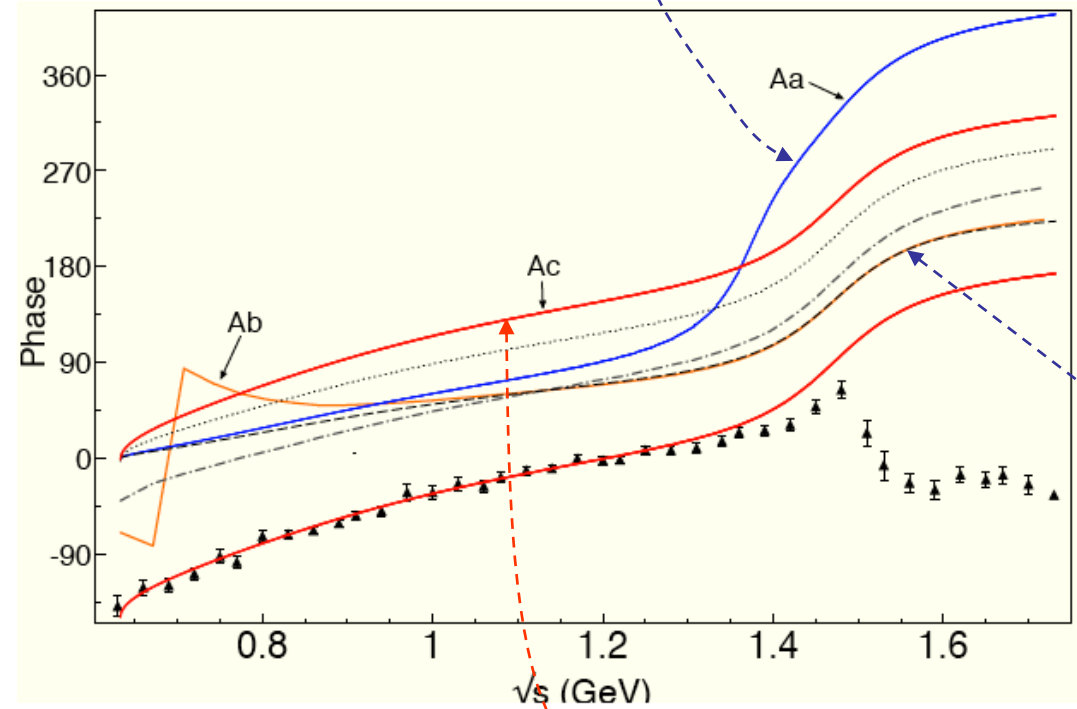


# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

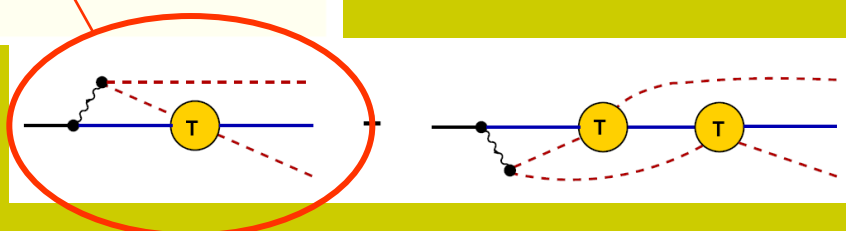
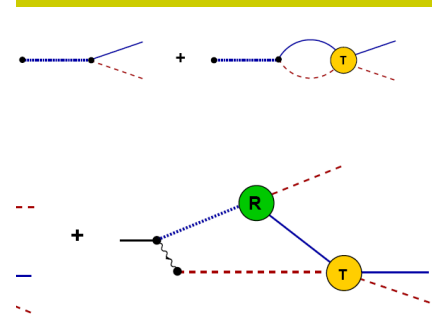
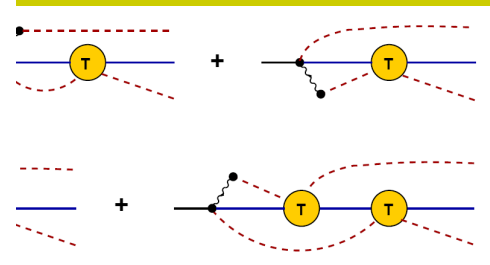
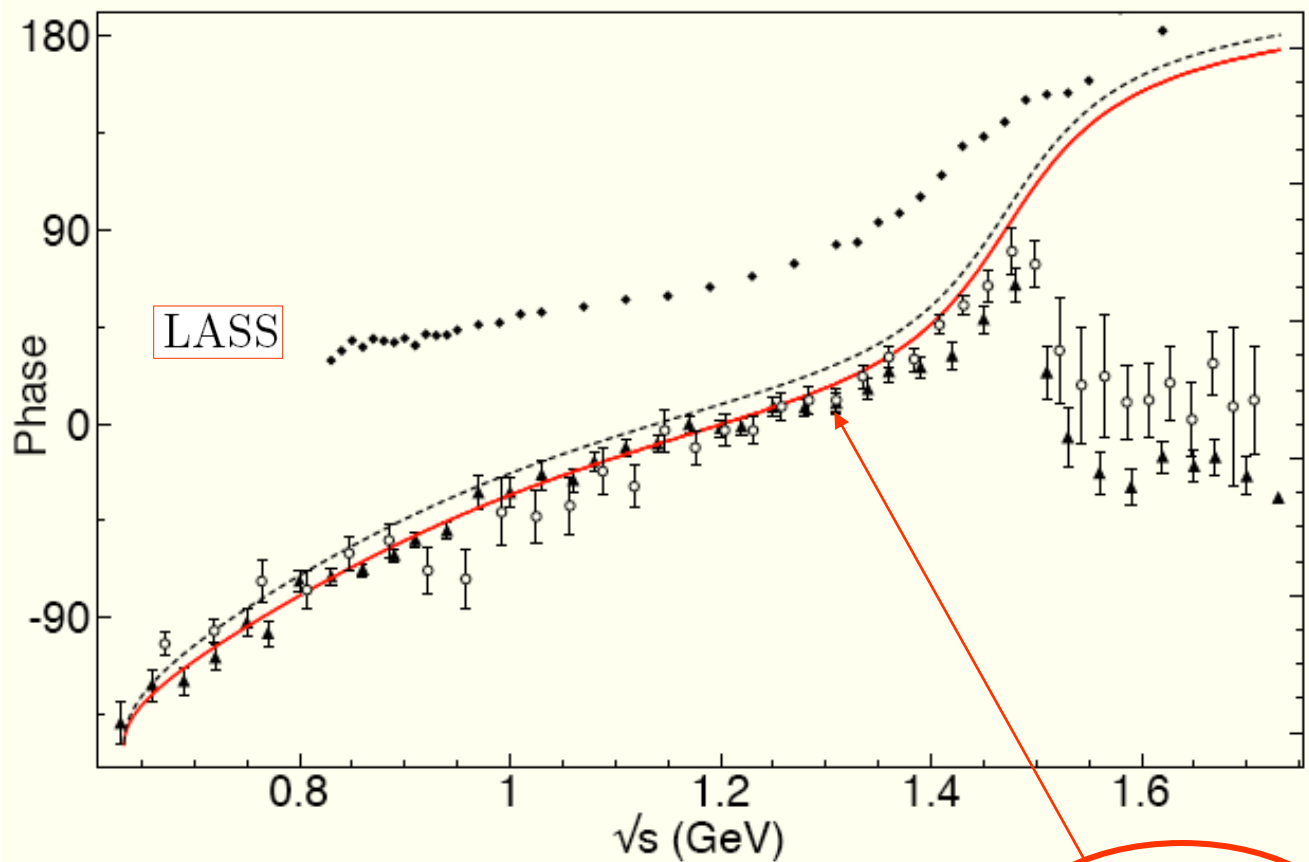




# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$



# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$





# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

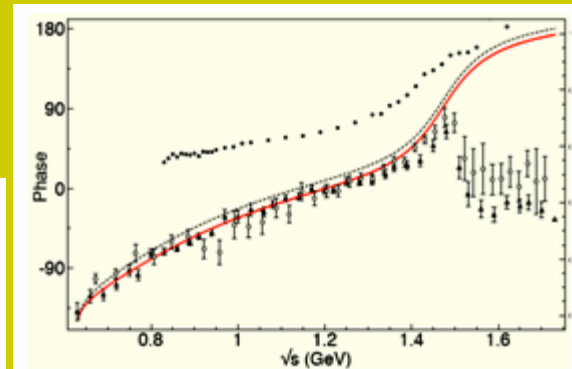
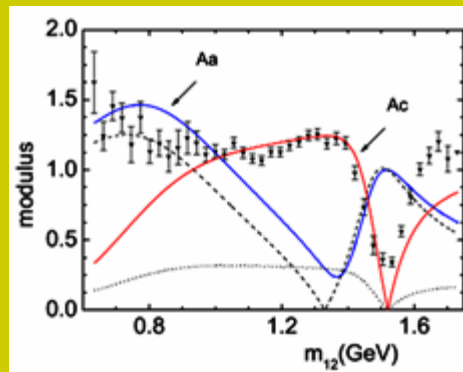
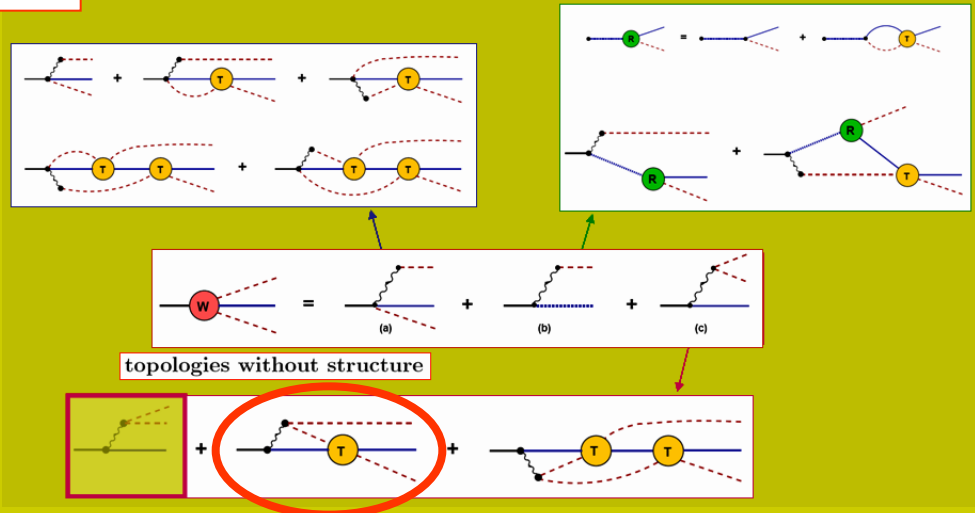
## conclusions

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- strong series converges rapidly
- dynamics:
  - leading mechanism identified
  - vector weak vertex
  - no tree contribution
- Watson's theorem does not apply
  - two phases around





# Towards three-body unitarity in $D^+ \rightarrow K^- \pi^+ \pi^+$

## conclusions

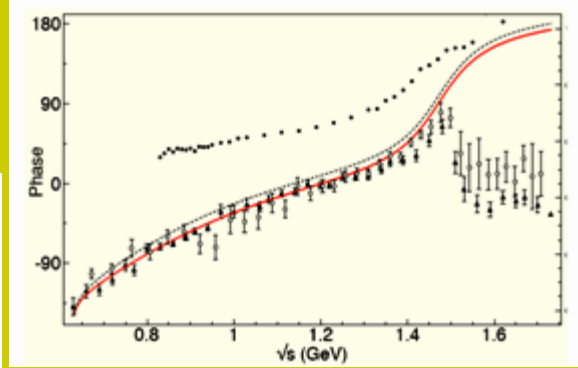
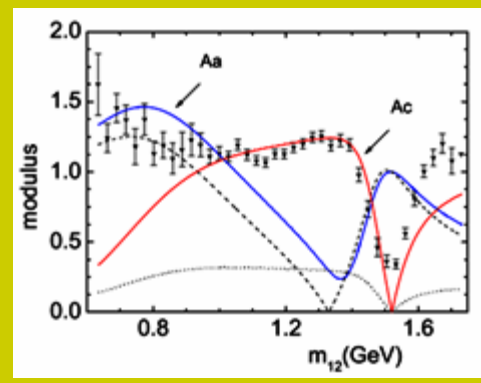
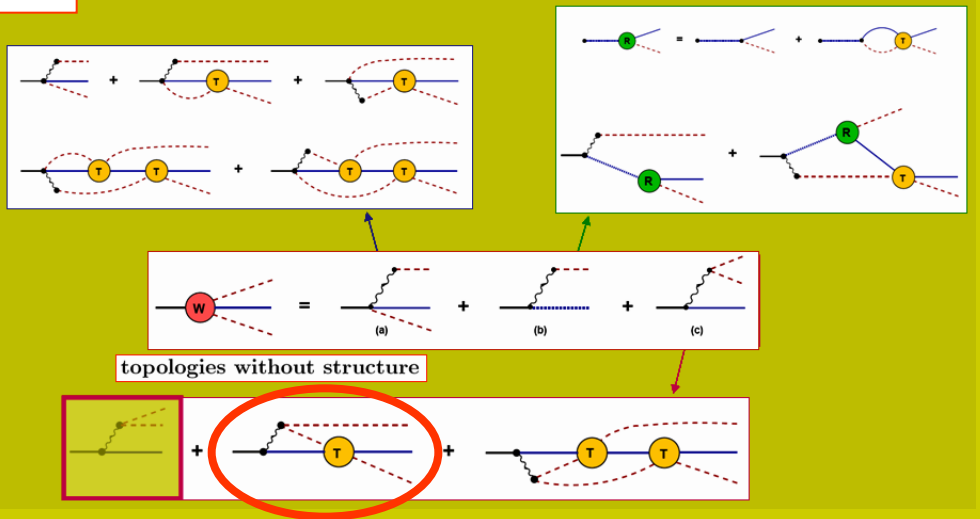
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**BUT...**



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- dynamics:
  - leading mechanism identified
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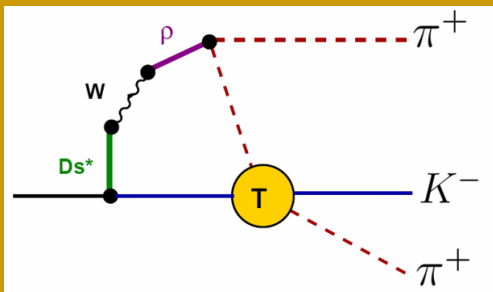


# research programme

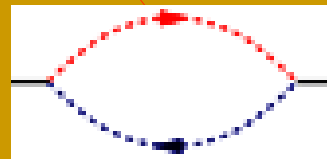
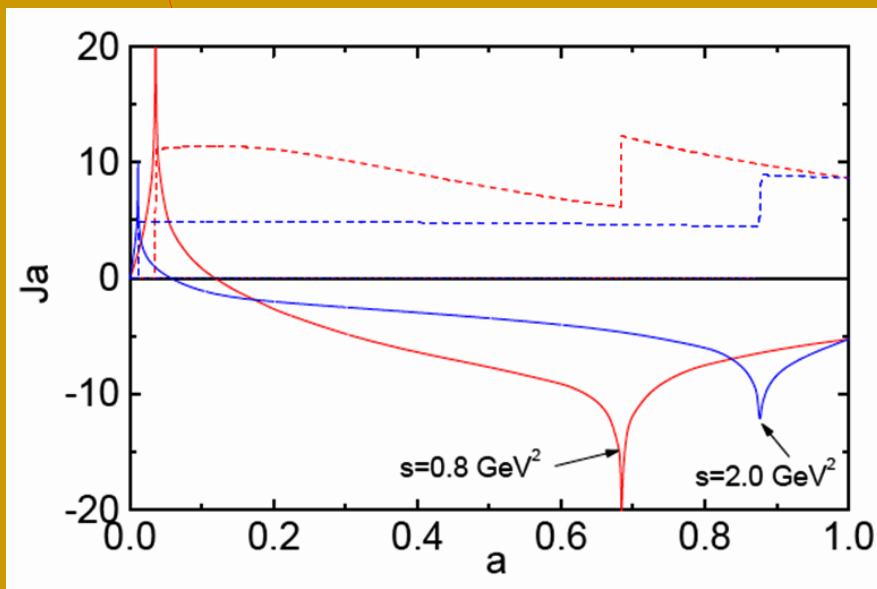
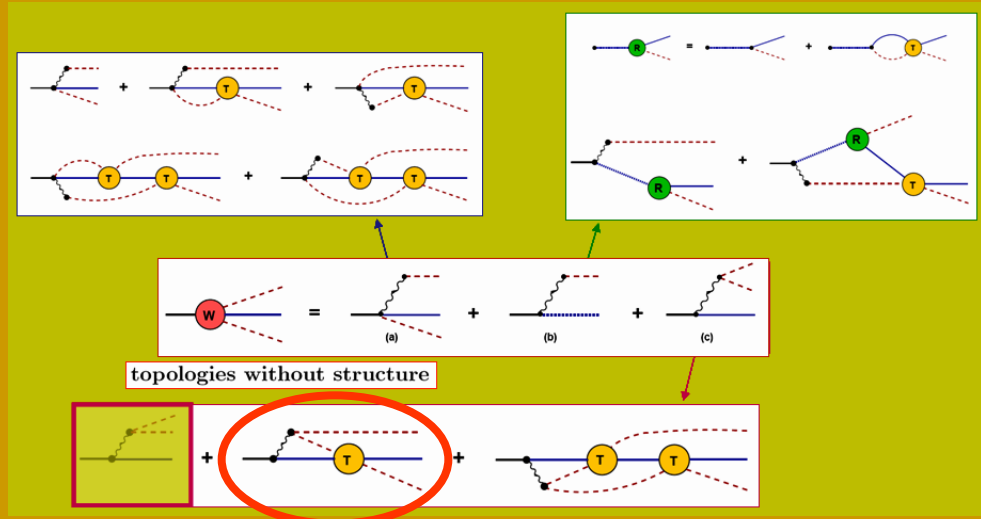
$$D^+ \rightarrow K^- \pi^+ \pi^+$$

## recent progress

Patricia C. Magalhães → Ph.D.



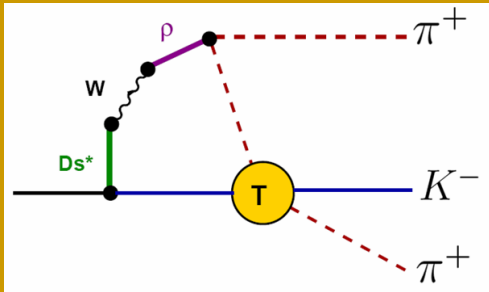
rich phase structure



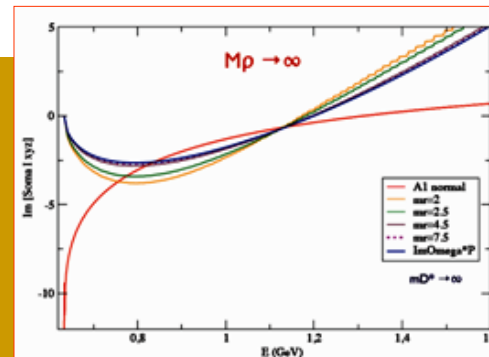
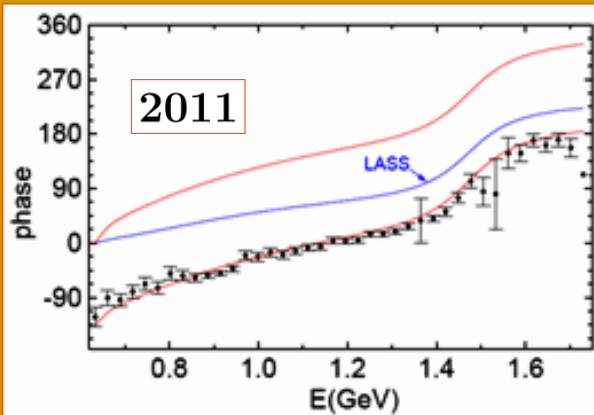
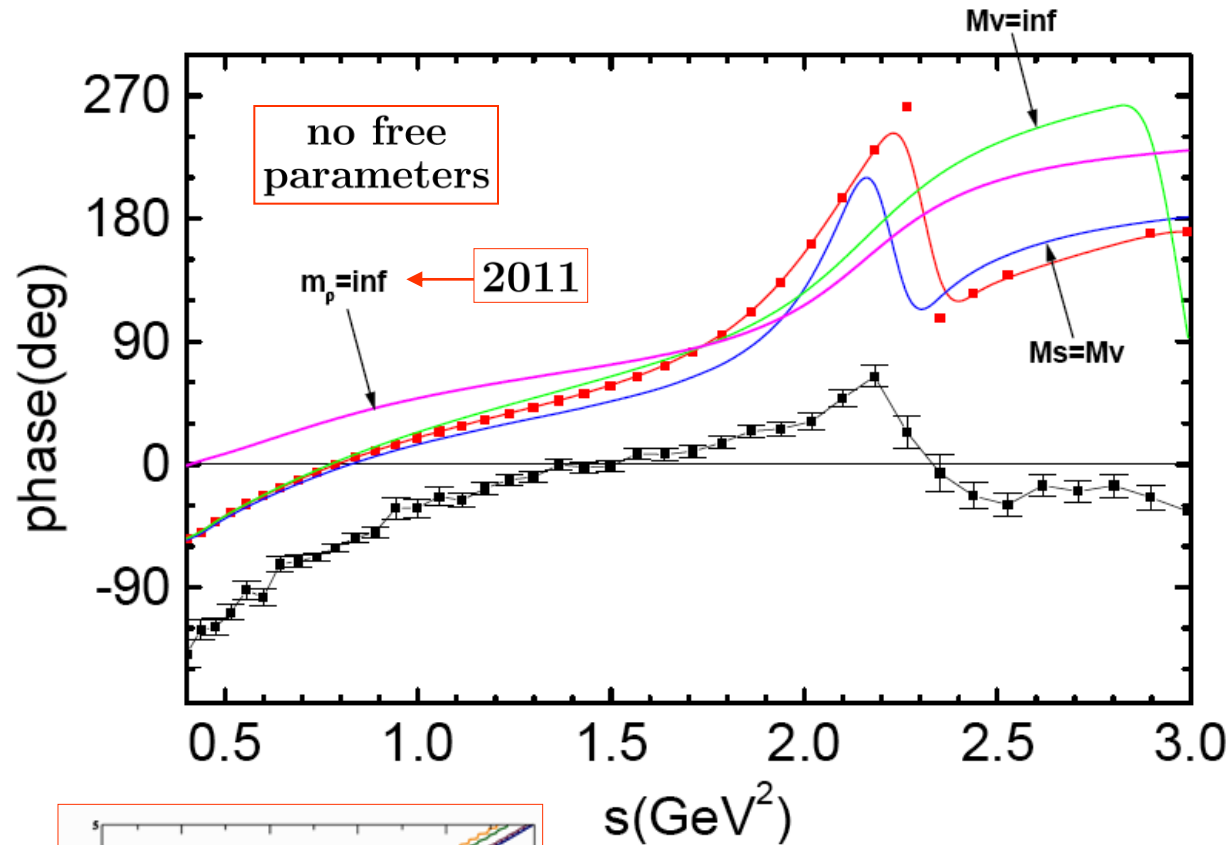
# research programme

## recent progress

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



rich phase structure

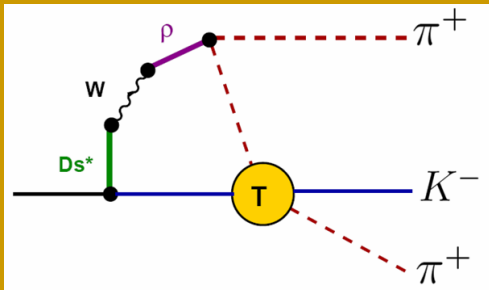




# research programme

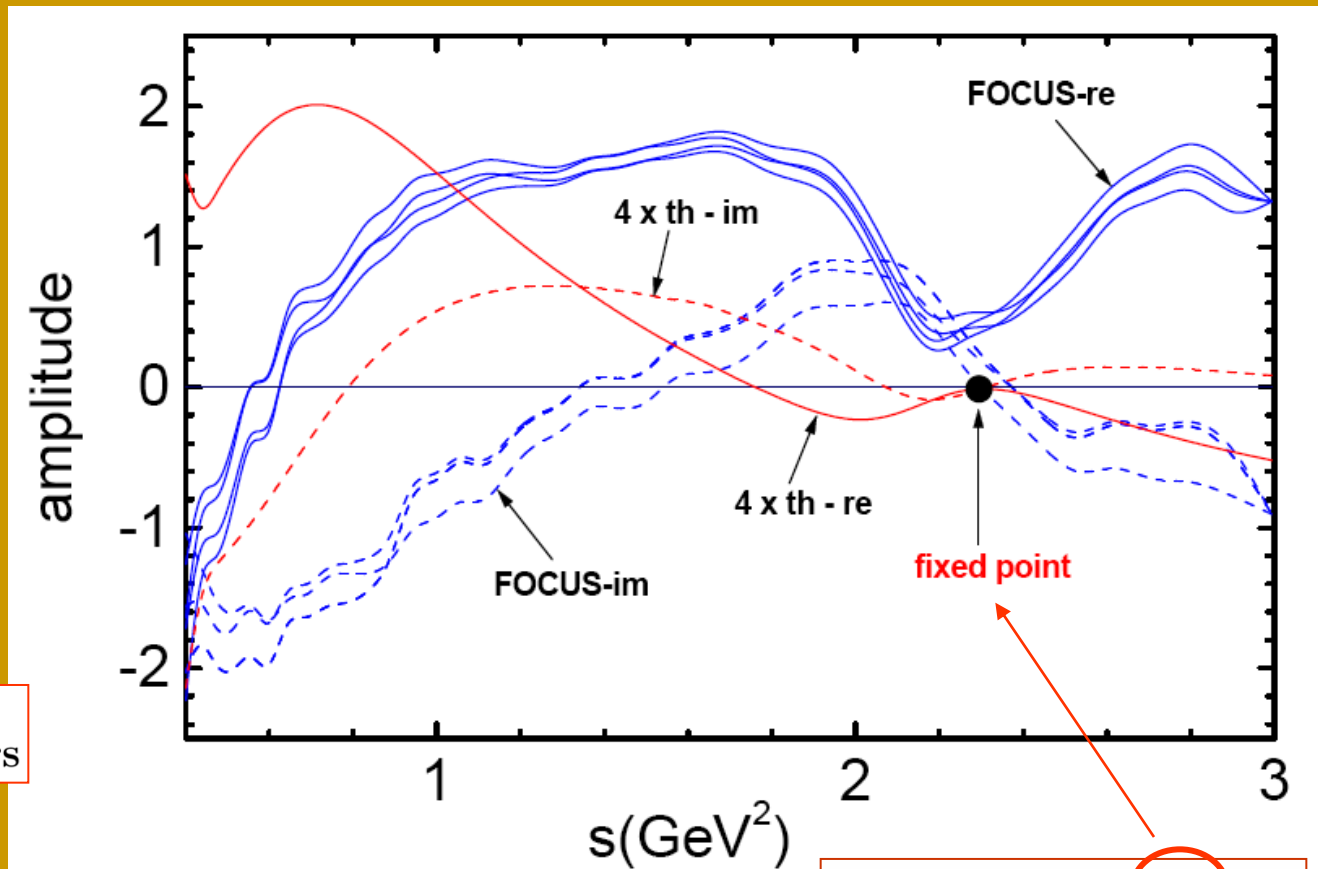
## recent progress

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



rich phase structure

no free parameters



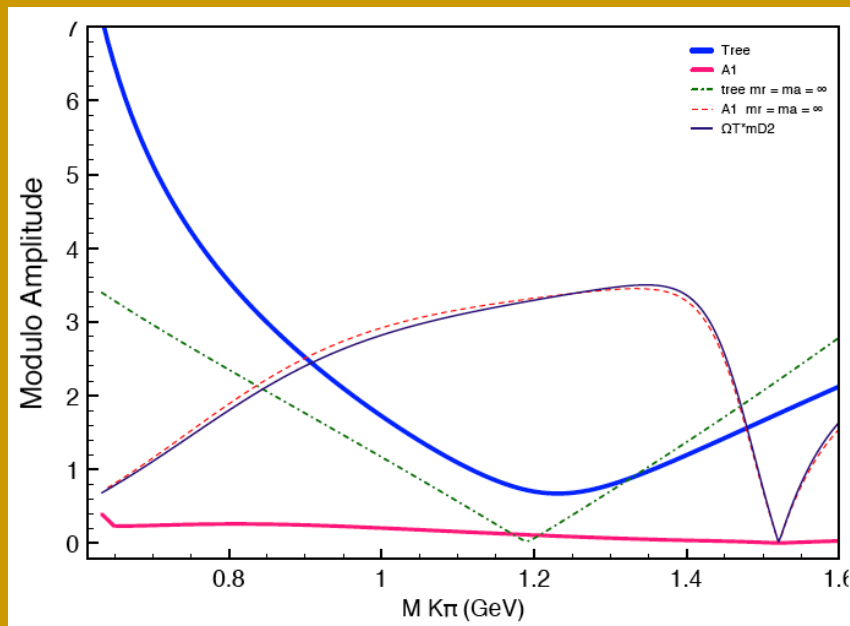
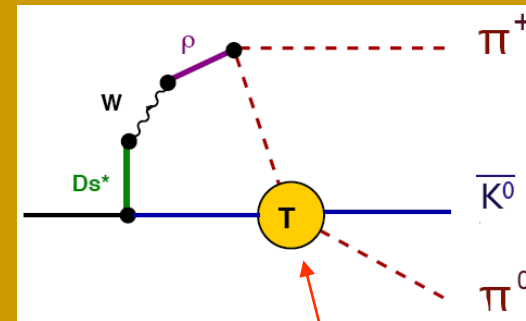
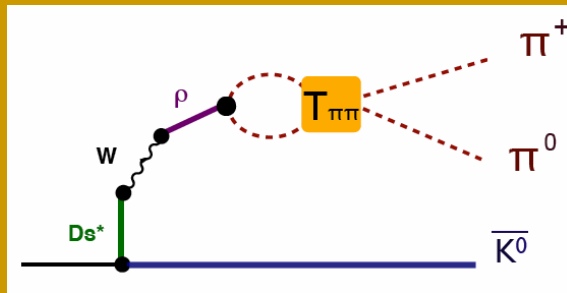
$$\text{[Yellow Square Box]} = \frac{\mathcal{K}}{1 + \bar{\Omega} \mathcal{K}}$$



# research programme

## recent progress

$$D^+ \rightarrow K^0 \pi^0 \pi^+$$



empirical  
amplitude

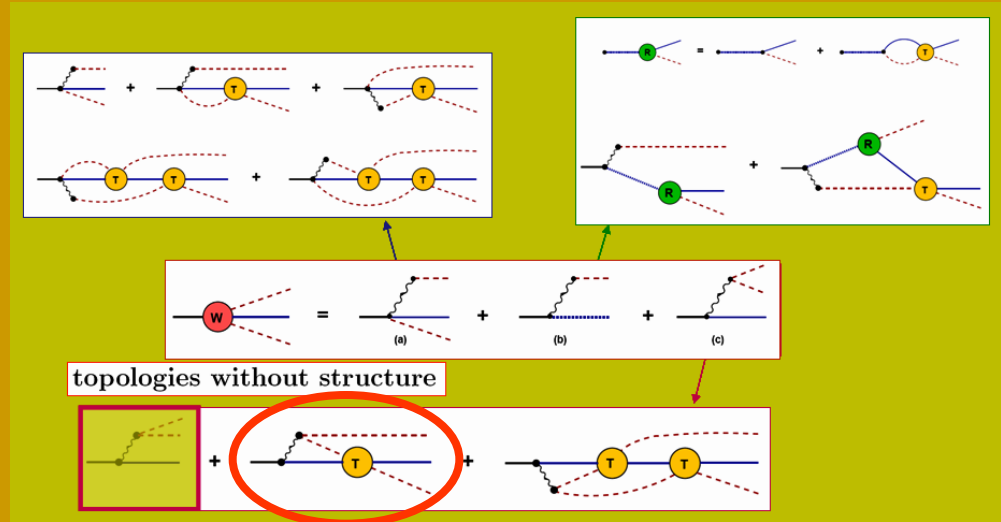
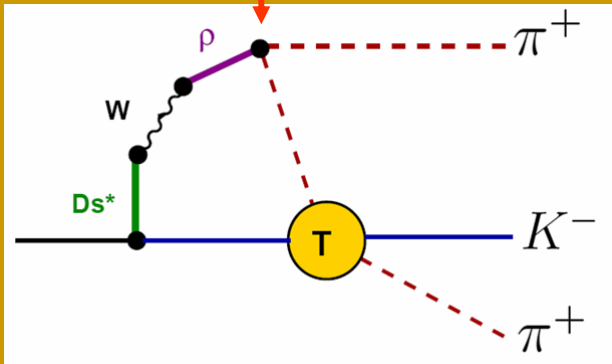
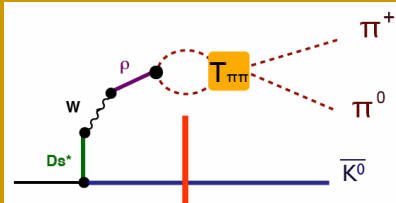
get rid of  
model  
dependence



# research programme

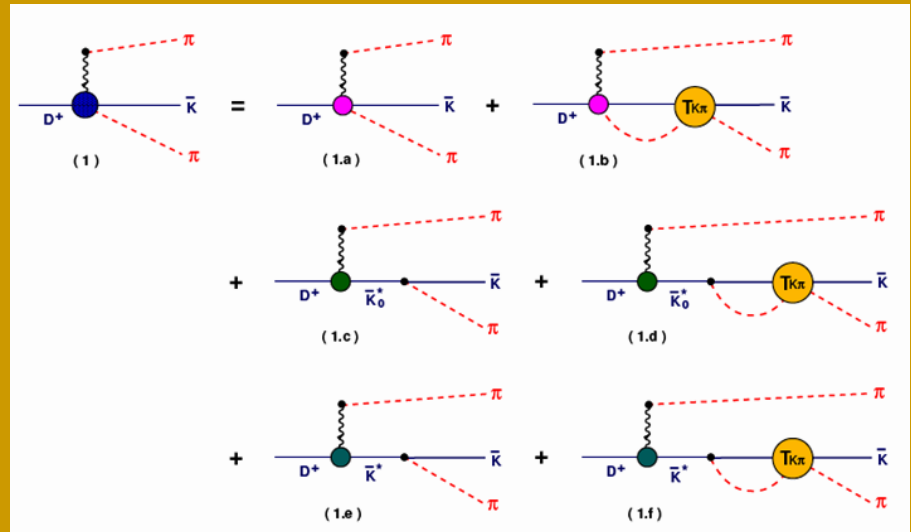
## next steps

$$D^+ \rightarrow K^- \pi^+ \pi^+$$



topologies without structure

QM - everything  
that is not forbidden,  
is compulsory





subject is crawling ...

