

# Measurement of the angle $\gamma$ ( $\phi_3$ ) at the $e^+e^-$ B-factories

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# The angle $\gamma$

- CKM is 3x3 unitary matrix  $\rightarrow$  4 parameters (after ad hoc choice of quark field phases): 3 real and 1 CP violating phase

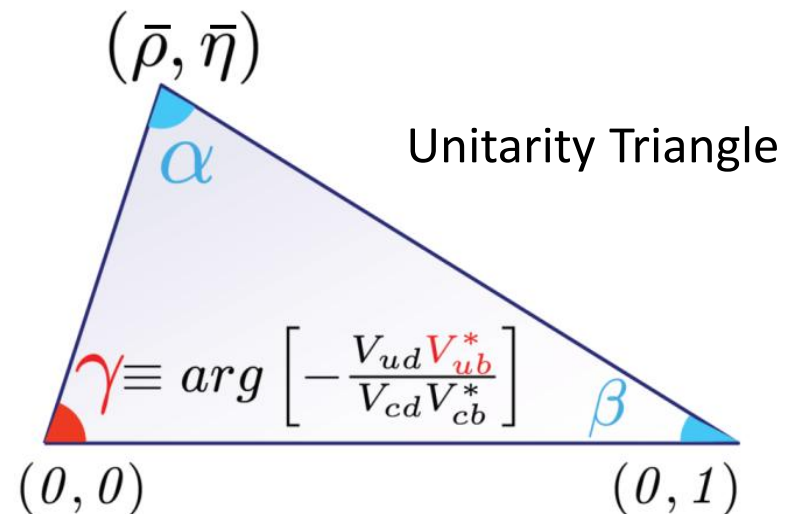
[Wolfenstein parametrization]

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

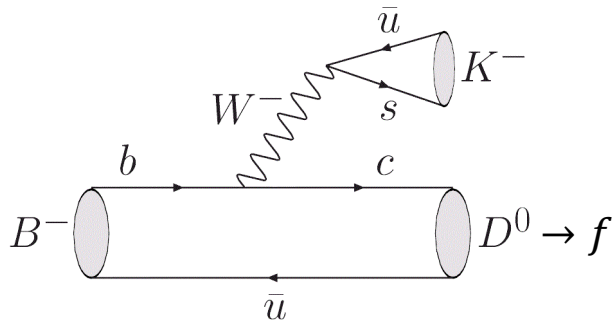
$\rightarrow |V_{ub}|e^{-i\gamma}$

- From the unitarity of V:

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

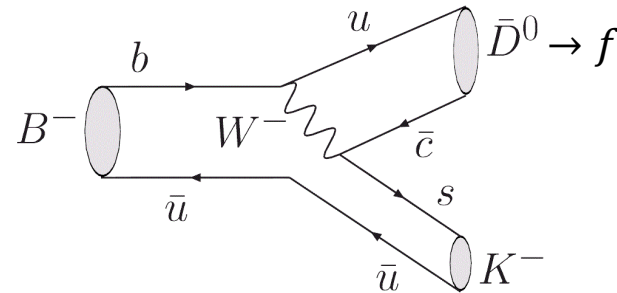


# Measurement of $\gamma$ with $B \rightarrow DK$



color allowed

$$A_1 \propto V_{cb} V_{us}^* \sim A\lambda^3$$



color suppressed

$$A_2 \propto V_{ub} V_{cs}^* \sim A\lambda^3(\rho + i\eta)$$

- $\gamma$  is measured in the interference of the two amplitudes

$$|A_{tot}|^2 = |A_1 + A_2|^2 = |A_1|^2 + |A_2|^2 + \underline{2|A_1||A_2| \cos \phi}$$

$$r_b \equiv \left| \frac{A(B^+ \rightarrow D^0 K^+)}{A(B^+ \rightarrow \bar{D}^0 K^+)} \right| \sim 0.1$$

$$\phi = -\gamma + \delta$$

- unknowns:  $\gamma, r_b, \delta_b + \delta_D$

B hadronic parameters  
extracted with  $\gamma$

measured at charm  
factories

$\delta_b + \delta_D =$  strong phase  
from B and D decay

- theoretically clean

- most sensitive method to constrain  $\gamma$  at present

- similar principle applies to several other processes:  $B^0 \rightarrow D^{(*)+} \pi^-$ ,  
 $B^0 \rightarrow D^0 K^{(*)0}$ ,  $B_s \rightarrow D_s K$ , ...

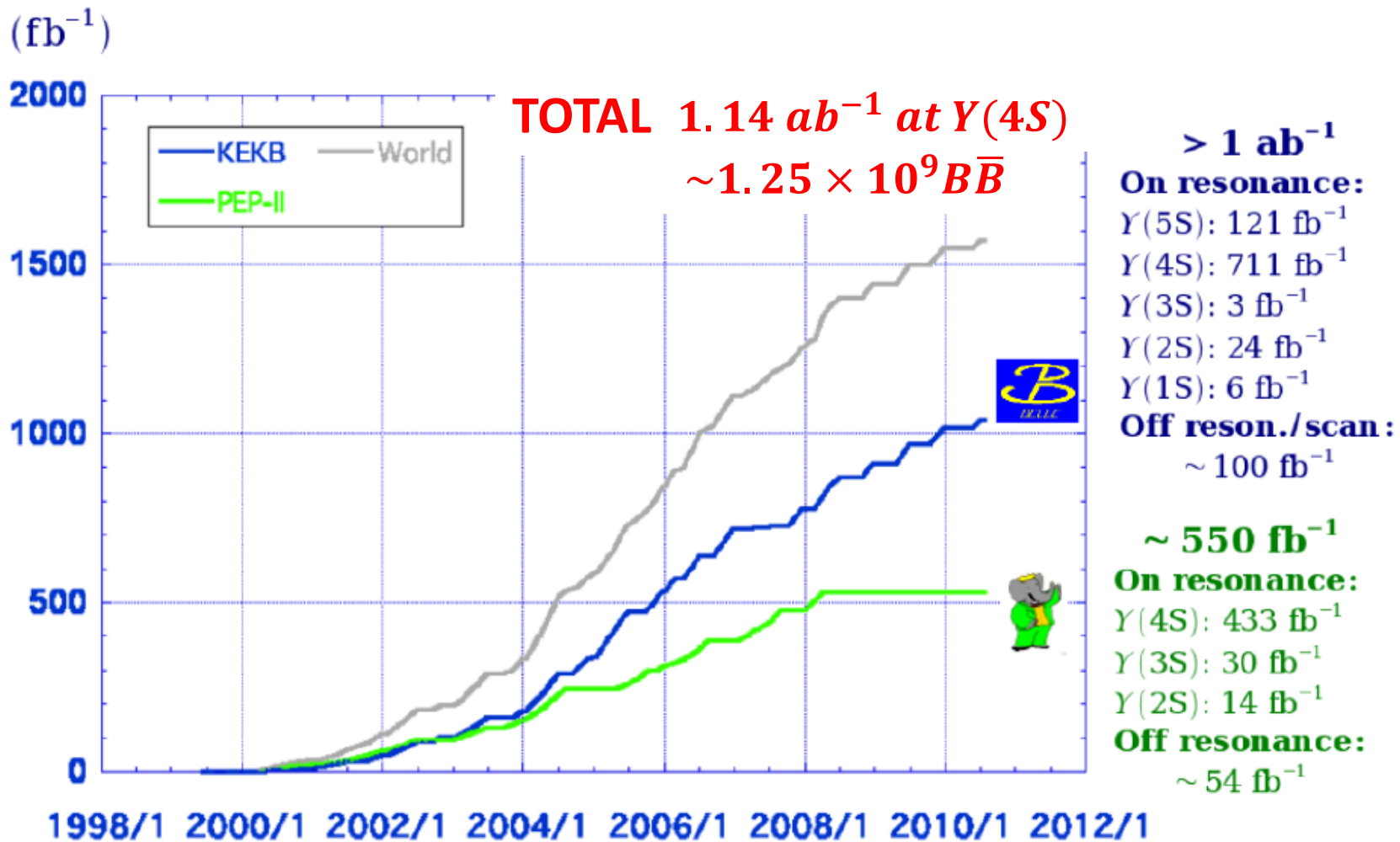
# Extraction of $\gamma$ using $B \rightarrow DK$ decays

## Different methods depending on D final state:

- **GLW** [M. Gronau, D. London, D. Wyler, PLB253,483 (1991); PLB 265, 172 (1991)]
  - $D^0$  to two-body **CP eigenstates**  $K^+K^-$ ,  $\pi^+\pi^-$  (even),  $K_S\pi^0$ ,  $K_S\omega$  (odd)
- **ADS** [D. Atwood, I. Dunietz, A. Soni, PRL 78, 3357 (1997)]
  - $D^0$  to **doubly Cabibbo suppressed decays**  $K^+\pi^-$ ,  $K^+\pi^-\pi^0$ , ...
- **GGSZ (Dalitz)** [D. Atwood et al., PRL78, 3257 (1997); A. Giri et al., PRD68, 054018 (2003)]
  - $D^0$  to **3-body decays**  $K_S\pi^+\pi^-$ ,  $K_S K^+K^-$ ,  $\pi^+\pi^-\pi^0$ , etc.
    - Dalitz plot fitted to determine how the strong phase of  $D^0$  decay amplitude varies over the Dalitz plane
    - model independent analysis
- Different B decays  $D^0K^\pm$ ,  $D^{*0}K^\pm$ ,  $D^0K^{*\pm}$  and flavour-tagged  $D^0K^{*0}$ . They depend on mode-dependent hadronic factors ( $r_b$ ,  $\delta_b$ )
- **Strategy:** combine as many channels as possible to improve the overall sensitivity

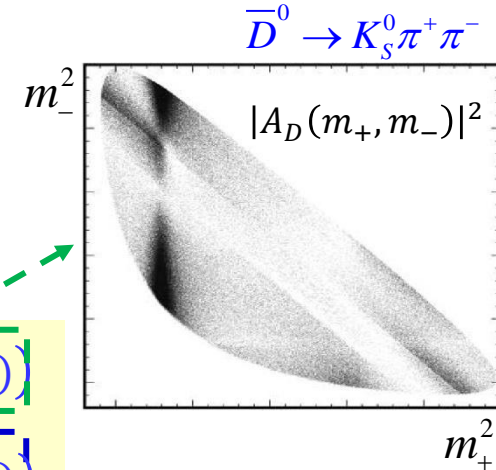
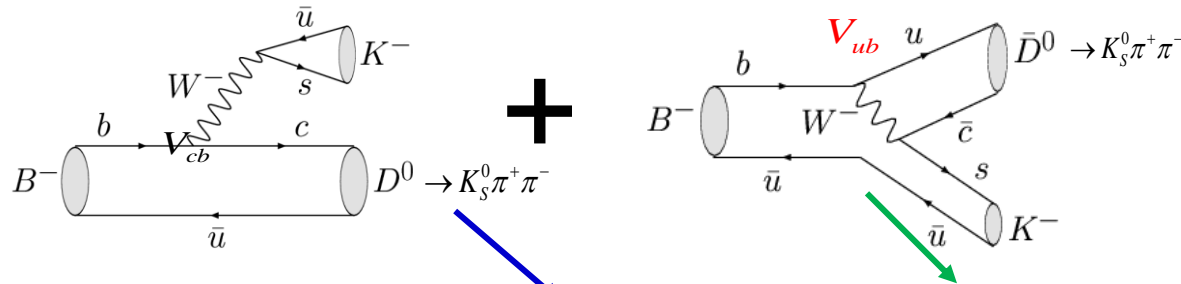
most  
powerful  
method  
nowadays

# The B-factories dataset



# The GGSZ method

- The interference varies as function of the position in the  $D^0$  Dalitz plot

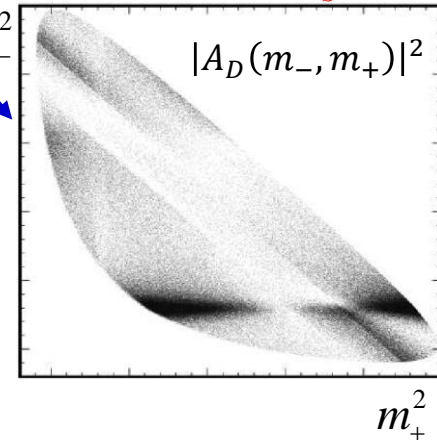


$$A_-(m_-^2, m_+^2) = |A(B^- \rightarrow D^0 K^-)| \left( A_D(m_-^2, m_+^2) + r_b e^{i(\delta_b - \gamma)} A_D(m_+^2, m_-^2) \right)$$

$$A_+(m_-^2, m_+^2) = |A(B^+ \rightarrow \bar{D}^0 K^+)| \left( A_D(m_+^2, m_-^2) + r_b e^{i(\delta_b + \gamma)} A_D(m_-^2, m_+^2) \right)$$

$$m_{\pm}^2 \equiv m^2(K_S^0 \pi^{\pm})^2$$

$D^0 \rightarrow K_S^0 \pi^+ \pi^-$

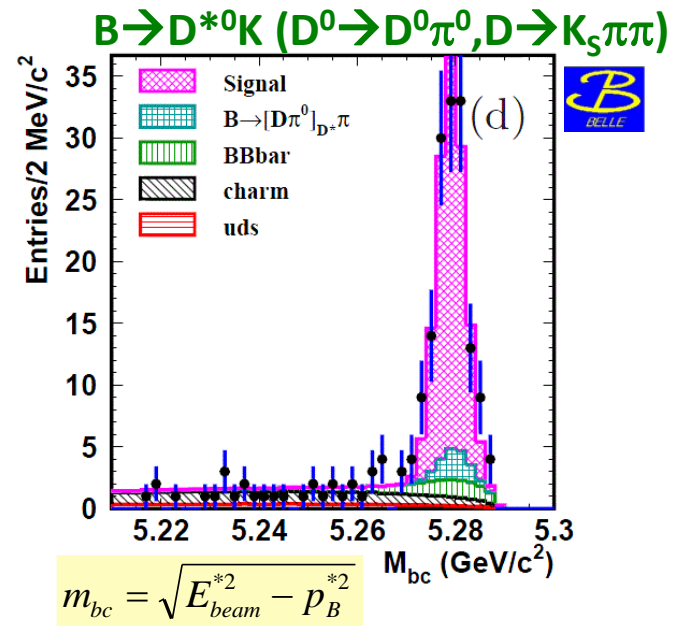
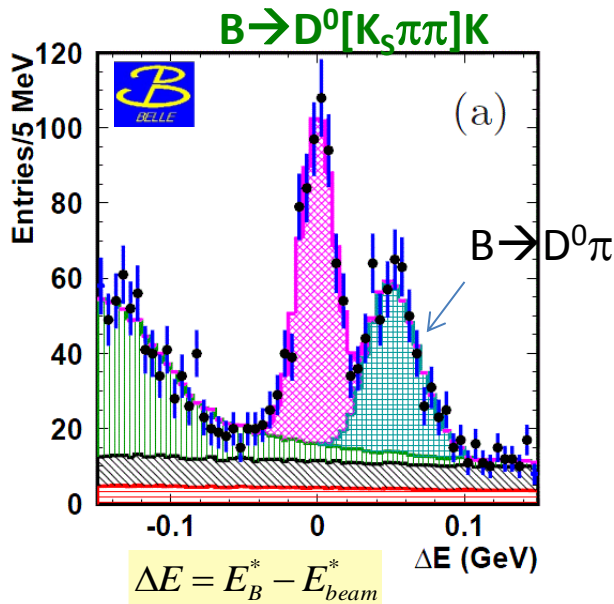


- $A_D(m_-^2, m_+^2)$  is measured with a Dalitz plot analysis of high statistics samples of flavour-tagged  $D^0$  and  $\bar{D}^0$
- The  $B^+$  and  $B^-$  yields are measured as a function of the position in the  $D^0$  Dalitz plot (ML fit)
- Unknowns:  $\gamma$ ,  $r_b$  and  $\delta_b$

# Reconstructed decay modes

Decays	$D^0 K^+$	$D^{*0} [D^0 \pi^0] K^+$	$D^{*0} [D^0 \gamma] K^+$	$D^0 K^{*+} [K_S \pi]$
$D^0 \rightarrow K_S \pi \pi$	 468M  657M	 468M  657M	 468M  657M	 468M  386M
$D^0 \rightarrow K_S K K$	 468M	 468M	 468M	 468M

Signal region defined by  $\Delta E$  and  $M_{bc}$



both BaBar and Belle fit signal vs Dalitz plot position using likelihood( $\Delta E, M_{bc}$ , event shape vars)

# Measurement of $x_{\pm}, y_{\pm}$

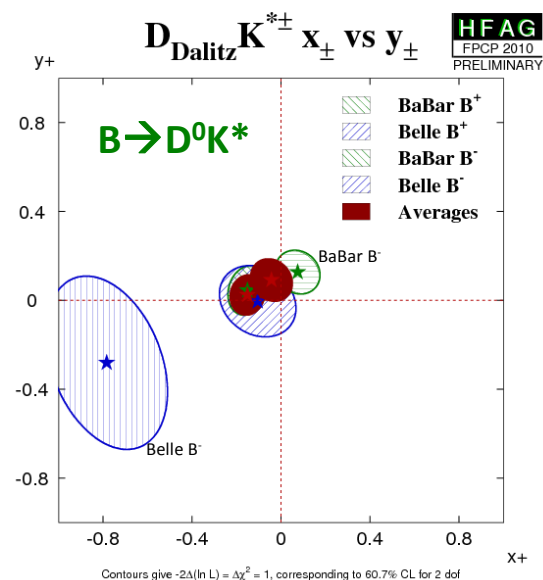
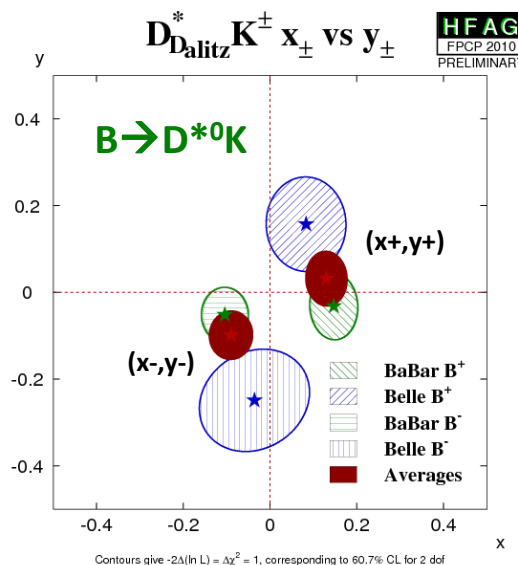
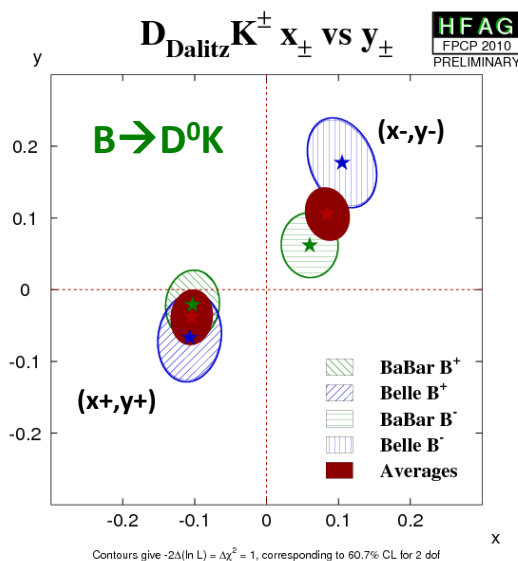
- Extract the *cartesian coordinates* instead of  $\gamma, r_b, \delta_b$  (likelihood unbiased and Gaussian-shaped using  $x, y$ )

$$\begin{aligned} x_{\mp} &= r_b \cos(\delta_b \mp \gamma) \\ y_{\mp} &= r_b \sin(\delta_b \mp \gamma) \end{aligned}$$

$$\begin{cases} \Gamma(B^+) \propto |f_+|^2 + (x_+^2 + y_+^2)|f_-|^2 + 2x_+ \operatorname{Re}(f_+ f_-^*) + 2y_+ \operatorname{Im}(f_+ f_-^*) \\ \Gamma(B^-) \propto |f_-|^2 + (x_-^2 + y_-^2)|f_+|^2 + 2x_- \operatorname{Re}(f_- f_+^*) + 2y_- \operatorname{Im}(f_- f_+^*) \end{cases}$$

$$f_{\mp} \equiv A_D(m_{\mp}^2, m_{\pm}^2)$$

$(x_{\pm}, y_{\pm})$  4 variables, 3 indep.  
 $x_+^2 + y_+^2 = x_-^2 + y_-^2$

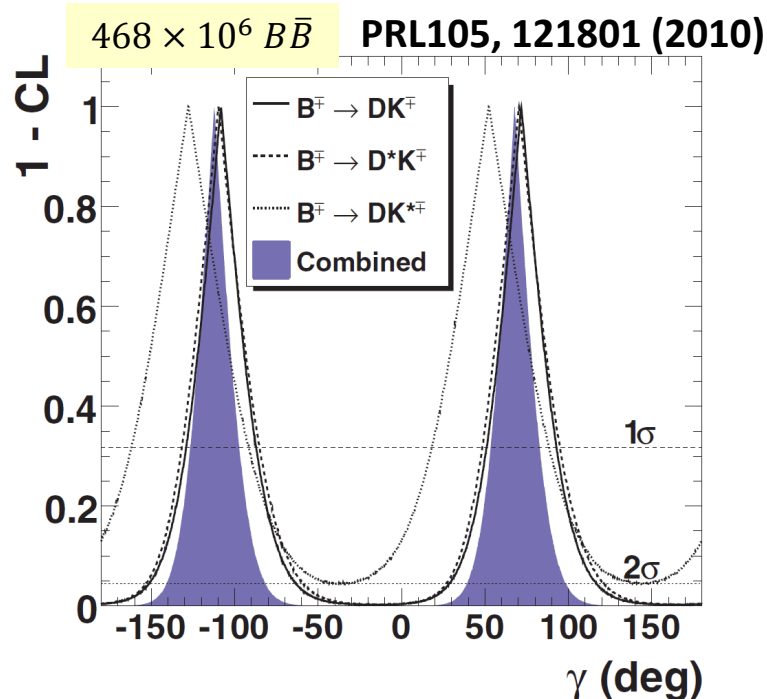


the contours do not include the Dalitz model errors



# From $x_{\pm}, y_{\pm}$ to $\gamma$ : BaBar

- Combine  $B \rightarrow DK, B \rightarrow D^*K, B \rightarrow DK^*$  with  $D \rightarrow K_S \pi^+ \pi^-$  and  $D \rightarrow K_S K^+ K^-$
- Use frequentist method to derive the physical parameters  $\gamma, r_b, \delta_b$  from  $(x_{\pm}, y_{\pm})$



3.5 $\sigma$  stat significance of CPV

Parameter	total error	exp. sys.	Dalitz model sys.
	68.3% C.L.		95.4% C.L.
$\gamma$ ( $^\circ$ )	$68_{-14}^{+15}$	{4, 3}	[39, 98]
$r_B$ (%)	$9.6 \pm 2.9$	{0.5, 0.4}	[3.7, 15.5]
$r^* B$ (%)	$13.3_{-3.9}^{+4.2}$	{1.3, 0.3}	[4.9, 21.5]
$\kappa r_s$ (%)	$14.9_{-6.2}^{+6.6}$	{2.6, 0.6}	<28.0
$\delta_B$ ( $^\circ$ )	$119_{-20}^{+19}$	{3, 3}	[75, 157]
$\delta_B^*$ ( $^\circ$ )	$-82 \pm 21$	{5, 3}	[-124, -38]
$\delta_s$ ( $^\circ$ )	$111 \pm 32$	{11, 3}	[42, 178]

$$\gamma = (68 \pm 14 \pm 4 \pm 3)^\circ \pmod{180^\circ}$$

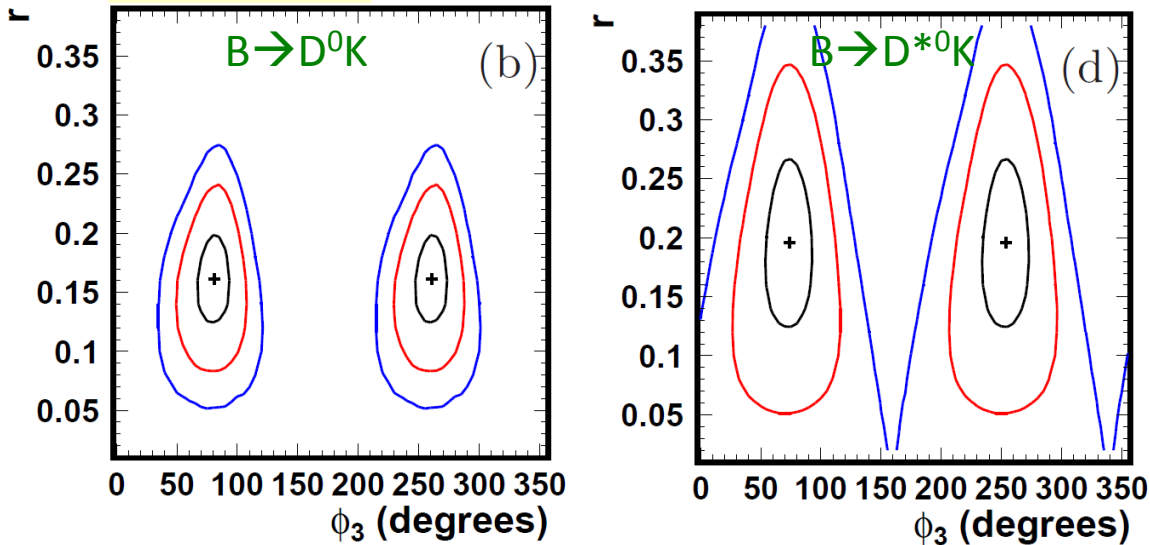
stat    syst    model



# From $x_{\pm}, y_{\pm}$ to $\gamma$ : Belle

- Combine  $B \rightarrow D^0 K$  and  $B \rightarrow D^{*0} K$  with  $D \rightarrow K_S \pi^+ \pi^-$
- Use frequentist method to derive the physical parameters  $\gamma$ ,  $r_b$ ,  $\delta_b$  from  $(x_{\pm}, y_{\pm})$

$657 \times 10^6 B\bar{B}$  PRD81, 112002 (2010)



combined  $3.5\sigma$  stat significance of CPV

	$B^+ \rightarrow DK^+$ mode
$\gamma$	$(80.8^{+13.1}_{-14.8} \pm 5.0 \pm 8.9)^\circ$
$r_b$	$0.161^{+0.040}_{-0.038} \pm 0.011^{+0.050}_{-0.010}$
$\delta_b$	$(137.4^{+13.0}_{-15.7} \pm 4.0 \pm 22.9)^\circ$

	$B^+ \rightarrow D^* K^+$ mode
$\gamma$	$(73.9^{+18.9}_{-20.2} \pm 4.2 \pm 8.9)^\circ$
$r_b^*$	$0.196^{+0.073}_{-0.072} \pm 0.013^{+0.062}_{-0.012}$
$\delta_b^*$	$(341.7^{+18.6}_{-20.9} \pm 3.2 \pm 22.9)^\circ$

$$\gamma = (78.4^{+10.8}_{-11.6} \pm 3.6 \pm 8.9)^\circ \pmod{180^\circ}$$

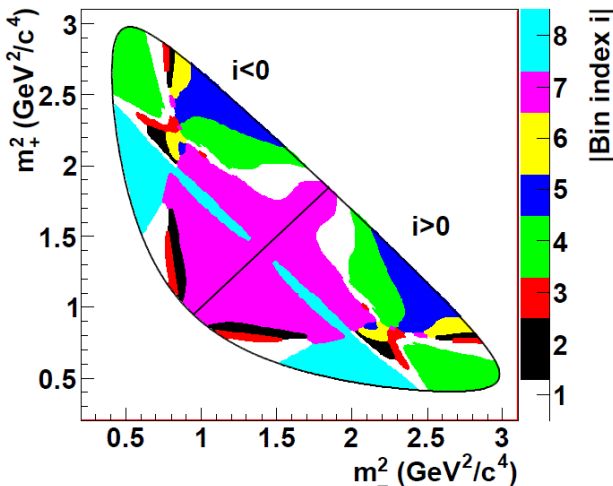
stat            syst            model



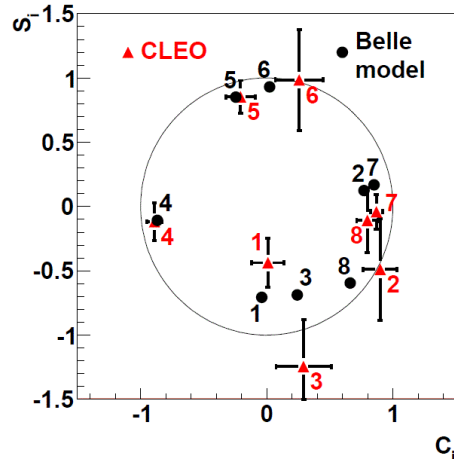
# Model independent analysis

Model-independent measurement of  $\gamma$ .

Proposed by A. Giri et al. [Phys Rev. D68 054018 (2003)]. Pioneered by Belle.



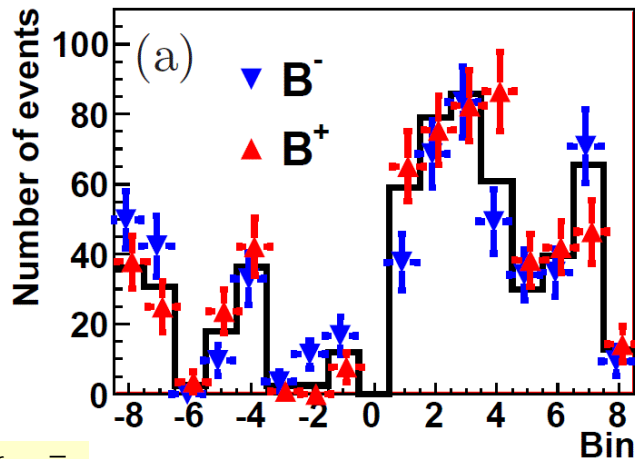
- divide the  $D \rightarrow K_S \pi \pi$  Dalitz plot in 2k bins (symmetric w.r.t. the  $m_+^2$  vs  $m_-^2$  axis)
- express the  $B^\pm \rightarrow DK^\pm$  yields in each bin  $i$  in terms of  $x_\pm, y_\pm$  and 2 parameters  $c_i, s_i$
- $c_i, s_i$  are measured by CLEO exploiting the quantum coherence in  $\psi(3770) \rightarrow \bar{D}^0 D^0$
- extract  $x_\pm, y_\pm$  from ML fit to  $B^\pm \rightarrow DK^\pm$  yields in all bins



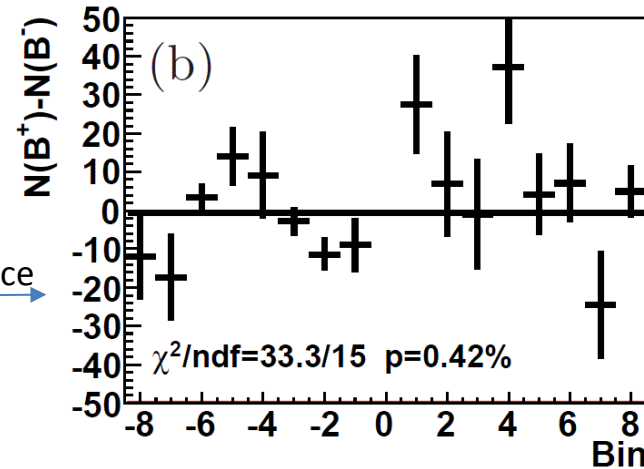
$$N_i^\pm = h_B \left[ K_{\pm i} + r_b^2 K_{\mp i} + 2\sqrt{K_i K_{-i}} (x_\pm c_i \pm y_\pm s_i) \right]$$

$N_i^\pm$  is labeled as  $B^\pm \rightarrow DK^\pm$  yields.  
 $K_{\pm i}$  is labeled as from flav.-tagged  $D \rightarrow K_S \pi \pi$ .  
 $x_\pm c_i$  is labeled as extracted from fit to the  $B^\pm$  yields.  
 $y_\pm s_i$  is labeled as measured by CLEO [PRD82, 112006 (2010)].

# Model independent measurement

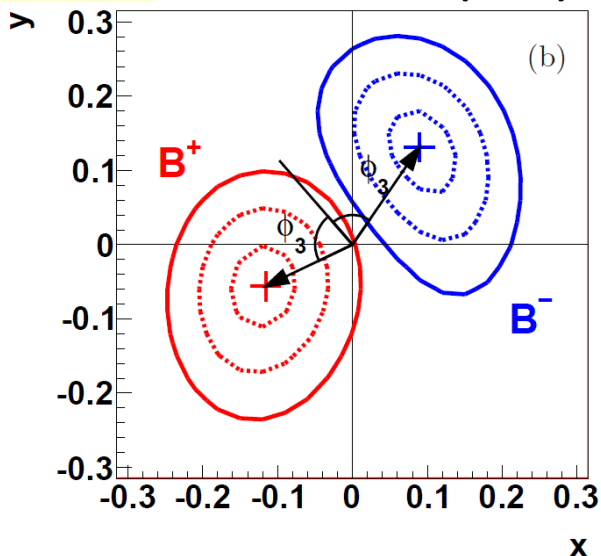


difference



$772 \times 10^6 B\bar{B}$

PRD85, 112014 (2012)



$$\gamma = \left( 77.3_{-14.9}^{+15.1} \pm 4.1 \pm 4.3 \right)^\circ \pmod{180^\circ}$$

stat      exp sys       $c_i, s_i$  errors

$$r_b = 0.145 \pm 0.030 \pm 0.010 \pm 0.011$$

$$\delta_b = (129.9 \pm 15.0 \pm 3.8 \pm 4.7)^\circ$$

uncertainty in  $c_i, s_i$  can be reduced at BESIII and possible next generation charm-tau factories

# GLW method

- $D^0$  to  $K^+K^-$ ,  $\pi^+\pi^-$  (CP+) and  $K_s\pi^0$ ,  $K_s\omega$ ,  $K_s\phi$  (CP-)
- measure  $B^+$  and  $B^-$  yields to determine the GLW observables:

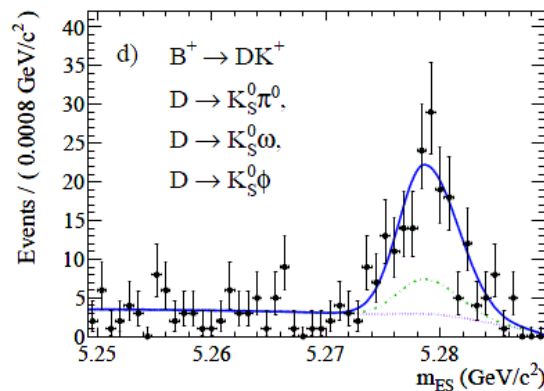
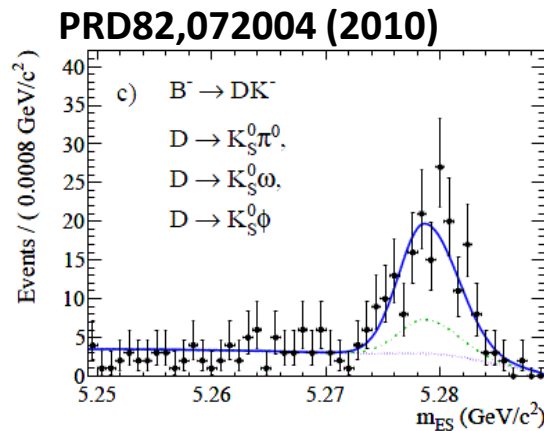
$$R_{CP\pm} \equiv \frac{\Gamma(B^- \rightarrow D_{CP\pm}^0 K^-) + \Gamma(B^+ \rightarrow D_{CP\pm}^0 K^+)}{2\Gamma(B^- \rightarrow D^0 K^-)} = 1 \pm 2r_b \cos \gamma \cos \delta_b + r_b^2$$

$$A_{CP\pm} \equiv \frac{\Gamma(B^- \rightarrow D_{CP\pm}^0 K^-) - \Gamma(B^+ \rightarrow D_{CP\pm}^0 K^+)}{\Gamma(B^- \rightarrow D_{CP\pm}^0 K^-) + \Gamma(B^+ \rightarrow D_{CP\pm}^0 K^+)} = \pm 2r_b \sin \gamma \sin \delta_b / R_{CP\pm}$$

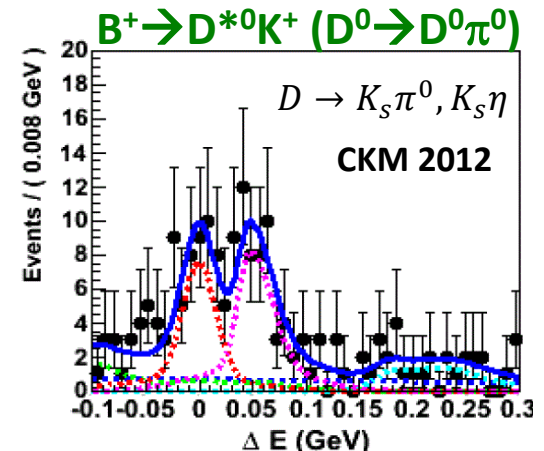
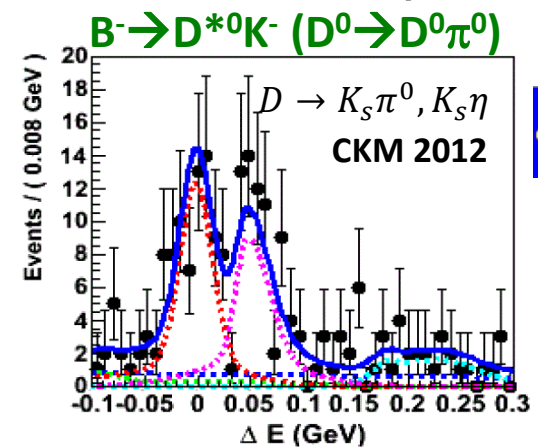
- 4 observables, 3 independent unknowns:  $\gamma$ ,  $\delta_b$ ,  $r_b$

# GLW reconstructed decay modes

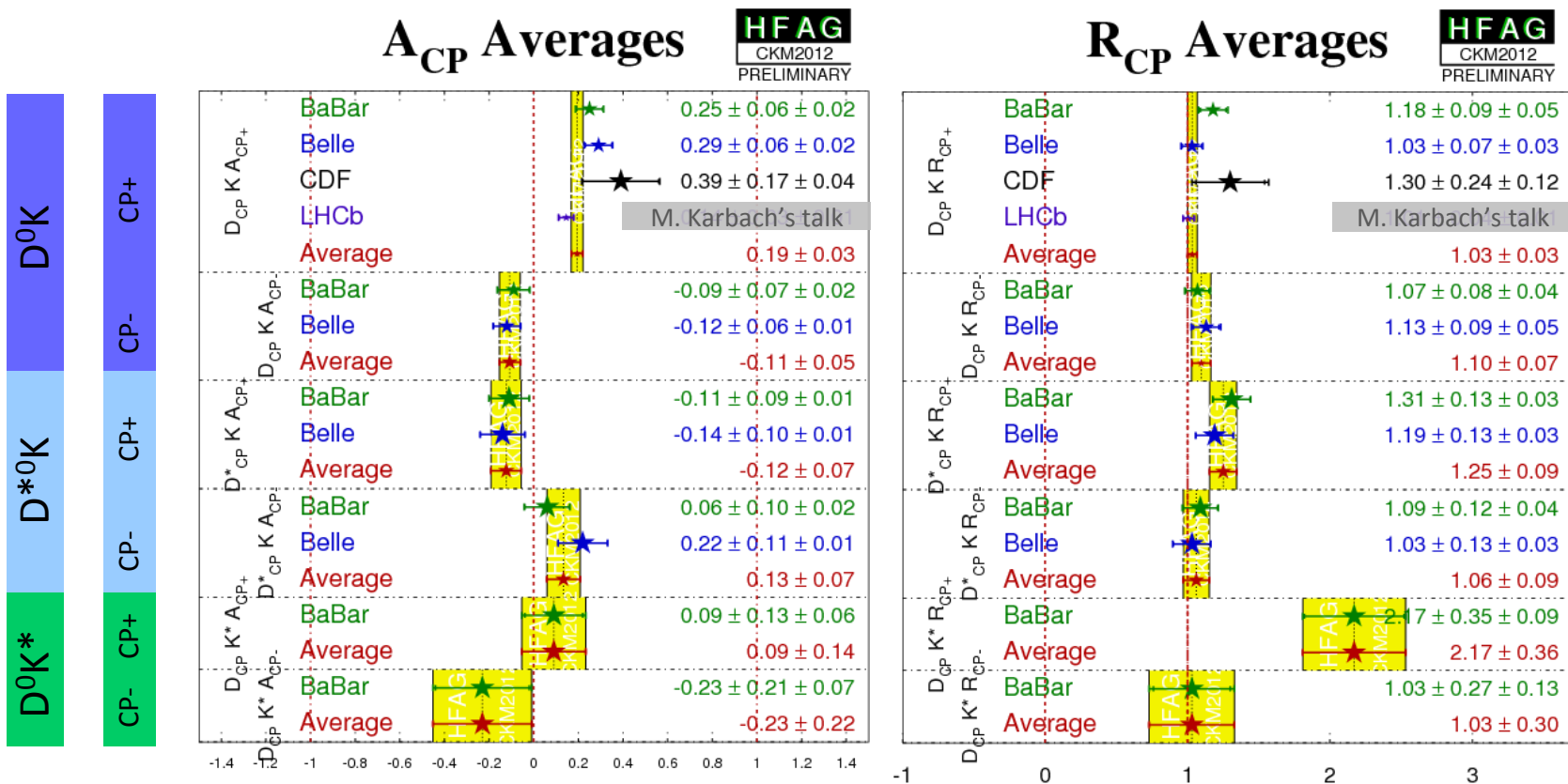
- $B \rightarrow D^0 K$ ,  $B \rightarrow D^{*0} K$  ( $D^{*0} \rightarrow D^0 \pi^0$  and  $D^{*0} \rightarrow D^0 \gamma$ ),  $B \rightarrow D^0 K^*$
- $D^0 \rightarrow K^+ K^-$ ,  $\pi^+ \pi^-$  (CP+),  $D^0 \rightarrow K_S \pi^0$ ,  $K_S \omega$ ,  $K_S \phi$ ,  $K_S \eta$  (CP-)



very challenging at hadronic colliders



# GLW results



# Why “not” $\gamma$ from GLW alone

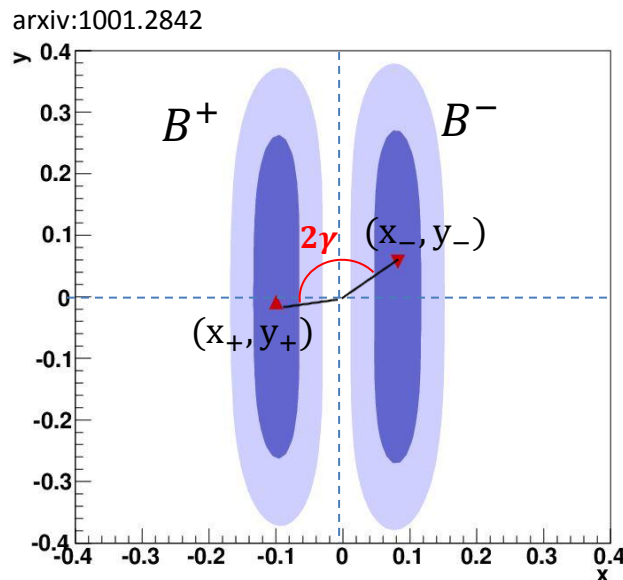
- The GLW observables can be expressed in terms of the cartesian coordinates  $x_{\pm}, y_{\pm}$  :

$$x_{\pm} = (R_{CP+}(1 \mp A_{CP+}) - R_{CP-}(1 \mp A_{CP-}))/4$$

$$r_b^2 = x_{\pm}^2 + y_{\pm}^2 = (R_{CP+} + R_{CP-})/2$$

good constraint on  $x_{\pm}$   
(comparable to GGSZ method  
when in same size dataset)

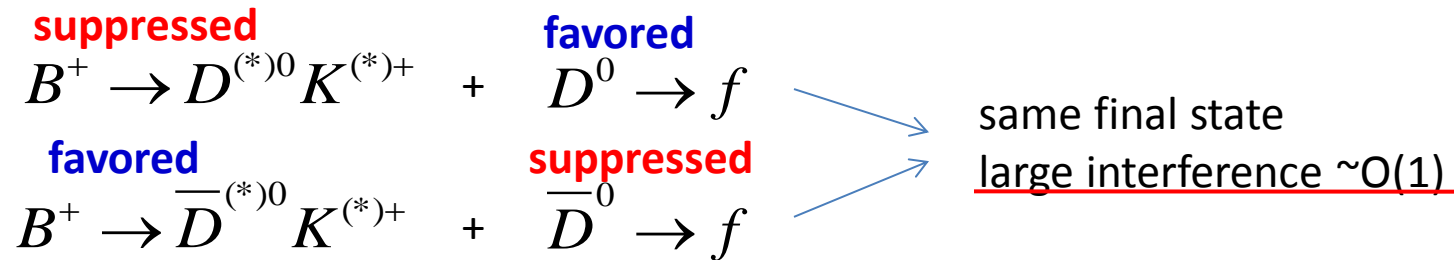
very loose constraint on  $y_{\pm}$





# ADS method

- $D^0$  to  $K^+\pi^-$ ,  $K^+\pi^-\pi^0$ ,  $K^+\pi^+\pi^+\pi^-$ , ... (doubly-Cabibbo-supp.)



- Measures  $B^+$  and  $B^-$  yields to determine the ADS observables:

$$R_{ADS} \equiv \frac{\Gamma(B^- \rightarrow D[\rightarrow f]K^-) + \Gamma(B^+ \rightarrow D[\rightarrow \bar{f}]K^+)}{\Gamma(B^- \rightarrow D[\rightarrow \bar{f}]K^-) + \Gamma(B^+ \rightarrow D[\rightarrow f]K^+)} = r_b^2 + r_D^2 + 2r_b r_D \cos(\delta_b + \delta_D) \cos \gamma$$

$$A_{ADS} \equiv \frac{\Gamma(B^- \rightarrow D[\rightarrow f]K^-) - \Gamma(B^+ \rightarrow D[\rightarrow \bar{f}]K^+)}{\Gamma(B^- \rightarrow D[\rightarrow f]K^-) + \Gamma(B^+ \rightarrow D[\rightarrow \bar{f}]K^+)} = 2r_b r_D \sin(\delta_b + \delta_D) \sin \gamma / R_{ADS}$$

$$r_D = \left| \frac{A(\bar{D}^0 \rightarrow f)}{A(D^0 \rightarrow f)} \right|$$





$$(r_D(K^+\pi^-) = 0.06)$$

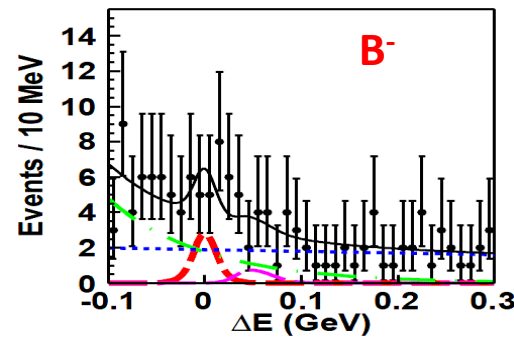
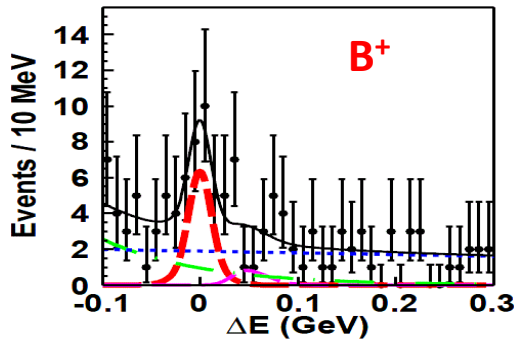
$$\delta_D = \arg \left[ \frac{A(\bar{D}^0 \rightarrow f)}{A(D^0 \rightarrow f)} \right]$$

measured at CLEOc/BESIII

# ADS reconstructed decay modes

$h=K,\pi$

Decays	$D^0 h^+$	$D^{*0} [D^0 \pi^0] h^+$	$D^{*0} [D^0 \gamma] h^+$	$D^0 K^{*+} [K_s \pi]$
$D^0 \rightarrow K^+ \pi^-$	 467M  772M	 467M  772M	 467M  772M	 379M ---
$D^0 \rightarrow K^+ \pi^- \pi^0$	 474M	---	---	---

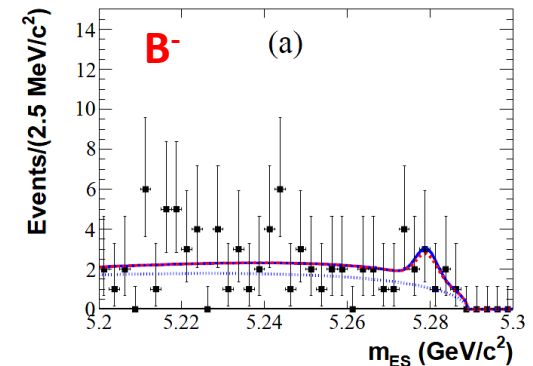
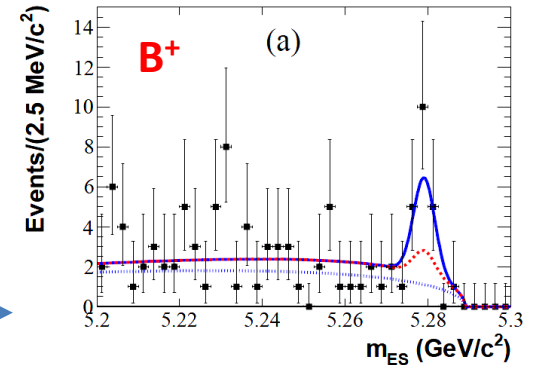


example:  $B \rightarrow D^0 K, D^0 \rightarrow K^+ \pi^-$

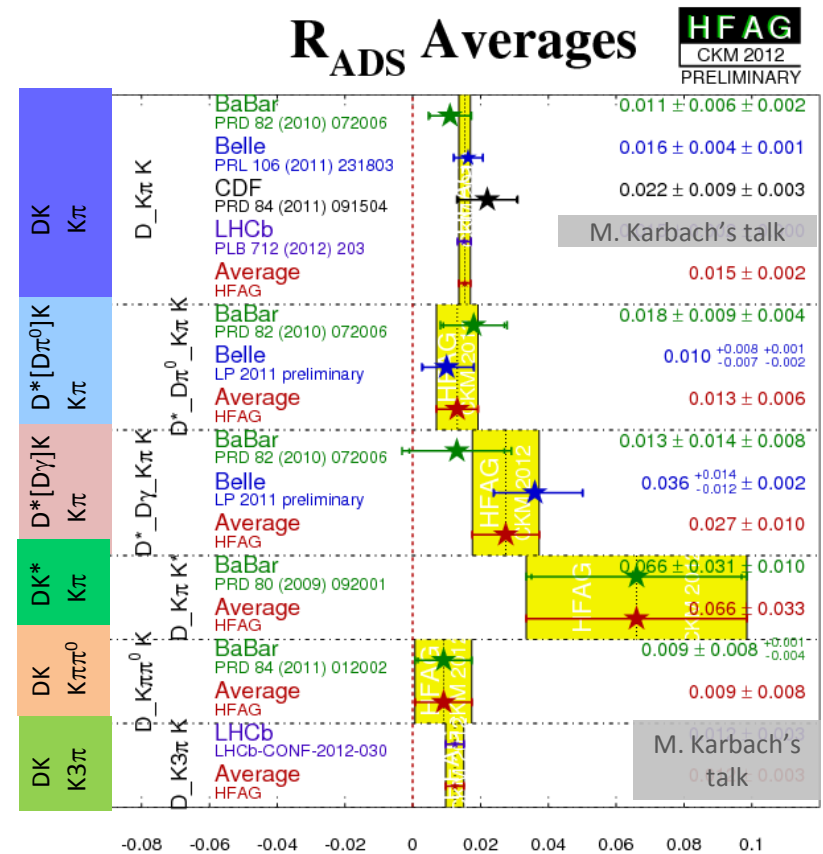
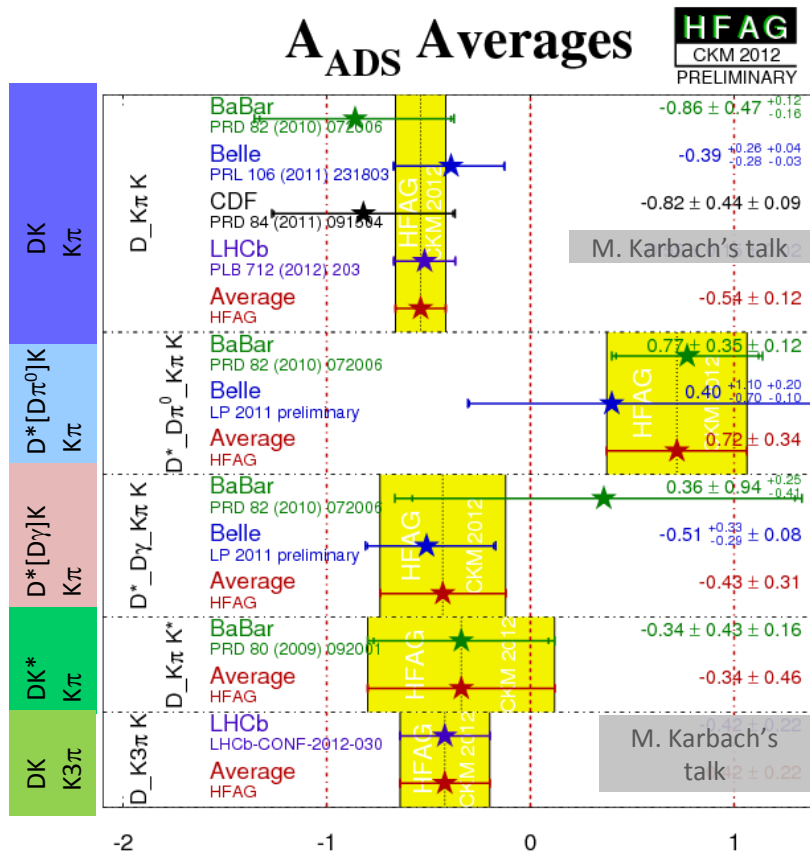
$$A_{ADS} = -0.82 \pm 0.46^{+0.12}_{-0.16}$$



$$A_{ADS} = -0.39^{+0.26+0.04}_{-0.28-0.03}$$



# ADS $B \rightarrow D^{(*)}K^{(*)}$ results



\* D<sup>(\*)</sup>π results available in backup slide

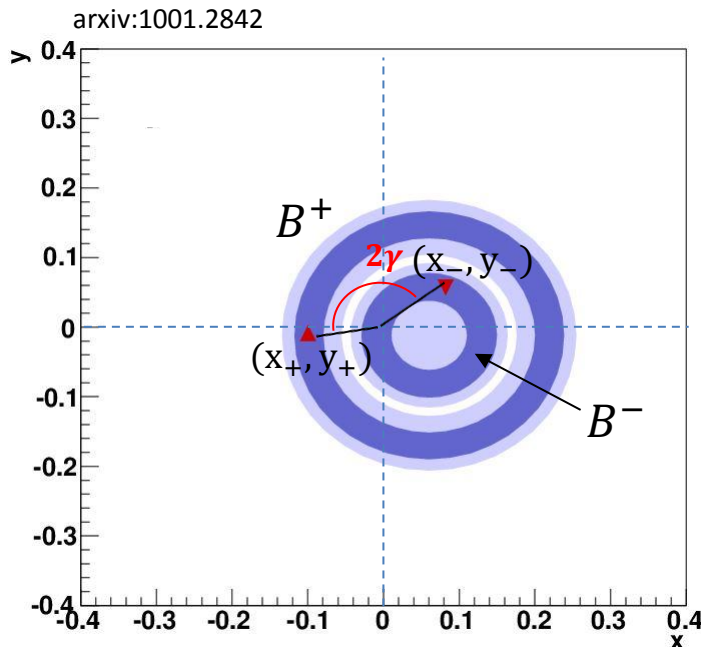
LHCb dominates the  $B \rightarrow DK$ ,  $D \rightarrow K\pi$  mode.

Final states with neutrals difficult in hadronic environment

# Why “not” $\gamma$ from ADS alone

- The ADS observables can be expressed in terms of the cartesian coordinates  $x_{\pm}, y_{\pm}$  :

$$(x_{\mp} + r_D \cos \delta_D)^2 + (y_{\mp} - r_D \sin \delta_D)^2 = \frac{\Gamma(B^{\mp} \rightarrow [K^{\pm} \pi^{\mp}]_D K^{\mp})}{\Gamma(B^{\mp} \rightarrow [K^{\mp} \pi^{\pm}]_D K^{\mp})}$$



$x_{\pm}, y_{\pm}$  are delocalized over two circles

Note: for  $B \rightarrow D^{*0} K$  the circles associated to  $D^{*0} \rightarrow D^0 \pi^0$  and  $D^{*0} \rightarrow D^0 \gamma$  are centered at opposite points  $(\mp r_D \cos \delta_D, \pm r_D \sin \delta_D)$  and  $\gamma$  can be extracted in principle up to discrete ambiguities.

GLW + ADS can constrain  $\gamma$

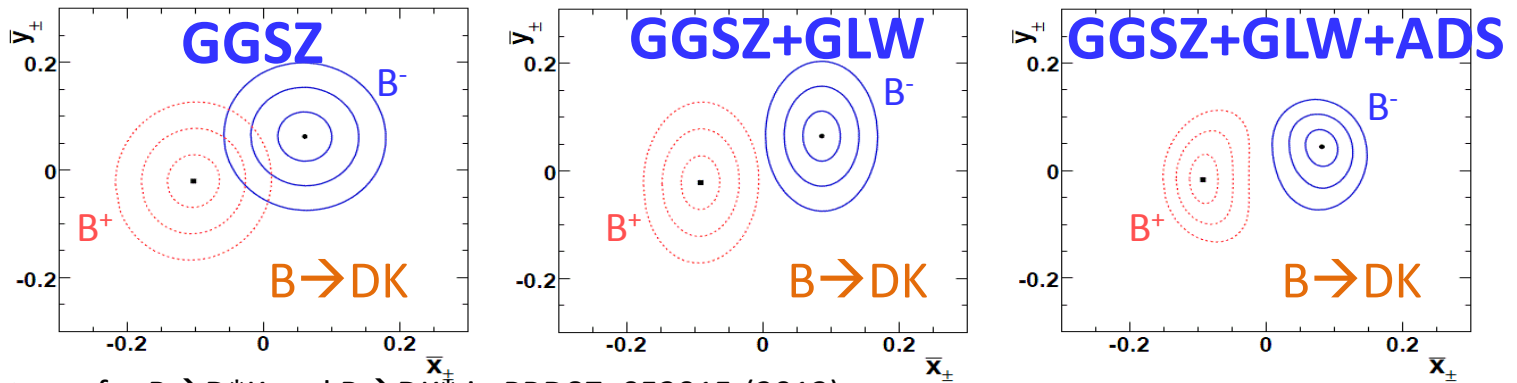
Note: uncertainty on the values of  $r_D$  and  $\delta_D$  are neglected

# BaBar GLW+ADS+GGSZ combination

## Combination of GGSZ+GLW+ADS in two stages

- I. combine the GGSZ, GLW and ADS  $B^\pm \rightarrow D^{(*)}K^{(*)\pm}$  observables (34 in total) to extract the combined  $x_\pm, y_\pm$  (4 for each B mode)
- II. transform the combined  $x_\pm, y_\pm$  into the physically relevant quantities  $\gamma, \{r_b, \delta_b\}_{D^{(*)}K^{(*)}}$

### stage I



contours for  $B \rightarrow D^*K$  and  $B \rightarrow DK^*$  in PRD87, 052015 (2013)

results:

$$\bar{z} \equiv x + iy$$

	Real part (%)	Imaginary part (%)	
DK	$\bar{z}_-$	$8.1 \pm 2.3 \pm 0.7$	$4.4 \pm 3.4 \pm 0.5$
	$\bar{z}_+$	$-9.3 \pm 2.2 \pm 0.3$	$-1.7 \pm 4.6 \pm 0.4$
D*K	$\bar{z}_-$	$-7.0 \pm 3.6 \pm 1.1$	$-10.6 \pm 5.4 \pm 2.0$
	$\bar{z}_+$	$10.3 \pm 2.9 \pm 0.8$	$-1.4 \pm 8.3 \pm 2.5$
DK*	$\bar{z}_{s-}$	$13.3 \pm 8.1 \pm 2.6$	$13.9 \pm 8.8 \pm 3.6$
	$\bar{z}_{s+}$	$-9.8 \pm 6.9 \pm 1.2$	$11.0 \pm 11.0 \pm 6.1$

external input required for the D hadronic parameters  $r_{K\pi}, \delta_{K\pi}, r_{K\pi\pi^0}, \delta_{K\pi\pi^0}, k_{K\pi\pi^0}$

# BaBar GLW+ADS+GGSZ combination

## stage II

- $(X_{\pm}, Y_{\pm})_{D^{(*)}K^{(*)}} \rightarrow \gamma, \{r_b, \delta_b\}_{D^{(*)}K^{(*)}}$   
with frequentist stat procedure

$$\gamma = (69^{+17}_{-16})^\circ \text{ (modulo } 180^\circ)$$

exp+DP model sys =  $\pm 4^\circ$

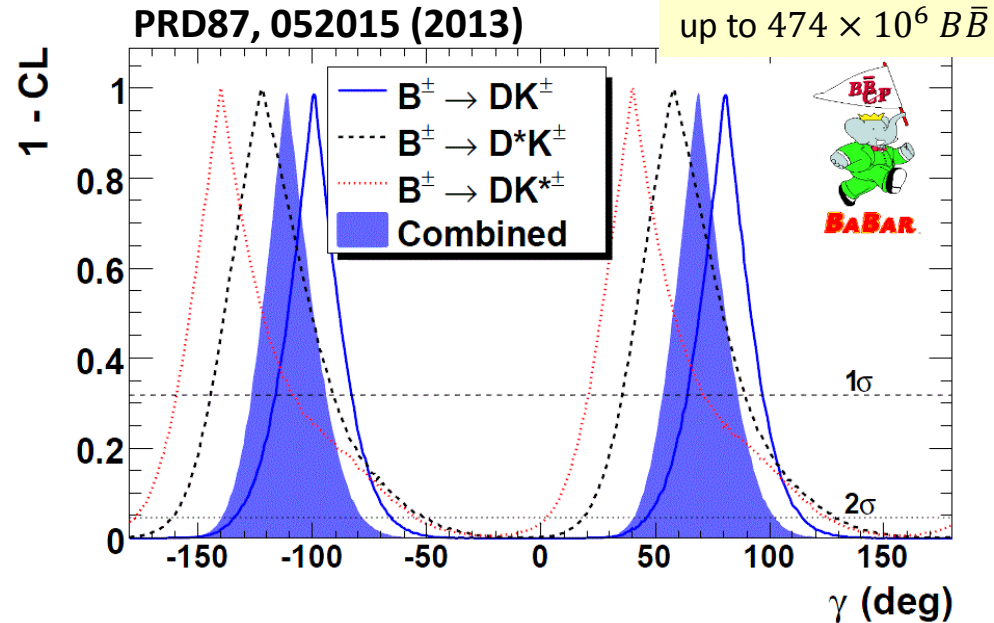
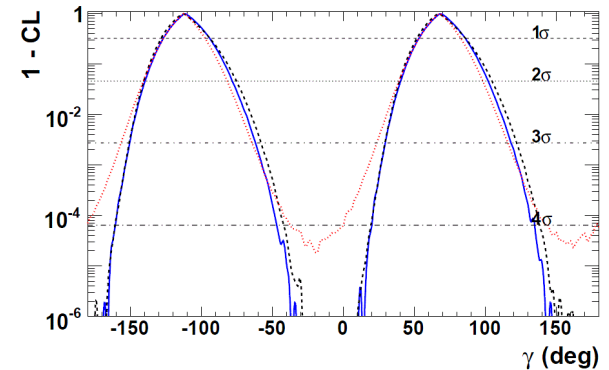
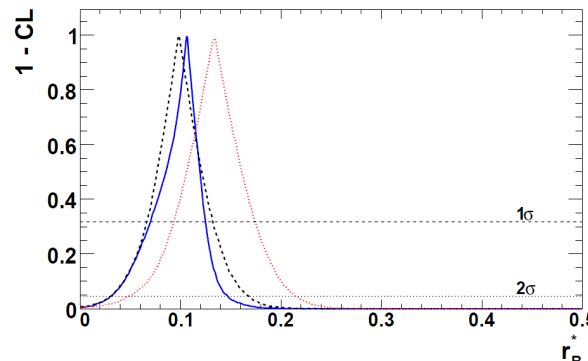
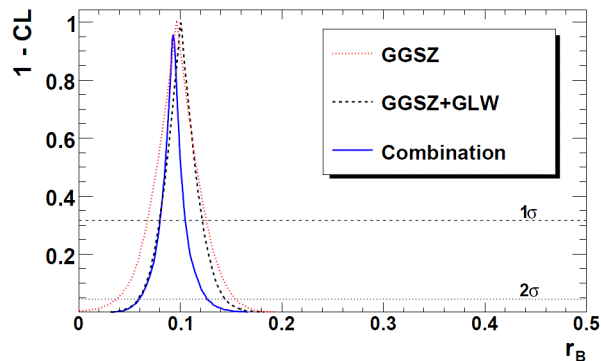
CPV significance: **5.9 $\sigma$**

(GGSZ alone 4.0 $\sigma$ . GGSZ+GLW 5.4 $\sigma$ )



$$r_b(\%) = 9.2^{+1.3}_{-1.2}$$

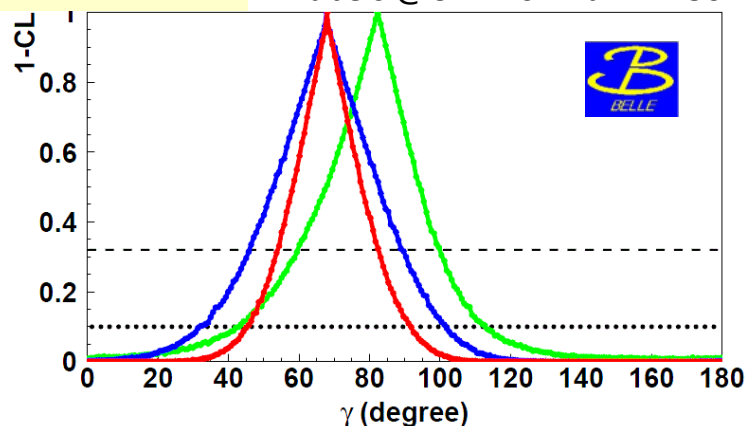
$$r_b^*(\%) = 10.6^{+1.9}_{-3.6}$$



# Belle GLW+ADS+GGSZ combination

Combination of GGSZ (mod. dep.)+GLW+ADS  $B \rightarrow D^0 K + B \rightarrow D^{*0} K$  (8+8+6 observables)  
 - frequentist stat procedure

up to  $772 \times 10^6 B\bar{B}$  K. Trabelsi@CKM2012 arxiv:1301.2033

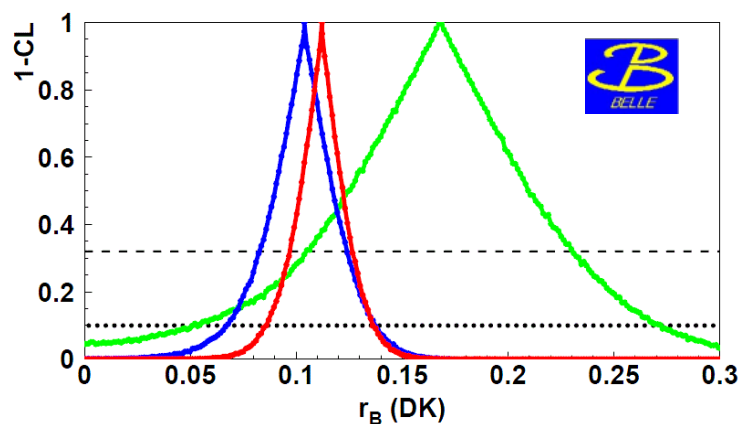


**GGSZ:**  $\gamma = (82^{+18}_{-23})^\circ$

**GGSZ + ADS:**  $\gamma = (68 \pm 22)^\circ$

**GGSZ + ADS + GLW:**

$\gamma = (68^{+15}_{-14})^\circ$



$B \rightarrow D^0 K:$

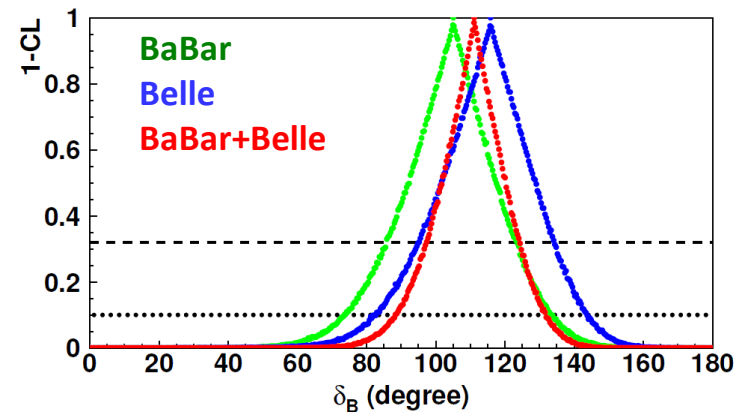
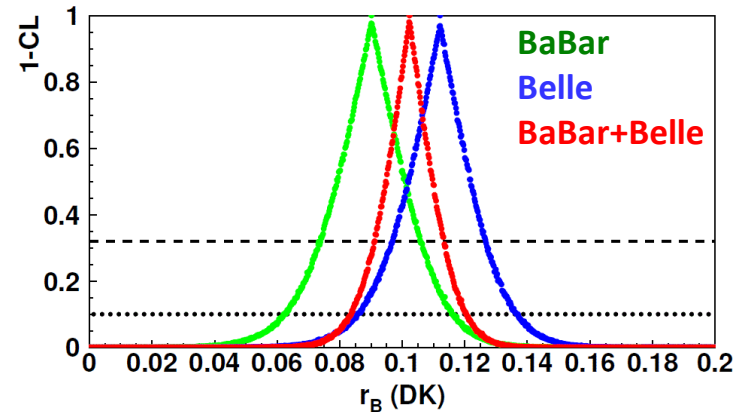
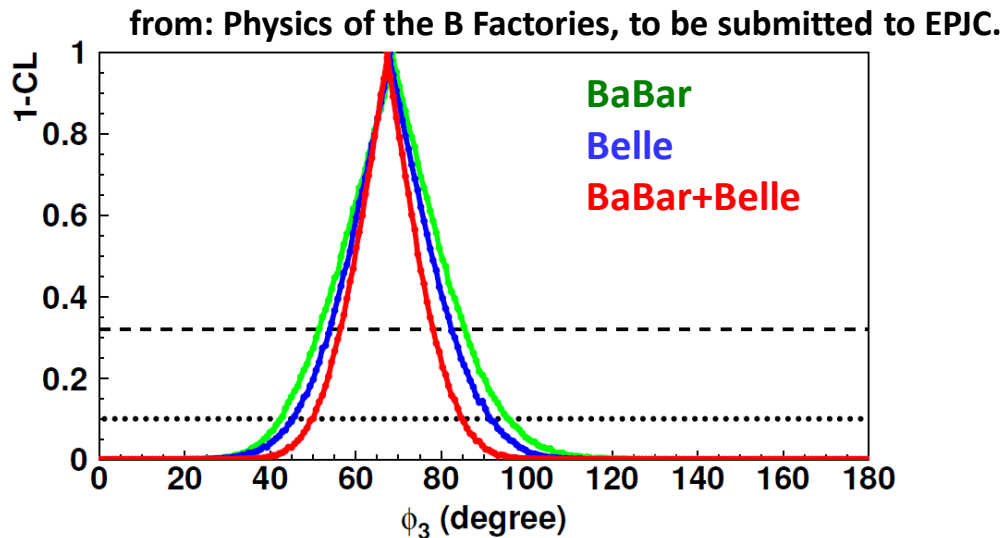
$r_b = 0.168^{+0.063}_{-0.064}$

$r_b = 0.104^{+0.020}_{-0.021}$

$r_b = 0.112^{+0.014}_{-0.015}$

# Belle and BaBar combination

Combination of BaBar and Belle: GGSZ mod. dep.+GLW+ADS with  $B \rightarrow D^0 K$ ,  $B \rightarrow D^{*0} K$  and  $B \rightarrow D^0 K^*$  (BaBar only). Frequentist stat procedure. BaBar and Belle model errors assumed uncorrelated.



**BaBar + Belle:**

$$\gamma = (67 \pm 11)^\circ \pmod{180^\circ}$$

$$r_b(DK) = 0.102 \pm 0.011$$

$$\delta_b(DK) = (111^{+13}_{-14})^\circ \pmod{180^\circ}$$



# D mixing and CPV in $B \rightarrow DK$ (and $B \rightarrow D\pi$ )

Several interesting studies on the effect of D mixing and CPV in the extraction of  $\gamma$  with  $B \rightarrow D^{(*)}K^{(*)}$  (and  $B \rightarrow D^{(*)}\pi$ )

- Effect of D mixing
  - Y. Grossman, Z. Ligeti, A. Soffer, PRD67, 071301 (2003); PRD72, 031501 (2005)
  - A. Bondar, A. Poluektov, V. Vorobiev, PRD82, 034033 (2010) (Dalitz mod ind)
- Effect of CPV in D decays
  - W. Wang, PRL110, 061802 (2013) (GLW)
  - M. Martone and J. Zupan, arXiv:1212.0165 (GLW)
  - B. Bhattacharya, D. London, M. Gronau, J. L. Rosner, arXiv:1301.5631 (GLW)
  - A. Bondar, A. Dolgov, A. Poluektov, V. Vorobiev arxiv:1303.6305 (Dalitz)

Corrections for D mixing and CPV not considered in BaBar and Belle  $B \rightarrow D^{(*)}K^{(*)}$  combinations

- effects expected to be small at present B-factories (although some may not be completely negligible)
- effects more and more important at LHCb and at Belle2

# Summary

- Sensitivity to  $\gamma$  ( $\phi_3$ ) dominated by the  $B^\pm \rightarrow D^{(*)0} K^{(*)\pm}$  decays so far
- BaBar and Belle have reconstructed all the most sensitive decay modes using all or almost all their final datasets
- BaBar GLW+ADS+GGSZ combination:

$$\gamma = (69_{-16}^{+17})^\circ \pmod{180^\circ}$$

exp+DP model sys =  $\pm 4^\circ$

- Belle GLW+ADS+GGSZ combination:

$$\gamma = (68_{-14}^{+15})^\circ \pmod{180^\circ}$$

- B-factories average:

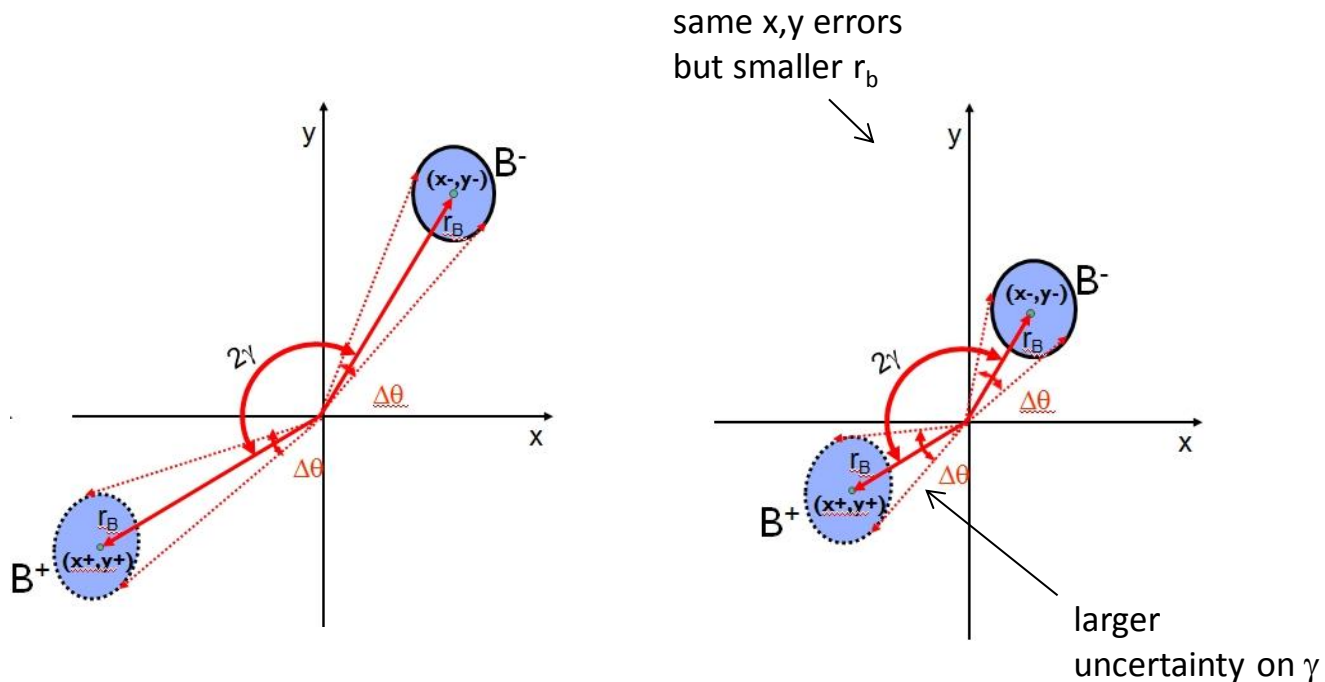
$$\gamma = (67 \pm 11)^\circ \pmod{180^\circ}$$

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BACKUP

# Dependence of $\sigma(\gamma)$ on $r_b$

- the error on  $\gamma$  (at fixed  $x, y$  uncertainty) scales roughly as  $1/r_b$



# $B \rightarrow D^{(*)} \pi$ ADS measurements

