

# Challenges related to the design, commissioning and operation of 3rd generation light sources

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### Synchrotron SOLEIL



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# Content

### 1. Introduction

synchrotron radiation storage ring synchrotron radiation sources

- 2. Accelerator Physics challenges emittance, brilliance, stability, time structure
- 3. Technical challenges
- **4.** Conclusion future trends



What is Synchrotron Radiation?

*Ultra relativistic electrons can be deviated by the constant magnetic field of bending magnets in which their trajectory is an arc of circle* 



### They emit photons in a direction tangent to their trajectory **This is synchrotron radiation**

### Such conditions are met in electron Storage Rings

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The Radiated Power :

$$P_{rad} = \frac{e^2 c}{6\pi\varepsilon_0 (m_0 c^2)^4} \frac{E^4}{\rho^2}$$

♦ Compare the radiation power from an electron ( $m_ec^2 = 0.511$  MeV) and from a proton ( $m_pc^2 = 939.19$  MeV) with the same energy :

$$\frac{P_{rad,e}}{P_{rad,p}} = \left(\frac{m_p c^2}{m_e c^2}\right)^4 = 1.13 \cdot 10^{13}$$

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### Why <u>Relativistic</u> Electrons?

The axially-symmetric radiation distribution in the moving frame K' (a.) transforms into a sharply forward peaked distribution in the laboratory frame (b.), with a half opening-angle  $\theta=1/\gamma$ .



#### This is one of the most useful features of synchrotron radiation.

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### Life science examples: DNA and myoglobin



Photograph 51 Franklin-Gosling DNA (form B) 1952 Franklin and Gosling used a X-ray tube:

Brightness was 10<sup>8</sup> (ph/sec/mm<sup>2</sup>/mrad<sup>2</sup>/0.1BW)

Exposure times of 1 day were typical (10<sup>5</sup> sec)

e.g. Diamond provides a brightness of 10<sup>20</sup>

100 ns exposure would be sufficient

Nowadays pump probe experiment in life science are performed using 100 ps pulses from storage ring light sources: e.g. ESRF myoglobin in action



(Courtesy of R. Bartolini))

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### All beamlines get beam simultaneously

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The successive generation of storage ring based sources

# <u>1<sup>st</sup> generation</u>

exploitation of the light from the bending magnets of e+/e- colliders originally built for elementary particle physics

# 2<sup>nd</sup> generation

exploitation of the light from the bending magnets and a few insertion devices at a later stage of dedicated storage rings ! lower e beam emittances and ! better optimisation of light extraction gain of about 1 000

# <u>3<sup>rd</sup> generation</u>

brilliance is the figure of merit rings designed for very low emittances priority to undulators large number of straight sections gain of about 10 000 moderate gain of about 100 on bending magnet light

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### Existing 3rd Generation Light Sources

1992	ESRF, France (EU)	6 GeV
	ALS, US	1.5-1.9 GeV
1993	TLS, Taiwan	1.5 GeV
1994	<b>ELETTRA</b> , Italy	2.4 GeV
	PLS, Korea	2 GeV
	MAX II, Sweden	1.5 GeV
1996	APS, US	7 GeV
	LNLS, Brazil	1.35 GeV
1997	Spring-8, Japan	8 GeV
1998	<b>BESSY II</b> , Germany	1.9 GeV
2000	ANKA, Germany	2.5 GeV
	SLS, Switzerland	2.4 GeV
2004	SPEAR3, US	3 GeV
	CLS, Canada	2.9 GeV
2006	SOLEIL, France	2.8 GeV
	DIAMOND, UK	3 GeV
	ASP, Australia 3 GeV	
	MAX III, Sweden	700 MeV
	Indus-II, India	2.5 GeV
2008	SSRF, China	3.5 GeV
2009	PETRAIII, Germany	6.0 GeV
2010	ALBA, Spain	3.0 GeV





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### **3rd GLS under construction or planned**

Under construction

2014	NSLS-II, US	3 GeV
2015	SOLARIS, Poland	1.5 GeV
2015	<b>TPS</b> , Taiwan	3 GeV
2016	SESAME, Jordan	2.5 GeV
2016	MAX-IV, Sweden	3 GeV

Planned

CANDLE, Armenia	3 GeV
SIRIUS, Brazil	3 GeV



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### 3<sup>rd</sup> generation Synchrotron light sources

- Machines optimised for High Brilliance
- Smaller source sizes, higher current
- Highly performing <u>insertion devices</u> matched to the beamline needs
- Beamlines much more accurate (specific scientific use).



### **Photon Brilliance at SOLEIL**





### **Photon Brilliance at Diamond**









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A parameter of prime importance in experiments with synchrotron radiation sources is the spectral brilliance (brightness) defined as :

$$B = \frac{dN_{ph}}{dA \, d\Omega \, dt \, d\lambda / \lambda}$$

 $\frac{photons \ per \ sec \ ond}{mm^2 mrad^2 0.1\% b.w}$ 

- High intensity, small beam cross section, good focusing of photon beam and high monochromaticity are looked for in order to maximize the photon flux.
- Apart from *diffraction effects*, we have :

$$dA \ d\Omega \approx \mathcal{E}_x \mathcal{E}_z$$

 $\mathcal{E}_x$  and  $\mathcal{E}_z$  are the transverse beam emittances, which are the areas occupied by the beam in horizontal (x,x') and vertical (z,z') phase planes.

**HIGH** PHOTON BEAM BRILLIANCE ⇒ **LOW** ELECTRON BEAM EMITTANCE and **HIGH** electron beam current.

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# Horizontal Emittance

₿⊗

ρ-Δρ

### Generation of emittance by radiation

Two e<sup>-</sup> with same E and same trajectory at the entrance of the bending magnet (they cannot be distinguished : no e- beam size)

They both radiate  $\Delta E$  and exit with  $E - \Delta E$ 

1 emits  $\Delta E$  at the origin, smaller  $\rho$ 2 emits  $\Delta E$  but at the end They follow different trajectories

$$dx(s) = \eta(s) \frac{\Delta E}{E}$$
  $dx'(s) = \eta'(s) \frac{\Delta E}{E}$ 

With  $\eta(s)$  and  $\eta'(s)$  the dispersion function and its derivative

Due to the random aspects of the emission (location, energy), the different electrons exit the bending magnet with different trajectories and different energies. The beam is heated up by the radiation :

# $\Rightarrow$ increase of H beam size, H beam divergence and energy spread.

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# Horizontal Emittance

### Damping by the RF



The heating in the H plane induced by radiation in every bending magnets would lead to a continuous blow up of the H emittance.

However, the restoring of the lost energy by the RF cavity cool down the transverse oscillations in the H plane

The damping time is about the time it would require for the electron to lose its full energy.  $\tau \sim E/U_{0}$ . Memory of initial conditions is totally lost when the RF has restored the total energy of the electron.  $\longrightarrow$  Emittance independent of injection conditions

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### Equilibrium Beam Emittance

The natural horizontal emittance for an isomagnetic ring
i.e. all bending magnets having same bending radius is :

 $J_x$  the horizontal damping partition number.  $J_x \sim 1$  (zero field gradient in bending magnet)  $J_x < 2$  (vertical focusing in bending magnet) : (potentially) emittance reduction of a factor two  $\varepsilon_{x} = \frac{C_{q} \gamma^{2} \langle H \rangle_{dipole}}{J_{x} \rho}$ 

 $C_q = 3.83 \times 10^{-13} \text{ m}$  and  $\gamma$  is the Lorentz factor.  $\rho$  is the bending radius.

H is the so called lattice invariant or dispersion's emittance or H-function

$$H(s) = \gamma_x(s)\eta^2(s) + 2\alpha_x(s)\eta(s)\eta'(s) + \beta_x(s)\eta'^2(s)$$

 $\langle ... \rangle$  average taken only in the part of the circumference where photons are emitted, (BM and IDs)

 $In practical units, \varepsilon_x$  is given by :

$$\mathcal{E}_{x}[nm.rad] = 1470 E[GeV]^{2} \frac{\langle H \rangle_{dipole}}{\rho J_{x}}$$

 $\mathcal{E}_x$  is completely determined by the energy, bending field and lattice functions.

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### Equilibrium Beam Emittance

✤ After calculation of the <H> value, the natural horizontal emittance for an isomagnetic ring i.e. all bending magnets having same bending radius can be expressed as :

- $J_x$  the horizontal damping partition number.  $\theta$  deviation angle of one bending magnet  $\rho$  bending radius  $l_b$  bending magnet length N number of bending magnets
- $C_q = 3.83 \times 10^{-13} \text{ m}$  and  $\gamma$  is the Lorentz factor.

$$\varepsilon_x = F \frac{C_q \gamma^2 \theta_b^3}{J_x}$$

$$F = \frac{1}{3} \left[ \frac{\beta_0}{l_b} - \frac{1}{4} \alpha_0 + \frac{1}{20} \gamma_0 l_b \right]$$

 $\beta_0, \alpha_0, \gamma_0$  twiss parameters at the entry of the BM

### It is a general lattice property, there is no assumption on the lattice type.

$$\mathcal{E}_x \propto \frac{1}{N_b^3}$$

⇒ Should use many short Bending Magnets to get low emittance

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### Minimum Emittance (with Achromatic Condition)

The minimum equilibrium beam emittance in an isomagnetic ring with an Achromatic Arc Condition,  $\eta_0 = \eta'_0 = 0$ , at the entrance of the **BM** :

$$\varepsilon_{x,\min} = \frac{C_q \gamma^2 \Theta^3}{4\sqrt{15} J_x}$$

 $\Rightarrow$ DBA or TBA lattices (Double/Triple Bend Achromat)







# **Minimum Emittance (without Achromatic Condition)**

By breaking the achromatic condition (non-zero dispersion in straight sections) we can obtain the configuration in which the emittance becomes the smallest.

$$\varepsilon_{x,\min} = \frac{C_q \gamma^2 \Theta^3}{12 \sqrt{15} J_x}$$



⇒It is smaller **by a factor 3** than in the achromatic arc configuration

Example of Machines which move from Achromatic conditions to non zero dispersion in SS

ESRF	7 nm	$\rightarrow$ 3.8 nm
APS	7.5 nm	$\rightarrow 2.5 \text{ nm}$
SPring8	4.8 nm	$\rightarrow$ 3.0 nm
SPEAR3	18.0 nm	$\rightarrow$ 9.8 nm
ALS (SB)	10.5 nm	$\rightarrow 6.7 \text{ nm}$

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**Emittance** Achieved

### The minimum emittance can't be easily achieved :

Ideal values for SOLEIL and Design values (without achromatic condition) :

$\alpha_{0,min}$ = $\sqrt{15}$ = 3.873	$\alpha_0 = 1.8$
$\beta_{0,min} = 2.17m$	$\beta_0 = 1.5 m$
$\mathcal{E}_{x0,min} = 1.8 nm.rad$	$\mathcal{E}_{x0}$ = 3.7 nm.rad

The ideal value  $\alpha_{0,min} = \sqrt{15}$ 

causes the betatron function to reach a sharp minimum inside the **BM** and then to increase from there on to large values in the quadrupoles, leading to extremely high <u>chromaticity</u>.



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# Emittance achieved

 $\Rightarrow$  Chromaticity has to be corrected for two reasons :

**Momentum acceptance :** some variation of energy deviation has to be accepted by the storage ring for reasons of **beam lifetime**.

**Head tail instability:** collective oscillation of electrons in head and tail of the bunch leading to very fast beam loss.  $\longrightarrow$  damped by operating with positive chromaticity







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Circumference = **354m** 

24 straight sections (variable length)

4 x 12 m 12 x 7 m 8 x 3.6m

Strong focusing

Different  $\beta$ -functions in straight sections

Bunch length = 4mm

### Low Emittance Lattice: SOLEIL example



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# **SOLEIL Lattice Modified**



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### Low Emittance Lattice: Diamond example



### nominal, non-zero dispersion lattice

(Courtesy of R. P. Walker)

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#### PETRA III 72/72 WITH DAMPING WIGGLERS Win32 version 8.51/15 05/11/06 13.42.25 50. **PETRA 3 (DESY) : in operation** β (m) B's 45. PETRA 3 40. 7 old octants FODO + 1 New octant DBA 35. 30. 6 GeV 25 20. 2304 m 15. 10. **Damping Wigglers** 5. 8.8 D (m) 6 SS for a total of 45 m (14 IDs) 0.8 0.7 0.6 => Ex = 1 nm.rad 0.5 0.4 0.3 0.2 uuu 0.1 0.0 -0. 5*0*0. 1500. 1000. 0.0 18 □ MAX-IV (Sweden): under construction 16 7-BA (zero dispersion in SS) 20 cells 14 3 GeV Beta Functions [m] 12 528 m 10 => Ex = 0.34 nm.rad 8 6 4 2 0 5 10 15 20 0 s [m]

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2500. s (m)

Dispersion [m]

2000.

0.08

0.07

0.06

0.05

0.04

0.03

0.02

0.01

-0.01

0

25

Ultra low Emittance rings



It is a complex process where the Accelerator Physicist is guided by

• (semi-)analytical formulae for the computation of nonlinear maps, detuning with amplitude and off-momentum, resonance driving terms

• numerical tracking: direct calculation of non linear tuneshifts with amplitude and off-momentum, 6D dynamics aperture and the frequency analysis of the betatron oscillations

Many iterations are required to achieve a good solution that guarantees a good dynamic aperture for injection and a good Touschek lifetime



# Working point: Tune Diagram





### Frequency Map Analysis (NAFF)

**Off**-momentum

Launched particles over a fine X-Z grid plotting Numerical tunes Highlighting, nonlinearity (diffusion rate)



#### **On**-momentum

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#### Bare Machine- $v_x$ =18.202 $v_z$ =10.317 - chromaticities 2/2



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## Non linear Optimization Strategy







To reach a high brilliance you need to operate with a high current. There are 3 main limitations to overcome :

**First limitation** : the coupled bunch instabilities, which are excited by Higher Order Modes in the RF cavities (HOM).  $\Rightarrow$  increase of the energy spread and bunch length => brilliance reduction  $\Rightarrow$  transverse beam oscillations =>beam losses

# Solutions to push the instability thresholds have been experienced and successfully implemented :

Shifting the mode frequencies to avoid interaction with the bunch spectrum (temperature control of the cavities)
 Landau damping : partial filling or harmonic cavity
 Reducing the number and impedances of parasitic modes (superconducting cavities, parasitic mode dampers)
 Feedback systems



<u>Second limitation</u> : the transverse instabilities, resistive wall, ions trapping in multibunch / mode detuning and TMCI in single bunch  $\Rightarrow$  Beam blow-up => brilliance reduction

 $\Rightarrow$  transverse beam oscillations => beam losses

Resistive wall is enhanced by the presence of low gap ID vessels

Solutions to push the instability thresholds have been experienced and successfully implemented :

Shifting the mode spectrum with positive chromaticities
 Landau damping : partial filling (mainly against ions)
 Reducing machine impedance (smooth tapers, no gaps, RF fingers..)=> at design stage
 Bunch by bunch transverse Feedback systems (required to achieve high current)



<u>**Third limitation</u></u> : the thermal load on the vacuum system and front-end \Rightarrow High heat load with high power density on beam absorbers \Rightarrow Heating due to RF losses in transitions</u>** 

#### **Solutions :**

Distributed lump absorbers

➢Geometry and material of the absorbers (Glidcop, cooling circuits,..)

Reduce trapping mode geometry (smooth tapers, no gaps, RF fingers..)=> at design stage

Require efficient and reliable Machine protection system (thermocouples, flowswitches, Beam position interlocks, instability interlock,..)

<u>Note</u> : the danger is not coming from the power in the electron beam (Thousands of Joules lost in few microsec) but from the photons beam  $\Rightarrow$  In case of electron beam mis-steering (without losing it) the photons may reach a part which is not able to sustain the power.



#### **Few figures (reached):**

APS,	100 mA
SPRING8	100 mA
ESRF	200 mA
SLS	400 mA
DIAMOND	300 mA
SOLEIL	430 mA => 500 mA (reached in machine
	development)



## SLS achievement: minimum vertical emittance

#### Measures for minimization

Vertical realignment of storage ring girders with stored beam and orbit feedback running.

Model dependent optimization: measurements of response matrix and vertical dispersion and calculation of skew quadrupole currents.

Model independent optimization: random walk of skew quadrupole currents with beam height measurement as target function.



Fine tuning of emittance monitor to extend its range of measurement .



#### **Results**

Record-Value
 (Dec. 6, 2011)
 0.9± 0.4 pm rad

Design value15 pm rad

Quantum limit0.2 pm rad

⇒ Nuclear Instruments and Methods in Physics Research A 694 (2012) 133–139



## **Beam Position Stability**



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# Short Term Beam Position Stability





## **SOLEIL Beamlines**





# **5 Modes of Operation for the Users** All in Top-up injection

Mode of operation	User Operation	Ultimate performance achieved		
Multibunch	<b>430 mA</b>	500 mA		
Hybrid	425 mA + 5 mA	425 mA + 10 mA		
8 bunch	<b>100 mA</b>	100 mA		
1 bunch	<b>12 mA</b>	20 mA		
Low <b>Q: bunch length and</b> current	4.7 ps RMS and 65 μA per bunch	2.5 ps RMS and 10 μA per bunch		

5 feedbacks: TFB, SOFB, FOFB, BTUNE-FB, Coupling-FB

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# *Time* structure

2

#### Many facilities provide specific filling pattern to enable time structure experiments

ESRF : Single bunch, 16 bunch, Hybrid mode (Camshaft),..

SOLEIL : 1 bunch, 8 bunch, hybrid mode Bunch length are of 10-20 ps (rms)

$$\sigma_z = \frac{\alpha c}{2\pi f_s} \sigma_\varepsilon \propto \sqrt{\frac{\alpha \gamma^3}{d V_{RF} / dz}}$$

#### **Request to get shorter bunches**

⇒Low alpha lattices provide **ps** bunch length at very low current per bunch (BESSY II, ANKA, ELETTRA, SPEAR3, SOLEIL)  $\alpha = \frac{1}{L} \oint \frac{D_x}{\rho} ds \approx 10^{-6}$ 

 $\Rightarrow$  Femto slicing enables to produce **100fs** bunch "slices" using an energy modulation induced by a laser (ALS, SLS, BESSY, and soon SOLEIL). The photon flux is rather weak.





## SLS

# **Top-up Operation**



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## SLS beam current



Typical 24 hours of user operation

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Machine





SOLEIL









soit 96,1 % du temps de faisceau programmé

SOLEIL



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## Machine statistics

ESRF Record				
Machine statistics	2009	2010	2011	2012
Availability (%)	99.04	98.78	98.91	98.6
Mean time between	75.8	67.50	107.8	60.0
failures (hrs)	0.72	0.82	1 10	0.95
failure (hrs)	0.73	0.82	1.18	
A lot of activities and modifications during the long winter 2011-12 shutdown, the longer summer shutdown, the October shutdown and the longer winter shutdown.				
Larger number of failures results direc	in 2012 but tly from the	only a sma upgrade pr	ll part of the o ogramme	down time
r NADJI				



Répartition des 135 heures de pannes avec incidence faisceau stocké (réinjections incluses) par groupe et par RUN, durant les sessions Lignes et RP en 2012, au RUN5





**\*** 

ESRF

## Machine statistics

A Light for Science





# MAX IV Magnets

• Pre-series for each type of magnet has been built and approved for mechanical tolerances.







# **SOLEIL Dipole Vacuum**



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# **RF SYSTEMS : cavities**

#### Superconducting RF cavities : SOLEIL





- $\rightarrow$  2 cavities @ 352 MHz
- $\rightarrow$  150 kW per cavity (LEP type coupler)
- $\rightarrow$  2MV per cavity
- $\rightarrow$  2 cryomodules installed
- →few problems with the cold tuning system

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(Collab CEA/CERN/SOLEIL/ESRF)



## Single cell NC HOM damped cavity A Light for Science



#### Goal:

RF distribution to create new experimental stations

➢Prepare future upgrades

→ 3 Prototypes under test



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#### IOT's : ELETTRA, DIAMOND, ALBA,...

Solid state amplifier : SOLEIL and more recently ESRF

→ 4x190 kW @ 352 MHz
→ Smooth operation

→ Excellent reliability



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# **BPM Electronic Boards**

# LIBERA BPM digital electronics provide sub micron resolution, and also :

- $\Rightarrow$  simultaneous turn by turn measurement
- $\Rightarrow$  slow orbit data
- $\Rightarrow$  fast orbit feedback data, and processing
- $\Rightarrow$  with low current dependence



In use at Diamond, SOLEIL, ELETTRA, ASP, ALBA, ESRF, PETRA3,...



### Insertion Devices: In-Vacuum undulator U20

SYNCHROTE		A				
Period	20 mm		1-	A TON		
Nbr of Periods	98				12	
Length	2.0 m				L.Em	
Туре	Hybrid In-Vacuum			آلے	1.27	
Min. gap (mm)	5.5 to 30		1		a sal	E. March
Polarisation	Linear H			1		A Par
Bzmax	<b>0.97 T</b>			11		Par I
Photon Energy	3 – 20 keV		N.			

# **Electromagnet Helical Undulator HU640**



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# Apple-II Type Helical Undulator HU80

Period	80 mm		
Number of Periods	19		
Length	1.65 m		
Туре	Apple-II		
Minimum gap	15 to 250		
( <b>mm</b> )			
Polarisation	Circ./Lin.		
Bxmax	<b>0.76 T</b>		
Bzmax	0.85 [1.0] T		
Photon Energy	80 [35] – 1500 eV		

SYNCHROTRON





# Superconducting Wigglers



Superconducting wigglers are used when a high magnetic field is required 3 -10 T

They need a cryogenic system to keep the coil superconductive

Nb<sub>3</sub>Sn and NbTi wires

SCMPW60 at Diamond 3.5 T coils cooled at 4 K 24 period of 64 mm gap 10 mm Undulator K = 21

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## Further decrease of the horizontal emittance



	Electrons beam dimensions		Photons beam dimensions	
	<i>σ</i> e μm	$\sigma'_{ m e}~\mu$ rad	<i>o</i> tot μm	$\sigma'_{ m tot}$ $\mu$ rad
Plan H	182	30.4	183	34
Plan V	5.1	2.9	6.5	16







#### The ESRF low emittance lattice See: Poster MOPEA008



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ESRE

A Light for Science

### Brilliance at lower horizontal emittance






## Storage ring performance (current and future sources) horizontal emittance

- ESRF 2BA
- PETRA III 2BA
- NSLS II 2BA
- MAX IV 7BA
- Sirius 5BA
- Spring-8 6BA
- ESRF 7BA

- 4000 pm 6 GeV, operational
  - 1000 pm 6 GeV, operational
  - ~350 pm 3 GeV, construction
- ~300 pm 3 GeV, construction
- ~250 pm 3 GeV, in planning
  - ~70 pm 6 GeV, in planning
- ~150 pm 6 GeV, in planning





SYNCHROTRON • Horizontal (natural) emittance:  $\mathcal{E}_{x0} = \frac{C_q \gamma^2 \langle H \rangle_{dipole}}{J_x \rho_x}$ 

✤ J<sub>x</sub> is related to the damping partition D by: J<sub>x</sub> = 1 − D

$$D = \frac{\oint \frac{\eta_x}{\rho_x^3} ds + \frac{2}{B^2 \rho_x^2} \oint \eta_x B \frac{dB}{dx} ds}{\oint \frac{ds}{\rho_x^2}}$$
 If  $D = -1 \implies J_x = 2$   
$$\int \frac{\int \frac{ds}{\rho_x^2}}{\int \frac{ds}{\rho_x^2}}$$
 The horizontal emittance can be divided by a factor 2

But with an increase of energy spread by a factor  $\sqrt{2}$ 

Such as: 
$$J_x + J_z + J_s = 4$$
 where  $J_z = 1$  and  $J_s = 1$ 

A magnetic element introducing the product  $B^*dB/dx$  in a straight section where the dispersion  $\eta_x$  is **non zero** contributes to the modification of **D**.

If this magnetic element is such as the field (*B*) and the gradient (dB/dx) are of opposite sign, this contribution is *negative* and consequently we can try to find the <u>conditions to make D =-1</u>!

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## **Robinson Wiggler** Application to SOLEIL: (work <u>in progress</u> by a PhD student)

↔ In short straight sections of the SOLEIL Storage Ring:  $\eta_x$  = 28 cm.



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## **CONCLUSIONS**

- $\Rightarrow$  Third generation light sources are very reliable sources of high brilliance, which provide very stable photon beam from IR to hard X-rays.
- ⇒ They operate ~5000 hours/year mostly in Top-up, with > 95% beam time availability.
- ⇒ No evidence of under subscription: the demand from the User's community and the number of beamlines per facility is still increasing.
- ⇒ The agreement with model is excellent for the linear optics and improvements can be foreseen for the nonlinear optics.

## $\Rightarrow$ Future developments will target :

- higher brilliance and more coherence => ultra low emittance ~ 10 500 pm
- Technological progress is needed to reach the challenging goal