The ATLAS detector ϕ_{η}^{*} , an alternative variable to study the physics of ρ_{T}^{2} Differential cross section measurements of $Z/\gamma^{*} \rightarrow \ell \ell$

ATLAS ϕ_{η}^* measurement of Z/γ^* at 7 TeV Phys. Lett. B 720 (2013) 32-51

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On behalf of the ATLAS collaboration





Laboratoire d'Annecy-le-Vieux de Physique des Particules The ATLAS detector ϕ^*_{η} , an alternative variable to study the physics of p_T^Z Differential cross section measurements of $Z/\gamma^* \to \ell\ell$

Contents

The ATLAS detector

- φ^{*}_η, an alternative variable to study the physics of p^Z_T
 p^Z_T theoretical points of view
 φ^{*}_η variable
- 3 Differential cross section measurements of $Z/\gamma^*
 ightarrow \ell\ell$
 - $Z/\gamma^* \to \ell \ell$ event selection
 - Differential cross section measurement as a function of ϕ_{η}^{*}
 - Differential cross section measurement as a function of ϕ_η^* and y_Z

The ATLAS detector

 ϕ_{η}^{*} , an alternative variable to study the physics of p_{T}^{Z} Differential cross section measurements of $Z/\gamma^{*} \to \ell \ell$

The ATLAS detector



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p_T^Z - theoretical points of view

- High p_T^Z :
 - Fixed order perturbative QCD predictions (like FEWZ).
- Low p_T^Z :
 - Resummation of leading logarithms to all orders in α_S (+ non-perturbative parameterization like RESBOS).
 - OR modeling by parton shower generators (like PYTHIA).
- \rightarrow Motivations of the p_T^Z measurement:
 - Provide a test of pQCD.
 - Study the low p_T^Z region where non-perturbative effects may play a role.
 - Help in improving the modelling of *W* boson production needed for a precise measurement of the *W* mass.
 - Help in understanding the low p_T spectrum of the Higgs boson.

The ATLAS detector ϕ^*_η , an alternative variable to study the physics of p^Z_T Differential cross section measurements of $Z/\gamma^* \to \ell \ell$ p_T^Z - theoretical points of vie $\phi^{*}_{oldsymbol{\eta}}$ variable

Why new variables?

The measurement at low p_T^Z is limited by the experimental resolution rather than the event statistics in particular for the electron channel.

Main systematic uncer-10 20 tainties of the latest p_T^Z ∆ơ/∆P_T (pb / GeV/c) 4⊽/2⊽ 10 measurement: 10⁻¹ Calorimeter response 10⁻² 20 modeling \oplus Unfolding 10⁻³ $\sim 1\%$ 10^{-4} CDF $\int Ldt = 2.1$ fb \sim 140K $Z \rightarrow ee$ events

10⁻⁵

Phys. Rev. D 86 (2012) 052010

250

200

e*e Pair P_T (GeV/c)

300

350

 \implies Optimize new variables:

- less sensitive to the effects of experimental resolution.
- probe the same physical effects as p_T^Z .

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ϕ_n^* definition

P T Eur. Phys. J. C 71 (2011) 1600
a⊤ ↓
$\frac{a_T/M}{a_T/M}$
$\downarrow (p_T^* \approx p_T^*)$ $a_T/M \approx \tan(\phi_{acop}/2)\sin(\theta^*)$
$\downarrow \cos(heta_\eta^*) = anh(rac{\eta^ \eta^+}{2})$
$\phi^*_\eta \equiv \tan(\phi_{acop}/2)\sin(\theta^*_\eta)$



- θ^* is the scattering angle of the leptons relative to the beam direction in the dilepton rest frame and is sensitive to the effects of lepton momentum resolution.

- θ^*_η is an alternative way to estimate the scattering angle.

$\rightarrow \phi_{\eta}^{*}$ is based entirely on the measured track directions

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Relation between p_T^Z and ϕ_n^*



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$Z/\gamma^* \rightarrow \ell \ell$ event selection

Differential cross section measurement as a function of ϕ^*_η Differential cross section measurement as a function of ϕ^*_η and y_Z

2011 Data and Monte Carlo samples



(*) with the p_T^Z spectrum reweighted to RESBOS prediction

The ATLAS detector ϕ^*_η , an alternative variable to study the physics of ρ^-_T Differential cross section measurements of $Z/\gamma^* \to \ell \ell$

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Event selection

Collision event selection				
${\sf Stable \ beams} + {\sf working \ detectors} + \geq 1 \ {\sf good \ vertex} + {\sf single \ lepton \ trigger}$				
Good lepton selection				
Electron	$p_{\mathrm{T}}^{e\!\prime 1}>25$ GeV, $p_{\mathrm{T}}^{e\!\prime 2}>20$ GeV			
	$ \eta_{trk} < 2.4$ excluding $1.37 < \eta_{cl} < 1.52$			
	Medium identification			
Muon	$p_{ m T}^{\mu1}>$ 20 GeV, $p_{ m T}^{\mu2}>$ 20 GeV			
	$ \eta_{trk} < 2.4$			
	Isolated muons $+$ impact parameter requirements			
$Z/\gamma^* ightarrow \ell\ell$ event selection				
Charge	Opposite sign			
Invariant mass	66 GeV $< M_{\ell\ell} < 116$ GeV			
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Selection results



	Z ightarrow ee	$Z ightarrow \mu \mu$
# events selected in data	$1.22\cdot 10^{6}$	$1.69\cdot 10^6$
Fraction of background events	$(0.61 \pm 0.31)\%$	$(0.56 \pm 0.28)\%$

The ATLAS detector ϕ_n^* , an alternative variable to study the physics of p_T^Z Differential cross section measurements of $Z/\gamma^* \to \ell \ell$ $Z/\gamma^* \rightarrow \ell \ell$ event selection Differential cross section measurement as a function of ϕ^*_{η} Differential cross section measurement as a function of ϕ^*_{η} and y_ℓ

Methodology of the measurement

• The differential cross section:

$$rac{1}{\sigma} \cdot rac{d\sigma}{d\phi_\eta^*}$$

- Measured within the fiducial acceptance: $p_T^{\ell 1} > 20$ GeV, $p_T^{\ell 2} > 20$ GeV, $|\eta_\ell| < 2.4$, 66 GeV $< M_{\ell\ell} < 116$ GeV.
- Extrapolated to the acceptance: 66 GeV $< M_{\ell\ell} < 116$ GeV.

• The reconstructed ϕ_{η}^* spectrum is corrected for detector and QED final state radiation (FSR) effects using an unfolding technique (bin-by-bin correction) to the underlying "true" spectrum.

• The "true" spectrum is defined at different reference points referring to the amount of QED FSR corrections considered at the generator level: "Born", "Dressed", "Bare".



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 $Z/\gamma^* \rightarrow \ell \ell$ event selection Differential cross section measurement as a function of ϕ^*_{η} Differential cross section measurement as a function of ϕ^*_{η} and y_Z

Systematic uncertainties in ϕ_n^* bins

Source	Sys. unc.
Background	< 0.3%
Tracking	< 0.3%
MC statistics	0.13-0.9%
Unfolding	< 0.1%
Energy/momentum scale and resolution	< 0.1%
Efficiencies	< 0.05%
Pile-up	< 0.05%
QED final state radiation (*)	0.3%

(*) ϕ_{η}^{*} is the first result in LHC where the experimental uncertainty is smaller than the QED FSR uncertainty \rightarrow require dedicated work together with the theoreticians (See Z. Was's talk). Some of these studies can be found in http://annapurna.ifj.edu.pl/~wasm/phi-star.html, CERN-THESIS-2013-001 and arXiv:1303.2220.

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The fiducial differential cross section result



• Results of electron and muon channels were combined using a χ^2 minimisation method [Eur. Phys. J. C 63 (2009) 625, JHEP 01 (2010) 109].

- Statistic uncertainty: 0.3 1.6%.
- Systematic uncertainty: $\sim 0.5\% \sim 2$ times smaller than the latest p_T^Z measurement.

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Comparison with QCD predictions



• Best prediction by RESBOS CT10 (Non-perturbative+NNLL+NLO-scaled to NNLO) within 2-5%. Full theoretical uncertainty was not supported.

• Calculation from A. Banfi *et al.* CTEQ6m (NNLL+NNLO - Phys. Lett. B 715 (2012) 152-156) agrees with data within full theoretical uncertainty (pdf+scale+resummation).

• FEWZ CT10 (NNLO) undershoots data by 10%, as already observed in Phys. Lett. B 705 (2011) 415-434.

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Double differential cross section measurement



• The Differential cross section measurement as a function of ϕ_{η}^* is measured in 3 bins of y_Z .

• SHERPA and ALPGEN provide good descriptions.

• POWHEG+PYTHIA8 has better description than POWHEG+PYTHIA6 and POWHEG+HERWIG.

• POWHEG+HERWIG and MC@NLO (interfaced with HERWIG) give worse descriptions.

\rightarrow valuable information for the tuning of MC generators

Note: CT10 is used in all Monte Carlo samples except ALPGEN with CTEQ6L1

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K.O. DOAN

15/16

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Conclusion

The first measurement of the ϕ_{η}^* spectrum of Z/γ^* bosons at $\sqrt{s} = 7$ TeV of pp collisions using the full 2011 ATLAS data sample corresponding with 4.6 fb⁻¹ integrated luminosity.

- At low ϕ_{η}^{*} ($\phi_{\eta}^{*} < 0.1$): Data agree with RESBOS CT10 within 2%.
- At high φ^{*}_η: RESBOS CT10 provides better prediction than FEWZ CT10 and A. Banfi *et al.* CTEQ6m but can not reproduce the detailed shape of data better than 4%.
- Among MC generators: the best description by SHERPA. MC@NLO, POWHEG+HERWIG, POWHEG+PYTHIA6 underestimate the data at higher ϕ_{η}^* .
- The typical experimental precision ($\sim 0.5\%$) is now ten times better than the typical theoretical precision and therefore is valuable to constrain the theoretical predictions further.

The ATLAS detector ϕ^*_η , an alternative variable to study the physics of p_T^2 Differential cross section measurements of $Z/\gamma^* \to \ell\ell$

THANK YOU!

The ATLAS detector ϕ_{B}^{*} , an alternative variable to study the physics of p_{T}^{Z} Differential cross section measurements of $Z/\gamma^{*} \to \ell \ell$

Backup

The ATLAS detector $\phi^*_{\mathcal{H}}$, an alternative variable to study the physics of p_T^2 Differential cross section measurements of $Z/\gamma^* \to \ell\ell$

Low p_T^Z calculation

Dominant contributions to the differential cross section ($Q^2 = M_Z^2$):

$$\frac{d\sigma}{dp_T^2} \sim \frac{\alpha_s}{p_T^2} \ln\left(\frac{Q^2}{p_T^2}\right) \left[v_1 + v_2 \alpha_s \ln^2\left(\frac{Q^2}{p_T^2}\right) + v_3 \alpha_s^2 \ln^4\left(\frac{Q^2}{p_T^2}\right) + \ldots \right]$$

For $p_T
ightarrow$ 0, $\alpha_s \ln^2(Q^2/p_T^2)$ is large even when α_s is small

 \rightarrow the p_T distribution diverges.

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Collins Soper Dilepton rest frame



The Collins-Soper frame (in the dilepton rest frame):

When we boost from the hadron center-of-mass frame to the dilepton rest frame, the collinearity of the hadron momenta P_1 and P_2 is lost and they span a plane which we identify with the **X-Z**-plane.

We are still free to fix Z within this plane. In the Collins-Soper frame Z is chosen to bisect the angle between P_1 and $-P_2$. The angle between Z and P_1 and between Z and P_2 is called γ_{CS} .

 θ and ϕ define the lepton direction in CS frame.

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Multi-jet background: template fit

The template background sample: is selected from data requiring:

- Electron: failing the medium identification criteria (based on shower shape and track-quality variables)

- Muon: inverting the isolation requirement

Isolation = {scalar sum of p_T of tracks within $\Delta R = 0.2$ around muon < 10% p_T^{μ} }

Method: using a binned likelihood fit on $Z \rightarrow ee$ mass distributions.

 \rightarrow The real multi-jet background (QCD) is estimated by scaling the template background sample from data with parameter P_1 from fitting results.

The ATLAS detector ϕ^*_n , an alternative variable to study the physics of p_T^Z Differential cross section measurements of $Z/\gamma^* \to \ell \ell$

Selection results: muon channel



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Bin-by-bin Correction factors

To fiducial acceptance:

$$\begin{pmatrix} \frac{1}{\sigma} \end{pmatrix}_{fid} \begin{pmatrix} \frac{\Delta \sigma_i}{\Delta \phi_{\eta,i}^*} \end{pmatrix}_{fid} = \frac{C}{C^i} \cdot \frac{\left(N_{\text{data}}^i - N_{\text{bg}}^i \right)}{\Delta \phi_{\eta,i}^* \cdot \left(N_{\text{data}}^{tot} - N_{\text{bg}}^{tot} \right)}$$

$$C^i = \frac{N_{\text{MC,rec,cuts}}^i}{N_{\text{MC,gen,fid}}^i} \quad , \quad C = \frac{N_{\text{MC,rec,cuts}}^{tot}}{N_{\text{MC,gen,fid}}^{tot}}$$

To full acceptance:

$$\begin{pmatrix} \frac{1}{\sigma} \end{pmatrix}_{tot} \begin{pmatrix} \frac{\Delta \sigma_i}{\Delta \phi_{\eta,i}^*} \end{pmatrix}_{tot} = \frac{A}{A^i} \cdot \begin{pmatrix} \frac{1}{\sigma} \end{pmatrix}_{fid} \begin{pmatrix} \frac{\Delta \sigma_i}{\Delta \phi_{\eta,i}^*} \end{pmatrix}_{fid}$$
$$A^i = \frac{N_{\rm MC,gen,fid}^i}{N_{\rm MC,gen}^i} \quad , \quad A = \frac{N_{\rm MC,gen,fid}^{tot}}{N_{\rm MC,gen}^{tot}}$$

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QED FSR uncertainty

(

Compare POWHEG+PYTHIA6 (interfaced with PHOTOS $\sim 2^{nd}$ QED FSR simulation) with SHERPA (YFS formalism - 1st QED FSR simulation).

 \rightarrow To extract pure QED FSR effect:

$$\begin{split} \mathsf{QED}_{\mathsf{Bare}} &= \left(\left(\frac{1}{\sigma}\right)_{\mathit{fid}} \left(\frac{\Delta\sigma_{\mathit{i}}}{\Delta\phi_{\mathit{\eta}\,\mathit{i}}^*}\right)_{\mathit{fid}} \right)_{\mathsf{Bare}} \middle/ \left(\left(\frac{1}{\sigma}\right)_{\mathit{fid}} \left(\frac{\Delta\sigma_{\mathit{i}}}{\Delta\phi_{\mathit{\eta}\,\mathit{i}}^*}\right)_{\mathit{fid}} \right)_{\mathsf{Born}} \\ \mathsf{QED}_{\mathsf{Dressed}} &= \left(\left(\frac{1}{\sigma}\right)_{\mathit{fid}} \left(\frac{\Delta\sigma_{\mathit{i}}}{\Delta\phi_{\mathit{\eta}\,\mathit{i}}^*}\right)_{\mathit{fid}} \right)_{\mathsf{Dressed}} \middle/ \left(\left(\frac{1}{\sigma}\right)_{\mathit{fid}} \left(\frac{\Delta\sigma_{\mathit{i}}}{\Delta\phi_{\mathit{\eta}\,\mathit{i}}^*}\right)_{\mathit{fid}} \right)_{\mathsf{Born}} \end{split}$$

 \rightarrow To extract systematic uncertainty:

$$\left(\mathsf{QED}_{\mathsf{Bare}}\right)_{\mathsf{POWHEG}+\mathsf{PYTHIA6}} / \left(\mathsf{QED}_{\mathsf{Bare}}\right)_{\mathsf{SHERPA}} - 1$$

 $\left(\mathsf{QED}_{\mathsf{Dressed}}\right)_{\mathsf{POWHEG}+\mathsf{PYTHIA6}} / \left(\mathsf{QED}_{\mathsf{Dressed}}\right)_{\mathsf{SHERPA}} - 1$

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Effect of PDF on RESBOS predictions

