

# Amplitude detuning measurement with AC dipole

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## **Amplitude detuning measurement**

*motivations, method and past measurements*

## **Amplitude detuning with driven oscillations**

*theory, simulations*

## **LHC experimental data**

*results of 2012 MD*

## **Summary and outlook**

*possible improvements, experimental proof of  
theoretical approach*

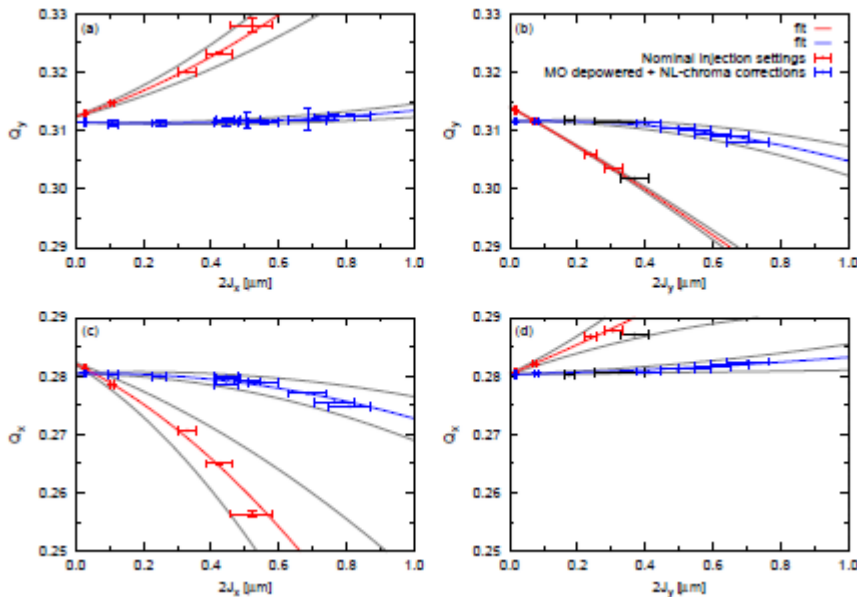
- Amplitude detuning measurement provides a handle on how well the non-linear model is understood
- Lots of discussion last year regarding Landau damping and beam stability:
  - How does the machine detuning contribute to the overall detuning (long-range beam-beam, Landau octupoles)?
  - Is there any compensation effect?
- The tune kicker is effective at injection. Large amplitude can be reached and fresh beams can be injected when the emittance becomes too large
- This is not the case at top energy. The AC dipole provides a non-destructive alternative allowing for large amplitude excitation at any energy

In the presence of non-linear fields the tune becomes amplitude dependent. Its behavior can be modeled by a polynomial of the form:

$$Q_x = Q_0 + \frac{\partial Q_x}{\partial \epsilon_x} \epsilon_x + \frac{\partial Q_x}{\partial \epsilon_y} \epsilon_y + \frac{1}{2} \frac{\partial^2 Q_x}{\partial \epsilon_x^2} \epsilon_x^2 + \frac{1}{2} \frac{\partial^2 Q_x}{\partial \epsilon_y^2} \epsilon_y^2 + \frac{\partial^2 Q_x}{\partial \epsilon_x \partial \epsilon_y} \epsilon_x \epsilon_y + \dots$$

where  $\epsilon = 2J$  and J is the action (normalized oscillation amplitude)

Measuring the tune as function of oscillation amplitude and applying a polynomial fit allows for a direct measurement of the detuning coefficients

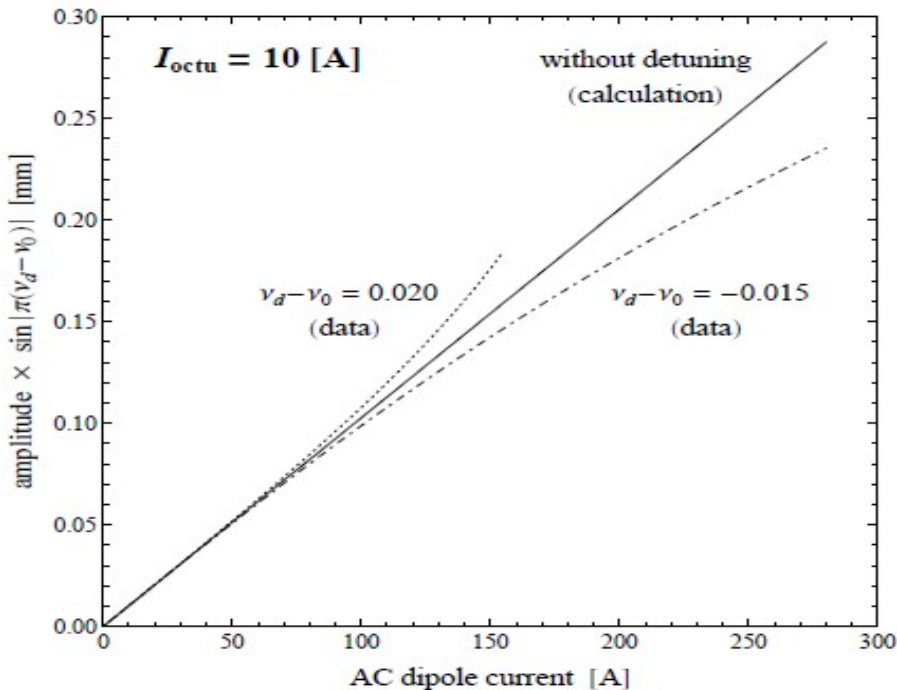
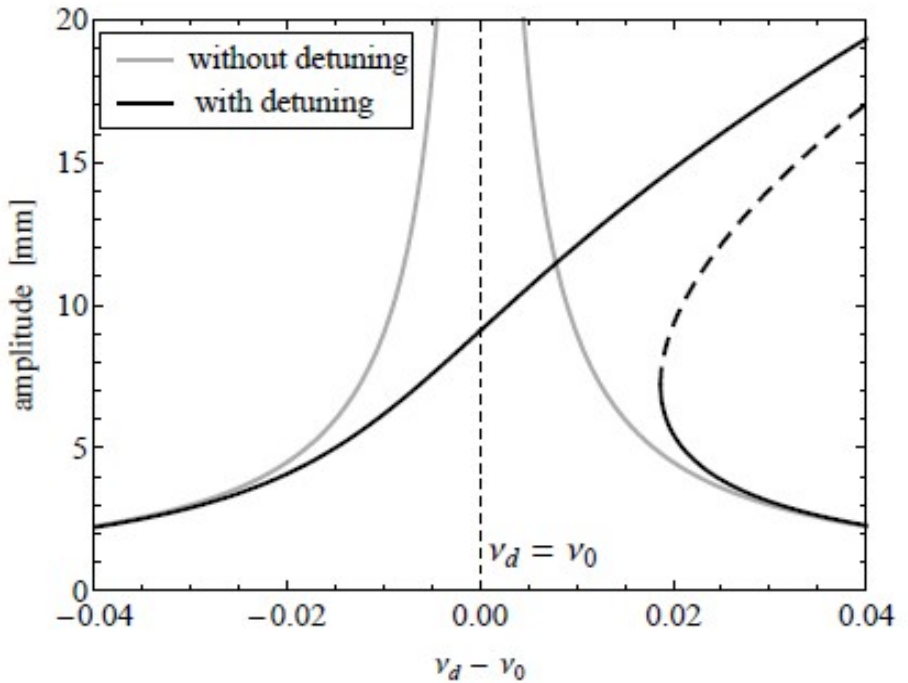


→ Example of amplitude detuning measurement at the LHC

→ Done at injection energy using the tune kicker

→ Reference: *E. Maclean et al. "Non-linear beam dynamics tests in the LHC: LHC dynamic aperture MD on Beam 2"*

- Reference: R. Miyamoto, PhD thesis
- In the presence of detuning the beam response is non-linear with driven oscillation amplitude
- This behavior provide information on detuning coefficient



- Measurements taken with strong octupolar field
- Plain curve is the linear case
- Dashed curves represent a fit of the AC dipole ramp: could be used to derive the detuning – **indirect measurement**

The Hamiltonian describing the linear motion of a single particle in the presence of AC dipole can be expressed as:

$$H_0(x, p_x, s, t) = \frac{1}{2}p_x^2 + \frac{1}{2}K_x(s)x^2 + \delta(s, t)x$$

$K_x(s)$  is the focusing strength and  $\delta(s, t)$  is the time dependent kick from the AC dipole

In the case of free oscillations the transverse coordinates can be parametrized as:

$$\begin{aligned} x(s) &= \sqrt{2J_x\beta_x(s)} \cos \phi_x \\ y(s) &= \sqrt{2J_y\beta_y(s)} \cos \phi_y \end{aligned}$$

For horizontal AC dipole excitation  $x(s)$  becomes:

$$x_D(s) = \sqrt{2J_x\beta_x(s)} \cos \phi_x + \sqrt{2A\beta'_x(s)} \cos \phi_D$$

→ The term  $J_x$  is driven by the non-adiabaticity of the ramping process of the AC dipole: in the presence of chromaticity or non-linearity it should be possible to observe the natural tune (*R. Tomas "Adiabaticity of the ramping process of an AC dipole" PRSTAB 8, 024401*).

Using the perturbative approach, the Hamiltonian in the presence of non-linear field becomes  $H_0 + H_1$ , where  $H_1$  (perturbation Hamiltonian) is given by:

$$H_1 = \frac{q}{p} \text{Re} \left[ \sum_{m=3}^{\infty} \frac{1}{m} [B_m(s) + iA_m(s)] (x + iy)^m \right]$$

$B_m(s)$  and  $A_m(s)$  are the normal and skew coefficients of the expansion of the magnetic field of a multipole of order  $m$ . Considering the case of a normal octupole:

$$H_4 = \frac{q}{p} \frac{B_4(s)}{4} (x^4 - 6x^2y^2 + y^4)$$

The detuning is given by: 
$$\Delta Q_{x,y} = \frac{1}{2\pi} \oint \frac{\partial}{\partial J_{x,y}} \langle H_4 \rangle ds$$

Using the parametrized coordinates in the presence of AC dipole we get:

$$\Delta Q_x = \frac{q}{p} \frac{3B_4}{8\pi} (\beta_x^2 J_x + 2\beta_x \beta'_x A) \quad \text{and} \quad \Delta Q_y = -\frac{q}{p} \frac{3B_4}{8\pi} (2\beta'_x \beta_y A + 2\beta_x \beta_y J_x)$$

→ The **direct term of the detuning measured with a single kick ( $A=0$ ) is a factor 2 smaller than the one measured with AC dipole ( $A \gg J$ )**. Effect of the AC dipole on optics are assumed negligible ( $\beta' = \beta$ ). **The cross term is not affected.**

The perturbation Hamiltonian of a multipole of order  $2n$  is expressed as:

$$H_{2n} = \frac{qB_{2n}}{2np} \operatorname{Re} \left[ (x + iy)^{2n} \right]$$

This gives for the detuning coefficient in the case of free oscillations:

$$\Delta Q_x = \frac{qB_{2n}}{2np} \frac{2^{-n}}{2\pi} \binom{2n}{0} \binom{2n}{n} \beta_x^n n J_x^{n-1}$$

$$\Delta Q_y = -\frac{qB_{2n}}{2np} \frac{2^{-n+1}}{2\pi} \binom{2n}{2} \binom{2n-2}{n-1} \beta_x^{n-1} \beta_y J_x^{n-1}$$

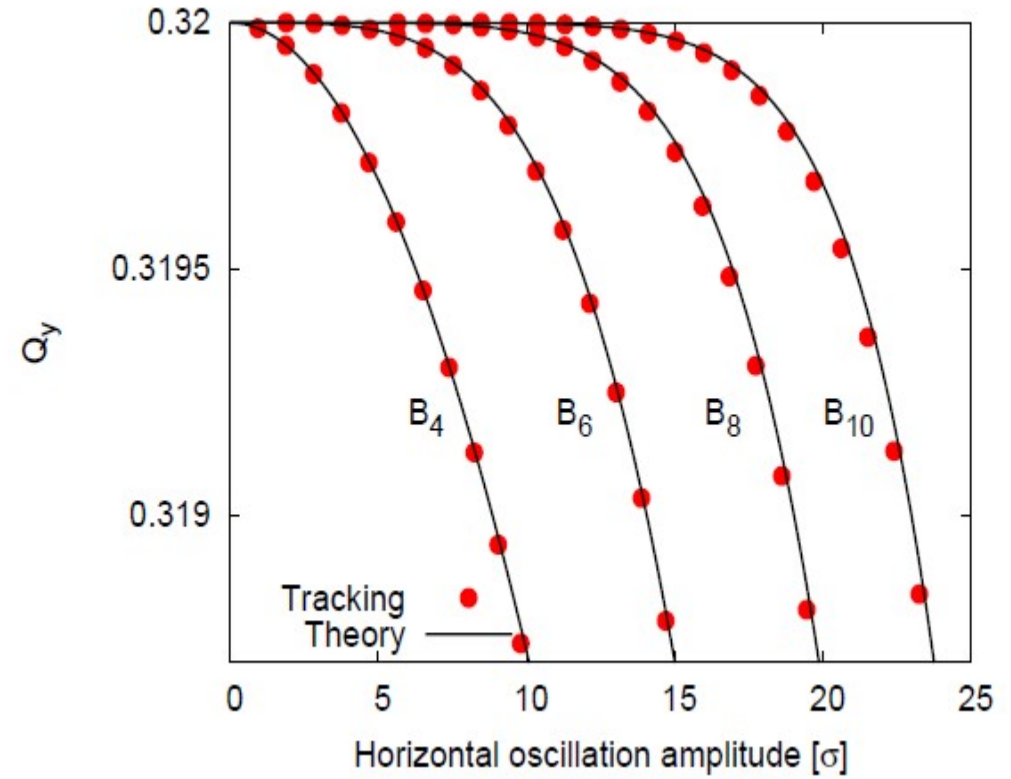
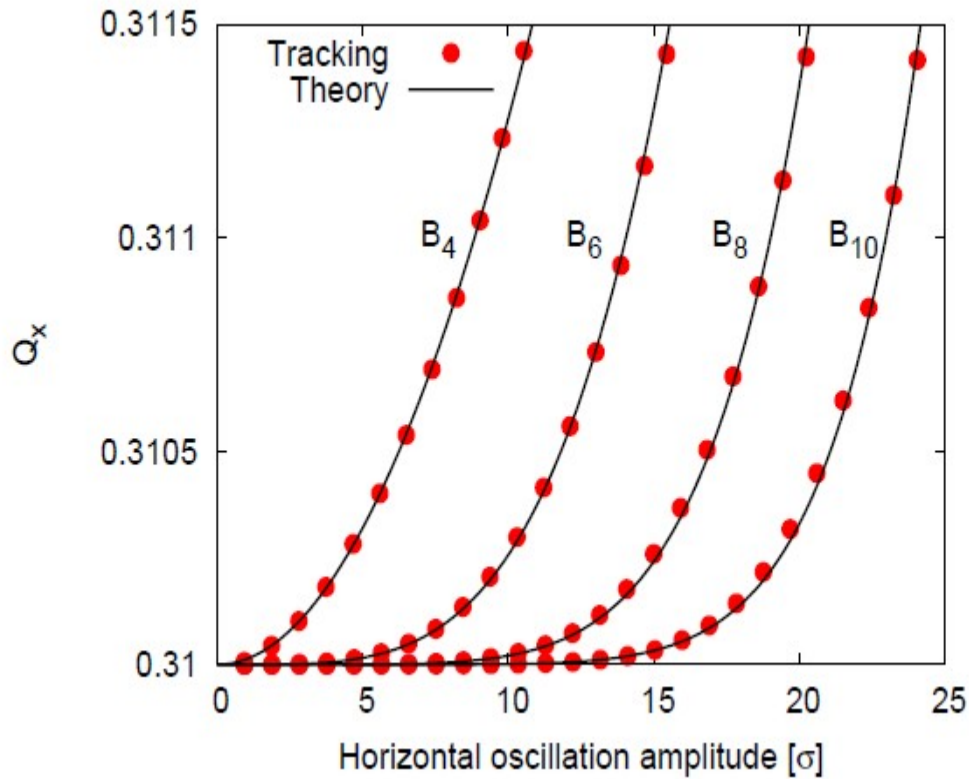
And in the case of driven oscillations

$$\Delta Q_x = \frac{qB_{2n}}{2np} \frac{2^{-n+1}}{2\pi} \binom{2n}{2n-2} \binom{2n-2}{n-1} \beta_x \beta_x'^{n-1} A^{n-1}$$

$$\Delta Q_y = -\frac{qB_{2n}}{2np} \frac{2^{-n+1}}{2\pi} \binom{2n}{2} \binom{2n-2}{n-1} \beta_y \beta_x'^{n-1} A^{n-1}$$

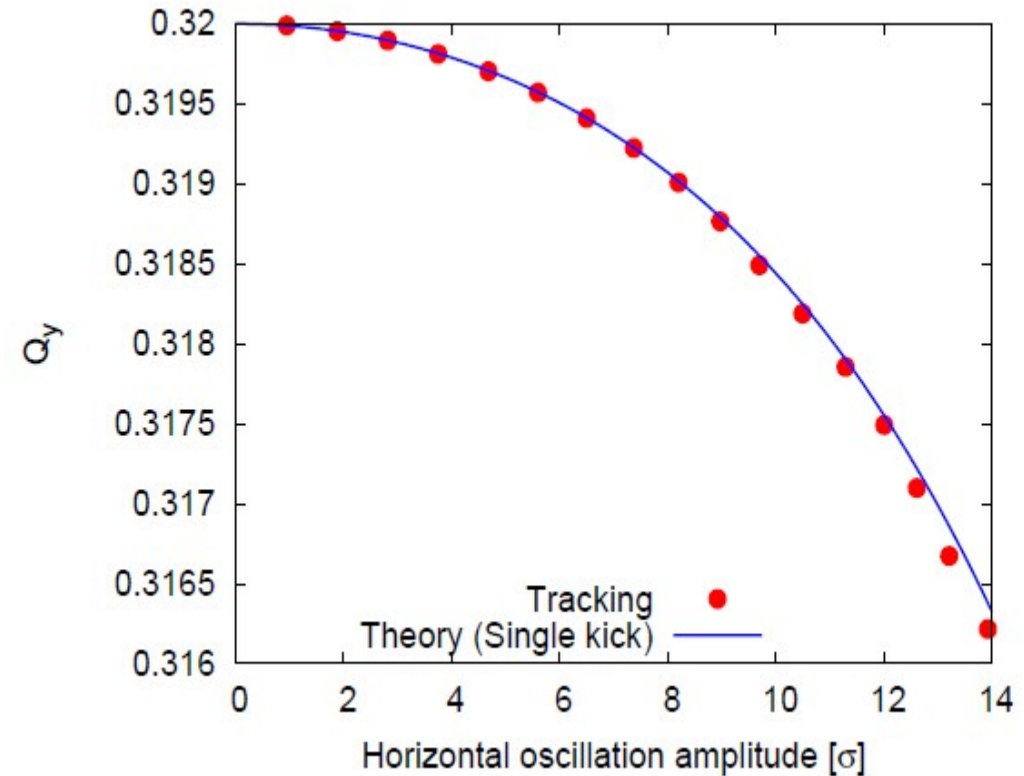
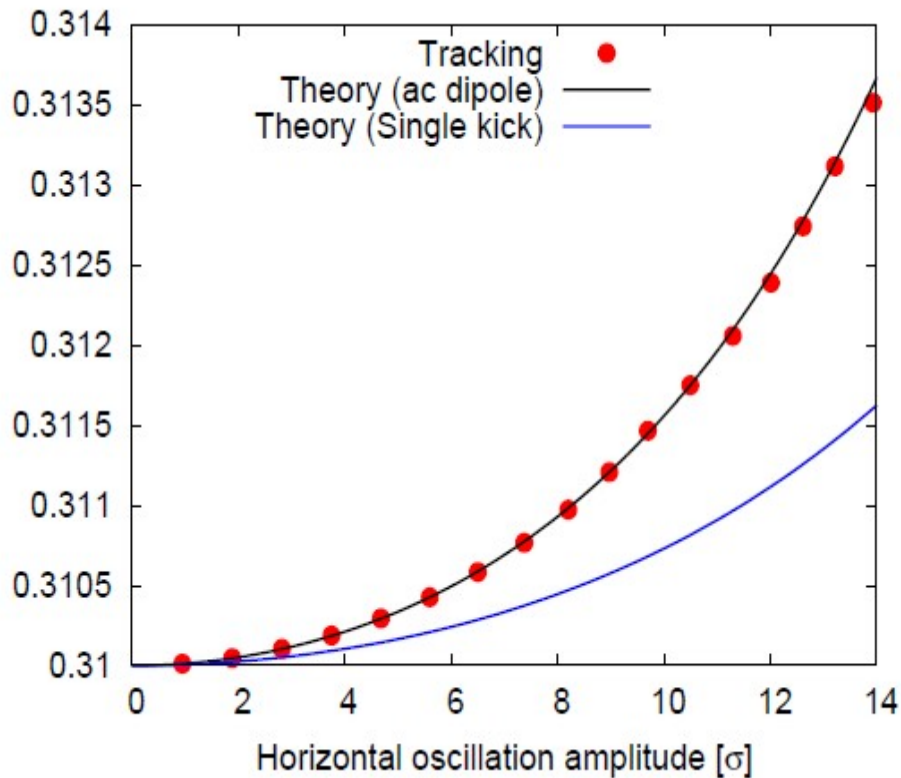
→ **Factor  $n$  difference between free and driven oscillation for the direct term**





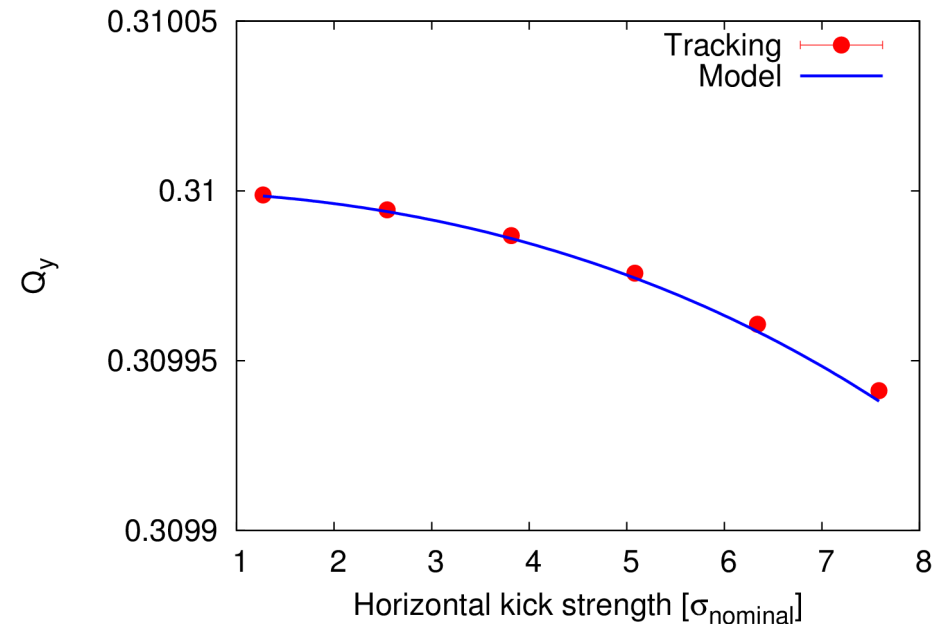
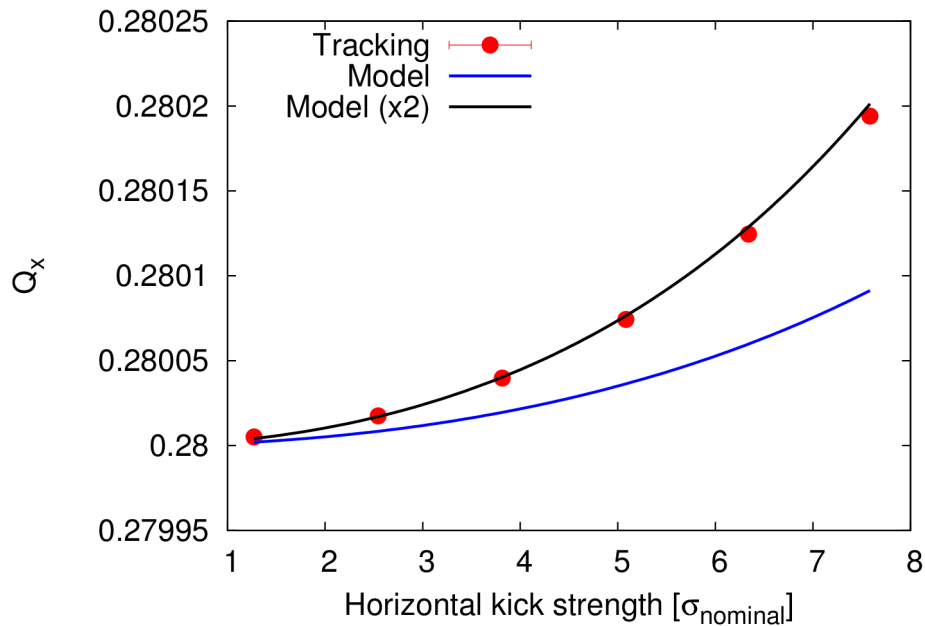
→ Built a simple tracking code allowing to study driven oscillation in the presence of arbitrary non-linear field components (behavior also checked with MADX, much slower)

→ Start with a single multipole: **excellent agreement between theory and tracking** up to a  $B_{10}$  magnetic field component



→ Repeat the same study with the combination of several multipoles from  $B_4$  to  $B_{10}$

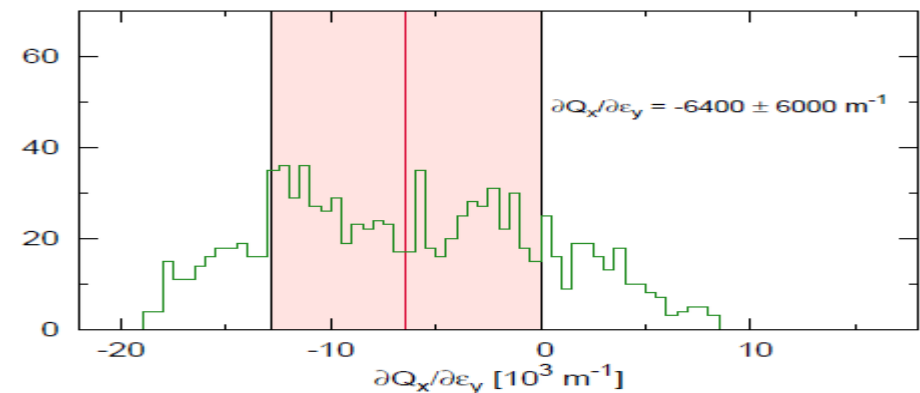
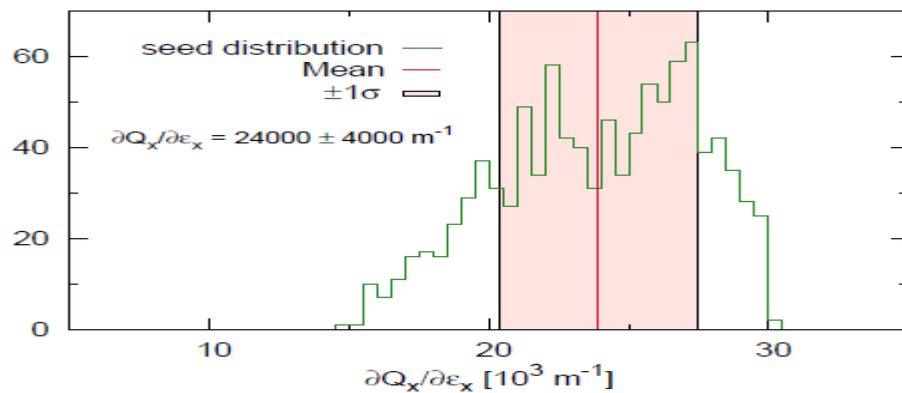
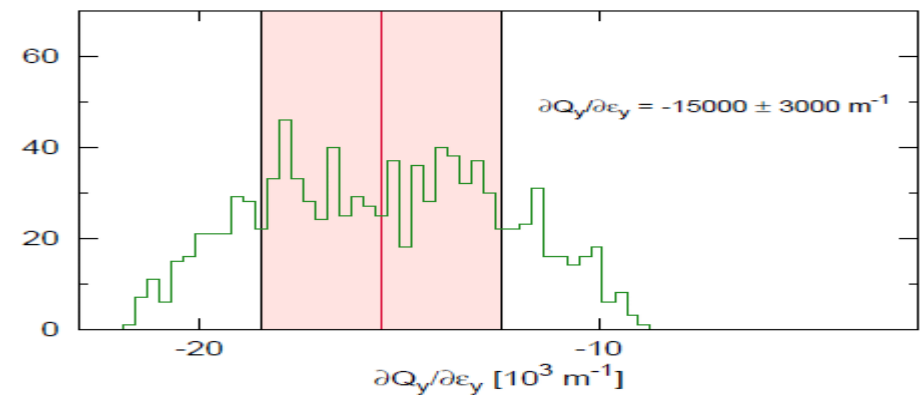
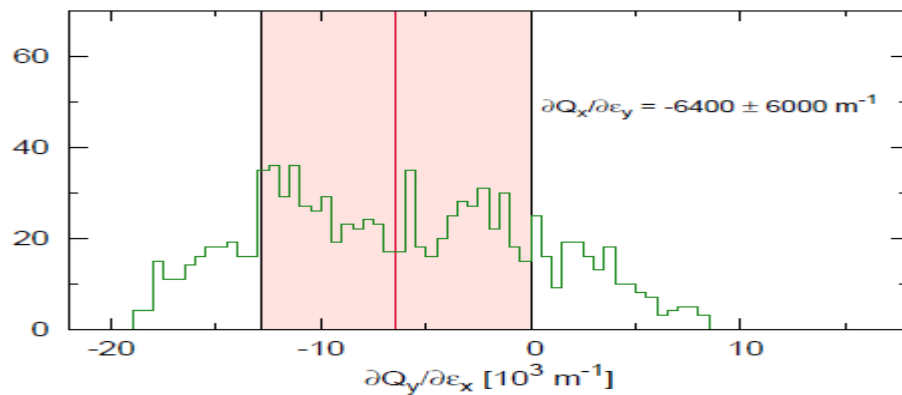
→ Again an excellent agreement is observed between theory and tracking, **the under-estimation using free oscillation model is clearly observed**



→ MADX model constructed using a thin lattice with magnetic field errors from WISE of order from  $(B_3, A_3)$  up to  $(B_{15}, A_{15})$ . Includes second order effects. Example of flat top optics in this case

→ The model detuning coefficients are computed using PTC, the tracking is performed using the AC dipole module in MADX

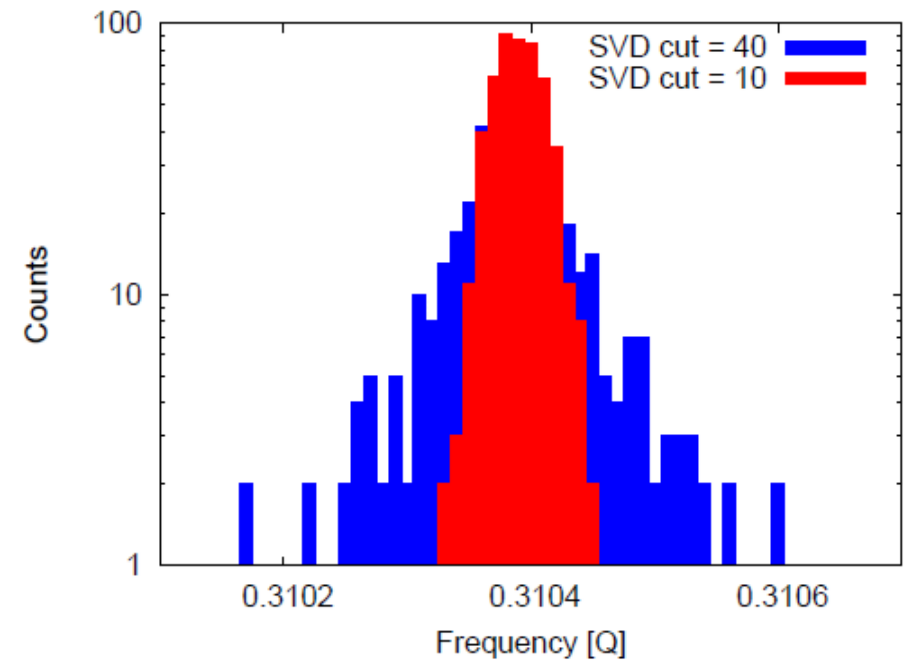
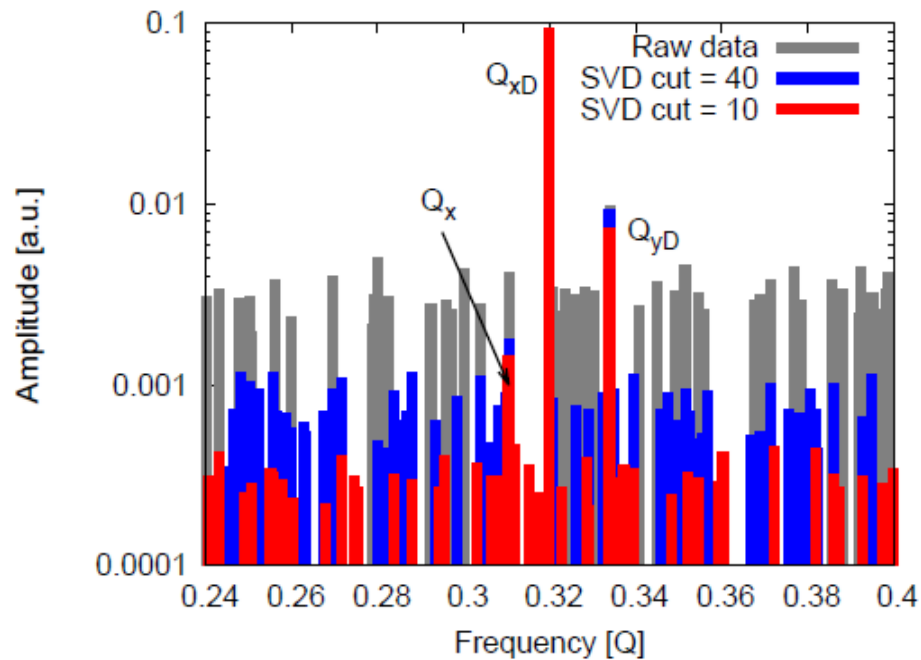
→ The detuning is dominated by the octupolar field components: **a factor 2 difference is observed between PTC and AC dipole tracking – consistent with theory**



→ Uncertainties on the model apply in the presence of coupling or misalignments. **Due to technical limitation in MADX and PTC only the effects of coupling could be studied**

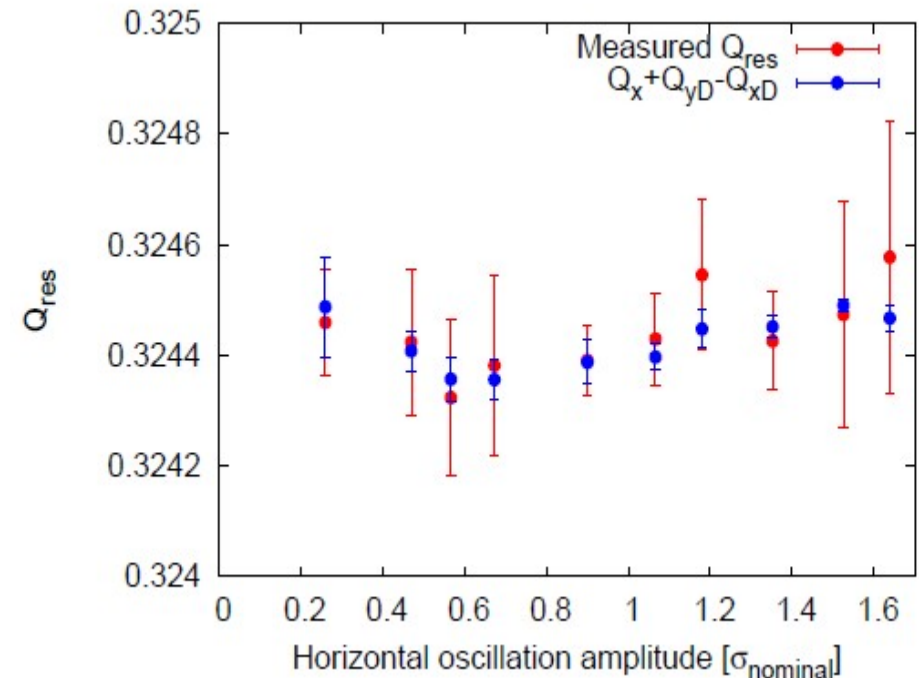
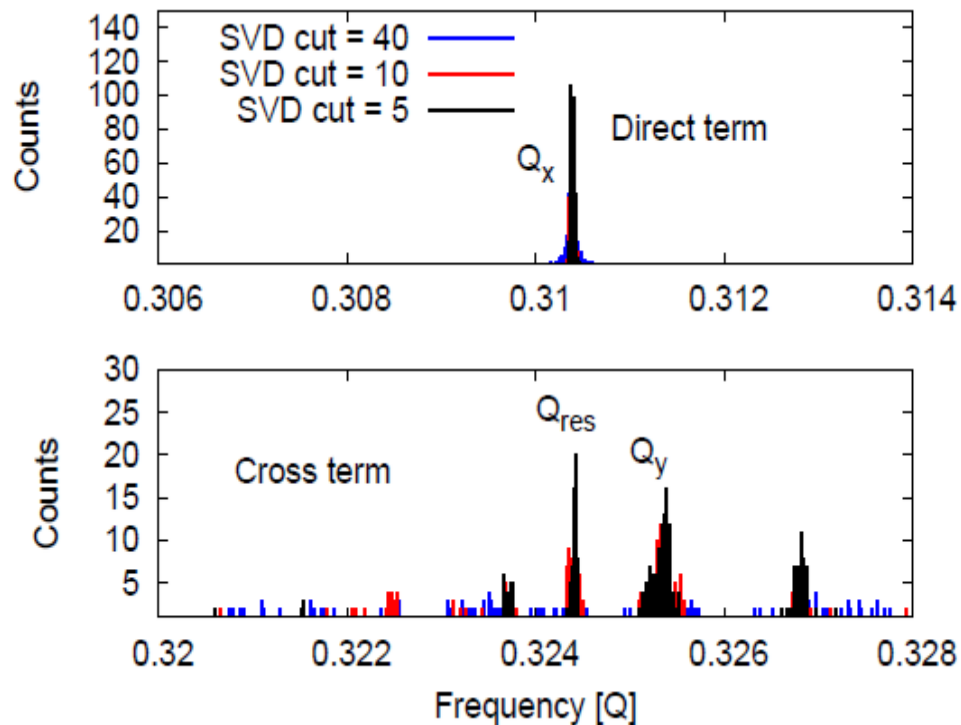
→ Distribution generated from random distribution of coupling amplitude and phase

→ **Uncertainties of the order of 20% for the direct term and 100% for cross term were derived from these simulations**



→ Although the natural tune spectral should be excited during the ramping process due to the non-linearities its amplitude remains much lower than the drive spectral line and most of the time below the noise floor

→ Noise removal is possible using SVD post-processing of the data keeping only the modes with physical meaning for data analysis: **SVD “cleaning” of the raw data allows to clearly observe the natural tune spectral line** (left) and reduces the spread from all BPMs (right)



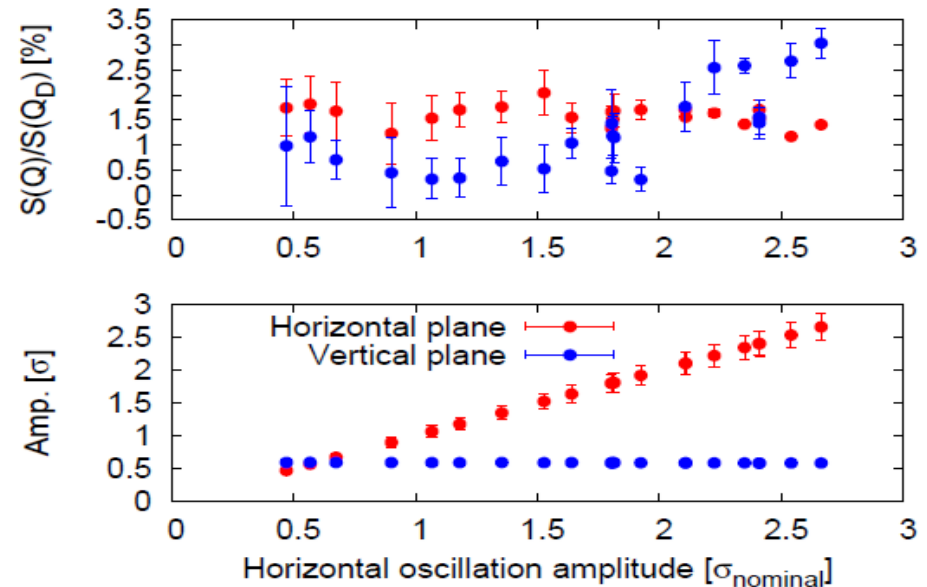
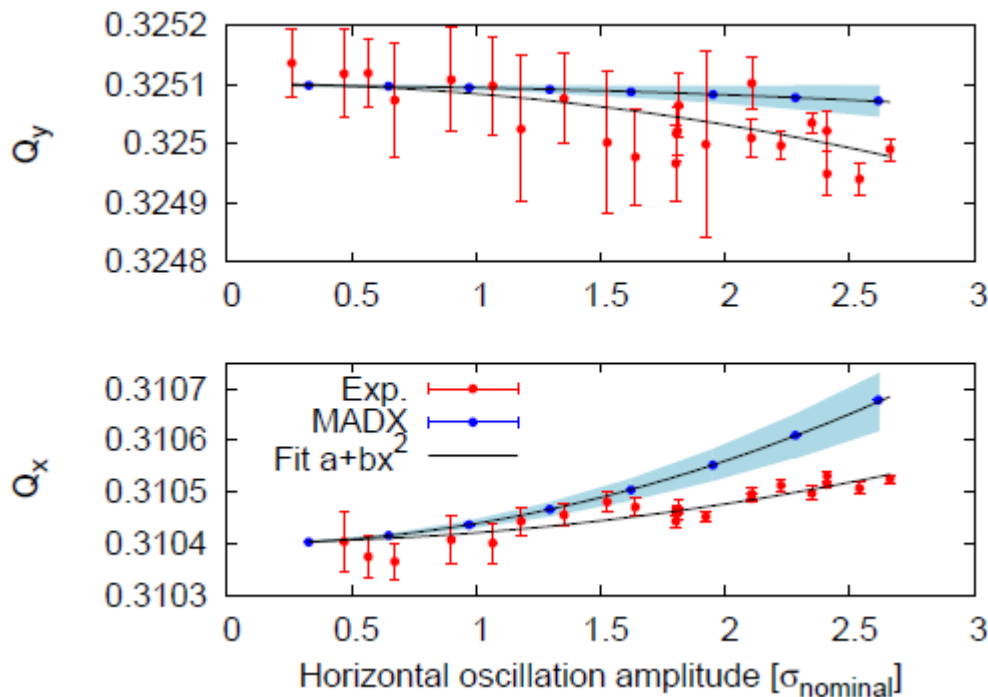
→ **Non-linearities not only allow to observe the natural tune but also excite non-linear resonances: should be avoided if possible**

→ This was strongly observed in the cross plane data and degraded the data quality and our ability to cleanly measure the natural tune

→ Here an example of an octupolar resonance which frequency was measured during the amplitude scan. Expected and measured frequencies are in good agreement

→ The importance of non-adiabaticity can be estimated by looking at the ratio of drive and natural tune spectral line

→ Horizontal plane is constant. Increase in the vertical plane is attributed to coupling



|                        | $\partial Q_x / \partial \epsilon_x$ $[\text{m}^{-1}]$ | $\partial Q_y / \partial \epsilon_x$ $[\text{m}^{-1}]$ |
|------------------------|--|--|
| Data                   | $18528 \pm 10\%$                                       | $-17387 \pm 23\%$                                      |
| AC dipole model        | $45757 \pm 1\%$  | $-4492 \pm 2\%$  |
| Free oscillation model | 22666  | -5493  |

→ Disagreement of a factor 2-3 between the data and the model

→ Measurement taken after non-linear correction: only very small errors left: How well are they understood?

- **Amplitude detuning equations in the presence of AC dipole were derived using the perturbation Hamiltonian:**
  - Valid for multipoles of arbitrary order
  - Driven oscillations cannot be approximated by the free oscillation model :correction factors are required
  - Direct amplitude detuning measurement with AC dipole is possible
- **Experimental protocol well defined:**
  - SVD cleaning essential for natural tune measurement
  - Some optimization required (working point, excitation of resonances,...) - see next slide
- **2012 MD results:**
  - First direct amplitude detuning measurement using AC dipole performed at the LHC
  - Disagreement of a factor 2-3 between measurement and model
  - Done after non-linear corrections: remaining errors are very small and possibly not well understood: difficulties to build an accurate model
  - Uncertainties associated to misalignments could not be assessed
- Very nice achievement for a first try, lots of lessons learned

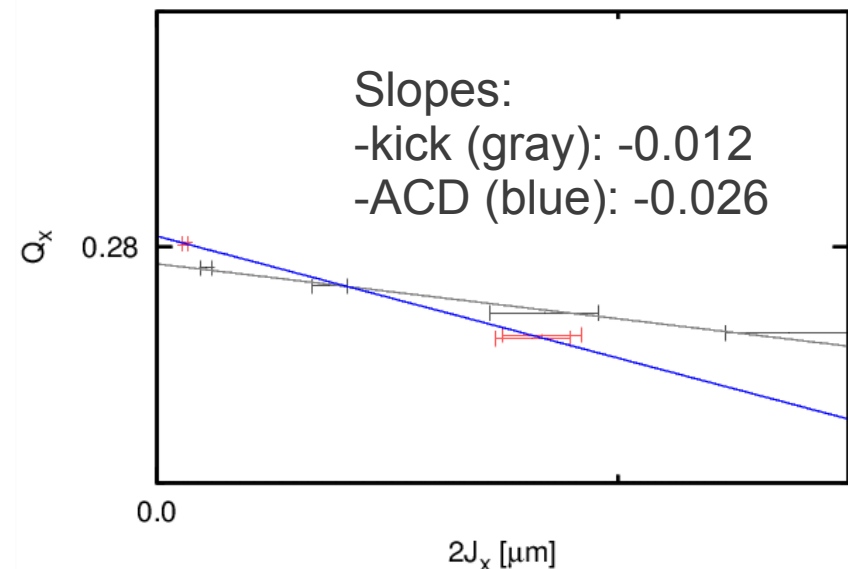


## • Possible improvements:

- Coupling: keep the tunes well separated avoid driving in between  $Q_x$  and  $Q_y$  (as was done in 2012)
- Non-linear resonances: single plane excitation – carefully choose the working point
- Drive to larger amplitudes: have a single bunch in the machine would reduce the losses
- Model: perform the measurement without non-linear corrections as a first check-effect of misalignments?

## • Experimental proof of theoretical derivations:

- Repeat at injection with strong octupoles – scan octupole current
- Compare single kick data with AC dipole data: model independent



- Maybe a first hint of experimental proof at the LHC?
- Data taken during the non-linear MD at injection

**Thank you for your attention!**