

BFKL Pomeron, reggeized gluons and Bern-Dixon-Smirnov amplitudes

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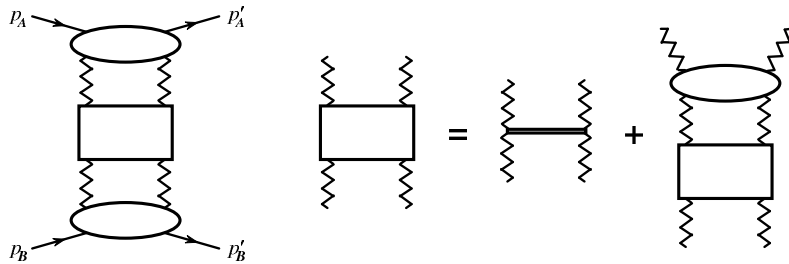
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1 BFKL equation (1975)

Diagrams for the total cross-section



Balitsky-Fadin-Kuraev-Lipatov equation

$$E \Psi(\vec{\rho}_1, \vec{\rho}_2) = H_{12} \Psi(\vec{\rho}_1, \vec{\rho}_2), \quad \sigma_t \sim s^\Delta, \quad \Delta = -\frac{\alpha_s N_c}{2\pi} E$$

BFKL Hamiltonian

$$H_{12} = \ln |p_1 p_2|^2 + \frac{1}{p_1 p_2^*} (\ln |\rho_{12}|^2) p_1 p_2^* \\ + \frac{1}{p_1^* p_2} (\ln |\rho_{12}|^2) p_1^* p_2 - 4\psi(1), \quad \rho_{12} = \rho_1 - \rho_2$$

Möbius invariance and conformal weights (L. (1986))

$$\rho_k \rightarrow \frac{a\rho_k + b}{c\rho_k + d},$$

$$m = \gamma + n/2, \quad \tilde{m} = \gamma - n/2, \quad \gamma = 1/2 + i\nu$$

2 Integrability at $N_c \rightarrow \infty$

Bartels-Kwiecinski-Praszalowicz equation (1980)

$$E \Psi(\vec{\rho}_1, \dots, \vec{\rho}_n) = H \Psi(\vec{\rho}_1, \dots, \vec{\rho}_n), \quad H = \sum_{k < l} \frac{\vec{T}_k \vec{T}_l}{-N_c} H_{kl}$$

Holomorphic factorization at large N_c (L. (1988))

$$\Psi(\vec{\rho}_1, \vec{\rho}_2, \dots, \vec{\rho}_n) = \sum_{r,s} a_{r,s} \Psi_r(\rho_1, \dots, \rho_n) \Psi_s(\rho_1^*, \dots, \rho_n^*)$$

Transfer and monodromy matrices (L. (1993))

$$T(u) = \text{tr } t(u), \quad t(u) = L_1 L_2 \dots L_n = \sum_{r=0}^n u^{n-r} q_r,$$

$$L_k = \begin{pmatrix} u + \rho_k p_k & p_k \\ -\rho_k^2 p_k & u - \rho_k p_k \end{pmatrix}, \quad \hat{l} = u \hat{1} + i \hat{P}$$

Yang-Baxter equation (L. (1993))

$$t_{r'_1}^{s_1}(u) t_{r'_2}^{s_2}(v) l_{r_1 r_2}^{r'_1 r'_2}(v-u) = l_{s'_1 s'_2}^{s_1 s_2}(v-u) t_{r_2}^{s'_2}(v) t_{r_1}^{s'_1}(u)$$

Heisenberg spin model (L. (1994); F.,K. (1995))

3 Pomeron in $N = 4$ SUSY

BFKL kernel in two loops (F., L. (1998))

$$\omega = 4\hat{a} \chi(n, \gamma) + 4\hat{a}^2 \Delta(n, \gamma), \quad \hat{a} = g^2 N_c / (16\pi^2),$$

$$\chi(n, \gamma) = 2\Psi(1) - \Psi(\gamma + |n|/2) - \Psi(1 - \gamma + |n|/2)$$

Non-analytic terms (K.,L. (2000))

$$\Delta_{QCD}(n, \gamma) = c_0 \delta_{|n|,0} + c_2 \delta_{|n|,2} + \dots, \quad c_0^{N=4} = c_2^{N=4} = 0$$

Hermitian separability in $N = 4$ SUSY (K.,L. (2000))

$$\Delta(n, \gamma) = \phi(M) + \phi(M^*) - \frac{\rho(M) + \rho(M^*)}{2\hat{a}/\omega}, \quad M = \gamma + \frac{|n|}{2},$$

$$\rho(M) = \beta'(M) + \frac{1}{2}\zeta(2), \quad \beta'(z) = \frac{1}{4} \left[\Psi'\left(\frac{z+1}{2}\right) - \Psi'\left(\frac{z}{2}\right) \right]$$

Maximal transcendentality (K.,L. (2002))

$$\phi(M) = 3\zeta(3) + \Psi''(M) - 2\Phi(M) + 2\beta'(M) \left(\Psi(1) - \Psi(M) \right)$$

$$\Phi(M) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k+M} \left(\Psi'(k+1) - \frac{\Psi(k+1) - \Psi(1)}{k+M} \right)$$

4 Three loops for the BFKL kernel

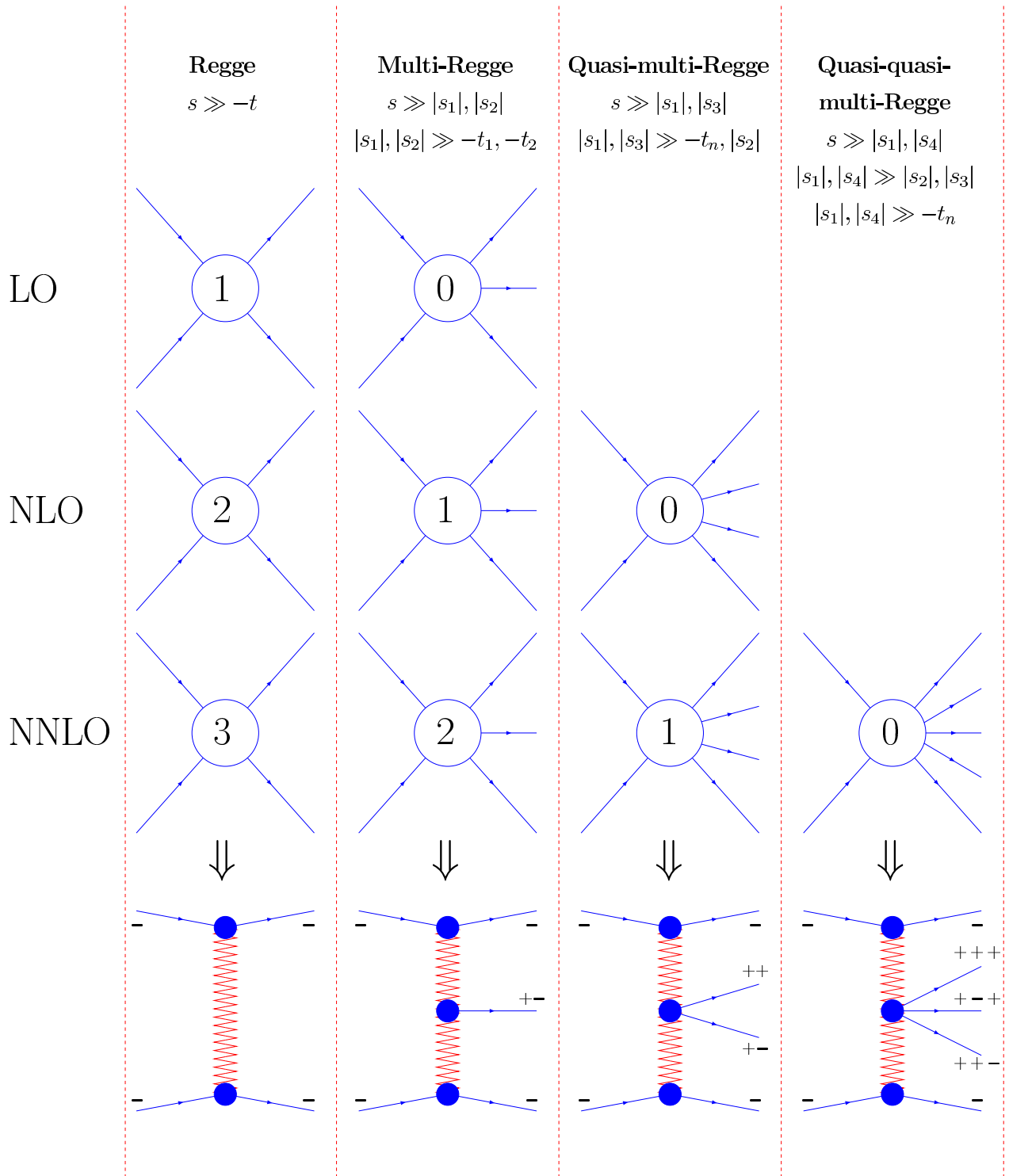


Figure 1: Diagrams contributing to the BFKL kernel

5 Effective action approach

Locality in the rapidity

$$y = \frac{1}{2} \ln \frac{\epsilon_k + |k|}{\epsilon_k - |k|}, \quad |y - y_0| < \eta, \quad \eta \ll \ln s$$

Gluon and Reggeized gluon fields

$$v_\mu(x) = -iT^a v_\mu^a(x), \quad A_\pm(x) = -iT^a A_\pm^a(x)$$

Local gauge transformations

$$\delta v_\mu(x) = \frac{1}{g} [D_\mu, \chi(x)], \quad \delta \psi(x) = -\chi(x) \psi(x), \quad \delta A_\pm(x) = 0$$

Effective action for reggeized gluons (L., 1995)

$$S = \int d^4x (L_0 + L_{ind}), \quad L_0 = i\bar{\psi} \hat{D} \psi + \frac{1}{2} \text{Tr} G_{\mu\nu}^2$$

$$L_{ind} = \text{Tr} (L_{ind}^k + L_{ind}^{GR}), \quad L_{ind}^k = -\partial_\mu A_+^a \partial_\mu A_-^a,$$

$$L_{ind}^{GR} = -\frac{1}{g} \partial_+ P \exp \left(-g \frac{1}{2} \int_{-\infty}^{x^+} v_+(x') d(x')^+ \right) \partial_\sigma^2 A_-$$

$$-\frac{1}{g} \partial_- P \exp \left(-g \frac{1}{2} \int_{-\infty}^{x^-} v_-(x') d(x')^+ \right) \partial_\sigma^2 A_+$$

6 Maximal helicity violation

BDS amplitudes in $N = 4$ SUSY at $N_c \gg 1$ (2005)

$$A^{a_1, \dots, a_n} = \sum_{\{i_1, \dots, i_n\}} \text{Tr} T^{a_{i_1}} T^{a_{i_2}} \dots T^{a_{i_n}} f(p_{i_1}, p_{i_2}, \dots, p_{i_n}),$$

$$f(p_{i_1}, p_{i_2}, \dots, p_{i_n}) = f_B(p_{i_1}, p_{i_2}, \dots, p_{i_n}) M_n(p_{i_1}, p_{i_2}, \dots, p_{i_n})$$

Invariant amplitudes

$$\ln M_n = \sum_{l=1}^{\infty} a^l \left(f^{(l)}(\epsilon) \left(\hat{I}_n^{(l)}(l\epsilon) + F_n^{(1)}(0) \right) + C^{(l)} \right),$$

$$C^{(1)} = 0, \quad C^{(2)} = -\zeta_2^2/2, \quad f^{(l)}(\epsilon) = f_0^{(l)} + \epsilon f_1^{(1)} + \epsilon^2 f_2^{(l)}$$

Cusp anomalous dimension

$$f_0^{(l)} = \frac{1}{4} \gamma_K^{(l)}, \quad f_1 = -a\zeta_3/2 + a^2(2\zeta_5 + 5\zeta_2\zeta_3/3) + \dots$$

Infraredly divergent contribution

$$\hat{I}_n^{(1)}(\epsilon) = -\frac{1}{2\epsilon^2} \sum_{i=1}^n \left(\frac{\mu^2}{-s_{i,i+1}} \right)^\epsilon, \quad a = \frac{\alpha N_c}{2\pi} (4\pi e^{-\gamma})^\epsilon$$

Finite remainder

$$F_n^{(1)}(0) = \lim_{\epsilon \rightarrow 0} F_n^{(1)}(\epsilon)$$

7 BDS amplitude at $s/t \rightarrow \infty$

Elastic amplitude

$$M_{2 \rightarrow 2} = \Gamma(t) \left(\frac{-s}{\mu^2} \right)^{\omega(t)} \Gamma(t)$$

Reggeized gluon trajectory

$$\begin{aligned} \omega(t) &= -\frac{\gamma(a)}{4} \ln \frac{-t}{\mu^2} + \int_0^a \frac{da'}{a'} \left(\frac{\gamma_K(a')}{4\epsilon} + \beta(a') \right) \\ &= \left(-\ln \frac{-t}{\mu^2} + \frac{1}{\epsilon} \right) a + \left(\zeta_2 \left(\ln \frac{-t}{\mu^2} - \frac{1}{2\epsilon} \right) - \frac{\zeta_3}{2} \right) a^2 \\ &\quad + \left(-\frac{11}{2} \zeta_4 \left(\ln \frac{-t}{\mu^2} - \frac{1}{3\epsilon} \right) + \frac{6\zeta_5 + 5\zeta_2\zeta_3}{3} \right) a^3 + \dots \end{aligned}$$

Reggeon residues

$$\begin{aligned} \ln \Gamma(t) &= \ln \frac{-t}{\mu^2} \int_0^a \frac{da'}{a'} \left(\frac{\gamma(a')}{8\epsilon} + \frac{\beta(a')}{2} \right) + \frac{C(a)}{2} + \frac{\gamma(a)}{2} \zeta_2 \\ &\quad - \int_0^a \frac{da'}{a'} \ln \frac{a}{a'} \left(\frac{\gamma(a')}{4\epsilon^2} + \frac{\beta(a')}{\epsilon} + \delta(a') \right) \end{aligned}$$

Dual representation

$$M_{2 \rightarrow 2} = \Gamma(s) \left(\frac{-t}{\mu^2} \right)^{\omega(s)} \Gamma(s)$$

8 Production amplitude

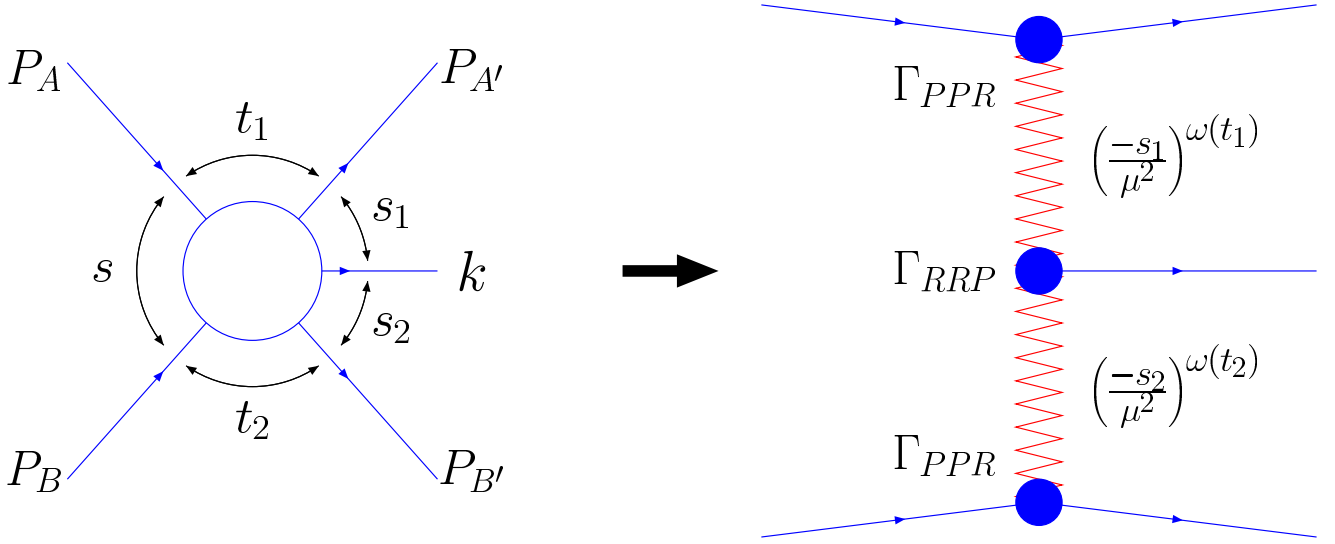


Figure 2: Production in the multi-Regge regime

Reggeon-Reggeon-gluon vertex

$$\ln \Gamma(t_2, t_1, \ln -\kappa) = -\frac{\gamma_K(a)}{16} \left(\ln^2 \frac{-\kappa}{\mu^2} - \ln^2 \frac{-t_1}{-t_2} - \zeta_2 \right) - \frac{1}{2} \int_0^a \frac{da'}{a'} \ln \frac{a}{a'} \left(\frac{\gamma_K(a')}{4\epsilon^2} + \frac{\beta(a')}{\epsilon} + \delta(a') \right) - \frac{1}{2} \left(\omega(t_1) + \omega(t_2) - \int_0^a \frac{da'}{a'} \left(\frac{\gamma_K(a')}{4\epsilon} + \beta(a') \right) \right) \ln \frac{-\kappa}{\mu^2},$$

$$-\kappa = \frac{(-s_1)(-s_2)}{(-s)}, \quad |\kappa| = \vec{k}_\perp^2$$

9 Two particle production

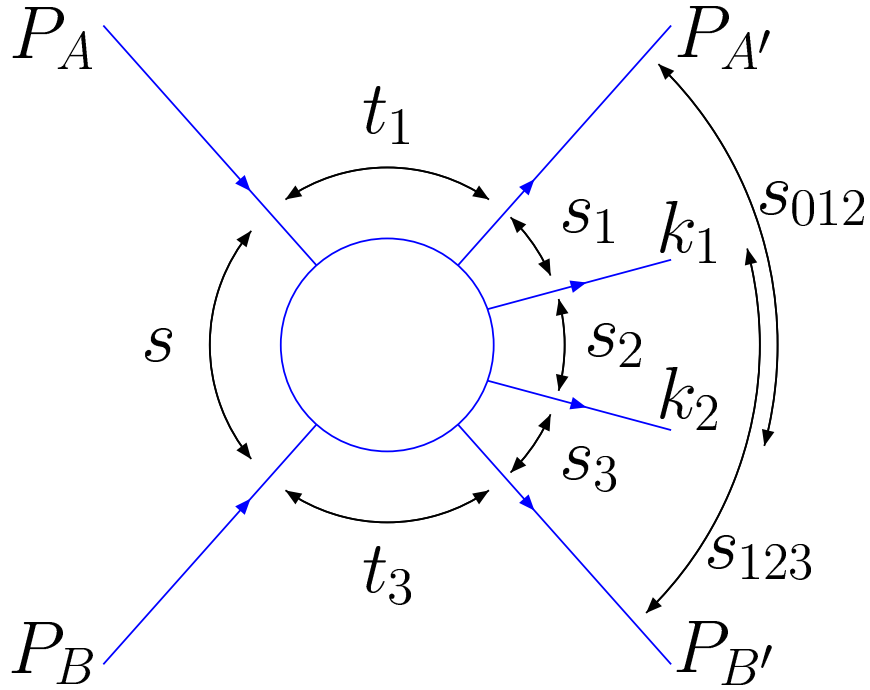


Figure 3: Multi-Regge kinematics

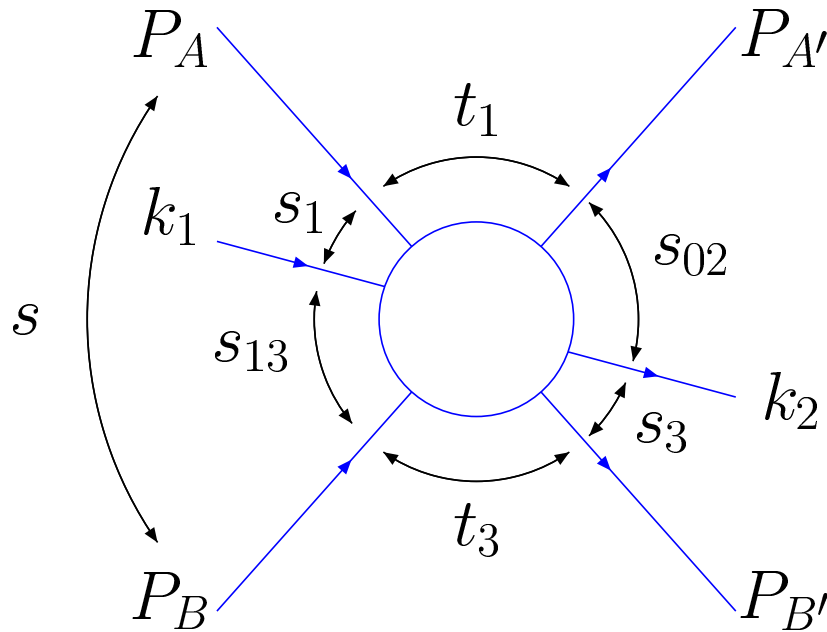
Regge factorization violation at

$$s, s_2 > 0; \quad s_1, s_3, s_{012}, s_{123} < 0$$

$$\frac{M_{2 \rightarrow 4}}{\Gamma(t_1)\Gamma(t_3)} = \exp \left[\frac{\gamma_K(a)}{4} i\pi \left(\ln \frac{t_1 t_2}{(\vec{k}_1 + \vec{k}_2)^2 \mu^2} - \frac{1}{\epsilon} \right) \right]$$

$$\times \left(\frac{-s_1}{\mu^2} \right)^{\omega(t_1)} \Gamma(t_2, t_1) \left(\frac{-s_2}{\mu^2} \right)^{\omega(t_2)} \Gamma(t_3, t_2) \left(\frac{-s_3}{\mu^2} \right)^{\omega(t_3)}$$

10 Tree particle scattering



Regge factorization violation at
 $s, t'_2 > 0; s_1, s_3, s_{13}, s_{02} < 0$

$$\frac{M_{3 \rightarrow 3}}{\Gamma(t_1)\Gamma(t_3)} = \exp \left[\frac{\gamma_K(a)}{4} i\pi \ln \frac{(\vec{q}_1 + \vec{q}_3 - \vec{q}_2)^2 \vec{q}_2^2}{(\vec{q}_1 - \vec{q}_2)^2 (\vec{q}_2 - \vec{q}_3)^2} \right]$$

$$\times \left(\frac{-s_1}{\mu^2} \right)^{\omega(t_1)} \Gamma(t_2, t_1, \ln \kappa_{12} + i\pi) \left(\frac{t'_2 e^{-i\pi}}{\mu^2} \right)^{\omega(t_2)}$$

$$\times \Gamma(t_2, t_1, \ln \kappa_{23} + i\pi) \left(\frac{-s_3}{\mu^2} \right)^{\omega(t_3)}, \quad \kappa_{12} = \vec{k}_1^2, \quad \kappa_{23} = \vec{k}_2^2$$

11 Regge cuts in j_2 -plane

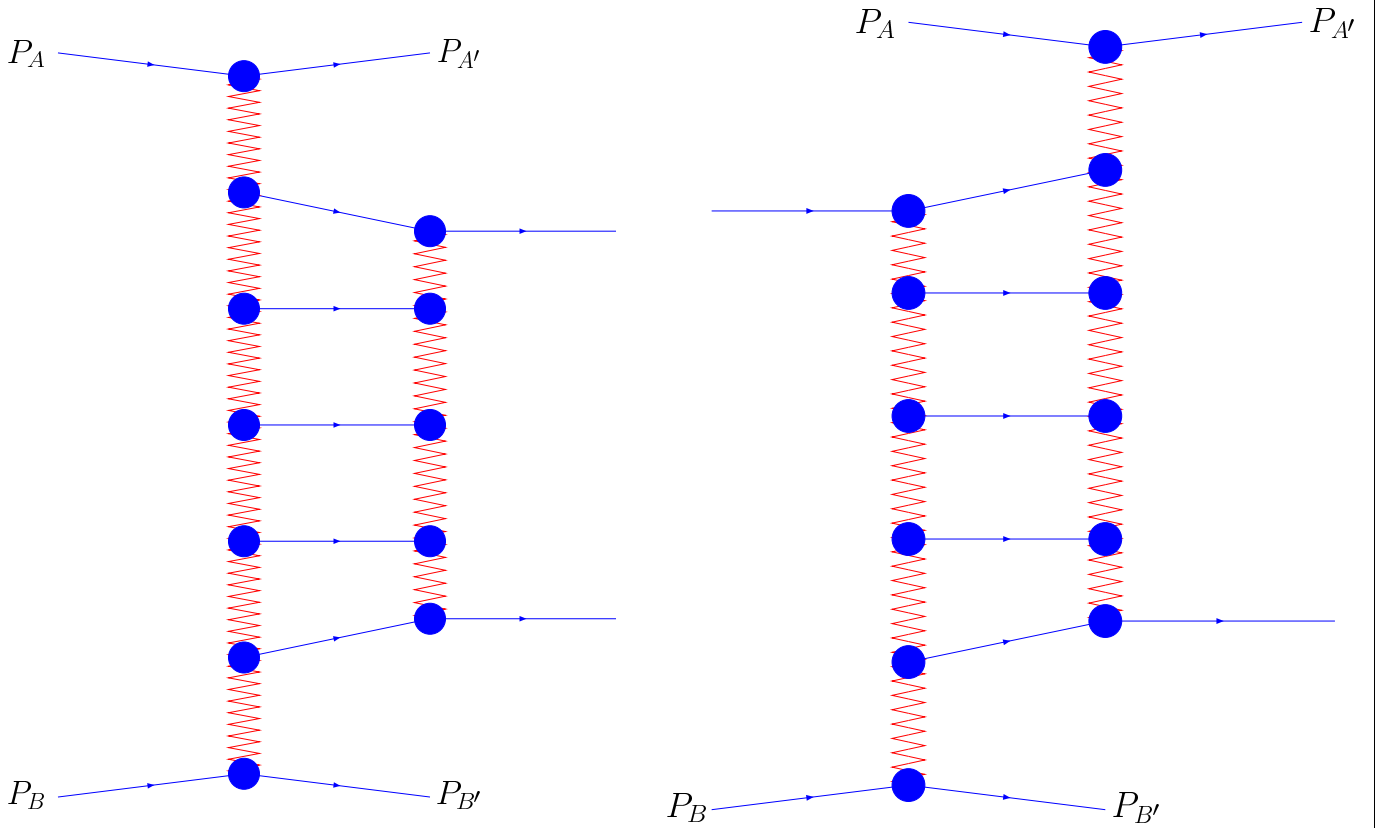


Figure 4: BFKL ladders in $M_{2 \rightarrow 4}$ and $M_{3 \rightarrow 3}$

BFKL equation in octet channels

$$\Delta_2 = -a \left(E - \frac{1}{\epsilon} \right), \quad E\Psi = -a \left(\frac{1}{2} \ln \frac{|p_1 p_2|^2}{\mu^4} + \frac{H}{2} \right) \Psi,$$

$$H = \ln |p_1 p_2|^2 + \frac{1}{p_1 p_2^*} \ln |\rho_{12}|^2 p_1 p_2^* + \frac{1}{p_1^* p_2} \ln |\rho_{12}|^2 p_1^* p_2 + 4\gamma$$

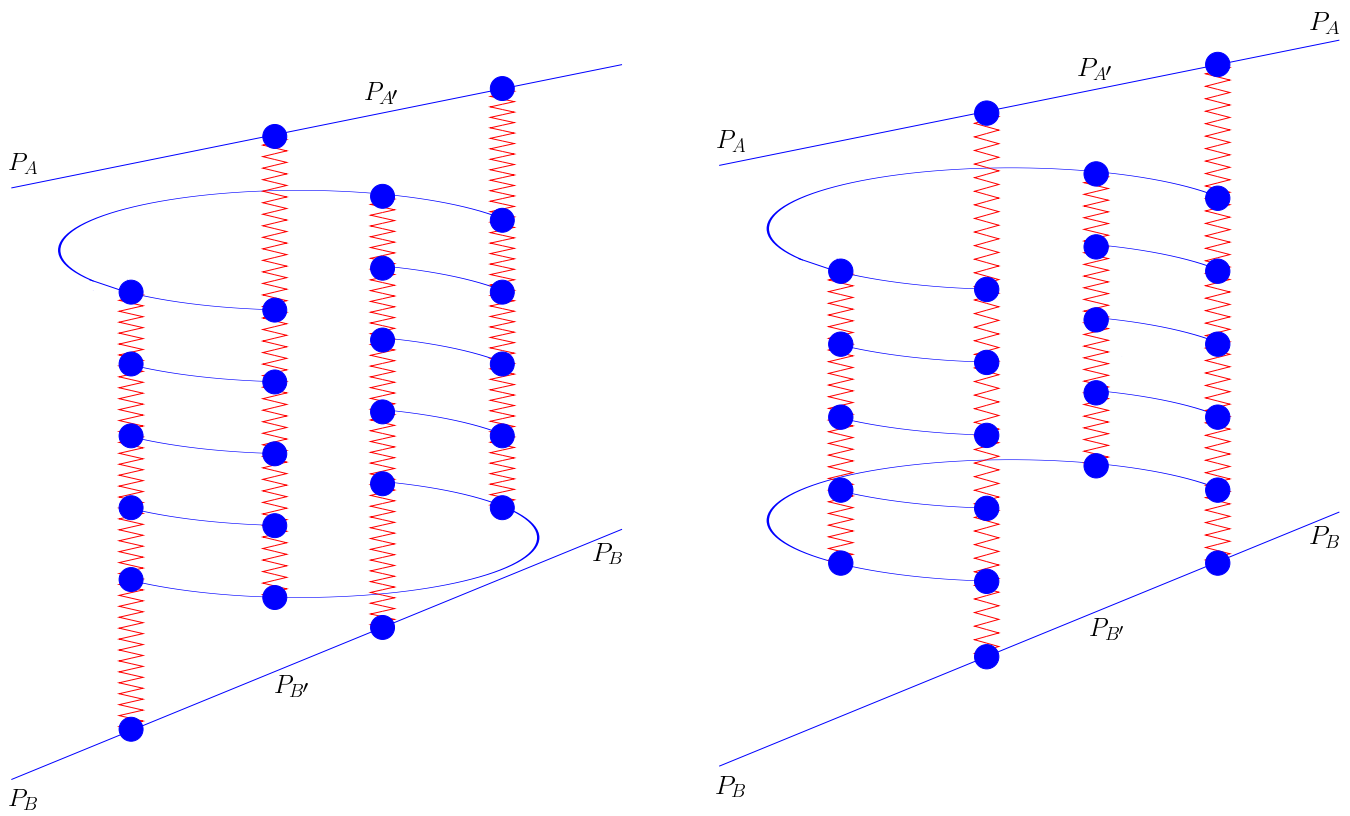


Figure 5: Cylinder-type topologies in the unitarity sums for the total cross section: the intermediate states (discontinuity cuts) are obtained by slicing the cylinders in all possible ways across the intermediate momenta $p_{A'}$ and $p_{B'}$

12 Discussion

1. Pomeron as a composite state of reggeized gluons.
2. Integrability of the BFKL dynamics.
3. Properties of the BFKL kernel in $N = 4$ SUSY.
4. Next-to-next-to-leading corrections.
5. Effective action for Reggeized gluon interactions.
6. BDS amplitudes in the multi-Regge kinematics.
7. Breakdown of the Regge factorization.
8. BFKL cuts in the planar amplitudes M_n for $n > 5$.