Early signals of breakdown through Stochastic modeling

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Dislocation mediated – self organized criticality

Plasticity of Micrometer-Scale Single Crystals in Compression

Michael D. Uchic,¹ Paul A. Shade,² and Dennis M. Dimiduk¹

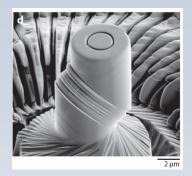
bursts (~20 micron).

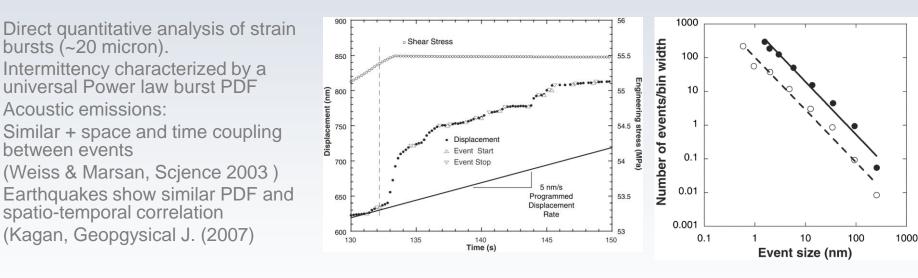
Acoustic emissions:

between events

Single crystal micro-pillar compression:

Dislocation mediated intermittent flow - size effects, hardening. Dislocation density inside a plane as a controlling parameter.

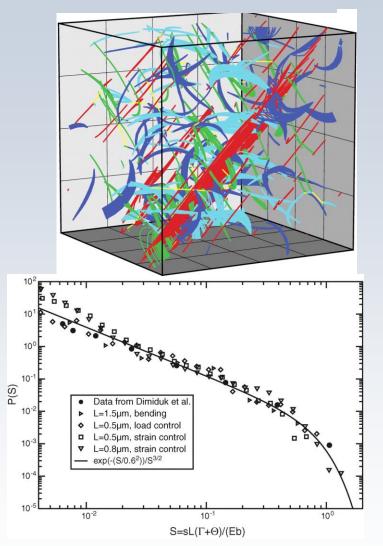




Uchic, Shade & Dimiduk, Annual Review of Materials Research (2009). Dimiduk, Woodward, LeSar & Uchic: "Scale-Free Intermittent Flow in Crystal lasticity." Science (2006) 1188.

Using dislocation dynamics to reproduce PDF

- 3D dislocation dynamics reproduce strain burst scaling $P(s) = Cs^{-t} \exp\left[-(s/s_0)^2\right]$
- where C is a normalization constant, τ is a scaling exponent, and s_0 is the characteristic strain of the largest avalanches.
- Intermittency as a result of dislocation Interactions. Stochastic nature a result of varying initial conditions.
- Avalanche is a 2D event, with an upper cutoff due to structure and work-hardening. Strain is limited to about 10^-6 in a cm size sample.
- Recently (Chen, choi, papanikolaou & Sethna 2010 to 2013): scaling of structures using an advanced CDD code.



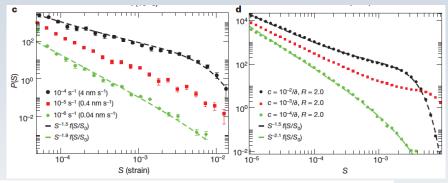
Csikor, Motz, Weygand, Zaiser & Zapperi, "Dislocation Avalanches, Strain Bursts, and the Problem of Plastic Forming at the Micrometer Scale". Science (2007)

Mean field models for critical depining

Quasi-periodic events in crystal plasticity and the self-organized avalanche oscillator

Stefanos Papanikolaou¹, Dennis M. Dimiduk², Woosong Choi³, James P. Sethna³, Michael D. Uchic², Christopher F. Woodward² & Stefano Zapperi^{4,5}

• Reproduce strain rate variation by modifying the mean field picture to include a competing relaxation mechanism. This lead to oscillation in avalanche size. (nature, 2012)



PRL 109, 105702 (2012)

PHYSICAL REVIEW LETTERS

week ending 7 SEPTEMBER 2012

Distribution of Maximum Velocities in Avalanches Near the Depinning Transition

Michael LeBlanc,¹ Luiza Angheluta,^{1,2} Karin Dahmen,¹ and Nigel Goldenfeld¹ • Using a mean field model for interface depining and by solving Fokker-Planck eq. reproduced the power law decay of avalanche size and maximal velocity $\frac{dV}{dt} = -kt$

$$\frac{dV}{dt} = -kV + F_c + \sqrt{V}X(t)$$

What are we trying to do...

- Use stochastic theory to allow for:
 - o transferability of failure scenario analysis (across drive conditions)
 - Identify controlling mechanisms
 - Define critical experiments model development / verification

Such models serve as a link between the microscopic, short time scale problem which is accessible via simulation to the measured system to the real life scenario.

For now – demonstrate the basic method using a "spherical horse" model.

Not trying to do (at this stage):

Create a comprehensive consistent microscopic model Describe the "real" mechanism at work Solve the full field - structure – current response function DC vs RF

Defect model for the dependence of breakdown rate on external electric fields

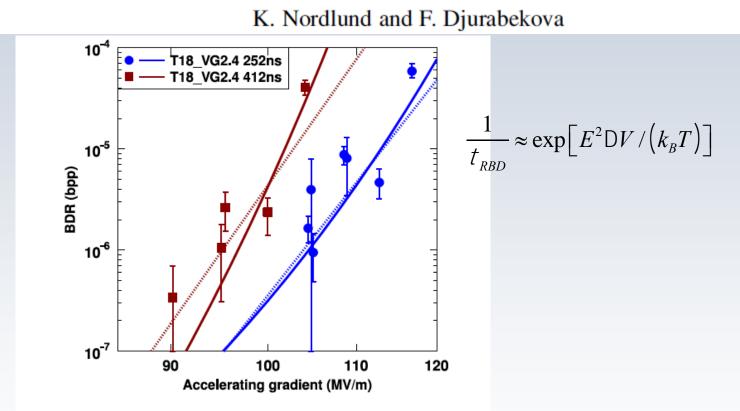


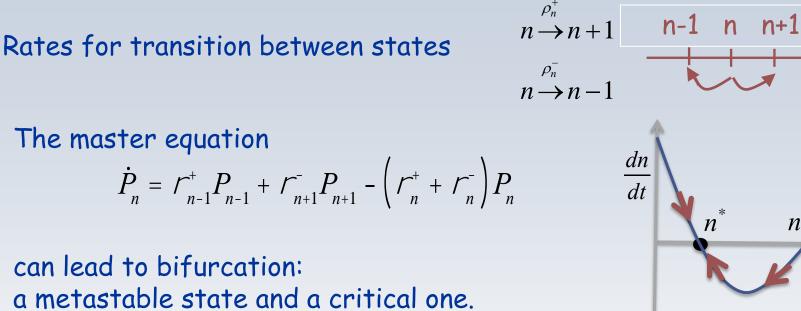
FIG. 4. Measured dependences of $R_{\rm BD}$ (in units of breakdown per pulse, bpp) versus electric field for the T18 accelerating structure [33,43] and fits of our model (solid lines) as well as power laws (dashed lines) to the data.

Formulation of a "well-mixed" 0d model

• Assumptions:

- Breakdown currents are driven by formation of surface extrusion/intrusions.
- Surface protrusions are formed due to multiple dislocation reaction leading to local geometric features
- Sub-breakdown surface protrusion are not identified (true?).
 Breakdown rates do not go up with time (BDR even goes down...).
 Therefor we assume that gradual protrusion accumulation does not control breakdown:
 - surface relaxation, interaction between various slip systems, protrusiondislocation interaction...
 - Field conditions are translated to an applied stress (AC thermal gradients ~ 100 Mpa, dc?)
- Suggested controlling parameter the number of mobile dislocations inside a band.
 - If large amount of dislocations reach the surface in unison an instant extrusion/intrusion may lead to breakdown.
 - We avoid spatial interaction and assume gain-loss dynamics inside a specific band.

General gain-loss type Markovian processes



We look for the quasi-stationary probability distribution function And the probability to cross the critical point (reach extinction)

Approximate solution based on WKB theory with 1/N being the small parameter. $\dot{P} = 0 \quad \triangleright \quad P(n) \circ P(rN) \sim e^{-N[S(r)+O(1/N)]}$

.

"Minimal" model

- Define the "in-plane" density (in units of 1/nm).
- External stress (due to temp gradient on surface), range of 0.1 Gpa.
- Mobile dislocations can increase in number due to stress gradient (the driving force) as well as thermal activation of the multiplication reaction dr^+

$$\frac{dT}{dt} = n_0 GA(S(\Gamma))^k e^{-(f_0 - SW)/k_BT}$$

Moving dislocations can become sessile at:

 Pre-existing barriers (concentration - C)
 "collisions" with other moving dislocations
 dr⁻

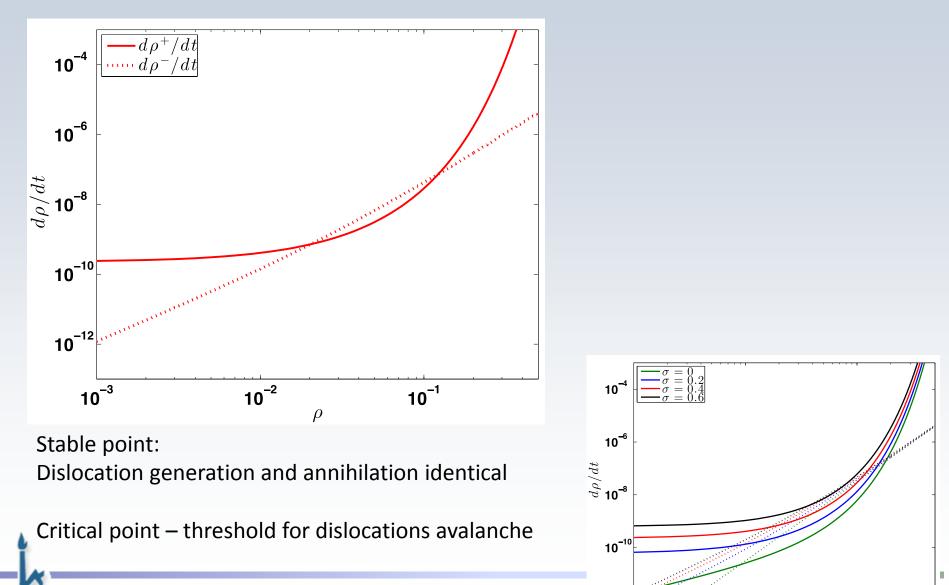
$$\frac{dr^{-}}{dt} = rV(s)c + r2V(s)r$$

• Properties dependence:

 \circ Velocity increase with stress, independent of the number of moving dislocations \circ Stress increase with dislocation content

$$S = S_E + G r b \qquad V = B n_0 b^2 S$$

Model characteristics



10

10

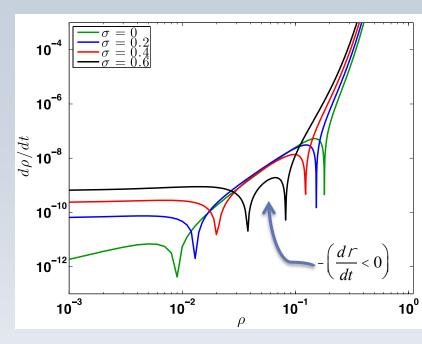
10⁻²

10⁻¹

Low stresses:

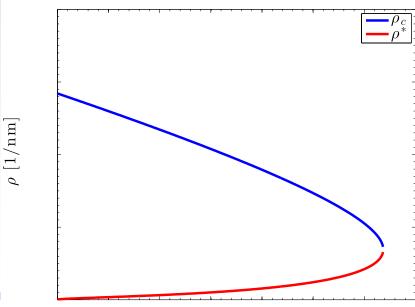
Mobile dislocation density remains in Metastable solution.

Dynamic barrier decreases with increasing stress.



Up to a critical stress – bifurcation to two solutions.

Above it - no stable solution.



Define :
$$s(\rho) = -\int_{\rho} \ln \frac{\rho^+(x)}{\rho^-(x)} dx$$

For (k=2):
$$s(\Gamma) = -\Gamma \ln \left[\frac{\Gamma^+(\Gamma)}{\Gamma^-(\Gamma)} \right] - \frac{S_E}{bm} \ln \left[S(\Gamma) \right] + \frac{c}{2} \ln(c+2\Gamma) - \Gamma \left(1 - \frac{\Gamma bmW}{2k_B T} \right)$$

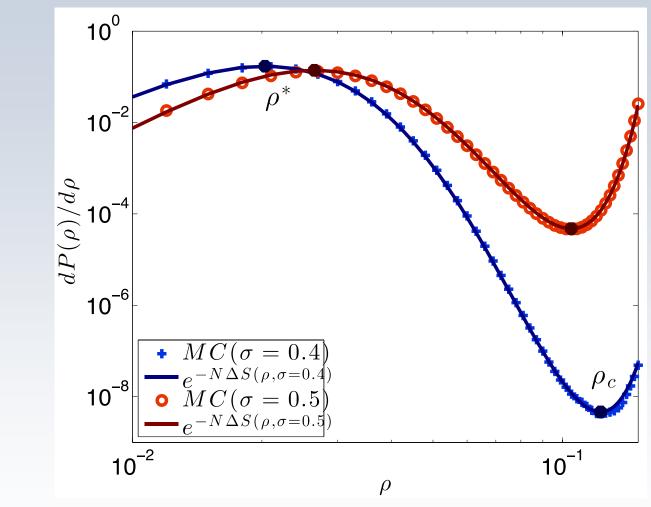
Using*:
$$P(n = N \Gamma) = P(\Gamma) \sim e^{-Ns(\Gamma)}$$

Leads to
$$P(\Gamma) \propto \left[\frac{\Gamma^+(\Gamma)}{\Gamma^-(\Gamma)}\right]^{N\Gamma} \frac{\left[S(\Gamma)\right]^{NS_E/bm}}{(c+2\Gamma)^{Nc/2}} e^{N\Gamma\left(1-\frac{rbmW}{2k_BT}\right)}$$

And the normalized PDF
$$P(r) = \sqrt{\frac{S''(r^*)}{2\rho N}} e^{-N[s(r)-s(r^*)]}$$

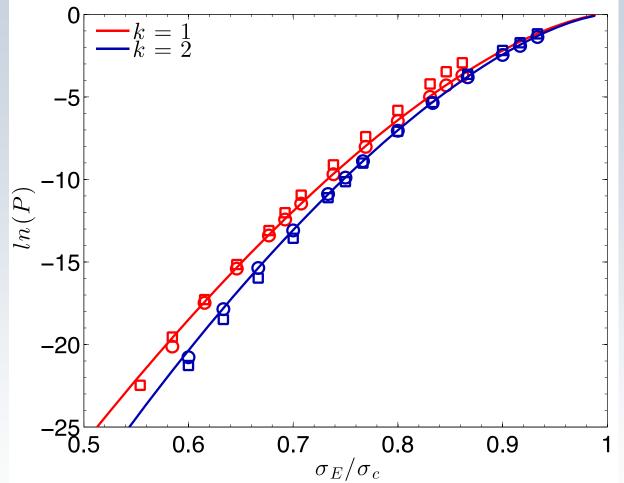
*Assaf and Meerson, Phys Rev. E 21116 (2010) •

Analytical vs MC



• Analytical analysis reproduces full PDF.

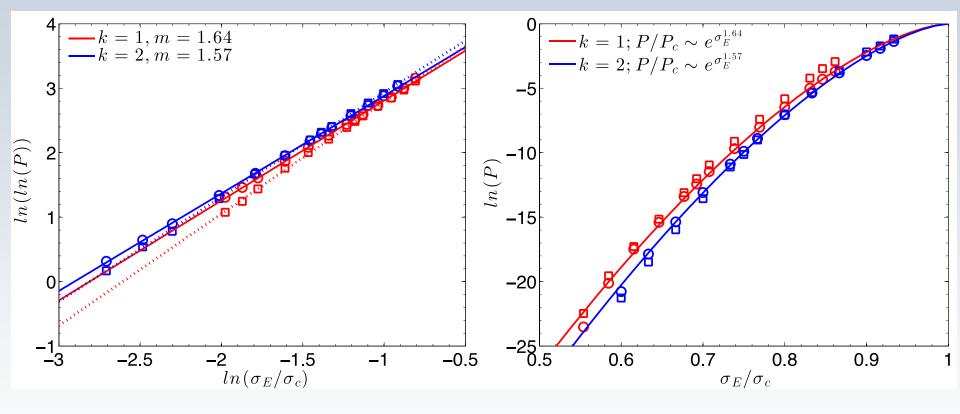
Getting BDR curves



Analytical solution for relative probability to reach critical point.

• Normalized probability and rate for reaching the critical state.

"Universal" BDR exponent

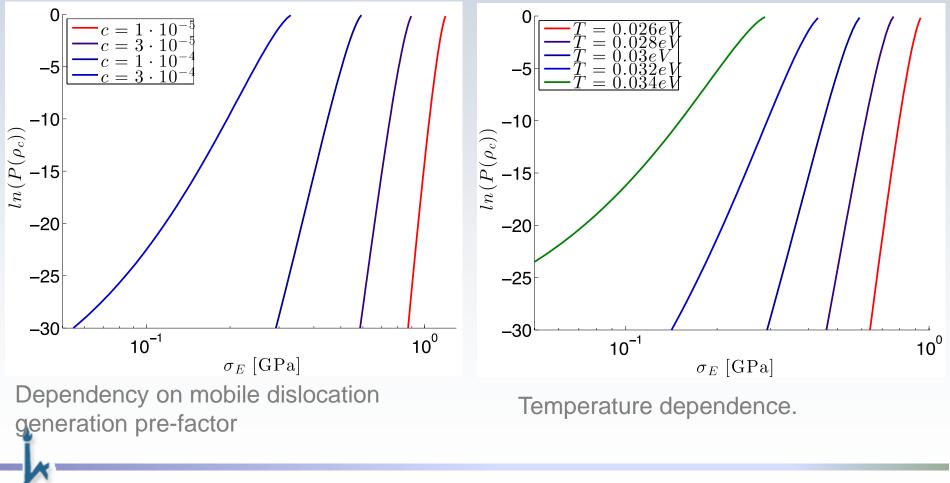


• fit to experimental results

 $P \sim \frac{1}{\tau} \sim \exp(\sigma_E^{1.6})$

Other dependencies

Employ analytical solution to various scenarios

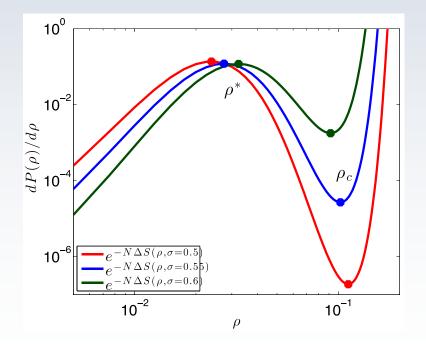


Signs of criticality

- Adiabatically moving between quasistationary PDF: Change in pdf moments with field

 identify threshold
- At specific conditions, probe time dependencies of the QS pdf: Identify large fluctuations time dependency
 - -> identify time constants
 - -> mechanism

$$P_{c}(S, \Gamma, t) = \check{\mathfrak{g}}_{0}^{t} P(S, \Gamma > \Gamma, t) dt$$



REVIEWS

PRE-breakdown

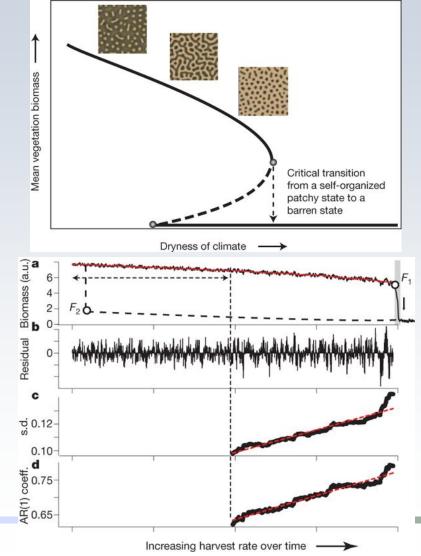
- As the system approaches the critical point.
 Fluctuation diverge.
- Observable through standard deviation of the time correlation

$$SD(t) = \frac{\int_{t-D}^{t+D} (I(t) - \langle I \rangle)^2 dt}{(\langle I \rangle)^2}$$

• Or, more generally, autocorrelation in the signal $R(k) = \frac{\hat{0}_{0}^{t-k} (I(t) - \langle I \rangle) (I(t+k) - \langle I \rangle) dt}{\hat{0}_{0}^{t-k} (I(t) - \langle I \rangle)^{2} dt}$

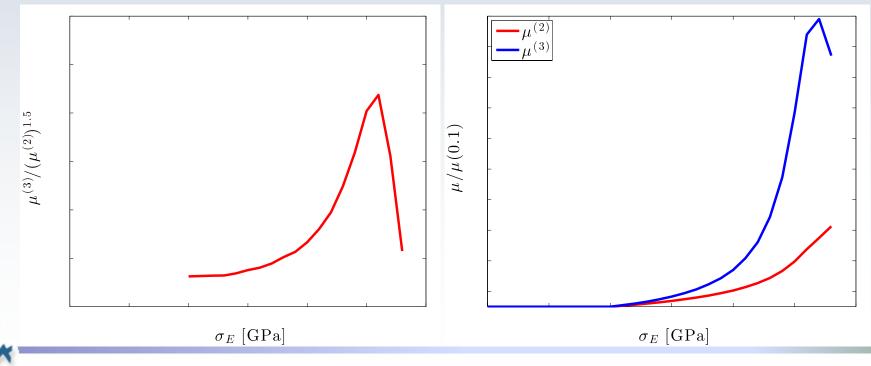
Early-warning signals for critical transitions

Marten Scheffer¹, Jordi Bascompte², William A. Brock³, Victor Brovkin⁵, Stephen R. Carpenter⁴, Vasilis Dakos¹, Hermann Held⁶, Egbert H. van Nes¹, Max Rietkerk⁷ & George Sugihara⁸



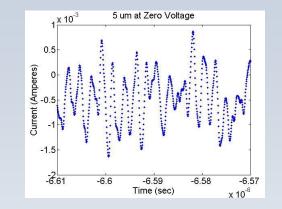
Early warning signals

- Using moments $m^{(n)} = a^{(r-\langle r \rangle)^n}$
- Variance and third moment indicate to Strong variation of quasi-stationary PDF - in breakdown states

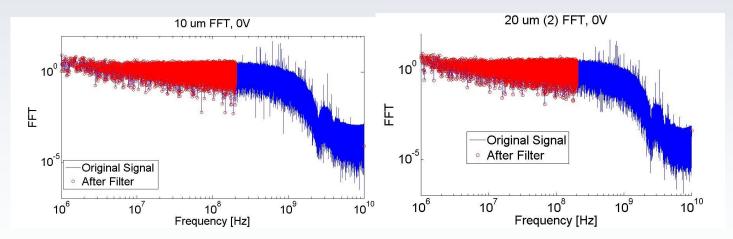


Initial try – DC measurements

 Dark current measurements (Varying field and gap distance)



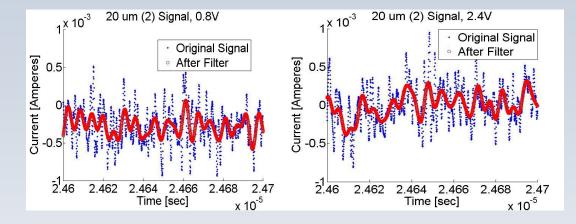
 Low pass filter is clearly needed (applied 0.2 GhZ)



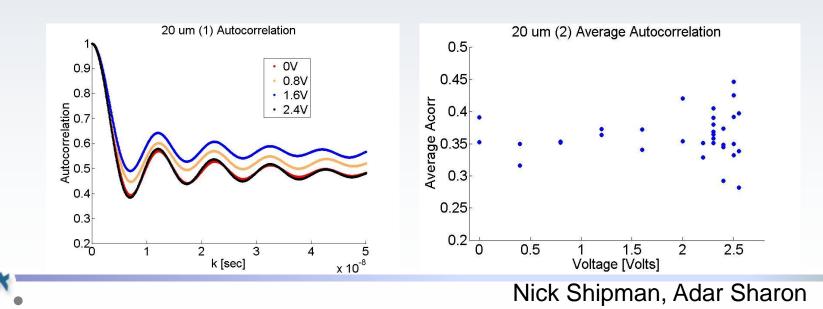
Nick Shipman, Adar Sharon

Initial try – DC measurmanets

Filtered signal

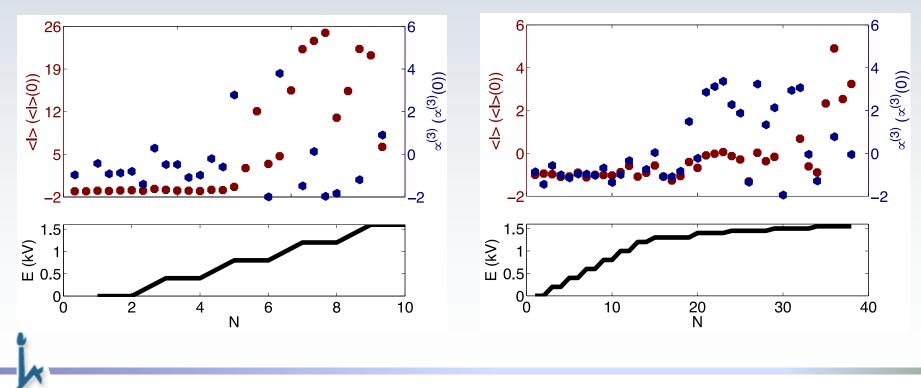


• Sadly ... autocorrelation shows nothing...



Normalized average and third moment

- Mean and third moments variatons with field.
- Need for better measurements ...



Summary

- Intermittency due to collective dislocation response is well established.
 - Experimental scenarios: acoustic emission, micro-compression.
 - Universal behavior, earth quakes, other non local bifurcating systems.
- Proposed a simple stochastic model to describe breakdown
 phenomena
 - Using a minimal model MANY simplifying assumption demonstrate critical behavior, bifurcation and reproduce observed BDR (E).
 - Analytically (or at least numerically) solvable
 - reproduce observed BCR exponent. Universality?
- Unique experimental scenarios:
 - PDF pre breakdown analysis: using PDF and PDF tail.
 - $\circ \quad \text{Pre breakdown fluctuation.}$
- Future directions
 - Effects of intrinsic noise?
 - o Join forces with experimental work to identify signs of criticality.
 - Suggest a modified conditioning scheme .
- Can serve to bridge microscopic mechanisms to experimental scenarios. need for an explicit response function.
- New opportunity for stochastic analysis...