

Introduction to Accelerator Physics

Bernhard Holzer,
CERN

A Real Introduction ...



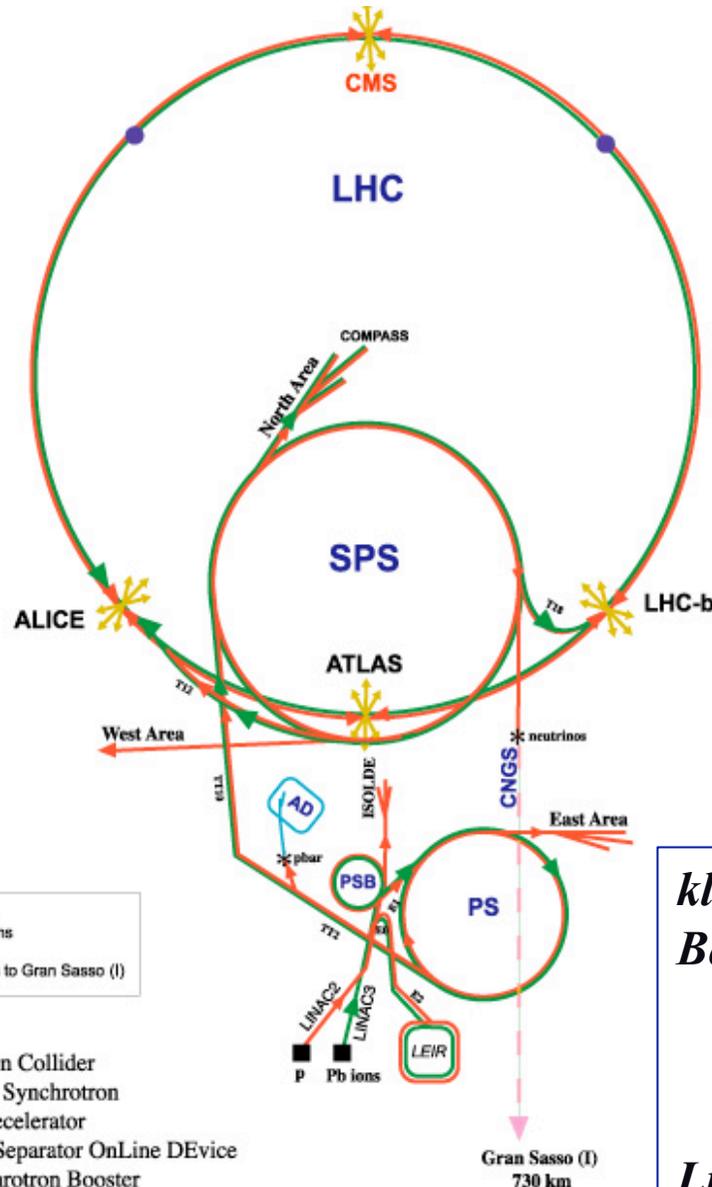
Teilchen Energien & Beschleuniger

Linac 2: 50 MeV
Booster: 1.4 GeV
PS: 26 GeV
SPS: 450 GeV

$$E = mc^2$$

Ruhemassen:

Proton: 938 MeV
Electron: 511 keV
Myon: 105 MeV



— protons
— antiprotons
— ions
— neutrinos to Gran Sasso (I)

LHC: Large Hadron Collider
 SPS: Super Proton Synchrotron
 AD: Antiproton Decelerator
 ISOLDE: Isotope Separator OnLine DEvice
 PSB: Proton Synchrotron Booster
 PS: Proton Synchrotron
 LINAC: LINEar ACcelerator
 LEIR: Low Energy Ion Ring
 CNGS: Cern Neutrinos to Gran Sasso

klassischer / relativistischer Bereich

$$\gamma = \frac{E}{mc^2} = \sqrt{\frac{1}{1 - (v/c)^2}}$$

Linac 2:

$$\gamma = 988 / 938$$

$$\beta = v/c = 0.3$$

Rudolf LEY, PS D
Revised and adapted
in collaboration with
D. Manglunki, PS

Example: Kernphysik & Isotope : ISOLDE

Find the produced isotopes from a given target

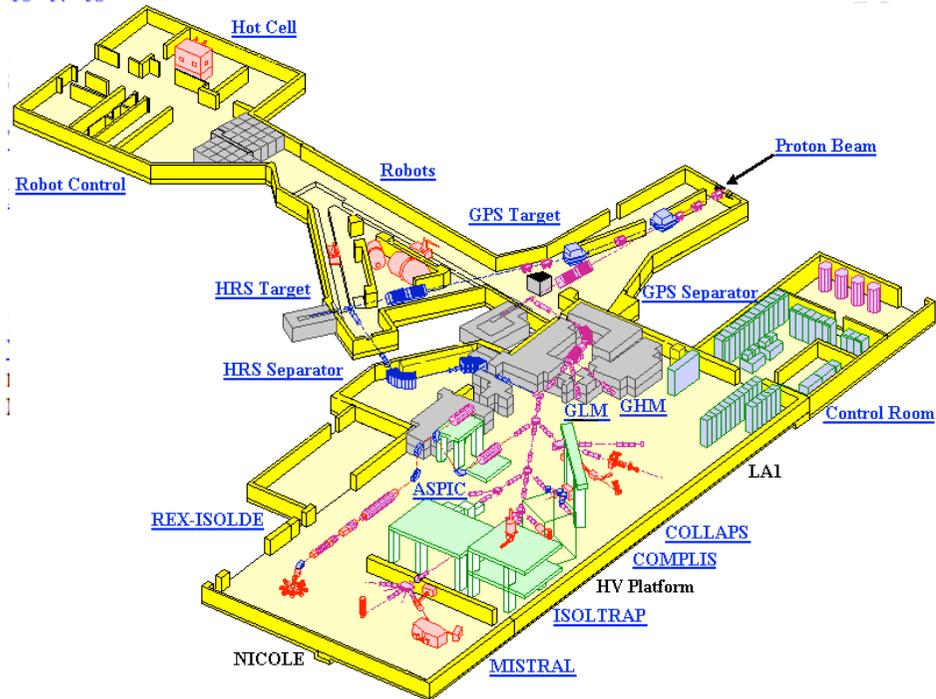
[Nuclear Chart for ISOLDE](#)

Find the produced isotope from an element independent on target

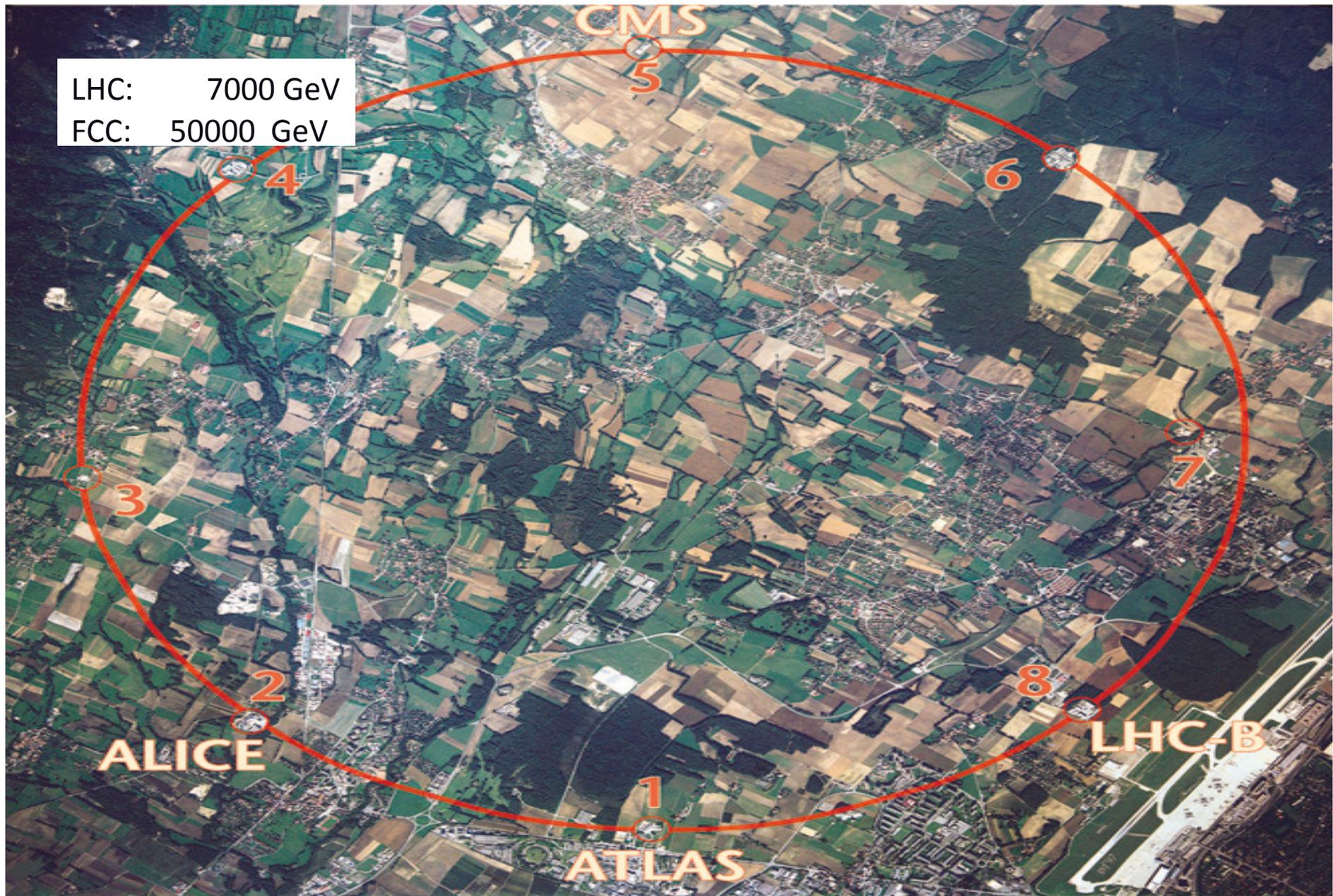
Group	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18											
	1A	2A	3B	4B	5B	6B	7B		8B		1B	2B	3A	4A	5A	6A	7A	8A											
Period	<p style="text-align: center;">Ion source: + Surface - hot Plasma cool Laser</p>																												
1	1 H																	2 He											
2	3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne											
3	11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar											
4	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr											
5	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe											
6	55 Cs	56 Ba	57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi
7	87 Fr	88 Ra	**	103 Lr	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg																	
* Lanthanides *			57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm														
** Actinides **			89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md														

68, 82m, 70, 71, 72, 74, 80, 69, 85, 73, 78, 81, 82 As

Element	A number	Half life	SC or PSB*	Yield at ISOLDE (ions/ μ C)	Target material
As	69	15.2 m	2 PSB	8.0E+05	ZrO ₂
As	70	52.6 m	3 PSB	6.0E+06	ZrO ₂
As	71	65.28 h	15 PSB	7.0E+07	ZrO ₂
As	72	26.0 h	1 PSB	3.0E+08	ZrO ₂
As	73	80.30 d	6 PSB	1.0E+09	ZrO ₂



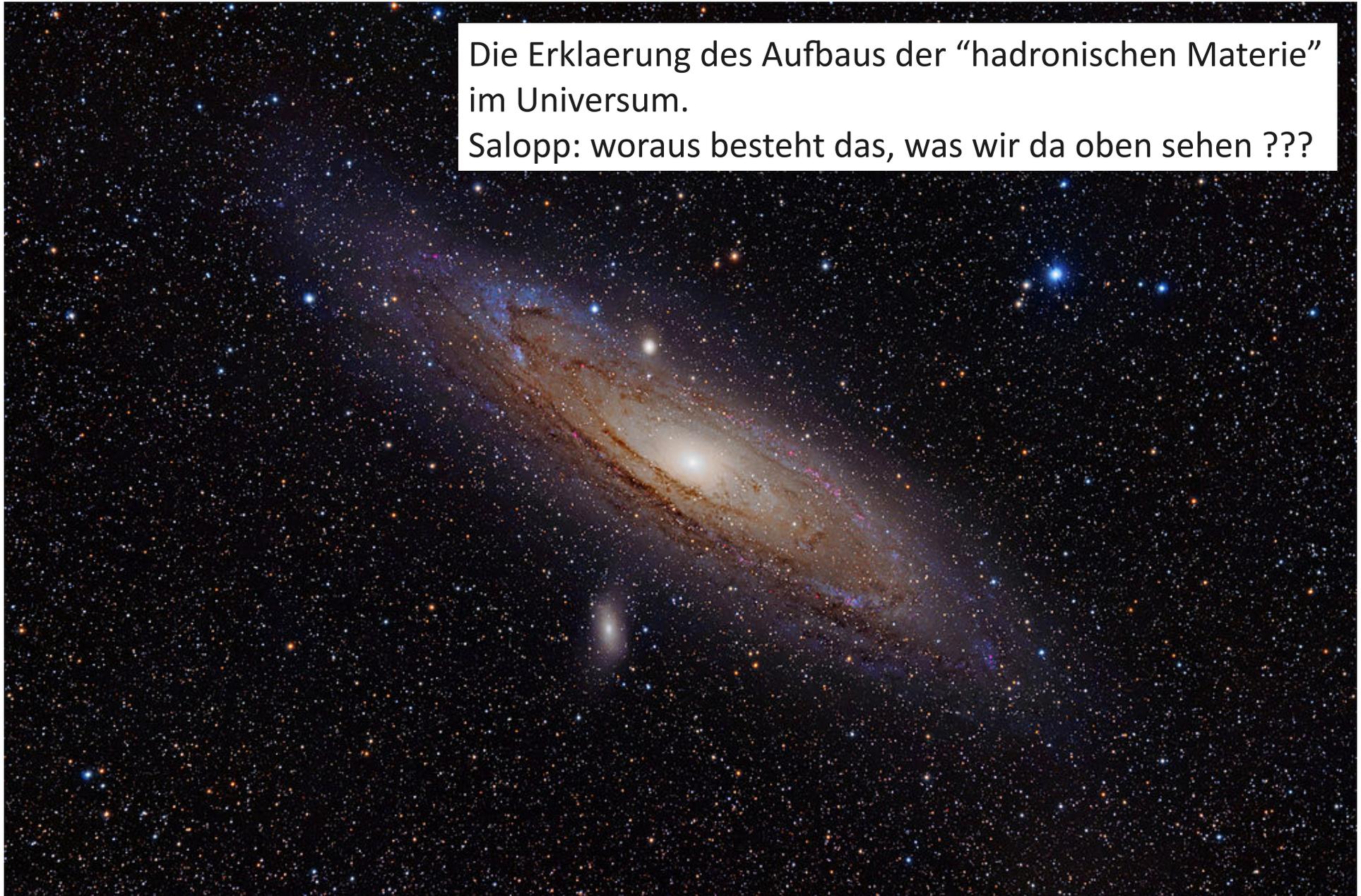
Und schlussendlich die Teilchen Physik



Standard Model: Worum gehts eigentlich ??

Die Erklärung des Aufbaus der “hadronischen Materie”
im Universum.

Salopp: woraus besteht das, was wir da oben sehen ???



Standard Model: Worum gehts eigentlich ??

um den Versuch Ordnung & Systematik zu erkennen

<p>GRUPPE</p> <p>PERIODEN</p> <p>1 1.007 H WASSERSTOFF</p> <p>2 6.94 Li LITHIUM</p> <p>3 22.99 Na NATRIUM</p> <p>4 39.09 K KALIUM</p> <p>5 85.46 Rb RUBIDIUM</p> <p>6 132.9 Cs CÄSIUM</p> <p>7 223 Fr FRANCIUM</p> <p>(1) Pure Appl. Chem. Die relative Atommasse ist angegeben. Für Elemente in Klammern ist die Atommasse in Klammern angegeben. Drei dieser Elemente sind in der Erdkruste in bedeutender Menge vorhanden.</p> <p>Redakteur: Marc</p>	<h3>Leptonen</h3> <table border="1"> <tr> <th>e-Neutrino</th> <th>μ-Neutrino</th> <th>τ-Neutrino</th> </tr> <tr> <td></td> <td></td> <td></td> </tr> <tr> <td>Elektron</td> <td>Myon</td> <td>Tauon</td> </tr> <tr> <td></td> <td></td> <td></td> </tr> </table>			e-Neutrino	μ -Neutrino	τ -Neutrino				Elektron	Myon	Tauon				<h3>Bosonen</h3> <table border="1"> <tr> <th>Photon</th> <th>Z^0</th> </tr> <tr> <td></td> <td></td> </tr> <tr> <td>W^+</td> <td>W^-</td> </tr> <tr> <td></td> <td></td> </tr> </table>		Photon	Z^0			W^+	W^-		
	e-Neutrino	μ -Neutrino	τ -Neutrino																						
	Elektron	Myon	Tauon																						
	Photon	Z^0																							
	W^+	W^-																							
	<h3>Quarks</h3> <table border="1"> <tr> <th>up</th> <th>charm</th> <th>top</th> </tr> <tr> <td></td> <td></td> <td></td> </tr> <tr> <th>down</th> <th>strange</th> <th>bottom</th> </tr> <tr> <td></td> <td></td> <td></td> </tr> </table>					up	charm	top				down	strange	bottom				<h3>Gluonen</h3>							
up	charm	top																							
down	strange	bottom																							

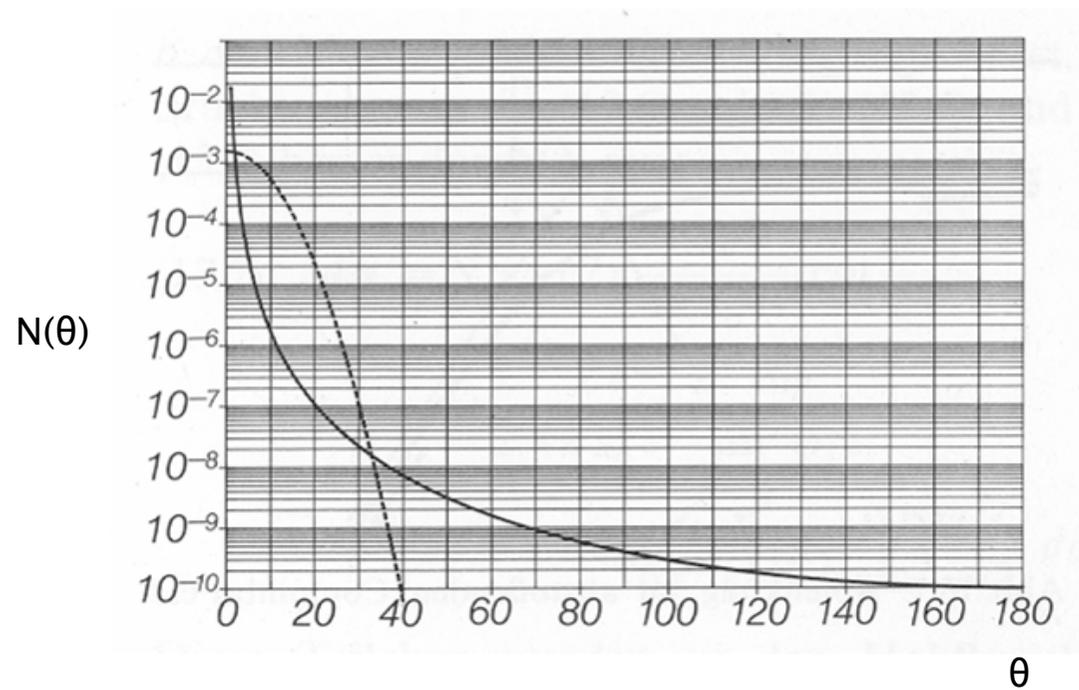
I.) A Bit of History



$$N(\theta) = \frac{N_i n t Z^2 e^4}{(8\pi\epsilon_0)^2 r^2 K^2} * \frac{1}{\sin^4(\theta/2)}$$

Rutherford Scattering, 1911

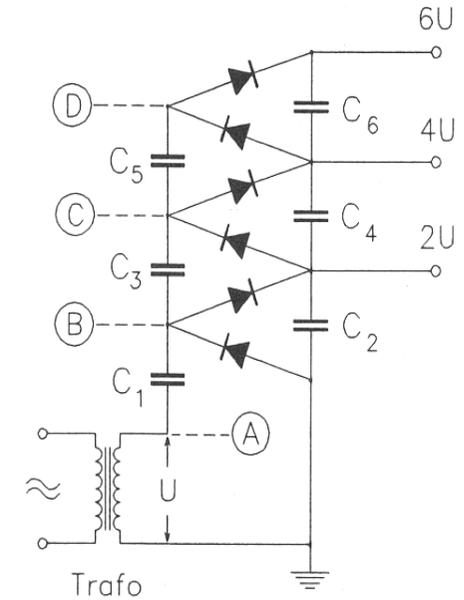
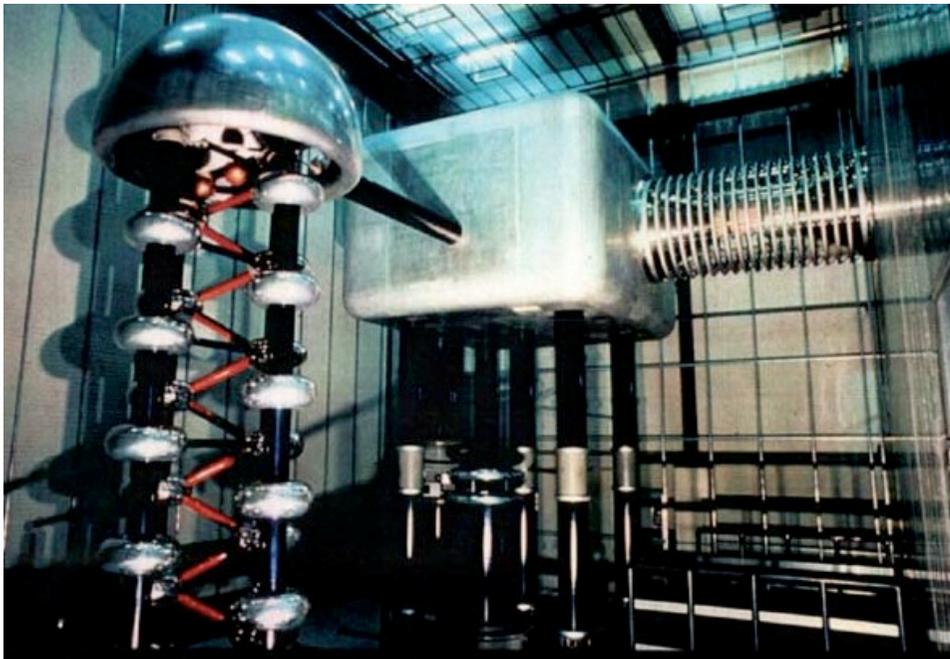
Using radioactive particle sources:
 α -particles of some MeV energy



1.) Electrostatic Machines: The Cockcroft-Walton Generator

1928: Encouraged by Rutherford Cockcroft and Walton start the design & construction of a high voltage generator to accelerate a proton beam

1932: First particle beam (protons) produced for nuclear reactions: splitting of Li-nuclei with a proton beam of 400 keV



Particle source: Hydrogen discharge tube

on 400 kV level

Accelerator: evacuated glass tube

Target: Li-Foil on earth potential
Technically: rectifier circuit, built of capacitors and diodes (Greinacher)

Problem:

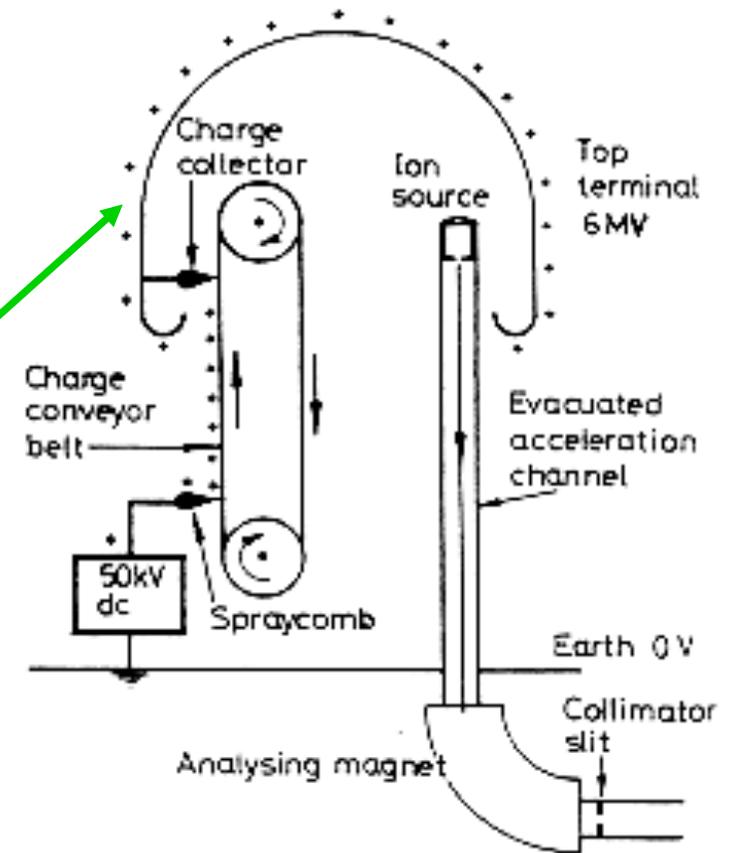
DC Voltage can only be used once

2.) Electrostatic Machines: (Tandem -) van de Graaff Accelerator (1930 ...)

creating high voltages by **mechanical**
transport of charges

- * Terminal Potential: $U \approx 12 \dots 28 \text{ MV}$
using high pressure gas to suppress discharge (SF_6)

- Problems:**
- * Particle energy limited by high voltage discharges
 - * high voltage **can only be applied once per particle ...**
... or twice ?



The „Tandem principle“: Apply the accelerating voltage twice ...
... by working with **negative ions (e.g. H^-)** and
stripping the electrons in the centre of the
structure

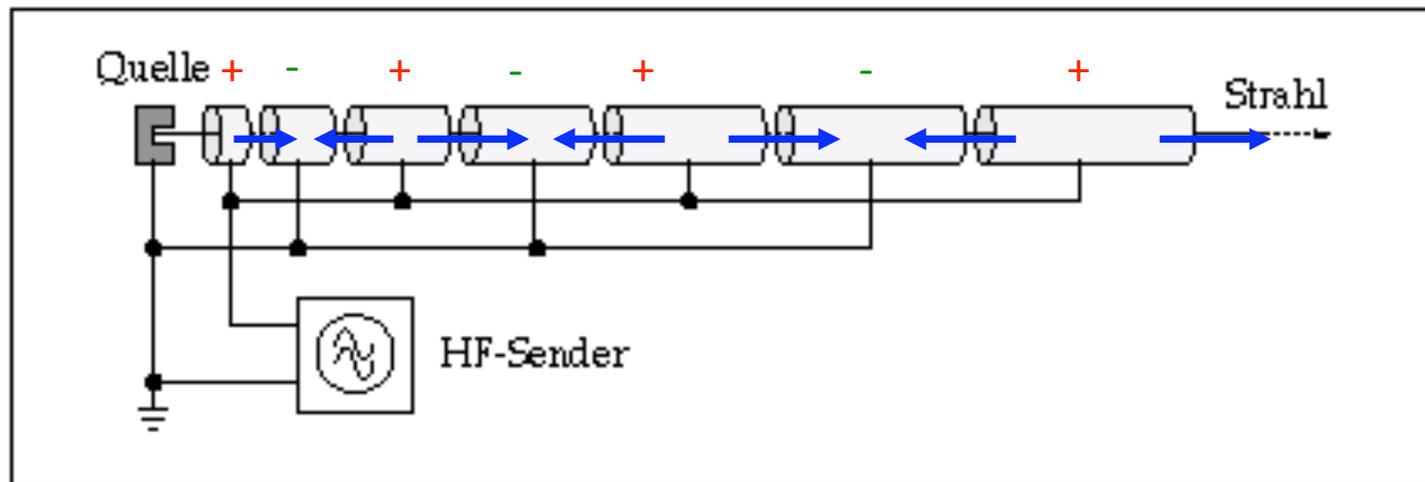
Example for such a „steam engine“: 12 MV-Tandem van de Graaff
Accelerator at MPI Heidelberg



3.) The first RF-Accelerator: „Linac“

1928, Wideroe: how can the acceleration voltage be applied several times to the particle beam

schematic Layout:



Energy gained after n acceleration gaps

$$E_n = n * q * U_0 * \sin \psi_s$$

n number of gaps between the drift tubes

q charge of the particle

U_0 Peak voltage of the RF System

ψ_s synchronous phase of the particle

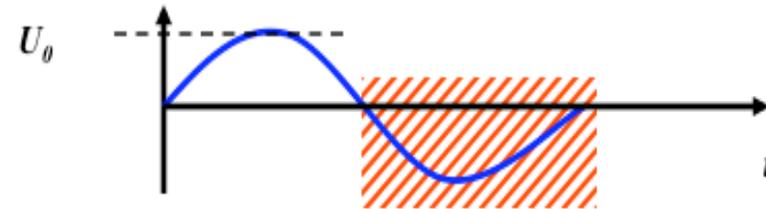
* acceleration of the proton in the first gap

* voltage has to be „flipped“ to get the right sign in the second gap → RF voltage

→ shield the particle in drift tubes during the negative half wave of the RF voltage

Wideroe-Structure: the drift tubes

shielding of the particles during the negative half wave of the RF



Time span of the negative half wave: $\tau_{RF}/2$

Length of the Drift Tube: $l_i = v_i * \frac{\tau_{rf}}{2}$

Kinetic Energy of the Particles

$$E_i = \frac{1}{2}mv^2$$



$$\rightarrow v_i = \sqrt{2E_i/m}$$

$$l_i = \frac{1}{v_{rf}} * \sqrt{\frac{i * q * U_{0 * \sin \psi_s}}{2m}}$$

valid for **non relativistic** particles ...

Alvarez-Structure: 1946, surround the whole structure by a rf vessel

Energy: ≈ 20 MeV per Nucleon $\beta \approx 0.04 \dots 0.6$, Particles: Protons/Ions

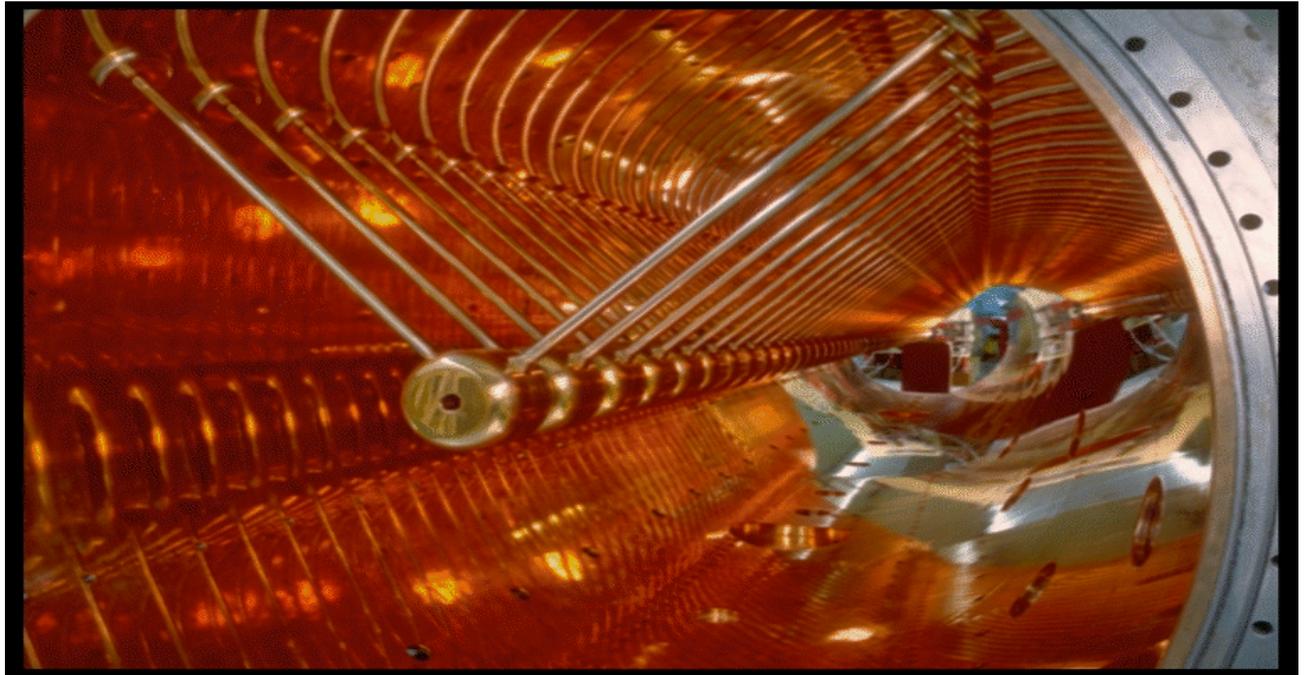
Accelerating structure of a Proton Linac (DESY Linac III)

$$E_{total} = 988 \text{ M eV}$$

$$m_0 c^2 = 938 \text{ M eV}$$

$$p = 310 \text{ M eV} / c$$

$$E_{kin} = 50 \text{ M eV}$$



Beam energies

1.) reminder of some relativistic formula

rest energy $E_0 = m_0 c^2$

total energy $E = \gamma * E_0 = \gamma * m_0 c^2$

kinetic energy $E_{kin} = E_{total} - m_0 c^2$

Energy Gain per „Gap“:

$$W = q U_0 \sin \omega_{RF} t$$

momentum $E^2 = c^2 p^2 + m_0^2 c^4$

4.) The Cyclotron: (Livingston / Lawrence ~1930)

Idea: Bend a Linac on a Spiral
Application of a constant magnetic field
keep $B = \text{const}$, $RF = \text{const}$

→ Lorentzforce

$$\vec{F} = q * (\vec{v} \times \vec{B}) = q * v * B$$

circular orbit

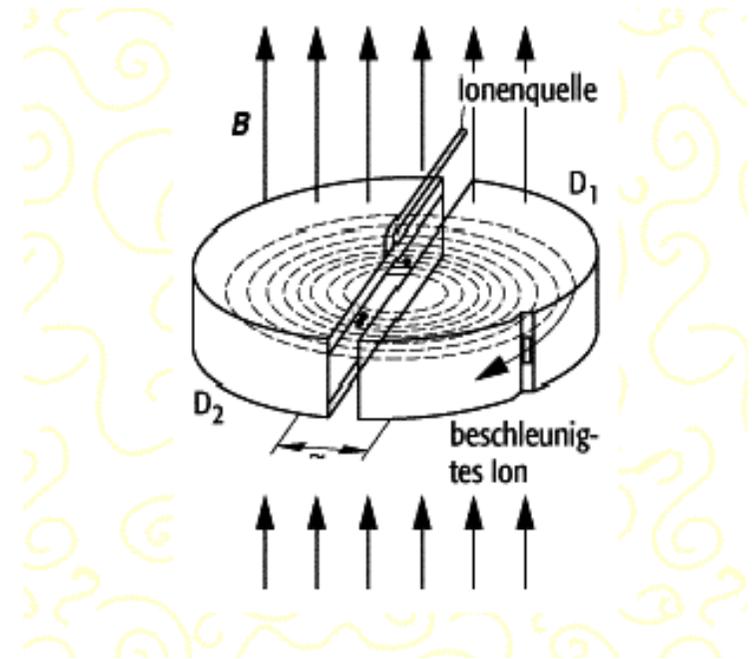
$$q * v * B = \frac{m * v^2}{R} \quad \rightarrow \quad B * R = p / q$$

increasing radius for
increasing momentum
→ Spiral Trajectory

revolution frequency

$$\omega_z = \frac{q}{m} * B_z$$

the cyclotron (rf-) frequency
is independent of the momentum

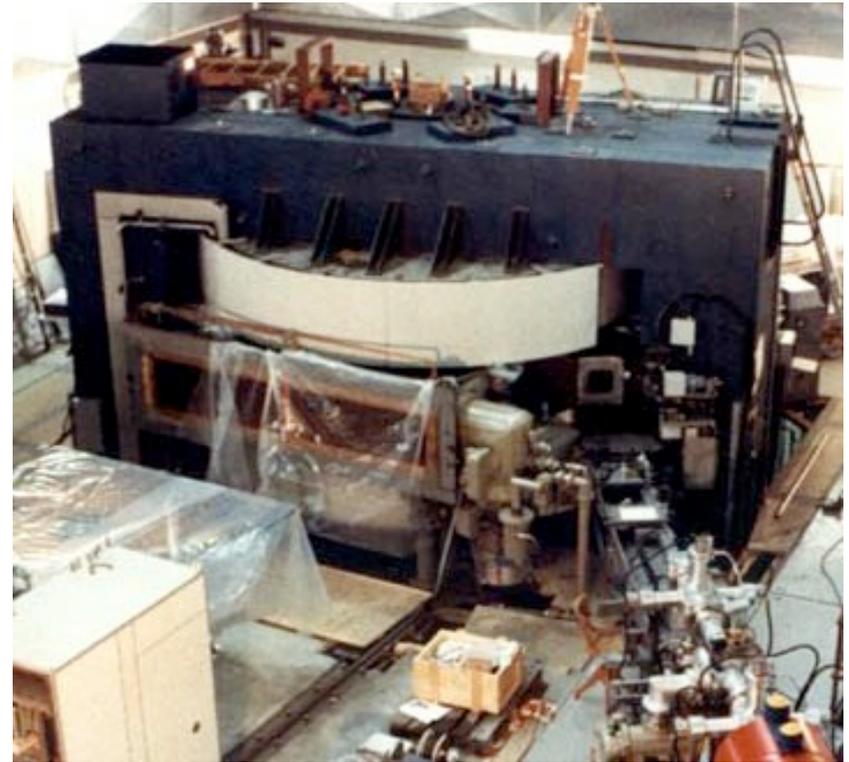


Cyclotron:

! ω is constant for a given q & B

!! $B \cdot R = p/q$
large momentum \rightarrow huge magnet

!!!! $\omega \sim 1/m \neq \text{const}$ works properly only for
non relativistic particles



PSI Zurich

Application:

Work horses for medium energy protons

Proton / Ion Acceleration up to ≈ 60 MeV (proton energy)

nuclear physics

radio isotope production, proton / ion therapy

1.) Introduction and Basic Ideas

„ ... in the end and after all it should be a kind of circular machine“

→ need transverse deflecting force

Lorentz force $\vec{F} = q * (\cancel{\vec{E}} + \vec{v} \times \vec{B})$

typical velocity in high energy machines: $v \approx c \approx 3 * 10^8 \text{ m/s}$

Example:

$$B = 1 \text{ T} \quad \rightarrow \quad F = q * 3 * 10^8 \frac{\text{m}}{\text{s}} * 1 \frac{\text{Vs}}{\text{m}^2}$$

$$F = q * \underbrace{300 \frac{\text{MV}}{\text{m}}}$$

equivalent E
electrical field:

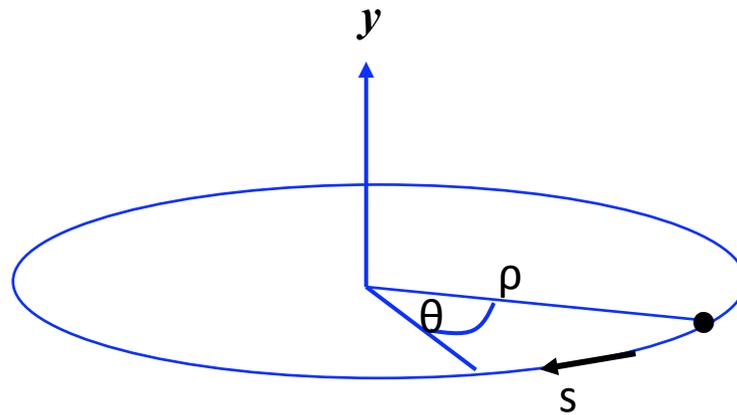
technical limit for el. field:♪

$$E \leq 1 \frac{\text{MV}}{\text{m}}$$

old greek dictum of wisdom:

if you are clever, you use magnetic fields in an accelerator wherever it is possible.

The ideal circular orbit



circular coordinate system

condition for circular orbit:

Lorentz force

$$F_L = e v B$$

centrifugal force

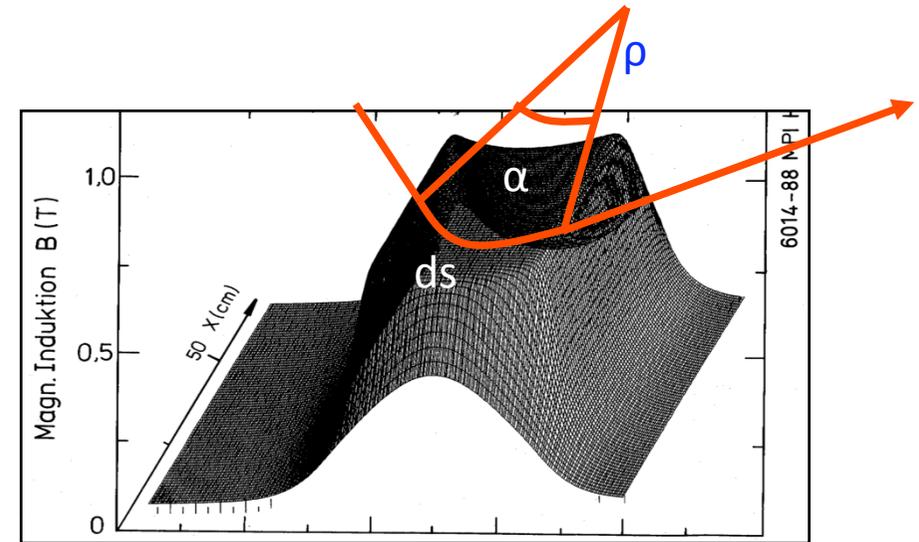
$$F_{centr} = \frac{\gamma m_0 v^2}{\rho}$$

$$\frac{\gamma m_0 v^2}{\rho} = e v B$$

$$\frac{p}{e} = B \rho$$

$B \rho =$ "beam rigidity"

The Magnetic Guide Field



field map of a storage ring dipole magnet

$$\rho = 2.53 \text{ km} \quad \longrightarrow \quad 2\pi\rho = 17.6 \text{ km} \approx 66\%$$

$$B \approx 1 \dots 8 \text{ T}$$

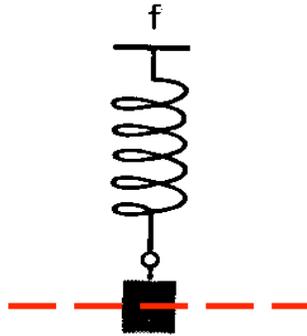
rule of thumb:

$$\frac{1}{\rho} \approx 0.3 \frac{B [T]}{p [GeV/c]}$$

„normalised bending strength“

2.) Focusing Properties - Kurzer Ausflug in die klassische Mechanik

classical mechanics:
pendulum



there is a **restoring force**, proportional to the elongation x :

$$m^* \frac{d^2 x}{dt^2} = -c^* x$$

general solution: free harmonic oscillation

$$x(t) = A^* \cos(\omega t + \varphi)$$

Storage Ring: we need a **Lorentz force** that rises as a function of the **distance to** ?

..... the design orbit

$$F(x) = q^* v^* B(x)$$

Quadrupole Magnets:

required: **focusing forces** to keep trajectories in vicinity of the ideal orbit

linear increasing Lorentz force

linear increasing magnetic field

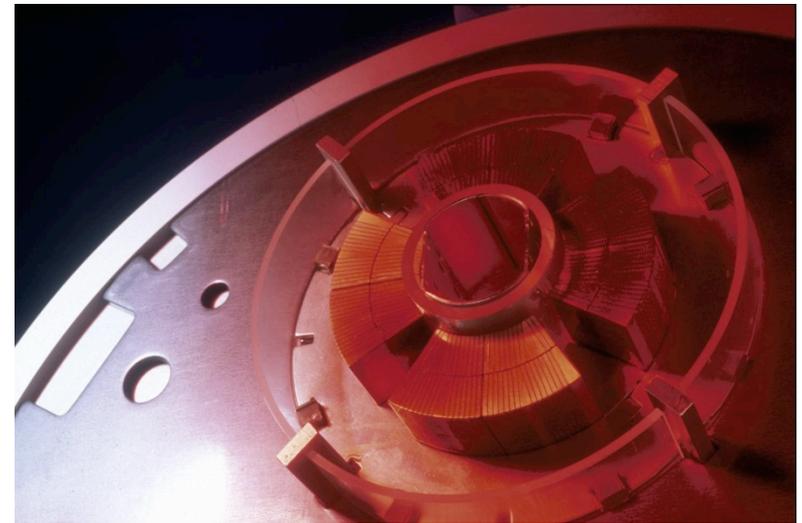
$$B_y = g x \quad B_x = g y$$

normalised quadrupole field:

→ $k = \frac{g}{p/e}$

simple rule:

$$k = 0.3 \frac{g(T/m)}{p(GeV/c)}$$



LHC main quadrupole magnet

$$g \approx 25 \dots 220 \text{ T/m}$$

what about the vertical plane:
... Maxwell

$$\vec{\nabla} \times \vec{B} = \cancel{\vec{j}} + \frac{\partial \cancel{\vec{E}}}{\partial t} = 0$$

$$\Rightarrow \frac{\partial B_y}{\partial x} = \frac{\partial B_x}{\partial y} = g$$

Focusing forces and particle trajectories:

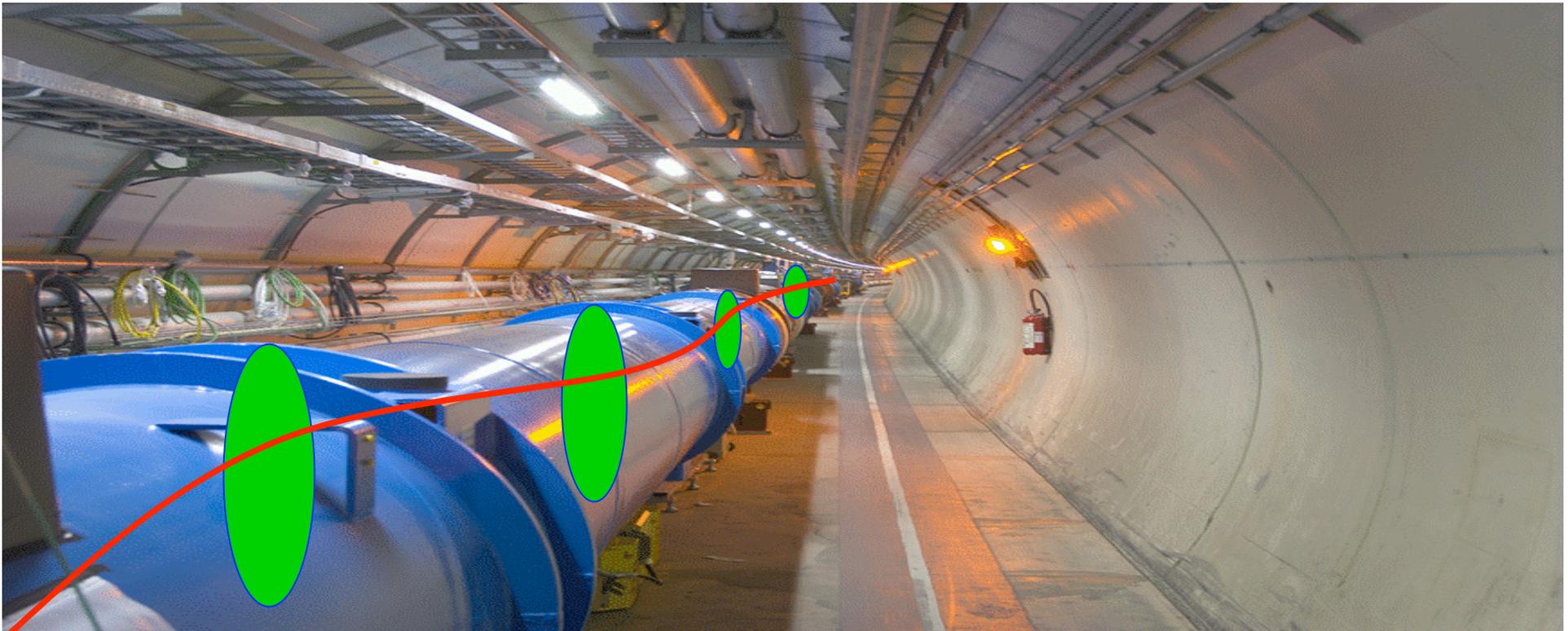
normalise magnet fields to momentum
(remember: $B \cdot \rho = p / q$)

Dipole Magnet

$$\frac{B}{p/q} = \frac{B}{B\rho} = \frac{1}{\rho}$$

Quadrupole Magnet

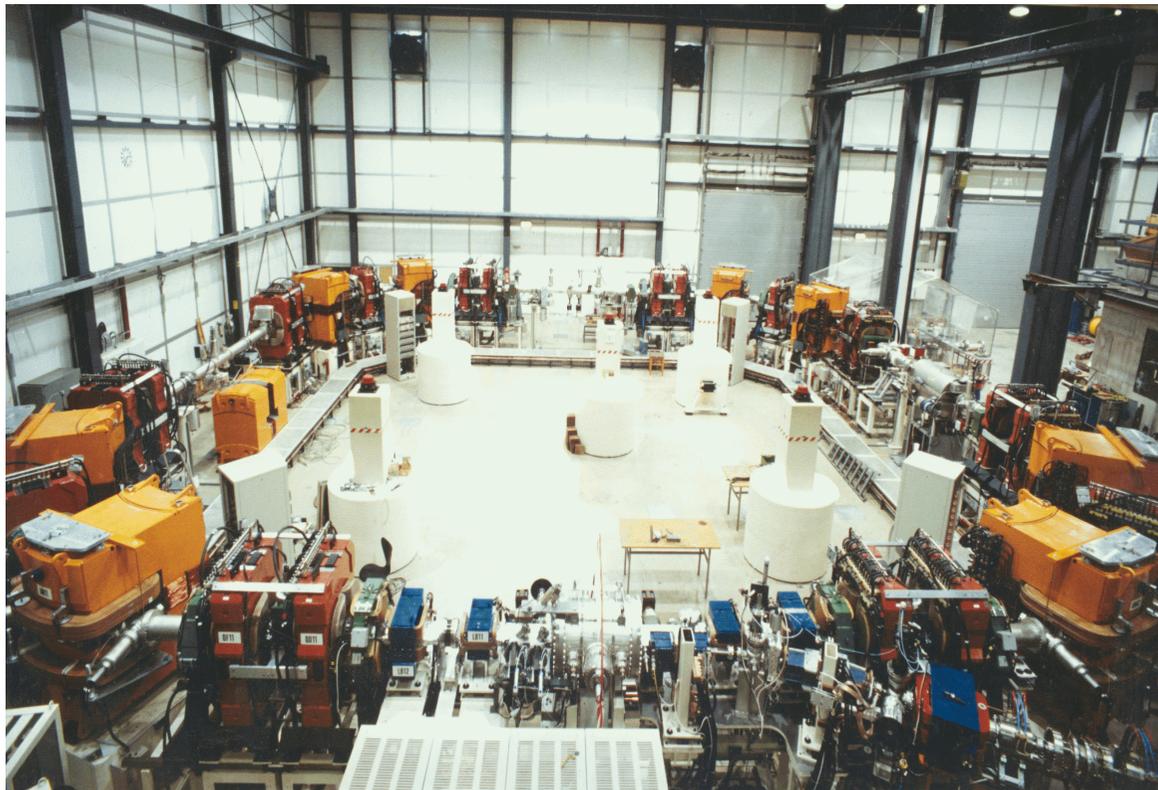
$$k := \frac{g}{p/q}$$



3.) The Equation of Motion:

$$\frac{B(x)}{p/e} = \frac{1}{\rho} + kx + \frac{1}{2!} \cancel{m} x^2 + \frac{1}{3!} \cancel{n} x^3 + \dots$$

only terms linear in x, y taken into account dipole fields
quadrupole fields



Separate Function Machines:

Split the magnets and optimise them according to their job:

bending, focusing etc

Example:
heavy ion storage ring TSR

*
man sieht nur
dipole und quads → linear

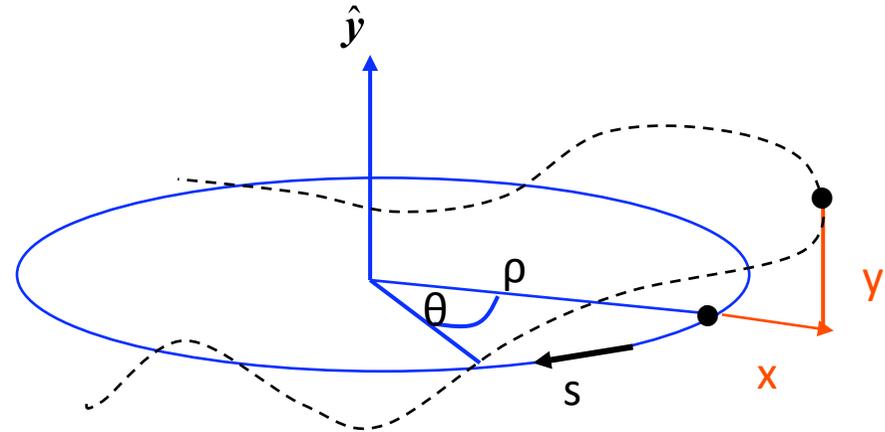
The Equation of Motion:

- * Equation for the horizontal motion:

$$x'' + x \left(\frac{1}{\rho^2} + k \right) = 0$$

x = particle amplitude

x' = angle of particle trajectory (wrt ideal path line)



- * Equation for the vertical motion:

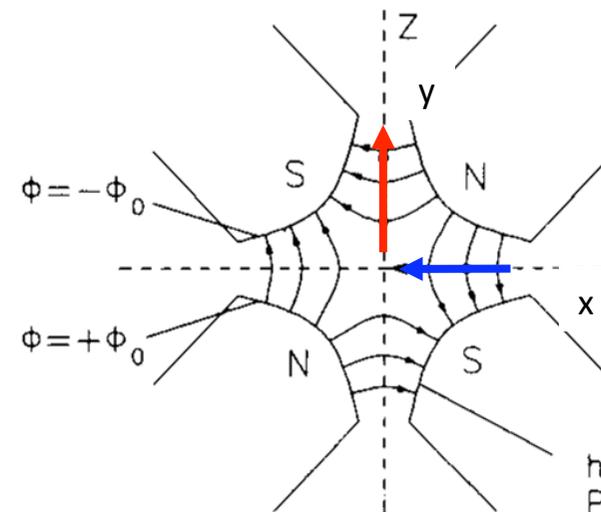
$$\frac{1}{\rho^2} = 0$$

no dipoles ... in general ...

$$k \Leftrightarrow -k$$

quadrupole field changes sign

$$y'' - k y = 0$$



4.) Solution of Trajectory Equations

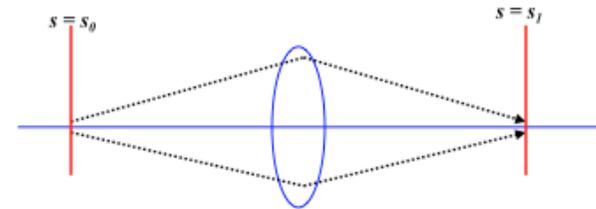
$$\left. \begin{array}{l} \text{Define ... hor. plane: } K = 1/\rho^2 + k \\ \text{... vert. Plane: } K = -k \end{array} \right\} \mathbf{x'' + K x = 0}$$

Differential Equation of harmonic oscillator ... with **spring constant K**

Ansatz: **Hor. Focusing Quadrupole K > 0:**

$$x(s) = x_0 \cdot \cos(\sqrt{|K|}s) + x'_0 \cdot \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}s)$$

$$x'(s) = -x_0 \cdot \sqrt{|K|} \cdot \sin(\sqrt{|K|}s) + x'_0 \cdot \cos(\sqrt{|K|}s)$$



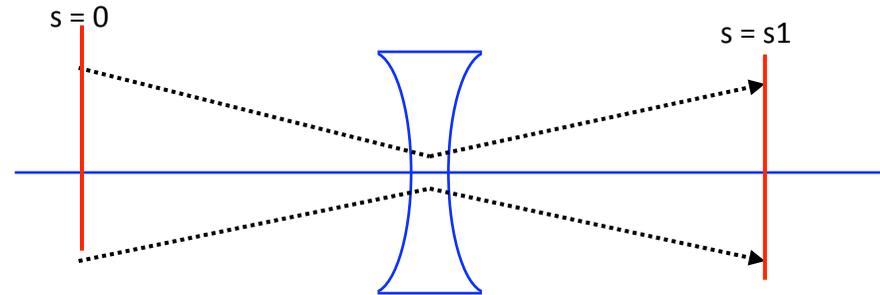
For convenience expressed in matrix formalism:

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s1} = M_{foc} * \begin{pmatrix} x \\ x' \end{pmatrix}_{s0}$$

$$M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}l) \\ -\sqrt{|K|} \sin(\sqrt{|K|}l) & \cos(\sqrt{|K|}l) \end{pmatrix}$$

hor. defocusing quadrupole:

$$x'' - K x = 0$$



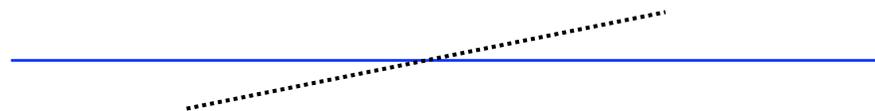
Ansatz: Remember from school

$$x(s) = a_1 \cdot \cosh(\omega s) + a_2 \cdot \sinh(\omega s)$$

$$M_{defoc} = \begin{pmatrix} \cosh \sqrt{|K|} l & \frac{1}{\sqrt{|K|}} \sinh \sqrt{|K|} l \\ \sqrt{|K|} \sinh \sqrt{|K|} l & \cosh \sqrt{|K|} l \end{pmatrix}$$

drift space:

$$K = 0$$



$$x(s) = x'_0 * s$$

$$M_{drift} = \begin{pmatrix} 1 & l \\ 0 & 1 \end{pmatrix}$$

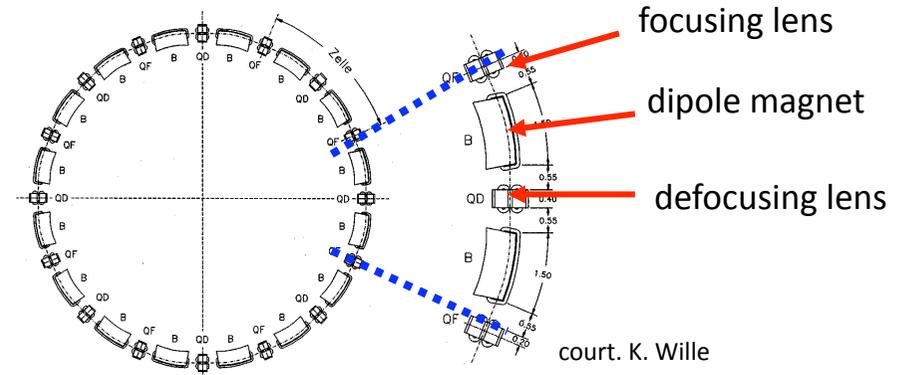
! with the assumptions made, the motion in the horizontal and vertical planes are independent „ ... the particle motion in x & y is uncoupled“

Transformation through a system of lattice elements

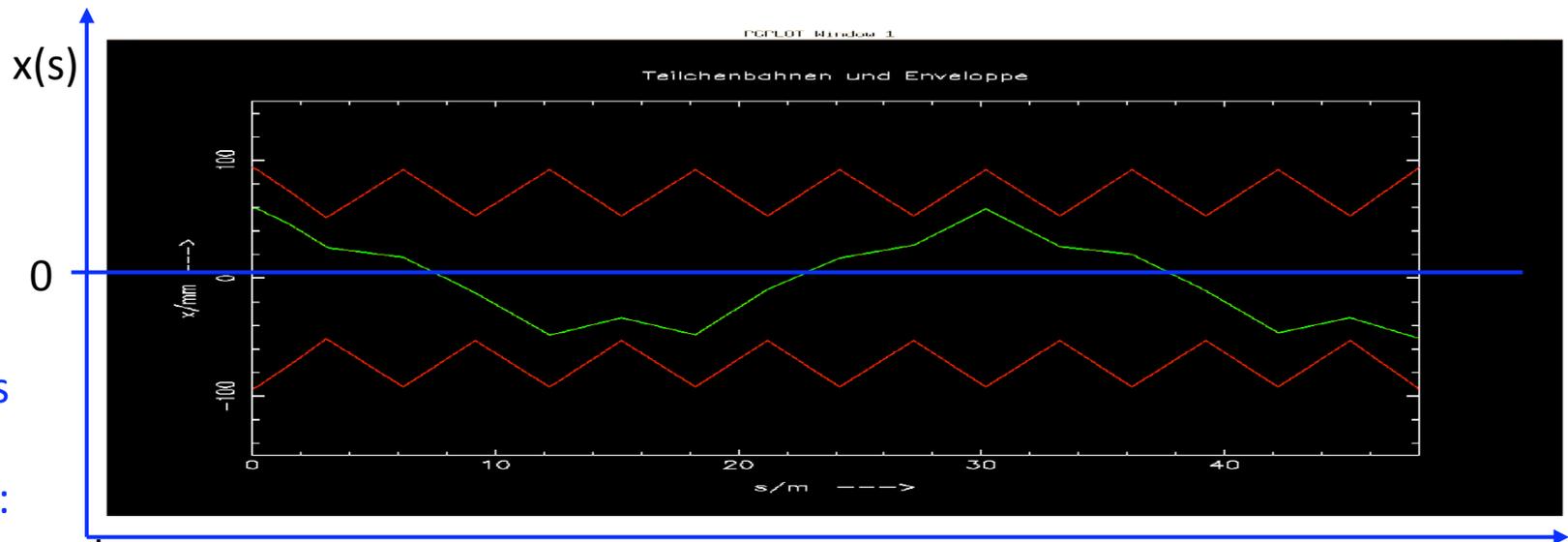
combine the single element solutions by multiplication of the matrices

$$M_{total} = M_{QF} * M_D * M_{QD} * M_{Bend} * M_D * \dots$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_2} = M(s_2, s_1) * \begin{pmatrix} x \\ x' \end{pmatrix}_{s_1}$$



in each accelerator element the particle trajectory corresponds to the movement of a harmonic oscillator „



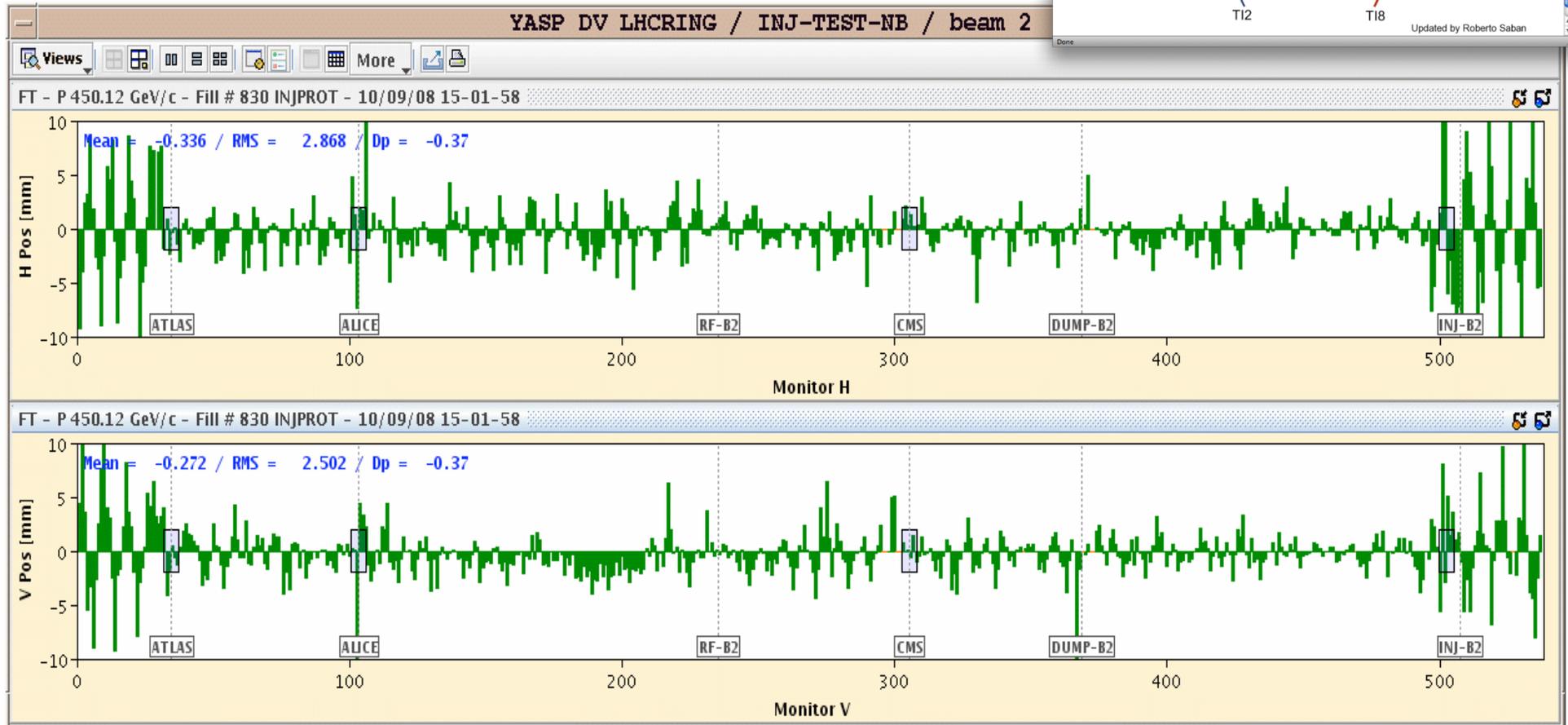
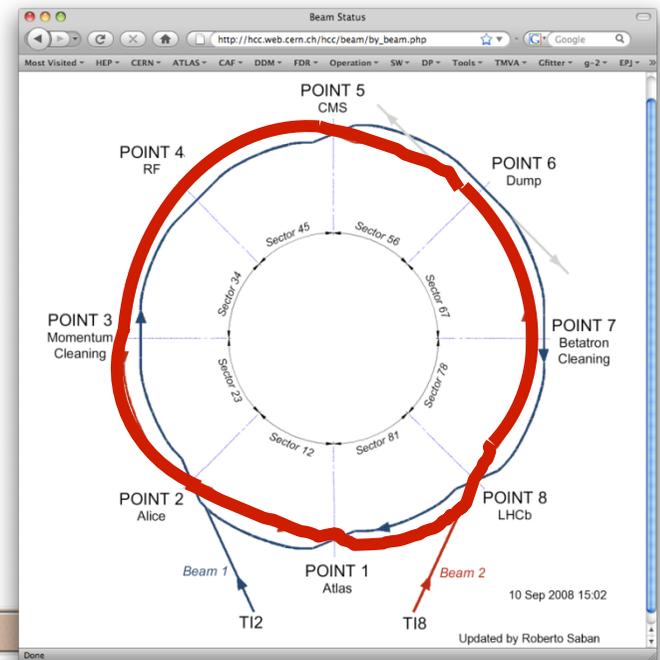
typical values
in a strong
foc. machine:

$x \approx \text{mm}, x' \leq \text{mrad}$

LHC Operation: Beam Commissioning

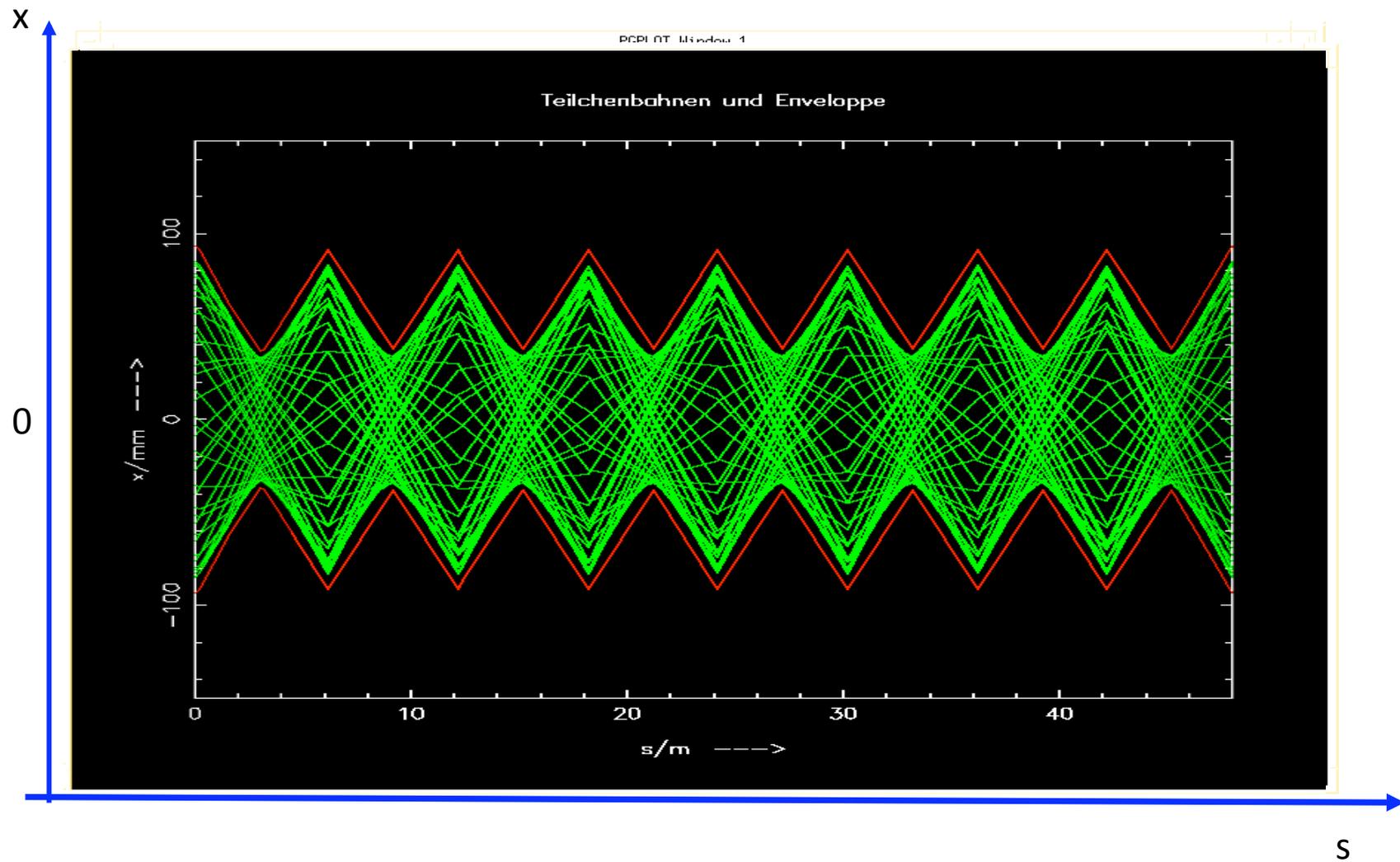
The transverse focusing fields create a harmonic oscillation of the particles with a well defined "Eigenfrequency" which is called **tune**

First turn steering "by sector:"



Question: what will happen, if the particle performs a second turn ?

... or a third one or ... 10^{10} turns



Astronomer Hill:

differential equation for motions with periodic focusing properties
„Hill's equation“

Example: particle motion with
periodic coefficient

equation of motion: $x''(s) - k(s)x(s) = 0$



restoring force \neq const,
 $k(s)$ = depending on the position s
 $k(s+L) = k(s)$, periodic function

} we expect a kind of quasi harmonic
oscillation: amplitude & phase will depend
on the position s in the ring.

Amplitude of a particle trajectory:

$$x(s) = \sqrt{\varepsilon} * \sqrt{\beta(s)} * \cos(\psi(s) + \varphi)$$

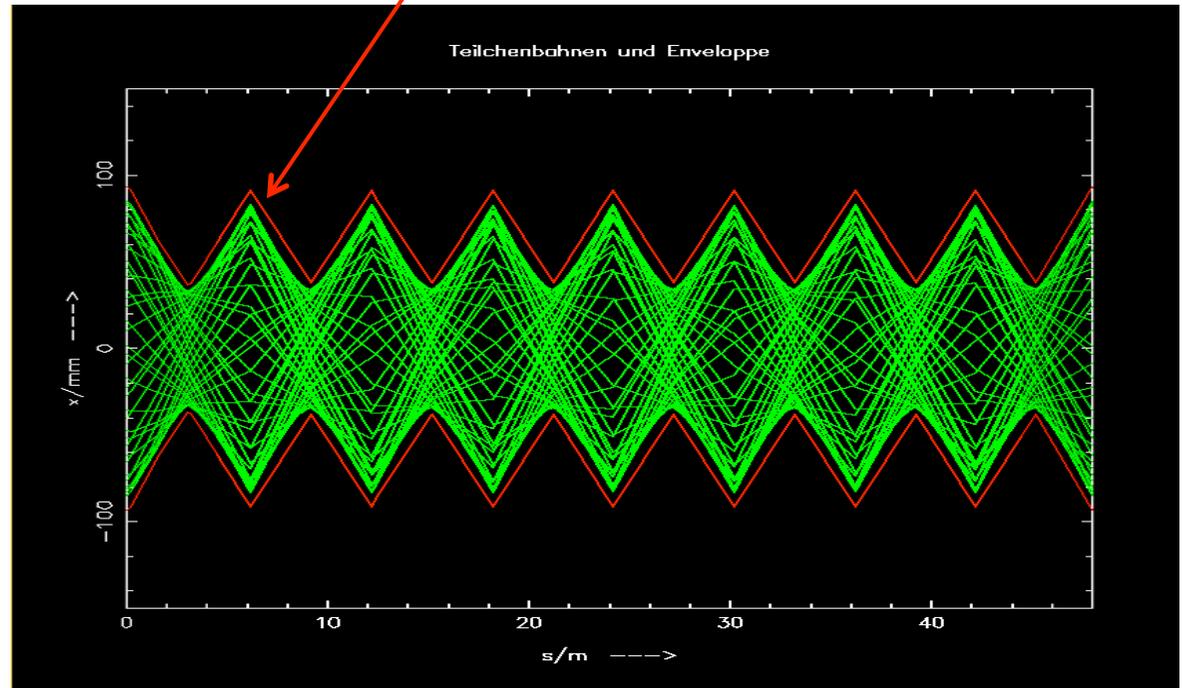
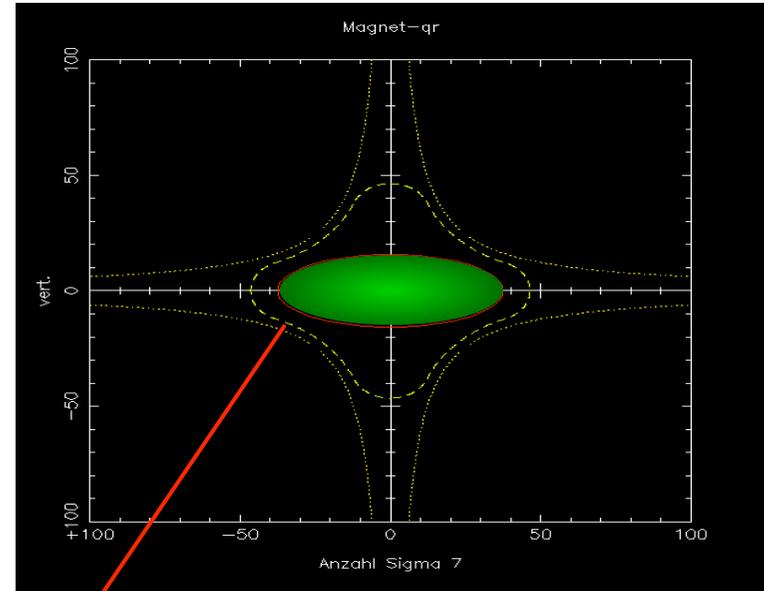
Maximum size of a particle amplitude

$$\hat{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)}$$

The Beta Function

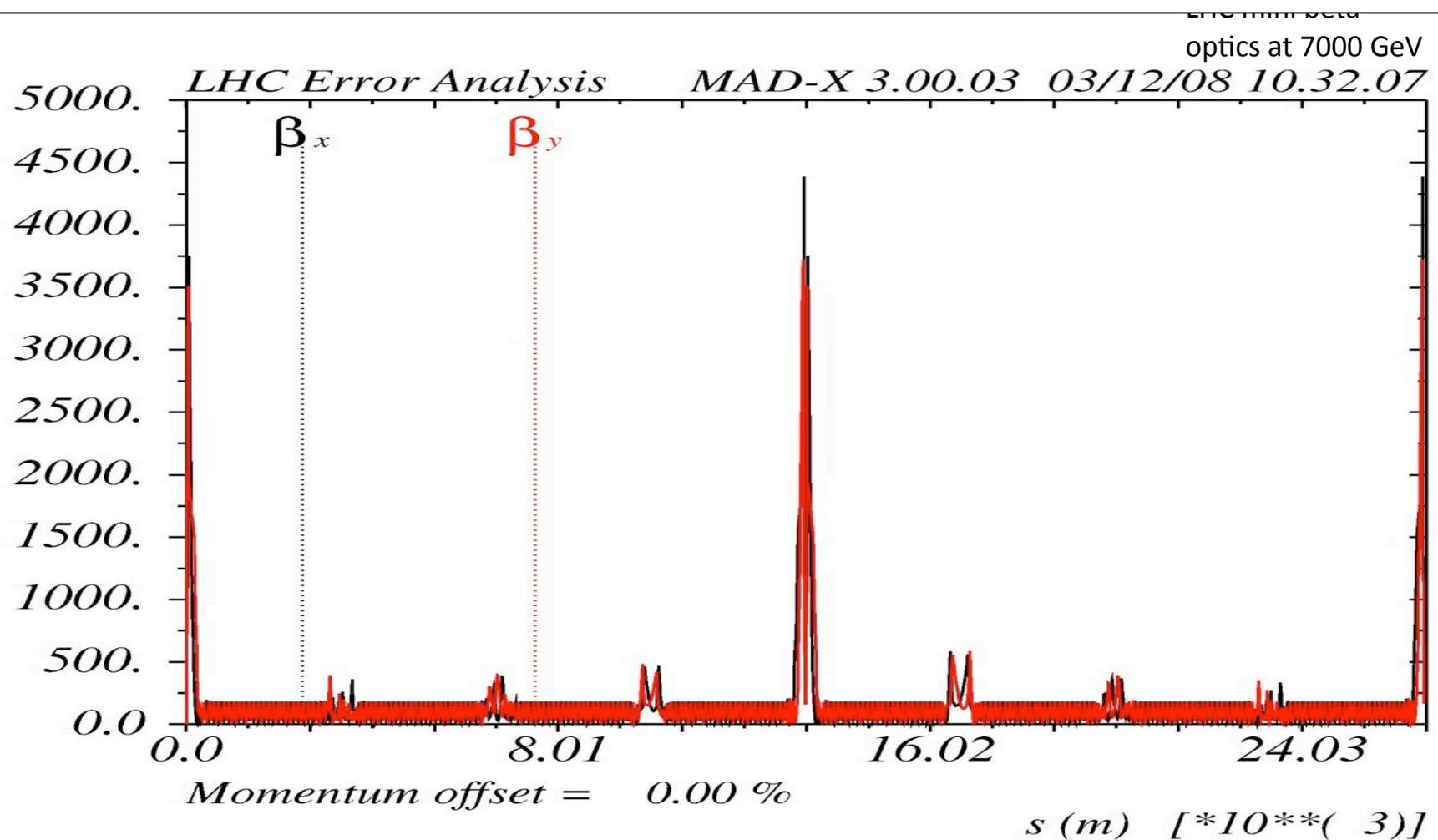
β determines the beam size
... the envelope of all particle trajectories at a given position
“s” in the storage ring under the influence of all (!) focusing fields.

It reflects the periodicity of the magnet structure.



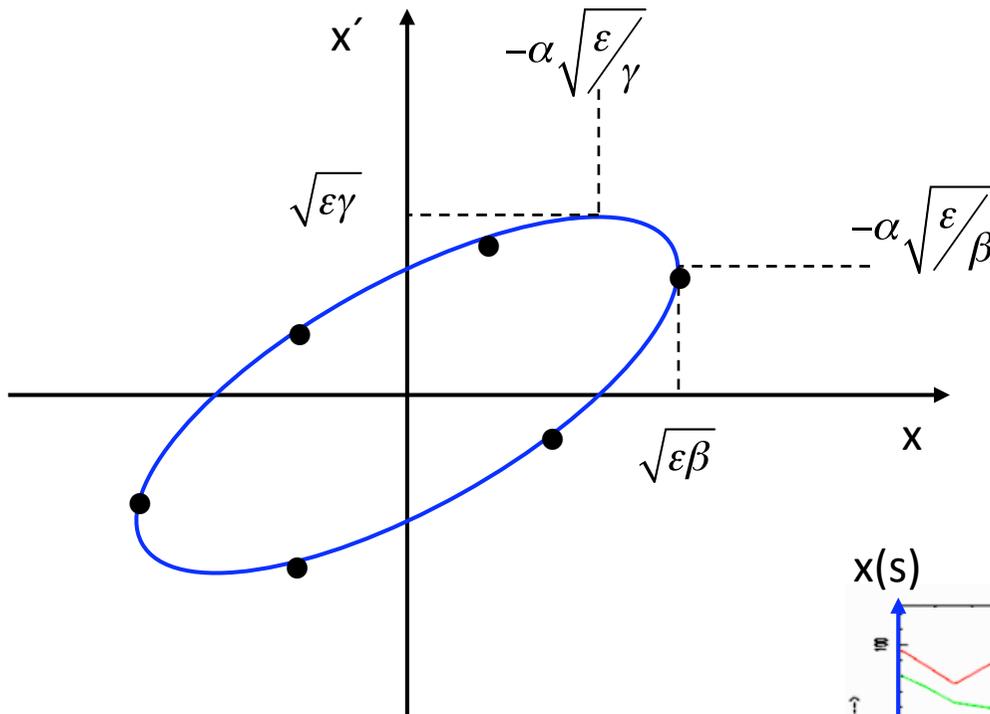
The Beta Function: Lattice Design & Beam Optics

The beta function determines the maximum amplitude a single particle trajectory can reach at a given position in the ring. It is determined by the focusing properties of the lattice and follows the periodicity of the machine.



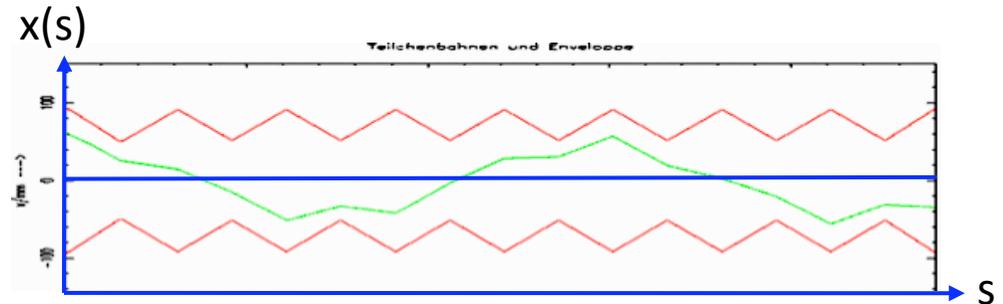
Beam Emittance and Phase Space Ellipse

$$\varepsilon = \gamma(s) * x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s)x'(s)^2$$



Liouville: in reasonable storage rings
area in phase space is constant.

$$A = \pi * \varepsilon = \text{const}$$



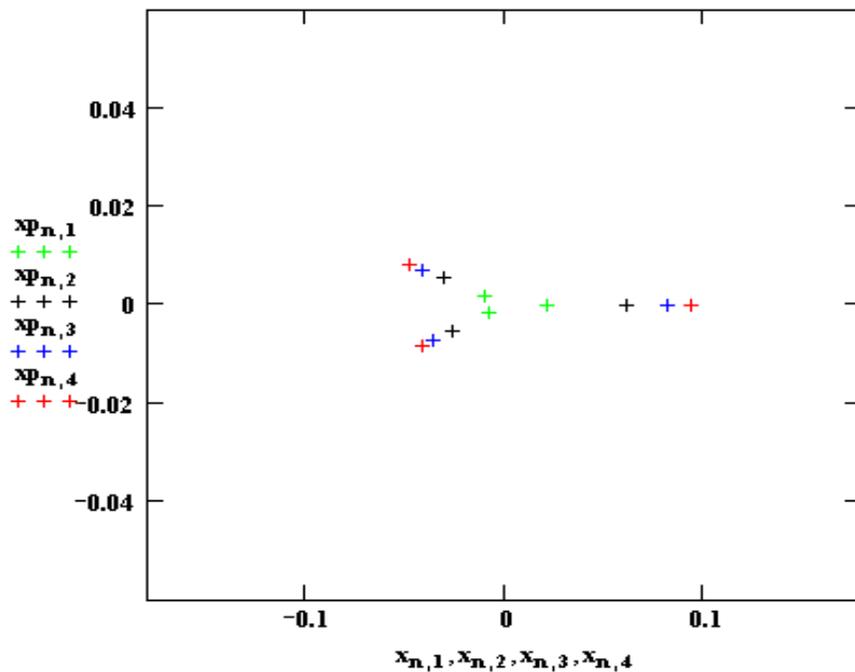
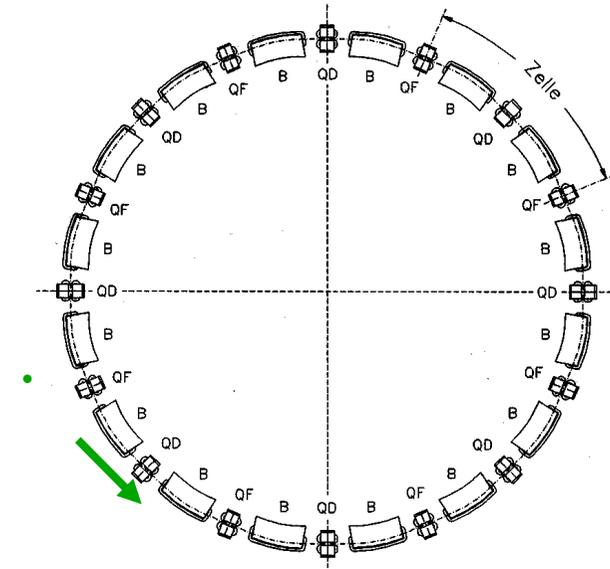
ε beam emittance = **woozilycity** of the particle ensemble, **intrinsic beam parameter**, cannot be changed by the foc. properties.

Scientificquely spoken: area covered in **transverse x, x' phase space** ... and it is constant !!!

Particle Tracking in a Storage Ring

Calculate x, x' for each accelerator element according to matrix formalism and plot x, x' at a given position „s“ in the phase space diagram

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s1} = M_{turn} * \begin{pmatrix} x \\ x' \end{pmatrix}_{s0}$$



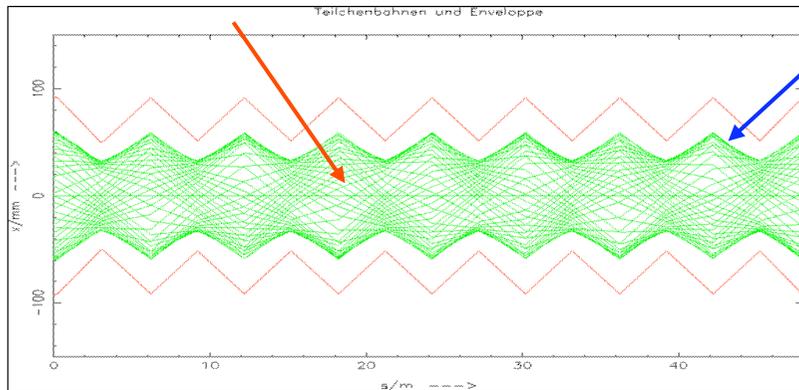
A beam of 4 particles

– each having a slightly different emittance:

Emittance of the Particle Ensemble:

$$x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\Psi(s) + \phi)$$

$$\hat{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)}$$



single particle trajectories, $N \approx 10^{11}$ per bunch

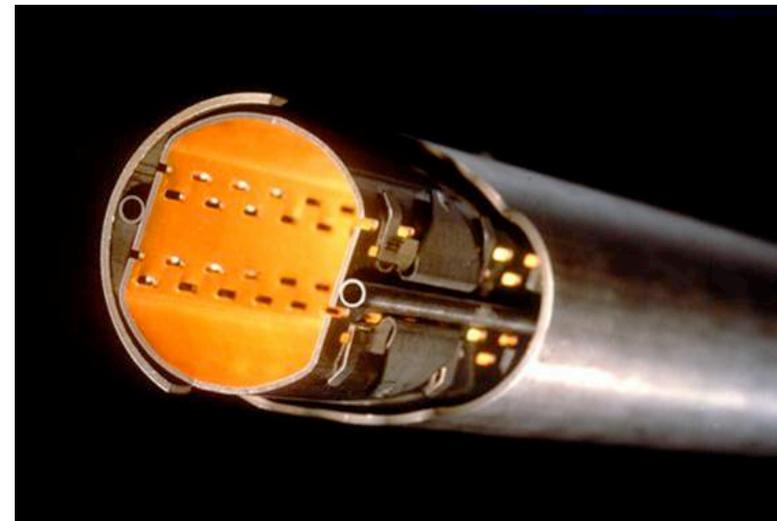
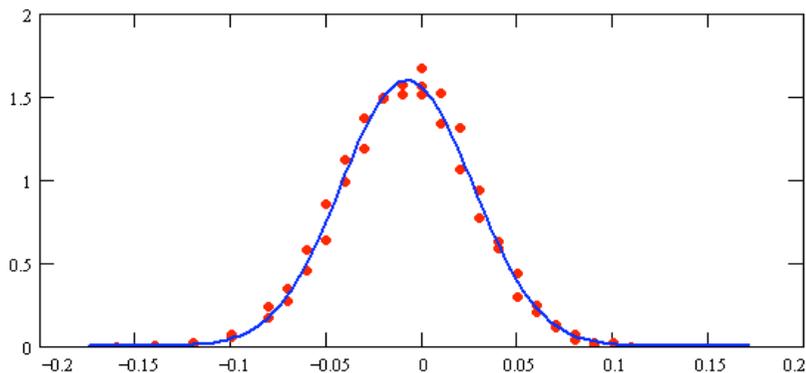
Gauß
Particle Distribution:

$$\rho(x) = \frac{N \cdot e}{\sqrt{2\pi}\sigma_x} \cdot e^{-\frac{1}{2}\frac{x^2}{\sigma_x^2}}$$

particle at distance 1σ from centre
 \leftrightarrow 68.3 % of all beam particles

LHC: $\beta = 180 m$
 $\varepsilon = 5 * 10^{-10} m rad$

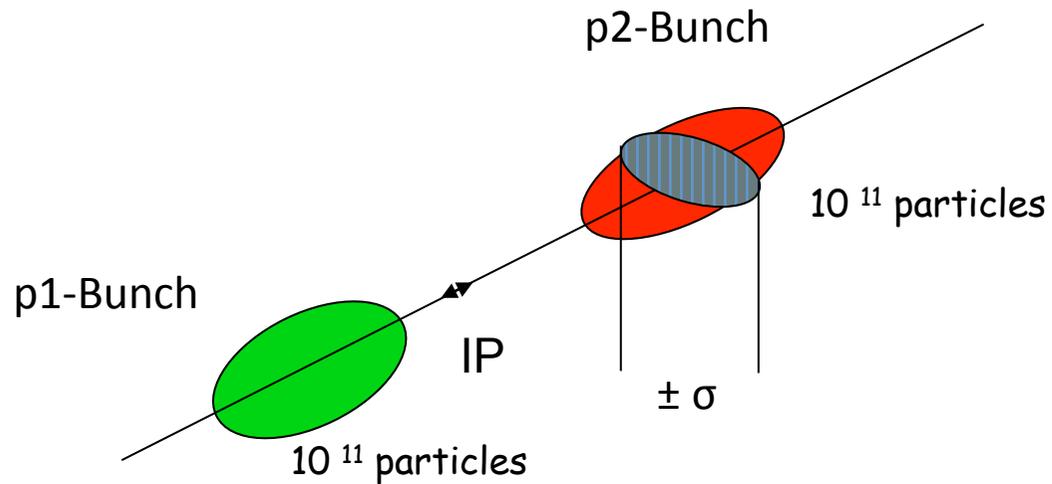
$$\sigma = \sqrt{\varepsilon * \beta} = \sqrt{5 * 10^{-10} m * 180 m} = 0.3 mm$$



aperture requirements: $r_0 = 17 * \sigma$

5.) Luminosity

$$R = L * \Sigma_{react}$$



Example: Luminosity run at LHC

$$\beta_{x,y} = 0.55 \text{ m}$$

$$f_0 = 11.245 \text{ kHz}$$

$$\varepsilon_{x,y} = 5 * 10^{-10} \text{ rad m}$$

$$n_b = 2808$$

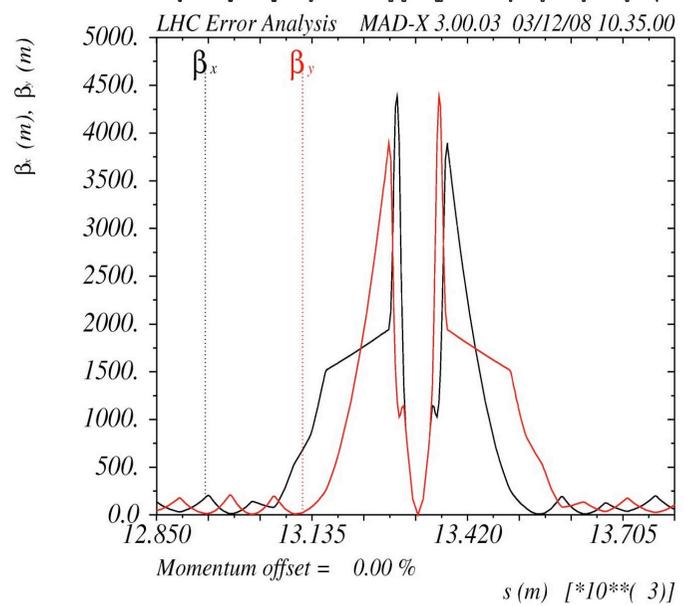
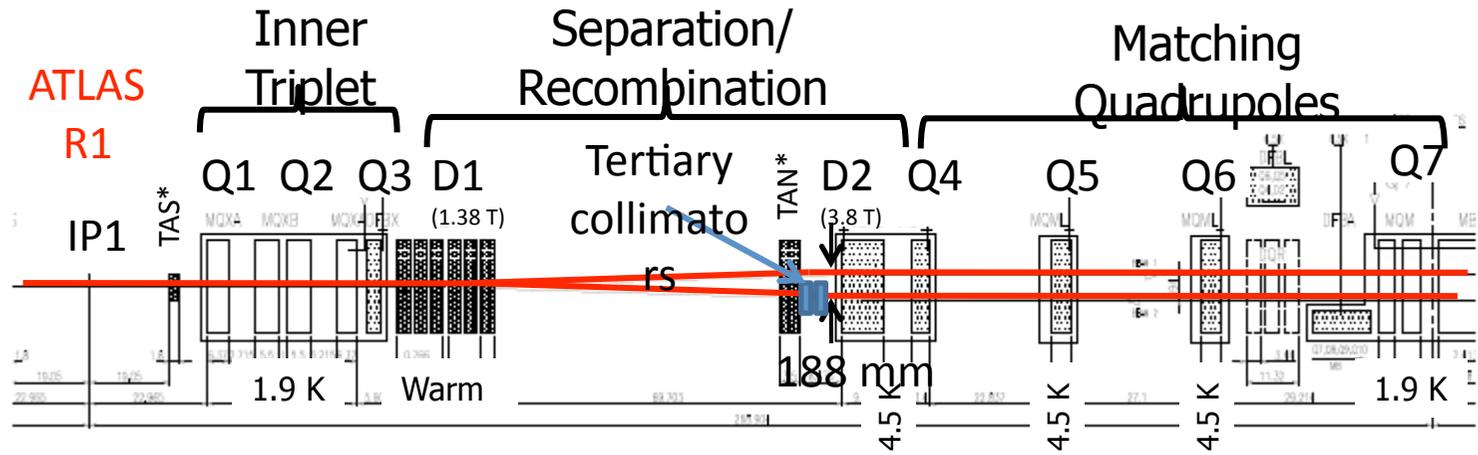
$$\sigma_{x,y} = 17 \text{ } \mu\text{m}$$

$$L = \frac{1}{4\pi e^2 f_0 n_b} * \frac{I_{p1} I_{p2}}{\sigma_x \sigma_y}$$

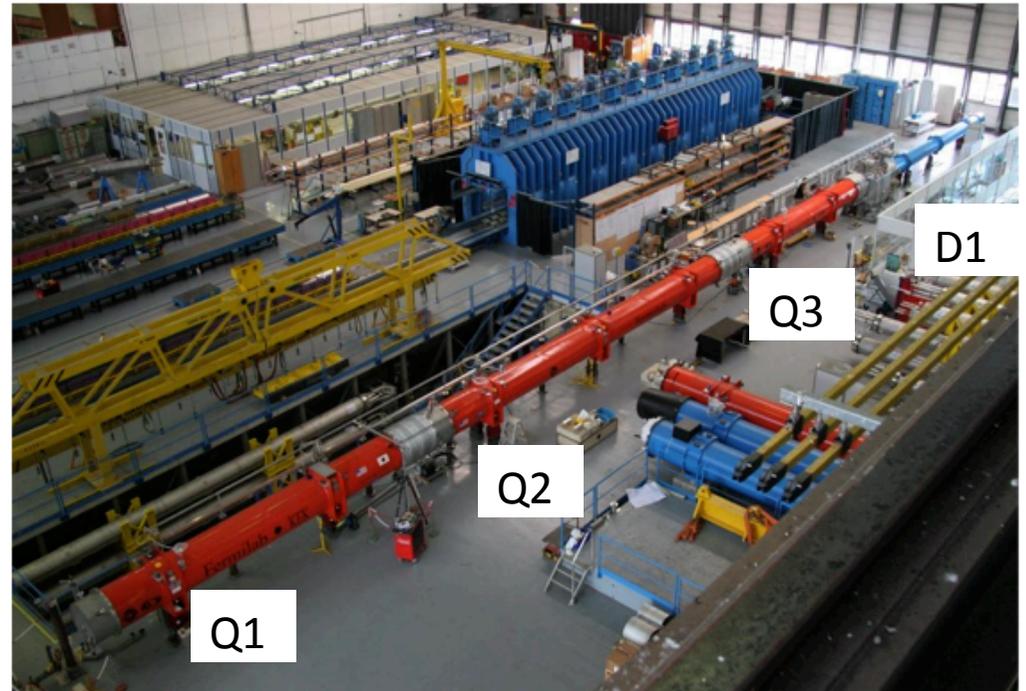
$$I_p = 584 \text{ mA}$$

$$L = 1.0 * 10^{34} \text{ } 1/\text{cm}^2 \text{ s}$$

The LHC Mini-Beta-Insertions



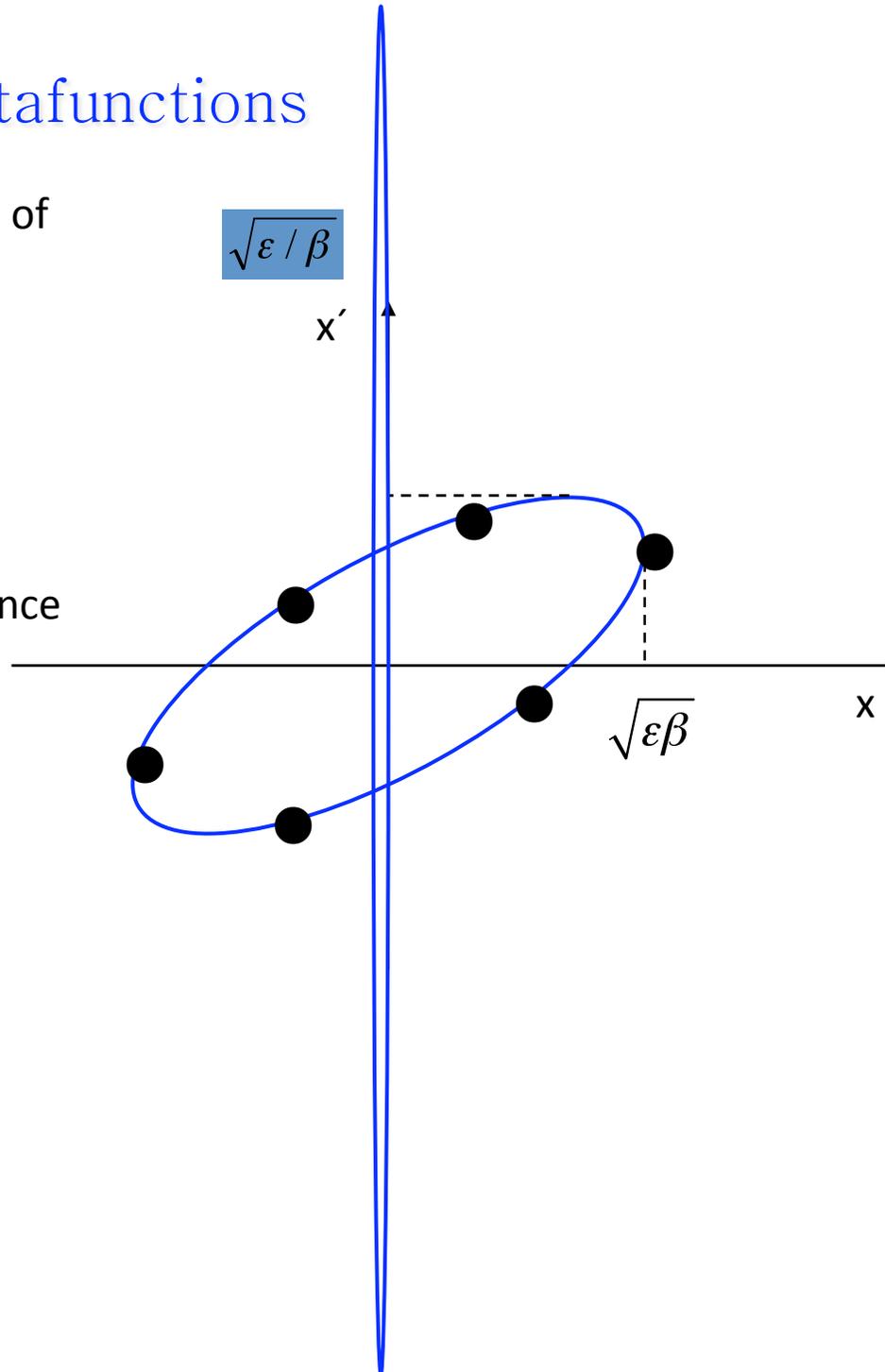
mini β optics



Mini- β Insertions: Betafunctions

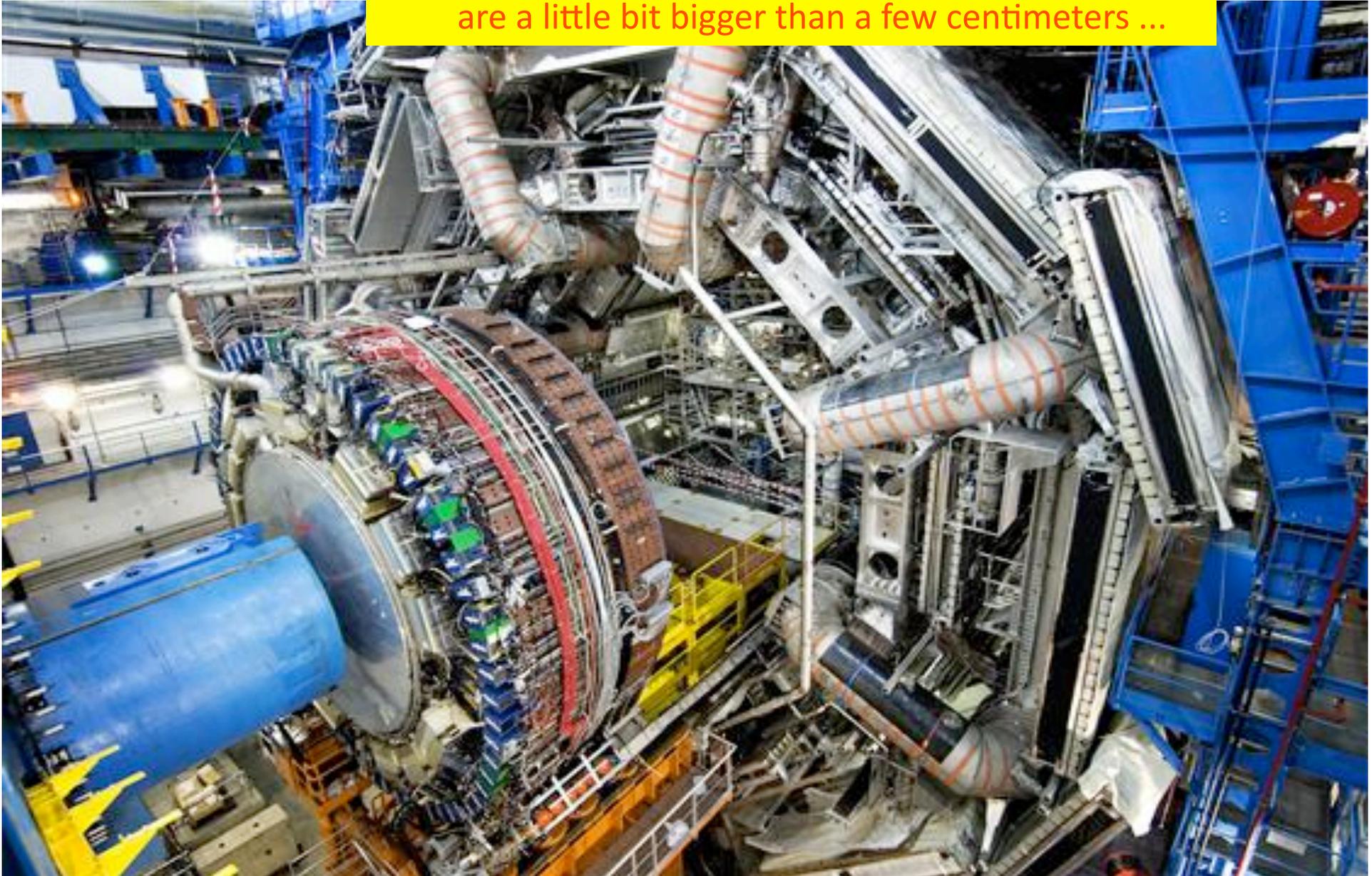
A mini- β insertion is always a kind of
special symmetric drift space.
→ greetings from Liouville

the smaller the beam size
the larger the beam divergence



... clearly there is ano

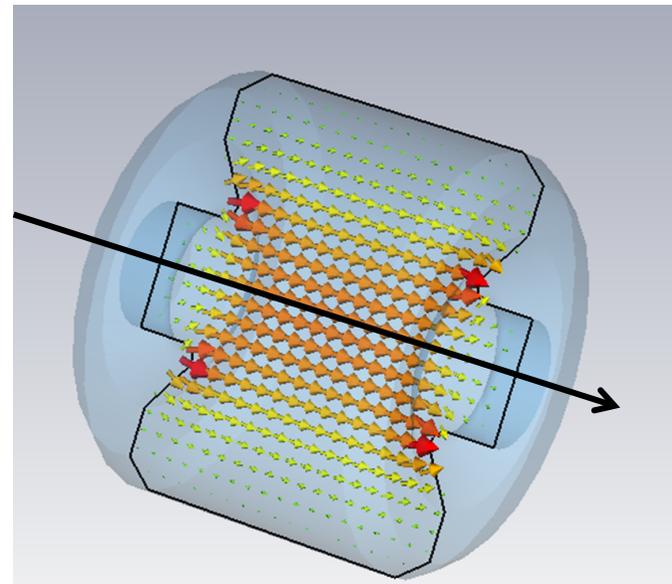
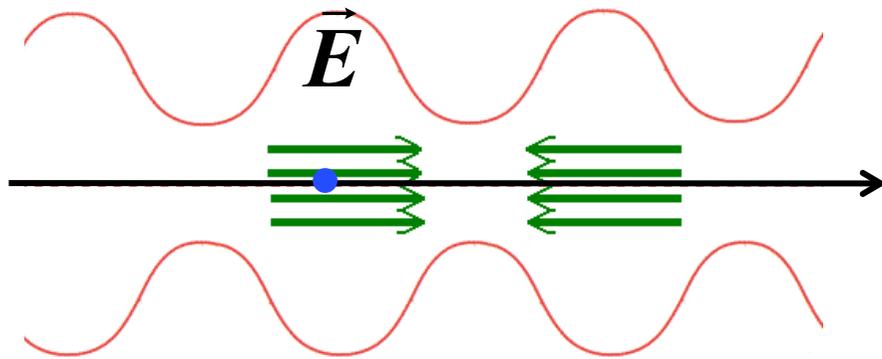
... unfortunately ... in general
high energy detectors that are
installed in that drift spaces
are a little bit bigger than a few centimeters ...



III. The Acceleration

Where is the acceleration?

Install an RF accelerating structure in the ring:



B. Salvant
N. Biancacci

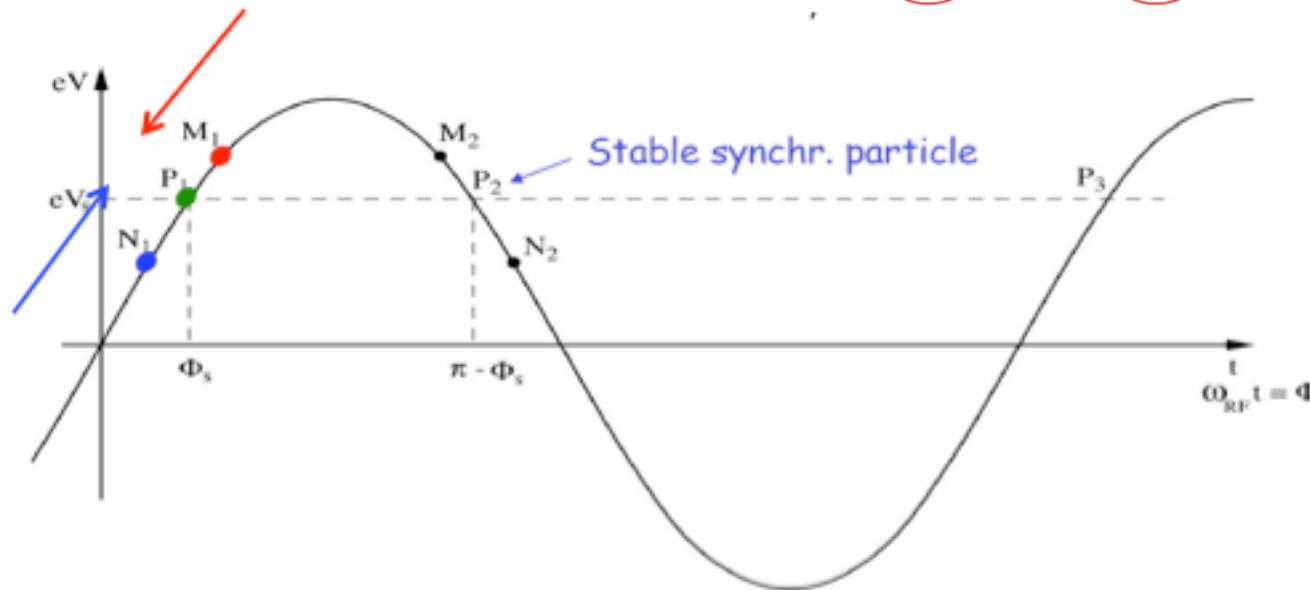
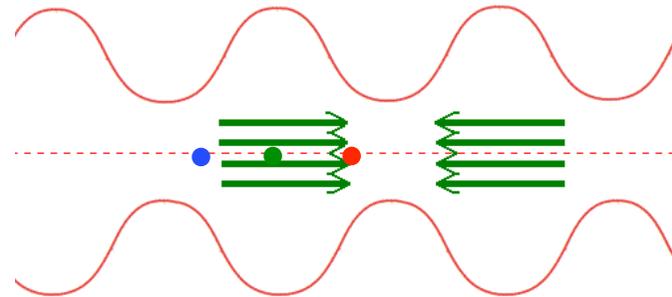
The Acceleration & "Phase Focusing"

$\Delta p/p \neq 0$ below transition

ideal particle •

particle with $\Delta p/p > 0$ • faster

particle with $\Delta p/p < 0$ • slower



Focussing effect in the longitudinal direction keeping the particles close together ... forming a "bunch"

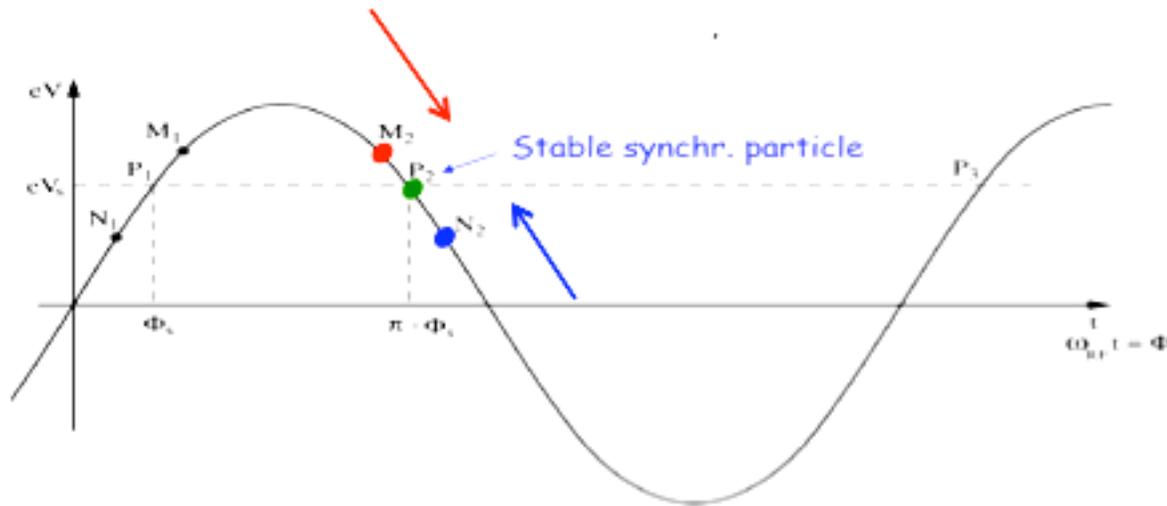
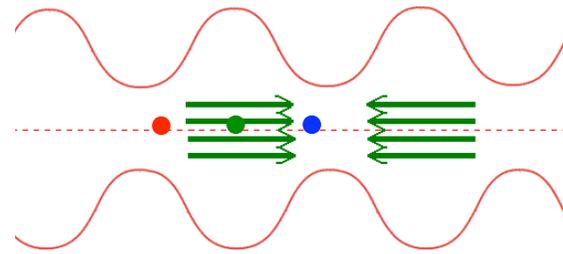
oscillation frequency: $f_s = f_{rev} \sqrt{-\frac{h\alpha_s * qU_0 \cos \phi_s}{2\pi E_s}} \approx \text{some Hz}$

The Acceleration above transition

ideal particle •

particle with $\Delta p/p > 0$ • heavier

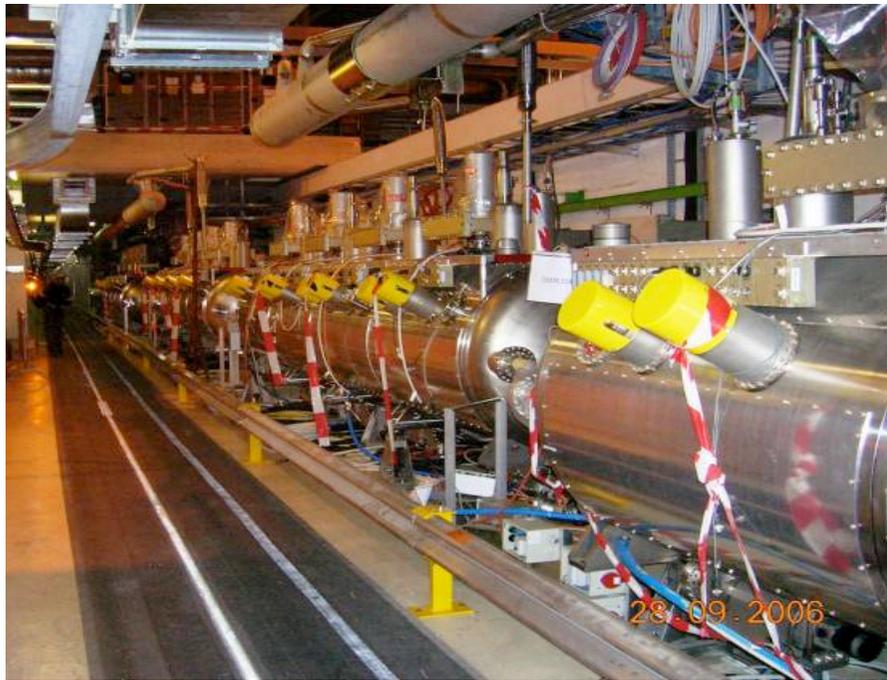
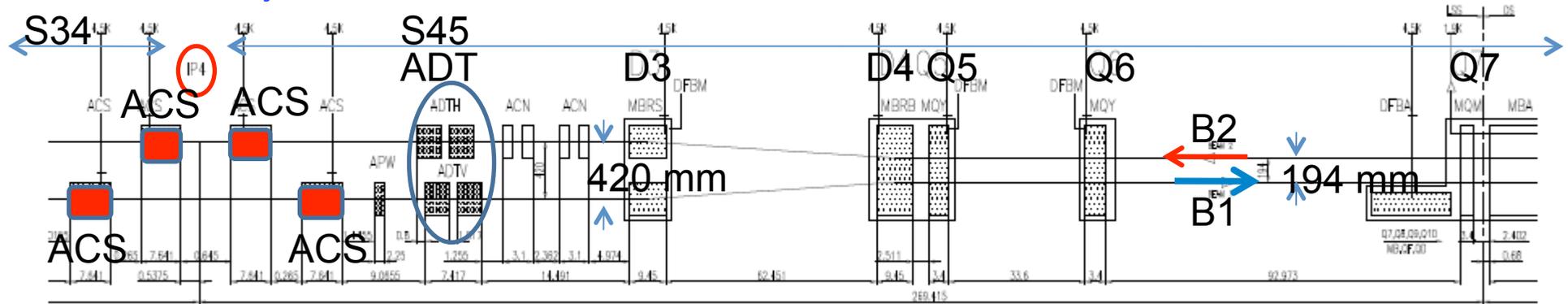
particle with $\Delta p/p < 0$ • lighter



Focussing effect in the longitudinal direction
 keeping the particles close together ... forming a “bunch”

... and how do we accelerate now ???
 with the dipole magnets !

The RF system: IR4

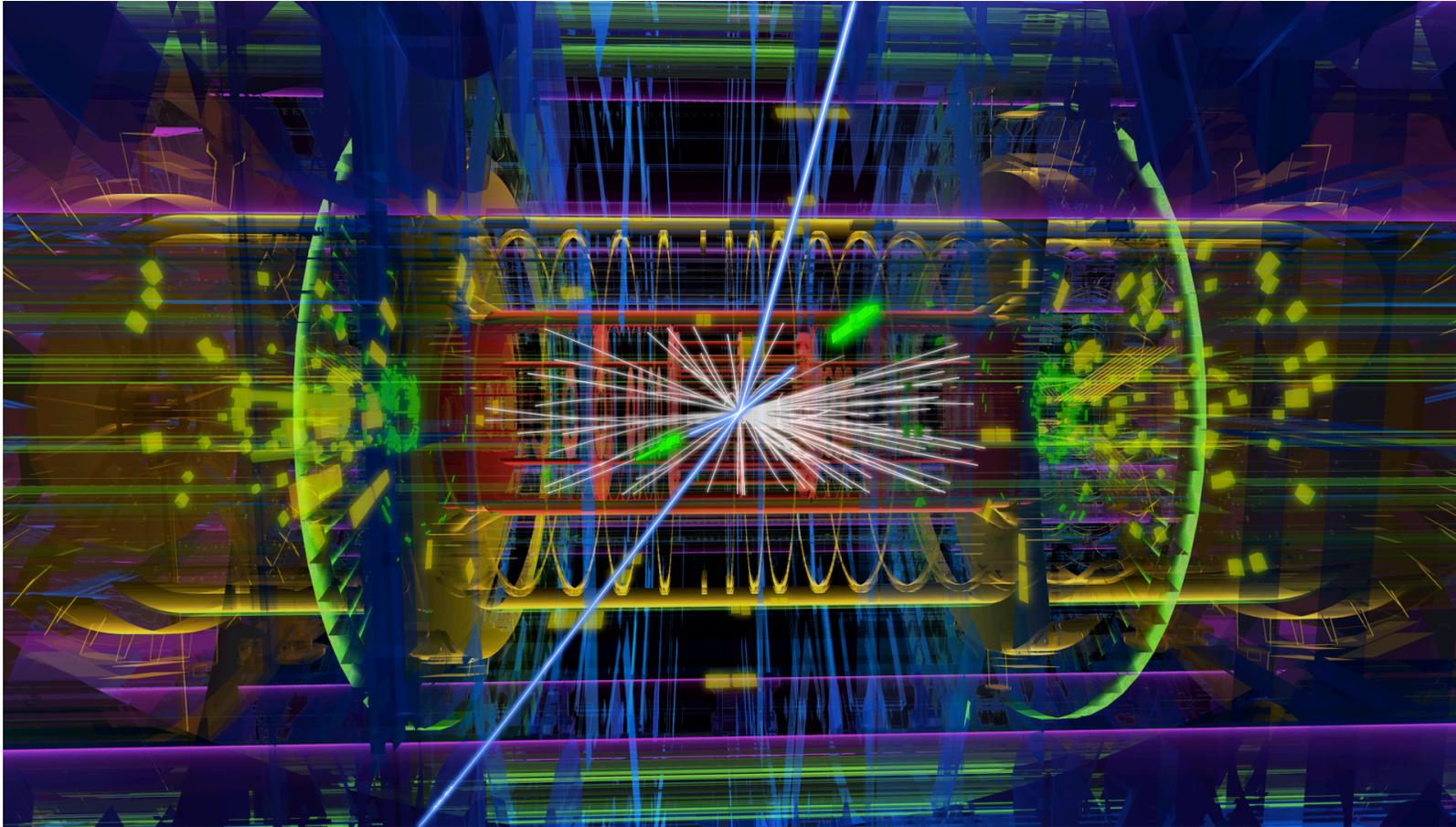


Nb on Cu cavities @4.5 K (=LEP2)
Beam pipe diam.=300mm

<i>Bunch length (4σ)</i>	<i>ns</i>	<i>1.06</i>
<i>Energy spread (2σ)</i>	<i>10^{-3}</i>	<i>0.22</i>
<i>Synchr. rad. loss/turn</i>	<i>keV</i>	<i>7</i>
<i>Synchr. rad. power</i>	<i>kW</i>	<i>3.6</i>
<i>RF frequency</i>	<i>M</i>	<i>400</i>
	<i>Hz</i>	
<i>Harmonic number</i>		<i>35640</i>
<i>RF voltage/beam</i>	<i>MV</i>	<i>16</i>
<i>Energy gain/turn</i>	<i>keV</i>	<i>485</i>
<i>Synchrotron frequency</i>	<i>Hz</i>	<i>23.0</i>

... und wozu das alles ??

High Light of the HEP-Year natuerlich das HIGGS

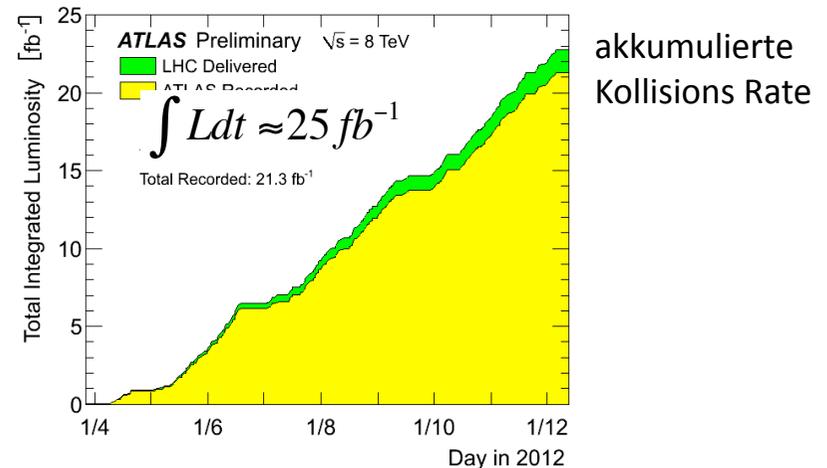
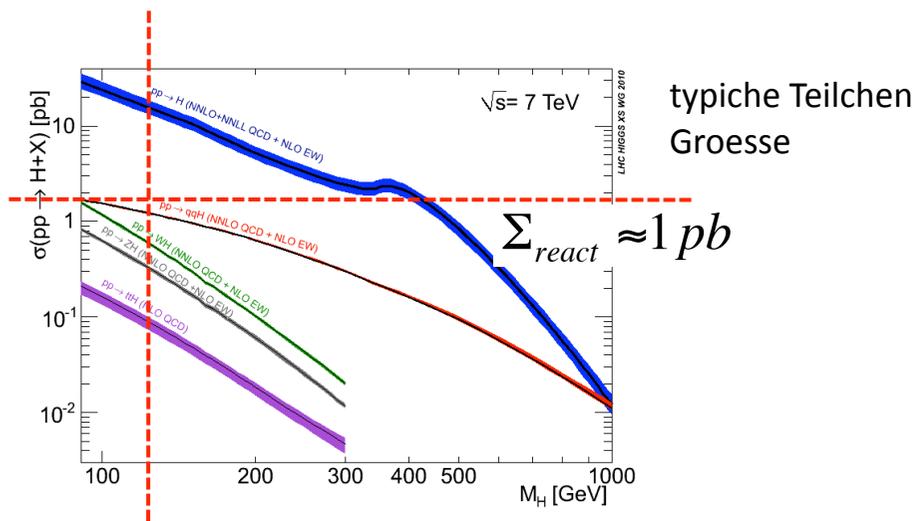


ATLAS event display: Higgs => two electrons & two muons

The High light of the year

production rate of events is determined by the cross section Σ_{react} and a parameter L that is given by the design of the accelerator:
 ... the luminosity

$$R = L * \Sigma_{react} \approx 10^{-12} b \cdot 25 \frac{1}{10^{-15} b} = \text{some } 1000 H$$



$$1b = 10^{-24} \text{ cm}^2 = 1 \text{ mio} * 1 \text{ mio} * 1 \text{ mio} * \frac{1}{10000} \text{ mm}^2$$

Die Teilchen sind "sehr klein"

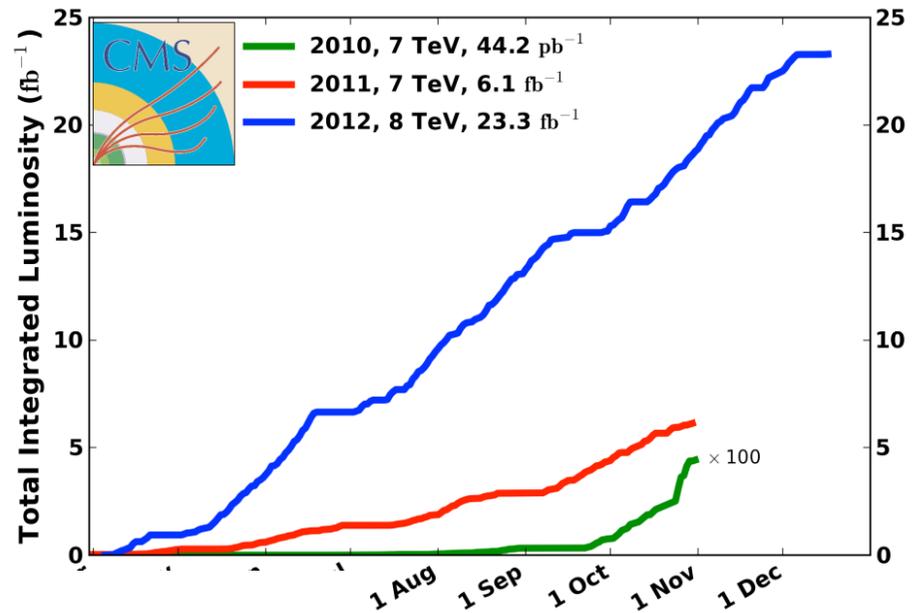
The luminosity is a storage ring quality parameter and depends on beam size (β !!) and stored current

$$L = \frac{1}{4\pi e^2 f_0 b} * \frac{I_1 * I_2}{\sigma_x^* * \sigma_y^*}$$

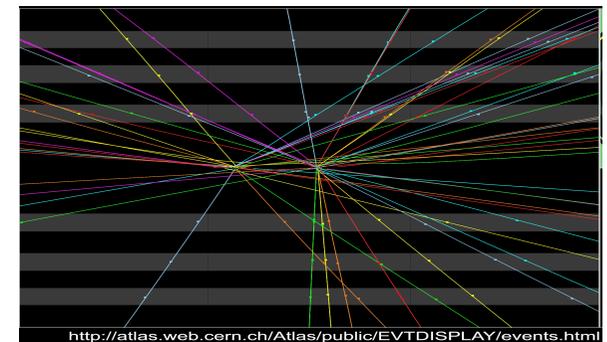
$\Delta p/p = 5 * 10^{-4}$

And still...

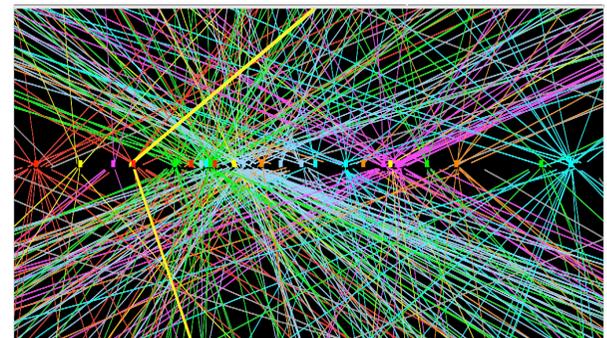
The LHC Performance in Run 1



Design	2010	2012
Momentum at collision	7 TeV /c	4 TeV/c
Luminosity	$10^{34} \text{ cm}^{-2} \text{ s}^{-1}$	$7.7 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
Protons per bunch	1.15×10^{11}	1.50×10^{11}
Number of bunches/beam	2808	1380
Nominal bunch spacing	25 ns	50ns
Normalized emittance	3.75 μm	2.5 μm
beta *	55 cm	60 cm
rms beam size (arc)	300 μm	350 μm
rms beam size IP	17 μm	20 μm



2 vertices



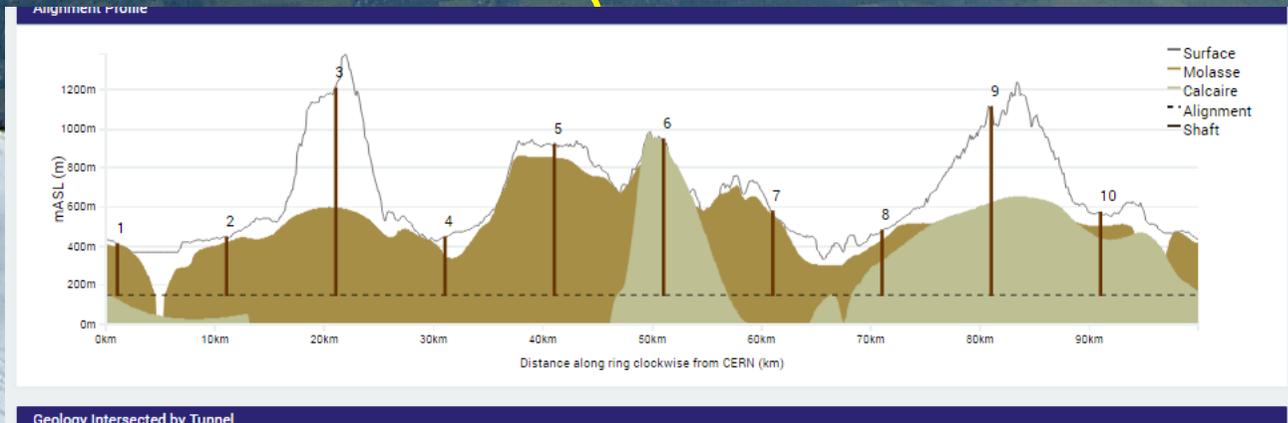
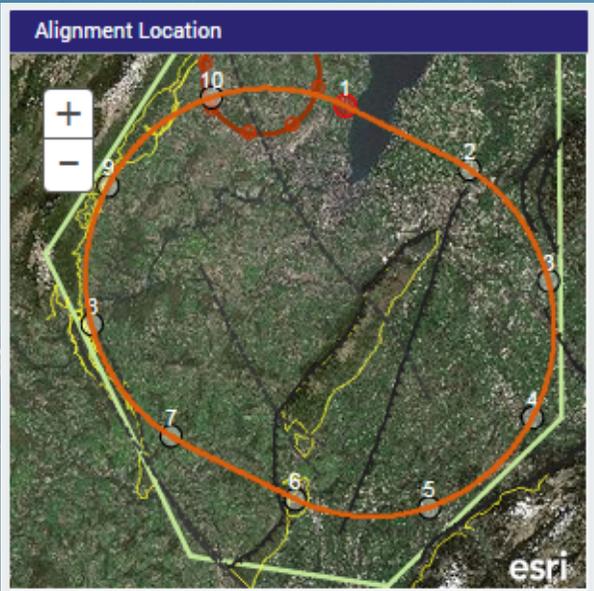
20 vertices

European Strategy Recommendation for the future of particle physics: FCC

LHC: 7000 GeV
FCC: 50000 GeV

$$\frac{p}{e} = B \rho$$





Geology Intersected by Tunnel