

FEMTOSCOPY AND HADROCHEMISTRY

MIKE LISA

OHIO STATE UNIVERSITY

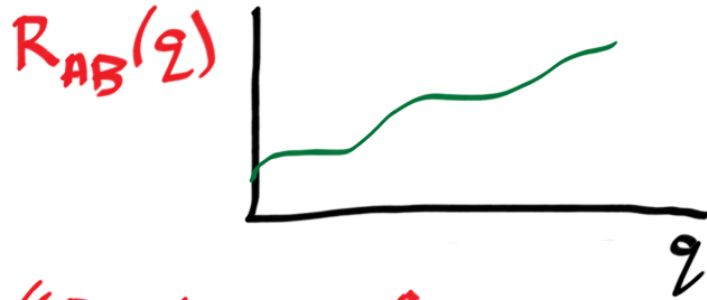
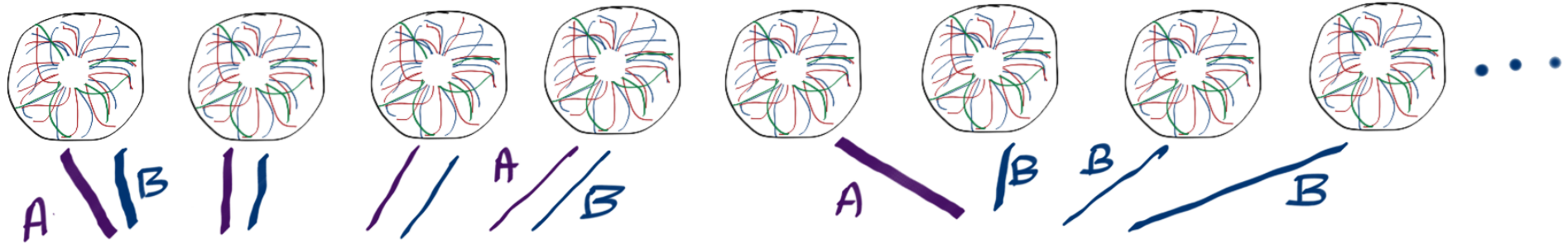
* in the spirit of a femtoscopy workshop

OUTLINE

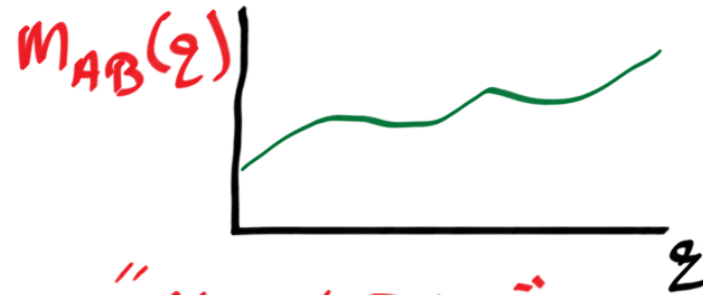
- TWO VIEWS OF A 2-PARTICLE CORRELATION
- STAR π - π PUZZLE
- FEMTOSCOPY - NUCLEO/HADRO CHEMISTRY "DUALITY"
 - IN SIMPLEST MODEL
- "DOUBLE COUNTING" ?
- MANY QUESTIONS
- NO CONCLUSION SLIDE

* in the spirit of a femtoscopy workshop

"MIXED EVENT" APPROACH:



"Real pairs"



"Mixed Pairs"

$q = \text{difference measure: } k^*, M_{inv}, Q, \dots$

RESONANCE $D^* \rightarrow A+B$ (e.g. $\Delta^+ \rightarrow p \pi^0$)

MANIFESTS AS AN ENHANCEMENT IN R_{AB}

- D^* • ATTRACTIVE STRONG FSI b/t $A+B$?
- "PARTICLE"?

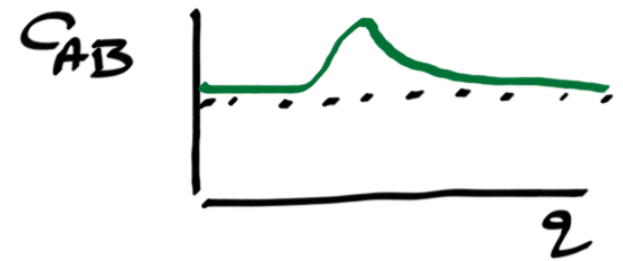
$R_{AB}(z)$ - "real" pairs $M_{AB}(z)$ - "mixed" pairs

FEMTOSCOPY

$$C_{AB}(z) = N \frac{R_{AB}(z)}{M_{AB}(z)} = \int d^3r S_{AB}(r') |\Phi(z, r')|^2$$

KOONIN - PRATT EQ.

- sensitive to spatial scales
("HBT radii", homo volume...)



- insensitive to "Temperature, μ_B ..."

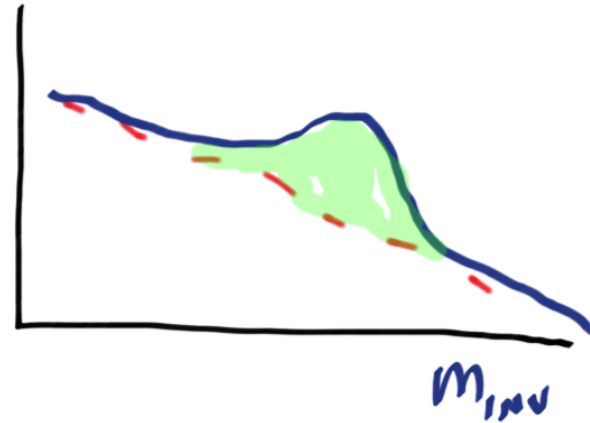
D^* - ATTRACTIVE STRONG F.S.I. b/t A, B
- included in $\Phi_{AB}(r, z)$

$R_{AB}(z)$ - "real" pairs

$M_{AB}(z)$ - "mixed" pairs

HADROCHEMISTRY

• D^* IS A PARTICLE



$$N_{D^*} = \int dM_{inv} [R_{AB}(M_{inv}) - N M_{AB}(M_{inv})]$$

$$N_A = \int d^3 p_A \frac{d^3 N}{d^3 p_A} = V \frac{g_A}{2\pi^2} e^{\mu_A/T} \int_0^\infty p^2 dp e^{-E_A(p)/T}$$

• PARTICLE RATIOS

- sensitive to T, μ^*

- insensitive to V

$$\bullet \mu_A = B_A \mu_B + S_A \mu_S + I_{3A} \mu_I + \dots$$

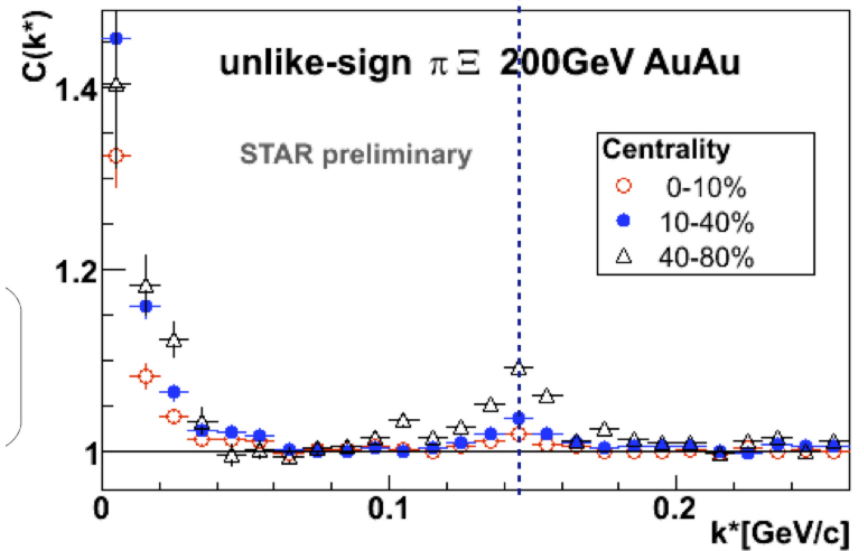
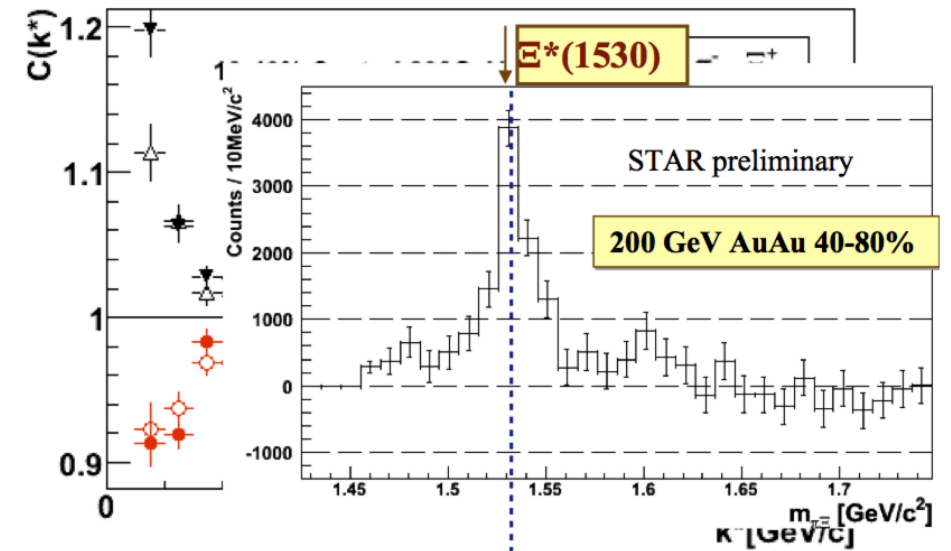
$$\bullet \mu_{D^*} = \mu_A + \mu_B$$

STAR COLLAB ~ 2005
P. CHALOUPKA (Prague)

$\Xi - \pi$ femtoscopy

• Coulomb

• $\Xi^* \rightarrow \Xi + \pi$ for
opposite sign

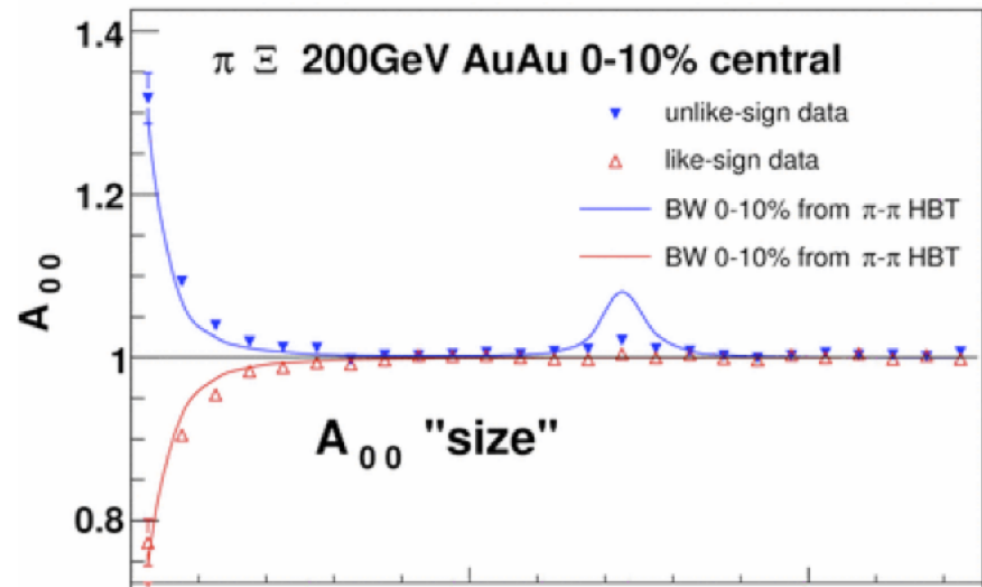


STAR COLLAB ~ 2005
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$\Xi - \pi$ femtoscopy

- Coulomb

- $\Xi^* \rightarrow \Xi + \pi$ for opposite sign



Calculation w/ Koonin-Pratt + STAR BW source:

- ✓ Coulomb
- ✗ Strong F.S.I. (Ξ^*)

Space-time source

STAR COLLAB ~ 2005

P. CHALOUPKA (Prague)

$\Xi - \pi$ femtoscopy

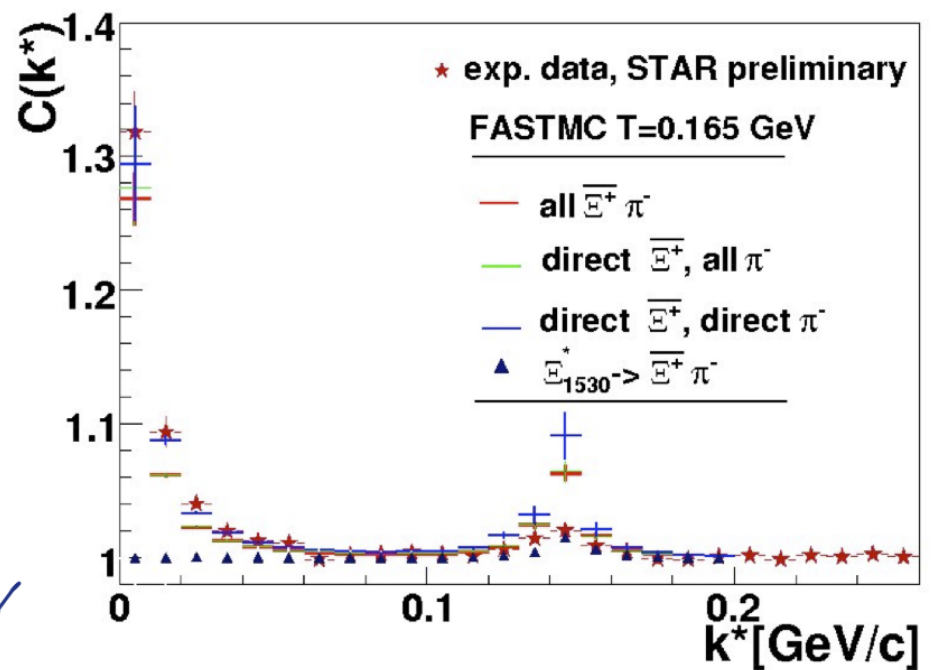
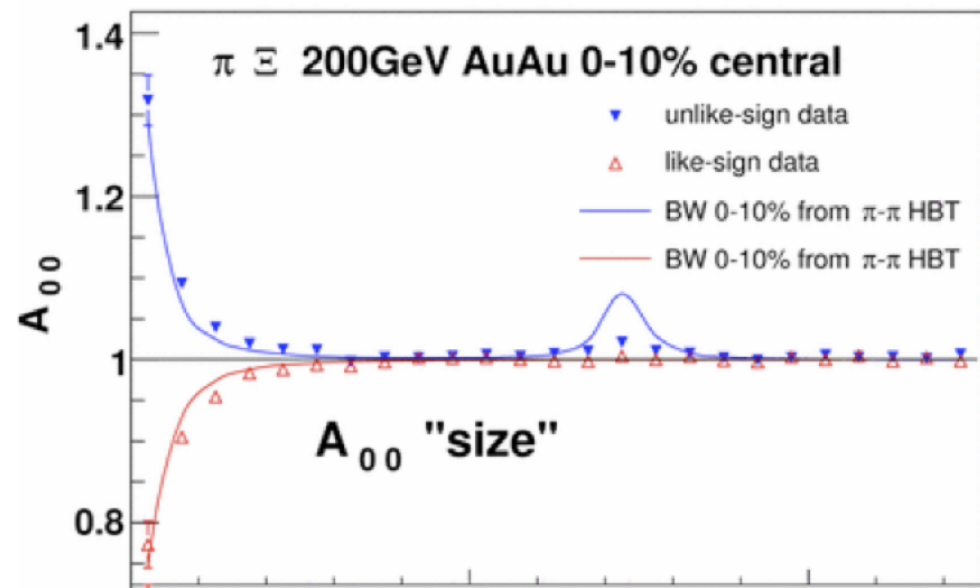
- Coulomb

- $\Xi^* \rightarrow \Xi + \pi$ for opposite sign

- Treatment as strong FSI with Koonin-Pratt
OVER-PREDICTS enhancement

- Neglecting $\Xi - \pi$ FSI, "directly-produced" Ξ^* decay
UNDER-PREDICTS enhancement

→ try a little of each?

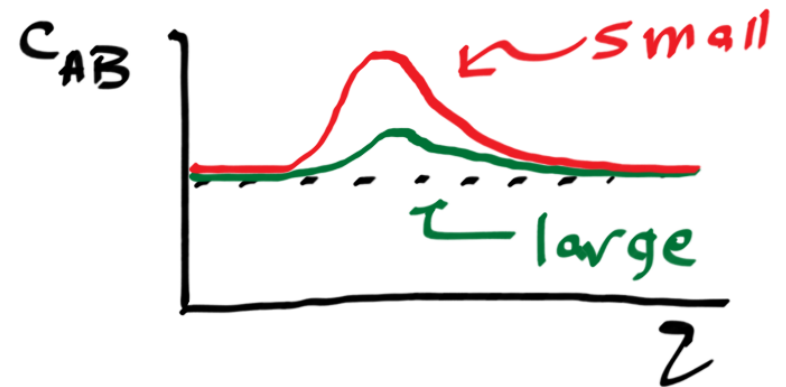


FOR A FIXED N_A, N_B , A LARGE N_{D^*} MEANS...

FEMTOSCOPIST:

SMALL SOURCE!

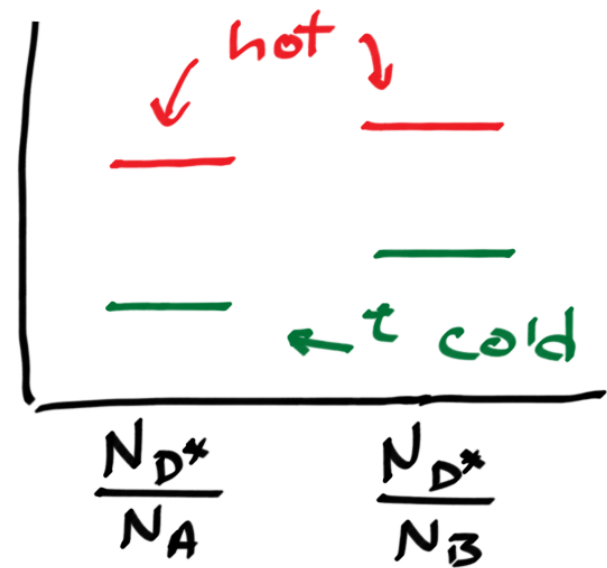
(Hot or cold)



HADROCHEMIST:

HOT SOURCE!

(large or small)

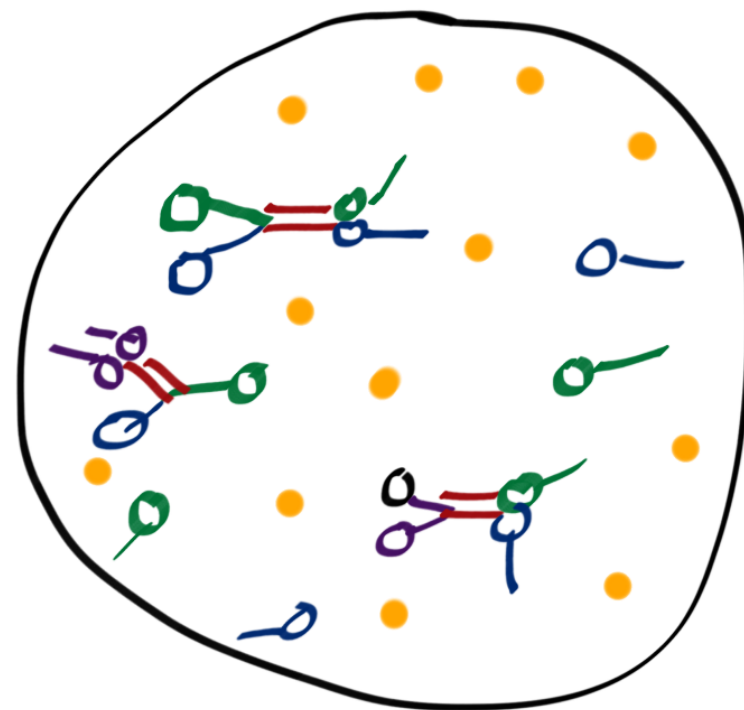


...clearly, we need to be more specific...

FIRST: SIMPLEST THERMAL MODEL ASSUMPTIONS

- SOURCE IN THERMOCHEMICAL EQUILIB $A+B \rightleftharpoons D^*$
- Volume V , TEMP T , C.P. μ^*
- UNIFORM DENSITY
- NO FLOW
- INSTANT FREEZEOUT

CONSIDER D^*
A PARTICLE



○ A = D^*
○ B ○, ● other

FEMTOSCOPIC

$$C_{AB}(M_{D^*}) = \frac{\frac{d^6 N_{AB}}{d^3 P_A d^3 P_B}}{\frac{d^3 N_A}{d^3 P_A} \cdot \frac{d^3 N_B}{d^3 P_B}} = \frac{R_{AB}(M_{D^*})}{M_{AB}(M_{D^*})} = \frac{D^* \text{ contrib} + M_{AB}(M_B)}{M_{AB}(M_{D^*})}$$

FEMTOSCOPIC

$$C_{AB}(M_{D^*}) = \frac{\frac{d^6 N_{AB}}{d^3 P_A d^3 P_B}}{\frac{d^3 N_A}{d^3 P_A} \cdot \frac{d^3 N_B}{d^3 P_B}} = \frac{R_{AB}(M_{D^*})}{M_{AB}(M_{D^*})} = \frac{D^* \text{ contrib} + M_{AB}(M_B)}{M_{AB}(M_{D^*})}$$

$$= \frac{v e^{\mu_{D^*}/T} \frac{g_{D^*}}{2\pi^2} \int_0^\infty p^2 e^{-E_{D^*}(p)/T} f(p_{D^*} \rightarrow \vec{P}_A, \vec{P}_B)}{(v e^{\mu_A/T} \frac{g_A}{2\pi^2} e^{-E_A(p_A)/T}) (v e^{\mu_B/T} \frac{g_B}{2\pi^2} e^{-E_B(p_B)/T})} + 1$$

FEMTOSCOPIC

$$C_{AB}(M_{D^*}) = \frac{\frac{d^6 N_{AB}}{d^3 P_A d^3 P_B}}{\frac{d^3 N_A}{d^3 P_A} \cdot \frac{d^3 N_B}{d^3 P_B}} = \frac{R_{AB}(M_{D^*})}{M_{AB}(M_{D^*})} = \frac{D^* \text{ contrib} + M_{AB}(M_B)}{M_{AB}(M_{D^*})}$$

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$$(\mu_{D^*} = \mu_A + \mu_B)$$

FEMTOSCOPIC

$$C_{AB}(M_{D^*}) = \frac{\frac{d^6 N_{AB}}{d^3 P_A d^3 P_B}}{\frac{d^3 N_A}{d^3 P_A} \cdot \frac{d^3 N_B}{d^3 P_B}} = \frac{R_{AB}(M_{D^*})}{M_{AB}(M_{D^*})} = \frac{D^* \text{ contrib} + M_{AB}(M_B)}{M_{AB}(M_{D^*})}$$

$$= \frac{V e^{\mu_{D^*}/T} \frac{g_{D^*}}{2\pi^2} \int_0^\infty p^2 e^{-E_{D^*}(p)/T} f(p_{D^*} \rightarrow \vec{p}_A, \vec{p}_B)}{(V e^{\mu_A/T} \frac{g_A}{2\pi^2} e^{-E_A(p_A)/T}) (V e^{\mu_B/T} \frac{g_B}{2\pi^2} e^{-E_B(p_B)/T})} + 1$$

$$(\mu_{D^*} = \mu_A + \mu_B)$$

$$(E_{D^*} = E_A + E_B)$$

μ, T dependence gone ✓

FEMTOSCOPIC

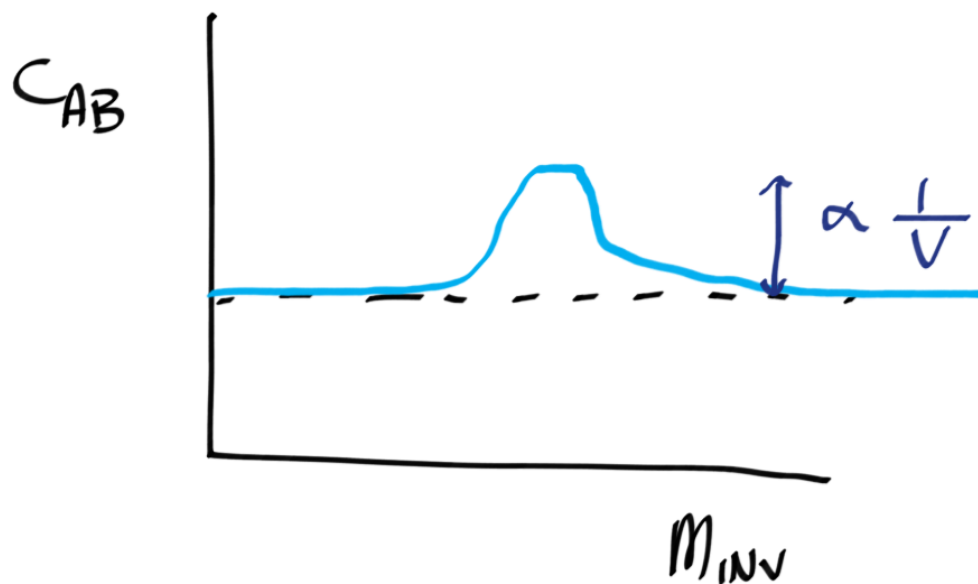
$$C_{AB}(M_{D^*}) = \frac{\frac{d^6 N_{AB}}{d^3 P_A d^3 P_B}}{\frac{d^3 N_A}{d^3 P_A} \cdot \frac{d^3 N_B}{d^3 P_B}} = \frac{R_{AB}(M_{D^*})}{M_{AB}(M_{D^*})} = \frac{D^* \text{ contrib} + M_{AB}(M_B)}{M_{AB}(M_{D^*})}$$

$$= \frac{1}{V} \frac{2\pi^2 g_{D^*}}{g_A g_B} \int_0^\infty P_{D^*}^2 dP_{D^*} f(\vec{P}_{D^*} \rightarrow \vec{P}_A \vec{P}_B) + 1$$


FEMTOSCOPIC

$$C_{AB}(M_{D^*}) = \frac{\frac{d^6 N_{AB}}{d^3 P_A d^3 P_B}}{\frac{d^3 N_A}{d^3 P_A} \cdot \frac{d^3 N_B}{d^3 P_B}} = \frac{R_{AB}(M_{D^*})}{M_{AB}(M_{D^*})} = \frac{D^* \text{ contrib} + M_{AB}(M_{D^*})}{M_{AB}(M_{D^*})}$$

$$= \frac{1}{V} \underbrace{\frac{2\pi^2 g_{D^*}}{g_A g_B} \int_0^\infty P_{D^*}^2 dP_{D^*} f(\vec{P}_{D^*} \rightarrow \vec{P}_A \vec{P}_B)}_{\text{resonance-dependent constant}} + 1$$



HADROCHEMICAL

non-rel. approx. 

$$\frac{N_{D^*}}{N_A} = \frac{v e^{(\mu_{D^*} - m_{D^*})/T} \frac{g_{D^*}}{2\pi^2} \sqrt{\frac{\pi}{2}} (m_{D^*} T)^{3/2}}{v e^{(\mu_A - m_A)/T} \frac{g_A}{2\pi^2} \sqrt{\frac{\pi}{2}} (m_A T)^{3/2}}$$

$$= \frac{g_{D^*}}{g_A} \left(\frac{m_{D^*}}{m_A}\right)^{3/2} e^{(\mu_{D^*} - \mu_A - (m_{D^*} - m_A))/T} \quad \checkmark$$

"Femto-like": $\frac{N_{D^*}}{N_A N_B} = \frac{g_{D^*}}{g_A g_B} \left(\frac{2\pi m_{D^*}}{m_A m_B T}\right)^{3/2} \frac{1}{v} e^{-Q/T}$

$\bullet \sim \frac{1}{v} \quad \checkmark$

\bullet residual T-dep (b/c integrate over all m^*)

$$Q = m_{D^*} - m_A - m_B$$

LOW-ENERGY EXAMPLES

Coulomb + ${}^6\text{Li}^*$ in $d-\alpha$

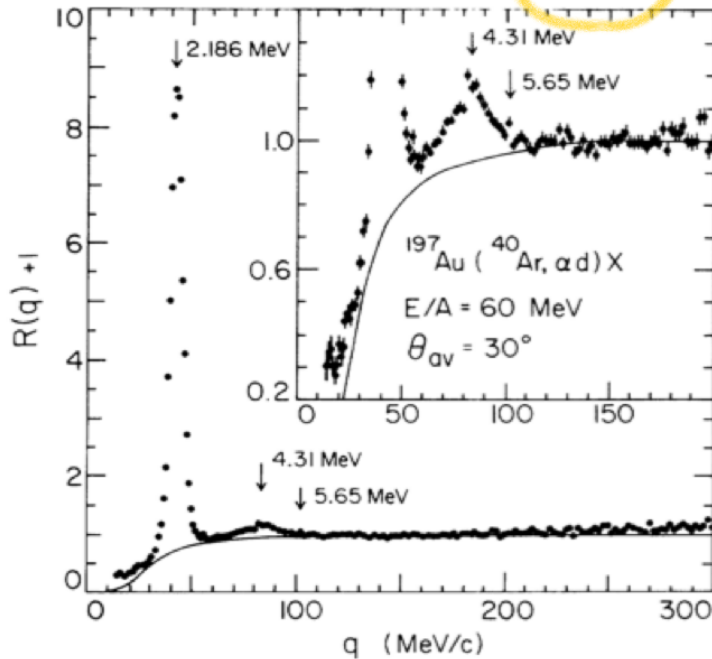


FIG. 8. $d-\alpha$ correlation function. Excited states in ${}^6\text{Li}$ are indicated. The curve represents the background correlation function.

${}^{40}\text{Ar} + {}^{197}\text{Au}$ $E = 60 \text{ A MeV}$

J. Pochodzalla et al

PRC 35 1695 (1987)

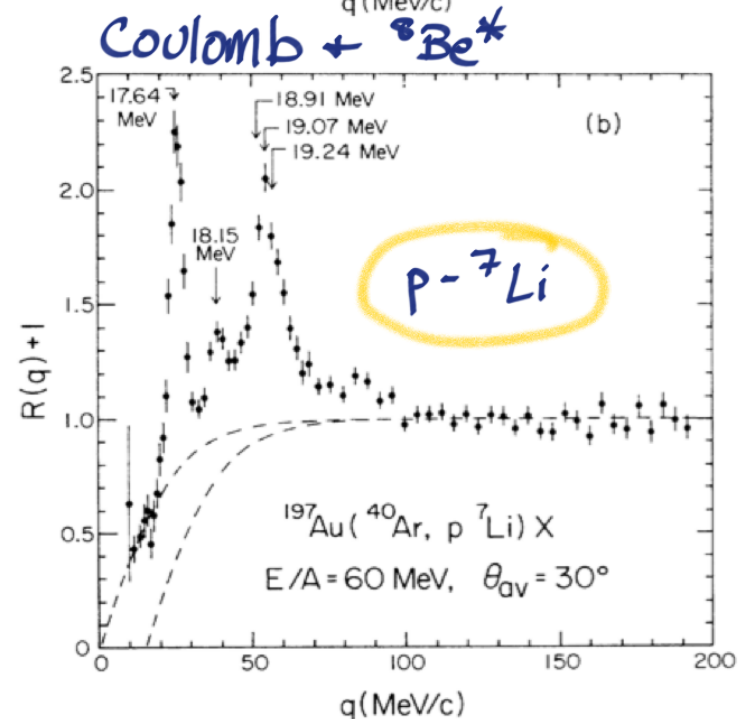
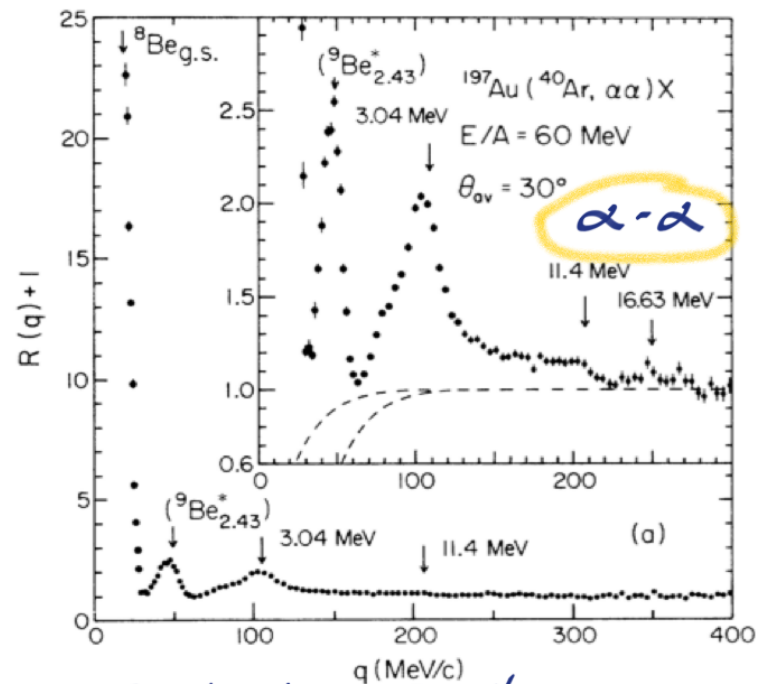


FIG. 11. (a) $\alpha-\alpha$ and (b) $p-{}^7\text{Li}$ correlation functions. Excited states in ${}^8\text{Be}$ are indicated. The dashed lines are extreme bounds for the background correlation function.

TWO-PARTICLE CORRELATION FUNCTIONS IN THE THERMAL MODEL AND NUCLEAR INTERFEROMETRY DESCRIPTIONS

B.K. JENNINGS, D.H. BOAL, J.C. SKILLCOCK

PHYS REV C33, 1303 (1986)

"The equivalence of thermal model and conventional zero-lifetime HBT descriptions is demonstrated"

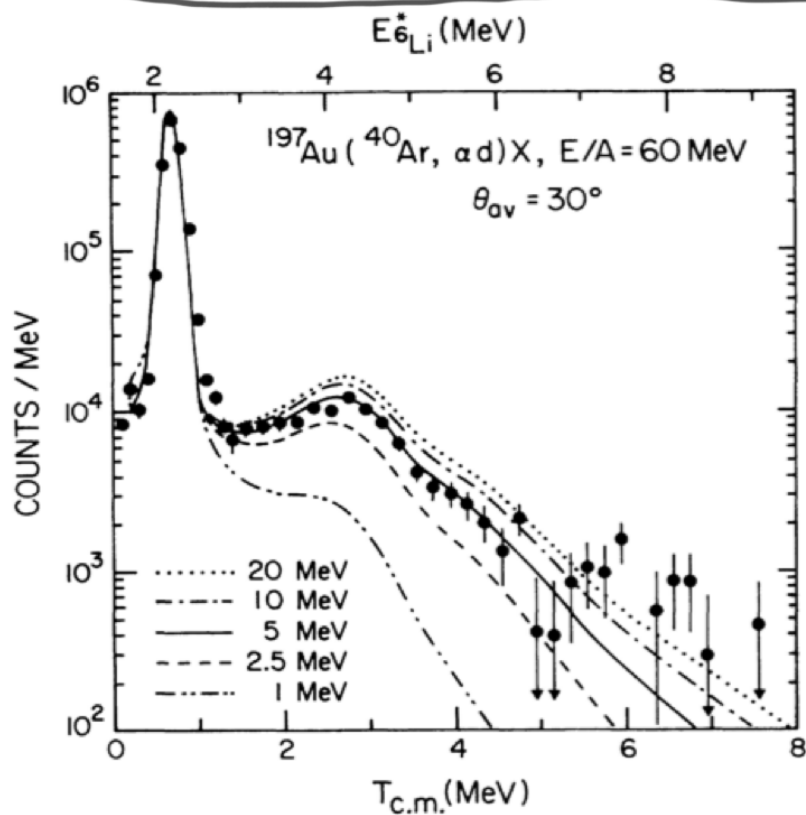


FIG. 20. Energy spectrum resulting from the decay of particle-unstable states in ${}^6\text{Li}$. The curves correspond to thermal distributions, Eq. (25), with $T=1, 2.5, 5, 10,$ and 20 MeV.

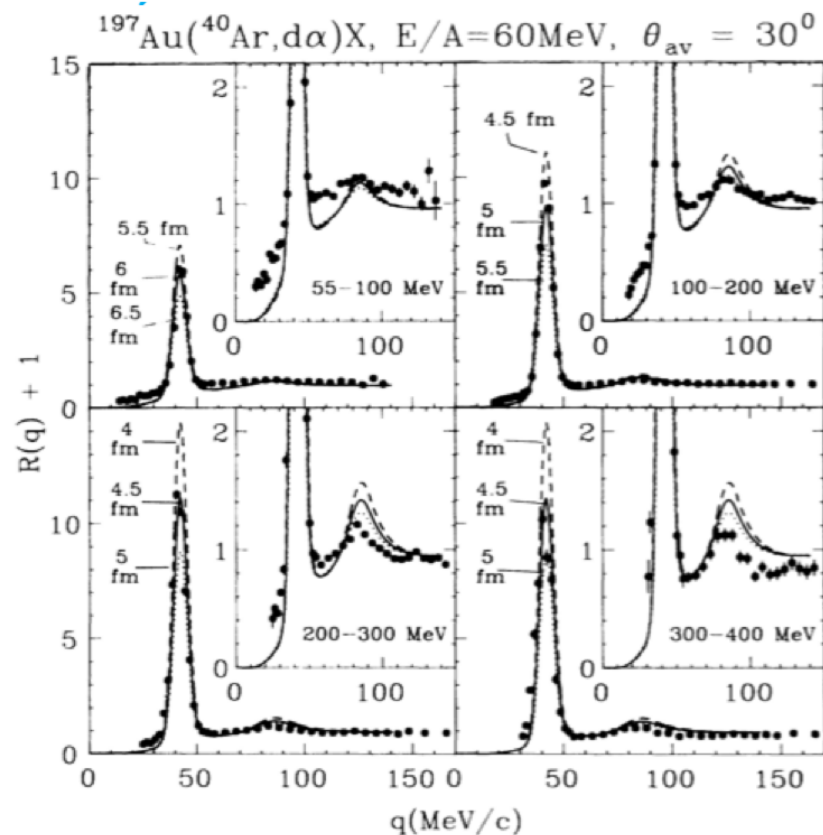


FIG. 16. $d-\alpha$ correlation functions for different constraints on the sum energy, $E_1 + E_2$. The curves show calculations with the final state interaction model, Eq. (18).

↑
KONIN-PRATT

$d-\alpha$ data from Pochodzalla (1987)

Statistical Emission of ^2He from Highly Excited Nuclear Systemsp-p femtoscopy

FOR REACTIONS FORMING
 COMPOUND NUCLEUS
" ^2He emission" ✓

M. Bernstein, ... M.B. Tsang, et al.

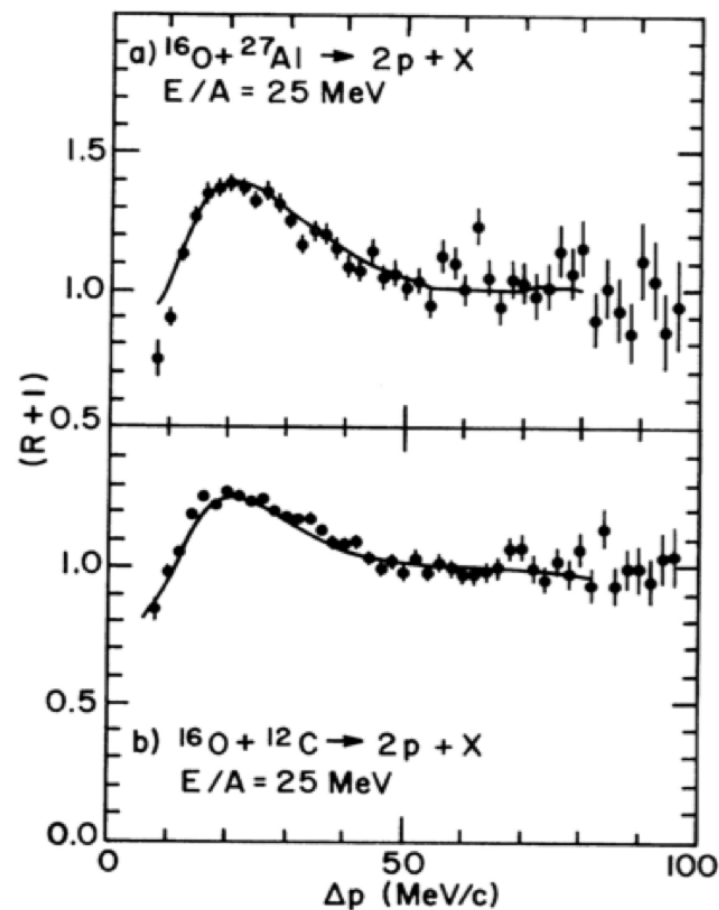


FIG. 2. (a) The experimental correlation function $1 + R(\Delta p)$ plotted as a function of the relative momentum for protons produced in reactions on the ^{27}Al target. The theoretical correlation function is drawn as the solid line. (b) Corresponding experimental and theoretical correlation functions for reactions on the ^{12}C target.

Statistical Emission of ^2He from Highly Excited Nuclear Systemsp-p femtoscopy

FOR REACTIONS FORMING
COMPOUND NUCLEUS
" ^2He emission" ✓

REACTIONS **NOT** RESULTING
IN EQUILIBRATED SOURCE

(forward $^{16}\text{O} + ^{197}\text{Au}$)

→ thermal treatment
does not work

⇒ requires "Koonin
formalism"

M. Bernstein, ... M.B. Tsang, et al.

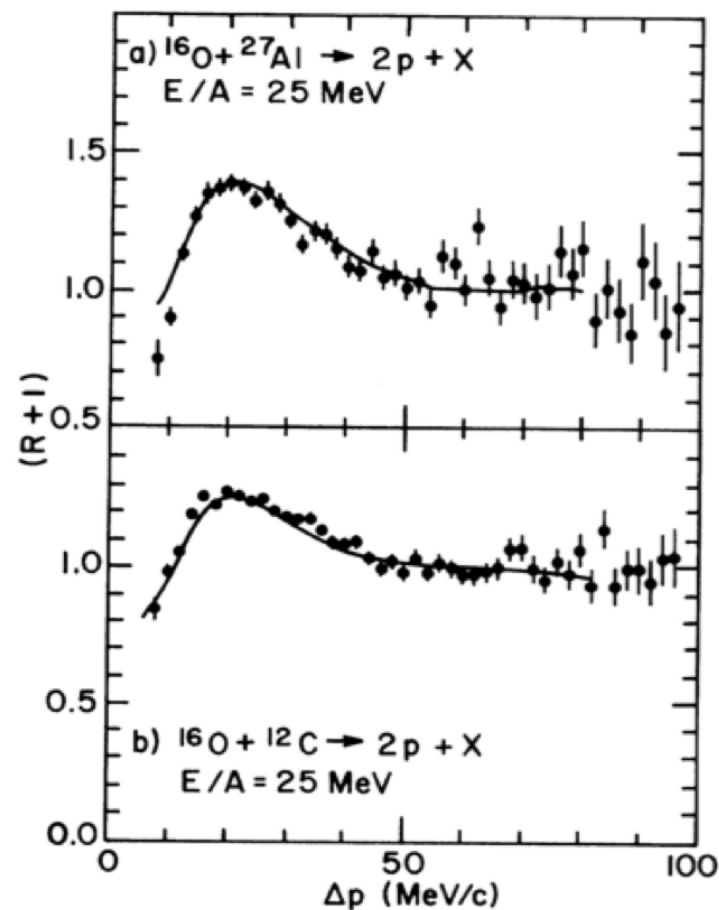
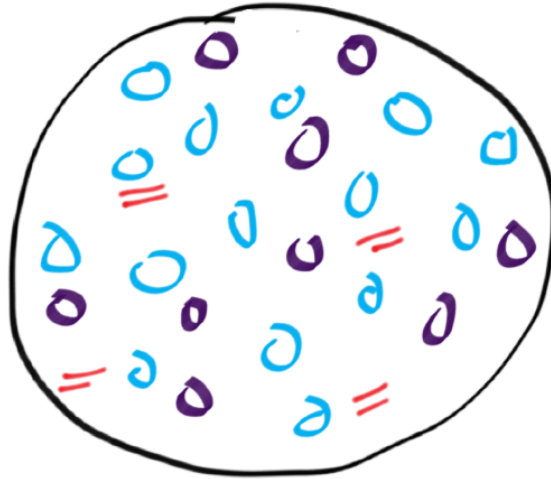
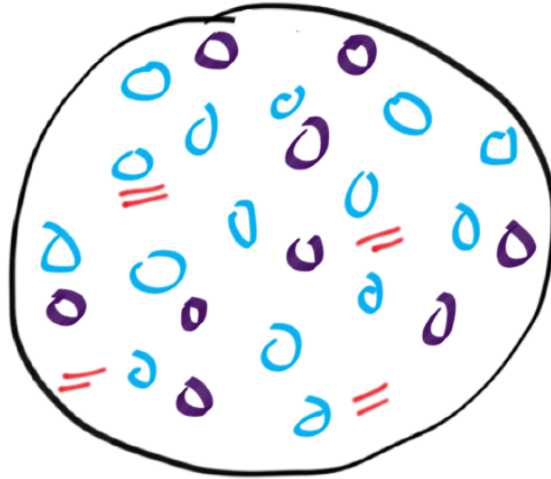


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INSTANT DECAY FROM EQUILIBRIUM: TWO TREATMENTS:

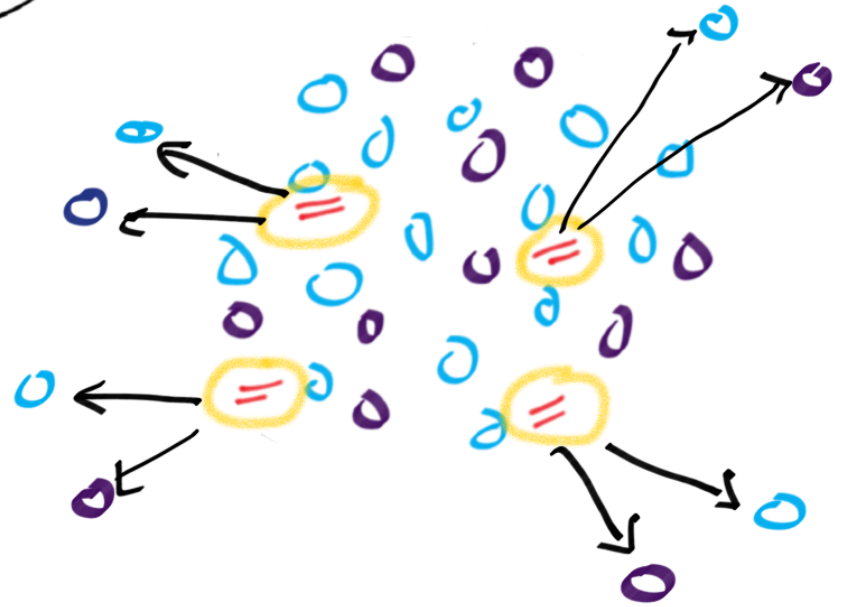
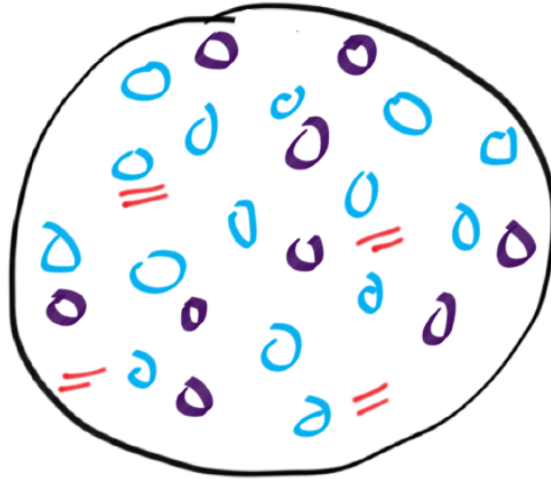


INSTANT DECAY FROM EQUILIBRIUM: TWO TREATMENTS:



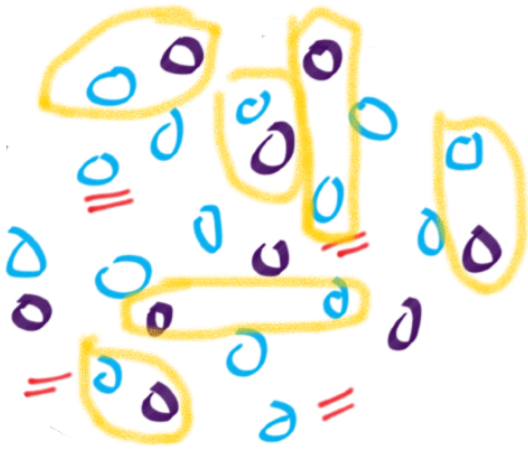
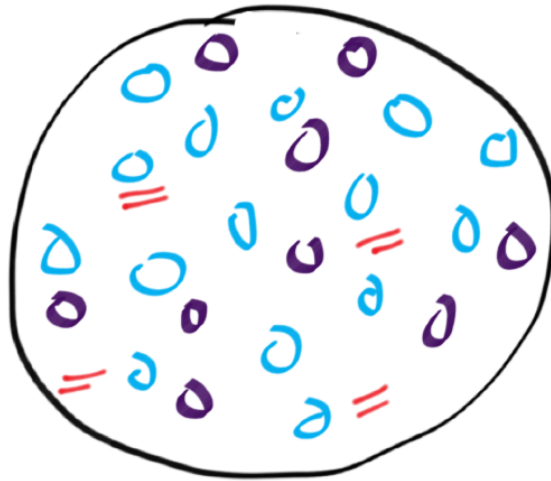
- pairwise $S(r)$
- $C(k^*) = \int d^3r S(r) |\psi(k^*, r)|^2$

INSTANT DECAY FROM EQUILIBRIUM: TWO TREATMENTS:

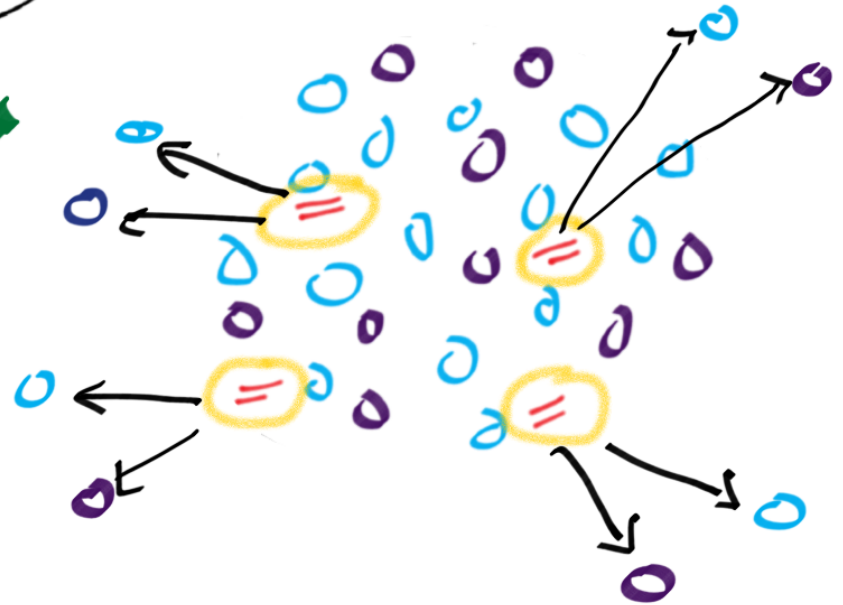


- pairwise $S(r)$
- $C(k^*) = \int d^3r S(r) |\psi(k^*, r)|^2$

INSTANT DECAY FROM EQUILIBRIUM: TWO TREATMENTS:



PICK
ONE

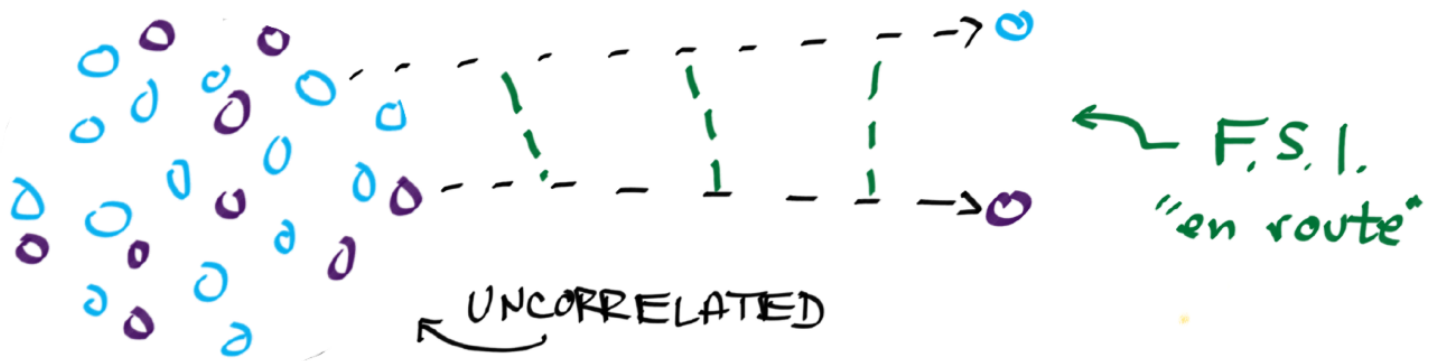


- pairwise $S(r)$

- $C(k^*) = \int d^3r S(r) |\psi(k^*, r)|^2$

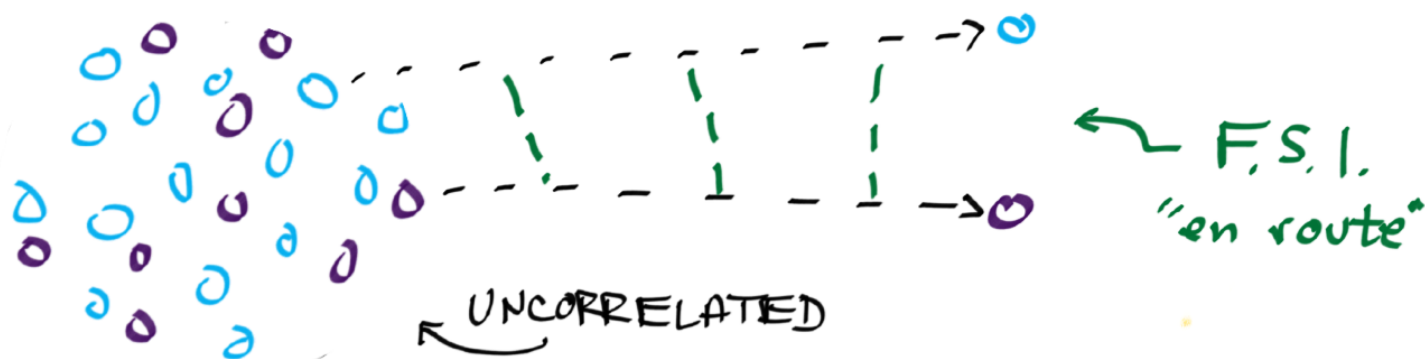


FEMTOSCOPIST'S COMMON PICTURE



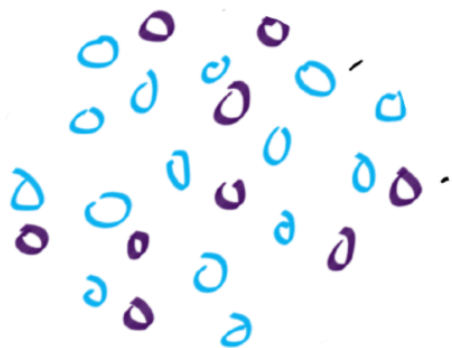
$$C(k^*) = \int d^3r S_{AB}(r) |\psi(k^*, r)|^2$$

FEMTOSCOPIST'S COMMON PICTURE



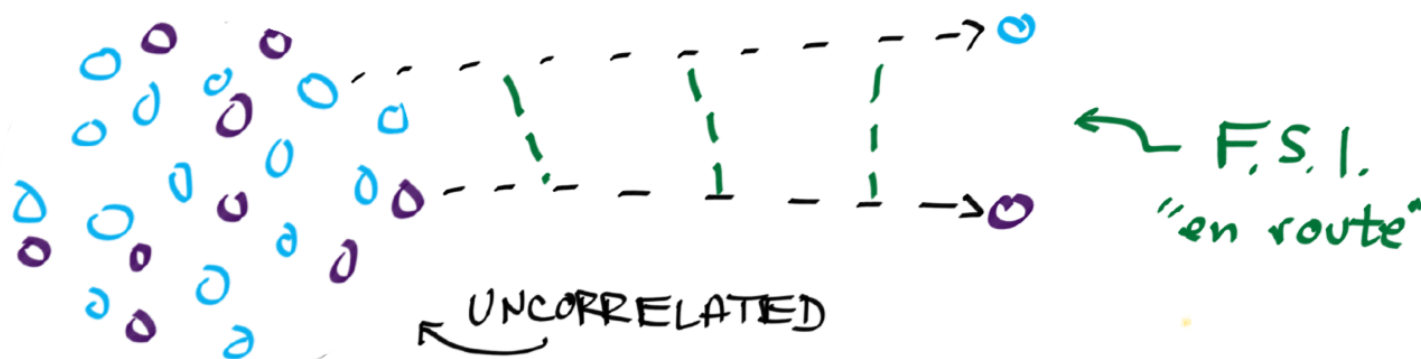
$$C(k^*) = \int d^3r S_{AB}(r) |\psi(k^*, r)|^2$$

CORRELATED



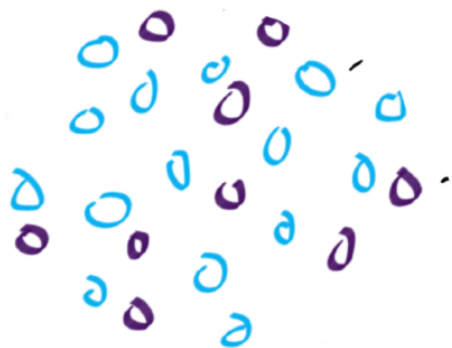
$|\psi(r, k^*)|^2$ quantifies
modified density
of states (including D^*)
 \Rightarrow correlation

EQUIVALENT? ONLY FOR EQUILIBRIUM?



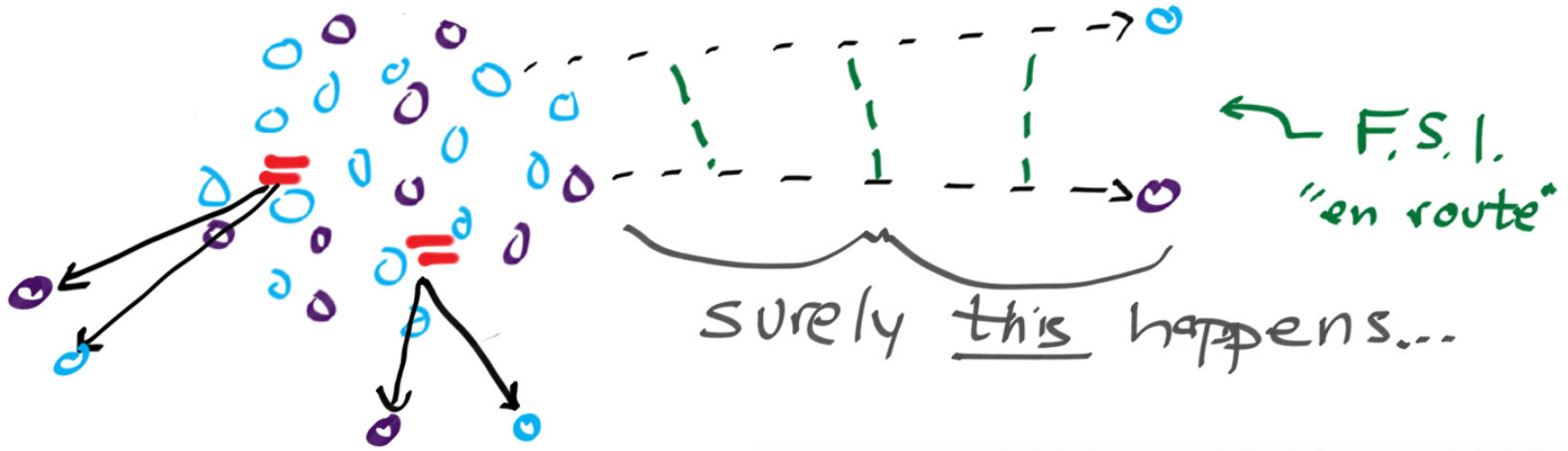
$$C(k^*) = \int d^3r S_{AB}(r) |\psi(k^*, r)|^2$$

CORRELATED

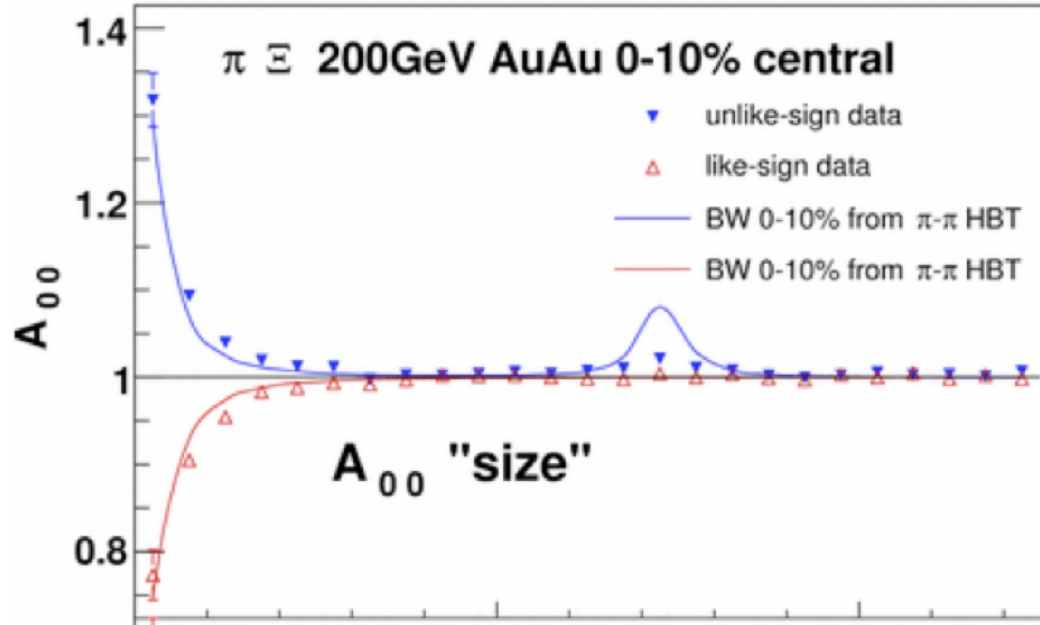


$|\psi(r, k^*)|^2$ quantifies
modified density
of states (including D^*)
 \Rightarrow correlation

IF THERE ARE "EXTRA" D^* 'S (Ξ^* 'S)...



Doesn't that worsen the discrepancy?



THAT'S ALL FOR NOW - THANKS

