



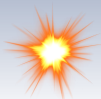
The influence of initial conditions on the final observables for heavy-ion collisions at RHIC energies

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in collaboration with:

E. Bratkovskaya, W. Cassing and J. Aichelin

WPCF2013, Acireale (Italy)



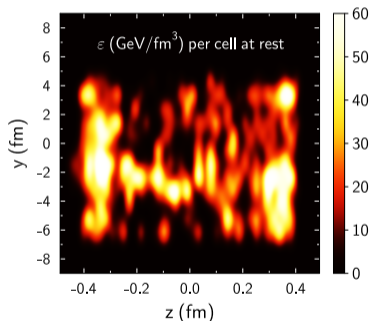
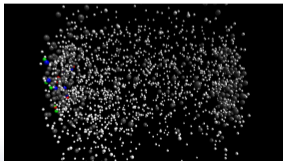
Motivations

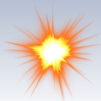
We want a **microscopic description** of partonic phase and phase transition (cross over).

→ Transport codes : **PHSD** & **RSP**

What is the influence of **initial conditions** using two different approaches on the observables for RHIC energies ?

→ Effective models : **DQPM** & **NJL**





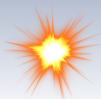
The Parton Hadron String Dynamics

Features:

- Description of heavy-ion collisions,
- **Non-equilibrium** approach,
- **Strings** formation and decay to pre-hadrons,
- Pre-hadrons fragmentation into partons,
- Dynamical Quasi-Particle Model (**DQPM**) for describing partons masses and widths,
- **Off-shell transport** of hadrons and partons with mean fields and scattering,
- Dynamical **hadronization** with cross over.



PRC 78, 034919 (2008)
NPA 831, 215 (2009)
EPJ ST 168, 3 (2009)
NPA 856, 162 (2011)



The Dynamical Quasi-Particle Model

Quasi-partons:

Masses:

$$M_g^2(T, \mu_q) = \frac{g^2}{6} \left((N_c + \frac{N_f}{2}) T^2 + \frac{3}{2} \sum_q \frac{\mu_q^2}{\pi^2} \right)$$

$$M_{q/\bar{q}}^2(T, \mu_q) = \frac{N_c^2 - 1}{8N_c} g^2 \left(T^2 + \sum_q \frac{\mu_q^2}{\pi^2} \right)$$

Widths:

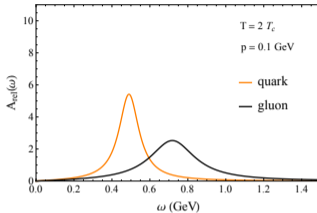
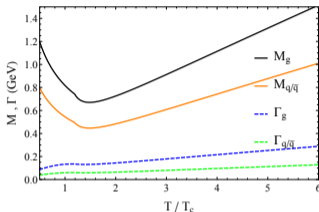
$$\Gamma_g(T) = \frac{1}{3} N_c \frac{g^2 T}{8\pi} \ln \left(\frac{2c}{g^2} + 1 \right),$$

$$\Gamma_{q/\bar{q}}(T) = \frac{1}{3} \frac{N_c^2 - 1}{2N_c} \frac{g^2 T}{8\pi} \ln \left(\frac{2c}{g^2} + 1 \right),$$

Coupling constant:

$$g^2(T/T_c) = \frac{48\pi^2}{(11N_c - 2N_f) \ln[\lambda^2(T/T_c - T_s/T_c)^2]}$$

Based on EPJ ST 168, 3 (2009)



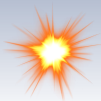
Off-shellness:

Breit-Wigner spectral function:

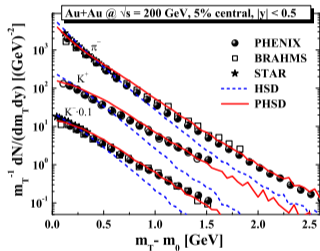
$$A(\omega, \mathbf{p}) = \frac{\Gamma}{E} \left(\frac{1}{(\omega - E)^2 + \Gamma^2} - \frac{1}{(\omega + E)^2 + \Gamma^2} \right)$$

with $E^2 = \mathbf{p}^2 + M^2 - \Gamma^2$ and

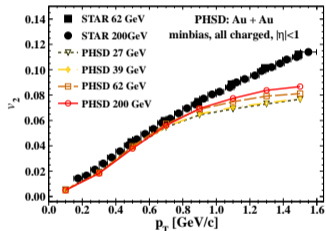
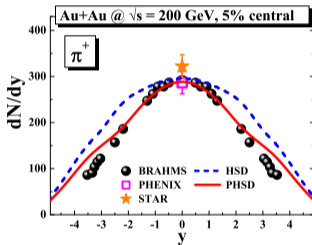
$$\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega A(\omega, \mathbf{p}) = \int_0^{\infty} \frac{d\omega}{2\pi} 2\omega A(\omega, \mathbf{p}) = 1$$



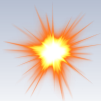
Some results



(NPA 856, 162 (2011))



(PRC 85, 044922 (2012))

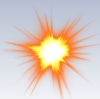


A new code on the market

Relativistic quantum molecular dynamics for
Strongly interacting matter with
Phase transition or crossover

Features:

- **C++** code ~5000 lines (modern and fully parallelizable),
- New relativistic quantum molecular dynamics (which is **causal** and **conserves energy**),
- **(P)NJL model** based dynamics with q and \bar{q} degrees of freedom (no gluons) and pseudoscalar mesons (π, K, η),
- All masses m_i and cross sections $\sigma_{2 \rightarrow 2}$ at finite (T, μ) for dynamical **cross over** or **phase transition**,
- **Local mean field** and relativistic quantum collisions (for fluctuations),
- Different possible initial conditions: box, toy model for heavy ion collisions, external input (e.g. PHSD).



The Nambu-Jona-Lasinio model

Lagrangian:

$$\begin{aligned} \mathcal{L}_{NJL} = & \bar{\psi} (i\partial - m_0) \psi \\ & + G \sum_{a=0}^8 [(\bar{\psi} \lambda^a \psi)^2 + (\bar{\psi} i\gamma_5 \lambda^a \psi)^2] \\ & - K [\det \bar{\psi} (1 - \gamma_5) \psi + \det \bar{\psi} (1 + \gamma_5) \psi] \end{aligned}$$

Quark mass:

$$m_i = m_{0i} - 4G \langle \bar{\psi}_i \psi_i \rangle + 2K \langle \bar{\psi}_j \psi_j \rangle \langle \bar{\psi}_k \psi_k \rangle$$

Chiral condensate:

$$\langle \bar{\psi}_i \psi_i \rangle = -2N_c \int_0^\Lambda \frac{d^3p}{(2\pi)^3} \frac{m_i}{E_{ip}} [1 - f_q - f_{\bar{q}}]$$

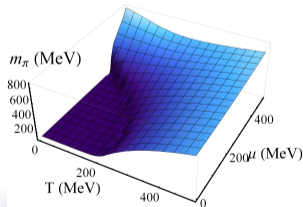
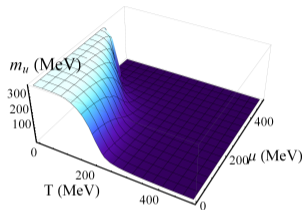
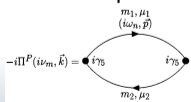
Based on [PRC 87, 034912 \(2013\)](#)

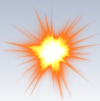
- Chiral model for q/\bar{q} ,
- QCD symmetries,
- hadrons construction,
- Finite (T, μ) .

Meson mass:

$$\frac{-ig^2 \pi q \bar{q}}{k^2 - M^2} = \frac{2iG}{1 - 2G \Pi(k)}$$

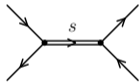
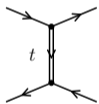
Polar. loop:



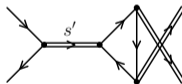
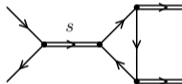
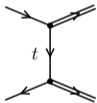


NJL cross sections

Leading order processes (NJL \rightarrow no gluons !):

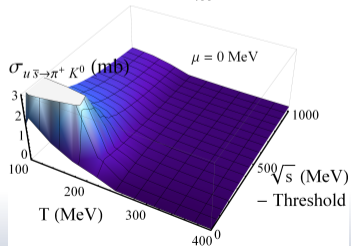
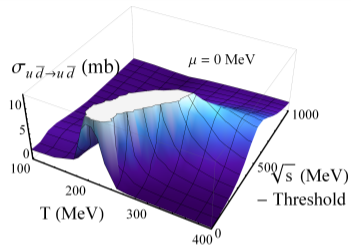


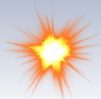
$q \bar{q} \rightarrow q \bar{q}$



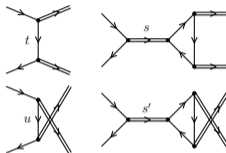
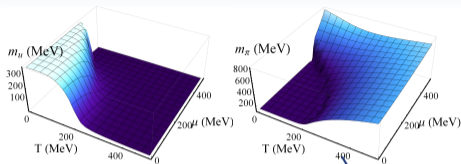
$q \bar{q} \rightarrow M M$

Based on PRD 51, 3728 (1995)





New relativistic dynamics



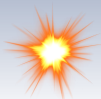
- $qq \leftrightarrow qq$,
- $q\bar{q} \leftrightarrow q\bar{q}$,
- $q\bar{q} \leftrightarrow MM$,
- $M \rightarrow q\bar{q}$.

$$\frac{dq_i^\mu}{d\tau} = \frac{p_i^\mu}{E_i}$$

$$\frac{dp_i^\mu}{d\tau} = - \sum_{k=1}^N \frac{m_k}{E_k} \frac{\partial m_k}{\partial q_{i,\mu}} + \langle \text{collisions} \rangle$$

NJL masses and cross sections enter in the propagation equations.
Wigner (Gaussian) distribution in phase-space for particles $f(\vec{q}_i, \vec{p}_i, \tau)$.

(PRC 87, 034912 (2013))



Local mean field

Local densities :

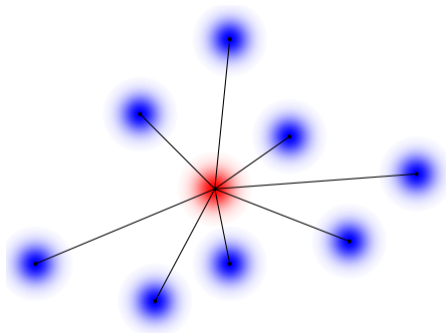
3D-probability that i feels j

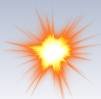
$$R_{ij} = \left(\frac{1}{\sqrt{\pi}L} \right)^3 \exp\left(-\frac{\Delta r_{ij}^2}{L^2} \right)$$

we define

$$\rho_{Fi} = \int (f_q + f_{\bar{q}}) d^3p \equiv \sum_{i \neq j} R_{ij}$$

$$\rho_{Bi} = \int (f_q - f_{\bar{q}}) d^3p \equiv \sum_{i \neq j} R_{ij} \text{Sign}(j)$$





Local mean field

Local densities :

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Local potentials :

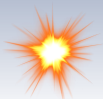
Thermodynamics gives :

$$T_i = (\hbar c) \left(\frac{\pi^2}{g\kappa} \right)^{1/3} \rho_{Fi}^{1/3} \quad (\text{for } \mu \rightarrow 0)$$

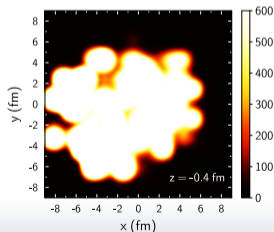
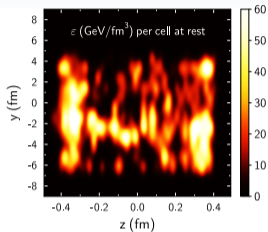
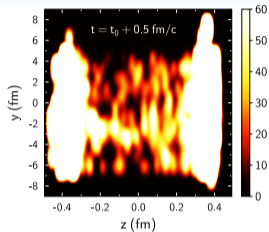
$$\mu_i = (\hbar c) \left(\frac{6\pi^2}{g\kappa} \right)^{1/3} \rho_{Bi}^{1/3} \quad (\text{for } T \rightarrow 0)$$

$\kappa \approx 0.9$ comes from Fermi integral.

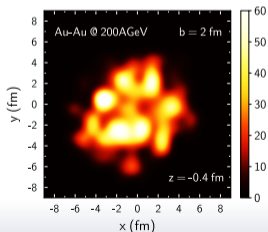
In equilibrium (box + large N) \rightarrow not sensitive to L anymore !



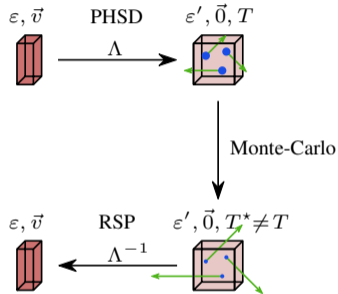
PHSD initial conditions



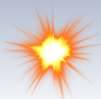
global system



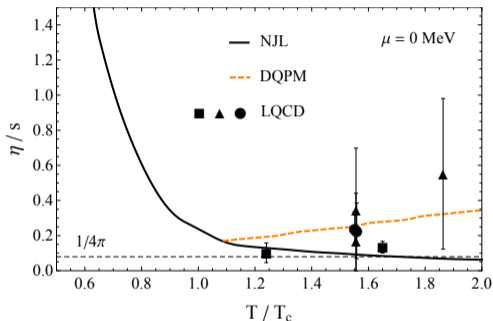
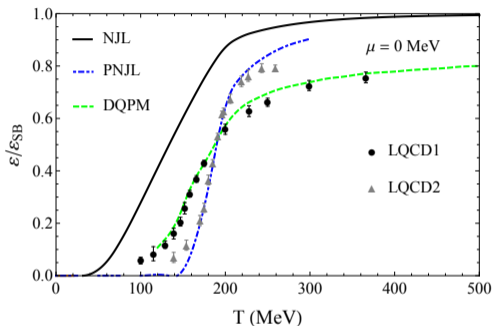
cells at rest



DQPM partons: heavy
NJL quarks: light



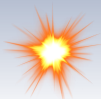
Equations of state and viscosity



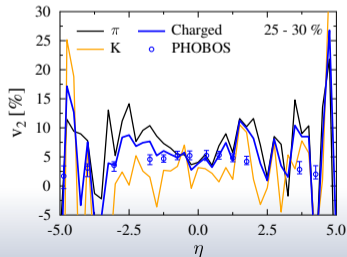
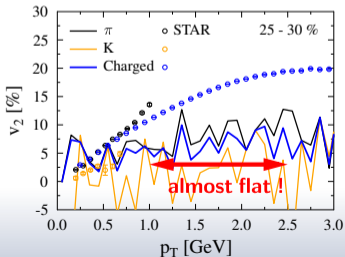
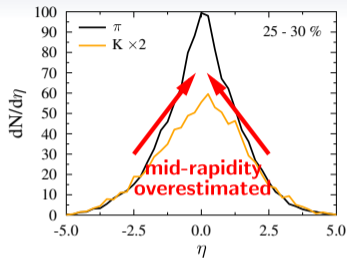
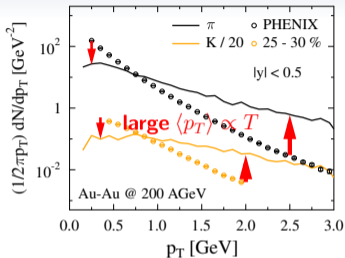
Many models = many equations of state
 ($T_c \text{ NJL} \neq T_c \text{ DQPM}$ and $\epsilon_{SB} \text{ NJL} \neq \epsilon_{SB} \text{ DQPM}$)

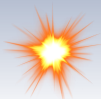
η/s in NJL is close to LQCD data !

(PRC 88, 045204 (2013))

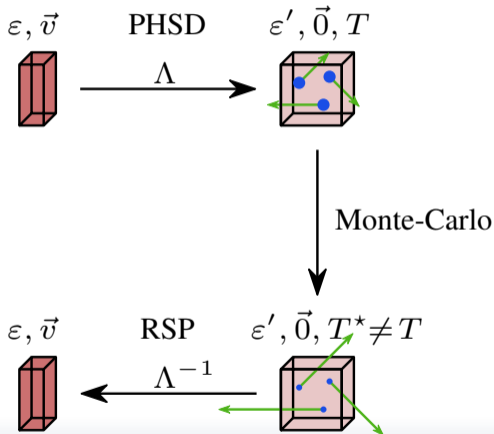


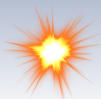
First results



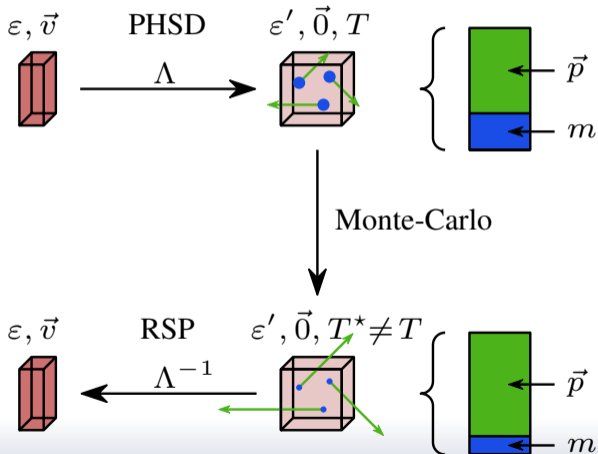


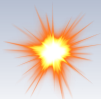
Out of equilibrium conversion



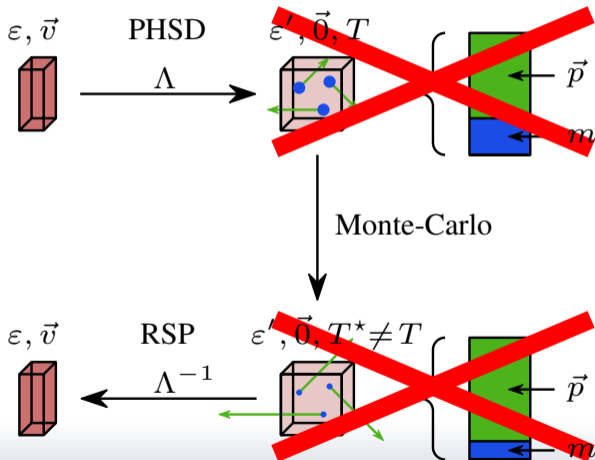


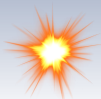
Out of equilibrium conversion



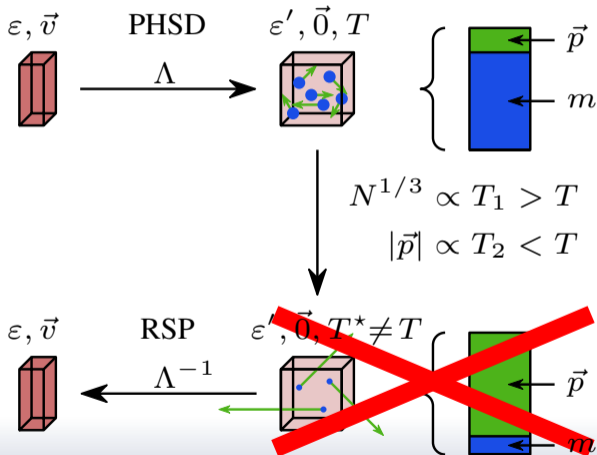


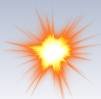
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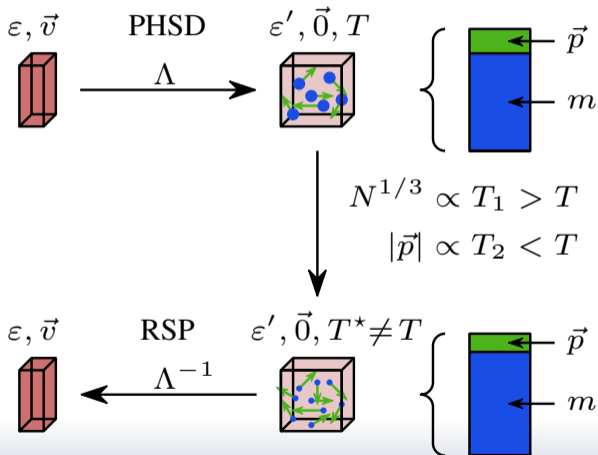


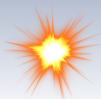
Out of equilibrium conversion



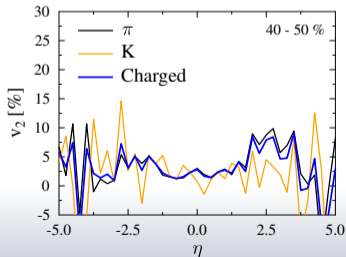
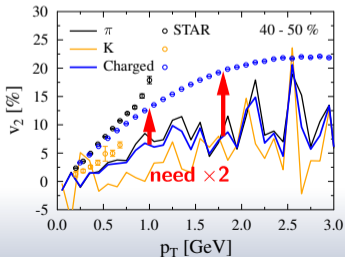
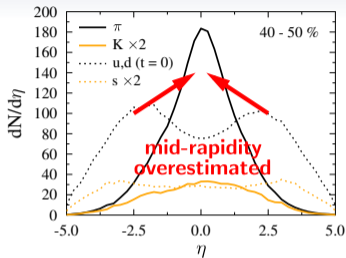
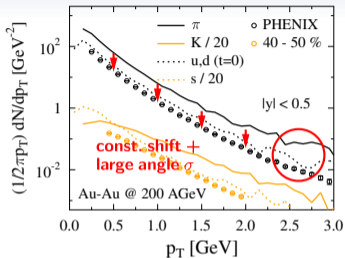


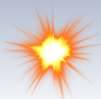
Out of equilibrium conversion





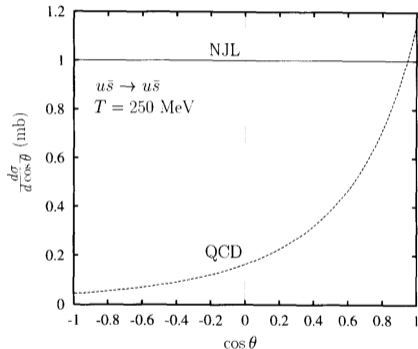
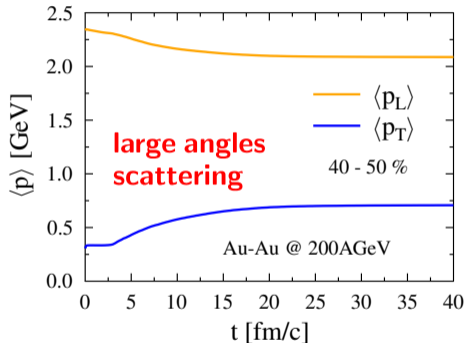
Results



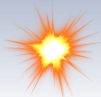


Small angles scattering

(NPA 608, 356 (1996))



We need pQCD-like **small angles scattering** !



Conclusion

The main messages:

- The **NJL model** provides a good framework to describe QGP around the critical temperature and allows for a **dynamical description** of the phase transition from a partonic medium to a hadron gas,
- Using the **PHSD initial conditions** (knowing that they reproduce RHIC data) gives us a good starting point with enough **granularity and fluctuations** in order to test event-by-event simulations,
- Out of equilibrium conversion from one model to another is **possible** knowing the equations of state in equilibrium,
- Initial cells at rest are **not in thermal equilibrium** nor in chemical equilibrium ($\langle m \rangle / \langle E \rangle$ ratio, $T(\varepsilon) \neq T^*(\langle |\vec{p}| \rangle)$).

Then ?

- Try first order phase transition for large baryonic densities (FAIR/NICA),
- Use Polyakov extended NJL model for better equation of state,
- ...



PHSD Group:

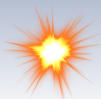
- Elena Bratkovskaya,
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- Olena Linnyk,
- Volodya Konchakovski,
- Hamza, Berrehrah,
- Daniel Cabrera,
- Taesoo Song.

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THANK YOU FOR YOUR ATTENTION !



System evolution

